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Retention of radiation damage in zircon xenocrysts from kimberlites, Northern Yakutia



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ARTICLE INFO

Article history: Received 17 March 2014 Accepted 6 August 2014 Available online 13 August 2014

Keywords: Xenocrystic zircon Kimberlite Yakutia Annealing SHRIMP U-Th-Pb ages

ABSTRACT

We have studied zircon xenocrysts from Mesozoic kimberlites from the Kuoika and Ary-Mastakh fields in Northern Yakutia. Zircon xenocrysts are assumed to originate from crustal rocks. Our SHRIMP (Sensitive High mass Resolution Ion MicroProbe) analyses yielded predominantly concordant U-Th-Pb ages (up to ~3570 Ma; Paleoarchean) that clearly predate kimberlite formation. The general U-Th-Pb concordance observed excludes notable disturbance of the zircon xenocrysts U-Th-Pb isotope system during kimberlite ascent and emplacement. In addition, zircon xenocrysts were found to be significantly more radiation-damaged than would correspond to damage accumulation only since the time of kimberlite formation. This observation first indicates that zircon crystals were sampled by the kimberlite magma at comparably shallow depths not exceeding 10-12 km. If, in contrast, zircon crystals originated from deeper levels of the Earth's crust, they would have been exposed to temperatures of 250-300 °C or more. This in turn would have caused long-term thermal annealing of the radiation damage, which was however not observed in our study. Second, our observation contradicts the hypothesis that high temperatures experienced by zircon xenocrysts during kimberlite ascent will cause notable structural reconstitution by short-term thermal annealing. Consequently, zircon crystals cannot have spent more than a few hours at temperatures exceeding ca. 700-800 °C, or more than a few days at temperatures exceeding ca. 500-600 °C. This in turn suggests that (i) temperatures of the ascending kimberlite magmas were rather moderate, and (ii) kimberlite ascent is a comparably short process followed by rapid cooling.

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1. Introduction

Kimberlite pipes in the Siberian Platform commonly contain zircon crystals, which belong to two groups. First, there are large grains with sizes ranging between ca. 300 µm and several centimetres, whose growth is related to kimberlite formation. Second, there are numerous smaller grains (sizes typically <300 µm) that are xenocrysts. The latter originate from rocks of the basement and/or lithospheric mantle, which were crushed and captured by the ascending kimberlite magma. In the Siberian Platform, zircon xenocrysts derived from crustal material originate predominantly from granulite xenoliths in the lower crust of the Siberian Craton (Buzlukova et al., 2004; Neymark et al., 1992; Rosen et al., 2002; Shatsky et al., 1990). Favoured by the physical stability and chemical and thermal resistance of zircon, such xenocrysts may yield a wealth of information on their primary rocks, and the milieu of conditions of kimberlite—magma formation (Koreshkova et al., 2009; Kostrovitsky et al., 2013; Vladykin and Lepekhina, 2009).

Here we report results of a study on zircon xenocrysts from kimberlite pipes in the Kuoika and Ary-Mastakh fields, Northern

Yakutia. The Kuoiksky field is located in the northeast of the Yakutian kimberlite province, in the lower section of the Olenek River (Fig. 1). It belongs to the Birehtinsk granite-greenstone terrane (Rosen, 2003). The Kuoika field is particularly rich in kimberlites; about one hundred pipes and dykes have been found thus far. The Ary-Mastakh field is located at the southeastern boundary of the Anabar shield (Fig. 1). Within the field there are about 40 pipes composed of kimberlites and alkaline picrites. Mesozoic ages for kimberlite pipes have been reported in the literature, based on U-Pb analyses of perovskite and zircon (Griffin et al., 1999; Kinny et al., 1997; Zaitsev and Smelov, 2010). Recent analyses yielded 170-150 Ma for Kuoika kimberlites, and two formation intervals (230–210 Ma and 170–160 Ma) for Ary–Mastakh kimberlites (Sun et al., 2014). Almost all pipes of the Kuoika and Ary-Mastakh fields were found to contain notable amounts of zircon xenocrysts in their heavy mineral fractions (Kostrovitsky, 1986; Kostrovitsky et al., 2013).

This present paper addresses the question as to what degree zircon xenocrysts are affected by the thermal impact of being sampled and transported to the Earth's surface by the hot kimberlite magma. In particular, possible thermal-annealing effects on the radiation damage that was accumulated prior to the kimberlite-formation event were focused on.

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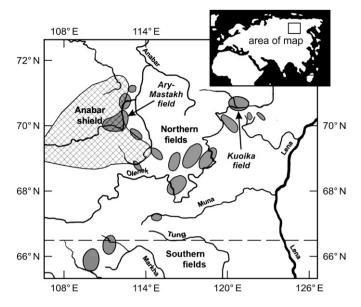


Fig. 1. Simplified map of a part of the Yakutian province, showing the locations of the Ary–Mastakh and Kuoika fields. Kimberlite fields are highlighted in grey. Figure from Kharkiv et al. (1998), substantially modified.

2. Background information: evaluation of damage retention

The cause of natural radiation damage, i.e. alpha-decay events experienced by zircon, is usually quantified for geological samples by the time-integrated alpha dose $D_{\alpha}.$ It is calculated (Murakami et al., 1991) according to

$$\begin{split} D_{\alpha} &= 8 \cdot \frac{c_{U} \cdot N_{A} \cdot 0.9928}{M_{238} \cdot 10^{6}} \cdot \left(e^{\lambda_{238}t} - 1\right) + 7 \cdot \frac{c_{U} \cdot N_{A} \cdot 0.0072}{M_{235} \cdot 10^{6}} \cdot \left(e^{\lambda_{235}t} - 1\right) \\ &+ 6 \cdot \frac{c_{Th} \cdot N_{A}}{M_{232} \cdot 10^{6}} \cdot \left(e^{\lambda_{232}t} - 1\right) \end{split} \tag{1}$$

with c_U and c_{Th} = present actinide concentrations (in $\mu g/g$), NA = Avogadro's number, M_{238} , M_{235} , and M_{232} = molecular weights of the parent isotopes, λ_{238} , λ_{235} , and λ_{232} = decay constants of parent isotopes, and t = integration time. The equation above assumes a present-day uranium composition of ~99.28% ²³⁸U and ~0.72% ²³⁵U (Cowan and Adler, 1976).

The calculated alpha dose, however, does in fact rarely correspond to the alpha-event damage present. This is because radiation damage may have been annealed, depending on the sample's thermal history (see detailed discussion by Nasdala et al., 2001). On the contrary, the comparison of the theoretical maximum damage, indicated by calculated alpha doses, and the damage actually present, provides information on the degree of damage storage by a given zircon.

The degree of the radiation damage present can be estimated on a micrometre-scale using Raman spectroscopy (Nasdala et al., 1995). For this purpose, the FWHM (full width at half-maximum) of the intense $\nu_3(\text{SiO}_4)$ Raman band of zircon (B_{1g} mode near 1000 cm⁻¹ assigned to anti-symmetric stretching of SiO_4 tetrahedrons; Dawson et al., 1971; Syme et al., 1977) is measured. This parameter changes from ~1.8 cm⁻¹ for well-crystallised zircon (Nasdala et al., 2002) to well above 30 cm⁻¹ for strongly damaged zircon.

Quantitative estimates of the degree of damage retention in unknowns are made by comparing the damage actually present with the theoretical maximum damage. This is possible by comparing FWHM– D_{α} data pairs of unknowns with those of unannealed zircon samples (Nasdala et al., 2001). To elucidate how such comparison may provide an estimate of the damage retention, and hence yield valuable conclusions on the thermal history of unknowns, we present two examples

in Fig. 2a–c. In the present study, we apply this method to study whether or not the temperature increase in the kimberlite magma has caused major annealing of the previously accumulated radiation damage in zircon xenocrysts.

Data pairs for metamorphic zircon from a sillimanite-biotite gneiss from Val Strona, Ivrea Zone, Italian Alps (Vavra et al., 1999) and those of unannealed zircon (Nasdala et al., 2001) show a clear mismatch (Fig. 2a) if D_{α} values for the Ivrea zircon are calculated assuming selfirradiation since the time of closure of the U-Th-Pb system. Consequently, the Ivrea zircon is much less radiation-damaged as would correspond to complete damage accumulation since the time of the high-grade metamorphic zircon growth. This in turn indicates a significant annealing history. However, FWHM– D_{α} data pairs of Ivrea zircon and unannealed references match reasonably well (Fig. 2b) if D_{α} values for the Ivrea zircon are calculated based on the ~31 Ma age of the Cenozoic uplift (Hurford et al., 1989). Both observations together suggest that after metamorphic growth, the Ivrea zircon may have resided at considerable depths below the Earth's surface for an extended period of time, undergoing continuous thermal annealing. Damage accumulation started only at the time of the Cenozoic uplift of the Alps (Nasdala et al., 2001).

In contrast, FWHM– D_{α} data pairs for zircon from Plešovice (well-known as a U–Pb and Hf isotope reference material; Sláma et al., 2008) correlate with FWHM– D_{α} data pairs of unannealed zircon, with time-integrated alpha doses of the Plešovice zircon being calculated based on this sample's U–Th–Pb age (Fig. 2c). This observation suggests that the Plešovice zircon has an accumulated level of radiation damage that excludes any major thermal event after magmatic growth and closure of the U–Th–Pb system (for details see Sláma et al., 2008).

3. Samples and experimental techniques

Xenocrysts were extracted by conventional mechanical separation from the heavy mineral fractions of crushed kimberlite material that was sampled in near-surface outcrops. Zircon grains were embedded in araldite epoxy, oriented roughly along their crystallographic c-axes, and ground and polished to about one half of their thicknesses, to expose their internal textures. These sample mounts were first imaged in reflected light under an optical microscope and then subjected to Raman microprobe analysis. After being carbon-coated, samples were subjected to BSE (back-scattered electrons) imaging and chemical analysis using EPMA (electron probe micro-analyser). Isotopic analyses by SHRIMP (Sensitive High mass Resolution Ion MicroProbe) were done after removal of the surficial carbon layer and coating with gold. After gentle re-polishing to remove the shallow SHRIMP analysis pits (depths of ca. 0.7–1.1 µm; cf. Nasdala et al., 1998), some samples were attached to a glass slide, and doubly polished thin sections were produced for transmitted-light optical microscopy.

The degree of radiation damage accumulated by zircon samples was first assessed from the birefringence observed in cross-polarised transmitted light (Chakoumakos et al., 1987; Sahama, 1981). Second, it was estimated quantitatively from the Raman-band broadening (Nasdala et al., 1995, 2001). Raman spectra were obtained by means of a Horiba LabRam–HR 800 system. This system was equipped with an Olympus BX41 optical microscope, a diffraction grating with 1800 grooves/mm, and a Peltier-cooled charge-coupled device (CCD) detector. Spectra were excited using He-Ne 632.8 nm emission (8 mW at the sample surface). The laser beam was focused at the sample surface using a $100 \times$ objective (numerical aperture 0.9). With the system operated in the confocal mode, the lateral resolution was ~1 µm. The volume resolution is assessed better than 5 μ m³ (compare Markwort et al., 1995). The spectral resolution was ~0.8 cm⁻¹. A minimum of three replicate measurements per sampling point were done to check for potential sample heterogeneity on a scale of within SHRIMP spots and to improve the reliability of fit results. Spectra were subjected to background correction and fitting was done assuming Lorentzian-Gaussian band shapes. Real

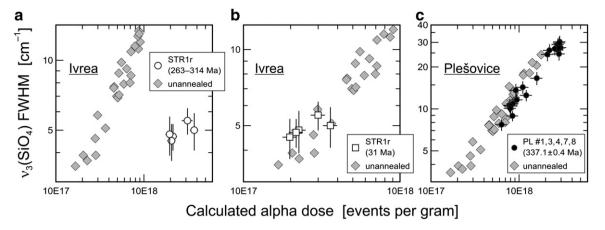


Fig. 2. Plot of the broadening of the main Raman band of zircon (FWHM = full width at half maximum; indicating the degree of radiation damage) of two zircon samples against calculated alpha doses, in comparison to unannealed zircon samples (reference data from Nasdala et al., 2001). (a) Zircon from the Ivrea Zone, alpha doses calculated based on U-Th-Pb zircon ages (Vavra et al., 1999). (b) Same sample, alpha doses calculated for the time period since the Cenozoic uplift (Hurford et al., 1989). (c) Zircon from Plešovice, alpha doses calculated based on U-Th-Pb age (Sláma et al., 2008).

FWHM values were calculated by correcting measured FWHMs for the apparatus function of the Raman instrument (Dijkman and van der Maas, 1976; Váczi, 2013). The wavenumber accuracy was better than $0.5\,\mathrm{cm}^{-1}$, and total uncertainties of corrected FWHMs (i.e. including uncertainties due to micrometre-scale heterogeneity) are assessed to be smaller than 10%.

The general chemical composition of the zircon samples was determined by means of a JEOL 8900 RL electron microprobe. The electron beam was focused to a 2 μm spot. The accelerating voltage was 20 kV and the beam current was 100 nA. The usual metal, oxide, and mineral standards were used, including ZrSiO₄ for Zr and Si, synthetic HfSiO₄ for Hf, synthetic ThSiO₄ for Th, and synthetic UO₂ for U. Counting times were varied between 15 s (peak) and 2 \times 5 s (lower and upper backgrounds) for the K α line of Si, and 300 s (peak) and 2 \times 150 s (backgrounds) for the M α lines of U and Th. Data were reduced using the CITZAF routine in the JEOL software, which is based on the $\Phi(\rho Z)$ correction method (Armstrong, 1991, 1995).

Analyses of U and Th concentrations and the Pb isotopic composition were made on a SHRIMP II at Curtin University of Technology, Perth. A primary, mass-filtered $(O_2)^-$ beam with ~1 nA current was focused to a ~10 µm spot to sputter the sample surface. Data for each spot were collected in sets of eight scans through the mass range of Zr_2O^+ , $^{204}Pb^+$, background, $^{206}Pb^+$, $^{207}Pb^+$, $^{208}Pb^+$, $^{238}U^+$, $^{248}ThO^+$, and $^{254}UO^+$. The mass resolution, M/ Δ M, was better than 5000. For more experimental details see De Laeter and Kennedy (1998). Results were calibrated against M257, a 561.3 Ma old reference zircon (Nasdala et al., 2008). Common Pb was assumed to have a Broken Hill Pb composition and corrected for using the ^{204}Pb method (Compston et al., 1984). The Excel-based programme Squid (Ludwig, 2002) was used for data processing. Data were plotted using the Isoplot programme (Ludwig, 2003).

4. Results and discussion

4.1. Internal textures

Zircon xenocrysts from both the Ary–Mastakh and the Kuoika fields have predominantly medium brown to yellowish brown (rarely reddish to pinkish brown) colours, and elongated to short-prismatic shapes that often appear rounded. Their sizes range from below 50 µm to occasionally several hundred micrometres in length. The internal textures show quite some diversity, ranging from more or less uniform growth zoning (Fig. 3a, c) to complex resorption/overgrowth and alteration patterns (Fig. 3g). Many of the grains contain an inherited component, either small euhedral cores (Fig. 3b, c) or irregularly shaped regions with

discordantly cut zoning (Fig. 3f, g). The inherited cores or interior regions themselves may show complex zoning (Fig. 3b, inset; g). Crusts or overgrowth rims of ZrO_2 , which typically form as a result of zircon alteration and decomposition in the kimberlite melt (Page et al., 2007; Robles-Cruz et al., 2012), were not observed in the present study.

Interference colours observed in cross-polarised transmitted light show remarkable scatter. Some interior regions yield "fresh" 3rd order interference colours whereas others appear isotropic (both can be seen in Fig. 3g). This indicates that the structural state of samples is characterised by strong internal diversity, ranging from mildly radiation-damaged to metamict (i.e. amorphised; Chakoumakos et al., 1987; Morgan and Auer, 1941; Sahama, 1981). Strong differences in the degree of radiation damage among zones and interior regions are also indicated by common fractures in the less damaged zones and regions (Fig. 3e-g). Fractures are typically oriented more or less perpendicular to zone boundaries; they are caused by strongly differential damage-induced volume expansion (Chakoumakos et al., 1987; Lee and Tromp, 1995; Nasdala et al., 2010; Smith et al., 1991). Interference colours show inverse correlation with BSE intensities (with the exception of the very low BSE discussed below). Interior regions and zones with lowered interference colours (and hence elevated degrees of radiation damage) show higher BSE intensity, and vice versa (Fig. 3d, g). This reconfirms again that the electron back-scatter coefficient of zircon is predominantly controlled by this mineral's structural state (Nasdala et al., 2006), whereas BSE Z-contrast caused by chemical variations is mostly insignificant in the case of natural zircon. The latter was shown experimentally by Nasdala et al. (2006) who found that (i) annealed halves of cut zircon crystals are much lower in BSE than their unannealed, chemically equivalent counterparts, and (ii) the BSE intensity of zircon correlates with the degree of radiation damage measured by Raman spectroscopy.

Another remarkable textural feature is the minor presence of interior regions with extremely low intensities in BSE images (compare Kempe et al., 2000; Pointer et al., 1988; Smith et al., 1991). This phenomenon is a reliable indicator of secondary alteration of zircon, caused by sub-micrometre porosity of the altered material and possibly the presence of water in pores (Nasdala et al., 2009; Sorbier et al., 2004). In the zircon samples studied here, such alteration patches are mostly small in size and typically occur within more radiation-damaged zones, near the ends of fractures that emanate from less radiation-damaged zones (Fig. 3e, g). Obviously less-damaged zones resisted the alteration (cf. Krogh and Davis, 1974; Medenbach, 1976). Fractures in the less-damaged zones evidently served as pathways for alteration fluids which then attacked and altered the neighbouring, more damaged areas (see BSE images in Fig. 3e, g; compare Lumpkin, 2001;

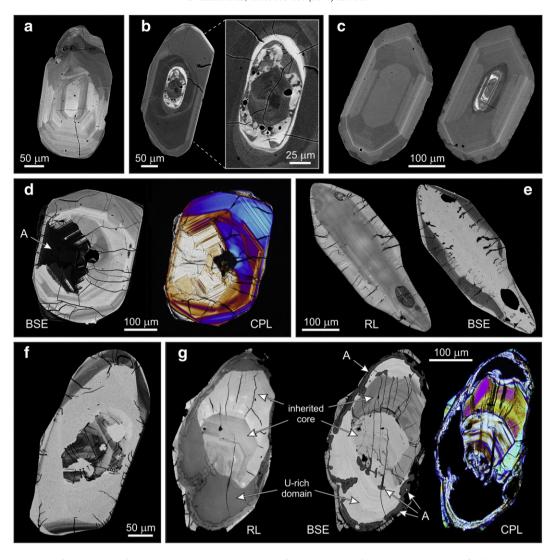


Fig. 3. Images of zircon xenocrysts from kimberlites from Northern Yakutia. (a–c) BSE images of three zircon grains from the Rubin pipe, Kuoika field (sample 7-51; a – grain #2, b – grain #5, c – grain #25). Note that the grain shown in (c) was imaged twice. The small inherited core was only detected after being exposed by gentle re-grinding (right image). (d) BSE and cross-polarised transmitted light (CPL) image of a zircon crystal from the Bumerang pipe, Ary–Mastakh field (sample 90-2-1; grain #2). An area affected by young alteration (marked "A") shows very low BSE and lowered interference colour. (e) Reflected light (RL) and BSE images of another zircon (grain #7) from the Bumerang pipe. (f) BSE image of a complex grain from the Bargidamalah pipe, Ary–Mastakh field (sample 90-22; grain #6). The central interior region shows strong primary zoning whereas the majority of the grain appears uniformly bright. (g) Series of RL, BSE, and CPL images of another grain (#4) from sample 90-22. The main, U-rich interior region is metamict; it is characterised by lowered reflectance, bright BSE and absence of any birefringence. The inherited, zoned core is lower in U and shows low to moderate interference colours. Small amounts of altered zircon (very dark BSE; marked "A") is observed near the rim of the grain and at the ends of large internal fractures.

Nasdala et al., 2010). In contrast, narrow near-surface alteration rims (Fig. 3g) or large alteration patches were observed only rarely (Fig. 3d).

4.2. Chemical compositions

Results of chemical analyses are presented in Table 1. Except for the alteration patches, zircon grains have rather "regular" chemical compositions. Hafnium (HfO₂ 0.73–1.34 wt.%) and uranium are the only two major non-formula constituents. Uranium concentrations are mostly several hundreds of ppm by weight or even less. Only some high-BSE interior regions contain elevated uranium with UO₂ up to 0.43 wt.% (Table 1). The partially high U content of unaltered zircon supports its xenocrystic origin, as kimberlite zircon typically shows low U concentrations of below 100 ppm (e.g. Belousova et al., 1998; Haggerty et al., 1983; Spetsius et al., 2002; Valley et al., 1998; Zartmann and Richardson, 2005).

Altered areas, in contrast, show elevated concentrations of Ca and Fe, with up to several wt.% oxides. Note that these two elements are typically excluded during primary zircon growth and incorporated in

low-temperature hydrothermal alteration processes (Geisler et al., 2003; Pérez-Soba et al., 2007; Törnroos, 1985). The analytical totals of altered areas are clearly deficient (93.1 and 94.9 wt.%, respectively, for the two analyses in altered areas listed in Table 1). The analytical EPMA shortfall may be assigned to the porous sub-micrometre texture of altered material (Nasdala et al., 2009; Putnis, 2002). Such porosity first causes charge trapping and energy losses at interfaces (Sorbier et al., 2004). Second, pores may in addition host chemical species that are not analysed in the EPMA, especially water (Geisler et al., 2003; Johan and Johan, 2005; Smith et al., 1991; Törnroos, 1985; Utsunomiya et al., 2007).

4.3. U-Th-Pb geochronology

Results of U–Th–Pb analyses are listed in Table 2, and selected Concordia plots are presented in Fig. 4. Some of the analyses with high U content show minor reverse discordance (up to -9%; cf. Table 2), which is probably an analytical artefact assigned to matrix effects that

Table 1Results of EPMA chemical analyses (wt.%).

Grain	Interior region	Al_2O_3	SiO_2	CaO	FeO	$Y_{2}O_{3}$	ZrO_2	Yb_2O_3	HfO_2	PbO	ThO_2	UO_2	Total
Xenocrysts from p	ipe Ruby, Kuoika field												
7-51 #2	High BSE	(bdl)	32.8	(bdl)	(bdl)	0.14	65.9	(bdl)	0.97	(bdl)	0.04	0.04	100.0
7-51 #2	Intermediate BSE	(bdl)	32.5	(bdl)	(bdl)	(bdl)	66.1	(bdl)	1.08	(bdl)	(bdl)	0.03	99.8
7-51 #2	Low BSE	(bdl)	32.9	(bdl)	(bdl)	(bdl)	65.8	(bdl)	1.03	0.04	(bdl)	(bdl)	100.0
7-51 #5	Bulk	(bdl)	32.5	(bdl)	(bdl)	(bdl)	65.9	(bdl)	1.14	(bdl)	(bdl)	0.03	99.8
7-51 #5	Core	(bdl)	32.7	(bdl)	(bdl)	0.09	65.7	(bdl)	1.23	(bdl)	(bdl)	0.07	100.0
7-51 #5	Core, high BSE	(bdl)	32.8	(bdl)	(bdl)	0.12	65.6	(bdl)	1.10	0.04	(bdl)	0.14	99.9
7-51 #8	-	(bdl)	32.8	(bdl)	(bdl)	(bdl)	66.2	(bdl)	1.13	(bdl)	(bdl)	0.02	100.3
7-51 #11		(bdl)	32.7	(bdl)	(bdl)	(bdl)	65.7	(bdl)	1.34	(bdl)	(bdl)	0.07	100.0
7-51 #25		(bdl)	32.7	(bdl)	(bdl)	0.13	66.3	(bdl)	0.73	(bdl)	(bdl)	(bdl)	100.0
Xenocrysts from p	ipe Bumerang, Ary–Masto	ıkh field											
90-2-1 #2	Bulk	(bdl)	32.7	(bdl)	(bdl)	(bdl)	65.0	(bdl)	1.01	(bdl)	(bdl)	(bdl)	98.8
90-2-1 #2	Altered area	(bdl)	29.8	3.60	0.95	0.12	57.5	(bdl)	0.92	(bdl)	0.03	0.07	93.1
90-2-1 #7	Central region	(bdl)	32.5	(bdl)	(bdl)	0.10	65.8	(bdl)	1.07	0.03	(bdl)	0.08	99.8
90-2-1 #7	Rim	0.04	32.6	(bdl)	0.14	(bdl)	65.5	(bdl)	1.11	(bdl)	(bdl)	(bdl)	99.6
90-2-1 #8	High BSE	(bdl)	32.8	(bdl)	(bdl)	(bdl)	65.6	(bdl)	0.96	(bdl)	(bdl)	0.06	99.7
90-2-1 #8	Low BSE	(bdl)	32.7	(bdl)	(bdl)	(bdl)	65.7	(bdl)	1.07	(bdl)	(bdl)	0.02	99.6
90-2-1 #9		(bdl)	32.9	(bdl)	(bdl)	(bdl)	65.9	(bdl)	0.97	(bdl)	(bdl)	0.03	99.9
90-2-1 #10		(bdl)	32.8	(bdl)	(bdl)	(bdl)	66.2	(bdl)	1.03	(bdl)	(bdl)	(bdl)	100.1
Xenocrysts from p	ipe Bargidamalah, Ary–M	astakh field											
90-22 #2		(bdl)	32.7	(bdl)	(bdl)	(bdl)	66.3	(bdl)	0.89	(bdl)	(bdl)	(bdl)	100.1
90-22 #3		(bdl)	32.8	(bdl)	(bdl)	(bdl)	66.2	(bdl)	0.84	(bdl)	(bdl)	0.03	99.9
90-22 #4	Core, high BSE	(bdl)	32.7	(bdl)	(bdl)	0.17	66.0	0.11	1.03	0.06	0.04	0.06	100.4
90-22 #4	Core, low BSE	(bdl)	32.6	(bdl)	(bdl)	(bdl)	66.6	(bdl)	1.09	(bdl)	(bdl)	(bdl)	100.3
90-22 #4	Bulk	(bdl)	32.4	(bdl)	(bdl)	0.19	65.5	0.14	1.23	0.30	(bdl)	0.43	100.3
90-22 #4	Alteration rim	(bdl)	29.5	0.85	3.51	0.09	59.4	(bdl)	1.06	0.03	(bdl)	0.14	94.9
90-22 #5		(bdl)	32.7	(bdl)	(bdl)	(bdl)	66.3	(bdl)	0.89	(bdl)	(bdl)	(bdl)	100.1
90-22 #6	High BSE	(bdl)	32.4	(bdl)	(bdl)	(bdl)	66.0	(bdl)	1.13	(bdl)	(bdl)	0.07	99.7
90-22 #6	Low BSE	(bdl)	32.6	(bdl)	(bdl)	(bdl)	66.4	(bdl)	1.13	(bdl)	(bdl)	0.02	100.2
Detection limit ^a		0.04	0.1	0.04	0.08	0.08	0.1	0.11	0.18	0.03	0.03	0.02	-

bdl = not detected, or below the detection limit quoted in the bottom row.

commonly occur in SIMS analyses of high-U zircon (for instance see McLaren et al., 1994).

The U–Th–Pb ages determined are predominantly Paleoproterozoic, with individual $^{207}\text{Pb}/^{206}\text{Pb}$ age values between 2078 \pm 31 and 1843 \pm 98 Ma (Table 2; Fig. 4a, b). These results seem to concur well with Paleoproterozoic ages determined for collision-related metamorphism and mafic magmatism events within the Anabar province (Rosen et al., 2006). They are assigned to multiple collision stages of accreting independent microcontinents (Rosen et al., 2006).

The ancient microcontinents (compare Rosen, 2002; Rosen et al., 2005; and references therein) are perhaps represented by the Siderian $^{207}\text{Pb}/^{206}\text{Pb}$ age of 2389 \pm 107 Ma that was determined for grain 7-51 #25 (Ruby pipe, Kuoika; Fig. 3a), and the Mesoarchaean $^{207}\text{Pb}/^{206}\text{Pb}$ age of 2832 \pm 12 Ma determined for the U-rich, bright-BSE domain of grain 90-22 #4 (Bargidamalah pipe, Ary–Mastakh; Figs. 3g, 4c). For instance, the age of the former seems to concur very well with the ~2.4 Ga age of the Birehtinsk granite–greenstone terrane (Rosen, 2003). The core of grain 90-22 #4 yielded a Paleoarchean $^{207}\text{Pb}/^{206}\text{Pb}$ age of 3569 \pm 33 Ma (Fig. 4c), which is assigned to an inherited component of unknown primary origin.

Note, however, that the present study was not focused on primary ages of zircon xenocrysts; for such kind of petrogenetic study a much higher number of grains needs to be analysed (compare Batumike et al., 2007; Donatti-Filho et al., 2013; Zheng et al., 2009). Nevertheless our results indicate a non-uniform provenance; zircon xenocrysts obviously were derived from several source rocks. This conclusion is also supported by the diversity of internal textures (Fig. 3), which we interpret to reflect different formation and post-growth histories of the different source rocks.

All U–Th–Pb ages clearly predate the Mesozoic kimberlite formation. This supports again the assignment of zircon grains as xenocrysts. The majority of ages determined are concordant or only slightly discordant. There is consequently no indication of significant Pb loss from zircon

xenocrysts during kimberlite ascent. This is in accordance with earlier findings that the U–Th–Pb system in zircon xenocrysts – in contrast to the (U–Th)/He system – often remains unaffected during the short kimberlite ascent and emplacement (McInnes et al., 2009; Vladykin and Lepekhina, 2009). Possible Pb loss during kimberlite ascent (indicated by lower intercepts that concur within errors with kimberlite eruption ages) has for instance been discussed by Zheng et al. (2009) for zircon xenocrysts from kimberlites in the North China Craton. In the present study, strong U–Pb discordance was found only in one case (altered area in grain 90-2-1 #2; Fig. 3d). We assign this observation (Fig. 4b) to recent Pb loss that was most probably caused by near-surface weathering.

4.4. Estimation of the radiation damage

The $\nu_3({\rm SiO_4})$ Raman band of zircon (i.e. crystalline ZrSiO₄) disappears entirely when the metamictisation process is completed and amorphous ZrSiO₄ has formed (Nasdala et al., 2001). The latter was found for the high-BSE domain of grain 90-22 #4, which is in agreement with the observations of strongly lowered reflectivity and optical isotropy of this interior region (Fig. 3g).

All other zircon micro-areas studied were found to be incompletely radiation-damaged. Mild damage, indicated by small Raman-band FWHMs (<5 cm⁻¹; Table 2), was for instance detected in grain 7-51 #25 (Fig. 3c), the low-BSE zone of grain 7-51 #2 (Fig. 3a), and grain 90-22 #2 (not shown). Such grains, or zones within grains, typically show low BSE (and high CL) intensity, and high interference colours. Strong radiation damage, indicated by FWHMs exceeding 20 cm⁻¹ (Table 2), was determined for instance for the central interior region of grain 90-2-1 #7 (Fig. 3e), the central, high-BSE region of the core of grain 90-22 #4 (Fig. 3g) and also for the large alteration patch in grain 90-2-1 #2 (Fig. 3d). Strongly damaged grains or interior regions are

^a Average detection limit (wt.% oxide; 2σ error by counting statistics of the background signal).

Table 2 SHRIMP and Raman microprobe results for zircon xenocrysts.

Measurement		SHRIMP ion microprobe results								α dose ^a		Raman		
Grain	Interior region	U (ppm)	Th (ppm)	Th/U	f 206 ^b (%)	$^{206}\text{Pb}/^{238}\text{U}$	²⁰⁶ Pb/ ²³⁸ U age (Ma)	²⁰⁷ Pb/ ²³⁵ U	²⁰⁷ Pb/ ²⁰⁶ Pb	²⁰⁷ Pb/ ²⁰⁶ Pb age (Ma)	Disc.c (%)	U-Pb age	Kimberlite age	FWHM d (cm ⁻¹
(enocrysts	from pipe Ruby, Kuoil	ka field												
7-51 #2	High BSE	406	353	0.86	0.09	0.0897 ± 0.0025	554 ± 15	0.721 ± 0.023	0.0583 ± 0.0009	538 ± 34	-3	0.91	0.26	6.9
7-51 #2	Intermediate BSE	224	161	0.72	0.24	0.0894 ± 0.0022	552 ± 13	0.714 ± 0.031	0.0580 ± 0.0021	526 ± 79	-5	0.49	0.14	4.8
7-51 #2	Low BSE	80	33	0.41	0.00	0.0837 ± 0.0028	518 ± 17	0.698 ± 0.030	0.0604 ± 0.0016	617 ± 57	+17	0.16	0.05	2.7
7-51 #5	Bulk	222	82	0.37	0.10	0.3600 ± 0.0089	1982 ± 42	6.25 ± 0.16	0.1258 ± 0.0008	2073 ± 12	+3	1.9	0.13	7.9
7-51 #5	Bulk	162	55	0.34	0.00	0.3548 ± 0.0093	1957 ± 44	5.86 ± 0.19	0.1199 ± 0.0024	1951 ± 36	-0	1.4	0.09	8.2
'-51 #5	Bulk	329	123	0.38	0.04	0.3441 ± 0.0077	1906 ± 37	5.56 ± 0.13	0.1173 ± 0.0006	1911 ± 10	+0	2.8	0.19	10.8
-51 #5	Core	658	56	0.09	0.02	0.3751 ± 0.0076	2053 ± 36	6.32 ± 0.13	0.1221 ± 0.0007	1984 ± 11	-4	5.2	0.35	(n.a.)
-51 #5	Core, high BSE	1574	16	0.01	0.00	0.3576 ± 0.0067	1971 ± 32	5.86 ± 0.11	0.1189 ± 0.0008	1936 ± 4	-2	12.4	0.83	(n.a.)
-51 #8		207	51	0.25	-0.02	0.3567 ± 0.0091	1966 ± 43	5.82 ± 0.15	0.1184 ± 0.0008	1929 ± 13	-2	1.7	0.12	10.2
-51 #8		88	30	0.34	0.05	0.3392 + 0.0111		5.35 ± 0.19	0.1144 ± 0.0013		-1	0.74	0.05	6.1
-51 #11		579	55	0.09	0.00	0.3528 ± 0.0148	1948 + 70	5.96 ± 0.26	0.1226 ± 0.0008		+2	4.6	0.31	13.0
-51 #25		155	119	0.77	0.18	0.4383 ± 0.0123		9.60 ± 0.29	0.1588 ± 0.0018		+5	1.4	0.10	3.4
-51 #25		50	50	0.99	-0.09	0.4731 ± 0.0191		11.26 ± 0.49	0.1727 ± 0.0028		+4	0.47	0.03	2.7
-51 #25		20	12	0.60	0.44	0.4309 ± 0.0285		9.08 ± 0.64	0.1527 ± 0.0035		+3	0.18	0.01	2.6
enocrysts	from pipe Bumerang,	Arv_Mastak	h field											
0-2-1 #2	Bulk	137	37	0.27	0.06	0.3346 ± 0.0095	1860 + 46	5.70 ± 0.17	0.1235 ± 0.0010	2003 + 15	+8	1.1	0.09	11.5
0-2-1 #2	Altered area	545	204	0.38	4.67	0.1431 ± 0.0030	862 ± 17	2.48 ± 0.07	0.1255 ± 0.0021		+61	4.7	0.38	20.7
0-2-1 #2	Altered area	621	221	0.36	4.21	0.1370 ± 0.0028	828 ± 16	2.43 ± 0.07	0.1288 ± 0.0023		+64	5.3	0.43	21.6
)-2-1 #7	Central region	882	284	0.32	0.02	0.3542 ± 0.0069		5.94 ± 0.12	0.1200 ± 0.0025 0.1217 ± 0.0004		+1	7.5	0.61	30.7
0-2-1 #7	Central region	929	318	0.34	0.02			5.92 ± 0.12	0.1217 ± 0.0004 0.1211 ± 0.0004		+1	7.9	0.65	20.7
0-2-1 #7	Rim	141	34	0.24	0.07	0.3728 ± 0.0100		6.18 ± 0.17	0.1203 ± 0.0004		-5	1.2	0.10	9.9
0-2-1 #7	Rim	135	35	0.24	0.07			5.78 ± 0.17	0.1203 ± 0.0010 0.1192 ± 0.0010		-0	1.1	0.09	10.6
0-2-1 #7		530	129	0.24	0.11			5.78 ± 0.17 5.99 ± 0.13	0.1192 ± 0.0010 0.1209 ± 0.0005		-0 -1	4.4	0.36	24.7
	High BSE		63	0.24		0.3596 ± 0.0074							0.15	
0-2-1 #8	Low BSE	212			0.02	0.3485 ± 0.0086		5.78 ± 0.15	0.1202 ± 0.0008		+2	1.8		16.6
0-2-1 #9 0-2-1 #10	ı	267 194	106 45	0.40 0.23	1.87 0.08	0.3159 ± 0.0074 0.3515 ± 0.0090		5.38 ± 0.15 5.94 ± 0.16	0.1236 ± 0.0018 0.1226 ± 0.0009	2005 ± 26 1990 ± 12	+13 +3	2.3 1.6	0.19 0.13	13.1 18.9
				0.23	0.00	0.5515 ± 0.0050	1312 ± 13	3.3 T ± 0.10	0.1220 ± 0.0003	1330 ± 12	13	1.0	0.13	10.5
enocrysts j 0-22 #2	from pipe Bargidama	ah, Ary–Ma 11	stakh field 5	0.50	0.52	0.3544 ± 0.0163	1056 79	E E2 0.20	0.1129 ± 0.0061	1042 00	- 7	0.01	0.01	2.4
10-22 #2 10-22 #3		276		0.50	0.52			5.52 ± 0.39				2.3	0.01	12.3
			112			0.3478 ± 0.0069	1924 ± 33	5.73 ± 0.12	0.1195 ± 0.0009	1945 ± 13	+1			
0-22 #3	Come hint DCE	419	193	0.46	0.17	0.3585 ± 0.0069		5.94 ± 0.12	0.1202 ± 0.0007		-1	3.6	0.30	15.6
0-22 #4	Core, high BSE	551	279	0.51	0.00	0.7331 ± 0.0139		32.6 ± 0.62	0.3220 ± 0.0006		+1	7.8	0.40	20.2
0-22 #4	Core, high BSE	415	194	0.47	0.03	0.7090 ± 0.0136		30.5 ± 0.60	0.3119 ± 0.0013		+3	5.9	0.30	22.7
0-22 #4	Core, high BSE	453	213	0.47	0.02	0.7207 ± 0.0138		31.9 ± 0.61	0.3208 ± 0.0007		+3	6.4	0.32	16.4
0-22 #4	Core, high BSE	447	232	0.52	0.01	0.7090 ± 0.0136		30.8 ± 0.60	0.3146 ± 0.0012		+3	6.4	0.32	27.9
0-22 #4	Core, low BSE	130	49	0.38	0.02	0.7681 ± 0.0161		35.1 ± 0.82	0.3313 ± 0.0035		-2	1.8	0.09	12.0
0-22 #4	Core, low BSE	141	54	0.38	0.08	0.7846 ± 0.0165		35.1 ± 0.75	0.3247 ± 0.0013		-5	1.9	0.10	10.1
0-22 #4	High BSE	2751	106	0.04	0.01	0.5565 ± 0.0103		15.2 ± 0.28	0.1987 ± 0.0035		-2	36.1	1.8	(Metamict)
0-22 #4	High BSE	3406	149	0.04	0.00	0.5608 ± 0.0106		15.6 ± 0.30		2840 ± 4	-1	44.7	2.2	(Metamict)
0-22 #4	High BSE	3708	144	0.04	0.00	0.5735 ± 0.0107		16.0 ± 0.30	0.2023 ± 0.0004		-4	48.7	2.4	(Metamict)
0-22 #4	High BSE	2103	78	0.04	0.00	0.5999 ± 0.0111	3030 ± 45	16.7 ± 0.31	0.2017 ± 0.0006	2835 ± 5	-9	27.6	1.4	(Metamict)
0-22 #4	High BSE	3490	143	0.04	0.00	0.5719 ± 0.0105	2916 ± 43	15.9 ± 0.29	0.2011 ± 0.0002	2831 ± 2	-4	45.8	2.3	(Metamict)
0-22 #5		25	24	0.97	0.53	0.5561 ± 0.0169	2850 ± 70	16.8 ± 0.58	0.2189 ± 0.0036	2968 ± 26	+5	0.4	0.02	5.5
0-22 #6	High BSE	772	74	0.10	0.06	0.3387 ± 0.0066	1880 ± 32	5.49 ± 0.11	0.1175 ± 0.0007	1915 ± 11	+2	5.9	0.51	17.0
0-22 #6	High BSE	531	141	0.27	0.04	0.3379 ± 0.0064	1877 ± 31	5.37 ± 0.11	0.1153 ± 0.0005	1881 ± 8	+0	4.2	0.36	23.0
0-22 #6	Low BSE	300	84	0.28	0.04	0.3471 ± 0.0068	1921 + 32	5.48 ± 0.11	0.1146 ± 0.0007	1870 + 11	-3	2.4	0.21	19.1

Isotopic ratios are corrected for common Pb (204Pb method). Isotopic ratios and ages are quoted with 1σ uncertainties.

a Time-integrated alpha dose, calculated twice (first, assuming damage accumulation since the time of closure of the U–Th–Pb system; second, assuming damage accumulation since the time of kimberlite formation).

b f 206 = common 206Pb/total 206Pb, calculated from the observed 204Pb.

c Disc. = U–Pb discretainties between the 204Pb-corrected 206Pb/238U and 207Pb/206Pb ages.

d Average of the full width at half maximum of the v_3 (SiO₄) Raman band of zircon, determined from multiple analyses. n.a. = not analysed (micro-area not existent anymore after re-polishing). metamict = amorphous micro-areas not yielding any measurable Raman signal of crystalline ZrSiO₄.

characterised by increased BSE and low (and partially anomalous) interference colours.

4.5. Degree of damage retention

Time-integrated alpha doses D_{α} of the zircon xenocrysts under investigation are listed in Table 2. These doses were calculated twice. First, they were calculated for the time period since closure of the U–Th–Pb system (for inherited cores, the age of the corresponding bulk was used). Second, doses were calculated for the time period since Mesozoic kimberlite formation (Table 2). Data pairs are plotted in Fig. 5.

Alpha doses based on U-Th-Pb ages quantify the entire radioactive self-irradiation since zircon formation (or since the hypothetical annealing of cores upon being overgrown, respectively). The corresponding FWHM-D_α data pairs plot below the trend defined by data pairs for unannealed zircon (Fig. 5a, b). This indicates that zircon xenocrysts did not completely accumulate the radiation damage caused by radioactive self-irradiation since the time of their formation. This conclusion is supported by geological evidence. For instance, in a study of Neoproterozoic crustal xenoliths (ages about 2000-1960 Ma) in kimberlites from the Anabar shield, Rosen et al. (2006) found that cooling from 700 to ~300 °C lasted approximately 300 m.y. On the other hand, it is well known that at temperatures exceeding 200–300 °C, radiation damage is not accumulated but, in contrast, moderately radiation-damaged zircon undergoes slow structural reconstitution (e.g., Meldrum et al., 1998, 1999). Consequently, accumulation of radiation damage in zircon xenocrysts cannot have started immediately after closure of the U-Th-Pb system (roughly 600-650 °C under metamorphic conditions; Mezger and Krogstad, 1997).

The extents of mismatches of FWHM– D_{α} data pairs observed in the present study (Fig. 4a, b) indicate that zircon xenocrysts are roughly half as radiation-damaged as corresponds to their ages. This high level of damage storage in zircon xenocrysts implies that previous host rocks must have resided at temperatures well below 300 °C, for at least one half or their entire history. Assuming a geothermal gradient of roughly 25–30 °C/km for the Northern Yakutian kimberlite province, this suggests that zircon xenocrysts sampled by the ascending kimberlite magmas originate from comparably shallow levels of the Earth's crust, not exceeding 10–12 km. In view of the local crustal thickness of 40–45 km (Suvorov et al., 2006), this indicates that the zircon xenocrysts studied here were probably not derived from lower but upper crustal rocks

Data pairs of Raman FWHMs and alpha doses based on the kimberlite-formation ages (Fig. 5c, d), in contrast, plot above and to the left of the unannealed zircon trend. In fact, the Raman-band broadening of xenocrysts exceeds several times the band broadening of unannealed references with the same calculated alpha doses. Consequently, zircon xenocrysts are much more radiation-damaged than it would correspond to damage accumulation only since the time of kimberlite formation. Our observations hence contradict the hypothesis that high temperatures during kimberlite ascent will cause complete structural reconstitution of zircon xenocrysts, by thermal annealing of the accumulated radiation damage.

Zhang et al. (2000) have shown that 1 h of dry heating of moderately radiation-damaged zircon (FWHMs in the range 5–25 cm⁻¹) at 800–1100 °C results in major but incomplete reconstitution. Initial annealing was observed already at temperatures as low as 500 °C by these authors. Zhang et al. (2000) demonstrated also that not only the annealing temperature but also the duration of the heat treatment has strong effects on the extent of the recovery. In an annealing study on Sri Lanka gem zircon it was found that thermal treatment of moderately damaged zircon at 600 °C (4 days) resulted in moderate reconstitution (C.L. Lengauer, personal communication), and heating at 1000 °C (4 days) caused major but still incomplete recovery (Mattinson et al., 2007). In view of these previous results, our present observations

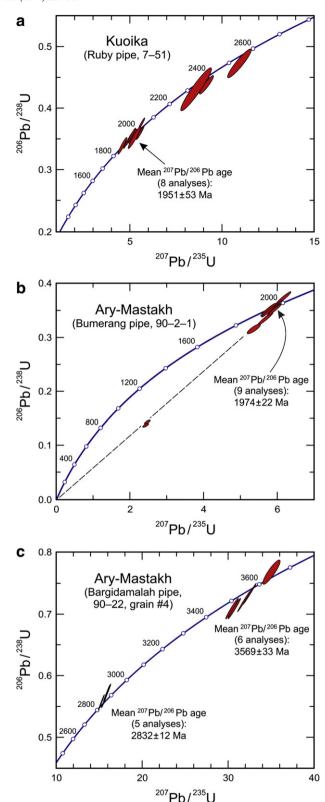


Fig. 4. Three plots showing selected results of SHRIMP U-Th-Pb analyses. Error ellipses represent 2σ uncertainties. Strong discordance of the U-Th-Pb isotopic system was found only in one case (see middle plot).

suggest that the zircon xenocrysts under investigation cannot have experienced temperatures above $600-800\,^{\circ}\text{C}$ (with the maximum temperature depending strongly on the duration of the thermal event).

Some authors have concluded that kimberlite magma may ascend with high velocities (faster than 10 m/s), especially along narrow

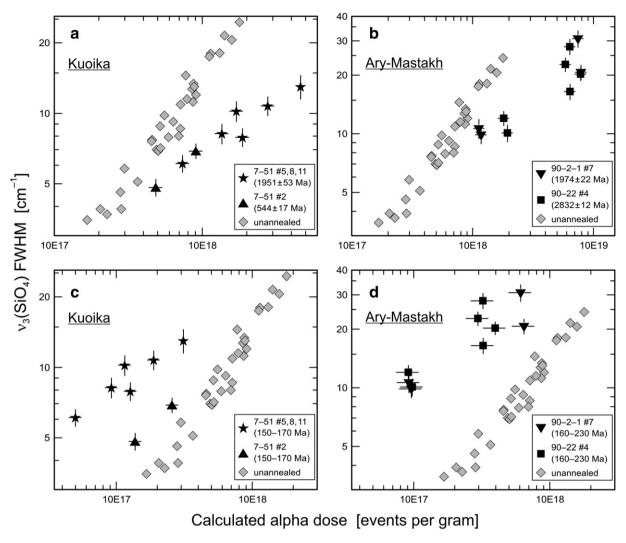


Fig. 5. Plot of the broadening of the main Raman band of zircon xenocrysts (FWHM = full width at half maximum; indicating the degree of radiation damage) against alpha doses, in comparison to unannealed zircon samples (reference data from Nasdala et al., 2001). Alpha doses for xenocrysts were calculated for assumed damage-accumulation periods since the closure of their U-Th-Pb system (a, b) and since kimberlite formation (c, d). Samples are decidedly more radiation-damaged than it would correspond to damage accumulation since kimberlite formation only.

dykes; hence rock formation would take only a few hours (Mattsson, 2012; Sparks et al., 2006; Wilson and Head, 2007). If such a time-regime should apply to the kimberlites studied here, the temperature reached by zircon crystals in the kimberlite magma cannot have exceeded ~700–800 °C. Alternatively, Kostrovitsky (1986) has postulated slow magma-ascent velocities on the order of 1 m/s, and slow subsequent cooling, leading to comparably slow kimberlite formation. Sparks et al. (2006) have proposed that kimberlite formation may extend over several days or even weeks in the case of large-diameter dykes. In such longer-term rock-formation processes, the temperature of zircon xenocrysts must have remained even well below 500–600 °C, to facilitate major damage retention.

Another noteworthy observation is that the U-rich domain of grain 90-22 #4 (Fig. 3g) consists solely of amorphous $ZrSiO_4$. Alpha doses calculated for the time since kimberlite formation $(1.4-2.4\times10^{18}~events/g; Table 2)$ would correspond merely to moderate stages of damage accumulation. Note that amorphisation requires close to 10^{19} alphadecay events/g (Murakami et al., 1991; Nasdala et al., 2001). It therefore seems reasonable to assume that the U-rich domain of grain 90-22 #4 was amorphised, or at least close to amorphous, by the time of the kimberlite ascent, without having experienced major annealing. Also, it is remarkable that neither tetragonal nor monoclinic ZrO_2 was not observed. It is well known that (dry) heating of amorphous $ZrSiO_4$ in an

intermediate temperature range of approximately 800–1150 °C for several hours results in chemical decomposition, that is, formation of a multitude of (crystalline) $\rm ZrO_2$ crystals embedded in an amorphous silica matrix (Capitani et al., 2000; Nasdala et al., 2002). This was not detected here. We may conclude again that, between being sampled by the hot, ascending kimberlite magma, and cooling near the Earth surface, zircon xenocrysts cannot have experienced temperatures in excess of 800 °C.

5. Concluding remarks

Our results suggest that major structural reconstitution of zircon in the hot kimberlite magma, through thermal annealing of previously accumulated radiation damage, can be discounted. Notable weathering or hydrothermal overprinting following the kimberlite emplacement can also be excluded. First, in spite of the occasionally high radiation damage, the majority of the zircon appears "fresh", with only a tiny volume fraction of alteration patches (Fig. 3). Second, U–Pb discordance assigned to recent loss of radiogenic Pb was only observed in one case (Fig. 4b). The complex observations made in the present study suggest that sampling by the hot kimberlite magma, kimberlite ascent, and subsequent kimberlite cooling near the Earth's surface, had very minor effects on zircon composition and physical parameters. Zircon

xenocrysts from kimberlites in the Kuoika and Ary–Mastakh fields are hence very suitable for the unbiased investigation of precursor rocks.

It is clear, however, that our results cannot be generalised. The ratio of damage retention and damage annealing experienced by zircon xenocrysts in the kimberlite magma will always depend on depth of sampling by the kimberlite magma, ascent velocity (both controlling the dwell time in the hot magma), and magma temperature. It seems worthwhile to subject future studies of zircon xenocrysts from kimberlites to a direct comparison of short-term, low-temperature effects, namely fractional structural recovery of previously accumulated radiation damage, and fractional loss of radiogenic helium.

Acknowledgements

We thank A. Wagner for sample preparation and N. Groschopf, E. Libowitzky, and M. Müller for experimental help. C.L. Lengauer and C.-Z. Liu are thanked for providing unpublished data. The SHRIMP reference zircon M257 was made available by W. Hofmeister (Institute of Gemstone Research, Idar-Oberstein and Mainz, Germany). The SHRIMP II ion microprobe used in this study is operated by a consortium consisting of Curtin University of Technology, the University of Western Australia, and the Geological Survey of Western Australia. Constructive comments by F. Corfu, an anonymous reviewer, and editor N. Eby are gratefully acknowledged. This research was supported partially by the Austrian Science Fund (FWF) through grant P24448-N19 to L.N. Co-authors S.K. and S.E. are thankful to the Russian Academy of Sciences funding through grants N 59, 151, and 27.1.

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