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Research Article

Logarithmical Regularity Criteria of the Three-Dimensional Micropolar Fluid Equations in terms of the Pressure

Yan Jia, Jing Zhang, and Bo-Qing Dong

School of Mathematical Sciences, Anhui University, Hefei 230601, China

Correspondence should be addressed to Bo-Qing Dong, bqdong@ahu.edu.cn

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This paper is devoted to the regularity criterion of the three-dimensional micropolar fluid equations. Some new regularity criteria in terms of the partial derivative of the pressure in the Lebesgue spaces and the Besov spaces are obtained which improve the previous results on the micropolar fluid equations.

1. Introduction and Main Result

In this paper, we consider a three-dimensional micropolar fluid equations model described by the following equations:

$$\begin{aligned}\nabla \cdot u &= 0, \\ \partial_t u - \Delta u - \nabla \times w + \nabla \pi + u \cdot \nabla u &= 0, \\ \partial_t w - \Delta w - \nabla \nabla \cdot w + 2w - \nabla \times u + u \cdot \nabla w &= 0,\end{aligned}\tag{1.1}$$

associated with the initial condition

$$u(x, 0) = u_0, \quad w(x, 0) = w_0,\tag{1.2}$$

where $u(x, t) = (u_1(x, t), u_2(x, t), u_3(x, t))$ is the unknown velocity vector field, $\pi(x, t)$ is the unknown scalar pressure field, and $w(x, t) = (w_1(x, t), w_2(x, t), w_3(x, t))$ is the unknown

microrotation vector field with $\nabla \cdot u_0 = 0$ in the sense of distribution. u_0 and w_0 represent the prescribed initial data for the velocity and microrotation fields.

Here,

$$\begin{aligned} \nabla &= (\partial_{x_1}, \partial_{x_2}, \partial_{x_3}), & \Delta &= \nabla \cdot \nabla, \\ \nabla \cdot u &= \sum_{i=1}^3 \partial_i u_i, & u \cdot \nabla w &= \sum_{i=1}^3 u_i \partial_i w_j, \quad (j = 1, 2, 3). \end{aligned} \quad (1.3)$$

Micropolar fluid equations introduced by Eringen [1] in 1960' are a special model of the non-Newtonian fluids which is coupled with the viscous incompressible Navier-Stokes model, microrotational effects and microrotational inertia. Physically, micropolar fluid may represent fluids that consist of rigid, randomly oriented (or spherical particles), suspended in a viscous medium, where the deformation of fluid particles is ignored. It can describe many phenomena appeared in a large number of complex fluids such as the suspensions, animal blood, and liquid crystals which cannot be characterized appropriately by the Navier-Stokes system [2]. When the microrotation effects are neglected or $w = 0$, the micropolar fluid equations reduces to the classic Navier-Stokes flows (see [3]).

Due to the importance of both physics and mathematics, the questions on the well posedness and asymptotic behaviors of solutions for micropolar fluid equations have been attracted more and more attention. Galdi and Rionero [4], Łukaszewicz [5] considered the existence of weak solutions of the micropolar fluid flows (1.1)-(1.2). the existence of strong solutions with either local for large data or global for small data is considered by many authors [6–9]. One may also refer to the interesting results on the large time behaviors of micropolar fluid equations in [10–13].

However, similar to the classic Navier-Stokes equations, the question of global regularity of the weak solutions of the 3D micropolar fluid equations still remain a big open problem. It is of importance on the study of regularity of the weak solutions under additional critical growth conditions on the velocity or the pressure. Based on some analysis technique, some regularity criteria via the velocity of weak solutions in the Lebesgue spaces, multiplier spaces and Besov spaces have been obtained in [14–16]. In particular, as for the pressure criterion, Dong et al. [17] (see also Yuan [18]) showed that the weak solution becomes regular if the pressure satisfies

$$\pi \in L^p\left(0, T; L^{q, \infty}(\mathbb{R}^3)\right), \quad \text{for } \frac{2}{p} + \frac{3}{q} \leq 2, \quad \frac{3}{2} < q \leq \infty, \quad (1.4)$$

or

$$\pi \in L^1\left(0, T; B_{\infty, \infty}^0(\mathbb{R}^3)\right), \quad (1.5)$$

where $L^{q, \infty}$ and $B_{\infty, \infty}^0$ denote the Lorentz space and Besov space. Later on, those results of pressure regularity criteria were further improved to the largest critical Besov spaces by Jia et al. [19] as

$$\pi(x, t) \in L^p\left(0, T; B_{q, \infty}^r(\mathbb{R}^3)\right), \quad (1.6)$$

where

$$\frac{2}{p} + \frac{3}{q} = 2 + r, \quad \frac{3}{2+r} < q < \infty, \quad -1 < r \leq 1. \tag{1.7}$$

One may also refer to some interesting results on the pressure criteria for Navier-Stokes equations [20–23] and MHD equations [24, 25].

The aim of this paper is to study the logarithmical regularity criterion for the three-dimensional micropolar fluid equations in terms of the partial derivative of pressure in the Lebesgue space or pressure in Besov spaces. Before giving the main results, we first recall the definition of weak solutions for micropolar fluid equations (1.1)-(1.2).

Definition 1.1 (Lukaszewicz [5]). Let $(u_0, w_0) \in L^2(\mathbb{R}^3)$ and $\nabla \cdot u_0 = 0$ in the sense of distribution. A pair vector field $(u(x, t), w(x, t))$ is termed as a weak solution of (1.1)-(1.2) on $(0, T)$, if (u, w) satisfies the following properties:

- (i) $(u, w) \in L^\infty(0, T; L^2(\mathbb{R}^3)) \cap L^2(0, T; H^1(\mathbb{R}^3))$;
- (ii) $\nabla \cdot u = 0$ in the sense of distribution;
- (iii) (u, w) verifies (1.1) in the sense of distribution.

Now our result reads as follows.

Theorem 1.2. *Suppose $(u_0, w_0) \in L^4(\mathbb{R}^3) \cap L^2(\mathbb{R}^3)$ and (u, w) is a weak solution of the 3D micropolar fluid equations (1.1)-(1.2). If the partial derivative of the pressure, for example, $\partial_3 \pi$ satisfies the logarithmical Serrin’s type condition*

$$\int_0^T \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4})} ds < \infty, \quad \frac{2}{p} + \frac{3}{q} = \frac{7}{4}, \quad \frac{12}{7} < q \leq \infty, \tag{1.8}$$

then the weak solution (u, w) becomes a regular solution on $(0, T)$.

Theorem 1.3. *On substitution of logarithmical Serrin’s type condition (1.8) by the following growth condition*

$$\int_0^T \frac{\|\pi\|_{\dot{B}_{\infty, \infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^4})} ds < \infty, \tag{1.9}$$

the conclusion of Theorem 1.2 holds true.

Remark 1.4. It is easy to see that (1.8) on the partial derivative of pressure does not seem comparable with (1.4)–(1.6) on the total pressure at least there is no inclusion relation between them. Moreover, (1.9) can be viewed as an extension of (1.4)–(1.6). Since the growth condition (1.8) is about $\partial_3 \pi$ and (1.9) is on the Besov space $\dot{B}_{\infty, \infty}^{-1}$, the methods in previous results in [17–19] are not available any more. In order to come over the additional difficulty, the rigorous analysis due to the new structure in nonlinear terms of the system (1.1)-(1.2) is made and some anisotropic function inequalities are also employed.

2. Preliminaries

Throughout this paper, we denote by C the constant which may change from line to line. We also use the following usual function spaces $L^p(\mathbb{R}^3)$, $1 \leq p \leq \infty$, denote the Lebesgue space of all L^p integral functions associated with the norm

$$\|f\|_{L^p} = \begin{cases} (\int_{\mathbb{R}^3} |f(x)|^p dx)^{1/p}, & 1 \leq p < \infty, \\ \text{ess sup}_{x \in \mathbb{R}^3} |f(x)|, & p = \infty. \end{cases} \quad (2.1)$$

$H^k(\mathbb{R}^3)$ denotes the Hilbert space $\{u \in L^2(\mathbb{R}^3); \|\nabla^k u\|_{L^2} < \infty\}$.

$\dot{B}_{p,q}^s(\mathbb{R}^3)$ denotes the homogeneous Besov space which is defined by the full-dyadic decomposition such as

$$\dot{B}_{p,q}^s(\mathbb{R}^3) = \left\{ f \in \frac{\mathcal{S}'(\mathbb{R}^3)}{\mathcal{P}(\mathbb{R}^3)} : \|f\|_{\dot{B}_{p,q}^s} < \infty \right\}, \quad (2.2)$$

where

$$\|f\|_{\dot{B}_{p,q}^s} = \begin{cases} \left(\sum_{j=-\infty}^{\infty} 2^{jsq} \|\Delta_j f\|_{L^p}^q \right)^{1/q}, & 1 \leq q < \infty, \\ \sup_{j \in \mathbb{Z}} 2^{js} \|\Delta_j f\|_{L^p}, & q = \infty, \end{cases} \quad (2.3)$$

and $\mathcal{S}'(\mathbb{R}^3)$, $\mathcal{P}(\mathbb{R}^3)$ are the spaces of all tempered distributions on \mathbb{R}^3 and the set of all scalar polynomials defined on \mathbb{R}^3 , respectively. It is worthy to note that the homogeneous Besov space $\dot{B}_{2,2}^s(\mathbb{R}^3)$ is equivalent to the Hilbert space $H^s(\mathbb{R}^3)$.

In order to prove the main results, we need to recall some anisotropic inequalities. The following Sobolev inequality is due to Cao and Wu [26].

Lemma 2.1 (Cao and Wu [26]). *Let the three constants $1 \leq \mu, \lambda, \gamma < \infty$ satisfy*

$$1 + \frac{3}{\gamma} = \frac{1}{\mu} + \frac{2}{\lambda}. \quad (2.4)$$

Assume $\phi \in H^1(\mathbb{R}^3)$, $\partial_1 \phi, \partial_2 \phi \in L^\lambda(\mathbb{R}^3)$ and $\partial_3 \phi \in L^\mu(\mathbb{R}^3)$, then there exists a constant $C = C(\mu, \lambda)$ such that

$$\|\phi\|_{L^\gamma} \leq C \|\partial_1 \phi\|_{L^\lambda}^{1/3} \|\partial_2 \phi\|_{L^\lambda}^{1/3} \|\partial_3 \phi\|_{L^\mu}^{1/3}. \quad (2.5)$$

We also recall an interpolation inequality in Besov space due to Meyer [27].

Lemma 2.2 (Meyer [27]). *Let $\alpha > 0$, $\beta > 0$ and $1 \leq p < q < \infty$, we have for all $f \in \mathcal{S}(\mathbb{R}^3)$,*

$$\|f\|_{L^q} \leq C \|f\|_{\dot{B}_{\infty,\infty}^{-\alpha}}^{1-\theta} \|f\|_{\dot{B}_{p,p}^{\beta}}^{\theta} \quad \text{for } \beta = \alpha \left(\frac{q}{p} - 1 \right), \quad \theta = \frac{q}{p}. \quad (2.6)$$

The following local existence and uniqueness of the strong solution for 3D micropolar fluid equations obtained by Dong et al. [17].

Lemma 2.3 (Dong et al. [17]). *Assume $3 < p < \infty$ and $(u_0, w_0) \in L^p(\mathbb{R}^3)$ with $\nabla \cdot u_0 = 0$ in the sense of distributions. Then there exist a constant $T > 0$ and a unique strong solution (u, w) of the 3D micropolar fluid equations (1.1)-(1.2) such that*

$$u \in BC\left([0, T); L^p\left(\mathbb{R}^3\right)\right), t^{1/2}\nabla u \in BC\left([0, T); L^p\left(\mathbb{R}^3\right)\right). \tag{2.7}$$

3. Proof of Main Results

According to the local existence result in Lemma 2.3, it allows us to construct a unique L^4 -strong solution (u, w) of micropolar fluid equations (1.1)-(1.2) under the initial data $(u_0, w_0) \in L^2(\mathbb{R}^3) \cap L^4(\mathbb{R}^3)$. Moreover, the strong solution can be proved to exist on a maximal time interval using the standard local solution extension technique. For the notation simplicity, we may suppose that the maximal time interval is $[0, T)$. Thus, to prove the main results, it suffices to show that

$$\lim_{t \rightarrow T} (\|u(t)\|_{L^4} + \|w(t)\|_{L^4}) < \infty. \tag{3.1}$$

This boundedness will be obtained by the following steps.

3.1. Uniform Estimates of Solutions

In order to prove the main results, we first need some basic estimates of the solutions for the 3D micropolar fluid equations (1.1)-(1.2).

Taking the inner product of the second equation of (1.1) with $u|u|^2$ and integrating by parts, it follows that

$$\frac{1}{4} \frac{d}{dt} \|u\|_{L^4}^4 + \| |u| \nabla u \|_{L^2}^2 + \frac{1}{2} \| \nabla |u|^2 \|_{L^2}^2 = \int_{\mathbb{R}^3} (\nabla \times w) \cdot u|u|^2 dx - \int_{\mathbb{R}^3} u \cdot \nabla \pi |u|^2 dx, \tag{3.2}$$

where we have used the following facts:

$$\begin{aligned} \int_{\mathbb{R}^3} (u \cdot \nabla u) \cdot u|u|^2 dx &= 0, \\ (\nabla u) \cdot \left(\nabla (u|u|^2) \right) &= |\nabla u|^2 |u|^2 + \frac{1}{2} \left| \nabla |u|^2 \right|^2. \end{aligned} \tag{3.3}$$

Similarly, taking the inner product of the second equation of (1.1) with $|w|^2 w$ and integrating by parts yield

$$\frac{1}{4} \frac{d}{dt} \|w\|_{L^4}^4 + \| |w| \nabla w \|_{L^2}^2 + \frac{1}{2} \| |w| \nabla \cdot w \|_{L^2}^2 + \| \nabla |w|^2 \|_{L^2}^2 = \int_{\mathbb{R}^3} (\nabla \times u) \cdot w|w|^2 dx - 2\|w\|_{L^4}^4, \tag{3.4}$$

where we have used

$$\begin{aligned}
& \int_{\mathbb{R}^3} (\mathbf{u} \cdot \nabla w) \cdot w |w|^2 dx = 0, \\
& (\nabla w) \cdot (\nabla (w |w|^2)) = |\nabla w|^2 |w|^2 + \frac{1}{2} |\nabla |w|^2|^2, \\
& \int_{\mathbb{R}^3} (\nabla \cdot w) \nabla \cdot (w |w|^2) dx = \int_{\mathbb{R}^3} |\nabla \cdot w|^2 |w|^2 dx + \int_{\mathbb{R}^3} (\nabla \cdot w) w \cdot \nabla |w|^2 dx \\
& \geq \frac{1}{2} \int_{\mathbb{R}^3} |\nabla \cdot w|^2 |w|^2 dx - \frac{1}{2} \int_{\mathbb{R}^3} |\nabla |w|^2|^2 dx.
\end{aligned} \tag{3.5}$$

Combination of (3.2) and (3.4) derives

$$\begin{aligned}
& \frac{1}{4} \frac{d}{dt} (\|u\|_{L^4}^4 + \|w\|_{L^4}^4) + \| |u| \nabla u \|_{L^2}^2 + \frac{1}{2} \| |\nabla |u|^2 | \|_{L^2}^2 + \| |w| \nabla w \|_{L^2}^2 + 2 \| w \|_{L^4}^4 \\
& \leq \int_{\mathbb{R}^3} (\nabla \times w) \cdot u |u|^2 dx + \int_{\mathbb{R}^3} (\nabla \times u) \cdot w |w|^2 dx - \int_{\mathbb{R}^3} u \cdot \nabla \pi |u|^2 dx.
\end{aligned} \tag{3.6}$$

With the use of Hölder's inequality, Young's inequality, and integration by parts, the first two terms on the right-hand side of (3.6) is bounded by

$$\begin{aligned}
& \int_{\mathbb{R}^3} (\nabla \times w) \cdot u |u|^2 dx + \int_{\mathbb{R}^3} (\nabla \times u) \cdot w |w|^2 dx \\
& \leq \|w\|_{L^4} \| |u| \nabla u \|_{L^2} \|u\|_{L^4} + \|u\|_{L^4} \| |w| \nabla w \|_{L^2} \|w\|_{L^4} \\
& \leq C (\|u\|_{L^4}^4 + 2 \|w\|_{L^4}^4) + \frac{1}{2} (\| |w| \nabla w \|_{L^2}^2 + \| |u| \nabla u \|_{L^2}^2).
\end{aligned} \tag{3.7}$$

Inserting (3.7) into (3.6), we obtain

$$\begin{aligned}
& \frac{d}{dt} (\|u\|_{L^4}^4 + \|w\|_{L^4}^4) + 2 \| |u| \nabla u \|_{L^2}^2 + 2 \| |\nabla |u|^2 | \|_{L^2}^2 \\
& \leq C (\|u\|_{L^4}^4 + \|w\|_{L^4}^4) + C \left| \int_{\mathbb{R}^3} u \cdot \nabla \pi |u|^2 dx \right| \\
& := C (\|u\|_{L^4}^4 + \|w\|_{L^4}^4) + I.
\end{aligned} \tag{3.8}$$

In order to estimate I , we first establish some estimates between the pressure and the velocity. Taking the operator div on both sides of the second equation of (1.1) gives

$$\begin{aligned} \pi &= (-\Delta)^{-1} \sum_{i,j=1}^3 \frac{\partial^2}{\partial x_i \partial x_j} (u_i u_j), \\ \nabla \pi &= (-\Delta)^{-1} \sum_{i,j=1}^3 \frac{\partial^2}{\partial x_i \partial x_j} (\nabla (u_i u_j)). \end{aligned} \tag{3.9}$$

Applying L^p ($1 < p < \infty$) boundedness of the singular integral operators yields

$$\|\pi\|_{L^p} \leq C \|u\|_{L^p}^2, \quad \|\nabla \pi\|_{L^p} \leq C \|u \cdot \nabla u\|_{L^p}. \tag{3.10}$$

Then employing the integration by parts and Hölder's inequality, we estimate I as

$$\begin{aligned} I &= C \left| \int_{\mathbb{R}^3} u \cdot \nabla \pi |u|^2 dx \right| \\ &= C \left| \int_{\mathbb{R}^3} \pi (u \cdot \nabla |u|^2) dx \right| \\ &\leq C \int_{\mathbb{R}^3} |\pi u| |\nabla |u|^2| dx \\ &\leq C \|\pi u\|_{L^2} \|\nabla |u|^2\|_{L^2} \\ &\leq C \|\pi\|_{L^4} \|u\|_{L^4} \|\nabla |u|^2\|_{L^2}. \end{aligned} \tag{3.11}$$

Now we begin to prove Theorems 1.2 and 1.3, respectively.

3.2. Proof of Theorem 1.2

Applying Lemma 2.1 to $\|\pi\|_{L^4}$ by choosing $\gamma = 4, \mu = q$ such that

$$\frac{1}{q} + \frac{2}{\lambda} = \frac{7}{4}. \tag{3.12}$$

Thus, we have

$$\|\pi\|_{L^4} \leq C \|\partial_1 \pi\|_{L^\lambda}^{1/3} \|\partial_2 \pi\|_{L^\lambda}^{1/3} \|\partial_3 \pi\|_{L^q}^{1/3} \leq C \|\nabla \pi\|_{L^\lambda}^{2/3} \|\partial_3 \pi\|_{L^q}^{1/3}. \tag{3.13}$$

Thanks to $12/7 < q \leq \infty$, then

$$\frac{8}{7} \leq \lambda < \frac{7}{6}. \tag{3.14}$$

Using (3.10) together with Hölder's inequality and Gagliardo-Nirenberg inequality derives that

$$\begin{aligned}
\|\nabla \pi\|_{L^\lambda} &\leq C\|u \cdot \nabla u\|_{L^\lambda} \leq C\|u\|_{L^{2\lambda/(2-\lambda)}} \|\nabla u\|_{L^2} \\
&= C\| |u|^2 \|_{L^{\lambda/(2-\lambda)}}^{1/2} \|\nabla u\|_{L^2} \\
&\leq C\| |u|^2 \|_{L^2}^{(3/\lambda)-(7/4)} \|\nabla |u|^2 \|_{L^2}^{(9/4)-(3/\lambda)} \|\nabla u\|_{L^2} \\
&\leq C\|u\|_{L^4}^{(6/\lambda)-(7/2)} \|\nabla |u|^2 \|_{L^2}^{(9/4)-(3/\lambda)} \|\nabla u\|_{L^2}.
\end{aligned} \tag{3.15}$$

Plugging (3.13) and (3.15) into (3.11) and employing Young's inequality yield that

$$\begin{aligned}
I &\leq C\|\partial_3 \pi\|_{L^q}^{1/3} \|\nabla u\|_{L^2}^{2/3} \|u\|_{L^4}^{(4/\lambda)-(4/3)} \|\nabla |u|^2 \|_{L^2}^{(5/2)-(2/\lambda)} \\
&\leq \frac{1}{2} \|\nabla |u|^2 \|_{L^2}^2 + C\|\partial_3 \pi\|_{L^q}^{4\lambda/3(4-\lambda)} \|\nabla u\|_{L^2}^{8\lambda/(12-7\lambda)} \|u\|_{L^4}^{16(3-\lambda)/3(4-\lambda)} \\
&\leq \frac{1}{2} \|\nabla |u|^2 \|_{L^2}^2 + C\left(\|\partial_3 \pi\|_{L^q}^{4\lambda/(12-7\lambda)} + \|\nabla u\|_{L^2}^2\right) \|u\|_{L^4}^{16(3-\lambda)/3(4-\lambda)} \\
&\leq \frac{1}{2} \|\nabla |u|^2 \|_{L^2}^2 + C\left(\|\partial_3 \pi\|_{L^q}^p + \|\nabla u\|_{L^2}^2\right) \left(\|u\|_{L^4}^4 + 1\right),
\end{aligned} \tag{3.16}$$

where we have used the following facts in the last line:

$$\frac{16(3-\lambda)}{3(4-\lambda)} < 4, \quad \frac{2}{p} + \frac{3}{q} = \frac{7}{4}, \quad \text{that is } p = \frac{4\lambda}{12-7\lambda}. \tag{3.17}$$

Then inserting (3.16) into (3.8) obtains

$$\begin{aligned}
&\frac{d}{dt} \left(\|u\|_{L^4}^4 + \|w\|_{L^4}^4 \right) + \|\nabla |u|^2 \|_{L^2}^2 \\
&\leq C\left(1 + \|\partial_3 \pi\|_{L^q}^p + \|\nabla u\|_{L^2}^2\right) \left(\|u\|_{L^4}^4 + 1\right) + C\left(\|u\|_{L^4}^4 + \|w\|_{L^4}^4\right),
\end{aligned} \tag{3.18}$$

or

$$\begin{aligned}
&\frac{d}{dt} \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4 \right) + \|\nabla |u|^2 \|_{L^2}^2 \\
&\leq C\left(1 + \|\partial_3 \pi\|_{L^q}^p + \|\nabla u\|_{L^2}^2\right) \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4\right) \\
&\leq C\|\partial_3 \pi\|_{L^q}^p \left(\|u\|_{L^4}^4 + \|w\|_{L^4}^4\right) + C\left(1 + \|\nabla u\|_{L^2}^2\right) \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4\right).
\end{aligned} \tag{3.19}$$

Denote

$$J = \|\partial_3 \pi\|_{L^q}^p \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4 \right), \tag{3.20}$$

and thanks to

$$\begin{aligned} 1 + \ln(e + \|w\|_{L^4}) &\leq 1 + \ln\left(e + \frac{1}{4}\|w\|_{L^4}^4 + \frac{4}{3}\right) \\ &\leq 1 + \ln\frac{1}{2}\left(e + \|w\|_{L^4}^4\right) \leq 1 + \ln\frac{1}{2} + \ln\left(e + \|w\|_{L^4}^4\right) \\ &\leq 1 + \ln\left(e + \|w\|_{L^4}^4\right), \end{aligned} \tag{3.21}$$

we have for J

$$\begin{aligned} J &= \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4})} \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4 \right) [1 + \ln(e + \|w\|_{L^4})] \\ &\leq \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4})} \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4 \right) [1 + \ln(e + \|w\|_{L^4}^4)] \\ &\leq \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4})} \left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4 \right) [1 + \ln(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4)]. \end{aligned} \tag{3.22}$$

Thus, plugging (3.22) into (3.19) and denoting

$$\mathcal{K} = e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4, \tag{3.23}$$

we have

$$\frac{d}{dt} \mathcal{K} \leq C \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4})} \{1 + \ln \mathcal{K}\} \mathcal{K} + C \left(1 + \|\nabla u\|_{L^2}^2\right) \mathcal{K}. \tag{3.24}$$

Taking Gronwall's inequality into consideration derives

$$\begin{aligned} \mathcal{K} &\leq \ln\left(e + \|u_0\|_{L^4}^4 + \|w_0\|_{L^4}^4\right) \\ &\quad \times \exp\left\{C \int_0^t \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4})} \{1 + \ln \mathcal{K}\} ds\right\} \\ &\quad \times \exp\left\{\int_0^t C \left(1 + \|\nabla u\|_{L^2}^2\right) ds\right\}. \end{aligned} \tag{3.25}$$

Thanks to (u, w) is a weak solution of the 3D micropolar equations (1.1)-(1.2), then

$$\int_0^T C(1 + \|\nabla u\|_{L^2}^2) ds < CT, \quad (3.26)$$

we rewrite (3.25) as

$$\begin{aligned} 1 + \ln \mathcal{K} &\leq CT + \ln(e + \|u_0\|_{L^4}^4 + \|w_0\|_{L^4}^4) \\ &+ C \int_0^t \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4}^4)} \{1 + \ln \mathcal{K}\} ds. \end{aligned} \quad (3.27)$$

Applying Gronwall's inequality again yields

$$\ln \mathcal{K} \leq C(u_0, w_0, T) \exp \left\{ C \int_0^T \frac{\|\partial_3 \pi\|_{L^q}^p}{1 + \ln(e + \|w\|_{L^4}^4)} ds \right\}, \quad (3.28)$$

which implies that

$$\lim_{t \rightarrow T} (\|u(t)\|_{L^4} + \|w(t)\|_{L^4}) < \infty. \quad (3.29)$$

Hence, we complete the proof of Theorem 1.2.

3.3. Proof of Theorem 1.3

Applying Lemma 2.2 to $\|\pi\|_{L^4}$ by choosing

$$q = 2, \quad p = 4, \quad \alpha = 1, \quad \beta = 1, \quad \theta = \frac{1}{2}, \quad (3.30)$$

we have

$$\begin{aligned} \|\pi\|_{L^4} &\leq C \|\pi\|_{\dot{B}_{\infty, \infty}^{-1}}^{1/2} \|\pi\|_{\dot{B}_{2, 2}^1}^{1/2} \\ &\leq C \|\pi\|_{\dot{B}_{\infty, \infty}^{-1}}^{1/2} \|\pi\|_{\dot{H}^1}^{1/2} = C \|\pi\|_{\dot{B}_{\infty, \infty}^{-1}}^{1/2} \|\nabla \pi\|_{L^2}^{1/2}. \end{aligned} \quad (3.31)$$

Plugging the above inequality into (3.11) together with Young's inequality implies

$$\begin{aligned} I &\leq C \|\pi\|_{L^4} \|u\|_{L^4} \left\| |\nabla u|^2 \right\|_{L^2} \\ &\leq C \|\pi\|_{L^4}^2 \|u\|_{L^4}^2 + \frac{1}{2} \left\| |\nabla u|^2 \right\|_{L^2}^2 \\ &\leq C \|\pi\|_{\dot{B}_{\infty, \infty}^{-1}} \|\nabla \pi\|_{L^2} \|u\|_{L^4}^2 + \frac{1}{2} \left\| |\nabla u|^2 \right\|_{L^2}^2 \end{aligned}$$

$$\begin{aligned}
 &\leq C\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}\| |u|\nabla u\|_{L^2}\|u\|_{L^4}^2 + \frac{1}{2}\|\nabla|u|^2\|_{L^2}^2 \\
 &\leq C\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2\|u\|_{L^4}^4 + \frac{1}{2}\| |u|\nabla u\|_{L^2}^2 + \frac{1}{2}\|\nabla|u|^2\|_{L^2}^2.
 \end{aligned}
 \tag{3.32}$$

Then inserting (3.32) into (3.8) derives

$$\frac{d}{dt}\left(\|u\|_{L^4}^4 + \|w\|_{L^4}^4\right) \leq C\left(\|u\|_{L^4}^4 + \|w\|_{L^4}^4\right) + C\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2\|u\|_{L^4}^4.
 \tag{3.33}$$

Thanks to

$$\begin{aligned}
 \|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2\|u\|_{L^4}^4 &= \frac{C\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^4})}\|u\|_{L^4}^4 [1 + \ln(e + \|w\|_{L^4})] \\
 &\leq \frac{C\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^2})}\|u\|_{L^4}^4 [1 + \ln(e + \|w\|_{L^4}^4)] \\
 &\leq \frac{C\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^2})}\left(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4\right) [1 + \ln(e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4)],
 \end{aligned}
 \tag{3.34}$$

inserting the above inequality into (3.33) and, similarly, denoting

$$\mathcal{K} = e + \|u\|_{L^4}^4 + \|w\|_{L^4}^4,
 \tag{3.35}$$

then we have

$$\frac{d}{dt}\mathcal{K} \leq C\frac{\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^4})}\{1 + \ln \mathcal{K}\}\mathcal{K}.
 \tag{3.36}$$

Taking Gronwall, inequality into consideration derives

$$\mathcal{K} \leq \ln\left(e + \|u_0\|_{L^4}^4 + \|w_0\|_{L^4}^4\right) \exp\left\{C\int_0^t \frac{\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^4})}\{1 + \ln \mathcal{K}\} ds\right\},
 \tag{3.37}$$

or

$$1 + \ln \mathcal{K} \leq \ln\left(e + \|u_0\|_{L^4}^4 + \|w_0\|_{L^4}^4\right) + C\int_0^t \frac{\|\pi\|_{\dot{B}_{\infty,\infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^4})}\{1 + \ln \mathcal{K}\} ds.
 \tag{3.38}$$

Likewise, applying Gronwall's inequality again implies

$$\ln \mathcal{K} \leq C(u_0, w_0) \exp \left\{ C \int_0^T \frac{\|\pi\|_{B_{\infty, \infty}^{-1}}^2}{1 + \ln(e + \|w\|_{L^4})} ds \right\}. \quad (3.39)$$

Hence, we derive that

$$\lim_{t \rightarrow T} (\|u(t)\|_{L^4} + \|w(t)\|_{L^4}) < \infty, \quad (3.40)$$

which completes the proof of Theorem 1.3.

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