

**EXPANSION OF A CLASS OF FUNCTIONS INTO AN INTEGRAL INVOLVING ASSOCIATED LEGENDRE FUNCTIONS**

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**ABSTRACT.** A theorem for expansion of a class of functions into an integral involving associated Legendre functions is obtained in this paper. This is a somewhat general integral expansion formula for a function  $f(x)$  defined in  $(x_1, x_2)$  where  $-1 < x_1 < x_2 < 1$ , which is perhaps useful in solving certain boundary value problems of mathematical physics and of elasticity involving conical boundaries.

**KEY WORDS AND PHRASES.** Integral expansion of a function, associated Legendre function, Mehler-Fok integral transform.

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## 1. INTRODUCTION.

Integral transforms are often used to solve the problems of mathematical physics involving linear partial differential equations and also other problems. Integral expansions involving spherical functions of a class of functions are known as Mehler-Fok type transforms. In these transform formulae, the subscript of the Legendre functions appear as the integration variable while its superscript is either zero or a fixed integer (see Sneddon [10]). There is another class of integral transforms involving associated Legendre functions somewhat related to the Mehler-Fok transforms, in which the superscript of the associated Legendre function appears in the integration formula while the subscript (complex) is kept fixed. Felsen [2] first developed this type of transform formulae involving  $P_{-1/2+i\tau}^{-\mu}(\cos \theta)$  as kernel where  $0 < \theta < \pi$  from a unique  $\delta$ -function representation. Later Mandal ([6], [7]) obtained somewhat similar types of two transform formulae from the solution of two appropriately designed boundary value problems. In the first type, the argument  $x$  of  $P_{-1/2+i\tau}^{-\mu}(x)$  ranges from  $-1$  to  $1$  while in the second, the argument  $z$  of  $P_{-1/2+i\tau}^{\mu}(z)$  ranges from  $1$  to  $\infty$ . Recently Mandal and Guha Roy [8] used a similar technique to establish another Mehler-Fok type integral transform formula involving  $P_{-1/2+i\tau}^{-\mu}(\cos \theta)$  as kernel ( $0 < \theta < \alpha$ ).

In the present paper, an integral expansion of a class of functions defined in  $(x_1, x_2)$  where  $-1 < x_1 < x_2 < 1$ , involving associated Legendre functions is obtained. Based on direct investigation of the properties of spherical functions, sufficient conditions which would establish the validity of this expansion formula for a wide class of functions are obtained in a manner

similar to the ideas used in ([3]-[5]). The main result is given in section 2 in the form of a theorem. Recently, we have used a similar technique to establish another type of integral representation [9] involving  $P_{-1/2+i\tau}^{-\mu}(\cosh \alpha)$  as kernel where  $0 < \alpha < \alpha_0$ .

**2. INTEGRAL EXPANSION OF A FUNCTION IN  $(x_1, x_2)$  WHERE  $-1 < x_1 < x_2 < 1$ .**

We present the main result of this paper in the form of the following theorem.

**THEOREM.** Let  $f(x)$  be a given function defined on the interval  $(x_1, x_2)$  where  $-1 < x_1 < x_2 < 1$  and satisfies the following conditions:

(1) The function  $f(x)$  is piecewise continuous and has a bounded variation in the open interval  $(x_1, x_2)$ .

(2) The function  $f(x)(1-x^2)^{-1} \ln(1-x^2)^{-1} \in L(x_1, x_2)$ ,  $-1 < x_1 < x_2 < 1$ .

Then we have

$$f(x) = \sum_k \sigma_k \Gamma\left(\frac{1}{2} + i\tau - i\sigma_k\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma_k\right) \frac{M(x, x_2; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)} F(\sigma_k) + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \sigma \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \frac{M(x, x_2; i\sigma)}{M(x_2, x_1; i\sigma)} F(\sigma) d\sigma \tag{2.1}$$

where

$$F(\sigma) = \int_{x_1}^{x_2} \frac{f(x)}{1-x^2} M(x, x_1; i\sigma) d\sigma, \tag{2.2}$$

$$-1 < x_1 < x_2 < 1, M(x, y; i\sigma) = P_{-1/2+i\tau}^{i\sigma}(x)P_{-1/2+i\tau}^{i\sigma}(-y) - P_{-1/2+i\tau}^{i\sigma}(-x)P_{-1/2+i\tau}^{i\sigma}(y)$$

and  $\sigma_k, \sigma, \tau$  are real. The equation (2.2) may be regarded as an integral transform of the function  $f(x)$  defined in  $(x_1, x_2)$  and (2.1) is its inverse. (2.1) and (2.2) together give the integral expansion of the function  $f(x)$ .

**PROOF OF THE EXPANSION THEOREM.** To prove this expansion theorem, we first note that the representation (cf. Erdélyi [1])

$$P_{-1/2+i\tau}^{i\sigma}(x) = \left(\frac{1+x}{1-x}\right)^{i\sigma/2} F\left(\frac{1}{2} + i\tau, \frac{1}{2} - i\tau; 1 - i\sigma; \frac{1-x}{2}\right) / \Gamma(1 - i\sigma),$$

$-1 < x_1 < x < x_2 < 1$ , where  $F(a, b; c; x)$  denotes the hypergeometric series, implies  $P_{-1/2+i\tau}^{i\sigma}(x)$  is continuous in the region defined by  $-1 < x_1 < x < x_2 < 1, -\infty < \sigma < \infty$  and satisfies the inequality

$$\left|P_{-1/2+i\tau}^{i\sigma}(x)\right| \leq \sqrt{sh\pi\sigma/\pi\sigma} P_{-1/2+i\tau}(x), \tag{2.3}$$

where the Legendre function  $P_{-1/2+i\tau}(x)$  is positive.

Using (2.3) it follows from (2.2) that

$$\int_{x_1}^{x_2} \left|\frac{f(x)}{1-x^2}\right| \left[P_{-1/2+i\tau}^{i\sigma}(x) P_{-1/2+i\tau}^{i\sigma}(-x_1) - P_{-1/2+i\tau}^{i\sigma}(-x) P_{-1/2+i\tau}^{i\sigma}(x_1)\right] dx \leq \sqrt{sh\pi\sigma/\pi\sigma} \int_{x_1}^{x_2} \frac{|f(x)|}{1-x^2} \left\{P_{-1/2+i\tau}(x) P_{-1/2+i\tau}(-x_1) - P_{-1/2+i\tau}(-x) P_{-1/2+i\tau}(x_1)\right\} dx,$$

and this shows that the conditions imposed on  $f(x)$  imply that the integral  $F(\sigma)$  is absolutely and uniformly convergent for  $\sigma \in [-T, T]$  where  $T$  is a positive large number. Hence  $F(\sigma)$  is continuous on  $[-T, T]$  and the repeated integral

$$J(x, T) = \frac{1}{2\pi i} \int_{-T}^T \sigma \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \frac{M(x, x_2; i\sigma)}{M(x_2, x_1; i\sigma)} d\sigma \cdot \int_{x_1}^{x_2} \frac{f(y)}{1-y^2} M(y, x_1; i\sigma) dy$$

is meaningful. Also, uniform convergence allows us to change the order of integration and write  $J(x, T)$  as

$$J(x, T) = \int_{x_1}^{x_2} \frac{f(y)}{1-y^2} K(x, y, T) dy. \quad (2.4)$$

where

$$K(x, y, T) = \frac{1}{2\pi i} \int_{-T}^T \sigma \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \frac{M(x, x_2; i\sigma)M(y, x_1; i\sigma)}{M(x_2, x_1; i\sigma)} d\sigma. \quad (2.5)$$

Now we shall show that the kernel  $K(x, y, T)$  is symmetric in the variables  $x$  and  $y$ . By definition, we have

$$K(x, y, T) - K(y, x, T) = \frac{1}{2\pi i} \int_{-T}^T \sigma \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \frac{1}{M(x_2, x_1; i\sigma)} \\ \cdot [M(x, x_2; i\sigma)M(y, x_1; i\sigma) - M(y, x_2; i\sigma)M(x, x_1; i\sigma)] d\sigma.$$

It follows from the properties of associated Legendre functions (cf. Erdélyi [1]) that the integrand in the above integral is an odd function of  $\sigma$ , hence the integral vanishes. Thus

$$K(y, x, T) = K(x, y, T). \quad (2.6)$$

To investigate the behavior of  $K(x, y, T)$  as  $T \rightarrow \infty$ , by writing  $\mu = -i\sigma$ , we write (2.5) as

$$K(x, y, T) = \frac{1}{2\pi i} \int_{-iT}^{iT} \mu \Gamma\left(\frac{1}{2} + i\tau + \mu\right) \Gamma\left(\frac{1}{2} - i\tau + \mu\right) \frac{M(x, x_2; -\mu)M(y, x_1; -\mu)}{M(x_2, x_1; -\mu)} d\mu. \quad (2.7)$$

Expression under the integral sign in (2.7) is analytic function for the complex variable  $\mu$  and it has no singularity in the semi-plane  $Re\mu \geq 0$ , except for simple poles at  $\mu = -i\sigma_k$  ( $k$  is positive integer) (cf. Felsen [2]), where

$$M(x_2, x_1; i\sigma_k) = 0, \quad \sigma_k > 0. \quad (2.8)$$

Completing the contour of integration on (2.7) with the arc  $\Gamma_T$  of radius  $T$  situated in the semi-plane  $Re\mu \geq 0$  and applying the residue theorem, we obtain

$$K(x, y, T) = K_1(x, y, T) - \sum_k \sigma_k \Gamma\left(\frac{1}{2} + i\tau - i\sigma_k\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma_k\right) \frac{M(x, x_2; i\sigma_k)M(y, x_1; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)}, \quad (2.9)$$

where

$$K_1(x, y, T) = \frac{1}{2\pi i} \int_{\Gamma_T} \mu \Gamma\left(\frac{1}{2} + i\tau + \mu\right) \Gamma\left(\frac{1}{2} - i\tau + \mu\right) \frac{M(x, x_2; -\mu)M(y, x_1; -\mu)}{M(x_2, x_1; -\mu)} d\mu. \quad (2.10)$$

Suppose that  $y \leq x$ . By virtue of the definition

$$P_{-1/2+i\tau}^{-\mu}(x) = \left(\frac{1+x}{1-x}\right)^{-\mu/2} \frac{1}{\Gamma(1+\mu)} [1 + O(|\mu|^{-1})], \\ P_{-1/2+i\tau}^{-\mu}(-x) = \left(\frac{1-x}{1+x}\right)^{-\mu/2} \frac{1}{\Gamma(1+\mu)} [1 + O(|\mu|^{-1})] \quad (2.11)$$

Using (2.11) and asymptotic properties of the gamma function for large  $\mu$ , we conclude that

$$\mu \Gamma\left(\frac{1}{2} + i\tau + \mu\right) \Gamma\left(\frac{1}{2} - i\tau + \mu\right) \frac{M(x, x_2; -\mu)M(y, x_1; -\mu)}{M(x_2, x_1; -\mu)}$$

$$= \frac{\left[ \left( \frac{1+x}{1-x} \frac{1-x_2}{1+x_2} \right)^{-\mu/2} - \left( \frac{1-x}{1+x} \frac{1+x_2}{1-x_2} \right)^{-\mu/2} \right] \left[ \left( \frac{1+y}{1-y} \frac{1-x_1}{1+x_1} \right)^{-\mu/2} - \left( \frac{1-y}{1+y} \frac{1+x_1}{1-x_1} \right)^{-\mu/2} \right]}{\left[ \left( \frac{1+x_2}{1-x_2} \frac{1-x_1}{1+x_1} \right)^{-\mu/2} - \left( \frac{1-x_2}{1+x_2} \frac{1+x_1}{1-x_1} \right)^{-\mu/2} \right]} \cdot [1 + O(|\mu|^{-1})]. \tag{2.12}$$

Now introduce the new variables

$$\xi = \frac{1}{2} \ell n \frac{1+x}{1-x}, \quad \eta = \frac{1}{2} \ell n \frac{1+y}{1-y}, \quad \alpha = \frac{1}{2} \ell n \frac{1+x_1}{1-x_1} \text{ and } \beta = \frac{1}{2} \ell n \frac{1+x_2}{1-x_2}.$$

Then, for large  $\mu$ , from (2.10) - (2.12) we obtain for  $y \leq x$

$$K_1(x, y, T) = \frac{1}{2\pi i} \int_T \left[ \exp\{-\mu(\xi - \eta)\} + \exp\{-\mu(2\beta - 2\alpha - \xi + \eta)\} \right. \\ \left. - \exp\{-\mu(\xi + \eta - 2\alpha)\} - \exp\{-\mu(2\beta - \xi - \eta)\} \right] d\mu \\ + O(1) \int_0^{\pi/2} \left[ \exp\{-\mu(\xi - \eta)\cos \varphi\} + \exp\{-\mu(2\beta - 2\alpha - \xi + \eta)\cos \varphi\} \right. \\ \left. - \exp\{-\mu(\xi + \eta - 2\alpha)\cos \varphi\} - \exp\{-\mu(2\beta - \xi - \eta)\cos \varphi\} \right] d\varphi,$$

for  $\alpha < \eta \leq \xi < \beta$ .

Using the identity

$$\frac{2}{\pi} \int_0^{\pi/2} \exp\{-\lambda T \cos \varphi\} d\varphi \leq \frac{1 - \exp(-\lambda T)}{\lambda T}, \quad \lambda \geq 0,$$

we obtain for  $y \leq x$ ,

$$K_1(x, y, T) = \frac{1}{\pi} \left[ \frac{\sin T(\xi - \eta)}{\xi - \eta} + \frac{\sin T(2\beta - 2\alpha - \xi + \eta)}{2\beta - 2\alpha - \xi + \eta} - \frac{\sin T(\xi + \eta - 2\alpha)}{\xi + \eta - 2\alpha} \right. \\ \left. - \frac{\sin T(2\beta - \xi - \eta)}{2\beta - \xi - \eta} \right] + O(1) \left[ \frac{1 - \exp\{-T(\xi - \eta)\}}{T(\xi - \eta)} + \frac{1 - \exp\{-T(2\beta - 2\alpha - \xi + \eta)\}}{T(2\beta - 2\alpha - \xi + \eta)} \right. \\ \left. - \frac{1 - \exp\{-T(\xi + \eta - 2\alpha)\}}{T(\xi + \eta - 2\alpha)} - \frac{1 - \exp\{-T(2\beta - \xi - \eta)\}}{T(2\beta - \xi - \eta)} \right], \quad \alpha < \eta \leq \xi < \beta, \tag{2.13}$$

where the factor  $O(1)$  is independent of  $y$ .

Again for  $y \geq x$ , we use the symmetry property (2.6) and the representation (2.10) of  $K_1(x, y, T)$  with the variables  $x, y$  replaced by  $y, x$ .

Now we write (2.4) as

$$J(x, T) = \int_{x_1}^x \frac{f(y)}{1-y^2} K_1(x, y, T) dy + \int_x^{x_2} \frac{f(y)}{1-y^2} K_1(x, y, T) dy \\ - \sum_k \sigma_k \left[ \left( \frac{1}{2} + i\tau - i\sigma_k \right) \Gamma\left( \frac{1}{2} - i\tau - i\sigma_k \right) \frac{M(x, x_2; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)} \int_{x_1}^{x_2} \frac{f(y)}{1-y^2} M(y, x_1; i\sigma_k) dy \right. \\ \left. = J_1(x, T) + J_2(x, T) - \sum_k \sigma_k \left[ \left( \frac{1}{2} + i\tau - i\sigma_k \right) \Gamma\left( \frac{1}{2} - i\tau - i\sigma_k \right) \frac{M(x, x_2; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)} \right. \right. \\ \left. \left. \times \int_{x_1}^{x_2} \frac{f(y)}{1-y^2} M(y, x_1; i\sigma_k) dy \right] \right] \tag{2.14}$$

Using (2.13) in  $J_1$ , we obtain

$$\begin{aligned}
 J_1(x, T) = & \frac{1}{\pi} \left[ \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(\xi - \eta)}{\xi - \eta} d\eta + \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(2\beta - 2\alpha - \xi + \eta)}{2\beta - 2\alpha - \xi + \eta} d\eta \right. \\
 & \left. - \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(\xi + \eta - 2\alpha)}{\xi + \eta - 2\alpha} d\eta - \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(2\beta - \xi - \eta)}{2\beta - \xi - \eta} d\eta \right] \\
 & + O(1) \left[ \int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(\xi - \eta)\}}{T(\xi - \eta)} d\eta \right. \\
 & + \int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(2\beta - 2\alpha - \xi + \eta)\}}{T(2\beta - 2\alpha - \xi + \eta)} d\eta \\
 & - \int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(\xi + \eta - 2\alpha)\}}{T(\xi + \eta - 2\alpha)} d\eta \\
 & \left. - \int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(2\beta - \xi - \eta)\}}{T(2\beta - \xi - \eta)} d\eta \right] \tag{2.15}
 \end{aligned}$$

The conditions satisfied by  $f(x)$  imply that  $f(\tanh \eta) \in L(\alpha, \beta)$ ; hence, by virtue of Dirichlet's theorem, for  $T \rightarrow \infty$

$$\begin{aligned}
 \frac{1}{\pi} \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(\xi - \eta)}{\xi - \eta} d\eta &= \frac{1}{2} f(\tanh \xi - o) + o(1) \\
 &= \frac{1}{2} f(x - o) + o(1),
 \end{aligned}$$

$$\frac{1}{\pi} \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(2\beta - 2\alpha - \xi + \eta)}{2\beta - 2\alpha - \xi + \eta} d\eta = o(1),$$

$$\frac{1}{\pi} \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(\xi + \eta - 2\alpha)}{\xi + \eta - 2\alpha} d\eta = o(1),$$

and

$$\frac{1}{\pi} \int_{\alpha}^{\xi} f(\tanh \eta) \frac{\sin T(2\beta - \xi - \eta)}{2\beta - \xi - \eta} d\eta = o(1).$$

Moreover, if the integral of integration is divided into the subintervals  $(\xi - \delta, \xi)$  and  $(\alpha, \xi - \delta)$  and if a sufficiently small positive  $\delta$  (implying a sufficiently large  $T$ ) is chosen, then we have

$$\begin{aligned}
 & \int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(\xi - \eta)\}}{T(\xi - \eta)} d\eta \\
 & \leq \frac{1}{\delta T} \int_{\alpha}^{\xi - \delta} |f(\tanh \eta)| d\eta + \int_{\xi - \delta}^{\xi} |f(\tanh \eta)| d\eta \\
 & = O(T^{-1}) + o(1) = o(1) \text{ for } T \rightarrow \infty, \\
 & \int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(2\beta - 2\alpha - \xi + \eta)\}}{T(2\beta - 2\alpha - \xi + \eta)} d\eta \leq \frac{1}{\xi T} \int_{\alpha}^{\xi} |f(\tanh \eta)| d\eta
 \end{aligned}$$

$$\begin{aligned}
 &= O(T^{-1}) = o(1) \text{ for } T \rightarrow \infty, \\
 &\int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(\xi + \eta - 2\alpha)\}}{T(\xi + \eta - 2\alpha)} d\eta \leq \frac{1}{\xi T} \int_{\alpha}^{\xi} |f(\tanh \eta)| d\eta \\
 &= O(T^{-1}) = o(1) \text{ for } T \rightarrow \infty, \\
 \text{and} \\
 &\int_{\alpha}^{\xi} |f(\tanh \eta)| \frac{1 - \exp\{-T(2\beta - \xi - \eta)\}}{T(2\beta - \xi - \eta)} d\eta \leq \frac{1}{\xi T} \int_{\alpha}^{\xi} |f(\tanh \eta)| d\eta \\
 &= O(T^{-1}) = o(1) \text{ for } T \rightarrow \infty. \tag{2.17}
 \end{aligned}$$

Thus (2.15) to (2.17) leads to

$$\lim_{T \rightarrow \infty} J_1(\tanh \xi, T) = \frac{1}{2} f(\tanh \xi - o) = \frac{1}{2} f(x - o). \tag{2.18}$$

Similarly,

$$\lim_{T \rightarrow \infty} J_2(\tanh \xi, T) = \frac{1}{2} f(\tanh \xi + o) = \frac{1}{2} f(x + o). \tag{2.19}$$

Hence,

$$\begin{aligned}
 \lim_{T \rightarrow \infty} J(x, T) &= \frac{1}{2} [f(x + o) + f(x - o)] - \sum_k \sigma_k \left[ \Gamma\left(\frac{1}{2} + i\tau - i\sigma_k\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma_k\right) \right. \\
 &\quad \left. \cdot \frac{M(x, x_2; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)} F(\sigma_k) \right]. \tag{2.20}
 \end{aligned}$$

Thus, at the points of continuity of  $f(x)$  we obtain (2.1). We note that (2.1) becomes a result in [5] when  $x_1 = -1$  and  $x_2 = 1$ .

It follows from the foregoing theorem that, at points of continuity of  $f(x)$ , we have

$$\begin{aligned}
 f(x) &= \sum_k \sigma_k \left[ \Gamma\left(\frac{1}{2} + i\tau - i\sigma_k\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma_k\right) \frac{R(x, x_2; i\sigma_k)}{(\partial^2/\partial x_2 \partial \sigma_k)R(x_2, x_1; i\sigma_k)} F(\sigma_k) \right. \\
 &\quad \left. + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \sigma \left[ \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \frac{R(x, x_2; i\sigma)}{(\partial/\partial x_2)R(x_2, x_1; i\sigma)} F(\sigma) d\sigma \right], \tag{2.21}
 \end{aligned}$$

where

$$F(\sigma) = \int_{x_1}^{x_2} \frac{f(x)}{1 - x^2} R(x, x_1; i\sigma) dx, \quad -1 < x_1 < x_2 < 1, \tag{2.22}$$

$$R(x, y; i\sigma) = P_{-1/2 + i\tau}^{i\sigma}(x) \frac{\partial}{\partial y} P_{-1/2 + i\tau}^{i\sigma}(-y) - P_{-1/2 + i\tau}^{i\sigma}(-x) \frac{\partial}{\partial y} P_{-1/2 + i\tau}^{i\sigma}(y)$$

and  $\sigma_k$ 's,  $\sigma, \tau$  are real.

The integrand in (2.21) has singularities at  $\sigma = \sigma_k$  ( $k$  is positive integers) which are simple poles along the positive  $\sigma$ -axis, where

$$\frac{\partial}{\partial x_2} R(x, x_1; i\sigma_k) = 0, \quad (\sigma_k > 0). \tag{2.23}$$

To prove (2.21) we use the following asymptotic formulas for large  $\mu$ :

$$\begin{aligned}
 \frac{\partial}{\partial x} P_{-1/2 + i\tau}^{-\mu}(x) &= -\frac{\mu}{\Gamma(1 + \mu)} \frac{1}{(1 - x)(1 + x)} \left(\frac{1 + x}{1 - x}\right)^{-\mu/2} [1 + O(|\mu|^{-1})], \\
 \frac{\partial}{\partial x} P_{-1/2 + i\tau}^{-\mu}(-x) &= -\frac{\mu}{\Gamma(1 + \mu)} \frac{1}{(1 + x)(1 - x)} \left(\frac{1 - x}{1 + x}\right)^{-\mu/2} [1 + O(|\mu|^{-1})], \tag{2.24}
 \end{aligned}$$

The proof of (2.21) is similar to the proof in the section 2, and we do not reproduce it. We note that (2.21) becomes a result in [5] when  $x_1 = -1$  and  $x_2 = 1$ .

**3. EXAMPLES.**

We now give examples of expansions of some functions.

$$(1) (1-x^2)^{\nu/2} = \sum_k \sigma_k \Gamma\left(\frac{1}{2} + i\tau - i\sigma_k\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma_k\right) \frac{M(x, x_2; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)} \\ \cdot \frac{2^\nu \Gamma(1+\nu)}{(\nu^2 + \sigma_k^2)} (\nu + i\sigma_k) [P_{\nu}^{-\nu}(x_1)M_1(x_1, x_1; i\sigma_k) - P_{\nu}^{-\nu}(x_2) M_1(x_2, x_1; i\sigma_k)] \\ + \frac{2^\nu \Gamma(1+\nu)}{2\pi i} \int_{-\infty}^{\infty} \frac{\sigma(\nu + i\sigma)}{\nu^2 + \sigma^2} \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \frac{M(x, x_2; i\sigma)}{M(x_2, x_1; i\sigma)} \\ \cdot [P_{\nu}^{-\nu}(x_1)M_1(x_1, x_1; i\sigma) - P_{\nu}^{-\nu}(x_2) M_1(x_2, x_1; i\sigma)] d\sigma,$$

$$(-1 < x_1 < x < x_2 < 1)$$

where

$$M(x, y; i\sigma) = P_{\nu}^{i\sigma}(x) P_{\nu}^{i\sigma}(-y) - P_{\nu}^{i\sigma}(-x) P_{\nu}^{i\sigma}(y),$$

$$M_1(x, y; i\sigma) = P_{\nu-1}^{i\sigma}(x) P_{\nu}^{i\sigma}(-y) - P_{\nu-1}^{i\sigma}(-x) P_{\nu}^{i\sigma}(y) \text{ and } \nu = -1/2 + i\tau.$$

$$(2) P_{\nu}^{\mu}(x) = \sum_k \sigma_k \Gamma\left(\frac{1}{2} + i\tau - i\sigma_k\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma_k\right) \frac{M(x, x_2; i\sigma_k)}{(\partial/\partial\sigma_k)M(x_2, x_1; i\sigma_k)} \frac{1}{(\mu^2 + \sigma_k^2)} \\ \cdot [(\nu + \mu) P_{\nu-1}^{\mu}(x_2)M(x_2, x_1; i\sigma_k) + (\nu + i\sigma_k) \{P_{\nu}^{\mu}(x_1)M_1(x_1, x_2; i\sigma_k) \\ - P_{\nu}^{\mu}(x_2)M_1(x_2, x_2; i\sigma_k)\}] + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\sigma}{\mu^2 + \sigma^2} \Gamma\left(\frac{1}{2} + i\tau - i\sigma\right) \Gamma\left(\frac{1}{2} - i\tau - i\sigma\right) \\ \cdot \frac{M(x, x_2; i\sigma)}{M(x_2, x_1; i\sigma)} [(\nu + \mu) P_{\nu-1}^{\mu}(x_2)M(x_2, x_1; i\sigma) + (\nu + i\sigma) \{P_{\nu}^{\mu}(x_1) \\ \cdot M_1(x_1, x_2; i\sigma) - P_{\nu}^{\mu}(x_2)M_1(x_2, x_1; i\sigma)\}] d\sigma.$$

In all these results the conditions under which the expansion theorem hold are satisfied.

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