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Research Article

Necessary and Sufficient Condition for Stability of Generalized Expectation Value

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A class of generalized definitions of expectation value is often employed in nonequilibrium statistical mechanics for complex systems. Here, the necessary and sufficient condition is presented for such a class to be stable under small deformations of a given arbitrary probability distribution.

Given a probability distribution $\{p_i\}_{i=1,2,\dots,W}$, that is, $0 \leq p_i \leq 1$ ($i = 1, 2, \dots, W$) and $\sum_{i=1}^W p_i = 1$, the ordinary expectation value of a quantity Q of a system under consideration is defined by $\sum_{i=1}^W p_i Q_i$, where W is the total number of accessible states and is enormously large in statistical mechanics, typically being $2^{10^{23}}$. In the field of generalized statistical mechanics for complex systems, on the other hand, discussions are often made about altering this definition. Among others, the so-called “escort average” is widely employed in the field of generalized statistical mechanics [1–3]. It is defined as follows:

$$\langle Q \rangle_{\phi} [p] = \sum_{i=1}^W P_i^{(\phi)} Q_i, \quad (1)$$

where $P_i^{(\phi)}$ stands for the escort probability distribution [4] given by

$$P_i^{(\phi)} = \frac{\phi(p_i)}{\sum_{j=1}^W \phi(p_j)}, \quad (2)$$

with a nonnegative function ϕ . In the special case when $\phi(x) = x$, $\langle Q \rangle_{\phi}$ is reduced to the ordinary expectation value mentioned above.

Consider measurements of a certain quantity of a system to obtain information about the probability distribution. Repeated measurements should be performed on the system, which is identically prepared each time. Suppose that two probability distributions, $\{p_i\}_{i=1,2,\dots,W}$ and $\{p'_i\}_{i=1,2,\dots,W'}$, are obtained through the measurements. They may slightly be different from each other, in general. If such measurements make sense, then the expectation values, $\langle Q \rangle[p]$ and $\langle Q \rangle[p']$, calculated from these two distributions should also be close to each other. This condition, which implies "experimental robustness," is represented as follows.

Definition (stability). An expectation value $\langle Q \rangle[p]$ is said to be stable, if the following predicate holds for any pair of probability distributions, $\{p_i\}_{i=1,2,\dots,W}$ and $\{p'_i\}_{i=1,2,\dots,W'}$:

$$(\forall \varepsilon > 0) (\exists \delta > 0) (\forall W) \quad (\|p - p'\|_1 < \delta \implies |\langle Q \rangle[p] - \langle Q \rangle[p']| < \varepsilon). \quad (3)$$

Here, $\|p - p'\|_1 = \sum_{i=1}^W |p_i - p'_i|$ is the l^1 -norm describing the distance between these two probability distributions. One might consider norms of other kinds, but what is physically relevant to discrete systems is the present l^1 -norm [5]. Equation (3) is analogous to Lesche's stability condition on entropic functionals [5], which has recently been revisited in the literature [6–11] (note that the discussion in [8] is corrected in [9]). This concept of stability is actually equivalent to that of uniform continuity.

In recent papers [12, 13], it has been shown that the generalized expectation value in (1) with a specific class, $\phi(x) = x^q$ ($q > 0$), (the associated expectation value being termed the q -expectation value), is not stable unless $q = 1$. This result needs the q -expectation-value formalism of nonextensive statistical mechanics [1, 2] be reconsidered. In addition, the result is supported by Boltzmann-like kinetic theory in an independent manner [14].

Here, it seems appropriate to make some comments on the latest situation of the problems concerning stabilities of entropic functionals and generalized expectation values. The authors of [15, 16] have presented discussions which aim to rescue the q -expectation values from the difficulties of their instability pointed out in [12]. Those authors insist that the q -expectation values can be stable in both the finite- W and continuous cases. Such possibilities are, however, fully refuted by the work in [13] both physically and mathematically, and the controversy seems to have been terminated with that work. The case of the continuous variables has further been carefully examined in a recent paper [17], where the so-called Tsallis q -entropies [1, 2] do not have the continuous limit in consistency with the physical principles such as the thermodynamic laws (see also [18, 19]). These controversies have led the researchers to give up the traditional form of nonextensive statistical mechanics based on the q -entropies and q -expectation values and to examine other entropic functionals combined with the ordinary definition of expectation values [20] (see also [21, 22]). Thus, it seems that nonextensive statistical mechanics has to be fully reexamined, theoretically.

In this paper, we present the necessary and sufficient condition for $\langle Q \rangle_\phi[p]$ in (1) to be stable.

Our main result is as follows.

Theorem. Let ϕ be nonnegative and continuous on $[0, 1]$, differentiable on $(0, 1)$, and satisfy the condition that $\phi(x) = 0 \iff x = 0$. And, let $Q = \{Q_i\}_{i=1,2,\dots,W}$ be a random variable. Then, $\langle Q \rangle_\phi[p]$ in (1) is stable, if and only if $\lim_{x \rightarrow +0} \phi(x)/x \in (0, \infty)$.

Proof. First, assume that $\lim_{x \rightarrow +0} \phi(x)/x = a > 0$. Then, there exists $\delta_1 > 0$ such that

$$a - \frac{a}{2} < \frac{\phi(x)}{x} < a + \frac{a}{2} \quad (\forall x \in (0, \delta_1]). \quad (4)$$

$\phi(x)/x$ does not vanish because of the condition $\phi(x) = 0 \Leftrightarrow x = 0$. Therefore, there exists $b > 0$ such that

$$\frac{\phi(x)}{x} \geq b \quad (\forall x \in (\delta_1, 1]). \quad (5)$$

Putting $c = \min\{a/2, b\}$ we have

$$cx \leq \phi(x) \quad (\forall x \in [0, 1]). \quad (6)$$

Consequently, for an arbitrarily large W and an arbitrary probability distribution $\{p_i\}_{i=1,2,\dots,W}$, we obtain

$$\frac{1}{\sum_{i=1}^W \phi(p_i)} \leq c. \quad (7)$$

From the mean value theorem, it follows that

$$|\phi(p_i) - \phi(p'_i)| \leq |p_i - p'_i| \cdot \sup_{x \in (0,1)} |\phi'(x)|, \quad (8)$$

where $\phi'(x)$ is the derivative of $\phi(x)$ with respect to x . For $\varepsilon > 0$, we put

$$\delta = \inf \left(\delta_1, \frac{c\varepsilon}{2|Q_{\max}| \cdot \left(\sup_{x \in (0,1)} |\phi'(x)| \right)} \right), \quad (9)$$

where $Q_{\max} = \max \{Q_i\}_{i=1,2,\dots,W}$. Now, for $\|p - p'\|_1 < \delta$, we have

$$\begin{aligned} & \left| \langle Q \rangle_{\phi} [p] - \langle Q \rangle_{\phi} [p'] \right| \\ &= \frac{1}{\sum_{i=1}^W \phi(p_i) \sum_{j=1}^W \phi(p'_j)} \left| \sum_{i=1}^W Q_i \left\{ \phi(p_i) \sum_{j=1}^W \phi(p'_j) - \phi(p'_i) \sum_{j=1}^W \phi(p_j) \right\} \right| \\ &\leq \frac{1}{\sum_{i=1}^W \phi(p_i) \sum_{j=1}^W \phi(p'_j)} \\ &\quad \times \left[\sum_{i=1}^W |Q_i| \left\{ |\phi(p_i) - \phi(p'_i)| \sum_{j=1}^W \phi(p'_j) + \phi(p'_i) \left| \sum_{j=1}^W \phi(p_j) - \sum_{j=1}^W \phi(p'_j) \right| \right\} \right] \end{aligned}$$

$$\begin{aligned}
&\leq \frac{1}{\sum_{j=1}^W \phi(p_j)} \sum_{i=1}^W |Q_i| |\phi(p_i) - \phi(p'_i)| \\
&\quad + \frac{\sum_{j=1}^W |\phi(p_j) - \phi(p'_j)|}{\sum_{i=1}^W \phi(p_i) \sum_{j=1}^W \phi(p'_j)} \sum_{i=1}^W |Q_i| \phi(p'_i) \\
&\leq \frac{2|Q_{\max}|}{\sum_{j=1}^W \phi(p_j)} \sum_{i=1}^W |\phi(p_i) - \phi(p'_i)| \\
&\leq \frac{2|Q_{\max}|}{\sum_{j=1}^W \phi(p_j)} \|p - p'\|_1 \cdot \sup_{x \in (0,1)} |\phi'(x)| \\
&\leq \frac{2|Q_{\max}|}{c} \|p - p'\|_1 \cdot \sup_{x \in (0,1)} |\phi'(x)| \\
&< \varepsilon.
\end{aligned} \tag{10}$$

Therefore, $\langle Q \rangle_\phi [p]$ is stable.

On the other hand, suppose that $\lim_{x \rightarrow +0} \phi(x)/x \notin (0, \infty)$. That is, (i) $\lim_{x \rightarrow +0} \phi(x)/x = 0$ or (ii) $\lim_{x \rightarrow +0} \phi(x)/x = \infty$. Below, we will examine these cases separately.

(i) Consider the following deformation:

$$\begin{aligned}
p_i &= \frac{1}{W-1} (1 - \delta_{i1}), \\
p'_i &= \left(1 - \frac{\delta}{2}\right) p_i + \frac{\delta}{2} \delta_{i1},
\end{aligned} \tag{11}$$

which are normalized and satisfy $\|p - p'\|_1 = \delta$. We have

$$\begin{aligned}
\sum_{i=1}^W \phi(p_i) &= (W-1) \phi\left(\frac{1}{W-1}\right), \\
\sum_{i=1}^W \phi(p'_i) &= \phi\left(\frac{\delta}{2}\right) + (W-1) \phi\left(\frac{1}{W-1} \left(1 - \frac{\delta}{2}\right)\right).
\end{aligned} \tag{12}$$

Difference of the expectation values is calculated as follows:

$$\begin{aligned}
&\langle Q \rangle_\phi [p] - \langle Q \rangle_\phi [p'] \\
&= -\frac{Q_1 \phi(\delta/2)}{\phi(\delta/2) + (W-1) \phi((1/(W-1))(1 - \delta/2))} \\
&\quad + \left(\sum_{i=2}^W Q_i \right) \left\{ \frac{1}{W-1} - \frac{\phi((1/(W-1))(1 - \delta/2))}{\phi(\delta/2) + (W-1) \phi((1/(W-1))(1 - \delta/2))} \right\}
\end{aligned}$$

$$\begin{aligned}
&= \frac{W}{W-1} (\bar{Q} - Q_1) \\
&\quad \times \frac{\phi(\delta/2)/(1-\delta/2)}{\phi(\delta/2)/(1-\delta/2) + \phi((1/(W-1))(1-\delta/2))/[(1/(W-1))(1-\delta/2)]} \\
&\xrightarrow{W \rightarrow \infty} \bar{Q} - Q_1,
\end{aligned} \tag{13}$$

since $\lim_{x \rightarrow +0} \phi(x)/x = 0$, where \bar{Q} is the arithmetic mean, $\bar{Q} = \sum_{i=1}^W Q_i/W$. Therefore, $\langle Q \rangle_\phi [p]$ is not stable.

(ii) Consider the following deformation:

$$\begin{aligned}
p_i &= \delta_{i1}, \\
p'_i &= \left(1 - \frac{\delta}{2} \frac{W}{W-1}\right) p_i + \frac{\delta}{2} \frac{1}{W-1},
\end{aligned} \tag{14}$$

which are also normalized and satisfy $\|p - p'\|_1 = \delta$. We have

$$\begin{aligned}
\sum_{i=1}^W \phi(p_i) &= \phi(1), \\
\sum_{i=1}^W \phi(p'_i) &= \phi\left(1 - \frac{\delta}{2}\right) + (W-1)\phi\left(\frac{\delta}{2} \frac{1}{W-1}\right).
\end{aligned} \tag{15}$$

Difference of the expectation values is calculated as follows:

$$\begin{aligned}
\langle Q \rangle_\phi [p] - \langle Q \rangle_\phi [p'] &= Q_1 \left\{ 1 - \frac{\phi(1-\delta/2)}{\phi(1-\delta/2) + (W-1)\phi((\delta/2)(1/(W-1)))} \right\} \\
&\quad - \left(\sum_{i=2}^W Q_i \right) \frac{\phi((\delta/2)(1/(W-1)))}{\phi(1-\delta/2) + (W-1)\phi((\delta/2)(1/(W-1)))} \\
&= \frac{W}{W-1} (Q_1 - \bar{Q}) \\
&\quad \times \frac{\phi((\delta/2)(1/(W-1)))/[(\delta/2)(1/(W-1))]}{\phi(1-\delta/2)/(\delta/2) + \phi((\delta/2)(1/(W-1)))/[(\delta/2)(1/(W-1))]} \\
&\xrightarrow{W \rightarrow \infty} Q_1 - \bar{Q},
\end{aligned} \tag{16}$$

since $\lim_{x \rightarrow +0} \phi(x)/x = \infty$. Therefore, $\langle Q \rangle_\phi [p]$ is not stable. \square

In the above proof, we have employed the specific deformations of the probability distributions as the counterexamples, which are considered in [5]. It is pointed out in [13] that these deformed distributions may experimentally be generated.

Finally, we mention a couple of simple stable examples.

Example 1.

$$\phi(x) = e^x - 1. \quad (17)$$

Example 2.

$$\phi(x) = \ln(1 + x^\alpha), \quad (18)$$

which yields a stable generalized expectation value, if and only if $\alpha = 1$.

On the other hand, as mentioned earlier, the q -expectation value is not stable, since $\phi(x) = x^q$ ($q > 0$, $q \neq 1$) does not satisfy the condition $\lim_{x \rightarrow +0} \phi(x)/x \in (0, \infty)$.

In conclusion, we have considered a class of generalized definitions of expectation value that are often employed in nonequilibrium statistical mechanics for complex systems, and have presented the necessary and sufficient condition for such a class to be stable under small deformations of a given arbitrary probability distribution.

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References

- [1] S. Abe and Y. Okamoto, Eds., *Nonextensive Statistical Mechanics and Its Applications*, vol. 560, Springer, Berlin, Germany, 2001.
- [2] C. Tsallis, *Introduction to Nonextensive Statistical Mechanics: Approaching a Complex World*, Springer, New York, NY, USA, 2009.
- [3] C. Tsallis and A. M. C. Souza, "Constructing a statistical mechanics for Beck-Cohen superstatistics," *Physical Review E*, vol. 67, no. 2, Article ID 026106, pp. 1–5, 2003.
- [4] C. Beck and F. Schlögl, *Thermodynamics of Chaotic Systems*, vol. 4, Cambridge University Press, Cambridge, UK, 1993, An Introduction.
- [5] B. Lesche, "Instabilities of Rényi entropies," *Journal of Statistical Physics*, vol. 27, no. 2, pp. 419–422, 1982.
- [6] S. Abe, "Stability of Tsallis entropy and instabilities of Rényi and normalized Tsallis entropies: a basis for q -exponential distributions," *Physical Review E*, vol. 66, no. 4, Article ID 046134, 2002.
- [7] S. Abe, G. Kaniadakis, and A. M. Scarfone, "Stabilities of generalized entropies," *Journal of Physics. A. Mathematical and General*, vol. 37, no. 44, pp. 10513–10519, 2004.
- [8] E. M. F. Curado and F. D. Nobre, "On the stability of analytic entropic forms," *Physica A*, vol. 335, no. 1-2, pp. 94–106, 2004.
- [9] A. El Kaabouchi, C. J. Ou, J. C. Chen, G. Z. Su, and Q. A. Wang, "A counterexample against the Lesche stability of a generic entropy functional," *Journal of Mathematical Physics*, vol. 52, Article ID 063302, 2011.
- [10] J. Naudts, "Continuity of a class of entropies and relative entropies," *Reviews in Mathematical Physics*, vol. 16, no. 6, pp. 809–822, 2004.
- [11] A. E. Rastegin, "Continuity and stability of partial entropic sums," *Letters in Mathematical Physics*, vol. 94, no. 3, pp. 229–242, 2010.
- [12] S. Abe, "Instability of q -averages in nonextensive statistical mechanics," *Europhysics Letters*, vol. 84, no. 6, Article ID 60006, 2008.

- [13] S. Abe, "Anomalous behavior of q -averages in nonextensive statistical mechanics," *Journal of Statistical Mechanics: Theory and Experiment*, vol. 2009, no. 7, Article ID P07027, 2009.
- [14] S. Abe, "Generalized molecular chaos hypothesis and the H theorem: problem of constraints and amendment of nonextensive statistical mechanics," *Physical Review E*, vol. 79, no. 4, Article ID 041116, 2009.
- [15] R. Hanel and S. Thurner, "Stability criteria for q -expectation values," *Physics Letters. A*, vol. 373, no. 16, pp. 1415–1420, 2009.
- [16] R. Hanel, S. Thurner, and C. Tsallis, "On the robustness of q -expectation values and Rényi entropy," *Europhysics Letters*, vol. 85, no. 2, Article ID 20005, 2009.
- [17] S. Abe, "Essential discreteness in generalized thermostatics with non-logarithmic entropy," *Europhysics Letters*, vol. 90, no. 4, Article ID 50004, 2010.
- [18] B. Andresen, "Comment on "Essential discreteness in generalized thermostatics with non-logarithmic entropy" by Abe Sumiyoshi," *Europhysics Letters*, vol. 92, no. 4, Article ID 40005, 2010.
- [19] S. Abe, "Reply to the comment by B. Andresen," *Europhysics Letters*, vol. 92, no. 4, Article ID 40006, 2010.
- [20] J. F. Lutsko, J. P. Boon, and P. Grosfils, "Is the Tsallis entropy stable?" *Europhysics Letters*, vol. 86, no. 4, Article ID 40005, 2009.
- [21] J. P. Boon and J. F. Lutsko, "Nonextensive formalism and continuous Hamiltonian systems," *Physics Letters A*, vol. 375, no. 3, pp. 329–334, 2011.
- [22] J. F. Lutsko and J. P. Boon, "Questioning the validity of non-extensive thermodynamics for classical Hamiltonian systems," *Europhysics Letters*, vol. 95, no. 2, Article ID 20006, 2011.



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