

RESEARCH

Open Access

Anisotropic curl-free wavelets with boundary conditions

Yingchun Jiang*

*Correspondence:
guilinjiang@126.com
School of Mathematics and
Computational Science, Guilin
University of Electronic Technology,
Guilin, 541004, P.R. China

Abstract

This paper deals with the construction of anisotropic curl-free wavelets that satisfy the tangent boundary conditions on bounded domains. Based on some assumptions, we first obtain the desired curl-free Riesz wavelet bases through the orthogonal decomposition of vector-valued L^2 . Next, the characterization of Sobolev spaces is studied. Finally, we give the concrete construction of wavelets satisfying the initial assumptions.

MSC: 42C20

Keywords: anisotropic; curl-free; wavelets; bounded domains; boundary conditions

1 Introduction

Due to their potential use in many physical problems, like the simulation of incompressible fluids or electromagnetism, curl-free wavelet bases have been advocated in several papers and all results focus on the cases of R^2 and R^3 [1–4]. Moreover, it is questionable whether they are appropriately called bases and whether they can be used to characterize Sobolev spaces. However, it is reasonable to study the corresponding wavelet bases on bounded domains because of some practical use. At the same time, the boundary conditions, the stability and the characterization of Sobolev spaces are also necessary in some applications such as adaptive wavelet methods. In references [5, 6], anisotropic divergence-free wavelets which satisfy the specific boundary conditions on the hypercube are studied. Inspired by the fact that a div-free space and a curl-free space form the orthogonal Helmholtz decomposition, we mainly study the anisotropic curl-free wavelet bases satisfying the tangent boundary conditions on bounded domains in this paper, which is organized as follows. In Section 2, based on some assumption, the desired curl-free wavelets are constructed through the orthogonal decomposition of vector-valued L^2 . Section 3 is devoted to studying the characterization of Sobolev spaces. We give the concrete construction of wavelets satisfying the initial assumption in the final section.

For two 2D vectors $\vec{u} = (u_1, u_2)^T$ and $\vec{v} = (v_1, v_2)^T$, $\vec{u} \times \vec{v}$ is defined as

$$\vec{u} \times \vec{v} =: u_1 v_2 - u_2 v_1.$$

Then for $\vec{u}(x, y) = (u_1(x, y), u_2(x, y))^T$, we define the 2D curl-operator by

$$\text{curl} \vec{u} =: (\partial_1, \partial_2) \times \vec{u} = \partial_1 u_2 - \partial_2 u_1$$

and for $\vec{u}(x, y, z) = (u_1, u_2, u_3)^T$, the 3D curl-operator is defined by

$$\text{curl}\vec{u} = (\partial_1, \partial_2, \partial_3) \times \vec{u} = (\partial_2 u_3 - \partial_3 u_2, \partial_3 u_1 - \partial_1 u_3, \partial_1 u_2 - \partial_2 u_1)^T.$$

2 Decomposition of $L^2(I^n)^n$

In this part, we will construct curl-free wavelets that satisfy tangent boundary conditions by the orthogonal decomposition of vector-valued L^2 .

Let $I = (0, 1)$. For $n = 2, 3$, we firstly define the following spaces:

$$\begin{aligned} H(\text{curl}; I^n) &=: \{\vec{u} \in L^2(I^n)^n : \text{curl}\vec{u} \in L^2(I^2) \text{ or } L^2(I^3)^3\}, \\ H_0(\text{curl}; I^n) &=: \{\vec{u} \in H(\text{curl}; I^n) : \vec{u} \times \vec{n} = 0 \text{ or } \vec{0} \text{ on boundary } \Gamma\}, \\ \mathcal{H}(I^n) &=: H_0(\text{curl}0; I^n) =: \{\vec{u} \in H_0(\text{curl}; I^n) : \text{curl}\vec{u} = 0 \text{ or } \vec{0}\}. \end{aligned}$$

For a scalar function $\phi(x, y)$, define $\overrightarrow{\text{curl}}\phi = (\partial_2\phi, -\partial_1\phi)^T$. Then integration by parts shows

$$\mathcal{H}(I^2) \perp \overrightarrow{\text{curl}}H^1(I^2) \quad \text{and} \quad \mathcal{H}(I^3) \perp \text{curl}H^1(I^3)^3.$$

Let $L^2 =: L^2(I)$, $L^{2,0} =: \{u \in L^2 : \int_0^1 u(x) dx = 0\}$. Furthermore, set

$$\begin{aligned} \widehat{L^2(I^2)}^2 &=: L^{2,0} \otimes L^2 \times L^2 \otimes L^{2,0}; \\ \widehat{L^2(I^3)}^3 &=: L^{2,0} \otimes L^2 \otimes L^2 \times L^2 \otimes L^{2,0} \otimes L^2 \times L^2 \otimes L^2 \otimes L^{2,0}; \\ \widehat{H^s(I^n)} &=: H^s(I^n) \cap (L^{2,0} \otimes \dots \otimes L^{2,0}), \quad n = 1, 2. \end{aligned}$$

For $n = 3$, we define $\widehat{H}_1^s(I^3) =: H^s(I^3) \cap (L^2 \otimes L^{2,0} \otimes L^{2,0})$ and

$$\widehat{H}_2^s(I^3) =: H^s(I^3) \cap (L^{2,0} \otimes L^2 \otimes L^{2,0}), \quad \widehat{H}_3^s(I^3) =: H^s(I^3) \cap (L^{2,0} \otimes L^{2,0} \otimes L^2).$$

Finally, let $\widehat{\mathcal{H}}(I^n) =: \mathcal{H}(I^n) \cap \widehat{L^2(I^n)^n}$, $n = 2, 3$.

The following result will be proved in Section 4:

Assumption 2.1 There exist bi-orthogonal Riesz bases $\Psi^{(n)} = \Psi_{\text{curl}}^{(n)} \cup \Psi_{\text{comp}}^{(n)}$ and $\widetilde{\Psi}^{(n)} = \widetilde{\Psi}_{\text{curl}}^{(n)} \cup \widetilde{\Psi}_{\text{comp}}^{(n)}$ for $\widehat{L^2(I^n)^n}$ (of wavelet type) such that

$$\begin{aligned} \Psi_{\text{curl}}^{(n)} &\subset \widehat{\mathcal{H}}(I^n) \quad (n = 2, 3), \quad \widetilde{\Psi}_{\text{comp}}^{(2)} \subset \overrightarrow{\text{curl}}\widehat{H}^1(I^2) \quad \text{or} \\ \widetilde{\Psi}_{\text{comp}}^{(3)} &\subset \text{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3)). \end{aligned}$$

Proposition 2.1 It holds that $\Psi_{\text{curl}}^{(n)}$, $\widetilde{\Psi}_{\text{comp}}^{(2)}$ and $\widetilde{\Psi}_{\text{comp}}^{(3)}$ are Riesz bases for $\widehat{\mathcal{H}}(I^n)$ ($n = 2, 3$), $\overrightarrow{\text{curl}}\widehat{H}^1(I^2)$ and $\text{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3))$, respectively.

Proof For any $\vec{u} \in \widehat{\mathcal{H}}(I^n)$ ($n = 2, 3$), we know

$$\widehat{\mathcal{H}}(I^2) \perp \overrightarrow{\text{curl}}\widehat{H}^1(I^2) \quad \text{and} \quad \widehat{\mathcal{H}}(I^3) \perp \text{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3)),$$

then $\vec{u} = \langle \vec{u}, \tilde{\Psi}^{(n)} \rangle_{L^2(I^n)^n} \Psi^{(n)} = \langle \vec{u}, \tilde{\Psi}_{curl}^{(n)} \rangle_{L^2(I^n)^n} \Psi_{curl}^{(n)}$ with $\|\vec{u}\|_{L^2(I^n)^n} \simeq \|\langle \vec{u}, \tilde{\Psi}_{curl}^{(n)} \rangle_{L^2(I^n)^n}\|_{\ell^2}$. Finally, it is easy to verify by the definition of \overrightarrow{curl} and $curl$ that

$$\overrightarrow{curl}\widehat{H^1(I^2)} \subset \widehat{L^2(I^2)^2} \quad \text{and} \quad curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)}) \subset \widehat{L^2(I^3)^3},$$

the remaining results can be proved similarly. □

Proposition 2.2 *The following decompositions hold:*

$$\begin{aligned} \widehat{L^2(I^2)^2} &= \widehat{H(I^2)} \oplus^\perp \overrightarrow{curl}\widehat{H^1(I^2)}, \\ \widehat{L^2(I^3)^3} &= \widehat{H(I^3)} \oplus^\perp curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)}). \end{aligned}$$

Proof We only prove the case of $n = 3$, the others can be proved similarly. Since

$$\vec{u} = \langle \vec{u}, \Psi_{curl}^{(3)} \rangle_{L^2(I^3)^3} \tilde{\Psi}_{curl}^{(3)} + \langle \vec{u}, \Psi_{comp}^{(3)} \rangle_{L^2(I^3)^3} \tilde{\Psi}_{comp}^{(3)} = \langle \vec{u}, \Psi_{comp}^{(3)} \rangle_{L^2(I^3)^3} \tilde{\Psi}_{comp}^{(3)}$$

for any $\vec{u} \in \widehat{H(I^3)}^\perp$, then $\widehat{H(I^3)}^\perp \subseteq curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)})$. On the other hand, since $\widehat{H(I^3)} \perp curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)})$, then $curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)}) \subseteq \widehat{H(I^3)}^\perp$. Therefore, $\widehat{L^2(I^3)^3} = \widehat{H(I^3)} \oplus^\perp curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)})$. □

Now, we consider the orthogonal decomposition of $L^2(I^n)^n$. Let $L^2 = L^{2,0} \oplus^\perp \Lambda$. Then there are the following orthogonal decompositions:

$$\begin{aligned} L^2 \otimes L^2 &= L^{2,0} \otimes L^2 \oplus^\perp \Lambda \otimes L^2, & L^2 \otimes L^2 &= L^2 \otimes L^{2,0} \oplus^\perp L^2 \otimes \Lambda; \\ L^2 \otimes L^2 \otimes L^2 &= L^{2,0} \otimes L^2 \otimes L^2 \oplus^\perp \Lambda \otimes L^2 \otimes L^2; \\ L^2 \otimes L^2 \otimes L^2 &= L^2 \otimes L^{2,0} \otimes L^2 \oplus^\perp L^2 \otimes \Lambda \otimes L^2; \\ L^2 \otimes L^2 \otimes L^2 &= L^2 \otimes L^2 \otimes L^{2,0} \oplus^\perp L^2 \otimes L^2 \otimes \Lambda. \end{aligned}$$

Therefore, we obtain the following decomposition:

$$L^2(I^2)^2 = \widehat{L^2(I^2)^2} \oplus^\perp \begin{pmatrix} \Lambda \otimes L^2 \\ 0 \end{pmatrix} \oplus^\perp \begin{pmatrix} 0 \\ L^2 \otimes \Lambda \end{pmatrix}, \tag{2.1}$$

$$\begin{aligned} L^2(I^3)^3 &= \widehat{L^2(I^3)^3} \oplus^\perp \begin{pmatrix} \Lambda \otimes L^2 \otimes L^2 \\ 0 \\ 0 \end{pmatrix} \\ &\oplus^\perp \begin{pmatrix} 0 \\ L^2 \otimes \Lambda \otimes L^2 \\ 0 \end{pmatrix} \oplus^\perp \begin{pmatrix} 0 \\ 0 \\ L^2 \otimes L^2 \otimes \Lambda \end{pmatrix}. \end{aligned} \tag{2.2}$$

By Proposition 2.2, $\widehat{L^2(I^2)^2} = \widehat{H(I^2)} \oplus^\perp \overrightarrow{curl}\widehat{H^1(I^2)}$, $\widehat{L^2(I^3)^3} = \widehat{H(I^3)} \oplus^\perp curl(\widehat{H^1(I^3)} \times \widehat{H^1(I^3)} \times \widehat{H^1(I^3)})$. Moreover,

$$\begin{pmatrix} \Lambda \otimes L^2 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ L^2 \otimes \Lambda \end{pmatrix} \subset \overrightarrow{curl}\widehat{H^1(I^2)},$$

$$\begin{pmatrix} \Lambda \otimes L^2 \otimes L^2 \\ 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ L^2 \otimes \Lambda \otimes L^2 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ L^2 \otimes L^2 \otimes \Lambda \end{pmatrix} \subset \text{curl}H^1(I^3)^3.$$

Finally, we obtain $L^2(I^2)^2 = \mathcal{H}(I^2) \oplus^\perp \overrightarrow{\text{curl}}H^1(I^2)$ and $L^2(I^3)^3 = \mathcal{H}(I^3) \oplus^\perp \text{curl}H^1(I^3)^3$.

Now, we will construct Riesz bases for $\mathcal{H}(I^n)$ ($n = 2, 3$), $\overrightarrow{\text{curl}}H^1(I^2)$ and $\text{curl}H^1(I^3)^3$. For $n = 2$, we define the embedding $E_{\{1\}}^{(2)}, E_{\{2\}}^{(2)} : L^2(I) \rightarrow L^2(I^2)^2$ by

$$\begin{aligned} (E_{\{1\}}^{(2)}v)(x_1, x_2) &= v(x_2)\vec{e}_1, & (E_{\{2\}}^{(2)}v)(x_1, x_2) &= v(x_1)\vec{e}_2, \\ E_{\{1,2\}}^{(2)} : \widehat{L^2(I^2)^2} &\rightarrow L^2(I^2)^2 & \text{by } (E_{\{1,2\}}^{(2)}\vec{v})(x_1, x_2) &= \sum_{i=1}^2 v_i(x_1, x_2)\vec{e}_i. \end{aligned}$$

For $n = 3$, define $E_{\{1\}}^{(3)}, E_{\{2\}}^{(3)}, E_{\{3\}}^{(3)} : L^2(I) \rightarrow L^2(I^3)^3$ by

$$\begin{aligned} (E_{\{1\}}^{(3)}v) &= v(x_2)v(x_3)\vec{e}_1, & (E_{\{2\}}^{(3)}v) &= v(x_1)v(x_3)\vec{e}_2, & (E_{\{3\}}^{(3)}v) &= v(x_1)v(x_2)\vec{e}_3, \\ E_{\{1,2,3\}}^{(3)} : \widehat{L^2(I^3)^3} &\rightarrow L^2(I^3)^3 & \text{by } (E_{\{1,2,3\}}^{(3)}\vec{v})(x_1, x_2, x_3) &= \sum_{i=1}^3 v_i(x_1, x_2, x_3)\vec{e}_i. \end{aligned}$$

It is obvious that $E_{\{1,2\}}^{(2)} = I, E_{\{1,2,3\}}^{(3)} = I$. Moreover, the image satisfies

$$\begin{aligned} \text{Im}E_{\{1\}}^{(2)} &= (\Lambda \otimes L^2, 0)^T, & \text{Im}E_{\{2\}}^{(2)} &= (0, L^2 \otimes \Lambda)^T; \\ \text{Im}E_{\{1\}}^{(3)} &= (\Lambda \otimes L^2 \otimes L^2, 0, 0)^T, & \text{Im}E_{\{2\}}^{(3)} &= (0, L^2 \otimes \Lambda \otimes L^2, 0)^T, \\ \text{Im}E_{\{3\}}^{(3)} &= (0, 0, L^2 \otimes L^2 \otimes \Lambda)^T. \end{aligned}$$

Furthermore, we know from (2.1) and (2.2) that $L^2(I^2)^2 = \text{Im}E_{\{1\}}^{(2)} \oplus^\perp \text{Im}E_{\{2\}}^{(2)} \oplus^\perp \text{Im}E_{\{1,2\}}^{(2)}$,

$$L^2(I^3)^3 = \text{Im}E_{\{1\}}^{(3)} \oplus^\perp \text{Im}E_{\{2\}}^{(3)} \oplus^\perp \text{Im}E_{\{3\}}^{(3)} \oplus^\perp \text{Im}E_{\{1,2,3\}}^{(3)}.$$

Since $\text{Im}(E_{\{1,2\}}^{(2)}|_{\widehat{\mathcal{H}}(I^2)}) \subset \mathcal{H}(I^2)$, $\text{Im}(E_{\{1,2\}}^{(2)}|_{\overrightarrow{\text{curl}}\widehat{\mathcal{H}}^1(I^2)}) \subset \overrightarrow{\text{curl}}H^1(I^2)$ and

$$\text{Im}(E_{\{1,2,3\}}^{(3)}|_{\widehat{\mathcal{H}}(I^3)}) \subset \mathcal{H}(I^3), \quad \text{Im}(E_{\{1,2,3\}}^{(3)}|_{\text{curl}(\widehat{\mathcal{H}}_1^1(I^3) \times \widehat{\mathcal{H}}_2^1(I^3) \times \widehat{\mathcal{H}}_3^1(I^3))}) \subset \text{curl}H^1(I^3)^3,$$

we obtain $L^2(I^2)^2 = \text{Im}(E_{\{1,2\}}^{(2)}|_{\widehat{\mathcal{H}}(I^2)}) \oplus^\perp \text{Im}(E_{\{1,2\}}^{(2)}|_{\overrightarrow{\text{curl}}\widehat{\mathcal{H}}^1(I^2)}) \oplus^\perp \text{Im}E_{\{1\}}^{(2)} \oplus^\perp \text{Im}E_{\{2\}}^{(2)}$ and $L^2(I^3)^3 = \text{Im}(E_{\{1,2,3\}}^{(3)}|_{\widehat{\mathcal{H}}(I^3)}) \oplus^\perp \text{Im}(E_{\{1,2,3\}}^{(3)}|_{\text{curl}(\widehat{\mathcal{H}}_1^1(I^3) \times \widehat{\mathcal{H}}_2^1(I^3) \times \widehat{\mathcal{H}}_3^1(I^3))}) \oplus^\perp \text{Im}E_{\{1\}}^{(3)} \oplus^\perp \text{Im}E_{\{2\}}^{(3)} \oplus^\perp \text{Im}E_{\{3\}}^{(3)}$.

In view of Proposition 2.1, we obtain

Theorem 2.1 *In the situation of Assumption 2.1, the collections $\Psi_{\text{curl}} =: \Psi_{\text{curl}}^{(n)}$ ($n = 2, 3$) are Riesz bases for $\mathcal{H}(I^n)$ ($n = 2, 3$), $\widetilde{\Psi}_{\text{comp}}^{(2)} \cup E_{\{1\}}^{(2)}\widetilde{\Psi}^{(1)} \cup E_{\{2\}}^{(2)}\widetilde{\Psi}^{(1)}$ and $\widetilde{\Psi}_{\text{comp}}^{(3)} \cup E_{\{1\}}^{(3)}\widetilde{\Psi}^{(1)} \cup E_{\{2\}}^{(3)}\widetilde{\Psi}^{(1)} \cup E_{\{3\}}^{(3)}\widetilde{\Psi}^{(1)}$ are Riesz bases for $\overrightarrow{\text{curl}}H^1(I^2)$ and $\text{curl}H^1(I^3)^3$, respectively.*

Note 2.1 In fact, $\widehat{\mathcal{H}}(I^n) = \mathcal{H}(I^n)$ for $n = 2, 3$. $\widetilde{\Psi}^{(1)} =: \widetilde{\Psi}^- = \{\widetilde{\psi}_\lambda^- : \lambda \in \nabla\}$, which is defined in Section 4.

3 Characterization of $H^m(I^n)$

This part will show that the curl-free wavelets constructed above can be used to characterize Sobolev spaces. For $n = 2, 3$ and $m \in \mathbb{N}$, define the following Sobolev spaces:

$$\begin{aligned} \vec{H}_0^m(I^n) &=: \{ \vec{u} \in H^m(I^n)^n : \vec{u} \times \vec{n} = 0 \text{ or } \vec{0} \text{ on } \Gamma \}, \\ \vec{V}(I^n) &=: \vec{H}_0^m(I^n) \cap \mathcal{H}(I^n). \end{aligned}$$

The following result will be verified in Section 4:

Assumption 3.1 The collection $\Psi^{(n)}$ from Assumption 2.1 can be constructed so that, normalized in $H^m(I^n)^n$, it is a Riesz basis for

$$\widehat{\vec{H}_0^m(I^n)} =: \vec{H}_0^m(I^n) \cap \widehat{L^2(I^n)^n}.$$

Based on this assumption, we obtain:

Theorem 3.1 *In the situation of Assumptions 2.1 and 3.1, the collection $\Psi_{curl} =: \Psi_{curl}^{(n)}$, normalized in $H^m(I^n)^n$, is a Riesz basis for $\vec{V}(I^n)$.*

Proof Since $\mathcal{H}(I^n) = \widehat{\mathcal{H}(I^n)} = \mathcal{H}(I^n) \cap \widehat{L^2(I^n)^n}$, then for any $\vec{u} \in \vec{V}(I^n)$, we know $\vec{u} \in \widehat{\vec{H}_0^m(I^n)}$ and by Assumption 3.1,

$$\vec{u} = \langle \vec{u}, \tilde{\Psi}^{(n)} \rangle_{L^2(I^n)^n} \Psi^{(n)} \quad \text{in } H^m(I^n)^n$$

with $\|\vec{u}\|_{H^m(I^n)^n}^2 \simeq \sum_{\tilde{\psi} \in \tilde{\Psi}^{(n)}} |\langle \vec{u}, \tilde{\psi} \rangle_{L^2(I^n)^n}|^2 \cdot \|\psi_{\tilde{\psi}}\|_{H^m(I^n)^n}^2$, where $\psi_{\tilde{\psi}} \in \Psi^{(n)}$ denotes the primal wavelets corresponding to $\tilde{\psi}$. Furthermore, since $\vec{u} \in \mathcal{H}(I^n)$, then

$$\vec{u} = \langle \vec{u}, \tilde{\Psi}_{curl}^{(n)} \rangle_{L^2(I^n)^n} \Psi_{curl}^{(n)} \quad \text{in } H^m(I^n)^n$$

with $\|\vec{u}\|_{H^m(I^n)^n}^2 \simeq \sum_{\tilde{\psi} \in \tilde{\Psi}_{curl}^{(n)}} |\langle \vec{u}, \tilde{\psi} \rangle_{L^2(I^n)^n}|^2 \cdot \|\psi_{\tilde{\psi}}\|_{H^m(I^n)^n}^2$. □

4 Construction of wavelets

In this section, we will give the construction of wavelets satisfying Assumptions 2.1 and 3.1.

Lemma 4.1 ([5, Corollary 3.3]) *Suppose that the collections $\Psi = \{\psi_\lambda : \lambda \in \nabla\}$ and $\tilde{\Psi} = \{\tilde{\psi}_\lambda : \lambda \in \nabla\}$ are bi-orthogonal in $L^{2,0}(I)$. In addition, for some $m < \gamma < d \in \mathbb{N}$, $2 < \tilde{\gamma} < \tilde{d} \in \mathbb{N}$,*

$$\inf_{v \in \text{span}\{\psi_\lambda : |\lambda| \leq \ell\}} \|u - v\|_{L^2(I)} \leq 2^{-\ell d} \|u\|_{H^d(I)} \quad (u \in \widehat{H^d(I)}),$$

$$\inf_{v \in \text{span}\{\tilde{\psi}_\lambda : |\lambda| \leq \ell\}} \|u - v\|_{L^2(I)} \leq 2^{-\ell \tilde{d}} \|u\|_{\widehat{H^{\tilde{d}}(I)}} \quad (u \in \widehat{H^{\tilde{d}}(I)}),$$

$$\text{for } s < \gamma, \quad \|\cdot\|_{H^s(I)} \leq 2^{\ell s} \|\cdot\|_{L^2(I)} \quad \text{on } \text{span}\{\psi_\lambda : |\lambda| \leq \ell\},$$

$$\text{for } s < \tilde{\gamma}, \quad \|\cdot\|_{H^s(I)} \leq 2^{\ell s} \|\cdot\|_{L^2(I)} \quad \text{on } \text{span}\{\tilde{\psi}_\lambda : |\lambda| \leq \ell\}.$$

Define the collections $\Psi^+ = \{\psi_\lambda^+ : \lambda \in \nabla\}$ and $\tilde{\Psi}^- = \{\tilde{\psi}_\lambda^- : \lambda \in \nabla\}$ by

$$\psi_\lambda^+(x) = 2^{|\lambda|} \int_0^x \psi_\lambda(y) dy \quad \text{and} \quad \tilde{\psi}_\lambda^-(x) = -2^{-|\lambda|} \tilde{\psi}_\lambda.$$

Then it holds that

$$\begin{aligned} \{2^{-|\lambda|s} \psi_\lambda : \lambda \in \nabla\} & \text{ is a Riesz basis for } \widehat{H}^s(I), \quad s \in [0, \gamma), \\ \{2^{-|\lambda|s} \widetilde{\psi}_\lambda : \lambda \in \nabla\} & \text{ is a Riesz basis for } \widehat{H}^s(I), \quad s \in [0, \widetilde{\gamma}), \\ \{2^{-|\lambda|s} \psi_\lambda^+ : \lambda \in \nabla\} & \text{ is a Riesz basis for } \mathcal{H}_0^s(I), \quad s \in [0, \gamma + 1), \\ \{2^{-|\lambda|s} \widetilde{\psi}_\lambda^- : \lambda \in \nabla\} & \text{ is a Riesz basis for } H^s(I), \quad s \in [0, \widetilde{\gamma} - 1), \end{aligned}$$

where

$$\mathcal{H}_0^s(I) =: \begin{cases} [L^2(I), H_0^1(I)]_{s,2}, & s \in [0, 1]; \\ H^s(I) \cap H_0^1(I), & s \geq 1. \end{cases}$$

Moreover, Ψ^+ and $\widetilde{\Psi}^-$ are bi-orthogonal.

Note 4.1 It has been pointed out in [5] that such wavelet bases can be obtained by taking standard bi-orthogonal wavelet bases for $L^2(I)$ that satisfy the corresponding Jackson and Bernstein assumptions of d, \widetilde{d}, γ and $\widetilde{\gamma}$ with $\widehat{H}^d(I)$ and $\widehat{H}^{\widetilde{d}}(I)$ reading as $H^d(I)$ and $H^{\widetilde{d}}(I)$ (see [7]), and then removing those scaling functions without a vanishing moment.

The following result can be proved by the same method as Corollary 3.7 of [5].

Corollary 4.1 For $0 \leq s < \gamma$ and $0 \leq \widetilde{s} < \widetilde{\gamma} - 1$, the sets

$$\begin{aligned} & \left\{ \left(\sum_{i=1}^n 4^{|\lambda_i|} \right)^{-\frac{\widetilde{s}}{2}} \psi_{\lambda_1}^+ \otimes \cdots \otimes \psi_{\lambda_k} \otimes \cdots \otimes \psi_{\lambda_n}^+ : \lambda = (\lambda_1, \dots, \lambda_n) \in \widetilde{\nabla} = (\nabla)^n \right\}, \\ & \left\{ \left(\sum_{i=1}^n 4^{|\lambda_i|} \right)^{-\frac{\widetilde{s}}{2}} \widetilde{\psi}_{\lambda_1}^- \otimes \cdots \otimes \widetilde{\psi}_{\lambda_k} \otimes \cdots \otimes \widetilde{\psi}_{\lambda_n}^- : \lambda = (\lambda_1, \dots, \lambda_n) \in \widetilde{\nabla} = (\nabla)^n \right\} \end{aligned}$$

are Riesz bases for

$$\begin{aligned} & \downarrow \text{ kth position} \\ & \mathcal{H}_0^s \otimes L^2 \otimes \cdots \otimes L^2 \otimes L^{2,0} \otimes L^2 \otimes \cdots \otimes L^2 \cap \\ & \quad \vdots \\ & \text{kth position } L^2 \otimes \cdots \otimes L^2 \otimes \widehat{H}^s \otimes L^2 \otimes \cdots \otimes L^2 \cap \\ & \quad \vdots \\ & L^2 \otimes L^2 \otimes \cdots \otimes L^2 \otimes L^{2,0} \otimes L^2 \otimes \cdots \otimes \mathcal{H}_0^s \end{aligned}$$

and

$$\begin{aligned} & \downarrow \text{ kth position} \\ & H^{\widetilde{s}} \otimes L^2 \otimes \cdots \otimes L^2 \otimes L^{2,0} \otimes L^2 \otimes \cdots \otimes L^2 \cap \\ & \quad \vdots \\ & \text{kth position } L^2 \otimes \cdots \otimes L^2 \otimes \widehat{H}^{\widetilde{s}} \otimes L^2 \otimes \cdots \otimes L^2 \cap \\ & \quad \vdots \\ & L^2 \otimes L^2 \otimes \cdots \otimes L^2 \otimes L^{2,0} \otimes L^2 \otimes \cdots \otimes H^{\widetilde{s}}, \end{aligned}$$

respectively. For $s = \tilde{s} = 0$, the corresponding collections are bi-orthogonal in $L^2 \otimes \dots \otimes L^2 \otimes L^{2,0} \otimes L^2 \otimes \dots \otimes L^2$.

For $\lambda \in \vec{\nabla}$, we define the vector-valued wavelets

$$\underline{\psi}_{\lambda,k}^{(n)} =: \psi_{\lambda_1}^+ \otimes \dots \otimes \psi_{\lambda_k} \otimes \dots \otimes \psi_{\lambda_n}^+ \vec{e}_k, \quad \tilde{\underline{\psi}}_{\lambda,k}^{(n)} =: \tilde{\psi}_{\lambda_1}^- \otimes \dots \otimes \tilde{\psi}_{\lambda_k} \otimes \dots \otimes \tilde{\psi}_{\lambda_n}^- \vec{e}_k.$$

From Corollary 4.1 and the definition of $\widehat{H_0^1(I^n)}$, we obtain

Proposition 4.1 For $0 \leq s < \gamma$ and $0 \leq \tilde{s} < \tilde{\gamma} - 1$, the sets

$$\left\{ \left(\sum_{i=1}^n 4^{|\lambda_i|} \right)^{-\frac{s}{2}} \underline{\psi}_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla} \right\} \quad \text{and}$$

$$\left\{ \left(\sum_{i=1}^n 4^{|\lambda_i|} \right)^{-\frac{\tilde{s}}{2}} \tilde{\underline{\psi}}_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla} \right\}$$

are Riesz bases for the vector spaces

$$\begin{cases} [\widehat{L^2(I^n)^n}, \widehat{H_0^1(I^n)}]_{s,2}, & s \in [0, 1]; \\ \widehat{H_0^1(I^n)} \cap H^s(I^n)^n, & s \geq 1 \end{cases}$$

and $\widehat{L^2(I^n)^n} \cap H^{\tilde{s}}(I^n)^n$, respectively. For $s = \tilde{s} = 0$, the collections are bi-orthogonal Riesz bases for $\widehat{L^2(I^n)^n}$.

Now, we are in the position to apply the basis transform. Let A^λ be an orthogonal matrix with its 1st row given by

$$A_1^\lambda =: \frac{1}{\left(\sum_{i=1}^n 4^{|\lambda_i|}\right)^{\frac{1}{2}}} (2^{|\lambda_1|}, \dots, 2^{|\lambda_n|}) = (\alpha_1, \dots, \alpha_n) =: \alpha^T.$$

Such an example is known as the Householder transform

$$A^\lambda = I - \frac{2(\alpha - \vec{e}_1)(\alpha - \vec{e}_1)^T}{(\alpha - \vec{e}_1)^T(\alpha - \vec{e}_1)},$$

which is

$$A^\lambda = \begin{pmatrix} \alpha_1 & \alpha_2 \\ \alpha_2 & -\alpha_1 \end{pmatrix} \quad \text{and} \quad A^\lambda = \begin{pmatrix} \alpha_1 & \alpha_2 & \alpha_3 \\ \alpha_2 & 1 - \frac{\alpha_2^2}{1-\alpha_1} & -\frac{\alpha_2\alpha_3}{1-\alpha_1} \\ \alpha_3 & -\frac{\alpha_2\alpha_3}{1-\alpha_1} & 1 - \frac{\alpha_3^2}{1-\alpha_1} \end{pmatrix}$$

in the case $n = 2$ and $n = 3$. Defining

$$\begin{pmatrix} \underline{\psi}_{\lambda,1}^{(n)} \\ \vdots \\ \underline{\psi}_{\lambda,n}^{(n)} \end{pmatrix} =: A^\lambda \begin{pmatrix} \underline{\psi}_{\lambda,1}^{(n)} \\ \vdots \\ \underline{\psi}_{\lambda,n}^{(n)} \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} \tilde{\underline{\psi}}_{\lambda,1}^{(n)} \\ \vdots \\ \tilde{\underline{\psi}}_{\lambda,n}^{(n)} \end{pmatrix} =: A^\lambda \begin{pmatrix} \tilde{\underline{\psi}}_{\lambda,1}^{(n)} \\ \vdots \\ \tilde{\underline{\psi}}_{\lambda,n}^{(n)} \end{pmatrix}.$$

Set

$$\Psi^{(n)} =: \{\psi_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla}\} \quad \text{and} \quad \tilde{\Psi}^{(n)} =: \{\tilde{\psi}_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla}\}.$$

Applying the property of an orthogonal transform, we infer the following result.

Proposition 4.2 For $0 \leq s < \gamma$ and $0 \leq \tilde{s} < \tilde{\gamma} - 1$, the sets

$$\left\{ \left(\sum_{i=1}^n 4^{|\lambda_i|} \right)^{-\frac{s}{2}} \psi_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla} \right\} \quad \text{and}$$

$$\left\{ \left(\sum_{i=1}^n 4^{|\lambda_i|} \right)^{-\frac{\tilde{s}}{2}} \tilde{\psi}_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla} \right\}$$

are Riesz bases for the vector spaces

$$\begin{cases} [L^2(I^n)^n, \widehat{H}_0^1(I^n)]_{s,2}, & s \in [0, 1]; \\ \widehat{H}_0^1(I^n) \cap H^s(I^n)^n, & s \geq 1 \end{cases}$$

and $L^2(I^n)^n \cap H^{\tilde{s}}(I^n)^n$, respectively. In particular, for $s = \tilde{s} = 0$, the collections $\Psi^{(n)}$ and $\tilde{\Psi}^{(n)}$ are bi-orthogonal Riesz bases for $L^2(I^n)^n$.

In the following, we are mainly concerned with the cases $n = 2$ and $n = 3$ because of the complicated form of curl operators in $n > 3$.

Theorem 4.1 Let $\Psi_{curl}^{(n)} =: \{\psi_{\lambda,1}^{(n)} : \lambda \in \vec{\nabla}\}$ and $\tilde{\Psi}_{comp}^{(n)} =: \{\tilde{\psi}_{\lambda,k}^{(n)} : 2 \leq k \leq n, \lambda \in \vec{\nabla}\}$. Then

- (i) $\Psi_{curl}^{(n)} \subset \widehat{\mathcal{H}}(I^n)$ ($n = 2, 3$), $\tilde{\Psi}_{comp}^{(2)} \subset \overrightarrow{curl} \widehat{H}^1(I^2)$ or $\tilde{\Psi}_{comp}^{(3)} \subset \overrightarrow{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3))$.
- (ii) $\{(\sum_{i=1}^n 4^{|\lambda_i|})^{-\frac{m}{2}} \psi_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \vec{\nabla}\}$ is a Riesz basis for the vector valued space $\widehat{H}_0^m(I^n) =: \widehat{H}_0^m(I^n) \cap \widehat{L}^2(I^n)^n$.

Proof (i) It is easy to see that $\underline{\psi}_{\lambda,k}^{(n)} \in H_0(curl; I^n)$ for $1 \leq k \leq n$, then

$$\psi_{\lambda,1}^{(n)} = \alpha_1 \underline{\psi}_{\lambda,1}^{(n)} + \alpha_2 \underline{\psi}_{\lambda,2}^{(n)} + \dots + \alpha_n \underline{\psi}_{\lambda,n}^{(n)} \in H_0(curl; I^n).$$

Furthermore, $\overrightarrow{curl} \psi_{\lambda,1}^{(2)} = 0$ and $curl \psi_{\lambda,1}^{(3)} = \vec{0}$. Therefore, $\Psi_{curl}^{(n)} \subset \widehat{\mathcal{H}}(I^n)$ ($n = 2, 3$). In addition,

$$\begin{aligned} \tilde{\psi}_{\lambda,2}^{(2)} &= \alpha_2 \underline{\tilde{\psi}}_{\lambda,1}^{(2)} - \alpha_1 \underline{\tilde{\psi}}_{\lambda,2}^{(2)} = \begin{pmatrix} \alpha_2 \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2}^- \\ -\alpha_1 \tilde{\psi}_{\lambda_1}^- \otimes \tilde{\psi}_{\lambda_2} \end{pmatrix} \\ &= -\frac{1}{(4^{|\lambda_1|} + 4^{|\lambda_2|})^{\frac{1}{2}}} \overrightarrow{curl} \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2} \in \overrightarrow{curl} \widehat{H}^1(I^2). \end{aligned}$$

Therefore, we obtain $\tilde{\Psi}_{comp}^{(2)} \subset \overrightarrow{curl} \widehat{H}^1(I^2)$. Finally, suppose that a , b and c are the solutions of

$$\begin{cases} b2^{|\lambda_3|} - c2^{|\lambda_2|} = A_{21}, \\ -a2^{|\lambda_3|} + c2^{|\lambda_1|} = A_{22}, \\ -b2^{|\lambda_1|} + a2^{|\lambda_2|} = A_{23}, \end{cases}$$

whose existence can be guaranteed by the orthogonality of A^λ . Then

$$\begin{aligned} \tilde{\psi}_{\lambda,2}^{(3)} &= A_{21} \tilde{\psi}_{\lambda,1}^{(3)} + A_{22} \tilde{\psi}_{\lambda,2}^{(3)} + A_{23} \tilde{\psi}_{\lambda,3}^{(3)} = \begin{pmatrix} A_{21} \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2}^- \otimes \tilde{\psi}_{\lambda_3}^- \\ A_{22} \tilde{\psi}_{\lambda_1}^- \otimes \tilde{\psi}_{\lambda_2} \otimes \tilde{\psi}_{\lambda_3}^- \\ A_{23} \tilde{\psi}_{\lambda_1}^- \otimes \tilde{\psi}_{\lambda_2}^- \otimes \tilde{\psi}_{\lambda_3} \end{pmatrix} \\ &= \text{curl} \begin{pmatrix} a \tilde{\psi}_{\lambda_1}^- \otimes \tilde{\psi}_{\lambda_2} \otimes \tilde{\psi}_{\lambda_3} \\ b \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2}^- \otimes \tilde{\psi}_{\lambda_3} \\ c \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2} \otimes \tilde{\psi}_{\lambda_3}^- \end{pmatrix} \in \text{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3)). \end{aligned}$$

Similarly, if a , b and c are the solutions of the equation

$$\begin{cases} b2^{|\lambda_3|} - c2^{|\lambda_2|} = A_{31}, \\ -a2^{|\lambda_3|} + c2^{|\lambda_1|} = A_{32}, \\ -b2^{|\lambda_1|} + a2^{|\lambda_2|} = A_{33}, \end{cases}$$

then we can also obtain

$$\tilde{\psi}_{\lambda,3}^{(3)} = \text{curl} \begin{pmatrix} a \tilde{\psi}_{\lambda_1}^- \otimes \tilde{\psi}_{\lambda_2} \otimes \tilde{\psi}_{\lambda_3} \\ b \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2}^- \otimes \tilde{\psi}_{\lambda_3} \\ c \tilde{\psi}_{\lambda_1} \otimes \tilde{\psi}_{\lambda_2} \otimes \tilde{\psi}_{\lambda_3}^- \end{pmatrix} \in \text{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3)).$$

Therefore, $\tilde{\Psi}_{comp}^{(3)} \subset \text{curl}(\widehat{H}_1^1(I^3) \times \widehat{H}_2^1(I^3) \times \widehat{H}_3^1(I^3))$.

(ii) Since $\gamma > m$, taking $s = m$ in Proposition 4.2, we know the set $\{(\sum_{i=1}^n 4^{|\lambda_i|})^{-\frac{m}{2}} \psi_{\lambda,k}^{(n)} : 1 \leq k \leq n, \lambda \in \bar{\nabla}\}$ is a Riesz basis for $\widehat{H}_0^1(I^n) \cap H^m(I^n)$. Furthermore, it is easy to verify

$$\widehat{H}_0^1(I^n) \cap H^m(I^n) = \widehat{H}_0^m(I^n) \cap L^2(I^n) = \widehat{H}_0^m(I^n). \quad \square$$

Competing interests

The author declares that they have no competing interests.

Acknowledgements

The project is supported by the National Natural Science Foundation of China (No. 11201094, 11161014), the 863 Project of China (No. 2012AA011005), the project of Guangxi Innovative Team (No. 2012jjGAG0001), the fund of Education Department of Guangxi Province (No. 201012M9094, 201102ZD015, 201106LX172).

Received: 3 March 2012 Accepted: 6 September 2012 Published: 19 September 2012

References

1. Deriaz, E, Perrier, V: Towards a divergence-free wavelet method for the simulation of 2D/3D turbulent flows. *J. Turbul.* **7**(3), 1-37 (2006)
2. Deriaz, E, Perrier, V: Orthogonal Helmholtz decomposition in arbitrary dimension using divergence-free and curl-free wavelets. *Appl. Comput. Harmon. Anal.* **26**(2), 249-269 (2009)

3. Jiang, YC, Liu, YM: Interpolatory curl-free wavelets and applications. *Int. J. Wavelets Multiresolut. Inf. Process.* **5**, 843-858 (2007)
4. Urban, K: Wavelet bases in $H(\text{div})$ and $H(\text{curl})$. *Math. Comput.* **70**(234), 739-766 (2001)
5. Stevenson, R: Divergence-free wavelet bases on the hypercube: free-slip boundary conditions, and applications for solving the instationary Stokes equations. *Math. Comput.* **80**, 1499-1523 (2011)
6. Stevenson, R: Divergence-free wavelet bases on the hypercube. *Appl. Comput. Harmon. Anal.* **30**, 1-19 (2011)
7. Dahmen, W, Kunoth, A, Urban, K: Biorthogonal spline-wavelets on the interval-stability and moment conditions. *Appl. Comput. Harmon. Anal.* **6**, 132-196 (1999)

doi:10.1186/1029-242X-2012-205

Cite this article as: Jiang: Anisotropic curl-free wavelets with boundary conditions. *Journal of Inequalities and Applications* 2012 **2012**:205.

Submit your manuscript to a SpringerOpen[®] journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Immediate publication on acceptance
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at ► springeropen.com
