# ORIGINAL PAPER



# Comparison of AP-XPS and PLIF Measurements During CO Oxidation Over Pd Single Crystals

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**Abstract** The interaction between the gas-phase molecules and a catalyst surface is crucial for the surface structure and are therefore important to consider when the active phase of a catalyst is studied. In this study we have used two different techniques to study the gas phase during CO oxidation over Pd single crystals. Gas-phase imaging by planar laser-induced fluorescence (PLIF) shows that a spherical boundary layer with a decreasing gradient of CO<sub>2</sub> concentration out from the surface, is present close to the surface when the Pd crystal is highly active. Within this boundary layer the gas composition is completely different than that detected at the outlet of the chamber. The PLIF images of the gas-phase distribution are used to achieve a better understanding of the gas composition between the surface and the detector of a set-up for ambient pressure X-ray photoelectron spectroscopy (AP-XPS), a common technique for surface structure determination of model catalysts. The results show that also the gas-phase peaks present in the AP-XPS spectra truly represent the gas closest to the surface, which facilitates the interpretation of the AP-XPS spectra and thereby also the understanding of the mechanism behind the reaction process.

**Keywords** Ambient pressure X-ray photoelectron spectroscopy  $\cdot$  Planar laser-induced fluorescence  $\cdot$  Pd(100)  $\cdot$  Pd(110)  $\cdot$  Gas-phase

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#### 1 Introduction

CO oxidation, where a CO molecule interacts with an oxygen molecule to form CO<sub>2</sub>, is one of many reactions that occur in an automotive catalyst to clean the exhaust gases from the toxic CO molecules. Due to its simplicity, the reaction is often used as a model reaction [1] and because its wide application the reaction has been studied in great detail for many decades [2]. The reaction process is well-known under ultra-high vacuum (UHV) conditions but less is known about the surface structure and reaction mechanism under more realistic operating conditions for the catalyst. More realistic conditions in turn involve higher pressure, which implies that the number of molecules interacting with the surface increases significantly. The gas molecules interacting with the catalyst surface are essential for the surface structure and a change in the gas composition close to the surface may result in a change of the surface composition [3, 4]. It is therefore important to perform catalysis experiments in situ where knowledge of the gas molecules interacting with the surface can be achieved and used to understand the active surface structure of a catalyst in a better way [5].

There is a limited number of available techniques for probing the gas phase close to the surface of a model catalyst, and traditionally the gas composition has been analyzed using a mass spectrometer (MS) positioned at the reactor outlet. The MS data will therefore not reveal any direct information about the gas interacting with the surface and the actual connection between the MS signal and the gas distribution close to the surface is not clear. However, the development of traditional electron-based surface science techniques, enabling operation at higher pressures, has allowed for extraction of chemical information not only from the surface but also from the gas

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phase. Ambient pressure X-ray photoelectron spectroscopy (AP-XPS) [6, 7] is an example of a high-pressure set-up, common for catalysis studies, which is capable of monitoring surface reconstruction in situ as well as the gas phase close to the surface.

Oxidation of CO over transition metals is often mass transfer limited (MTL) by the minority reactants [8, 9] in the highly active regime of the catalyst. This means that the reaction is not kinetically controlled but, limited by the number of minority reactants that can penetrate a boundary layer of the product to reach the surface. The reaction is therefore completely controlled by diffusion of the reactants, which makes the gas-phase composition highly important for the catalytic process. Gas-phase studies using AP-XPS have been performed and reported in the literature [10–12], but very few have been performed for a highly active catalyst during CO oxidation [13-16]. As mentioned previously, the formation of a CO<sub>2</sub> boundary layer changes the gas composition close to the sample surface, but how this is reflected in the AP-XPS gas-phase signal is not obvious. Planar laser-induced fluorescence (PLIF) has been shown to be capable of imaging the gas-phase distribution above an active catalyst surface with high spatial resolution [17–20]. The imaged gas-phase region is similar as the one probed by AP-XPS and can therefore contribute to an improved understanding of the AP-XPS spectra and, as a consequence, a better understanding of the details of the catalytic processes on the surface.

This contribution presents a comparison of the gas phase composition detected by AP-XPS and PLIF close to a highly active Pd catalyst surface exposed to a 1:1 mixture of CO and  $O_2$ . By comparing the ratios of CO and  $CO_2$  for the two techniques, we show that AP-XPS probes the gas composition in the vicinity of the surface, and therefore provides a considerably better estimation of the gas composition as compared to an MS located at the outlet of the reactor.

# 2 Experimental

All experiments were performed in mixtures having a 1:1 ratio of CO and O<sub>2</sub>, but at different total pressures. In the AP-XPS experiment the total pressure was varied from 0.013 to 1.3 mbar [13]. The PLIF experiments performed for CO and CO<sub>2</sub> detection were carried out at a total pressure of 106 and 150 mbar, respectively, where a pressure controller at the outlet was used to keep the pressure constant in the chamber. The partial pressures in the CO experiment were set to 26.5 mbar (18 ml<sub>n</sub>/min) each for CO and O<sub>2</sub>, and to 53 mbar (36 ml<sub>n</sub>/min) for Ar whereas 6 mbar (4 ml<sub>n</sub>/min) each of CO and O<sub>2</sub> and 138 mbar (92 ml<sub>n</sub>/min) of Ar were used in the CO<sub>2</sub> experiment. This results in a gas residence time in the

chamber of 20 s for the CO experiment while in the  $\mathrm{CO}_2$  experiment the gas residence time is only 2 s. For the XPS experiment the inlet flow was controlled via a leak valve and the partial pressure of each gas was measured. No pressure controller was used and the inlet and outlet flow was constant throughout the whole experiment. In all experiments a standard BN heater [21, 22] was used to ramp the temperature of the crystal in order to start with an inactive crystal and then monitor the gas phase when the catalytic activity of the sample increases and reaches the highly active regime, in which the reaction is mass transfer limited.

# 2.1 Ambient Pressure X-Ray Photoelectron Spectroscopy

The AP-XPS measurements were carried out at beamline 9.3.2 at the ALS in Berkeley [10], CA, USA where differential pumping stages combined with electrostatic lenses enables reaction experiments in situ at pressures up to approximately 10 mbar. The sample was mounted on a special-made holder for catalysis experiments and a thermocouple was attached to the heating plate close to the single crystal. The set-up consists of two chambers between which the sample can be transferred without being exposed to air. This allows for sample preparation, such as sputtering, before transfer into the analysis chamber, where a hemispherical analyzer is attached. The gas pressures were controlled by individual leak valves. The only outlet from the analysis chamber is via the nozzle to the detector and in order to maintain the pressure at the surface to at least 98 % of the inlet pressure at room temperature, the working distance between the sample and nozzle should be larger than twice the diameter of the aperture at the first pumping stage [10]. To achieve this for the present beamline set-up, the working distance was approximately 2 mm, which is also the focus of the detector and implies that the number of detected electrons from the surface is optimized.

The C 1s spectra were measured with a photon energy of 435 eV and the energy was calibrated to the Fermi level. The decomposition of the spectra was made using a Doniach and Sunjic [23] lineshape convoluted with a Gaussian lineshape and subtraction of a linear background.

# 2.2 Planar Laser-Induced Fluorescence

PLIF allows for species-specific detection in the gas phase with high sensitivity and two-dimensional imaging. The technique is commonly employed for combustion diagnostics, but less used in the field of catalysis [22, 24]. Here, we have used PLIF to visualize the gas-phase CO and CO<sub>2</sub> above two different Pd single crystals. Two different laser



systems and reactors were used for probing CO and CO<sub>2</sub> respectively, but the layout of the experimental set-up was the same in both experiments and is shown in Fig. 1. In both set-ups an MS was attached to the gas outlet of the chamber for analysis of the overall gas composition. The Pd(100) crystal used for the CO<sub>2</sub> experiment was the same as used in the AP-XPS experiments but when CO was probed a Pd(110) crystal, with the same size, was used. In these studies, thin lasersheets arranged perpendicular to the crystal surface probed CO and CO2 across the center of the crystals with wavelengths of 230 nm and 2.7 µm, respectively. The resulting fluorescence emission at 440-660 nm for CO and 4.26 µm for CO2 was detected with a CCD camera and focal plane array detector, respectively, and images of an area of approximately  $5 \times 16 \text{ mm}^2$  of the specific gas distribution were obtained. The PLIF detection limit is at ppm levels and a spatial resolution better than 0.4 and 0.07 mm for the CO and CO2 measurements, respectively, were achieved in the presented measurements. The laser repetition rate was 10 Hz, also representing the number of images taken every second. However, for noise reduction, the images shown in Fig. 4, are averaged over ten acquired images generating a timescale of 1 image/s. The analysis process for the PLIF images, which for instance includes temperature corrections and calibrations, is described in detail in [22].

#### 2.3 Sample Preparation

Pd single crystals with surface orientation (100) and (110) were used as model catalysts. In the AP-XPS measurements the Pd(100) crystal was cleaned by sputtering and oxygen treatment in the preparation chamber. The crystal was then transferred into the analysis chamber without being exposed to air. The cleanness of the surface was then checked with AP-XPS before the reaction experiments

were performed. In the PLIF experiments, the crystals were cleaned in the same way, with sputtering and oxygen treatment, but in this set-up there is no transfer system between the preparation chamber and the reactor, and the crystal was exposed to air before it was mounted in the reactor. To reduce contaminations on the surface the crystal temperature was ramped up and down in a CO and  $O_2$  environment before the reaction experiments were performed.

# 3 Results

In Fig. 2, the C 1s spectra are shown, acquired during CO oxidation using a 1:1 ratio of CO and O2 at different total pressures [13] with a Pd(100) single crystal used as a model catalyst. The O 1s and Pd 3d<sub>5/2</sub> core levels were also measured but are not shown here. Starting with the experiment at a total pressure of 1.3 mbar (Fig. 2a), the sample is inactive at low temperature and CO in the gas phase is clearly observed at approximately 290 eV. CO adsorbed in a bridge site on the Pd(100) surface is also detected at 286 eV [25]. The temperature of the crystal is then stepwise increased and when a temperature of 340 °C is reached the sample is highly active and CO desorbs instantaneously from the surface simultaneously as the CO gas-phase peak disappears. In the spectrum measured at 340 °C a single peak, corresponding to CO<sub>2</sub> in the gas phase, is observed, showing that the sample is highly active. When the temperature is increased further, no obvious increase in the CO<sub>2</sub> gas-phase peak can be seen which shows that under these conditions, the reaction reaches its maximum conversion of CO to CO2 immediately after ignition. The conclusion is then that after ignition the reaction is not kinetically controlled but mass transfer limited by the CO diffusion. However, a highly

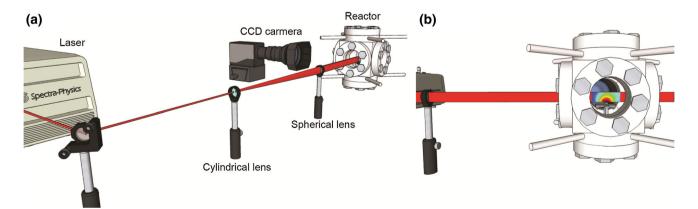


Fig. 1 An overview of the PLIF experimental set-up. a The laser-beam is shaped into a thin lasersheet by passing a cylindrical and spherical lens before it enters the reactor. The lasersheet can then

probe the gas in two dimensions. b The fluorescence is detected by a CCD camera perpendicular to the lasersheet and a PLIF image of the gas distribution over the Pd crystal can be detected



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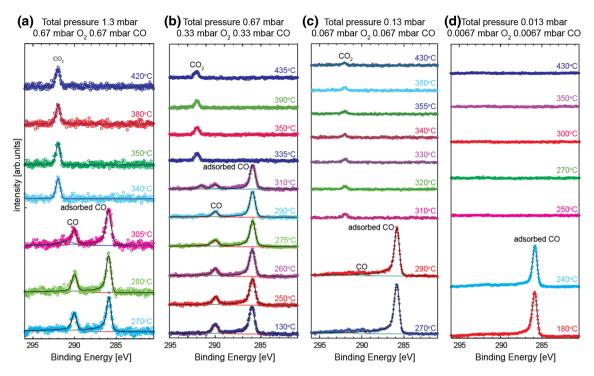


Fig. 2 AP-XPS C 1s spectra from CO oxidation using Pd(100) catalyst in 1:1 CO and O<sub>2</sub> atmospheres at different total pressures **a** 1.3 mbar, **b** 0.67 mbar, **c** 0.13 mbar, **d** 0.013 mbar. CO<sub>2</sub> and CO gas-phase peaks are observed as well as CO adsorbed on the surface

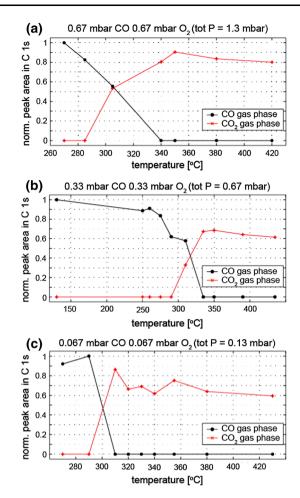
active sample producing a significant amount of CO<sub>2</sub> should still have CO in the gas phase close to the surface to maintain its high activity. Surprisingly, no trace of CO can be is seen in the spectra after ignition. When the total pressure is reduced to 0.67 mbar, shown in Fig. 2b, a similar behavior can be observed. CO is adsorbed on the surface until the ignition temperature is reached and after ignition, which occurs at slightly lower temperature compared with the higher pressure (1.3 mbar), only gas-phase CO<sub>2</sub> can be seen in the C 1s spectra. Interestingly, at 310 °C both a CO<sub>2</sub> and a CO gas-phase peak can be detected together with a peak originating from CO adsorbed on the surface, indicating that a complete MTL has not yet been reached, which in turn implies that only parts of the surface is fully active. We know from our previous studies [13] that chemisorbed oxygen starts to occupy the surface at this moment, but the exact nature of the surface is obviously still beyond our reach. To speculate, we could suggest that the surface consist of islands of atomic oxygen bordered by islands of CO, and the CO<sub>2</sub> production occurs at the border between the islands and/or at vacancies within the oxygen islands. Furthermore, it can also be seen that the CO<sub>2</sub> peak is shifting towards higher energy as the surface goes from partly to fully active. The reason is the change in the surface work function [26, 27] as the CO desorbs fully from the surface and the oxygen fully adsorb on the surface. Decreasing the total pressure even further to 0.13 mbar (Fig. 2c) the ignition temperature is decreased

to 310 °C, but the CO-trend is still the same. It is also clear that the ratio between peaks of adsorbed CO and gas-phase CO is increasing with decreasing pressure, even though a similar CO coverage of the surface is expected for the different total pressures. This is simply because the number of molecules in the gas phase decreases when the pressure is reduced. The spread of the raw data (circles in Fig. 2) is also lower when the pressure is reduced, which results in more well-defined spectra. However, at an even lower total pressure of 0.013 mbar (Fig. 2d), the partial pressures of CO and CO<sub>2</sub> are too low to be detected and no gas-phase peaks in the C 1s spectra can be observed. Still, the ignition temperature is concluded to be around 250 °C, due to the absence of the peak corresponding to adsorbed CO at this temperature. In fact, the absence of a peak corresponding to CO adsorbed on the surface demonstrates a very low coverage suggesting that the reaction is extremely fast, and that all CO reaching the surface instantaneously is converted to CO<sub>2</sub>.

To investigate the changes in the gas phase in more detail the areas of the CO and CO<sub>2</sub> gas-phase peaks are extracted from the C 1s spectra and plotted in Fig. 3. The advantage of extracting the CO and CO<sub>2</sub> peak areas from the same spectrum is that difficulties such as compensation for background and number-of-sweeps, are reduced. In addition, the cross sections for CO and CO<sub>2</sub> in the gas phase are similar [28], making it possible to compare the integrated intensities of the peaks. In the plots, the peak



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**Fig. 3** Normalized CO and CO<sub>2</sub> gas-phase peak areas, evaluated from the AP-XPS C 1s spectra, plotted versus catalyst temperature for different total pressures **a** 1.3 mbar, **b** 0.67 mbar, **c** 0.13 mbar. *Lines* have been inserted between the measurement points for guidance

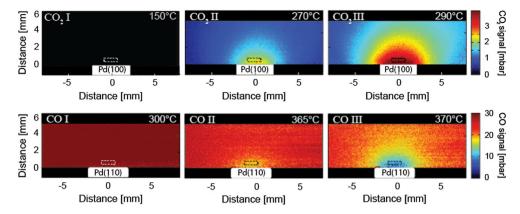
areas are normalized to the area of the CO peak in the beginning of the experiment but similar results are achieved by plotting the ratios between the gas-phase peaks and the Pd substrate peak. The error of the data can be estimated to be  $\pm 5$  %, based on the statistical error between consecutive scans. Figure 3a shows results for a total pressure of 1.3 mbar, where the area of the CO gasphase peak (black points) decreases when the temperature is increased while the CO<sub>2</sub> peak is not detectable until after ignition. At ignition the CO peak decreases to zero and the CO<sub>2</sub> peak area increases (red crosses) to an approximately constant level confirming that the MTL regime is reached. Figure 3b and c show the peak areas at the lower total pressures of 0.67 and 0.13 mbar. No obvious difference can be seen between the three plots, suggesting that the gas composition close to the surface is similar for all three total pressures. The absence of the CO gas-phase peak after ignition could be interpreted as no CO molecules are present in the gas phase close to the surface after ignition but instead immediately converted into CO<sub>2</sub>. However, for none of the three experiments in which the gas phase is detected, the CO<sub>2</sub> peak area reaches the same peak area as the CO prior to ignition (it only reaches 60–80 % of the CO peak). This would then suggest that CO is still present in the gas phase but below the detection limit of the AP-XPS set-up, however, the interpretation of the data is still not obvious. It is of particular interest to know how the gasphase peaks observed in the AP-XPS spectra are related to the actual gas composition close to the surface. In an attempt to answer this question we have investigated the composition of the gas-phase more thoroughly using PLIF.

PLIF images of CO and CO<sub>2</sub> acquired at three different temperatures, respectively, during the temperature ramp, are shown in Fig. 4. The images represent the gas composition of an inactive sample (image I), at ignition (image II), and a highly active sample (image III). Due to the different ignition temperatures of the crystals [365 °C for Pd(110) and 270 °C for the Pd(100)] the temperature of the samples when the images were acquired are different for the CO and CO<sub>2</sub> cases. Starting with images acquired at the lowest temperatures, the lack of CO<sub>2</sub> PLIF signal and the homogenous partial pressure of CO clearly indicates that the samples are inactive at these temperatures (image I). When the ignition temperatures of the samples are reached, local changes in the gas composition close to the surface are observed (image II). Immediately after ignition, a spherically shaped region with increased CO<sub>2</sub> concentration is formed around the surface (image III) which stays unchanged even if the temperature is increased further. This is a visualization of the boundary layer of the CO<sub>2</sub> that hinders CO from reaching the surface in the MTL regime, consistent with the reduced CO signal detected in the corresponding region. The images show that after the boundary layer is formed, the changes in the gas composition close to the surface are small, indicating steady state conditions in the gas phase. It is also clear that the gasphase evolution during the reaction is comparable for both crystals and that the surface structure does not affect the gas phase significantly, which is in agreement with the results reported by Toyoshima et al. [14, 15]. They also report a lower ignition temperature for Pd(110) compared to the Pd(100), which is not observed in the PLIF studies, possibly because different chambers with different gas flow properties were used for the two experiments. The partial pressures of CO and O<sub>2</sub> are also higher (26.5 mbar) in the experiment where the Pd(110) was used compared to the partial pressures used when the Pd(100) crystal was investigated (6 mbar). This may also contribute to a higher ignition temperature for the Pd(110) compared with the Pd(100) crystal.

Figure 5a shows MS data recorded during CO oxidation over Pd(110). When the sample temperature reaches



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**Fig. 4** PLIF images of CO<sub>2</sub> (*upper panel*) and CO (*lower panel*) gasphase distributions over Pd single crystals during CO oxidation in a 1:1 ratio of CO and O<sub>2</sub>. The gas flow is in the direction from *left to right* in all images. The *white filled rectangles* represent the single

crystals samples and the *dashed rectangles* above show the areas where the PLIF signals plotted in Fig. 5 are evaluated. The *Roman numerals* are correlated to time points and temperatures in Fig. 5

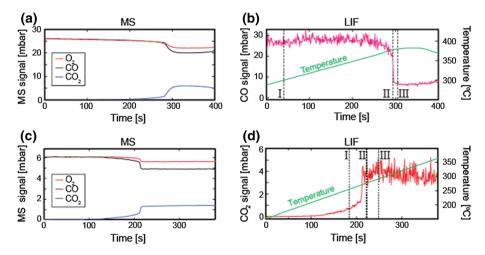


Fig. 5 Comparison of the MS and PLIF signals, a MS signal detected during CO oxidation over Pd(110). b Corresponding CO PLIF trend extracted 0.5 mm from surface and sample temperature. c MS signal

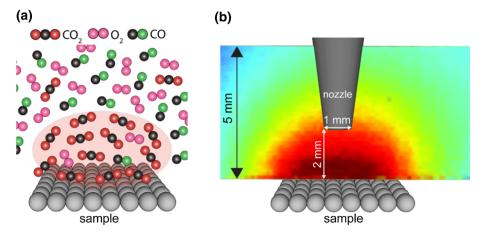
detected during CO oxidation over Pd(100). d Corresponding CO<sub>2</sub> PLIF trend extracted 0.5 mm from surface and sample temperature

ignition the CO signal decreases simultaneously as the CO<sub>2</sub> signal increases and both signals reach constant levels, even though the temperature is increased further. This is, as discussed above, a clear indication of an MTL reaction. In this regime the MS CO partial pressure detected at the outlet of the chamber has decreased by approximately 25 %. This can be compared to the CO PLIF signal plotted in Fig. 5b, extracted 0.5 mm from the surface (the dashed rectangle in the images in Fig. 4 shows the area where the signal is extracted) where the CO signal decreases by approximately 80 %. The Roman numerals in the plot indicate time points for the corresponding image of the gasphase distribution shown in Fig. 4. Similar divergence is also observed between the MS and CO<sub>2</sub> PLIF signal shown in Fig. 5c and d, respectively. The results from the MS and PLIF can in turn be compared to the C 1s peak areas plotted in Fig. 3. The C 1s peak areas show, just as the PLIF data, a higher ratio between the reactant and the product signal as compared to the MS signal. The reason for the different ratios detected by the different techniques can be understood by studying the PLIF images. The images reveal a completely different gas composition in the MTL regime, close to the surface where the AP-XPS and PLIF signals are extracted, which is not detected by the MS.

# 4 Discussion

The gas molecules interacting with the surface are obviously important for the surface composition, and the determination of an active site or phase of the catalyst. It is





**Fig. 6** Illustration of the boundary layer of CO<sub>2</sub> during CO oxidation over a highly active surface. **a** Model showing the high CO<sub>2</sub> and low CO concentration within the boundary layer. **b** Schematic view of the highly active sample surface with PLIF image showing the CO<sub>2</sub>

distribution above the catalyst. The AP-XPS nozzle is located at normal operating distance, which is approximately twice the diameter of the nozzle aperture. Electrons with a take-off angle of 23° will reach the detector

therefore important to gain knowledge about the gas distribution around the catalyst in order to understand to which extent relevant information can be deduced from gas-phase AP-XPS spectra. Due to the differences in the experimental set-ups, the extent and size of the gas phase distribution at MTL may be influenced by the different flow and conditions of the different set-ups, but as our results demonstrate, the main contribution to the gas-phase peaks at MTL in AP-XPS originates from the seriously CO-depleted zone, often referred to in the literature as the boundary layer. The very high conversion of CO that is observed in both the AP-XPS and in the PLIF signal is within this boundary layer. We do not observe any indication that the catalytic behavior at MTL is different because of the different experimental set-ups.

Both PLIF and AP-XPS show consistently a significantly different gas composition above the single crystal surfaces compared to the rest of the chamber when the MTL regime is reached, as illustrated in Fig. 6a. This behavior is observed for both the Pd(110) and Pd(100) crystals and is expected to follow for any highly active transition metal during CO oxidation, at least in the case of a single crystal positioned in a reactor in a similar way as in the present experiments. The concentration gradient of CO and CO<sub>2</sub> within the boundary layer is dependent on the gas exchange in the chamber and the activity of the sample. The gas flow in the AP-XPS set-up is complex due to the large volume of the reactor with the only outlet being the nozzle of the analyzer. The distance between the aperture of the nozzle and the surface is about 2 mm resulting in a semi-flow type of reactor, and the rate of the gas exchange and the gas composition above the surface are difficult to estimate.

Comparing ratios of the CO and CO<sub>2</sub> peak areas in the AP-XPS C 1s spectra, measured during the temperature ramp, as well as PLIF and MS signals, shows that AP-XPS

data are in considerably better agreement with PLIF data than with MS data. The peak areas, however, differ compared with ratios of the gas-phase peak heights shown in the AP-XPS spectra (Fig. 2). The reason for the larger area of the CO peak is the larger splitting of the vibrational fine structure in CO as compared with CO<sub>2</sub> [29, 30], which gives a broadening of the peak and generates a larger area for the CO peak. The stronger scattering of the electrons of the CO<sub>2</sub> molecules might also modify the peak area, but this contribution is relatively small and will not affect the high ratio between CO and CO<sub>2</sub> signals in the MTL regime. This observation indicates that most of the gas-phase molecules detected in the AP-XPS measurements originate from a region with a comparable relation between the CO and CO<sub>2</sub> concentrations as inside the boundary layer imaged in the PLIF experiments, and that the aperture of the AP-XPS detector nozzle is located at a distance from the surface where the gas-phase concentration gradient has not changed significantly but reveal a similar CO and CO<sub>2</sub> ratio as that close to the surface, schematically illustrated in Fig. 6b. Furthermore, the AP-XPS measurements are averaged over a volume above the surface. As a consequence, if the gas-phase distribution or the boundary layer were to be measured by AP-XPS, it is likely to appear smeared out as compared to the distribution observed by PLIF. This drastic reduction would result in a CO level just below the detection limit for AP-XPS, explaining the absence of the CO gas-phase peak for the highly active sample. It could be argued that the CO should still be detected in the AP-XPS measurements for the highest pressure, 0.13 mbar (20 % of 0.67 mbar). However, 0.13 mbar is close to the general detection limit of AP-XPS  $(\sim 0.05 \text{ mbar})$  [31], which is further affected by the photoelectron kinetic energy, the photoelectron cross-section, and the types of gases used in the experiments. The PLIF



measurements, however, confirm that CO is still present in the gas-phase above the crystal but the concentration decreases by approximately 80 % compared to the initial CO level. This is interpreted as the fast build-up of the boundary layer of CO<sub>2</sub> after ignition that hinders CO to reach the surface which decreases the CO concentration close to the surface significantly in the MTL regime.

# 5 Conclusions

The PLIF images of CO and CO<sub>2</sub> display a boundary layer with a spherical shape around the surface when the MTL is reached during CO oxidation. This boundary layer has a significantly different gas composition than that measured at the chamber gas inlet or outlet, with a CO concentration below the detection limit of the AP-XPS. By studying the gas-phase peak areas extracted from the C 1s AP-XPS spectra and compare with the PLIF and MS results, we confirm that the ratio of CO and CO2 concentrations detected by AP-XPS is comparable with the gas ratio inside the boundary layer. The results indicate that the gas-phase peaks in the AP-XPS spectra represent the gas phase interacting with the surface, and that the aperture of the nozzle is located at a distance from the surface where the CO and CO<sub>2</sub> concentrations are still similar as within the boundary layer. AP-XPS is therefore one of few techniques able to probe the gas phase just above the catalyst surface. This makes AP-XPS suitable for in situ catalysis studies where the spectra can provide information about the gas phase molecules relevant for the surface structure determination of a highly active catalyst.

Altogether, the presence of a boundary layer in the highly active phase of the catalyst implies that in situ studies are important when CO oxidation is investigated in order to achieve a representative picture of the surface structure and the gas molecules interacting with the surface.

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