Li and Li Journal of Inequalities and Applications (2015) 2015:356 DOI 10.1186/s13660-015-0881-3

 Journal of Inequalities and Applications a SpringerOpen Journal

# RESEARCH





# Density expansions of extremes from general error distribution with applications

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# Abstract

In this paper the higher-order expansions of density of normalized maximum with parent following general error distribution are established. The main results are applied to derive the higher-order expansions of the moments of extremes.

MSC: Primary 60G70; secondary 60F05

**Keywords:** density expansion; general error distribution; maximum; moment expansion

# 1 Introduction

In extreme value theory, the quality of convergence of normalized partial maximum of a sample has been studied in recent literature. For the convergence rate of distribution of normalized maximum, we refer to Smith [1], Leadbetter *et al.* [2], de Haan and Resnick [3] for general cases, and specific cases were studied by Hall [4], Nair [5], Peng *et al.* [6] and Jia and Li [7]. Nair [5] derived the higher-order expansions of moments of normalized maximum with parent following normal distribution. Liao *et al.* [8] and Jia *et al.* [9] extended Nair's results to skew-normal distribution and general error distribution, respectively.

The main objective of this paper is to derive the higher-order expansions of density of normalized maximum with parent following the general error distribution. To the best of our knowledge, there are few studies on the rate of convergence of density of normalized maximum except the work of de Haan and Resnick [10] for local limit theorems and Omey [11] for rates of convergence of densities with regular variation with remainders excluding the case we will study in this paper, *i.e.*, the general error distribution.

Let  $\{X_n, n \ge 1\}$  be a sequence of independent and identically distributed (i.i.d.) random variables with marginal cumulative distribution function (cdf)  $F_{\nu}$  following the general error distribution ( $F_{\nu} \sim \text{GED}(\nu)$  for short), and let  $M_n = \max_{1 \le k \le n} X_k$  denote its partial maximum. The probability density function (pdf) of the GED( $\nu$ ) is given by

$$f_{\nu}(x) = \frac{\nu \exp(-(1/2)|x/\lambda|^{\nu})}{\lambda 2^{1+1/\nu} \Gamma(1/\nu)}, \quad x \in \mathbb{R},$$

where  $\nu > 0$  is the shape parameter,  $\lambda = [2^{-2/\nu}\Gamma(1/\nu)/\Gamma(3/\nu)]^{1/2}$  and  $\Gamma(\cdot)$  denotes the gamma function (Nelson [12]). Note that the GED(2) reduces to the standard normal distribution.

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For the GED( $\nu$ ), the limiting distribution of maximum  $M_n$  and its associated higherorder expansions are given by Peng *et al.* [13] and Jia and Li [7]. Peng *et al.* [6] showed that

$$\lim_{n \to \infty} \mathbb{P}(M_n \le a_n x + b_n) = \Lambda(x) = \exp(-\exp(-x)), \quad x \in \mathbb{R}$$
(1.1)

provided the norming constants  $a_n$  and  $b_n$  satisfy the following equations:

$$1 - F_{\nu}(b_n) = n^{-1}, \qquad a_n = f(b_n), \tag{1.2}$$

where  $f(x) = 2\nu^{-1}\lambda^{\nu}x^{1-\nu}$ . In the sequel, let

$$g_n(x) = na_n F_v^{n-1}(a_n x + b_n) f_v(a_n x + b_n)$$
(1.3)

denote the density of normalized maximum, and

$$\Delta_n(g_n, \Lambda'; x) = g_n(x) - \Lambda'(x) \tag{1.4}$$

with  $\Lambda'(x) = e^{-x}\Lambda(x)$ . By Proposition 2.5 in Resnick [14],  $\Delta_n(g_n, \Lambda') \to 0$  as  $n \to \infty$ . For both applications and theoretical analysis, it is of interest to know the convergence rate of (1.4). This paper focuses on this topic and applies the main results to derive the high-order expansions of moments of extremes.

The paper is organized as follows. Section 2 provides the main results and all proofs are deferred to Section 4. Auxiliary lemmas with proofs are given in Section 3.

# 2 Main results

In this section, we present the asymptotic expansions of density for the normalized maximum formed by the  $GED(\nu)$  random variables and its applications to the higher-order expansions of moments of extremes.

**Theorem 2.1** Let  $F_{\nu}(x)$  denote the cdf of GED( $\nu$ ) with  $\nu > 0$ , then for  $\nu \neq 1$ , with norming constants  $a_n$  and  $b_n$  given by (1.2), we have

$$b_n^{\nu} \left[ b_n^{\nu} \Delta_n(g_n, \Lambda'; x) - k_{\nu}(x) \Lambda'(x) \right] \to \omega_{\nu}(x) \Lambda'(x)$$
(2.1)

as  $n \to \infty$ , where  $k_{\nu}(x)$  and  $\omega_{\nu}(x)$  are respectively given by

$$k_{\nu}(x) = k_{\nu 1}(x) + k_{\nu 2}(x), \qquad (2.2)$$

$$\omega_{\nu}(x) = (1 - \nu^{-1})\lambda^{2\nu} (\omega_{\nu 1}(x) + \omega_{\nu 2}(x)e^{-x} + \omega_{\nu 3}(x)e^{-2x})$$
(2.3)

with

$$\begin{split} k_{\nu 1}(x) &= -(1-\nu^{-1})\lambda^{\nu}(x^{2}+2x),\\ k_{\nu 2}(x) &= (1-\nu^{-1})\lambda^{\nu}(2+2x+(x^{2}+2x)e^{-x}),\\ \omega_{\nu 1}(x) &= -4-2(1-\nu^{-1})x^{2}-\frac{2}{3}(1-2\nu^{-1})x^{3}+\frac{1}{2}(1-\nu^{-1})x^{4}, \end{split}$$

$$\begin{split} \omega_{\nu 2}(x) &= -\frac{4}{\nu}x - \frac{2}{\nu}x^2 - \frac{2}{3}(5 - 4\nu^{-1})x^3 - \frac{3}{2}(1 - \nu^{-1})x^4,\\ \omega_{\nu 3}(x) &= \frac{1}{2}(1 - \nu^{-1})(x^2 + 2x)^2. \end{split}$$

**Remark 2.1** If we choose the norming constants  $a_n$  and  $b_n$  such that

$$\frac{1 - F_{\nu}(b_n)}{f_{\nu}(b_n)} \sim \frac{2\lambda^{\nu}}{\nu} b_n^{1-\nu}, \qquad a_n = f(b_n)$$
(2.4)

with  $\nu \neq 1$ , then

$$b_n^{\nu} \left[ b_n^{\nu} \Delta_n \left( g_n, \Lambda'; x \right) - \bar{k}_{\nu}(x) \Lambda'(x) \right] \to \bar{\omega}_{\nu}(x) \Lambda'(x)$$
(2.5)

as  $n \to \infty$ , where  $\bar{k}_{\nu}(x)$  and  $\bar{\omega}_{\nu}(x)$  are respectively given by

$$\bar{k}_{\nu}(x) = (1 - \nu^{-1})\lambda^{\nu} \left(-x^{2} + (2 + 2x + x^{2})e^{-x}\right)$$
(2.6)

and

$$\bar{\omega}_{\nu}(x) = \lambda^{2\nu} (1 - \nu^{-1}) (\bar{\omega}_{\nu 1}(x) + \bar{\omega}_{\nu 2}(x) e^{-x} + \bar{\omega}_{\nu 3}(x) e^{-2x})$$
(2.7)

with

$$\begin{split} \bar{\omega}_{\nu 1}(x) &= \frac{2(2\nu^{-1}-1)}{3}x^3 + \frac{1-\nu^{-1}}{2}x^4, \\ \bar{\omega}_{\nu 2}(x) &= 4\left(\nu^{-1}-2\right) + 4\left(\nu^{-1}-2\right)x + 2\left(2\nu^{-1}-3\right)x^2 + \frac{2}{3}\left(4\nu^{-1}-5\right)x^3 + \frac{3}{2}\left(\nu^{-1}-1\right)x^4, \\ \bar{\omega}_{\nu 3}(x) &= 2\left(1-\nu^{-1}\right)\left(1+x+\frac{x^2}{2}\right)^2. \end{split}$$

For the case of v = 1, we have the following results.

**Theorem 2.2** For v = 1, with norming constants  $a_n = 2^{-1/2}$  and  $b_n = 2^{-1/2} \log(n/2)$ , we have

$$n\left[n\Delta_n(g_n,\Lambda';x) - k_1(x)\Lambda'(x)\right] \to \left(\omega_1(x) + \frac{k_1^2(x)}{2}\right)\Lambda'(x)$$
(2.8)

as  $n \to \infty$ , where  $k_1(x)$  and  $\omega_1(x)$  are respectively given by

$$k_1(x) = -\frac{1}{2}e^{-2x}, \qquad \omega_1(x) = -\frac{1}{3}e^{-3x}.$$
 (2.9)

To end this section, we apply the higher-order expansions of densities to derive the asymptotic expansions of the moments of extremes. Methods used here are different from those in Nair [5] and Jia *et al.* [9].

In the sequel, for nonnegative integers *r*, let

$$m_r(n) = \int_{x \in \mathbb{R}} x^r g_n(x) \, dx, \qquad m_r = \int_{x \in \mathbb{R}} x^r \Lambda'(x) \, dx$$

denote respectively the *r*th moments of  $(M_n - b_n)/a_n$  and its limits.

**Theorem 2.3** Let  $\{X_n, n \ge 1\}$  be an iid sequence with marginal distribution  $F_v \sim \text{GED}(v)$ , then

(i) for  $v \neq 1$ , with norming constants  $a_n$  and  $b_n$  given by (1.2), we have

$$b_n^{\nu} \left[ b_n^{\nu} (m_r(n) - m_r) + (1 - \nu^{-1}) \lambda^{\nu} r(m_{r+1} + 2m_r) \right]$$
  

$$\rightarrow 2r \lambda^{2\nu} (1 - \nu^{-1}) \left[ \left( (1 - \nu^{-1})(r+1) + 2 \right) m_r + \left( (1 - \nu^{-1})(r+1) + 1 \right) m_{r+1} + \left( \frac{1}{4} (1 - \nu^{-1})(r-1) + \frac{1}{3} (2 - \nu^{-1}) \right) m_{r+2} \right]$$
(2.10)

as  $n \to \infty$ ;

(ii) for v = 1, with norming constants  $a_n = 2^{-1/2}$  and  $b_n = 2^{-1/2} \log(n/2)$ , we have

$$n\left[n\left(m_r(n) - m_r\right) + (-1)^r \frac{r}{2} \Gamma^{(r-1)}(2)\right] \to (-1)^{r-1} \frac{r}{24} \left[8\Gamma^{(r-1)}(3) - 3\Gamma^{(r-1)}(4)\right]$$
(2.11)

as  $n \to \infty$ , where  $\Gamma^{(r-1)}(t)$  denote the (r-1)th derivative of the gamma function at x = t.

# **3** Auxiliary lemmas

In this section we provide auxiliary lemmas which are needed to prove the main results.

**Lemma 3.1** Let  $F_v(x)$  and  $f_v(x)$  respectively denote the cdf and pdf of GED(v) with  $v \neq 1$ , for large x, we have

$$I - F_{\nu}(x) = f_{\nu}(x) \frac{2\lambda^{\nu}}{\nu} x^{1-\nu} \Big[ 1 + 2\left(\nu^{-1} - 1\right)\lambda^{\nu} x^{-\nu} + 4\left(\nu^{-1} - 1\right)\left(\nu^{-1} - 2\right)\lambda^{2\nu} x^{-2\nu} + O\left(x^{-3\nu}\right) \Big].$$
(3.1)

Furthermore, with the norming constants  $a_n$  and  $b_n$  given by (1.2), we have

(i) for  $v \neq 1$ ,

$$b_n^{\nu} \left[ b_n^{\nu} \left( F_{\nu}^n(a_n x + b_n) - \Lambda(x) \right) - \tilde{k}_{\nu}(x) \Lambda(x) \right] \to \left( \tilde{\omega}_{\nu}(x) + \frac{\tilde{k}_{\nu}^2(x)}{2} \right) \Lambda(x)$$
(3.2)

as  $n \to \infty$ , where  $\tilde{k}_{\nu}(x)$  and  $\tilde{\omega}_{\nu}(x)$  are respectively given by

$$\tilde{k}_{\nu}(x) = (1 - \nu^{-1})\lambda^{\nu}(x^{2} + 2x)e^{-x},$$
(3.3)

$$\tilde{\omega}_{\nu}(x) = \left(\nu^{-1} - 1\right)\lambda^{2\nu} \left(4x + 2x^2 + \frac{2}{3}\left(2 - \nu^{-1}\right)x^3 + \frac{1}{2}\left(1 - \nu^{-1}\right)x^4\right)e^{-x};$$
(3.4)

(ii) for v = 1, with norming constants  $a_n = 2^{-1/2}$  and  $b_n = 2^{-1/2} \log(n/2)$ , we have

$$n\left[n\left(F_1^n(a_nx+b_n)-\Lambda(x)\right)-k_1(x)\Lambda(x)\right] \to \left(\omega_1(x)+\frac{k_1^2(x)}{2}\right)\Lambda(x)$$
(3.5)

as  $n \to \infty$ , where  $k_1(x)$  and  $\omega_1(x)$  are those given by (2.9).

*Proof* See Lemma 1 and Theorem 1 in Jia and Li [7].

**Lemma 3.2** Let  $F_{\nu}(x)$  denote the cdf of GED( $\nu$ ) with  $\nu > 0$ , then (i) for  $\nu \neq 1$ , with norming constants given by (1.2), we have

$$F_{\nu}^{n-1}(a_n x + b_n) = \left(1 + \tilde{k}_{\nu}(x)b_n^{-\nu} + \left(\tilde{\omega}_{\nu}(x) + \frac{\tilde{k}_{\nu}^2(x)}{2}\right)b_n^{-2\nu}(1 + o(1))\right)\Lambda(x)$$
(3.6)

as  $n \to \infty$ , where  $\tilde{k}_{\nu}(x)$  and  $\tilde{\omega}_{\nu}(x)$  are respectively given by (3.3) and (3.4);

(ii) for v = 1, with norming constants  $a_n = 2^{-1/2}$  and  $b_n = 2^{-1/2} \log(n/2)$ , we have

$$F_1^{n-1}(a_n x + b_n) = \left(1 + \frac{1}{n}k_1(x) + \frac{1}{n^2}\left(\omega_1(x) + \frac{k_1^2(x)}{2}\right)\left(1 + o(1)\right)\right)\Lambda(x)$$
(3.7)

as  $n \to \infty$ , where  $k_1(x)$  and  $\omega_1(x)$  are given by (2.9).

Proof (i) It follows from (3.2) and (3.3) that

$$F_{\nu}^{n}(a_{n}x+b_{n}) = b_{n}^{-\nu} \left( b_{n}^{-\nu} \left( \tilde{\omega}_{\nu}(x) + \frac{\tilde{k}_{\nu}^{2}(x)}{2} \right) (1+o(1)) + \tilde{k}_{\nu}(x) \right) \Lambda(x) + \Lambda(x)$$
$$= \left( 1 + \tilde{k}_{\nu}(x) b_{n}^{-\nu} + \left( \tilde{\omega}_{\nu}(x) + \frac{\tilde{k}_{\nu}^{2}(x)}{2} \right) b_{n}^{-2\nu} (1+o(1)) \right) \Lambda(x).$$
(3.8)

Noting that

$$F_{\nu}^{n}(a_{n}x+b_{n}) \rightarrow \Lambda(x) = \exp(-\exp(-x)),$$

by taking logarithms, we have

$$n(1-F_{\nu}(a_nx+b_n))\to e^{-x}.$$

Thus

$$1 - F_{\nu}(a_n x + b_n) \sim n^{-1} e^{-x} = o(b_n^{-3\nu})$$

since  $b_n \sim 2^{1/\nu} \lambda (\log n)^{1/\nu}$ , which implies

$$\frac{1}{F_{\nu}(a_n x + b_n)} = 1 + \left(1 - F_{\nu}(a_n x + b_n)\right) \left(1 + o(1)\right) = 1 + o\left(b_n^{-3\nu}\right).$$
(3.9)

The desired result (3.6) follows by (3.8) and (3.9). The proof of (ii) is similar and details are omitted here.  $\hfill \Box$ 

**Lemma 3.3** Let  $f_v(x)$  denote the pdf of GED(v) with  $v \neq 1$ , then

$$f_{\nu}(x) = \left(1 - F_{\nu}(x)\right) \frac{\nu}{2\lambda^{\nu}} x^{\nu-1} \left(1 + 2\left(1 - \nu^{-1}\right)\lambda^{\nu} x^{-\nu} - 4\left(1 - \nu^{-1}\right)\lambda^{2\nu} x^{-2\nu} + O\left(x^{-3\nu}\right)\right)$$
(3.10)

for large x, and

$$C_{n}(x) := \frac{a_{n}f_{\nu}(a_{n}x + b_{n})}{1 - F_{\nu}(a_{n}x + b_{n})}$$
  
= 1 + (1 - \nu^{-1})2\lambda^{\nu}(x + 1)b\_{n}^{-\nu} + 4\lambda^{2\nu}(1 - \nu^{-1})\left(-1 - \frac{x}{\nu} + \frac{1 - 2\nu^{-1}}{2}x^{2}\right)b\_{n}^{-2\nu} + O(b\_{n}^{-3\nu}) (3.11)

as  $n \to \infty$ .

*Proof* The desired results follow directly by (3.1).

### Lemma 3.4 Let

$$H_{\nu}(b_n; x) = \frac{1 - F_{\nu}(a_n x + b_n)}{1 - F_{\nu}(b_n)} e^x - 1$$

and the norming constants  $a_n$  and  $b_n$  be given by (1.2), then

$$b_n^{\nu} \left( b_n^{\nu} H_{\nu}(b_n; x) - k_{\nu 1}(x) \right) \to \omega_{\nu}^{\circ}(x)$$
(3.12)

as  $n \to \infty$ , where  $k_{v1}(x)$  is given by (2.2) and  $\omega_v^{\circ}(x)$  is given by

$$\omega_{\nu}^{\circ}(x) = (1 - \nu^{-1})\lambda^{2\nu} \left( 4x + 2x^{2} + \frac{2}{3}(2 - \nu^{-1})x^{3} + \frac{1}{2}(1 - \nu^{-1})x^{4} \right).$$

Proof Let

$$B_n(x) = \frac{1 + 2(\nu^{-1} - 1)\lambda^{\nu}(a_n x + b_n)^{-\nu} + 4(\nu^{-1} - 1)(\nu^{-1} - 2)\lambda^{2\nu}(a_n x + b_n)^{-2\nu} + O((a_n x + b_n)^{-3\nu})}{1 + 2(\nu^{-1} - 1)\lambda^{\nu}(b_n)^{-\nu} + 4(\nu^{-1} - 1)(\nu^{-1} - 2)\lambda^{2\nu}(b_n)^{-2\nu} + O((b_n)^{-3\nu})},$$

it is easy to check that  $\lim_{n\to\infty} B_n(x) = 1$  and

$$B_n(x) - 1 = (1 + o(1)) \left[ -4(v^{-1} - 1)\lambda^{2\nu} b_n^{-2\nu} x + O(b_n^{-3\nu}) \right].$$

Hence,

$$\lim_{n \to \infty} b_n^{\nu} (B_n(x) - 1) = 0 \tag{3.13}$$

and

$$\lim_{n \to \infty} b_n^{2\nu} (B_{\nu}(x) - 1) = -4 (\nu^{-1} - 1) \lambda^{2\nu} x.$$
(3.14)

By (3.1), we have

$$\frac{1 - F_{\nu}(a_n x + b_n)}{1 - F_{\nu}(b_n)} e^x$$
  
=  $B_n(x) \exp\left[-\int_0^x \left(\frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1\right) dt\right]$ 

$$=B_{n}(x)\left\{1-\int_{0}^{x}\left(\frac{(\nu-1)a_{n}}{b_{n}+a_{n}t}+\frac{\nu a_{n}(b_{n}+a_{n}t)^{\nu-1}}{2\lambda^{\nu}}-1\right)dt +\frac{1}{2}\left[\int_{0}^{x}\left(\frac{(\nu-1)a_{n}}{b_{n}+a_{n}t}+\frac{\nu a_{n}(b_{n}+a_{n}t)^{\nu-1}}{2\lambda^{\nu}}-1\right)dt\right]^{2}(1+o(1))\right\}.$$
(3.15)

Combining (3.13)-(3.15) together, we have

$$\begin{split} \lim_{n \to \infty} b_n^{\nu} H_{\nu}(b_n; x) \\ &= \lim_{n \to \infty} b_n^{\nu} \left( B_n(x) \left\{ 1 - \int_0^x \left( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \right) dt \left( 1 + o(1) \right) \right\} - 1 \right) \\ &= \lim_{n \to \infty} \left( b_n^{\nu} (B_n(x) - 1) - b_n^{\nu} B_n(x) \int_0^x \left( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \right) dt \right) \\ &= -\lim_{n \to \infty} \int_0^x b_n^{\nu} \left( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \right) dt \\ &= -(1 - \nu^{-1})\lambda^{\nu} (x^2 + 2x) \\ &= k_{\nu 1}(x). \end{split}$$

Hence,

$$\begin{split} \lim_{n \to \infty} b_n^{\nu} \left( b_n^{\nu} H_{\nu}(b_n; x) - k_{\nu 1}(x) \right) \\ &= \lim_{n \to \infty} \left[ b_n^{2\nu} \left( B_n(x) - 1 \right) - B_n(x) b_n^{\nu} \right. \\ &\quad \times \left( \int_0^x b_n^{\nu} \left( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n(b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \right) dt + k_{\nu 1}(x) \right) \\ &\quad + \frac{1}{2} \left( 1 + o(1) \right) B_n(x) b_n^{2\nu} \left( \int_0^x \left( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n(b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \right) dt \right)^2 \right] \\ &= -4 \left( \nu^{-1} - 1 \right) \lambda^{2\nu} x + \frac{2\lambda^{2\nu} (1 - \nu^{-1})}{\nu} x^2 - \frac{2}{3} \left( 1 - \nu^{-1} \right) (1 - 2\nu^{-1}) \lambda^{2\nu} x^3 \\ &\quad + \frac{(1 - \nu^{-1})^2}{2} \lambda^{2\nu} (x^2 + 2x)^2 \\ &= \omega_{\nu}^{\circ}(x). \end{split}$$

The proof is complete.

**Lemma 3.5** Let  $C_n(x)$  be given by (3.11) and  $D_n(x)$  be denoted by

$$D_n(x) = 1 + \tilde{k}_{\nu}(x)b_n^{-\nu} + \left(\tilde{\omega}_{\nu}(x) + \frac{\tilde{k}_{\nu}^2(x)}{2}\right)b_n^{-2\nu}(1+o(1)).$$

For  $v \neq 1$  and  $-d \log b_n < x < c b_n^{\frac{3}{\nu}}$  with 0 < c, d < 1, we have

$$\begin{split} \left| C_n(x) D_n(x) \right| < 2, \\ \left| b_n^{\nu} \left( C_n(x) D_n(x) - 1 \right) \right| &\leq 1 + \left( 1 + \nu^{-1} \right) \lambda^{\nu} \left( 2 + 2|x| + \left( x^2 + 2|x| \right) e^{-x} \right), \end{split}$$

$$\begin{split} b_n^{\nu} \Big[ b_n^{\nu} \Big( C_n(x) D_n(x) - 1 \Big) - k_{\nu 2}(x) \Big] \Big| \\ &\leq 1 + (1 + \nu^{-1}) \lambda^{2\nu} \bigg[ 4 + \frac{4}{\nu} |x| + (2 + 4\nu^{-1}) x^2 + \bigg( \frac{4}{\nu} |x| + (4 + 6\nu^{-1}) x^2 \\ &+ \frac{2}{3} \big( 1 + 2\nu^{-1} \big) |x|^3 + \frac{1 + \nu^{-1}}{2} x^4 \bigg) e^{-x} + \frac{1}{2} \big( 1 + \nu^{-1} \big) \big( x^2 + 2|x| \big)^2 e^{-2x} \bigg] \end{split}$$

for large n, where  $k_{\nu 1}(x)$  and  $k_{\nu 2}(x)$  are given by (2.2).

Proof The desired results follow from Lemmas 3.2 and 3.3.

The following Mills' inequalities are from the GED(v) in Jia *et al.* [9], which will be used later.

**Lemma 3.6** Let  $F_{\nu}(x)$  and  $f_{\nu}(x)$  denote the cdf and pdf of GED( $\nu$ ), respectively. Then (i) for  $\nu > 1$  and all x > 0, we have

$$\frac{2\lambda^{\nu}}{\nu}x^{1-\nu}\left(1+\frac{2(\nu-1)\lambda^{\nu}}{\nu}x^{-\nu}\right)^{-1} < \frac{1-F_{\nu}(x)}{f_{\nu}(x)} < \frac{2\lambda^{\nu}}{\nu}x^{1-\nu};$$
(3.16)

(ii) for 0 < v < 1 and all  $x > \lambda [2(1/v - 1)]^{1/v}$ , we have

$$\frac{2\lambda^{\nu}}{\nu}x^{1-\nu} < \frac{1-F_{\nu}(x)}{f_{\nu}(x)} < \frac{2\lambda^{\nu}}{\nu}x^{1-\nu}\left(1+\frac{2(\nu-1)\lambda^{\nu}}{\nu}x^{-\nu}\right)^{-1}.$$
(3.17)

**Lemma 3.7** Let the norming constant  $b_n$  be given by (1.2), for any constant 0 < c < 1 and arbitrary nonnegative integers *i*, *j* and *k*, we have

$$\lim_{n \to \infty} b_n^i \int_{c b_n^{\frac{\nu}{3}}}^{\infty} |x|^j e^{-kx} \Lambda(x) \, dx = 0, \tag{3.18}$$

$$\lim_{n \to \infty} b_n^i \int_{cb_n^{\frac{y}{3}}}^{\infty} |x|^j g_n(x) \, dx = 0 \tag{3.19}$$

*if*  $v \neq 1$ .

*Proof* By arguments similar to Lemma 3.3 in Jia *et al.* [9], we can get (3.18). The rest is to prove (3.19). By (3.16) and (3.17), and Lemma 3.5, we have

$$b_{n}^{i} \int_{cb_{n}^{\frac{v}{3}}}^{\infty} |x|^{j} g_{n}(x) dx$$
  
=  $b_{n}^{i} \int_{cb_{n}^{\frac{v}{3}}}^{\infty} |x|^{j} na_{n} F_{v}^{n-1}(a_{n}x + b_{n}) f_{v}(a_{n}x + b_{n}) dx$   
=  $b_{n}^{i} \int_{cb_{n}^{\frac{v}{3}}}^{\infty} |x|^{j} \frac{1 - F_{v}(a_{n}x + b_{n})}{1 - F_{v}(b_{n})} e^{x} C_{n}(x) D_{n}(x) \Lambda'(x) dx$   
 $\leq 4 b_{n}^{i} \int_{cb_{n}^{\frac{v}{3}}}^{\infty} |x|^{j} e^{-x} \Lambda(x) dx$   
 $\rightarrow 0$ 

as  $n \to \infty$ . The proof is complete.

**Lemma 3.8** Assume that the shape parameter  $v \neq 1$ , then for any constant 0 < d < 1 and arbitrary nonnegative integers *i*, *j* and *k*, we have

$$\lim_{n \to \infty} b_n^i \int_{-\infty}^{-d \log b_n} |x|^j e^{-kx} \Lambda(x) \, dx = 0, \tag{3.20}$$

$$\lim_{n \to \infty} b_n^i \int_{-\infty}^{-d \log b_n} |x|^j g_n(x) \, dx = 0.$$
(3.21)

*Proof* By arguments similar to that of Lemma 3.7, we have

$$b_n^i \int_{-\infty}^{-d \log b_n} |x|^j e^{-kx} \Lambda(x) \, dx$$
  
$$\leq b_n^i \exp\left(-\frac{b_n^d}{2}\right) \int_1^\infty x^j e^{kx} \exp\left(-\frac{e^x}{2}\right) \, dx$$
  
$$\to 0$$

as  $n \to \infty$  since  $\int_1^\infty x^j e^{kx} \exp(-\frac{e^x}{2}) dx < \infty$ .

For assertion (3.21), we only consider the case of  $\nu > 1$  since the proof of the case of  $0 < \nu < 1$  is similar. Rewrite

$$b_n^i \int_{-\infty}^{-d\log b_n} |x|^j g_n(x) \, dx = b_n^i \int_{-\infty}^{-\frac{3\nu\lambda^{-\nu}b_n^{\nu}}{2}} |x|^j g_n(x) \, dx + b_n^i \int_{-\frac{3\nu\lambda^{-\nu}b_n^{\nu}}{2}}^{-\frac{\nu\lambda^{-\nu}b_n^{-2}}{2}} |x|^j g_n(x) \, dx$$
$$+ b_n^i \int_{-\frac{\nu\lambda^{-\nu}b_n^{-2}}{2}}^{-d\log b_n} |x|^j g_n(x) \, dx$$
$$= I_n + II_n + III_n.$$

First note that  $\int_{\mathbb{R}} |x|^j f_{\nu}(x) dx < \infty$  and the symmetry of  $f_{\nu}$  implies  $F_{\nu}(-x) + F_{\nu}(x) = 1$ . By using (1.2) and (3.16) we have

$$I_n < 2^{j} \nu \lambda^{-\nu} b_n^{i+j+\nu} (1 - F_{\nu}(2b_n))^{n-2} < 2^{j} \nu \lambda^{-\nu} b_n^{i+j+\nu} \exp\left( (n-2) \left( c - (\nu-1) \log b_n - \frac{(2b_n)^{\nu}}{2\lambda^{\nu}} \right) \right) \rightarrow 0$$
(3.22)

as  $n \to \infty$ , where  $c = (1 - \nu - 1/\nu) \log 2 + (\nu - 1) \log \lambda - \log \Gamma(1/\nu)$ .

To show  $II_n \rightarrow 0$  and  $III_n \rightarrow 0$ , we consider the case of  $\nu > 2$  first. By using the following inequalities

$$1 - \nu x < (1 - x)^{\nu} < 1 - \nu x + \frac{\nu(\nu - 1)}{2} x^{2}, \quad 0 < x < \frac{1}{2}, \nu > 2,$$
(3.23)

we can get

$$na_n f_{\nu} \left( b_n - b_n^{\frac{1-\nu}{2}} \right) < 2 \exp\left(\frac{\nu}{2\lambda^{\nu}} b_n^{\frac{\nu-1}{2}}\right)$$

$$\frac{1 - F_{\nu}(b_n - b_n^{\frac{1-\nu}{2}})}{1 - F_{\nu}(b_n)}$$

$$> \frac{\frac{2\lambda^{\nu}}{\nu}(b_n - b_n^{\frac{1-\nu}{2}})^{1-\nu}(1 + \frac{2(\nu-1)\lambda^{\nu}}{\nu}(b_n - b_n^{\frac{1-\nu}{2}})^{-\nu})^{-1}f_{\nu}(b_n - b_n^{\frac{1-\nu}{2}})}{\frac{2\lambda^{\nu}}{\nu}b_n^{1-\nu}f_{\nu}(b_n)}$$

$$> \frac{(1 - b_n^{-\frac{1+\nu}{2}})^{1-\nu}}{1 + \frac{2(\nu-1)\lambda^{\nu}}{\nu}(b_n - b_n^{\frac{1-\nu}{2}})^{-\nu}} \exp\left(\frac{\nu}{2\lambda^{\nu}}b_n^{\frac{\nu-1}{2}} - \frac{\nu(\nu-1)}{4\lambda^{\nu}}b_n^{-1}\right)$$

$$> \frac{1}{2}\exp\left(\frac{\nu}{2\lambda^{\nu}}b_n^{\frac{\nu-1}{2}}\right)$$

for large *n*, which implies that

$$\begin{aligned} H_{n} &< b_{n}^{i} \left(\frac{3\nu\lambda^{-\nu}b_{n}^{\nu}}{2}\right)^{j} na_{n} f_{\nu} \left(-a_{n} \frac{\nu\lambda^{-\nu}b_{n}^{\frac{\nu-1}{2}}}{2} + b_{n}\right) F_{\nu}^{n-1} \left(-a_{n} \frac{\nu\lambda^{-\nu}b_{n}^{\frac{\nu-1}{2}}}{2} + b_{n}\right) \frac{3\nu\lambda^{-\nu}b_{n}^{\nu}}{2} \\ &< b_{n}^{i} \left(\frac{3\nu\lambda^{-\nu}b_{n}^{\nu}}{2}\right)^{j+1} na_{n} f_{\nu} \left(-a_{n} \frac{\nu\lambda^{-\nu}b_{n}^{\frac{\nu-1}{2}}}{2} + b_{n}\right) \exp\left(-(n-1)\left(1 - F_{\nu}(b_{n} - b_{n}^{\frac{1-\nu}{2}})\right)\right) \\ &= 2b_{n}^{i} \left(\frac{3\nu\lambda^{-\nu}b_{n}^{\nu}}{2}\right)^{j+1} \frac{\lambda 2^{1+\frac{1}{\nu}}\Gamma(\frac{1}{\nu})}{\nu} \exp\left(\frac{b_{n}^{\nu} + \nu b_{n}^{\frac{\nu-1}{2}}}{2\lambda^{\nu}} - \frac{1}{2}\exp\left(\frac{\nu}{2\lambda^{\nu}}b_{n}^{\frac{\nu-1}{2}}\right)\right) \\ &\to 0 \end{aligned}$$
(3.24)

as  $n \to \infty$ . Similarly, for  $-\frac{\nu \lambda^{-\nu} b_n^{\frac{\nu-1}{2}}}{2} < x < -d \log b_n$ , we have

$$na_n f_{\nu}(a_n x + b_n)$$

$$< \left(1 + \frac{2(\nu - 1)\lambda^{\nu}}{\nu} b_n^{-\nu}\right) \frac{f_{\nu}(a_n x + b_n)}{f_{\nu}(b_n)}$$

$$< \left(1 + \frac{2(\nu - 1)\lambda^{\nu}}{\nu} b_n^{-\nu}\right) \exp\left(-\frac{b_n^{\nu}}{2\lambda^{\nu}} 2\lambda^{\nu} b_n^{-\nu} x\right)$$

$$< 2e^{-x}$$

and

$$\begin{aligned} &\frac{1-F_{\nu}(a_{n}x+b_{n})}{1-F_{\nu}(b_{n})} \\ &> \frac{\frac{2\lambda^{\nu}}{\nu}(a_{n}x+b_{n})^{1-\nu}(1+\frac{2(\nu-1)\lambda^{\nu}}{\nu}(a_{n}x+b_{n})^{-\nu})^{-1}f_{\nu}(a_{n}x+b_{n})}{\frac{2\lambda^{\nu}}{\nu}b_{n}^{1-\nu}f_{\nu}(b_{n})} \\ &> \frac{(1+\frac{2\lambda^{\nu}}{\nu}b_{n}^{-\nu}x)^{1-\nu}}{1+\frac{2(\nu-1)\lambda^{\nu}}{\nu}(a_{n}x+b_{n})^{-\nu}}\exp\left(-x-\frac{\nu-1}{\nu}\lambda^{\nu}b_{n}^{-\nu}x^{2}\right) \\ &> \frac{1}{2}e^{-x} \end{aligned}$$

and

for large *n* by using (3.23). Then

$$III_{n} < 2b_{n}^{i} \int_{-\frac{\nu\lambda^{-\nu}b_{n}^{2}}{2}}^{-d\log b_{n}} |x|^{j} e^{-x} \exp\left(-\frac{1}{2}e^{-x}\right) dx$$
  
$$< 2b_{n}^{i} \exp\left(-\frac{e^{d\log b_{n}}}{4}\right) \int_{-\frac{\nu\lambda^{-\nu}b_{n}^{2}}{2}}^{-d\log b_{n}} |x|^{j} e^{-x} \exp\left(-\frac{1}{4}e^{-x}\right) dx$$
  
$$\to 0$$
(3.25)

as  $n \to \infty$ .

Combining with (3.22)-(3.25), the assertion (3.21) is derived for  $\nu > 2$ . Similar proofs for the case of  $1 < \nu \le 2$  and details are omitted here. The proof is complete.

**Lemma 3.9** Let  $\alpha = \min(1, v)$  as  $v \neq 1$ . For large n and  $-d \log b_n < x < cb_n^{\frac{1}{3}}$ , both  $x^r b_n^v \Delta_n(g_n, \Lambda'; x)$  and  $x^r b_n^v [b_n^v \Delta_n(g_n, \Lambda'; x) - k_v(x)\Lambda'(x)]$  are bounded by integrable functions independent of n, with r > 0, 0 < c < 1 and  $0 < d < \alpha$ , where  $a_n$  and  $b_n$  are given by (1.2), and  $k_v(x)$  is given by (2.2).

*Proof* We only consider the case of v > 1. For the case of 0 < v < 1, the proofs are similar and details are omitted here. Rewrite

$$b_{n}^{\nu}\Delta_{n}(g_{n},\Lambda';x) = b_{n}^{\nu}H_{\nu}(b_{n};x)e^{-x}\Lambda(x) + b_{n}^{\nu}(C_{n}(x)D_{n}(x)-1)e^{-x}\Lambda(x),$$

where  $C_n(x)$  is given by (3.11),  $H_\nu(b_n; x)$  and  $D_n(x)$  are respectively defined in Lemma 3.4 and Lemma 3.5. Note that  $\int_{-\infty}^{\infty} x^k e^{-tx} \exp(-e^{-x}) dx = (-1)^k \Gamma^{(k)}(t)$  is finite for t > 0 and nonnegative integers k. Lemma 3.5 shows that  $b_n^\nu(C_n(x)D_n(x) - 1)e^{-x}\Lambda(x)$  is bounded by integrable function independent of n. The rest is to prove that  $b_n^\nu H_\nu(b_n; x)$  is bounded by m(x), where m(x) is a polynomial on x. Rewrite

$$b_n^{\nu} H_{\nu}(b_n; x) = I_n(x) + J_n(x), \qquad (3.26)$$

where

$$I_n(x) = b_n^{\nu} (B_n(x) - 1),$$
  

$$J_n(x) = b_n^{\nu} B_n(x) \int_0^x \left( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \right) dt$$

For  $-d \log b_n < x < c b_n^{\frac{y}{3}}$ , from Lemma 3.4 it follows that

$$\left|I_n(x)\right| < 1 \tag{3.27}$$

and

$$\left|J_{n}(x)\right| < 1 + \left[\frac{\nu^{-1}}{2^{-1}\nu\lambda^{-\nu} - d}|x| + (1 - \nu^{-1})\lambda^{\nu}x^{2} + \frac{4}{3}(1 - \nu^{-1})\left|1 - 2\nu^{-1}\right|\lambda^{2\nu}|x|^{3}\right].$$
 (3.28)

Hence, the desired result (3.26) follows by combining (3.16), (3.27) and (3.28) together.

Rewrite

$$\begin{split} b_n^{\nu} \Big[ b_n^{\nu} \Delta_n \big( g_n, \Lambda'; x \big) - k_{\nu}(x) \Lambda'(x) \Big] \\ &= b_n^{\nu} \Big[ b_n^{\nu} \bigg( \frac{1 - F_{\nu}(a_n x + b_n)}{1 - F_{\nu}(b_n)} e^x C_n(x) D_n(x) - 1 \bigg) - k_{\nu}(x) \Big] \Lambda'(x) \\ &= b_n^{\nu} \Big[ b_n^{\nu} C_n(x) D_n(x) \bigg( \frac{1 - F_{\nu}(a_n x + b_n)}{1 - F_{\nu}(b_n)} e^x - 1 \bigg) \\ &+ b_n^{\nu} \big( C_n(x) D_n(x) - 1 \big) - \big( k_{\nu 1}(x) + k_{\nu 2}(x) \big) \Big] \Lambda'(x) \\ &= b_n^{\nu} \Big[ b_n^{\nu} C_n(x) D_n(x) \big( H_{\nu}(b_n; x) - b_n^{-\nu} k_{\nu 1}(x) \big) + b_n^{\nu} \big( C_n(x) D_n(x) - 1 - b_n^{-\nu} k_{\nu 2}(x) \big) \Big] \\ &+ \big( C_n(x) D_n(x) - 1 \big) k_{\nu 1}(x) \Big] \Lambda'(x). \end{split}$$

By Lemma 3.5, we only need to estimate the bound of  $b_n^{\nu}[b_n^{\nu}C_n(x)D_n(x)(H_{\nu}(b_n;x) - b_n^{-\nu}k_{\nu 1}(x))]$ . Rewrite

$$b_n^{\nu} \left[ b_n^{\nu} C_n(x) D_n(x) \left( H_{\nu}(b_n; x) - b_n^{-\nu} k_{\nu 1}(x) \right) \right] = H_n(x) - K_n(x) + L_n(x),$$

where

$$\begin{split} H_n(x) &= b_n^{2\nu} \Big( B_n(x) - 1 \Big), \\ K_n(x) &= b_n^{2\nu} \bigg( B_n(x) \int_0^x \bigg( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \bigg) dt + b_n^{-\nu} k_{\nu 1}(x) \bigg), \\ L_n(x) &= \frac{1}{2} \Big( 1 + o(1) \Big) B_n(x) b_n^{2\nu} \bigg( \int_0^x \bigg( \frac{(\nu - 1)a_n}{b_n + a_n t} + \frac{\nu a_n (b_n + a_n t)^{\nu - 1}}{2\lambda^{\nu}} - 1 \bigg) dt \bigg)^2. \end{split}$$

For  $0 < x < c b_n^{\frac{y}{3}}$ , by using  $1 - \nu y < (1 + y)^{-\nu} < 1$  for  $\nu > 0$  and y > 0, we have

$$\left|H_n(x)\right| < 1 + 8\lambda^{2\nu} \left(1 + 8\lambda^{\nu}\right) x$$

due to Lemma 3.4. If  $-d \log b_n < x < 0$ , by using  $1 + vy < (1 + y)^v < 1$  for v > 1 and -1 < y < 0, Lemma 3.4 shows that

$$|H_n(x)| < 1 + 16\lambda^{2\nu} (1 + 8\lambda^{\nu})|x|$$

for large *n*. Similarly,

$$\left|K_{n}(x)\right| < 2\left(1-\nu^{-1}\right)\left|\frac{\nu}{2\lambda^{\nu}}-d\right|^{-1}\lambda^{\nu}x^{2}+\frac{4}{3}\left(1-\nu^{-1}\right)\left|1-2\nu^{-1}\right|\lambda^{2\nu}|x|^{3},$$
(3.29)

$$\left|L_{n}(x)\right| < 1 + \frac{1}{2} \left[\frac{\nu^{-1}}{2^{-1}\nu\lambda^{-\nu} - d}|x| + (1 - \nu^{-1})\lambda^{\nu}x^{2} + \frac{4}{3}(1 - \nu^{-1})\left|1 - 2\nu^{-1}\right|\lambda^{2\nu}|x|^{3}\right]^{2} \quad (3.30)$$

as  $-d \log b_n < x < c b_n^{\frac{y}{3}}$ . Hence, we derive the desired result by combining (3.16), (3.29) and (3.30) together.

The proof is complete.

For v = 1, note that the GED(1) is the Laplace distribution with pdf given by

$$f_1(x) = 2^{-\frac{1}{2}} \exp\left(-2^{\frac{1}{2}}|x|\right), \quad x \in \mathbb{R},$$

and its distributional tail can be written as

$$1 - F_1(x) = 2^{-\frac{1}{2}} f_1(x) = 2^{-1} \exp\left(-2^{\frac{1}{2}}\right) \exp\left(-\int_1^x \frac{1}{f(t)} dt\right), \quad x > 0$$

with  $f(t) = 2^{-\frac{1}{2}}$ . For the Laplace distribution, similar to the case of  $\nu > 1$ , we have the following two results.

**Lemma 3.10** For 0 < d < 1 and an arbitrary nonnegative real number *j*, we have

$$\lim_{n \to \infty} n^2 \int_{-\infty}^{-db_n^{\frac{1}{2}}} |x|^j e^{-kx} \Lambda(x) \, dx = 0, \quad k = 1, 2, \dots,$$
(3.31)

$$\lim_{n \to \infty} n^2 \int_{-\infty}^{-db_n^{\frac{1}{2}}} |x|^j F_1^n(a_n x + b_n) \, dx = 0.$$
(3.32)

**Lemma 3.11** For  $x > -db_n^{\frac{1}{2}}$ , both  $x^r n((F_1^n(a_nx + b_n))' - \Lambda'(x))$  and  $x^r n[n((F_1^n(a_nx + b_n))' - \Lambda'(x)) + \frac{1}{2}e^{-2x}\Lambda'(x)]$  are bounded by integrable functions independent of n, where r > 0 and 0 < d < 1.

# 4 Proofs of the main results

Proof of Theorem 2.1 From Lemma 3.3 and Lemma 3.5 it follows that

$$C_{n}(x)D_{n}(x) = 1 + k_{\nu2}(x)b_{n}^{-\nu} + (1 - \nu^{-1})\lambda^{2\nu} \bigg[ -4 - \frac{4}{\nu}x + 2(1 - 2\nu^{-1})x^{2} + \bigg(-\frac{4}{\nu}x + 2(2 - 3\nu^{-1})x^{2} + \frac{2(1 - 2\nu^{-1})}{3}x^{3} - \frac{1 - \nu^{-1}}{2}x^{4}\bigg)e^{-x} + \frac{(1 - \nu^{-1})}{2}(x^{2} + 2x)^{2}e^{-2x}\bigg]b_{n}^{-2\nu} + O(b_{n}^{-3\nu}), \quad (4.1)$$

where  $k_{\nu 2}(x)$  is given by (2.2). Note that (3.12) shows

$$\frac{1 - F_{\nu}(a_n x + b_n)}{1 - F_{\nu}(b_n)} e^x = 1 + k_{\nu 1}(x) b_n^{-\nu} + \omega_{\nu}^{\circ}(x) b_n^{-2\nu} + O(b_n^{-3\nu}).$$
(4.2)

Hence,

$$\begin{split} \Delta_n(g_n,\Lambda';x) \\ &= \left(\frac{1-F_\nu(a_nx+b_n)}{1-F_\nu(b_n)}e^xC_n(x)D_n(x)-1\right)\Lambda'(x) \\ &= \left[\left(1-\nu^{-1}\right)\lambda^\nu(2-x^2+(x^2+2x)e^{-x})b_n^{-\nu}+(1-\nu^{-1})\lambda^{2\nu}\left(-4-2(1-\nu^{-1})x^2\right)\right. \\ &\left. -\frac{2(1-2\nu^{-1})}{3}x^3+\frac{1-\nu^{-1}}{2}x^4+\left(-\frac{4}{\nu}x-\frac{2}{\nu}x^2-\frac{2(5-4\nu^{-1})}{3}x^3-\frac{3(1-\nu^{-1})}{2}x^4\right)e^{-x} \end{split}$$

implying that

$$\lim_{n\to\infty} b_n^{\nu} \Big[ b_n^{\nu} \Delta_n \big( g_n, \Lambda'; x \big) - k_{\nu}(x) \Lambda'(x) \Big] = \omega_{\nu}(x) \Lambda'(x).$$

The proof is complete.

*Proof of Theorem* 2.3 For  $\nu \neq 1$ , by Lemmas 3.5-3.9 and the dominated convergence theorem, we have

$$b_n^{\nu}(m_r(n) - m_r)$$

$$= b_n^{\nu} \int_{-\infty}^{-d \log b_n} x^r \Delta_n(g_n, \Lambda'; x) \, dx + b_n^{\nu} \int_{-d \log b_n}^{c b_n^{\frac{\nu}{3}}} x^r \Delta_n(g_n, \Lambda'; x) \, dx$$

$$+ b_n^{\nu} \int_{c b_n^{\frac{\nu}{3}}}^{\infty} x^r \Delta_n(g_n, \Lambda'; x) \, dx$$

$$\rightarrow \int_{-\infty}^{\infty} x^r k_{\nu}(x) \Lambda'(x) \, dx$$

$$= -(1 - \nu^{-1}) \lambda^{\nu} r(2m_r + m_{r+1})$$

and

$$\begin{split} b_n^{\nu} \left[ b_n^{\nu} (m_r(n) - m_r) + (1 - \nu^{-1}) \lambda^{\nu} r(2m_r + m_{r+1}) \right] \\ &= \int_{-\infty}^{\infty} x^r b_n^{\nu} \left[ b_n^{\nu} \Delta_n (g_n, \Lambda'; x) - k_{\nu}(x) \Lambda'(x) \right] dx \\ &= \int_{-\infty}^{-d \log b_n} x^r b_n^{2\nu} \Delta_n (g_n, \Lambda'; x) dx - \int_{-\infty}^{-d \log b_n} x^r b_n^{\nu} k_{\nu}(x) \Lambda'(x) dx \\ &- \int_{cb_n^{\frac{\nu}{3}}}^{\infty} x^r b_n^{\nu} k_{\nu}(x) \Lambda'(x) dx + \int_{-d \log b_n}^{cb_n^{\frac{\nu}{3}}} x^r b_n^{\nu} \left[ b_n^{\nu} \Delta_n (g_n, \Lambda'; x) - k_{\nu}(x) \Lambda'(x) \right] dx \\ &+ \int_{cb_n^{\frac{\nu}{3}}}^{\infty} x^r b_n^{2\nu} \Delta_n (g_n, \Lambda'; x) dx \\ &\rightarrow \int_{-\infty}^{\infty} x^r \omega_{\nu}(x) \Lambda'(x) dx \\ &= 2r \lambda^{2\nu} (1 - \nu^{-1}) \left[ \left( (1 - \nu^{-1})(r + 1) + 2 \right) m_r + \left( (1 - \nu^{-1})(r + 1) + 1 \right) m_{r+1} \\ &+ \left( \frac{1}{4} (1 - \nu^{-1})(r - 1) + \frac{1}{3} (2 - \nu^{-1}) \right) m_{r+2} \right] \end{split}$$

as  $n \to \infty$ .

For the case of v = 1, note that

$$\int_{-\infty}^{\infty} x^k e^{-mx} \Lambda(x) \, dx = (-1)^k \Gamma^{(k)}(m).$$

Combining with Lemmas 3.10 and 3.11 and the dominated convergence theorem, we can derive the desired results.

The proof is complete.

### **Competing interests**

The authors declare that they have no competing interests.

### Authors' contributions

CL drafted the manuscript and TL revised the whole paper critically. Both authors were involved in the proof of the main results of the paper. All authors read and approved the final manuscript.

### Acknowledgements

We would like to appreciate the reviewers for reading the paper and making helpful comments that improved the original paper. This work was supported by the National Natural Science Foundation of China (11171275), the Natural Science Foundation Project of CQ (cstc2012jjA00029), the Fundamental Research Funds for the Central Universities (XDJK2013C021) and the Doctoral Grant of Southwest University (SWU113012).

### Received: 12 May 2015 Accepted: 1 November 2015 Published online: 10 November 2015

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