Research Article

Exact Solutions of the Generalized Benjamin-Bona-Mahony Equation

Xun Liu, Lixin Tian, and Yuhai Wu

Faculty of Science, Jiangsu University, Zhenjiang, Jiangsu 212013, China

Correspondence should be addressed to Xun Liu, liuxundy@163.com

We apply the theory of Weierstrass elliptic function to study exact solutions of the generalized Benjamin-Bona-Mahony equation. By using the theory of Weierstrass elliptic integration, we get some traveling wave solutions, which are expressed by the hyperbolic functions and trigonometric functions. This method is effective to find exact solutions of many other similar equations which have arbitrary-order nonlinearity.

1. Introduction

The nonlinear phenomena in the scientific work or engineering fields are more and more attractive to scientists. To depict and analyze such nonlinear phenomena, the nonlinear evolutionary equations are playing an important role and their solitary wave solutions are the main interests of mathematicians and physicists.

To obtain the traveling wave solutions of these nonlinear evolution equations, many methods were attempted, such as the inverse scattering method, Hirota's bilinear transformation, the tanh-sech method, extended tanh method, sine-cosine method, homogeneous balance method, and exp-function method. With the aid of symbolic computation system, many explicit solutions are easily obtained, and many interesting works deeply promote the research of nonlinear phenomena.

The present work is interested in generalized Benjamin-Bona-Mahony (BBM) equation:

$$u_t + au_x + \left(bu^n + cu^{2n}\right)u_x + ku_{xxx} = 0. \tag{1.1}$$

In the above equation, the first term of left side represents the evolution term while parameters b and c represent the coefficients of dual-power law nonlinearity, a and k are the coefficients of dispersion terms, n is the power law parameter, and variable u is the wave profile. In [1], Biswas used the solitary wave ansatz and obtained an exact 1-soliton solution of (1.1). In order to find more exact solutions of some nonlinear evolutionary equations, the Weierstrass elliptic function was introduced. For example, Kuru [2, 3] discussed the BBM-like equation, and Estévez et al. [4] analyzed another type of generalized BBM equations. In [5], Deng et al. also applied the similar method to the study of a nonlinear variant of the PHI-four equation. In this paper, we will apply the method to the generalized Benjamin-Bona-Mahony equation.

The rest of this paper is organized as follows. In Section 2, we first outline the Weierstrass elliptic function method. In Section 3, we give exact expression of some traveling wave solutions of generalized Benjamin-Bona-Mahony (BBM) equation (1.1) by using the Weierstrass elliptic function method. Finally, some conclusions are given in Section 4.

2. Description of the Weierstrass Elliptic Function Method

When we search for the solutions of some evolutionary equations, we will meet the following ordinary differential equation:

$$\left(\frac{d\varphi}{d\theta}\right)^2 = P_4(\varphi) = a_0\varphi^4 + 4a_1\varphi^3 + 6a_2\varphi^2 + 4a_3\varphi + a_4. \tag{2.1}$$

In [6], Whittaker and Watson introduced two invariants:

$$g_2 = a_0 a_4 - 4a_1 a_3 + 3a_2^2, \qquad g_3 = a_0 a_2 a_4 + 2a_1 a_2 a_3 - a_2^3 - a_0 a_3^2 - a_1^2 a_4, \tag{2.2}$$

and a discriminant $\Delta = g_2^3 - 27g_3^2$; then the solutions of differential (2.1) have the following form:

$$\varphi(\theta) = \varphi_0 + \frac{1}{4}P_\varphi(\varphi_0)\left(\wp(\theta; g_2, g_3) - \frac{1}{24}P_{\varphi\varphi}(\varphi_0)\right)^{-1}, \tag{2.3}$$

where φ_0 is one of the roots of the polynomial $P_4(\varphi)$, and $P_\varphi(\varphi_0)$, $P_{\varphi\varphi}(\varphi_0)$, respectively, denote the first and second derivative of $P_4(\varphi)$ with respect to φ at the φ_0. Particularly, if $\Delta = g_2^3 - 27g_3^2 = 0$, then the Weierstrass elliptic function $\wp(\theta; g_2, g_3)$ satisfies these conditions:

$$\wp\left(\theta; 12b^2, -8b^3\right) = b + 3b\sinh^{-2}\left(\sqrt{3b}\theta\right),$$

$$\wp\left(\theta; 12b^2, 8b^3\right) = -b + 3b\sin^{-2}\left(\sqrt{3b}\theta\right). \tag{2.4}$$

Once we get the solution of (2.1) with the form of (2.3), we will get the exact expressions of solutions of many partial differential equations. Suppose that $\Delta = g_2^3 - 27g_3^2 = 0$; it is obvious that if $g_2 > 0$, $g_3 > 0$, then there exist period solutions of original evolutionary equation, otherwise if $g_2 > 0$, $g_3 < 0$, then there exist solitary solutions.

3. Discussion on Generalized BBM Equation (1.1)

Let $u(x, t) = u(\xi)$, $\xi = x - vt$, so (1.1) is carried to

$$(a - v)u + \frac{b}{n + 1}u^{n+1} + \frac{c}{2n + 1}u^{2n+1} + ku'' = D, \tag{3.1}$$

where D is an integral constant. We introduce the transformation that

$$u(\xi) = \left(\frac{k(n + 1)}{b}\right)^{1/n} W(\xi), \tag{3.2}$$

so that (3.1) is changed to

$$W'' - \frac{v - a}{k}W + W^{n+1} + \frac{ck(n + 1)^2}{b^2(2n + 1)}W^{2n+1} = \frac{D}{k}\left(\frac{b}{k(n + 1)}\right)^{1/n}. \tag{3.3}$$

Let $B = (v - a)/k$, $C = -ck(n + 1)^2/b^2(2n + 1)$, and $R_1 = (D/k)(b/k(n + 1))^{1/n}$; then (3.3) is transformed to the following form:

$$W'' = BW - W^{n+1} + CW^{2n+1} + R_1. \tag{3.4}$$

By multiplying each side of (3.4) by W', integrating once again, we get that

$$(W')^2 = BW^2 - \frac{2}{n + 2}W^{n+2} + \frac{C}{n + 1}W^{2n+2} + 2R_1W + R_2, \tag{3.5}$$

where R_2 is the other integration constant.

Let $W = \varphi^p$; then the above equation is changed to

$$\left(\frac{d\varphi}{d\xi}\right)^2 = \frac{B}{p^2}\varphi^2 - \frac{2}{(n + 2)p^2}\varphi^{np+2} + \frac{C}{(n + 1)p^2}\varphi^{2np+2} + \frac{2R_1}{p^2}\varphi^{2-p} + \frac{R_2}{p^2}\varphi^{2(1-p)}. \tag{3.6}$$

Only if all the power of φ appearing in the right side of (3.6) are the integers and are between 0 and 4, can we guarantee the integrability of (3.6). Therefore, if $R_1 = R_2 = 0$, then $p = \pm 1/n$. We choose the only case that if $p = -1/n$, then

$$\left(\frac{d\varphi}{d\xi}\right)^2 = n^2B\varphi^2 - \frac{2n^2}{n + 2}\varphi + \frac{n^2C}{n + 1}. \tag{3.7}$$

The polynomial $P(\varphi)$ has two roots: $\varphi_0 = ((n+1)\pm\sqrt{(n + 1)[BC(n + 2)^2 + (n + 1)]})/B(n+1)(n+2)$. For simplicity, let $\lambda = \sqrt{(n + 1)[BC(n + 2)^2 + (n + 1)]}$, then $\varphi_0 = ((n+1)\pm\lambda)/B(n+1)(n+2)$

$$g_2 = \frac{n^4B^2}{12}, \qquad g_3 = -\frac{n^6B^3}{216}, \qquad \Delta = 0. \tag{3.8}$$

Then at the point $\varphi_0 = ((n+1) \pm \lambda)/B(n+1)(n+2)$, the nonzero solutions of (3.7) are

$$\wp(\theta; g_2, g_3) = \frac{n^2 B}{12} + \frac{1}{4}n^2 B \csc h^2\left(\frac{\sqrt{B}}{2}n\theta\right), \quad \text{(if } B > 0),$$

$$\wp(\theta; g_2, g_3) = \frac{n^2 B}{12} - \frac{1}{4}n^2 B \csc^2\left(\frac{\sqrt{-B}}{2}n\theta\right), \quad \text{(if } B < 0).$$

(3.9)

The solutions of the ordinary differential equation are

$$\varphi(\xi) = \frac{(n+1) \pm \left[\lambda - 2\lambda \cosh^2\left(\left(\sqrt{B}/2\right)n\theta\right)\right]}{B(n+1)(n+2)}, \quad \text{(if } B > 0),$$

$$\varphi(\xi) = \frac{(n+1) \pm \left[\lambda - 2\lambda \cos^2\left(\left(\sqrt{-B}/2\right)n\theta\right)\right]}{B(n+1)(n+2)}, \quad \text{(if } B < 0),$$

(3.10)

where $B = (v-a)/k$, $C = -ck(n+1)^2/b^2(2n+1)$, $R_1 = (D/k)(b/k(n+1))^{1/n}$. Therefore, the solutions of partial differential equation are

$$u(x,t) = \left(\frac{k(n+1)}{b} \frac{B(n+1)(n+2)}{(n+1) \pm \left[\lambda - 2\lambda \cosh^2\left(\left(\sqrt{B}/2\right)n\theta\right)\right]}\right)^{1/n}, \quad \text{(if } B > 0),$$

$$u(x,t) = \left(\frac{k(n+1)}{b} \frac{B(n+1)(n+2)}{(n+1) \pm \left[\lambda - 2\lambda \cos^2\left(\left(\sqrt{-B}/2\right)n\theta\right)\right]}\right)^{1/n}, \quad \text{(if } B < 0),$$

(3.11)

where $B = (v-a)/k, C = -ck(n+1)^2/b^2(2n+1)$, and $R_1 = (D/k)(b/k(n+1))^{1/n}$.

Remark 3.1. We have checked that (3.11) is the solution of (1.1). From the above discussion, we find more solutions than the work of [1].

4. Conclusions

From the above discussion, we find the traveling wave solutions of the generalized BBM equation, which are expressed by the hyperbolic functions and trigonometric functions, without the aid of symbolic computations. In addition, the method is effective to find exact solutions of many other similar equations which have arbitrary-order nonlinearity.

Acknowledgments

References

[1] A. Biswas, "1-soliton solution of Benjamin-Bona-Mahoney equation with dual-power law nonlinearity," *Communications in Nonlinear Science and Numerical Simulation*, vol. 15, no. 10, pp. 2744–2746, 2010.

[2] S. Kuru, "Compactons and kink-like solutions of BBM-like equations by means of factorization," *Chaos, Solitons and Fractals*, vol. 42, no. 1, pp. 626–633, 2009.

[3] Ş. Kuru, "Traveling wave solutions of the BBM-like equations," *Journal of Physics A*, vol. 42, no. 37, Article ID 375203, 12 pages, 2009.

[4] P. G. Estévez, Ş. Kuru, J. Negro, and L. M. Nieto, "Travelling wave solutions of the generalized Benjamin-Bona-Mahony equation," *Chaos, Solitons and Fractals*, vol. 40, no. 4, pp. 2031–2040, 2009.

[5] X. Deng, M. Zhao, and X. Li, "Travelling wave solutions for a nonlinear variant of the PHI-four equation," *Mathematical and Computer Modelling*, vol. 49, no. 3-4, pp. 617–622, 2009.

[6] E. T. Whittaker and G. N. Watson, *A Course of Modern Analysis. An Introduction to the General Theory of Infinite Processes and of Analytic Functions*, Cambridge Mathematical Library, Cambridge University Press, Cambridge, UK, 1996.