

Research Article

A Postverification Method for Solving Forced Duffing Oscillator Problems without Prescribed Periods

Hong-Yen Lin,^{1,2} Chien-Chang Yen,³ Kuo-Ching Jen,⁴ and Kang C. Jea^{2,3}

¹ Department of Computer Science and Information Engineering, Fu Jen Catholic University, New Taipei City 24205, Taiwan

² Graduate Institute of Applied Science and Engineering, Fu Jen Catholic University, New Taipei City 24205, Taiwan

³ Department of Mathematics, Fu Jen Catholic University, New Taipei City 24205, Taiwan

⁴ General Education Center, St. John's University, New Taipei City 25135, Taiwan

Correspondence should be addressed to Hong-Yen Lin; hylin@mail.cc.fju.edu.tw

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This paper proposes a postverification method (PVM) for solving forced Duffing oscillator problems without prescribed periods. Comprising a postverification procedure and small random perturbation, the proposed PVM improves the sensitivity of the convergence of Newton's iteration. Numerical simulations revealed that the PVM is more accurate and robust than Kubiček's approach. We applied the PVM to previous research on bifurcation problems.

1. Introduction

A Duffing oscillator equation [1–10], which is an ordinary differential nonhomogeneous equation of the second order, can be expressed as

$$u'' + au' + bu + cu^3 = \alpha \cos \omega t, \quad (1)$$

where the dampening coefficient a , coefficients b and c in the restoring force, amplitude α , and frequency ω in the external oscillator force are constant. In [11], the constant coefficients a , b , and c were generalized to variable coefficients, and the existence of positive solutions was shown. However, researchers are interested in periodic and bifurcation solutions with the prescribed period $T = 2\pi/\omega$ [12–15]. The nonlinear equation shown in (1) can be rewritten as a nonlinear dynamic system as follows:

$$\frac{dx_i}{dt} = f_i(x_1, x_2, \dots, x_n, t), \quad i = 1, 2, \dots, n, \quad t \geq 0. \quad (2)$$

There are several applications for nonlinear dynamic systems with ambiguous or unknown force oscillators [16–18]. However, further research is necessary to solve (1) where the frequency of the external oscillator force is unknown.

Specifically, it is necessary to identify a state where $I_0 = (x_1(0), x_2(0), \dots, x_n(0), T)$, which is the initial condition and period T for (2).

A continuation method involving Newton's iterative method has frequently been applied to solve the branch of bifurcation points. The approach proposed by Kubiček and Marek [19] for solving the aforementioned problem does not guarantee that the obtained numerical solution is always periodic, potentially because fixing the specific direction or approach is sensitive to the convergence of Newton's iteration. First, we modified Kubiček's approach by applying an alternative method to eliminate the drawback of fixing a specific direction; hence, we term the proposed approach the modified Kubiček (MK) approach. A postverification process, called the postverification method (PVM), is then introduced into the MK approach to ensure periodicity. The PVM, which is based on the stability of solutions, has the following three characteristics: (i) an alternative direction for the initial state is fixed; (ii) a randomly small perturbation approach is used to confirm the stability of the solution; and (iii) a postverification test of periodicity is applied.

Forced Duffing oscillators, which lack nonlinear restoring force, were simulated using Kubiček's approach, the MK approach, and proposed PVM to evaluate the performances

of these methods. The numerical results indicated that the proposed approaches can be used to obtain periodic numerical solutions with greater stability. Furthermore, the PVM was applied to autonomous equations to obtain the periodic solutions.

The remainder of this paper is organized as follows. Section 2 introduces the proposed method, and the numerical results are reported in Section 3. Section 4 describes the application of the obtained results to a bifurcation calculation. Finally, Section 5 provides our conclusion.

2. The Postverification Method

Based on the assumption that the solution is stable, we developed the MK approach [19] and PVM. The proposed method is feasible because of the stable criterion and postverification process. The proposed method involves starting from the system shown in (2) by using a prescribed period T . Changing the variable $t = Ts$, where $0 \leq s \leq 1$, transforms (2) into

$$\frac{dx_i}{dt} = \frac{dx_i}{Ts} = f_i(x_1, x_2, \dots, x_n, Ts), \quad i = 1, 2, \dots, n \quad (3)$$

or

$$\frac{dx_i}{ds} = Tf_i(x_1, x_2, \dots, x_n, Ts), \quad i = 1, 2, \dots, n. \quad (4)$$

Denote $y_i(s) = x_i(Ts)$, where $0 \leq s \leq 1$. Thus, the solution x_i for period T is equivalent to y_i of the unit period, and (4) is transformed into

$$\frac{dy_i}{ds} = Tf_i(y_1, y_2, \dots, y_n, Ts), \quad i = 1, 2, \dots, n. \quad (5)$$

An initial value $y_i(0) = \eta_i$, where $i = 1, 2, \dots, n$ such that the solutions that are both periodic and have bifurcation points, can be identified because a periodic solution of (5) exists. Moreover, the value $y_i(1)$ depends on $\eta_1, \eta_2, \dots, \eta_n$, and let

$$y_i(1) \equiv \phi_i(\eta_1, \eta_2, \dots, \eta_n) \quad (6)$$

for some functions ϕ_i , $i = 1, \dots, n$. According to the assumption that solutions are periodic, $y_i(0) = y_i(1)$; therefore, (6) becomes

$$\eta_i - \phi_i(\eta_1, \eta_2, \dots, \eta_n) = 0, \quad i = 1, 2, \dots, n. \quad (7)$$

Define

$$\begin{aligned} F_i &= \eta_i - \phi_i(\eta_1, \eta_2, \dots, \eta_n), \quad i = 1, 2, \dots, n, \\ \eta &= (\eta_1, \eta_2, \dots, \eta_n)^t, \\ y &= (y_1, y_2, \dots, y_n)^t, \\ \phi &= (\phi_1, \phi_2, \dots, \phi_n)^t, \\ F &= (F_1, F_2, \dots, F_n)^t, \end{aligned} \quad (8)$$

where $(\cdot)^t$ denotes the transpose of a vector. Expressed in vector form, (7) can be rewritten as

$$F(\eta) = 0, \quad (9)$$

and (9) can be solved using Newton's iterative method

$$\eta^{(l+1)} = \eta^{(l)} - \left[\frac{\partial F(\eta^{(l)})}{\partial \eta} \right]^{-1} F(\eta^{(l)}), \quad l = 0, 1, 2, \dots \quad (10)$$

by applying an initial estimate of $\eta^{(0)}$. The iterative process is executable if the Jacobian matrix $F_\eta(\eta^{(l)}) = \partial F(\eta^{(l)})/\partial \eta$ is nonsingular, vector $F(\eta^{(l)})$ can be evaluated, and initial estimate $\eta^{(0)}$ is close to the solution. According to (8) and (6), the Jacobian matrix is

$$F_\eta = I_{n \times n} - \frac{\partial \phi}{\partial \eta} \quad \text{with} \quad \frac{\partial \phi_i}{\partial \eta_j} = \frac{\partial y_i(1)}{\partial \eta_j}, \quad (11)$$

where $I_{n \times n}$ is an identity matrix of dimension $n \times n$. Let

$$P_{ij}(s) = \frac{\partial y_i(s)}{\partial \eta_j}, \quad 0 \leq s \leq 1, \quad i, j = 1, 2, \dots, n. \quad (12)$$

Subsequently,

$$\begin{aligned} \frac{d}{ds} P_{ij}(s) &= \frac{d}{ds} \frac{\partial y_i(s)}{\partial \eta_j} = \frac{\partial}{\partial \eta_j} \left(\frac{dy_i(s)}{ds} \right) \\ &= \frac{\partial}{\partial \eta_j} (Tf_i(y_1, y_2, \dots, y_n, Ts)) = T \sum_{m=1}^n \frac{\partial f_i}{\partial y_m} \frac{\partial y_m}{\partial \eta_j}, \end{aligned} \quad (13)$$

and the preceding $n \times n$ equations can be written in matrix form as

$$\frac{dP}{ds} = T \frac{\partial f}{\partial y} P. \quad (14)$$

The matrix $\partial \phi/\partial \eta = P(1)$ is calculated by solving (14) with the initial conditions $P(0) = I_{n \times n}$. The vector $\phi = y(1)$ can be evaluated using Newton's iterative process, the Runge-Kutta-Fehlberg (RKF) method, and (5) and (6). We introduce both Kubíček's method and the proposed MK and PVM to solve (2) without the prescribed period T . Accordingly, (8) is modified as

$$F_i(\eta_1, \eta_2, \dots, \eta_n, T) = \eta_i - \phi_i(\eta_1, \eta_2, \dots, \eta_n, T), \quad i = 1, 2, \dots, n, \quad (15)$$

which have $n + 1$ unknowns, $\eta_1, \eta_2, \dots, \eta_n$, and T . But, there are only n equations $F_i = 0$, $i = 1, 2, \dots, n$. Kubíček proposed fixing a direction (e.g., k). The derivation is similar for the prescribed periods and is required to calculate the term

$$\frac{\partial F_i}{\partial T} = -\frac{\partial \phi_i}{\partial T} = -\frac{\partial y_i(1)}{\partial T}. \quad (16)$$

Let $r_i(s) = \partial y_i(s)/\partial T$ and

$$\begin{aligned} \frac{d}{ds} r_i(s) &= \frac{d}{ds} \frac{\partial y_i(s)}{\partial T} \\ &= \frac{d}{ds} \frac{\partial y_i}{\partial T} = \frac{\partial}{\partial T} \frac{dy_i(s)}{ds} \\ &= \frac{\partial}{\partial T} (Tf_i(y_1, y_2, \dots, y_n, Ts)) = f_i(y_1, y_2, \dots, y_n, Ts) \\ &\quad + T \left(\frac{\partial f_i(y_1, y_2, \dots, y_n, Ts)}{\partial T} + \sum_{m=1}^n \frac{\partial f_i}{\partial y_m} \frac{\partial y_m}{\partial T} \right). \end{aligned} \tag{17}$$

This can be abbreviated as

$$r' = f + T \frac{\partial f}{\partial T} + T \frac{\partial f}{\partial y} r. \tag{18}$$

With $r_i(0) = \partial y_i(0)/\partial T = \partial \eta_i/\partial T = 0$, for $i = 1, 2, \dots, n$. The RKF method is used again to obtain $r_i(1) = \partial F_i/\partial T = \partial \phi_i/\partial T$. Therefore,

$$\frac{\partial \phi_i}{\partial \eta_j} = \begin{cases} P_{ij}(1), & i \leq j \leq k-1, \\ P_{ij+1}(1), & k \leq j \leq n-1, \\ r_i(1), & j = n. \end{cases} \tag{19}$$

Denote $U = (y_1, \dots, y_{k-1}, y_{k+1}, \dots, y_n, T)$. When $F_i(U) = F_i(y_1, y_2, \dots, y_n, T) = \eta_i - \phi_i(\eta_1, \eta_2, \dots, \eta_n) = \eta_i - y_i(1)$,

$$\frac{dy_i}{ds} = Tf_i(y_1, y_2, \dots, y_n, Ts), \tag{20}$$

and $y_i(0) = \eta_i$ is initial estimated appropriately, $y(1)$ can be obtained using the RKF. The expression

$$U^{(l+1)} = U^{(l)} - (F_U(U^{(l)}))^{-1} (y(1) - \eta^{(l)}) \tag{21}$$

is repeated until $\|U^{(l+1)} - U^{(l)}\| < \varepsilon$.

There exist two numbers t_1 and t_2 in $[0, T]$ such that $y_k(t)$ is monotonic at each interval $[0, t_1]$ and $[t_2, T]$, where the intervals $[0, t_1]$ and $[t_2, T]$ are the largest. Define y_k^- and y_k^+ as the minimum and maximum of the union $[0, t_1] \cup [t_2, T]$. The integrals of

$$y' = f_i(y_1, y_2, \dots, y_n, T), \quad \text{for } i = 1, 2, \dots, n. \tag{22}$$

with the initial value $y(0) = (\eta_1, \eta_2, \dots, \eta_n)$ are determined by applying the RKF method from $t = 0$ to T . Furthermore,

$$\begin{aligned} \eta_k^- &= \frac{1}{2} ((1 + \omega) y_k^- + (1 - \omega) y_k^+), \\ \eta_k^+ &= \frac{1}{2} ((1 - \omega) y_k^- + (1 + \omega) y_k^+), \end{aligned} \tag{23}$$

where $0 < \omega < 1$, and ω can be set as $1/2$, which is acceptable if $\eta_k^- < \eta_k < \eta_k^+$; otherwise, η_k should be modified as follows.

The interval $[0, 1]$ is divided into subintervals, where y_k is monotonic and the minimum and maximum occur at the boundaries. The turning positions y_k^- and y_k^+ are the

minimum and maximum on the monotonic part of $y_k(t)$ with the maximal difference (e.g., $[t_i, t_{i+1}]$); that is,

$$y_k^- = \min_{t \in [t_i, t_{i+1}]} y_k(t), \quad y_k^+ = \max_{t \in [t_i, t_{i+1}]} y_k(t). \tag{24}$$

Therefore,

$$\eta_k = \frac{1}{2} (y_k^{++} - y_k^{--}). \tag{25}$$

Denote $\eta_{n+1} = T$ and calculate

$$F_U = \begin{bmatrix} \partial F_i \\ \partial \eta_j \end{bmatrix} = \begin{bmatrix} \partial \phi_i \\ \partial \eta_j \end{bmatrix} - \begin{bmatrix} \partial \eta_i \\ \partial \eta_j \end{bmatrix}, \tag{26}$$

where $i = 1, 2, \dots, n$, and $j = 1, 2, \dots, k-1, k+1, \dots, n+1$, yielding

$$\frac{\partial \eta_i}{\partial \eta_j} = \begin{cases} \delta_{ij}, & 1 \leq i, j \leq k-1, \\ 0, & i = k, 1 \leq j \leq n-1, \\ 1, & j = i-1, k+1 \leq i \leq n-1, \\ 0, & j = n, 1 \leq i \leq n. \end{cases} \tag{27}$$

The aforementioned argument can be summarized as the following algorithm, which is referred to as Kubíček's approach [12, 19].

2.1. Kubíček's Approach. The approach proposed by Kubíček and Marek [19] for solving (q) is to fix the specific direction during the whole computation. The corresponding algorithm is given in Algorithm 1.

2.2. The Modified Kubíček Approach. When using Kubíček's approach, a periodic solution is not always obtainable, even when the criterion of convergence is satisfied, potentially because the k th direction remains fixed throughout the process. To investigate this possibility, we propose an iterative approach that involves alternating the fixed k th direction from $k = 1$ to n . However, for the fixed k th direction, Newton's method is applied to obtain a convergent solution $(\eta_1, \dots, \eta_{k-1}, \eta_{k+1}, \dots, \eta_n, T)$. The same procedure is applied for each direction until the criterion for convergence is satisfied.

The algorithm for the modified Kubíček(MK) approach is given in Algorithm 2.

2.3. The Postverification Method. In simulations, the computational cost of evaluating y_k^- and y_k^+ is high. We eliminated the integration and replaced the f value of the k th direction by introducing a small random perturbation to η_k . In addition, we investigated the sensitivity of the convergence of Newton's iteration. The proposed PVM was employed to overcome the sensitivity and, thus, confirm the following three conditions: (i) the values at $0, T$, and $2T$ should be identical; (ii) the maximal (minimal) values at $[0, T]$ and $[T, 2T]$ should be equal (similar); and (iii) the time interval between the two maximal (minimal) values should be equal (similar) to T . The algorithm for the PVM is detailed as in Algorithm 3.

Let ε be given and fix k
 Set an initial estimate $U^{(0)} = (\eta_1, \dots, \eta_{k-1}, \eta_{k+1}, \dots, \eta_n, T)^t$.
 Repeat
 $U^{(l+1)} = U^{(l)} - (F_U(U^{(l)}))^{-1} F(U^{(l)})$, where $F = (F_1, \dots, F_n)^t$.
 Modify η_k by applying (25)
 until $\|U^{(l+1)} - U^{(l)}\| < \varepsilon$.

ALGORITHM 1: Kubíček method [19].

Let ε be a given value.
 Set the initial estimate $(\eta_1, \dots, \eta_n, T)$.
 Repeat
 For $k = 1$ to n // let η_k be fixed
 Let $U = (\eta_1, \dots, \eta_{k-1}, \eta_{k+1}, \dots, \eta_n, T)^t$.
 Repeat
 $\delta \leftarrow [F_U(U)]^{-1} F(U)$.
 $U \leftarrow U - \delta$.
 until $\|\delta\| < \varepsilon$.
 Modify η_k by applying (25).
 Next k
 until $(\eta_1, \dots, \eta_n, T)$ becomes constant.

ALGORITHM 2: The modified Kubíček method.

3. Numerical Results

The performance of the proposed methods was evaluated by simulating the Duffing oscillator equation with a nonlinear restoring force [20] and autonomous equations representing the case of the nonlinear restoring force of an unknown period.

Example 1. This example illustrates that the results obtained by using Kubíček's approach, the MK approach, and proposed PVM to solve (1), where $a = b = 0.1$, $c = 1$, $\alpha = 5.0259$, and $\omega = 1$ [20], are identical. For the initial estimate $(\eta_1^{(0)}, \eta_2^{(0)}, T^{(0)}) = (3.0, 1.2, 5.6)$, where $k = 2$, we obtained the solution $(\eta_1, \eta_2, T) = (2.6885114, 1.2, 6.28134528)$ at the convergence tolerance $\varepsilon = 1E - 6$. Figure 1 shows the numerical solutions and trajectories in the phase plane, which were simulated using Kubíček's approach, the MK approach, and proposed PVM. The solutions of the two periods were obtained quickly at the maximum, and a slow variation occurred near the minimum. The trajectories of the solutions in the phase plane show that the periodicity was preserved. The results of these three methods were identical.

Example 2. The values of the coefficients a , b , c , and ω were reset to those used in Example 1, but α was changed from 5.0259 to 5.0258827. This example illustrates that Kubíček's approach is sensitive to the parameter α . Figure 2 shows the profiles of the numerical solutions Figure 2 (top panels) and phase plane Figure 2 (bottom panels) for the initial estimate $(\eta_1^{(0)}, \eta_2^{(0)}, T^{(0)}) = (3.0, 1.67, 5.6)$ when simulated using Kubíček's approach (left panels), the MK approach (middle

panels), and the proposed PVM (right panels). Kubíček's approach with fixed $k = 2$ and MK converge to the initial values $(\eta_1, \eta_2, T) = (2.6630419, 1.4719480, 2.418418)$ and $(-2.3248997, 1.67, \text{and } 34.8045258)$, respectively. Unexpectedly, the solutions were not periodic because the trajectory of the solution in the phase plane (left-bottom panel) was not closed, potentially because the solutions are sensitive to the convergence of Newton's iteration. However, the proposed PVM yielded a periodic solution with the initial value $(\eta_1, \eta_2, T) = (2.69035153, 1.1774813, 6.2842755)$ where the convergence tolerance $\varepsilon = 1E - 6$.

Example 3. Because the period is unprescribed, the problem (1) allows alternative solutions for various periods. The parameters were set the same as they were set in Example 1, except $\alpha = 1$. Figure 3 shows the result for the initial value $(\eta_1, \eta_2, T) = (2.70149522, 1.00802403, 12.56646993)$ obtained using the PVM with the initial estimate $(\eta_1^{(0)}, \eta_2^{(0)}, T^{(0)}) = (3.0, 1.0, 5.6)$. The numerical solutions and trajectories in the phase plane are shown in the left and right panels of the figure, respectively. Compared with Example 2, the period of the solution was $12.56646993 \approx 4\pi$, which can be confirmed according to the trajectory in the phase plane. According to this simulation, the unprescribed period enables the period to differ from the forced oscillator. However, the periodicity of the PVM was still robust.

Example 4. This example illustrates an autonomous Duffing-type oscillator [8, 17] expressed as

$$\begin{aligned} x'' - bx' - x + x^3 + \kappa z &= 0, \\ z' &= \omega(x' - z), \end{aligned} \quad (28)$$

where ω is the characteristic rate and κ is the feedback coefficient. This example shows that the external force is determined implicitly. Our goal was to calculate a periodic solution that differs from the motivation of [9]. Let $x_1 = x$, $x_2 = x'$, and $x_3 = z$. The system described in (28) then becomes

$$\begin{aligned} \frac{dx_1}{dt} &= x_2, \\ \frac{dx_2}{dt} &= x_1 - x_1^3 + bx_2 - \kappa x_3, \\ \frac{dx_3}{dt} &= \omega(x_2 - x_3) \end{aligned} \quad (29)$$

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Let  $\varepsilon$  be given.
Set the initial estimate  $(\eta_1, \dots, \eta_n, T)$ .
Repeat
  For  $k = 1$  to  $n$  // let  $\eta_k$  be fixed.
    Let  $U = (\eta_1, \dots, \eta_{k-1}, \eta_{k+1}, \dots, \eta_n, T)^t$ .
    Repeat
       $\delta \leftarrow [F_U(U)]^{-1} F(U)$ .
       $U \leftarrow U - \delta$ .
    until  $\|\delta\| < \varepsilon$ .
     $\eta_k \leftarrow \eta_k + \text{random number}$ 
  Next  $k$ 
until  $(\eta_1, \dots, \eta_n, T)$  remains constant and the postverification is successful.
    
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ALGORITHM 3: The postverification method.

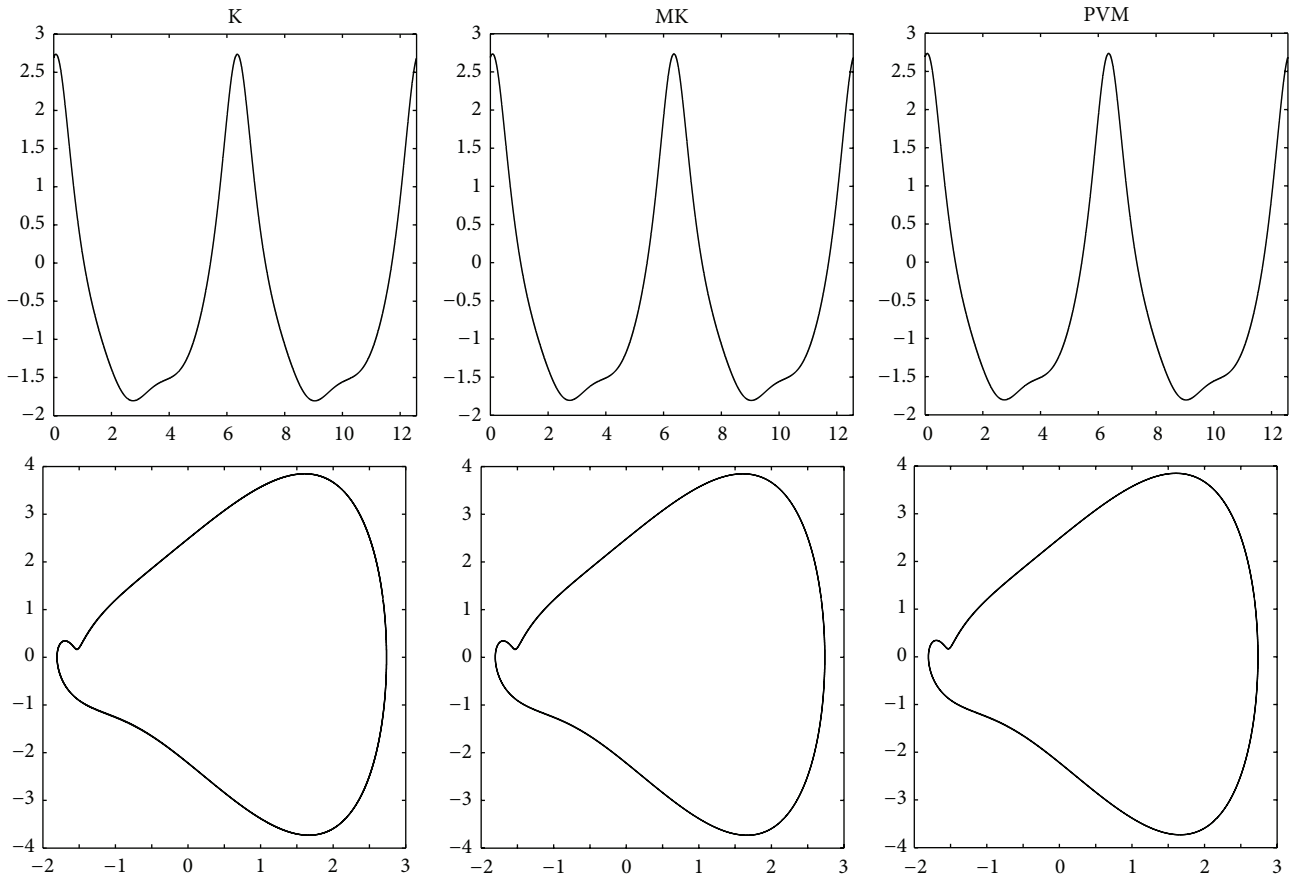


FIGURE 1: The numerical solution (top panels) and the trajectories in the phase planes (bottom panels) simulated by Kubíček's approach, MK, and PVM are shown from left to right panels, respectively. These results are identical for these three methods.

with $x_i(t + T) = x_i(t)$ for each i . The coefficients are set as $b = \kappa = 0.5$ and $\omega = 1$. The obtained initial conditions for a periodic solution are

$$\begin{aligned}
 x_1(0) &= -1.3241577E + 7, \\
 x_2(0) &= -6.9441420E + 13,
 \end{aligned}$$

$$\begin{aligned}
 x_3(0) &= -1.3241578E + 7, \\
 T &= 5.23147E - 7.
 \end{aligned}$$

(30)

Figure 4 shows the solution of two periods and the corresponding trajectory in three dimensions. Although the

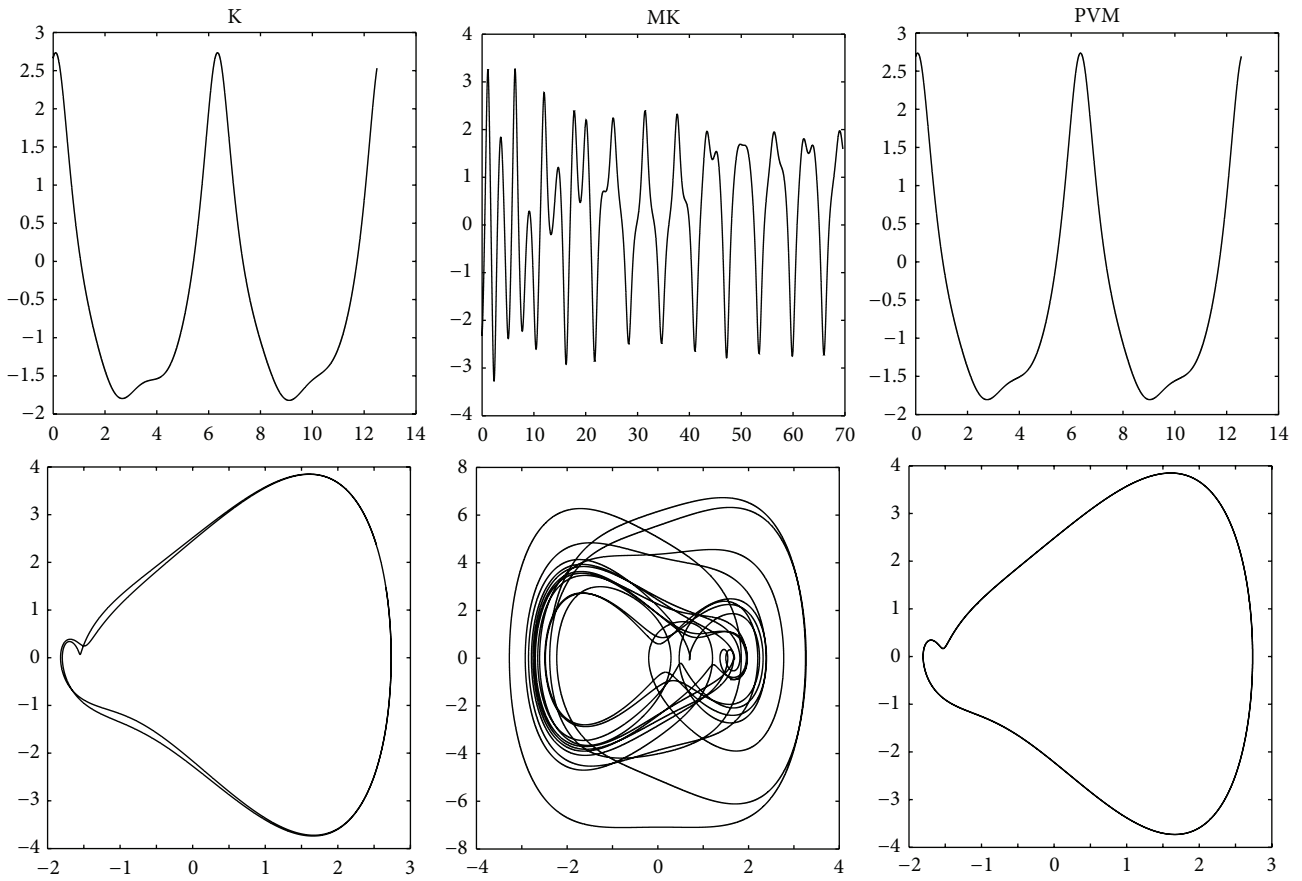


FIGURE 2: These results show the numerical solutions on $[0, 4\pi]$ (top panels) and the trajectories in the phase planes (bottom panels) simulated by Kubiček's approach, MK, and PVM are shown from left to right panels, respectively.

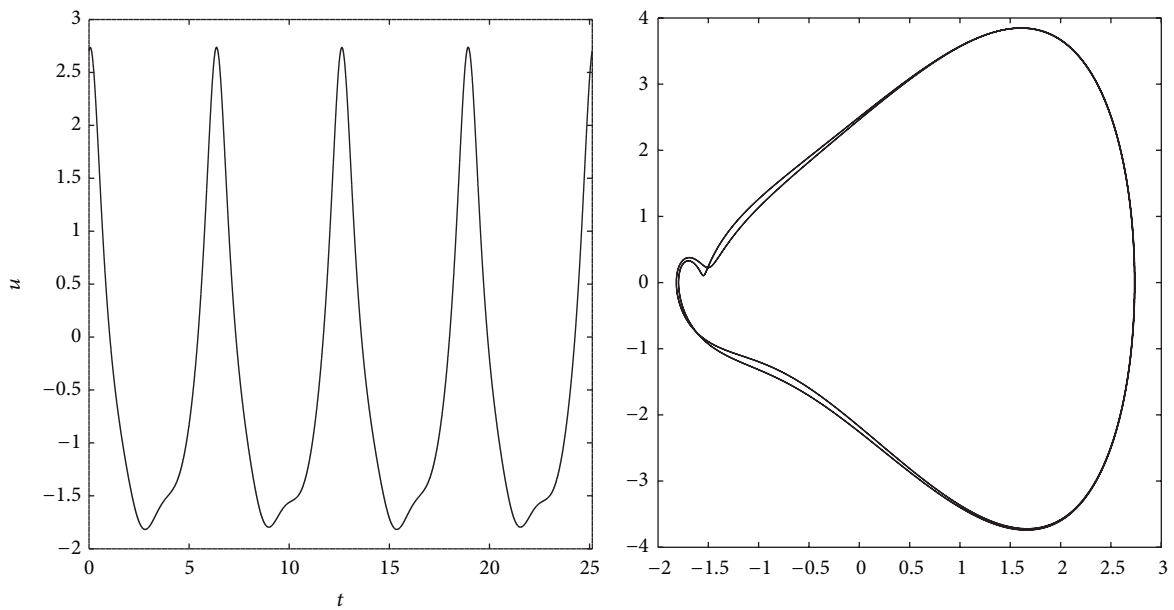


FIGURE 3: These results show the numerical solution on $[0, 8\pi]$ (left) and the trajectories in the phase plane (right) for Example 3. The periodicity can be observed by the profiles of the trajectories in the phase planes.

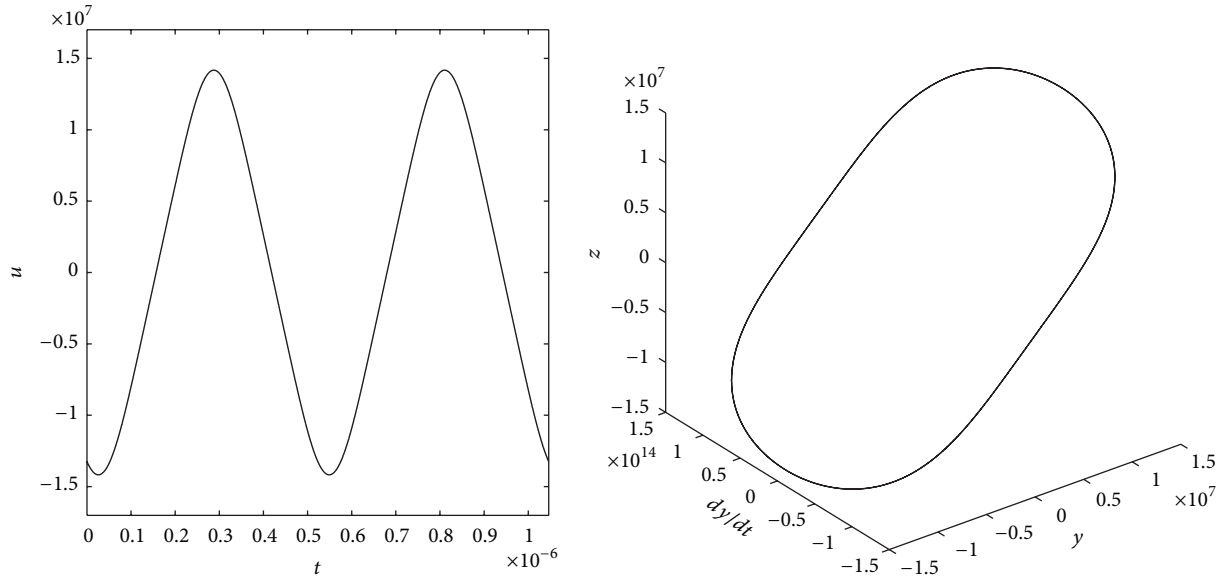


FIGURE 4: These results show the numerical solution x_1 (left) and the trajectories in 3D (right) for Example 4.

period was small and the amplitudes of the solutions were large, the solutions were stable. This stability enables identifying the periodic solutions. In this simulation, it was necessary to narrow the initial estimate to a value that was close to the solution.

4. Applications to Periodic-Doubling Bifurcation Problems

Kubíček and Marek [19] studied the system shown in (2) and replaced the source term $f_i(x_1, x_2, \dots, x_n, \alpha)$ depending on the value of parameter α . They set the eigenvalue of the monodromy matrix

$$M \equiv \frac{\partial \phi_i}{\partial \eta_j} \tag{31}$$

at -1 to identify the periodic-doubling bifurcation points and to extend the solution path with amplitudes from the periodic-doubling bifurcation points. There exist solutions $y_i(t)$ of period T where the function f_i has a period equal to T [21]. We concentrated on identifying the periodic-doubling bifurcation points and calculating the solution path for a prescribed initial value by using the quasi-arclength continuation method proposed in [8]. The other branch solution path of period $2T$ has been obtained based on the periodic-doubling bifurcation points.

The eigenvalues ρ_j of the monodromy matrix M are called Floquet multipliers, and the numbers $\gamma_i = \log \rho_i$ are called Floquet exponents. According to Floquet theory, a periodic solution is stable when $|\rho_j| < 1$ for each ρ_j (except $\rho_j = 1$). The stability of the periodic solutions on one branch may change according to variation in α ; one or two multipliers may be located outside the unit disk. Once a multiplier passes through the unit circle at -1 on the real axis, the periodic stable solution becomes unstable,

and another branch of periodic solutions can be identified at this point. In other words, the periodic-doubling bifurcation point exists where the two branches intersect at $(-1, 0)$ on the unit circle. Another method for determining the periodic-doubling bifurcation points is detailed as follows. The monodromy matrix M has an eigenvalue of -1 at the periodic-doubling bifurcation point; therefore, there exists a nonzero vector $v = (v_1, \dots, v_n)^T \neq 0$ such that

$$(M + I)v = 0 \tag{32}$$

is fulfilled. Because $y_i(0) = y_i(1)$, where $i = 1, 2, \dots, n$, these parameters satisfy (19) and

$$F_i(\eta_1, \eta_2, \dots, \eta_n, \alpha) = \phi_i(\eta_1, \eta_2, \dots, \eta_n, \alpha) - \eta_i = 0, \tag{33}$$

$i = 1, 2, \dots, n.$

The aforementioned system of $2n + 1$ variables $(\eta_1, \dots, \eta_n, \alpha, v_1, \dots, v_n)^T$ has only $2n$ equations. Furthermore, each nonzero multiplier of vector v is a solution of (19). A normalization approach,

$$v = v_k \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_k \\ \vdots \\ v_n \end{pmatrix}, \quad \text{for some } v_k \neq 0, \tag{34}$$

enables fixing the component v_k of $v = 1$ throughout the computation. Without loss of generality, we set $k = 1$. Let $P = (P_1, \dots, P_n, P_{n+1}, \dots, P_{2n})^T$ be the unique solution of the system shown in (19) and (20), which satisfy

$$P_i = F_i = 0, \tag{35}$$

$$P_{n+i} = \sum_{m=1}^n \frac{\partial \phi_i}{\partial \eta_m} v_m + v_i = 0,$$

where $i = 1, 2, \dots, n$. The solution w of $P(w) = 0$ (where $w_1 = 1$) is the periodic-doubling bifurcation point. Finally, the

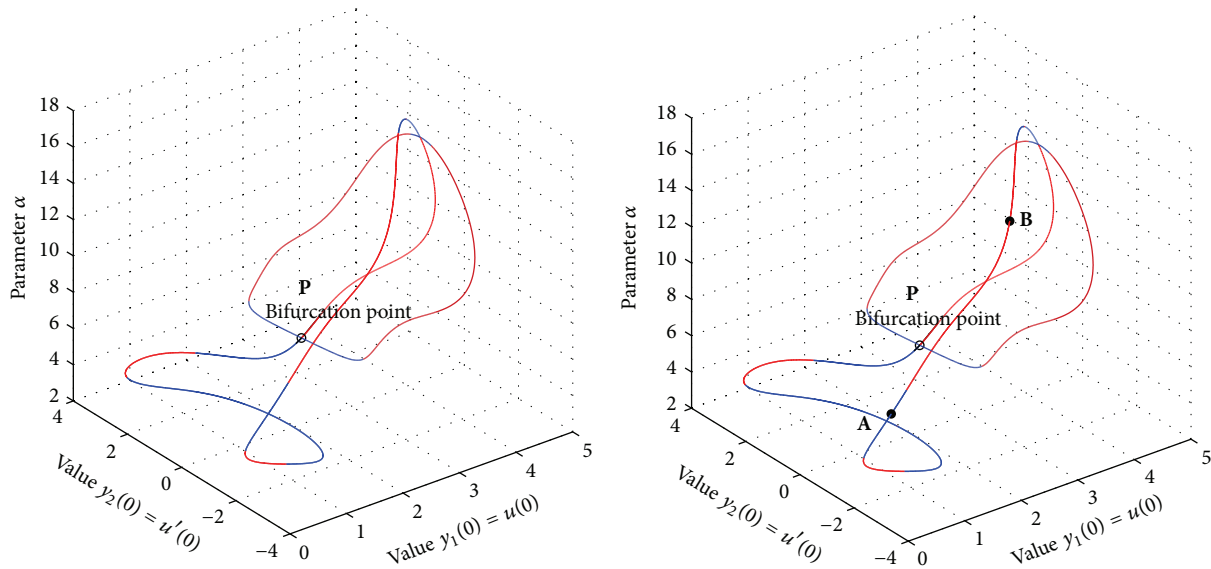


FIGURE 5: The calculation of bifurcation points depicts the periods of 2π (blue line) and 4π (red line) in the left panel and illustrates the stable (blue line) and unstable (red line) curves in the right panel.

solution w can be determined by applying Newton’s iterative method.

Example 5. The Duffing equation [20, 22] can be expressed as

$$u'' + 0.1u' + 0.1u + u^3 = \alpha \cos t, \tag{36}$$

where $y_1 = u, y_2 = u'$. Equation (36) then becomes

$$\begin{aligned} \frac{dy_1}{dt} &= y_2, \\ \frac{dy_2}{dt} &= -0.1y_2 - 0.1y_1 - y_1^3 + \alpha \cos t, \end{aligned} \tag{37}$$

$$y_1(0) = \eta_1, \quad y_2(0) = \eta_2.$$

We apply the PVM to solving (37) without prescribed periods, but for fixed parameter α . The obtained initial values represent the solutions of periods 2π and 4π , similarly for Examples 2 and 3. Furthermore, the quasi-arclength continuation method [8] can be applied to calculate the two curves that intersect at point **P**. Figure 5 shows the periodic-doubling bifurcation point $\mathbf{P}(\eta_1, \eta_2, \alpha) = (2.68964196, 1.18618911, 5.02588278)$.

The remaining of this section is to demonstrate the difficulty on calculation of the unstable periodic solutions. In the left panel of Figure 5, each point on the curve represents an initial value determined by (37). The blue curve denotes the initial values, which can be used to ascertain the solution of period 2π , and the red curve represents the initial values, which can be used to determine the solution of period 4π Figure 5 (left panel). The Floquet multipliers are used to evaluate the stability. According to the definition of stability, the initial values for the solution of periods 2π and 4π can be

divided into stable (blue line) and unstable (red line) regions. In the right panel of Figure 5, the markers

$$\begin{aligned} \mathbf{A}(y_1(0), y_2(0), \alpha) &= (1.42420721847224, -0.39906340734434, \\ &4.10971356401897), \\ \mathbf{B}(y_1(0), y_2(0), \alpha) &= (3.41661920410986, -0.61423445258396, \\ &12.69779142606720) \end{aligned} \tag{38}$$

denote the initial values of the solution of 2π , where the initial value **A** indicates the stable solution and the initial values **B** represents the unstable solution. Figure 6 depicts the phase plane and the unstable periodic solution requires a further study.

5. Conclusion

This study proposes a PVM for obtaining the periodic solutions of forced Duffing equations without a prescribed period. Specifically, to solve a system with $n + 1$ variables and n equations, we alternatively fixed a variable to improve Kubíček’s method and introduced a random small perturbation to verify the stability of the solutions. The postverification ensures that specific conditions are met whenever a solution is obtained. The proposed PVM is more robust on periodicity based on alternative convergence criterion, as demonstrated by our numerical results. We conclude that the most crucial application of the proposed PVM is determining the periodicity of solutions.

The application of the bifurcation problem demonstrated the branching of other periodic solutions that can be traced

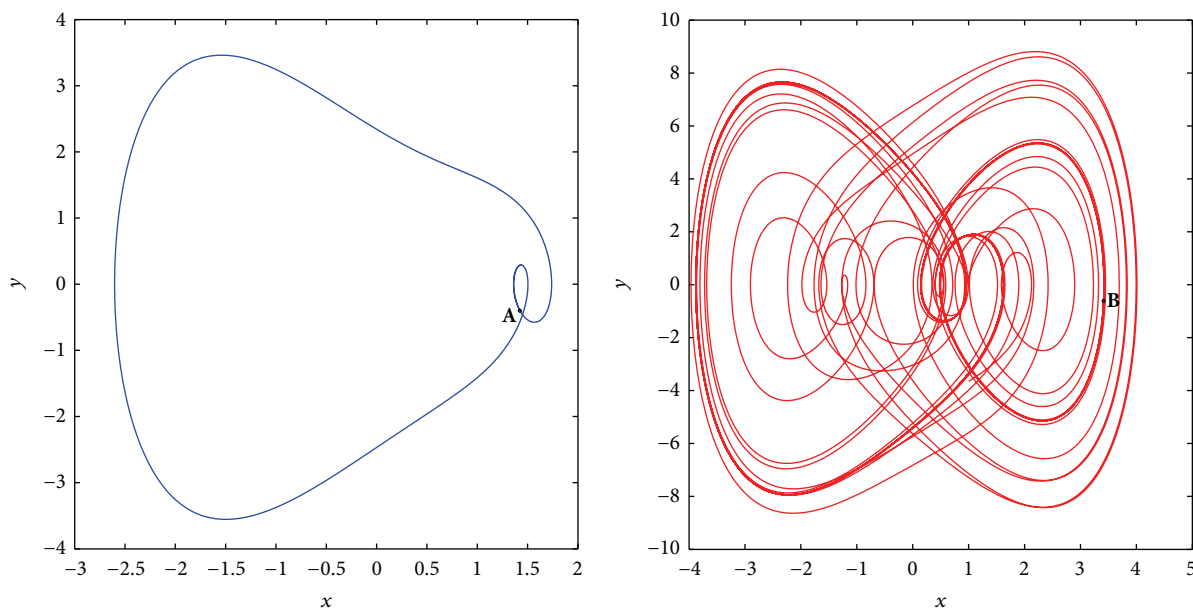


FIGURE 6: It illustrates the phase plane with the initial value point **A** (stable) in the left panel and shows the phase plane with the initial value point **B** (unstable) expressed in (37) in the right panel.

to the bifurcation points. However, the simulation and autonomous case revealed that the convergent initial values must be near to the initial estimate. Thus, calculating unstable periodic solutions remains challenging.

Conflict of Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

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