



Klaver, M., & Coath, C. (2019). Obtaining Accurate Isotopic Compositions with the Double Spike Technique: Practical Considerations. *Geostandards and Geoanalytical Research*, 43(1), 5-22. <https://doi.org/10.1111/ggr.12248>

Peer reviewed version

Link to published version (if available):  
[10.1111/ggr.12248](https://doi.org/10.1111/ggr.12248)

[Link to publication record in Explore Bristol Research](#)  
PDF-document

This is the author accepted manuscript (AAM). The final published version (version of record) is available online via Wiley at <https://onlinelibrary.wiley.com/doi/full/10.1111/ggr.12248> . Please refer to any applicable terms of use of the publisher.

## University of Bristol - Explore Bristol Research

### General rights

This document is made available in accordance with publisher policies. Please cite only the published version using the reference above. Full terms of use are available:  
<http://www.bristol.ac.uk/pure/about/ebr-terms>

1 **Obtaining accurate isotopic compositions with the double spike technique: practical**  
2 **considerations**

3 **Martijn Klaver\* and Christopher D. Coath**

4 Bristol Isotope Group, School of Earth Sciences, University of Bristol; Wills Memorial Building, Queen's Road,  
5 Bristol BS8 1RJ, United Kingdom

6 \*corresponding author: martijn.klaver@bristol.ac.uk; +44 (0) 117 954 5421

7

8 **ABSTRACT**

9 Ever-increasing precision in isotope ratio measurements requires a concomitant small bias and minimisation of  
10 inter-laboratory bias. The double spike technique is the most suitable method to obtain reliable isotope  
11 composition data that are accurately corrected for instrumental mass fractionation. Compared with other  
12 methods, such as sample-standard bracketing, only the double spike technique can correct for all sources of  
13 fractionation after equilibration of the sample with the double spike, such as that incurred during chemical  
14 separation and measurement. In addition, it is not dependent on *a priori* assumptions of perfect matrix  
15 matching of samples to reference materials or quantitative recovery of the sample through the chemical  
16 separation procedure to yield accurate results. In this review article, we present a detailed discussion of the  
17 merits of the double spike technique, how to design and calibrate a suitable double spike and analytical  
18 strategies. Our objective is to offer a step-by-step introduction to the use of the double spike technique in  
19 order to lower potential barriers that researchers new to the subject might face, such that double spiking will  
20 replace sample-standard bracketing as the measurement method of choice.

21 **Keywords:** double spike method; instrumental mass fractionation; mass-dependent isotope variation; spike  
22 calibration; measurement procedure optimisation

23

24 **1. Introduction**

25 Isotope geochemistry, the study of variation in the isotopic composition of elements, is a key aspect of the  
26 Earth Sciences. It has proven to be a crucial tool in establishing major concepts such as absolute dating of  
27 geological materials, the formation of the Earth in our solar system and the recycling of the Earth's crust into  
28 the mantle through plate tectonics. Traditionally, isotope geochemistry focused on isotopic anomalies arising  
29 from the decay of long-lived nuclides, such as the radiogenic  $^{87}\text{Rb}$ - $^{87}\text{Sr}$  and  $^{238}\text{U}$ - $^{206}\text{Pb}$  systems, and the mass-  
30 dependent isotopic fractionation of light elements (H, C, N, O, S) in the low-temperature geochemical realm.  
31 The latter, often but inappropriately called stable isotope fractionation, comprises fractionation of isotopes of  
32 an element through physicochemical processes as a function of the masses following a limited class of well-  
33 known functions (e.g., linear, power and exponential fractionation laws; Bigeleisen and Mayer 1947, Schauble  
34 2004, Young *et al.* 2015). Elements heavier than S were traditionally regarded not to show isotopic

35 fractionation in nature due to the small relative mass differences between their isotopes, but at present mass-  
36 dependent isotope fractionation has been detected for elements as heavy as Tl (Prytulak *et al.* 2013) and U  
37 (Andersen *et al.* 2015). The use of these “non-traditional stable isotopes” has taken flight to address a wide  
38 range of research questions across the Earth sciences. In particular, the advent of ICP-MS instruments heralded  
39 a new era of increasingly precise isotopic measurements that could, in principle, be applied to every non-  
40 monoisotopic, non-gaseous element in the periodic table. Improving the precision (repeatability) of isotopic  
41 measurements is relatively straightforward and relies for a large part on optimising counting statistics by longer  
42 measurements of higher-intensity ion beams (e.g., Albarede *et al.* 2004). The current levels of measurement  
43 precision, however, require a concomitant small bias, meaning the minimisation of (systematic) measurement  
44 errors both within and between laboratories, which has proven to be more of a challenge. A strict control on  
45 the accuracy of high-precision isotope ratio data is clearly required.

46 The double spike technique, where a tracer consisting of two artificially enriched isotopes is mixed with the  
47 sample, provides the most reliable method to obtain accurate isotopic compositions for elements with at least  
48 four isotopes (Albarède and Beard 2004, Rudge *et al.* 2009). Its fundamentals were proposed by Dodson  
49 (1963), but at first the double spike technique saw little use apart from the application to Pb (e.g., Compston  
50 and Oversby 1969, Cumming 1973, Hamelin *et al.* 1985, Galer and Abouchami 1998) and Ca isotope  
51 measurements (Russell *et al.* 1978). It took over three decades and a significant increase in measurement  
52 precision with new MC-ICP-MS instruments before the geochemical community became aware of the full  
53 potential of the double spike technique and the first doubly-spiked studies of mass-dependent fractionation of  
54 non-traditional isotopes appeared (e.g., Johnson and Beard 1999), thus vindicating Dodson’s (1963) statement  
55 that, with the double spike technique, “*it would be possible in principle to measure precisely the natural*  
56 *isotopic fractionation of a large number of polyisotopic elements, for example Mg, Ca, Ti, Cr, Ni, Fe, Mo, Sn*”.

57 Despite the clear advantages, Rudge *et al.* (2009) noted a general hesitation in the implementation of double  
58 spike measurement protocols, possibly related to the perceived difficulty of their practical use. Even though  
59 ready-to-use data reduction software and spreadsheets are available (e.g., Rudge *et al.* 2009, Creech and Paul  
60 2015), designing, making and calibrating a double spike might present a significant inertia barrier. This article  
61 aims to lower this barrier by providing a practical guide to double spiking. The solutions to the double spike  
62 equations have been presented in various forms in the literature (e.g., Dodson 1963, Russell 1971, Hamelin *et al.*  
63 *et al.* 1985, Galer 1999, Johnson and Beard 1999, Siebert *et al.* 2001, Albarede and Beard 2004). The most recent  
64 version by Rudge *et al.* (2009) is generally accepted by the geochemical community and we will follow their  
65 nomenclature and definitions throughout. As Rudge *et al.* (2009) have published the definitive approach for  
66 solving the double spike equation and its uncertainty propagation, and most researchers use publicly available  
67 data reduction programs such as the Double Spike Toolbox Matlab code (Rudge *et al.* 2009), we will not delve  
68 into a detailed discussion of the mathematics behind double spiking. Rather, we will focus on practical  
69 considerations when working with double spikes, aimed at researchers new to the subject. We aim to address  
70 questions arising when setting up a mass spectrometry protocol for a new element, including the advantages of  
71 double spiking over other corrections methods, how to choose a suitable double spike, ways to calibrate and

72 validate a double spike, etc. To aid the discussion, we will often refer to the new Ni double spike that we  
73 recently implemented at the University of Bristol as an example.

74

## 75 **2. Principles of the double spike technique**

### 76 **2.1. Instrumental mass fractionation**

77 Mass spectrometric data require extensive correction for analytical artefacts and specifically the effects of  
78 instrumentally-induced mass-dependent fractionation (IMF). Non-quantitative transmission of a sample in the  
79 mass spectrometer gives rise to mass-dependent fractionation of isotope ratios just as physicochemical  
80 processes do in nature. In TIMS, preferential evaporation of the lighter isotopes in a sample leads to relatively  
81 small but variable IMF, whereas in ICP-MS, space charge effects in the interface cause strong but relatively  
82 constant IMF (e.g., Albarède and Beard 2004). Albeit the magnitude of IMF is variable, it closely follows the  
83 exponential mass fractionation law (Russell *et al.* 1978) that, for a given pair of isotopes (denoted  $i$ ), can be  
84 written as

$$85 \text{ Eq. (1) } N_i = n_i e^{-\alpha P_i}$$

86 where  $N$  is the true isotopic ratio of a sample,  $n$  is the measured isotopic ratio of a sample,  $\alpha$  is the  
87 fractionation parameter and  $P$  is the natural logarithm of the exact mass ratio of the two isotopes (notation  
88 following (Rudge *et al.* 2009). The subscript  $i$  denotes the  $i$ th isotope ratio of the element in question, for  
89 example  $^{60}\text{Ni}/^{58}\text{Ni}$ ,  $^{61}\text{Ni}/^{58}\text{Ni}$ , etc., corresponding to  $i = 1, 2$ , etc. The only free parameter in equation (1) is the  
90 mass fractionation parameter,  $\alpha$ , whose value specifies the “degree of fractionation” with positive  $\alpha$  being an  
91 IMF favouring detection of heavy isotopes over light. Note that  $\alpha$  has no subscript and thus has the same value  
92 for all pairs of isotopes for a given element. Hence, to correct mass spectrometric data for IMF, it is necessary  
93 to determine  $\alpha$ . Accurate correction is critical as the magnitude of the IMF often exceeds natural mass-  
94 dependent variations by orders of magnitude in ICP-MS measurements. There are several ways to approach  
95 this problem, the addition of a double spike tracer being one of them. We discuss the four most common IMF  
96 correction procedures below.

97 For radiogenic isotope systems, such as Sr and Nd, internal normalisation eliminates all mass-dependent  
98 fractionation, both natural and IMF. In this case the sample is assumed to be fractionated (by a natural process)  
99 relative to some reference ratio for a non-radiogenic isotope pair to (e.g.,  $^{86}\text{Sr}/^{88}\text{Sr} = 0.1194$ ; Nier 1938). During  
100 measurement, IMF further fractionates this ratio and the combined effect is rolled into a single value of  $\alpha$ , the  
101 solution of equation (1), with  $N$  set equal to the reference ratio. Thus, the radiogenic isotope ratio (e.g.,  
102  $^{87}\text{Sr}/^{86}\text{Sr}$ ), devoid of all mass-dependent effects and leaving only the radiogenic excess or deficit, is readily  
103 determined from the measured ratio and the previously solved-for value of  $\alpha$ .

104 When a non-radiogenic isotope pair is not available, notably for Pb, or the magnitude of natural mass-  
105 dependent fractionation is investigated, an external method is required to correct for IMF. The most commonly  
106 used method is sample-standard bracketing. Here, a reference material (standard) with a known composition is

107 measured alongside the unknown samples and used to monitor IMF. A clear advantage of this approach is that  
108 the absolute composition of the reference material need not be known as long as results are expressed as the  
109 relative difference of an isotopic ratio ( $\delta$  notation; Coplen 2011). By alternately measuring the reference  
110 material and unknown samples (“bracketing”) and using the mean isotope ratios of the adjacent reference  
111 materials to correct the unknown, drift in IMF during the measurement session can be corrected for (Albarede  
112 *et al.* 2004). A critical prerequisite for sample-standard bracketing is that samples and reference materials  
113 fractionate to the same extent during measurement, which is very difficult to validate as we will discuss below.

114 A third correction method involves doping samples with an element that has at least two interference-free  
115 isotopes and a mass close to that of the element of interest to monitor IMF, such as Tl in case of Pb (e.g.,  
116 Longerich *et al.* 1987). The isotope composition of the dopant element is presumed to be known and, when  
117 one assumes that the magnitude of IMF is the same for elements of similar mass, can be used to correct the  
118 IMF of the element of interest. Whether IMF is indeed sufficiently similar for different elements is contentious  
119 and doping with another element can produce precise data that are nevertheless inaccurate (e.g., Thirlwall  
120 2002, Waight *et al.* 2002, Taylor *et al.* 2015).

121 Finally, the double spike technique uses a well-calibrated artificial tracer enriched in two isotopes (the double  
122 spike) that is added to a sample. The addition of this exotic tracer allows an internal correction to be made to  
123 obtain the isotopic composition of a sample that is corrected for all mass-dependent isotope fractionation  
124 incurred after homogenisation with the tracer.

## 125 **2.2. The double spike equation**

126 The principles of the double spike technique are shown schematically in Figure 1. In four-isotope space,  
127 measurement of a natural sample  $N$  yields three independent isotope ratios that are affected by IMF so that  
128 the measured composition ( $n$ ) is offset from the true value. The form of the IMF line is only dependent on the  
129 mass of the isotopes as given by equation (1) and is therefore known, but the displacement along the line,  
130 given by the mass fractionation parameter  $\alpha$ , is not. Hence, the true composition  $N$  cannot be constrained  
131 other than to lie somewhere on the IMF line that passes through  $n$ . If an aliquot of the natural sample  $N$  is  
132 mixed with the double spike tracer ( $T$ ), a mixture  $M$  will result with a composition that is dependent on the  
133 proportion (given by  $\lambda$ ) in which the sample and double spike are mixed. The measurement of this mixture ( $m$ )  
134 will also be affected by IMF, again along a line of known form but with unknown displacement ( $\beta$ ). It is now  
135 possible to construct one, and only one, mixing line that passes through the point  $T$  and intersects with the two  
136 IMF lines. The two intersection points are  $N$  and  $M$ ; thus a solution for the sample composition ( $N$ ) is found.  
137 The procedure can be described in algebraic terms as follows: the line  $N$ - $n$  is given by equation (1);  $M$ - $m$  and  
138 the mixing line  $N$ - $M$ - $T$  can be expressed as

139 Eq. (2)  $M_i = m_i e^{-\beta P_i}$

140 Eq. (3)  $M_i = \lambda T_i + (1 - \lambda) N_i$

141 Note that a linear mixing line, as described above, requires that the same denominator isotope is chosen for  
142 each of the three isotope ratios. Substituting equations (1) and (2) into equation (3) gives the exponential law  
143 double spike equation (Rudge *et al.* 2009):

144 Eq. (4)  $\lambda T_i + (1 - \lambda)n_i e^{-\alpha P_i} - m_i e^{-\beta P_i} = 0$

145 Since the composition of the double spike is known and the ratios of  $n$  and  $m$  are measured, equation (4) has  
146 three unknowns:  $\alpha$ ,  $\beta$  and  $\lambda$ . Writing out equation (4) for  $i = 1, 2$  and  $3$  provides three equations that can be  
147 solved numerically for the three unknowns (see Rudge *et al.* 2009), a procedure often referred to as inversion.  
148 Herein lies the main limitation of the double spike method: to solve the equations for three unknowns the  
149 isotope system requires at least four available isotopes (three ratios). We will refer to the four isotopes used in  
150 the double spike equation as the inversion isotopes. The true composition of the sample  $N$  can then be found  
151 by entering the value of  $\alpha$ , obtained from the inversion, and the measured ratios for  $n$  in equation (1). As such,  
152 the isotope composition of a sample can be obtained through the measurement of a pure ( $n$ ) and spiked ( $m$ )  
153 aliquot of the sample. This is used in the case of Pb isotopes where  $^{206}\text{Pb}/^{204}\text{Pb}$ ,  $^{207}\text{Pb}/^{204}\text{Pb}$  and  $^{208}\text{Pb}/^{204}\text{Pb}$  are  
154 used as geochemical tracers (e.g., Thirlwall 2000, Klaver *et al.* 2016). Note that, although four isotopes are  
155 required for the double spike inversion, all isotopes of an element with five or more isotopes can be measured  
156 and corrected for IMF following Equation (1) so that the double spike technique can be used to obtain the  
157 absolute isotopic composition of an element.

### 158 2.3. Double spiking applied to natural mass-dependent fractionation

159 When one is interested in the natural mass-dependent (“stable”) isotope fractionation relative to a reference  
160 material, a simpler approach can be followed. Assuming IMF and natural fractionation both follow the  
161 exponential fractionation law, it follows that the isotopic composition of the sample and reference material  
162 satisfy equation (1) (i.e., they lie on the same mass-dependent fractionation line). In this case, the isotopic  
163 composition of the reference material in the double spike inversion (Eq. 4) can take the places of the measured  
164 ratios of the unspiked sample ( $n$ ), thus obviating the need to measure the unspiked sample; a single  
165 measurement of the spiked sample ( $m$ ) suffices to calculate  $\alpha$  relative to the reference material. The value for  $\alpha$   
166 resulting from the double spike inversion represents the degree of natural fractionation (the displacement  
167 from reference material  $n$  to sample  $N$ ), which can be directly expressed in the  $\delta$  notation. Following the IUPAC  
168 convention (Coplen 2011) and taking Ni as an example:

169 Eq. (5)  $\delta^i \text{Ni} = e^{-\alpha P_i} - 1$

170 where  $i$  again indicates a specific pair of isotopes of the element (e.g.,  $^{60}\text{Ni}/^{58}\text{Ni}$ ) and  $P$  is the natural logarithm  
171 of the ratio of the exact isotopic masses. Note that  $\alpha$  is the quantity of interest and that the choice of isotope  
172 ratio in the  $\delta$  notation is completely arbitrary; there is no advantage in choosing one ratio over another other  
173 than convention. Consequently, for double spike data, consistency between, for example,  $\delta^{60/58}\text{Ni}$  and  $\delta^{62/58}\text{Ni}$   
174 cannot be used as an independent measure of data quality. This presents a marked contrast with other  
175 correction methods, such as sample-standard bracketing, where each isotope ratio is independently corrected  
176 for IMF through bracketing against the reference material.

## 177 **2.4. Two special cases**

178 The double spike technique is normally only applicable to elements with at least four stable isotopes but can be  
179 extended to 3-isotope elements using an ingenious adaptation (Hofmann 1971, Coath *et al.* 2017). This critical-  
180 mixture method employs the fact that  $\beta$  is an output parameter of the double spike equation but has little  
181 practical use. Thus, if  $\beta$  can be constrained through independent means, the number of unknowns decreases by  
182 one and only two independent isotope ratios are required to solve equation (4). Typically, any bias in  
183 estimating  $\beta$ , for instance through sample-standard bracketing, propagates directly into a bias in  $\alpha$  and its  
184 associated uncertainty, the parameter of interest. In the special case where the IMF line  $M$ - $m$  is tangential to  
185 the sample – double spike mixing line, however, a bias in estimating  $\beta$  has a marginal influence on  $\alpha$  as the  
186 displacement of the presumed position of  $M$  is parallel to the mixing line. The IMF vector and mixing line only  
187 become parallel when sample and double spiked are precisely combined in a “critical mixture”, typically  
188 requiring an iterative effort to optimise the sample – double spike proportion (Coath *et al.* 2017).

189 An alternative option to extend the double spike technique is the use of one or two non-naturally occurring  
190 radioactive isotopes to make a double spike. This is commonly used for high-precision U-Pb and U-series  
191 measurements of small samples where well-calibrated  $^{202}\text{Pb}$ - $^{205}\text{Pb}$  and  $^{233}\text{U}$ - $^{236}\text{U}$  spikes allow direct correction  
192 for IMF as these isotopes do not naturally occur in samples (e.g., Todt *et al.* 1996). The main problem of  
193 upscaling this technique to mass-dependent fractionation studies of a greater variety of elements is the use of  
194 large quantities of costly radioactive spike material and the consequent radiological hazard. To our knowledge  
195 this has not yet been undertaken for elements other than U.

196

## 197 **3. The advantages of double spiking**

### 198 **3.1. Matrix effects and variable mass fractionation**

199 The main advantage of double spiking is that it provides accurate correction for IMF as long as the double spike  
200 is properly calibrated. It does not rely on an *a priori* assumption of identical IMF behaviour of samples and  
201 reference material – in clear contrast to sample-standard bracketing (SSB). To illustrate this point, we present a  
202 comparison between double spike and SSB-corrected  $\delta^{60/58}\text{Ni}$  data, obtained by MC-ICP-MS (see Appendix A),  
203 for the same measurements of a pure Ni solution in Figure 2. Because the double spike inversion yields a value  
204 for  $\beta$  (the magnitude of the IMF) for every measurement, double spike corrected data can be re-corrected  
205 assuming constant IMF between samples and the reference material as in SSB. Repeated double spike  
206 measurements of a pure Ni solution (our in-house CPI Ni solution) yield results with a high intermediate  
207 measurement precision. Correcting the same data through SSB leads to more scatter of the data points and  
208 introduces a bias compared to the double spike results. This is particularly evident in the case of CPI standards  
209 processed through the chemical separation procedure: whereas double spike-corrected results are  
210 indistinguishable from unprocessed CPI solutions, SSB-corrected data can yield over 1 ‰ lower  $\delta^{60/58}\text{Ni}$  values.  
211 This is a significant effect considering that natural variations in Ni isotopes are often <0.1 ‰. Clearly, the  
212 requirement of identical IMF of the reference material and processed CPI standards is not met. In this case, the

213 likely cause is residual organic material (from the dimethylglyoxime used to elute Ni) in the processed Ni  
214 solutions. The presence of a small amount of matrix (any element or compound other than the one of interest)  
215 in samples can affect plasma conditions to such an extent that the IMF behaviour becomes significantly  
216 different compared to the matrix-free reference material. Such matrix effects are very common and very hard  
217 to systematically suppress (e.g., Barling and Weis 2008, Tipper *et al.* 2008, Shiel *et al.* 2009, van den Boorn *et*  
218 *al.* 2009, Nan *et al.* 2015, Peters *et al.* 2015), making it extremely difficult to obtain bias-free isotope data  
219 through sample-standard bracketing. Strikingly, the repeatability measurement precision (the 2 SE of the mean  
220 of 6-8 independent measurement results) is only a factor ~3 worse than for the DS-corrected data, indicating  
221 that SSB can induce a significant measurement bias yet yield relatively precise results, with the high precision  
222 potentially giving a false sense of security. The double spike method, on the other hand, has no trouble dealing  
223 with these matrix effects and produces accurate results.

### 224 **3.2. Correcting for fractionation during chemical purification**

225 As matrix effects can have a pronounced influence on IMF, care has to be taken to matrix-match samples to the  
226 elementally pure reference material, which typically involves the elimination of the sample matrix as much as  
227 possible. Purification of geological materials is mostly achieved through a single- or multistage ion-exchange  
228 chromatography (“column”) procedure (e.g., Schönbacher and Fehr 2013). Although sample purification is  
229 essential, it can have negative side effects such as the introduction of organic species as in the case of Ni  
230 (Figure 2). A second major complication is that isotopic fractionation can occur during ion-exchange  
231 chromatography, implying that non-quantitative recovery of the analyte can lead to mass-dependent  
232 fractionation. This is particularly pertinent to lighter elements that are more readily fractionated due to the  
233 larger relative mass difference between their isotopes (e.g., Russell and Papanastassiou 1978, Oi *et al.* 1991,  
234 Chernonozhkin *et al.* 2015), but with increasing precision of isotope measurement results it should not be  
235 neglected for heavier elements either. Quantitative recovery of the analyte is thus required to avoid  
236 introducing a bias but is difficult to demonstrate on a sample-to-sample basis and any fractionation incurred  
237 during chemical purification cannot be corrected for using sample-standard bracketing. On the other hand, the  
238 double spike technique corrects for any mass-dependent fractionation after sample-spike equilibration. Hence,  
239 if a sample and double spike are mixed and equilibrated prior to ion-exchange purification, the double spike  
240 inversion will correct for the combined mass-dependent fractionation during purification and measurement. As  
241 such, the double spike technique is the only method that is not dependent on quantitative recovery through  
242 the separation procedure and can produce accurate results even if a significant fraction of the sample is lost  
243 during purification (e.g., <50 % yield).

244

### 245 **4. Designing a double spike**

246 The aim of the double spike technique is to produce isotope data with a small uncertainty. The accuracy of  
247 double spike data relies on the quality of the calibration of the double spike and reference material, which is  
248 discussed below, but the precision of the results depends greatly on the choice of double spike composition.



249 Obviously, one would prefer a double spike that produces the best possible precision, either in a ratio of  
250 interest (e.g.,  $^{207}\text{Pb}/^{204}\text{Pb}$ ) or in  $\alpha$  when natural mass-dependent fractionation is investigated. Rudge *et al.*  
251 (2009) included a model in their Double Spike Toolbox that can be used to calculate the theoretical precision in  
252  $\alpha$  or an isotope ratio for a wide range of variable parameters. Briefly, their precision model is based on linear  
253 propagation of the variance associated with the measurement of ion beam intensities: i) thermal noise of the  
254 resistors fitted in the amplifier feedback loop of the Faraday detectors (Johnson-Nyquist noise) and ii) counting  
255 statistics of the ion beam intensities (Poisson or shot noise; see e.g., Albaredo *et al.* 2004). If the double spike is  
256 calibrated perfectly, the double spike inversion does not introduce a bias and only leads to magnification of the  
257 measurement precision. As such, the precision model provides a theoretical limit on the highest precision that  
258 can be achieved for a given total ion beam intensity and integration time.

259 A double spike that leads to minimal magnification of the measurement precision is considered optimal and a  
260 “cocktail list” of double spikes that yield the highest theoretical precision is provided by Rudge *et al.* (2009). We  
261 strongly prefer their approach of linear propagation of variances to assess the performance of a double spike  
262 over geometrical methods proposed by e.g., Galer (1999) and Johnson and Beard (1999). The latter predict that  
263 the precision of a double spike solution is dependent on the choice of denominator isotope in the ratios used in  
264 the double spike inversion. Because mixing and mass fractionation processes are coordinate-independent,  
265 there is no physical basis for the choice of denominator isotope to have any effect, as discussed by Rudge *et al.*  
266 (2009).

267 There are several factors that can influence the choice of the isotopes used in a double spike and, in the case of  
268 elements with five or more isotopes, the four isotopes used in the double spike inversion. Selecting the four  
269 inversion isotopes should be the first step in choosing an appropriate double spike: the same double spike can  
270 perform better or worse depending on the choice of inversion isotopes. The presence of isobaric interferences  
271 can make some isotopes less suitable as an inversion isotope. In the case of Ni, five stable isotopes are  
272 available:  $^{58}\text{Ni}$  (68.08 % abundance),  $^{60}\text{Ni}$  (26.22 %),  $^{61}\text{Ni}$  (1.14 %),  $^{62}\text{Ni}$  (3.63 %) and  $^{64}\text{Ni}$  (0.93 %; Gramlich *et al.*  
273 1989). We decided to avoid the use of  $^{64}\text{Ni}$  as it suffers from isobaric interference by a major isotope of Zn ( $^{64}\text{Zn}$   
274 has an abundance of 48.6 %; Rosman 1972), even though double spikes with  $^{64}\text{Ni}$  yield the best possible  
275 precision. It is difficult to quantitatively remove Zn from samples as it is a common contaminant in labware and  
276 the environment. Correction for Zn interference is possible but undesirable as monitoring  $^{66}\text{Zn}$  would require a  
277 magnet jump in the method and thus consume time and analyte. Hence, we have opted to exclude  $^{64}\text{Ni}$  and use  
278  $^{58}\text{Ni}$ ,  $^{60}\text{Ni}$ ,  $^{61}\text{Ni}$  and  $^{62}\text{Ni}$  as inversion isotopes. Isobaric interferences have also affected the choice of, for  
279 instance, Mo (Siebert *et al.* 2001) Cd (Ripperger and Rehkämper 2007) and Pt (Creech *et al.* 2013) inversion  
280 isotopes. In addition to isobaric interferences, the presence of nucleosynthetic anomalies in extra-terrestrial  
281 materials can also be taken into account when choosing the four inversion isotopes (e.g., Millet and Dauphas  
282 2014).

283 Upon choosing the inversion isotopes, the precision model of Rudge *et al.* (2009) can be used to explore a wide  
284 parameter space to find the most suitable double spike. Simply choosing the one with highest theoretical  
285 precision might not always be the best choice. Taking Ni as an example again, we used the precision model of

286 Rudge *et al.* (2009) to generate contour diagrams of the precision in  $\alpha$  as a function of double spike  
287 composition and the proportion of spike in the sample-spike mixture for two different double spikes (Figure 3).  
288 Excluding  $^{64}\text{Ni}$ , a  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike yields the best possible precision, albeit only marginally better than a  
289  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  spike. The performance of the  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  spike, however, is much more dependent on getting the  
290 sample-spike proportion exactly right: a small deviation from the optimal sample-spike proportion leads to a  
291 significant decrease in precision. If the double spike deviates from the optimal composition, this becomes even  
292 more pertinent. On the other hand, the  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  spike has a broad precision minimum with very little  
293 deterioration in precision over a wide range in the sample-spike ratio. Versatility of a double spike in terms of  
294 sample-spike proportion is desirable, in particular when the concentration of the element in the double spike  
295 or, more likely, the sample is poorly known. As such, we preferred the use of a  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  double spike.

296 The default precision model of the Double Spike Toolbox employs a total beam intensity for the four inversion  
297 isotopes of 100 pA (10 V), integration time of 8 s and a  $10^{11} \Omega$  amplifier feedback resistor. These measurement  
298 conditions are typical for MC-ICP-MS and TIMS measurements, but if the double spike method will be applied  
299 to much larger or smaller sample sizes it is worth investigating the effect of total beam intensity on the  
300 expected precision. Figure 4 shows the precision in  $\alpha$  as a function of sample-spike proportion for two Ni  
301 double spikes at a total beam intensity that varies by three orders of magnitude. The broad precision minimum  
302 of the  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  double spike, as seen in Figure 3, is maintained down to very low (0.1 V) beam intensities. In  
303 contrast, the deterioration in the precision when moving away from the optimal sample-spike proportion is  
304 more acute for the  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike although the best possible precision is always higher than for the  $^{61}\text{Ni}$ -  
305  $^{62}\text{Ni}$  double spike. Thus, use of a  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike would require more stringent control on mixing of  
306 sample and double spike in the optimal proportion, in particular for small sample sizes. If only a small amount  
307 of sample is available, one might not want to sacrifice an aliquot for concentration measurement prior to  
308 spiking and hence one would have to rely on published or estimated concentration data for the element of  
309 interest. In such a case, a versatile double spike with a wide precision minimum is preferable as it allows for  
310 significant deviation from the optimal sample-spike proportion without suffering a large decrease in precision.  
311 Another argument for choosing the  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  spike is that it boosts the intensity of the minor  $^{61}\text{Ni}$  isotope (1.14  
312 % abundance). Although the precision model of Rudge *et al.* (2009) incorporates the Johnson-Nyquist noise  
313 that is the dominant source of predictable variance at low beam intensities, it does not take into account  
314 unpredictable, non-proportional effects, such as the variable effects of background and blank corrections on  
315 minor isotopes, which diminish at higher beam intensities.

316 A final consideration is the required purity of the single spikes and their cost. Optimal double spikes commonly  
317 include at least one minor (<5 % abundance) isotope of the element (Rudge *et al.* 2009) and the cost of single  
318 spikes is typically higher for isotopes with a low abundance in nature and increases with the enrichment factor  
319 of an isotope. Although the double spike technique does not critically rely on high purity single spikes as do  
320 isotope dilution methods (e.g., Stracke *et al.* 2014), the enrichment factor of the single spikes does have some  
321 influence on the versatility of a double spike and the precision that can be obtained. If reliable IMF correction is  
322 required but not the highest possible precision, it might be cost-efficient to opt for less-enriched single spikes

323 of isotopes with a reasonably high natural abundance. Moreover, a double spike with a low proportion of spike  
324 in the optimal sample-spike mixture allows more measurements for a given quantity of double spike. The  
325 Double Spike Toolbox precision model allows the spike compositions to be varied so that the effects of single  
326 spike purity on the performance of double spikes can be investigated for a specific application.

327

## 328 **5. Calibrating a double spike**

### 329 **5.1. Calibration relative to a reference material**

330 The accuracy of data corrected for IMF using the double spike technique relies heavily on how well the  
331 composition of the double spike is known, although reporting data relative to a reference material that was  
332 spiked and measured together with the samples can correct a potential bias. Careful gravimetric mixing  
333 provides insufficient accuracy even in the rare case where the isotope composition of the single spikes is well  
334 known. The compositional certificate of single spikes provided by their distributor gives an indication only and  
335 cannot be trusted to be sufficiently accurate. Hence, proper calibration of the double spike composition is  
336 required. As it is extremely difficult to obtain an absolute isotope composition, double spikes are typically  
337 calibrated relative to a reference material: the  $\delta = 0$  reference material for mass-dependent fractionation  
338 studies. For Pb double spikes, SRM 982 is commonly used as calibrant, leaving SRM 981 as an independent  
339 check of the quality of Pb isotope data (e.g., Thirlwall 2000, Klaver *et al.* 2016).

340 The calibration of a double spike relative to a reference material is shown in Figure 5; a schematic flowchart  
341 with a step-by-step approach is provided in Figure 6. As the calibration of the double spike is made relative to  
342 the reference material, the absolute isotope composition of the latter does not need to be known and it is  
343 sufficient to constrain the mass fractionation line on which the reference material lies (Rudge *et al.* 2009). The  
344 first step of the double spike calibration thus constitutes a precise measurement of the (unspiked) reference  
345 material. A first-order IMF correction can be applied by either an internal normalisation to a given ratio (e.g., Ni  
346 reference material SRM 986 corrected to  $^{61}\text{Ni}/^{58}\text{Ni} = 0.016744$ ; Gramlich *et al.* 1989) or by doping with another  
347 element for an external estimate of the IMF. This ensures that the corrected composition of the reference  
348 material is close to its true value and only offset along the mass fractionation line, which does not introduce a  
349 bias in the double spike inversion (see Rudge *et al.* 2009 for a proof).

350 The next step is calibrating the composition of the double spike itself, which can be addressed in multiple ways  
351 but generally requires a measurement of the pure double spike. One can apply an external IMF correction to  
352 the pure double spike measurement by doping with another element (e.g., Pd for Mo or Cu for Ni; Siebert *et al.*  
353 2001, Gall *et al.* 2012). The same caveat as for samples (section 2.1) applies here: an external estimate of IMF  
354 critically relies on the assumption of identical IMF behaviour between double spike and reference materials. In  
355 the absence of a geological matrix in these materials this could be justifiable and a small difference in IMF will  
356 not directly have a disastrous effect on the accuracy of double spike data, as discussed below. The second and  
357 our preferred approach entails the construction of the mixing line  $N$ - $T$  between the double spike and a  
358 reference material through the measurement of one or more mixtures of the two ( $M_k$ ; Figure 5). In theory, only

359 a single mixture  $M_k$  is required (in addition to the measurement of the pure double spike) to solve Equation 5  
360 for the composition of the double spike: simply substitute the composition of the reference material for the  
361 double spike (i.e., use the reference material as a “quadruple spike”) and treat the pure double spike  
362 measurement as the unknown sample. A more accurate result can be obtained by measuring multiple mixtures  
363 with different proportions of sample to double spike. These mixtures serve to define the  $N$ - $T$  mixing line in  
364 four-isotope space; the composition of the double spike lies at the intersection of this mixing line and the mass  
365 fractionation line defined by the measurement of the pure double spike (Figure 5). Two mixtures are required  
366 to constrain the direction of the  $N$ - $T$  mixing line; the measurement of more mixtures leaves the system  
367 overdetermined and a least-squares approach can be used improve the precision with which the direction of  
368 the mixing line is determined (Rudge *et al.* 2009). Note that, in isotope systems where the total range of  
369 geological variation in mass-dependent isotope fractionation is small, the accuracy of the double spike  
370 inversion is largely dependent on how well the direction of the  $N$ - $T$  mixing line is known, rather than the  
371 location of the double spike composition on this line (Rudge *et al.* 2009).

## 372 **5.2. Validation of the calibration and possible biases**

373 Upon calibration of the double spike, it is important to validate the quality of the calibration. Any bias in the  
374 double spike calibration should become apparent when a series of mixtures of a quality control material  
375 (“secondary standard”) is measured at various sample-spike proportions (Figure 6). If the calibration is  
376 accurate, the quality control material mixtures should yield the correct value and not show any systematic  
377 variation with sample-spike proportion. A constant offset or deviations at high or low double spike proportions  
378 in the mixtures flag an issue with the double spike calibration. We will explore the effects of two common  
379 sources of uncertainty in the double spike calibration: the effect of blank correction on the pure double spike  
380 and reference material measurements and variable IMF between reference material and double spike when  
381 applying an external IMF correction. Again, we will use our Ni double spike as an example. This double spike  
382 was calibrated relative to reference material SRM 986 using multiple SRM 986-double spike mixtures as  
383 described above. The accuracy of the calibration was checked by measuring multiple spiked aliquots of a  
384 quality control material (our in-house CPI Ni solution); the results are shown in Figure 7.

385 Systematic errors in the measurement of the pure double spike are the most likely source of inaccuracies in the  
386 double spike calibration. Due to its extremely non-natural isotopic composition and generally low abundance of  
387 all but the spike isotopes, correction for the contribution of instrumental background has a pronounced effect  
388 on the measured isotopic composition of the double spike. Making blank corrections without introducing a bias  
389 can be problematic for isotopic measurement by MC-ICP-MS. Blanks corrections are usually made by  
390 subtracting the on-peak intensity measured in the pure solvent, the same as that used for dissolution and  
391 dilution of the analyte. The intensities so measured will include: i) trace levels of the analyte and other  
392 elements present in the solvent, ii) plasma-generated molecular species and iii) instrument “memory” of  
393 previously analysed solutions. In particular, memory effects (iii) are of special concern in the case of double  
394 spike measurements because of the variable isotopic composition of the double-spike sample mixtures in  
395 combination with a mass-fractionation of the instrumental memory, which is often distinct from the IMF of the

396 analyte. We have modelled, therefore, the effect of over- and under-correcting a) the double-spike  
397 composition with an IMF-affected reference material composition and b) the reference material isotopic  
398 composition with an IMF-affected double-spike composition. Figure 7 shows that an increasingly inaccurate  
399 double-spike inversion results as the proportion of double-spike in the mixture increases and decreases for  
400 cases (a) and (b) respectively. The sign and magnitude of the bias are a function of the composition of this  
401 blank and the error in correcting for it, but the overall form of the bias, as shown in Figure 7, is independent of  
402 these parameters. The magnitude of the bias at a given sample-spike proportion is constant and does not scale  
403 with  $\alpha$ , meaning that samples spiked in exactly the same proportion will yield results that are correct relative to  
404 one another, but will be wrong in an absolute sense; any difference in sample-spike proportion will lead to bias  
405 between samples.

406 For a double-spike calibrated by direct measurement of the pure double-spike and corrected for IMF by  
407 sample-standard bracketing, or by internal normalisation by addition of another element, such as Cu in the  
408 case of Ni (Gall *et al.* 2012) or Pd for a Mo double spike (Siebert *et al.* 2001), it is likely that the IMF correction  
409 will not be accurate. From the mixing equation (4) it is easily shown that the effect of this inaccuracy on the  
410 calibration is to offset the apparent sample fractionation  $\alpha$  by the same magnitude but with the opposite sign.  
411 This effect is independent of the proportion of double-spike in the mixture: all samples and all mixtures are  
412 offset the same. Therefore, an external correction may be applied, based on measurements of double-spiked  
413 reference materials, or, equivalently, the double-spike calibration can be adjusted.

414 The measured Ni CPI data in Figure 7 are identical within uncertainty of the mean for sample-spike proportions  
415 between  $\sim 0.15$  and  $\sim 0.75$ . The measurement at  $p \sim 0.8$  is marginally higher, potentially reflecting a small blank  
416 correction-induced inaccuracy in the double spike calibration, but these results give confidence that samples  
417 measured at sample-spike proportions  $< 0.7$  should yield results with a small uncertainty. Demonstrating  
418 constant results for a quality control material over a wide range in sample-spike proportions is vital in  
419 validating a double spike method.

### 420 **5.3. Improving a double spike calibration?**

421 From the discussion above it is evident that the measurement of the pure double spike is likely the weak spot  
422 of the calibration. Blank correction has a significant effect on the measured composition and it is questionable  
423 to what degree a blank or background correction can be bias-free. In particular for ICP-MS instruments,  
424 memory effects can be pronounced and such a background can have a distinctly non-natural composition due  
425 to complex mass fractionation effects (e.g., Albarede *et al.* 2004). Bias introduced by the blank correction on  
426 the sample-spike mixtures are less of an issue because it induces a translation largely along the sample-spike  
427 mixing line, thus hardly influencing the obtained direction of the mixing line. A way to circumvent the need to  
428 measure the pure double spike involves a second mixing line between the double spike and a secondary  
429 standard with a very different isotopic composition to the primary reference material (Rudge *et al.* 2009). This  
430 is illustrated in Figure 5. Multiple mixtures  $EM_k$  between the secondary, exotic standard  $E$  and the double spike  
431 will allow the construction of the secondary mixing line; the double spike composition lies at the intersection  
432 between the mixing lines  $N-T$  and  $E-T$ . This approach requires the use of a secondary standard with a

433 sufficiently exotic composition, potentially by doping with an isotopic spike that is not used in the double spike.  
434 Its composition need not be known as long as at least four mixtures (including an unspiked exotic standard) are  
435 measured. In theory, this method might yield a more precise double spike calibration as it does not rely on  
436 accurately blank-correcting a pure double spike measurement. Practical impedances include the availability of  
437 a suitable exotic secondary standard, introducing such a material in the mass spectrometer in the light of  
438 possible memory effects and the complicated mathematics required to solve for the intersection of the two  
439 mixing lines.

440

## 441 **6. Double spike measurement strategies**

### 442 **6.1. General considerations**

443 As long as the isotopic compositions of the double spike and reference material are properly calibrated and any  
444 spectral interferences are resolved or corrected for, the double spike method should yield data with a low  
445 uncertainty. In practice, however, there might be significant long- and short-term drift caused by external  
446 factors such as Faraday cup degradation, magnet instability and tailing from interfering species, clipping of the  
447 ion beam and other mass-independent effects (e.g., Albarede *et al.* 2004, Carlson 2014). These effects can be  
448 transient in time but, particularly in ICP-MS instruments, surprisingly straightforward to address by adopting  
449 sample-standard bracketing as a secondary correction. Bracketing spiked samples with the spiked reference  
450 material typically improves the reproducibility of data and counteracts long- and short-term drift: the double  
451 spike method accurately corrects for mass fractionation while bracketing resolves additional analytical  
452 artefacts. In addition, bracketing with a reference material spiked in a similar fashion as the samples will  
453 eliminate some of the potential biases introduced by inaccuracies in the double spike calibration and is hence  
454 recommended. The precision of double spike data can be enhanced by improving counting statistics by  
455 measuring more intense ion beams (higher analyte concentration) for a longer time. In order to offer the best  
456 correction for transient instrument drift, it might be preferable to measure the same sample multiple times  
457 bracketed by the spiked reference material, potentially in different measurement sessions, and pool the data of  
458 these measurements. In this case, there is a trade-off between analyte and time loss during sample uptake and  
459 washout and the improvement in precision.

460 There is no time-penalty in the use of double spiking as a correction method for IMF in the case of mass-  
461 dependent isotope fractionation studies. Once the double spike is calibrated, total integration time per sample  
462 need not be different to other correction methods (sample-standard bracketing, doping). Hence, the number of  
463 samples that can be measured per day is the same if bracketing is adopted as a secondary correction; if such a  
464 secondary correction is not deemed necessary, more samples can be measured per session by double spike  
465 than by sample-standard bracketing. In addition, the use of a double spike obviates the need to verify column  
466 yields and absence of matrix in the analyte on a sample-by-sample basis, which are clear prerequisites for  
467 sample-standard bracketing, and as such presents a significant decrease in analytical time and effort.

468 Similar to single spikes used in isotope dilution techniques, a double spike measurement can be used to obtain  
469 the concentration of the element of interest in a sample. As the sample – double spike mixing parameter ( $\lambda$ ) is  
470 computed as part of the double spike inversion, it can be used to calculate the amount of the natural element  
471 present in the mixture; see Rudge *et al.* (2009) regarding how to convert  $\lambda$  to a molar proportion. This requires  
472 that the double spike is calibrated for concentration and careful weighing of the sample and spike contributions  
473 to the mixture; the former is easily achieved using a gravimetric solution of the reference material for the  
474 double spike calibration.

## 475 **6.2. Data reduction**

476 Solutions to the double spike equation have been presented in various forms in the literature. The earliest  
477 forms are based on a linear mass fractionation law (e.g., Dodson 1963, Hamelin *et al.* 1985) for which the  
478 double spike equations can be solved algebraically. Natural and instrumental mass fractionation, however,  
479 most closely follow an exponential mass fractionation law (Russell *et al.* 1978). A geometric approach to  
480 approximate an exponential law solution was commonly used (e.g., Johnson and Beard 1999) until Albarède  
481 and Beard (2004) and Rudge *et al.* (2009) promoted numerical methods to solve the exponential law double  
482 spike equation (Equation 4). Rudge *et al.* (2009) presented a detailed guide on how to numerically solve  
483 Equation 4 using matrix algebra and a Newton-Raphson iteration, which, following their equations, is easily  
484 implemented into a spreadsheet program. Alternatively, the Double Spike Toolbox Matlab code or the Iolite  
485 add-in IsoSpike (Creech and Paul 2015) can be used for data reduction.

486 A mass spectrometric measurement is typically divided into multiple cycles with a discrete integration time  
487 (e.g., 50 cycles of 4.2 s integration time yielding a total measurement duration of 210 s). This means that there  
488 are two ways to approach data reduction: the double spike inversion can be carried out on a cycle-by-cycle  
489 basis after which an average can be calculated for the results of the individual cycles, or the average of the  
490 measured cycles can be used in the inversion. In other words, the difference lies in taking an average before or  
491 after performing the double spike inversion. If the input data are identical (i.e., no cycles are rejected at any  
492 stage), the results of the two approaches should be almost identical and thus we assert no preference for  
493 either one. Treating the data on a cycle-by-cycle basis has the advantage that it provides an easy way to assess  
494 the uncertainty of a single measurement by taking the standard error of the results for all cycles (Creech and  
495 Paul 2015) rather than having to rely on complex linear uncertainty propagation (formulated by Rudge *et al.*  
496 2009) or Monte Carlo approaches.

497 In the case of our Ni method, we perform all data reduction in an automated offline spreadsheet. Data are  
498 corrected for isobaric interference of  $^{58}\text{Fe}$  on a cycle-by-cycle basis after which the corrected ratios are  
499 screened for outliers (>4 interquartile ranges from the median), which typically originate from memory effects of  
500 the introduction system (see Appendix A). We then use the average ratios for the double spike inversion and  
501 apply no uncertainty propagation or estimate for individual measurements. Rather, we analyse every sample 6-  
502 10 times, bracketed by the spiked SRM 986 reference material, in at least two different measurement sessions  
503 and report the average and standard error for these repeat measurements.

### 504 6.3. Mass-independent anomalies and double spiking

505 So far, we have largely focused on the use of the double spike technique to derive the magnitude of natural  
506 mass-dependent isotope fractionation relative to a reference material. We reiterate that for this purpose a  
507 single measurement of the spiked sample is required; the calibrated reference material is used as the unspiked  
508 composition in the double spike inversion. The fundamental assumption of this approach is that the sample  
509 and reference material are distinct through mass-dependent fractionation *only* (i.e., they lie on the same mass  
510 fractionation line). The presence of mass-independent anomalies, i.e. any isotopic effects that do not obey the  
511 assumed fractionation law, will cause a bias in the calculated value of  $\alpha$  unless explicitly corrected for (e.g., Hu  
512 and Dauphas 2017). In practically all terrestrial applications this condition is met, and a single measurement of  
513 the spiked sample is required to solve the double spike equation. Mass-independent anomalies can, however,  
514 be present in various forms and particularly so in extra-terrestrial samples (e.g., Dauphas and Schauble 2016).  
515 Regardless of the origin of the mass-independent anomalies, however, measuring both a spiked and an  
516 unspiked aliquot will always yield the true isotopic composition of a sample. It must be noted that because the  
517 double spike is calibrated relative to a reference material, the true isotopic composition so measured is known  
518 only as a relative difference from the reference material, but without the restriction of a purely mass-  
519 dependent relationship between the two.

520 The most commonly encountered type of mass-independent anomaly results from the radiogenic ingrowth due  
521 to the decay of a long- or a short-lived radioactive parent and affects, for instance,  $^{40}\text{Ca}$ ,  $^{87}\text{Sr}$  and  $^{182}\text{W}$ . The  
522 magnitude of these radiogenic anomalies can be orders of magnitude larger than natural mass-dependent  
523 fractionation of these elements. If it is not possible to exclude the radiogenic isotope as an inversion isotope,  
524 separate measurements of the unspiked and spiked sample are required. This allows one to obtain the true  
525 isotopic composition and a non-radiogenic isotope pair can subsequently be expressed as a deviation from the  
526 reference material (e.g.,  $\delta^{88/86}\text{Sr}$ ; Krabbenhöft *et al.* 2009, Lewis *et al.* 2017). In the case of Pb where three out  
527 of four stable isotopes are radiogenic, two measurements are also required but it is impossible to disentangle  
528 the contribution of mass-dependent fractionation from radiogenic anomalies.

529 Other sources of deviation from mass-dependent fractionation include the nuclear field shift effect,  
530 cosmogenic or spallation effects and nucleosynthetic anomalies. Nuclear field shift anomalies result from  
531 variations in the shape and size of nuclei as a function of the number of neutrons and are commonly  
532 manifested as an odd-even isotope effect. It predominantly affects heavier elements and has been  
533 demonstrated for Hg, Tl and U in natural samples (e.g., Dauphas and Schauble 2016). In addition, it can occur  
534 during ion-exchange chemistry or measurement (e.g., Fujii *et al.* 2009). For instance, nuclear field shift effects  
535 have been invoked to explain the anomalous behaviour of  $^{207}\text{Pb}$  above certain filament temperatures in TIMS  
536 measurements (e.g., Thirlwall 2000, Amelin *et al.* 2005, Klaver *et al.* 2016). As nuclear field shift effects remain  
537 poorly understood, there is no straightforward way to correct potential fractionation during sample processing.

538 Extra-terrestrial materials often display mass-independent anomalies as a result of cosmogenic and/or  
539 nucleosynthetic effects. Meteorites display systematic mass-independent variations resulting from the  
540 heterogeneous distribution of the products of stellar nucleosynthesis in the Solar System, roughly as a function



541 of heliocentric distance (e.g., Dauphas 2017, Kruijer *et al.* 2017). In addition, materials exposed in space  
542 undergo neutron-capture reactions induced by interactions with cosmic rays, thus producing cosmogenic  
543 anomalies. Again, the simplest way to deal with these mass-independent anomalies is to perform a spiked and  
544 an unspiked measurement for a sample, but this might not be possible when the amount of sample is limited,  
545 such as in the case of precious meteoritic samples, or when contamination of the unspiked sample with double  
546 spike during sample preparation or measurement is an issue. In this case, it is still possible to perform some  
547 corrections to account for the mass-independent anomalies if additional data are available. For instance,  
548 measurements of another element have been successfully used as neutron dosimeter to correct cosmogenic  
549 anomalies in the element of interest (e.g., Kruijer *et al.* 2013). When internally-normalised mass-independent  
550 data are reported in the literature, ideally for the same meteorite/sample or at least the same meteorite class,  
551 these data can be used to correct for nucleosynthetic anomalies given that these are expressed relative to  
552 the same reference material. The most direct approach is to impose the nucleosynthetic anomaly of the sample  
553 on the composition of the reference material; this is equivalent to making an unspiked measurement and, in  
554 other words, synthesising an unspiked measurement from the reference ratios and mass-independent  
555 anomalies. The obtained alternative reference material is related to the sample only by mass-dependent  
556 fractionation and can then be used as an input parameter in the double spike inversion to solve for the natural  
557 mass-dependent fractionation. Alternatively, one can apply an *a posteriori* correction to double spike data  
558 given that the magnitude of the mass-independent anomalies is known (Burkhardt *et al.* 2014, Hu and Dauphas  
559 2017). The latter is likely easier to implement and we found no significant difference between the two methods  
560 correcting our mass-dependent Ni data.

#### 561 **6.4. Optimising measurements when sample is limited**

562 In our arguments for choosing an appropriate double spike above, we assumed that an unlimited amount of  
563 sample material is available, i.e. one can always hold the total beam intensity constant regardless of the double  
564 spike contribution. This is rarely the case and the amount of analyte available for measurement can be  
565 relatively small, thus restricting the presumed flexibility in choosing sample-spike proportions. John (2012)  
566 suggested that increasing the amount of double spike in excess of the optimum proportion can improve the  
567 precision, which might be beneficial when sample-limited. There is some merit to this argument: a higher total  
568 beam intensity leads to an improvement in counting statistics, which might exceed the decrease in precision  
569 caused by a suboptimal sample-spike proportion. That this is, however, not always the case is illustrated in  
570 Figure 8. We modelled the theoretical precision in  $\alpha$  as a function of sample-double spike proportion with the  
571 Double Spike Toolbox, but instead of adopting a constant total beam intensity we fixed the amount of natural  
572 Ni to 0.5 V and varied the amount of double spike added. For our  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  double spike, precision indeed  
573 improves marginally when double spike was added in excess of the optimal proportion (denoted by a star in  
574 Figure 8), as found by John (2012), but this is not the case for a hypothetical  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  spike (the optimal  
575 composition based on Oak Ridge single spikes). This makes sense when comparing the behaviour of these  
576 spikes at a constant total beam intensity (Figure 4): the  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike is characterised by a steeper  
577 decrease in precision away from the optimum that is not balanced by the improvement in counting statistics

578 obtained through over-spiking. Note that, for a  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike, over-spiking holds no benefit even when  
579 prepared from isotopically pure single spikes. Another important constraint is the trade-off with a potential  
580 decrease in accuracy at higher sample-spike proportions as a result of a bias in the double spike calibration  
581 (Figure 7). We argue that the marginal gain in precision at high sample-spike proportions is not sufficient to  
582 potentially jeopardize the robustness of the data. Hence, we are wary of over-spiking to improve precision  
583 unless the accuracy of the double spike data at high sample-spike proportions is explicitly demonstrated.

584

## 585 **7. Summary and outlook**

586 The aim of this review article was to demonstrate the usefulness of the double spike method in obtaining  
587 reliable mass-dependent isotope data and lowering potential inertia barriers in working with double spikes  
588 through a step-by-step discussion. We emphasize that the double spike technique is the only way to eliminate  
589 the effects of potential mass-dependent fractionation during sample processing and measurement, given that  
590 the double spike is added beforehand. In addition, it is superior to other methods, including sample-standard  
591 bracketing, in correcting for fractionation during measurement as it is not susceptible to matrix effects. An  
592 external normalisation to spiked reference materials measured alongside the samples can eliminate the effects  
593 of long- and short-term instrumental drift and potential inaccuracies in the double spike calibration. With the  
594 double spike technique, it has been proven to be possible to measure natural mass-dependent fractionation to  
595 an uncertainty better than 20 ppm for elements with at least four isotopes. In the light of the rapid increase in  
596 the application of the “non-traditional stable isotopes” in the high-temperature geochemical realm, the use of  
597 the double spikes will continue to grow. We envisage that double spiking will rapidly replace sample-standard  
598 bracketing as the method of choice for obtaining reliable mass-dependent isotope data.

599 Although the basics of the double spike technique are well established, there is still room for innovation and  
600 potential improvements. With ongoing developments in amplifier technology, measurements of low intensity  
601 ion beams will become increasingly precise. This might open possibilities for the use of radioactive spikes that is  
602 currently restricted by their cost and/or activity. Elements that are outside the reach of the regular double  
603 spike technique might have a radioactive isotope with a sufficiently long half live to allow its use as a spike. In  
604 addition, the regular double spike inversion requires three independent isotope ratios. For elements with five  
605 or more isotopes, additional independent isotope ratios are available for data reduction but in the present  
606 form of the inversion these are not used. There is no theoretical restriction to a 3-dimensional solution of the  
607 double spike equation; including more independent isotope ratios will leave the system overdetermined.  
608 Although difficult to visualise, a least-squares regression solution in a higher number of dimensions can  
609 potentially increase the precision and accuracy of the double spike method. Moreover, in three dimensions no  
610 triple spike has been found to be advantageous over a double spike (e.g., Rudge *et al.* 2009), but whether this  
611 holds in a higher number of dimensions remains to be investigated.

612

## 613 **Acknowledgements**

614 We would like to thank editor Thomas Meisel for the invitation to write this back-to-basics review article. MK is  
615 greatly indebted for his analytical training to Richard Smeets while this article draws heavily on discussions with  
616 other double spikers, including Tim Elliott, Remco Hin, Jamie Lewis, Ian Parkinson, Pieter Vroon and Bas van der  
617 Wagt. Dan Bevan, Max Jansen and Danny Stubbs are thanked for providing useful feedback on the clarity and  
618 presentation of the manuscript. We acknowledge two anonymous reviewers and editor Thomas Meisel for  
619 their insightful comments and suggestions. MK acknowledges funding from the STFC through grant no.  
620 ST/M007715/1.

621

## 622 **Appendix A – Analytical details**

623

624 All Ni isotope data shown in the figures were collected using a Thermo-Finnigan Neptune multi-collector ICP-  
625 MS and a CETAC Aridus introduction system at the University of Bristol. When applicable, samples were spiked  
626 prior to ion-exchange purification. All measurements were carried out in medium resolution ( $m/\Delta m > 6000$ , 5-  
627 95% peak edge definition) to resolve isobaric interferences, notably  $^{40}\text{Ar}^{18}\text{O}^+$  on  $^{58}\text{Ni}^+$  and trace argide and oxide  
628 interference on the other masses. Isobaric interference of  $^{58}\text{Fe}$  on  $^{58}\text{Ni}$  could not be resolved but was corrected  
629 by monitoring  $^{56}\text{Fe}^+$  and  $^{57}\text{Fe}^+$  and subtracting IMF-corrected  $^{58}\text{Fe}^+$  from  $^{58}\text{Ni}^+$ ; this correction was found to be  
630 robust to at least  $^{56}\text{Fe}/^{60}\text{Ni} \geq 0.6$  while samples had  $^{56}\text{Fe}/^{60}\text{Ni} < 0.15$ . Samples were measured at a concentration  
631 of total Ni (sample plus double spike) of  $\sim 1 \mu\text{g ml}^{-1}$  in  $0.3 \text{ mol L}^{-1} \text{ HNO}_3$ , yielding a total Ni ion beam intensity of  
632 50-85 V on  $10^{11} \Omega$  amplifiers compared to a background of  $< 20 \text{ mV}$ . Each sample was measured 6-8 times (50  
633 cycles of 4.2 s integration time) in at least two measurement sessions, always bracketed by measurements of  
634 double spiked SRM 986 reference material. Results are expressed as  $\delta^{60/58}\text{Ni}$  relative to this reference material  
635 (Equation 5) with a quoted precision that is 2 standard error of the mean of the replicate measurements.

636

## 637 **References**

638 **Albarède F. and Beard B. (2004)** Analytical methods for non-traditional isotopes **Reviews in Mineralogy and**  
639 **Geochemistry, 55**, 113-152.

640 **Albarede F., Telouk P., Blichert-Toft J., Boyet M., Agranier A. and Nelson B. (2004)** Precise and accurate  
641 isotopic measurements using multiple-collector ICPMS **Geochimica et Cosmochimica Acta, 68**, 2725-2744.

642 **Amelin Y., Davis D. and Davis W. (2005)** Decoupled fractionation of even- and odd-mass isotopes of Pb in TIMS  
643 **Geochimica et Cosmochimica Acta Supplement, 69**, 215.

644 **Andersen M.B., Elliott T., Freymuth H., Sims K.W., Niu Y. and Kelley K.A. (2015)** The terrestrial uranium  
645 isotope cycle **Nature, 517**, 356.

646 **Barling J. and Weis D. (2008)** Influence of non-spectral matrix effects on the accuracy of Pb isotope ratio  
647 measurement by MC-ICP-MS: implications for the external normalization method of instrumental mass bias  
648 correction **Journal of Analytical Atomic Spectrometry**, **23**, 1017-1025.

649 **Bigeleisen J. and Mayer M.G. (1947)** Calculation of equilibrium constants for isotopic exchange reactions **The**  
650 **Journal of Chemical Physics**, **15**, 261-267.

651 **Burkhardt C., Hin R.C., Kleine T. and Bourdon B. (2014)** Evidence for Mo isotope fractionation in the solar  
652 nebula and during planetary differentiation **Earth and Planetary Science Letters**, **391**, 201-211.

653 **Carlson R.W. (2014)** Thermal Ionization Mass Spectrometry. In: **Holland H. and Turekian K.K. (eds)**, **Treatise on**  
654 **Geochemistry**, 2nd edition. Elsevier, 337-354.

655 **Chernozhukhin S.M., Goderis S., Lobo L., Claeys P. and Vanhaecke F. (2015)** Development of an isolation  
656 procedure and MC-ICP-MS measurement protocol for the study of stable isotope ratio variations of nickel  
657 **Journal of Analytical Atomic Spectrometry**, **30**, 1518-1530.

658 **Coath C.D., Elliott T. and Hin R.C. (2017)** Double-spike inversion for three-isotope systems **Chemical Geology**,  
659 **451**, 78-89.

660 **Compston W. and Oversby V. (1969)** Lead isotopic analysis using a double spike **Journal of Geophysical**  
661 **Research**, **74**, 4338-4348.

662 **Coplen T.B. (2011)** Guidelines and recommended terms for expression of stable-isotope-ratio and gas-ratio  
663 measurement results **Rapid communications in mass spectrometry**, **25**, 2538-2560.

664 **Creech J., Baker J., Handler M., Schiller M. and Bizzarro M. (2013)** Platinum stable isotope ratio measurements  
665 by double-spike multiple collector ICPMS **Journal of Analytical Atomic Spectrometry**, **28**, 853-865.

666 **Creech J.B. and Paul B. (2015)** IsoSpike: Improved Double-Spike Inversion Software **Geostandards and**  
667 **Geoanalytical Research**, **39**, 7-15.

668 **Cumming G. (1973)** Propagation of experimental errors in lead isotope ratio measurements using the double  
669 spike method **Chemical Geology**, **11**, 157-165.

670 **Dauphas N. (2017)** The isotopic nature of the Earth's accreting material through time **Nature**, **541**, 521.

671 **Dauphas N. and Schauble E.A. (2016)** Mass fractionation laws, mass-independent effects, and isotopic  
672 anomalies **Annual Review of Earth and Planetary Sciences**, **44**, 709-783.

673 **Dodson M. (1963)** A theoretical study of the use of internal standards for precise isotopic analysis by the  
674 surface ionization technique: Part I-General first-order algebraic solutions **Journal of Scientific Instruments**, **40**,  
675 289.

676 **Fujii T., Moynier F. and Albarède F. (2009)** The nuclear field shift effect in chemical exchange reactions  
677 **Chemical Geology**, **267**, 139-156.

678 **Galer S.J.G. (1999)** Optimal double and triple spiking for high precision lead isotopic measurement **Chemical**  
679 **Geology, 157**, 255-274.

680 **Galer S.J.G. and Abouchami W. (1998)** Practical application of lead triple spiking for correction of instrumental  
681 mass discrimination **Mineral. Mag. A, 62**, 491-492.

682 **Gall L., Williams H., Siebert C. and Halliday A. (2012)** Determination of mass-dependent variations in nickel  
683 isotope compositions using double spiking and MC-ICPMS **Journal of Analytical Atomic Spectrometry, 27**, 137-  
684 145.

685 **Gramlich J., Machlan L., Barnes I. and Paulsen P. (1989)** Absolute isotopic abundance ratios and atomic weight  
686 of a reference sample of nickel **Journal of research of the National Institute of Standards and Technology, 94**,  
687 347.

688 **Hamelin B., Manhès G., Albarede F. and Allegre C.J. (1985)** Precise lead isotope measurements by the double  
689 spike technique: a reconsideration **Geochimica et Cosmochimica Acta, 49**, 173-182.

690 **Hofmann A. (1971)** Fractionation corrections for mixed-isotope spikes of Sr, K, and Pb **Earth and Planetary**  
691 **Science Letters, 10**, 397-402.

692 **Hu J.Y. and Dauphas N. (2017)** Double-spike data reduction in the presence of isotopic anomalies **Journal of**  
693 **Analytical Atomic Spectrometry, 32**, 2024-2033.

694 **John S.G. (2012)** Optimizing sample and spike concentrations for isotopic analysis by double-spike ICPMS  
695 **Journal of Analytical Atomic Spectrometry, 27**, 2123-2131.

696 **Johnson C.M. and Beard B.L. (1999)** Correction of instrumentally produced mass fractionation during isotopic  
697 analysis of Fe by thermal ionization mass spectrometry **International Journal of Mass Spectrometry, 193**, 87-  
698 99.

699 **Klaver M., Smeets R., Koornneef J.M., Davies G. and Vroon P. (2016)** Pb isotope analysis of ng size samples by  
700 TIMS equipped with a  $10^{13} \Omega$  resistor using a  $^{207}\text{Pb}$ - $^{204}\text{Pb}$  double spike **Journal of Analytical Atomic**  
701 **Spectrometry, 31**, 171-178.

702 **Krabbenhöft A., Fietzke J., Eisenhauer A., Liebetrau V., Böhm F. and Vollstaedt H. (2009)** Determination of  
703 radiogenic and stable strontium isotope ratios ( $^{87}\text{Sr}/^{86}\text{Sr}$ ;  $\delta^{88/86}\text{Sr}$ ) by thermal ionization mass spectrometry  
704 applying an  $^{87}\text{Sr}/^{84}\text{Sr}$  double spike **Journal of Analytical Atomic Spectrometry, 24**, 1267-1271.

705 **Kruijjer T.S., Burkhardt C., Budde G. and Kleine T. (2017)** Age of Jupiter inferred from the distinct genetics and  
706 formation times of meteorites **Proceedings of the National Academy of Sciences, 114**, 6712-6716.

707 **Kruijjer T.S., Fischer-Gödde M., Kleine T., Sprung P., Leya I. and Wieler R. (2013)** Neutron capture on Pt  
708 isotopes in iron meteorites and the Hf–W chronology of core formation in planetesimals **Earth and Planetary**  
709 **Science Letters, 361**, 162-172.

710 **Lewis J., Pike A., Coath C. and Evershed R. (2017)** Strontium concentration, radiogenic ( $^{87}\text{Sr}/^{86}\text{Sr}$ ) and stable  
711 ( $\delta^{88}\text{Sr}$ ) strontium isotope systematics in a controlled feeding study **STAR: Science & Technology of**  
712 **Archaeological Research**, **3**, 45-57.

713 **Longerich H., Fryer B. and Strong D. (1987)** Determination of lead isotope ratios by inductively coupled plasma-  
714 mass spectrometry (ICP-MS) **Spectrochimica Acta Part B: Atomic Spectroscopy**, **42**, 39-48.

715 **Millet M.-A. and Dauphas N. (2014)** Ultra-precise titanium stable isotope measurements by double-spike high  
716 resolution MC-ICP-MS **Journal of Analytical Atomic Spectrometry**, **29**, 1444-1458.

717 **Nan X., Wu F., Zhang Z., Hou Z., Huang F. and Yu H. (2015)** High-precision barium isotope measurements by  
718 MC-ICP-MS **Journal of Analytical Atomic Spectrometry**, **30**, 2307-2315.

719 **Nier A.O. (1938)** The isotopic constitution of strontium, barium, bismuth, thallium and mercury **Physical**  
720 **Review**, **54**, 275.

721 **Oi T., Kawada K., Hosoe M. and Kakihana H. (1991)** Fractionation of lithium isotopes in cation-exchange  
722 chromatography **Separation Science and Technology**, **26**, 1353-1375.

723 **Peters S.T., Münker C., Wombacher F. and Elfers B.-M. (2015)** Precise determination of low abundance  
724 isotopes ( $^{174}\text{Hf}$ ,  $^{180}\text{W}$  and  $^{190}\text{Pt}$ ) in terrestrial materials and meteorites using multiple collector ICP-MS equipped  
725 with  $10^{12} \Omega$  Faraday amplifiers **Chemical Geology**, **413**, 132-145.

726 **Prytulak J., Nielsen S., Plank T., Barker M. and Elliott T. (2013)** Assessing the utility of thallium and thallium  
727 isotopes for tracing subduction zone inputs to the Mariana arc **Chemical Geology**, **345**, 139-149.

728 **Ripperger S. and Rehkämper M. (2007)** Precise determination of cadmium isotope fractionation in seawater by  
729 double spike MC-ICPMS **Geochimica et Cosmochimica Acta**, **71**, 631-642.

730 **Rosman K. (1972)** A survey of the isotopic and elemental abundance of zinc **Geochimica et Cosmochimica**  
731 **Acta**, **36**, 801-819.

732 **Rudge J.F., Reynolds B.C. and Bourdon B. (2009)** The double spike toolbox **Chemical Geology**, **265**, 420-431.

733 **Russell R. (1971)** The systematics of double spiking **Journal of Geophysical Research**, **76**, 4949-4955.

734 **Russell W. and Papanastassiou D. (1978)** Calcium isotope fractionation in ion-exchange chromatography  
735 **Analytical Chemistry**, **50**, 1151-1154.

736 **Russell W., Papanastassiou D. and Tombrello T. (1978)** Ca isotope fractionation on the Earth and other solar  
737 system materials **Geochimica et Cosmochimica Acta**, **42**, 1075-1090.

738 **Schauble E.A. (2004)** Applying stable isotope fractionation theory to new systems **Reviews in Mineralogy and**  
739 **Geochemistry**, **55**, 65-111.

740 **Schönbächler M. and Fehr M.A. (2013)** Basics of ion exchange chromatography for selected geological  
741 applications. **Treatise on geochemistry. Vol. 15: Analytical geochemistry/inorganic instrument analysis.**  
742 **Elsevier**, 124-146.

743 **Shiel A.E., Barling J., Orians K.J. and Weis D. (2009)** Matrix effects on the multi-collector inductively coupled  
 744 plasma mass spectrometric analysis of high-precision cadmium and zinc isotope ratios **Analytica Chimica Acta,**  
 745 **633,** 29-37.

746 **Siebert C., Nägler T.F. and Kramers J.D. (2001)** Determination of molybdenum isotope fractionation by double-  
 747 spike multicollector inductively coupled plasma mass spectrometry **Geochemistry, Geophysics, Geosystems, 2.**

748 **Stracke A., Scherer E.E. and Reynolds B.C. (2014)** Application of Isotope Dilution in Geochemistry. In: **Turekian**  
 749 **K.K. (ed), Treatise on Geochemistry (Second Edition). Elsevier (Oxford),** 71-86.

750 **Taylor R.N., Ishizuka O., Michalik A., Milton J.A. and Croudace I.W. (2015)** Evaluating the precision of Pb  
 751 isotope measurement by mass spectrometry **Journal of Analytical Atomic Spectrometry, 30,** 198-213.

752 **Thirlwall M. (2000)** Inter-laboratory and other errors in Pb isotope analyses investigated using a  $^{207}\text{Pb}$ – $^{204}\text{Pb}$   
 753 double spike **Chemical Geology, 163,** 299-322.

754 **Thirlwall M. (2002)** Multicollector ICP-MS analysis of Pb isotopes using a  $^{207}\text{Pb}$ – $^{204}\text{Pb}$  double spike demonstrates  
 755 up to 400 ppm/amu systematic errors in TI-normalization **Chemical Geology, 184,** 255-279.

756 **Tipper E.T., Louvat P., Capmas F., Galy A. and Gaillardet J. (2008)** Accuracy of stable Mg and Ca isotope data  
 757 obtained by MC-ICP-MS using the standard addition method **Chemical Geology, 257,** 65-75.

758 **Todt W., Cliff R.A., Hanser A. and Hofmann A. (1996)** Evaluation of a  $^{202}\text{Pb}$ – $^{205}\text{Pb}$  Double Spike for High-  
 759 Precision Lead Isotope Analysis **Earth processes: reading the isotopic code,** 429-437.

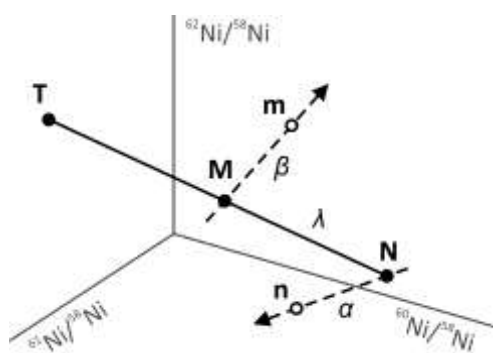
760 **van den Boorn S.H., Vroon P.Z. and van Bergen M.J. (2009)** Sulfur-induced offsets in MC-ICP-MS silicon-  
 761 isotope measurements **Journal of Analytical Atomic Spectrometry, 24,** 1111-1114.

762 **Waight T., Baker J. and Willigers B. (2002)** Rb isotope dilution analyses by MC-ICPMS using Zr to correct for  
 763 mass fractionation: towards improved Rb–Sr geochronology? **Chemical Geology, 186,** 99-116.

764 **Young E.D., Manning C.E., Schauble E.A., Shahar A., Macris C.A., Lazar C. and Jordan M. (2015)** High-  
 765 temperature equilibrium isotope fractionation of non-traditional stable isotopes: Experiments, theory, and  
 766 applications **Chemical Geology, 395,** 176-195.

767

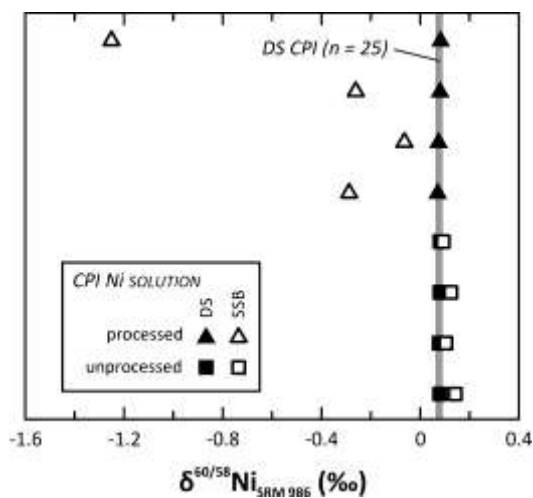
768 **FIGURE CAPTIONS**



769

770 **Figure 1.** Schematic diagram showing the principle of the double spike technique. In 4-isotope space, using Ni  
 771 as an example, a natural sample ( $N$ ) and double spike ( $T$ ) have very different compositions. A mixture between  
 772 the two ( $M$ ) lies on the  $N$ - $M$ - $T$  mixing line (mixing parameter  $\lambda$ ), but the measured composition of the mixture  
 773 ( $m$ ) and natural sample ( $n$ ) are displaced along instrumental mass fractionation (IMF) lines  $M$ - $m$  and  $N$ - $n$ ,  
 774 respectively. The form of these IMF lines is known, but the degree of fractionation (mass fractionation  
 775 parameters  $\alpha$  and  $\beta$ ) is not. As the composition of the double spike is known, measurement of  $n$  and  $m$  is  
 776 needed to solve the double spike equation to yield  $\alpha$ ,  $\beta$  and  $\lambda$ , where  $\alpha$  can subsequently be used to calculate  
 777 the true composition  $N$  of the sample.

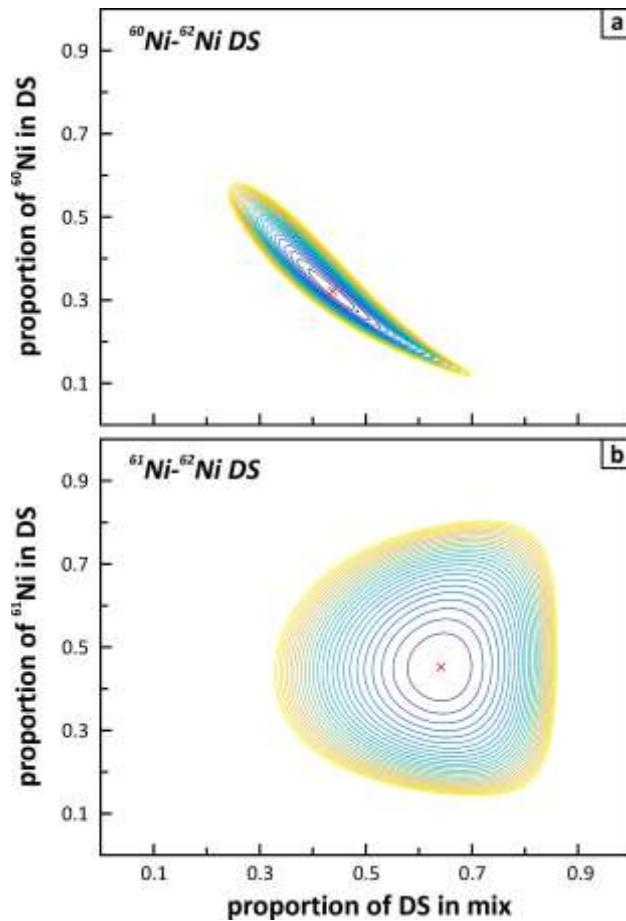
778



779 **Figure 2.** Comparison of double spike (DS) versus sample-standard bracketing (SSB) corrected data for our in-  
 780 house CPI Ni solution. Data points are the average of 6-8 repeats of the same solution measured in at least two  
 781 measurement sessions; repeatability precision (2 SE of 6-8 independent measurement results per sample) is  
 782 smaller than symbol size. The CPI Ni solution is elementally pure and should hence be expected to behave the  
 783 same as the SRM 986 Ni reference material. The CPI solution was measured without processing, but spiked  
 784 aliquots were also processed along with samples through the chemical purification protocol. Double spike-  
 785 reduced data are highly reproducible ( $\delta^{60/58}\text{Ni}_{\text{SRM 986}} = 0.078 \pm 0.015 \text{ ‰}$ , 2s,  $n = 25$ ; grey bar). The double spike  
 786 inversion yields a value for  $\beta$  (see section 2.2) for every measurement, allowing a correction to be made  
 787 assuming that IMF is identical for the CPI standard and SRM 986. These pseudo-corrected data (shown as open  
 788 symbols) therefore show the composition of the CPI standard as if they were corrected through sample-  
 789 standard bracketing. See text for discussion.

791

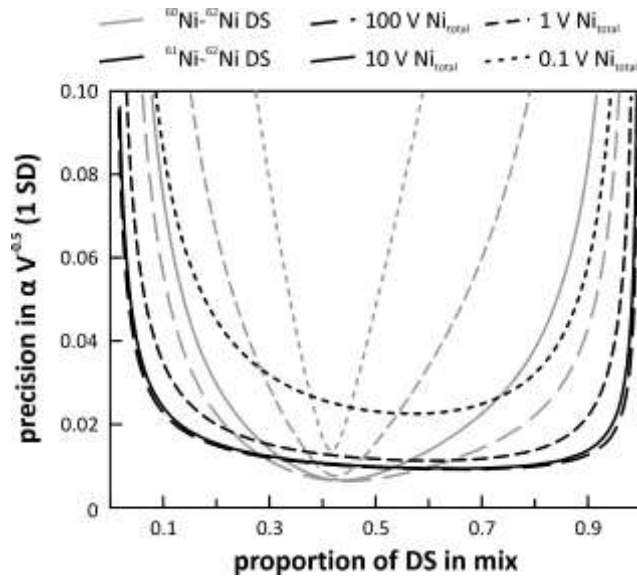




792

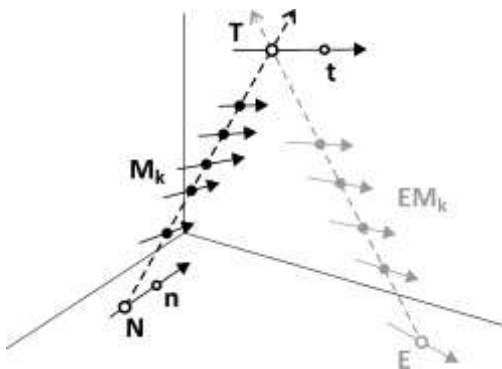
793 **Figure 3.** Contour diagram of the theoretical precision in  $\alpha$  (see Figure 1) as a function of double spike  
 794 composition and sample-spike ratio for two different Ni double spikes, generated using the default settings of  
 795 the precision model in the Double Spike Toolbox (Rudge et al. 2009) and  $^{58}\text{Ni}$ ,  $^{60}\text{Ni}$ ,  $^{61}\text{Ni}$  and  $^{62}\text{Ni}$  as inversion  
 796 isotopes. The precision decreases from blue to yellow; the optimal composition of the double spike is denoted  
 797 by the cross. Although a  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike will yield a slightly higher absolute precision in  $\alpha$  at the optimal  
 798 sample-spike proportion (see Figure 4), its magnification of measurement precision is much more dependent  
 799 on the sample-spike ratio than for a  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  double spike.

800



801  
802  
803  
804  
805  
806  
807  
808  
809  
810  
811  
812  
813  
814

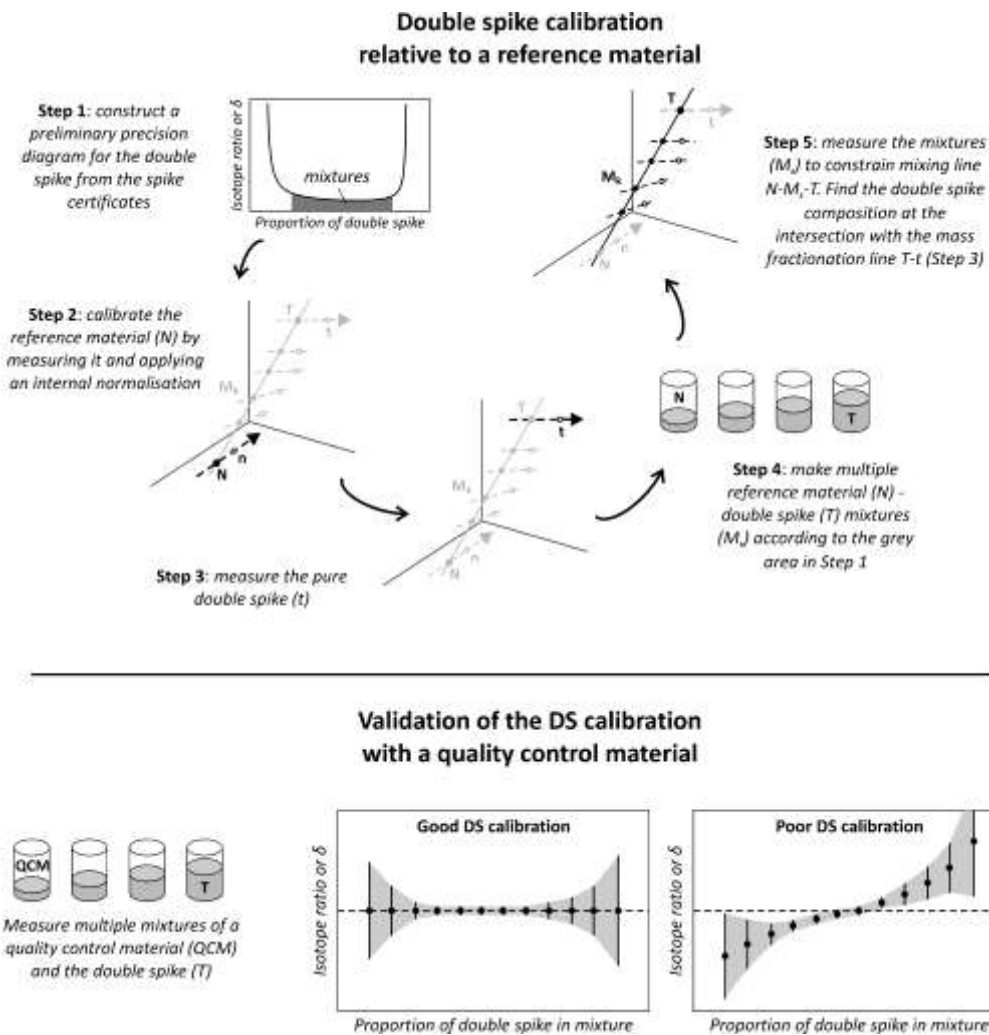
**Figure 4.** The theoretical precision in  $\alpha$  as a function of sample-spike ratio and total beam intensity (in V for  $10^{11} \Omega$  amplifier feedback resistors) for the two Ni double spikes (the optimal double spike composition shown in Figure 3), generated using the precision model in the Double Spike Toolbox (Rudge et al. 2009) and  $^{58}\text{Ni}$ ,  $^{60}\text{Ni}$ ,  $^{61}\text{Ni}$  and  $^{62}\text{Ni}$  as inversion isotopes; the default setting of the Toolbox is a 100 pA (10 V) total ion beam intensity, 8 s integration time and  $10^{11} \Omega$  amplifier feedback resistors. The precision in  $\alpha$  is divided by the square root of the beam intensity to eliminate the effect of shot noise and facilitate comparison of the two double spikes. A  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  double spike has a broad precision minimum at any total ion beam intensity and hence, near this minimum, the precision in  $\alpha$  is not very sensitive to the proportion in which the sample and double spike are mixed, which is ideal if the concentration of either is not well known. A  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  double spike on the other hand can produce a higher absolute precision for  $\alpha$  if double spike and sample are mixed accurately but suffers a lower precision when the sample-spike proportion is suboptimal. This behaviour is compounded at lower beam intensities.



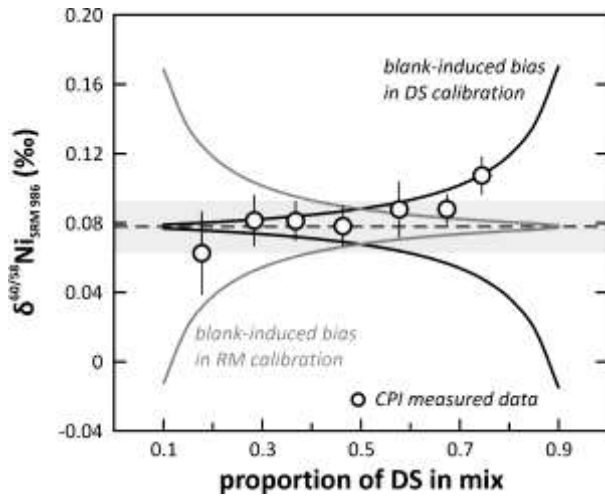
815  
816  
817  
818  
819

**Figure 5.** Schematic four-isotope diagram illustrating the calibration of a double spike. Solid arrows indicate mass fractionation lines; dashed arrows mixing lines. The measured composition of a reference material ( $n$ ) and double spike ( $t$ ) is offset from the true values  $N$  and  $T$  by IMF. A first-order IMF correction will yield the composition  $N$ . Multiple mixtures  $M_k$  ( $k = 1, 2, \dots$ ) between  $N$  and  $T$  will constrain the direction of the mixing

820 line  $N-T$ ; the calibrated double spike composition is the intersection between the mixing line and the mass  
 821 fractionation line  $t-T$ . Alternatively, a reference material with an exotic composition ( $E$ ) and mixtures  $EM_k$  ( $k = 1,$   
 822  $2, \dots$ ) can be used to construct a secondary mixing line  $E-T$  such that the calibrated double spike composition  
 823 lies at the intersection of mixing lines  $N-T$  and  $E-T$ , thus obviating the need of a pure double spike  
 824 measurement.  
 825



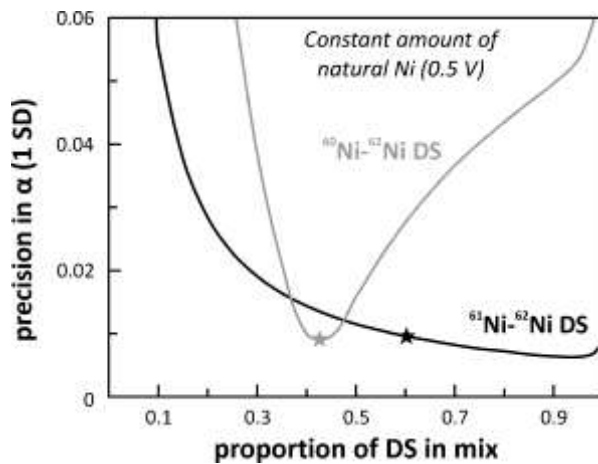
826  
 827 **Figure 6.** Flowchart diagram for the calibration procedure of a double spike ( $T$ ) against a reference material ( $N$ )  
 828 and its validation using a quality control material (QCM).  
 829



830

831 **Figure 7.** Effects of errors in blank corrections on the accuracy of double spike data. Modelled curves illustrate  
 832 the effect of under- and over-correcting for blank contribution on the pure reference material (RM; grey  
 833 curves) and double spike (DS; black curves) compositions. The bias introduced by a DS blank correction error  
 834 increases at higher double spike proportions; contamination of the pure RM measurement with double spike  
 835 introduces a bias at low double spike proportions; see main text for more discussion. The in-house CPI Ni  
 836 solution was measured at different sample-spike proportions and yield results within uncertainty of the  
 837 intermediate measurement precision (light grey bar) for  $0.15 < p < 0.75$ ; the measurement at  $p \sim 0.8$  is  
 838 marginally higher and might indicate a small bias in the double spike calibration.

839



840

841 **Figure 8.** The precision in  $\alpha$  as a function of sample-spike ratio and total beam intensity for the two Ni double  
 842 spikes shown in Figures 3 and 4. Here, the amount of sample is fixed at 0.5 V natural Ni while the amount of  
 843 double spike added to this sample is varied to obtain different sample-spike proportions. The optimal double  
 844 spike composition at a total beam intensity of 1 V is indicated by the stars. In the case of a  $^{61}\text{Ni}$ - $^{62}\text{Ni}$  spike,  
 845 increasing the amount of spike in the mixture past the optimal proportion at 1 V leads to a marginal  
 846 improvement in precision, but potentially at the cost of a bias introduced by inaccuracies in the double spike  
 847 calibration (compare to Figure 6). A  $^{60}\text{Ni}$ - $^{62}\text{Ni}$  spike does not offer any advantage for over-spiking.