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Symplectic difference systems with periodic coefficients: Krein's traffic rules for multipliers

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Abstract

We investigate symplectic difference systems with periodic coefficients. We show that the known traffic rules for eigenvalues of the monodromy matrix of linear Hamiltonian differential and difference systems can be extended also to discrete symplectic systems. **MSC:** 39A12; 15A18

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1 Introduction

A symplectic difference system is the discrete first-order system

$$z_{k+1} = \mathcal{S}_k z_k, \quad k \in \mathbb{Z},$$

where the matrices $S_k \in \mathbb{C}^{2n \times 2n}$ are *symplectic, i.e.*,

$$S_k^* \mathcal{J} S_k = \mathcal{J}, \quad \mathcal{J} = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}.$$
(1)

In particular, if $S_k \in \mathbb{R}^{2n \times 2n}$, then we have $S_k^T \mathcal{J} S_k = \mathcal{J}$. The terminology for symplectic matrices is not unique. Our terminology is taken from [1]. An alternative terminology \mathcal{J} -unitary, \mathcal{J} -orthogonal can be found in [2]. Also, complex-valued matrices satisfying (1) are sometimes called *complex symplectic*.

Symplectic difference systems are discrete counterparts of linear Hamiltonian differential systems

$$z' = \mathcal{JH}(t)z$$

with the Hermitean matrix \mathcal{H} , *i.e.*, $\mathcal{H}^* = \mathcal{H}$, in the sense that they are the most general firstorder systems whose fundamental matrix is symplectic. It is known that one can establish oscillation and Sturmian theory for real symplectic difference systems similar to those for Hamiltonian differential systems, even if passing from continuous to discrete is far from being trivial. We refer to [3, 4] and the references given therein. Note that symplectic difference systems cover, as a very special case, the second-order Sturm-Liouville difference equation

$$\Delta(r_k \Delta x_k) + c_k x_{k+1} = 0, \quad r_k \neq 0,$$

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© 2013 Došlý; licensee Springer. This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/2.0), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited. whose qualitative theory is deeply developed; see, e.g., [5].

In our paper we suppose that the matrices S_k depend on a (generally complex-valued) parameter λ , *i.e.*, we consider the system

$$z_{k+1} = \mathcal{S}_k(\lambda) z_k. \tag{2}$$

Based on the papers [6, 7], see also the parts (i), (ii) of the next section, our assumptions on the matrices $S_k(\lambda)$ are following:

(A) System (2) is *symplectic* for $\lambda \in \mathbb{R}$, *i.e.*,

$$\mathcal{S}_k^*(\lambda)\mathcal{J}\mathcal{S}_k(\lambda) = \mathcal{J}, \quad k \in \mathbb{Z}.$$

(B) The matrices $S_k(\lambda)$ are *periodic* with the period $N \in \mathbb{N}$, *i.e.*,

$$S_k(\lambda) = S_{k+N}(\lambda), \quad k \in \mathbb{Z}, \lambda \in \mathbb{C}.$$

(C) The matrices $S_k(\lambda)$, k = 0, ..., N - 1, are *analytic* in a neighborhood of $\lambda_0 = 0$, in particular,

$$\mathcal{S}_k(\lambda) = \sum_{j=0}^{\infty} S_k^{[j]} \lambda^j, \qquad S_k^{[0]} = I,$$

where $S_k^{[j]} \in \mathbb{C}^{2n \times 2n}$ satisfy the periodicity condition

$$S_{k+N}^{[j]} = S_k^{[j]}, \quad k \in \mathbb{Z}, j \in \mathbb{N}.$$

Our principal concern is to study the stability (for $k \in \mathbb{Z}$) of solutions of (2). Concerning the continuous counterpart of symplectic difference systems, linear Hamiltonian differential systems

$$z' = \lambda \mathcal{JH}(t)z,\tag{3}$$

the stability theory of periodic Hamiltonian differential systems is deeply developed since the fifties of the last century. In that period, Russian mathematicians Gelfand, Krein, Lidskii, Yakubovich, Starzhinskii and others [2, 8, 9] published fundamental papers on stability of (3), which were summarized in the book [10].

The requirement that our system is symplectic for real λ reads in terms of the coefficients $S_k^{[j]}$ as follows. We suppose that the below given formulas hold for all indices k = 0, ..., N - 1, so we suppress the index k:

$$(S^{[1]})^* \mathcal{J} + \mathcal{J}S^{[1]} = 0, \tag{4}$$

$$\left(S^{[2]}\right)^* \mathcal{J} + \mathcal{J}S^{[2]} + \left(S^{[1]}\right)^* \mathcal{J}S^{[1]} = 0, \tag{5}$$

:

$$(S^{[n]})^* \mathcal{J} + \dots + (S^{[n-j]})^* \mathcal{J}S^{[j]} + \dots + \mathcal{J}S^{[n]} = 0,$$
(6)

The following statement plays the crucial role in the later introduced concept of a symplectic system of *positive type*.

Proposition 1 Suppose that (4)-(6) hold. Then there exist Hermitean matrices $\mathcal{A}_k(\lambda, \overline{\lambda})$ such that for $\lambda \in \mathbb{C}$

$$\mathcal{S}_{k}^{*}(\lambda)\mathcal{J}\mathcal{S}_{k}(\lambda) = \mathcal{J} + (\bar{\lambda} - \lambda)\mathcal{A}_{k}(\lambda, \bar{\lambda}).$$
⁽⁷⁾

Proof Suppressing again the index k, using (4),

+ · · · .

$$S^{*}(\lambda)\mathcal{J}S(\lambda) = \left(I + \sum_{j=1}^{\infty} \bar{\lambda}^{j} \left(S^{[j]}\right)^{*}\right) \mathcal{J}\left(I + \sum_{j=1}^{\infty} \lambda^{j} S^{[j]}\right)$$
$$= \mathcal{J} + (\bar{\lambda} - \lambda) \left(-\mathcal{J}S^{[1]}\right)$$
$$+ (\bar{\lambda})^{2} \left(S^{[2]}\right)^{*} \mathcal{J} + \bar{\lambda}\lambda \left(S^{[1]}\right)^{*} \mathcal{J}S^{[1]} + \lambda^{2} \mathcal{J}S^{[2]} + \cdots$$
(8)
$$+ (\bar{\lambda})^{n} \left(S^{[n]}\right)^{*} \mathcal{J} + \cdots + (\bar{\lambda})^{n-j} \lambda^{j} \left(S^{[n-j]}\right)^{*} \mathcal{J}S^{[j]} + \cdots + \lambda^{n} \mathcal{J}S^{[n]}$$
(9)

Concerning the expression (8), using (5), we immediately have

$$(\bar{\lambda})^{2} \left(S^{[2]} \right)^{*} \mathcal{J} + \bar{\lambda} \lambda \left(S^{[1]} \right)^{*} \mathcal{J} S^{[1]} + \lambda^{2} \mathcal{J} S^{[2]} = (\bar{\lambda} - \lambda) \left[\bar{\lambda} \left(S^{[2]} \right)^{*} \mathcal{J} - \lambda \mathcal{J} S^{[2]} \right].$$

Generally, concerning expression (9), consider the matrices $X_1^{[n]}, \ldots, X_{n-2}^{[n]}, n \ge 3$, given by the formulas

$$X_{1}^{[n]} = (S^{[n]})^{*} \mathcal{J} + (S^{[n-1]})^{*} \mathcal{J} S^{[1]},$$

$$X_{2}^{[n]} = X_{1}^{[n]} + (S^{[n-2]})^{*} \mathcal{J} S^{[2]},$$

$$\vdots$$

$$X_{n-2}^{[n]} = X_{n-3}^{[n]} + (S^{[2]})^{*} \mathcal{J} S^{[n-2]}.$$
(10)

Then formula (6) implies that we also have

$$-X_{n-2}^{[n]} = \left(S^{[1]}\right)^* \mathcal{J}S^{[n-1]} + \mathcal{J}S^{[n]}.$$

Now, by a direct computation, using (10), we have that (9) for $n \ge 3$ equals

$$(\bar{\lambda}-\lambda)\big[(\bar{\lambda})^{n-1}\big(S^{[n]}\big)^*\mathcal{J}+\dots+(\bar{\lambda})^{n-j-1}\lambda^j X_j^{[n]}+\dots-\lambda^{n-1}\mathcal{J}S^{[n]}\big],\tag{11}$$

i.e., one can factor out the term $(\bar{\lambda} - \lambda)$ from the expression for $S^*(\lambda) \mathcal{J}S(\lambda) - \mathcal{J}$. It remains to prove that $\mathcal{A}(\lambda, \bar{\lambda})$ is Hermitean. This matrix is given by an infinite series, and the typical summand in this series is given by the expression in brackets in (11). To prove that this

expression is Hermitean, it suffices to show that $(X_j^{[n]})^* = X_{n-j-1}^{[n]}$. We have, using (6),

$$(X_j^{[n]})^* = [(S^{[n]})^* J + (S^{[n-1]})^* \mathcal{J}S^{[1]} + \dots + (S^{[n-j]})^* \mathcal{J}S^{[j]}]^*$$

= $-[\mathcal{J}S^{[n]} + (S^{[1]})^* \mathcal{J}S^{[n-1]} + \dots + (S^{[j]})^* \mathcal{J}S^{[n-j]}]$
= $(S^{[n]})^* \mathcal{J} + (S^{[n-1]})^* \mathcal{J}S^{[1]} + \dots + (S^{[j+1]})^* \mathcal{J}S^{[n-j-1]}$
= $X_{n,j-1}^{[n]}.$

Consequently, each term in the infinite series expressing $\mathcal{A}(\lambda, \overline{\lambda})$ is a Hermitean matrix, and hence this matrix is Hermitean as well.

2 Particular cases

In this section, we treat special cases of our periodic symplectic difference system as presented in the recent literature.

(i) Second-order matrix difference system.

The results concerning this system were published in [6]. We consider the second-order matrix difference system with $P_k \in \mathbb{C}^{n \times n}$

$$\Delta^2 x_{k-1} + \lambda^2 P_k x_k = 0, \qquad P_k^* = P_k, \qquad P_{k+N} = P_k.$$
(12)

In this case, we write the system in the form

$$\Delta\left(\frac{1}{\lambda}\Delta x_{k-1}\right)+\lambda P_k x_k=0.$$

If we put $u_k = \frac{1}{\lambda} \Delta x_{k-1}$, then $\Delta u_k = -\lambda P_k x_k$, *i.e.*, $u_{k+1} = -\lambda P_k x_k + u_k$ and $x_{k+1} - x_k = \lambda u_{k+1} = \lambda (-\lambda P_k x_k + u_k)$. Consequently,

$$\begin{pmatrix} x_{k+1} \\ u_{k+1} \end{pmatrix} = \begin{bmatrix} I + \begin{pmatrix} 0 & I \\ -P_k & 0 \end{pmatrix} \lambda + \begin{pmatrix} -P_k & 0 \\ 0 & 0 \end{pmatrix} \lambda^2 \end{bmatrix} \begin{pmatrix} x_k \\ u_k \end{pmatrix}.$$
 (13)

Denote by $S^{[1]}$ and $S^{[2]}$ the matrices by λ and λ^2 , respectively. By a direct computation, taking into account the form of the matrices $S^{[1]}$ and $S^{[2]}$, we obtain for $\lambda \in \mathbb{C}$

$$\begin{split} S^{*}(\lambda)\mathcal{J}S(\lambda) \\ &= \mathcal{J} + \bar{\lambda} \left(S^{[1]} \right)^{*} \mathcal{J} + \lambda \mathcal{J}S^{[1]} + (\bar{\lambda})^{2} \left(S^{[2]} \right)^{*} \mathcal{J} + \bar{\lambda}\lambda \left(S^{[1]} \right)^{*} \mathcal{J}S^{[1]} + \lambda^{2} \mathcal{J}S^{[2]} \\ &+ (\bar{\lambda})^{2} \lambda \left(S^{[2]} \right)^{*} \mathcal{J}S^{[1]} + \lambda^{2} \bar{\lambda} \left(S^{[1]} \right)^{*} \mathcal{J}S^{[2]} + (\bar{\lambda}\lambda)^{2} \left(S^{[2]} \right)^{*} \mathcal{J}S^{[2]} \\ &= \mathcal{J} + (\bar{\lambda} - \lambda) \left(-\mathcal{J}S^{[1]} \right) + (\bar{\lambda})^{2} \begin{pmatrix} 0 & -P^{*} \\ 0 & 0 \end{pmatrix} + \lambda \bar{\lambda} \begin{pmatrix} 0 & P^{*} \\ -P & 0 \end{pmatrix} \\ &+ \lambda^{2} \begin{pmatrix} 0 & 0 \\ P & 0 \end{pmatrix} + \left[(\bar{\lambda})^{2} \lambda - \bar{\lambda} \lambda^{2} \right] \begin{pmatrix} P^{*}P & 0 \\ 0 & 0 \end{pmatrix} \\ &= \mathcal{J} + (\bar{\lambda} - \lambda) \begin{pmatrix} P + |\lambda|^{2} P^{*}P & -\bar{\lambda}P \\ -\lambda P & I \end{pmatrix}. \end{split}$$

Therefore, (13) is really a symplectic system for $\lambda \in \mathbb{R}$, and to verify (5) and (6) with n = 3, 4 is a matter of an easy computation.

(ii) Linear Hamiltonian difference system.

The stability of this system was investigated in [7]; see also [11]. We consider the system

$$\Delta \begin{pmatrix} x_k \\ u_k \end{pmatrix} = \lambda \mathcal{J} \begin{pmatrix} -C_k & A_k \\ A_k^* & B_k \end{pmatrix} \begin{pmatrix} x_{k+1} \\ u_k \end{pmatrix}$$
(14)

with the Hermitean matrices *B*, *C*, *i.e.*, $B = B^*$, $C = C^*$. This system can be written in the form

$$\begin{pmatrix} x_{k+1} \\ u_{k+1} \end{pmatrix} = \begin{pmatrix} (I - \lambda A_k)^{-1} & \lambda (I - \lambda A_k)^{-1} B_k \\ \lambda C_k (I - \lambda A_k)^{-1} & \lambda^2 C_k (I - \lambda A_k)^{-1} B_k + I - \lambda A_k^* \end{pmatrix} \begin{pmatrix} x_k \\ u_k \end{pmatrix}.$$

If we denote by $S(\lambda)$ the matrix in the last system, we have, using the formula $(I - \lambda A_k)^{-1} = \sum_{i=0}^{\infty} \lambda^j A_k^j$ (suppressing the index *k*),

$$S(\lambda) = I + \lambda \begin{pmatrix} A & B \\ C & -A^* \end{pmatrix} + \sum_{j=2}^{\infty} S^{[j]} \lambda^j$$

with the matrices (for $j \ge 2$)

$$S^{[j]} = \begin{pmatrix} A^j & A^{j-1}B \\ CA^{j-1} & CA^{j-2}B \end{pmatrix}.$$

To show that the matrices $S^{[j]}$, $j \ge 2$, together with the matrix

$$S^{[1]} = \begin{pmatrix} A & B \\ C & -A^* \end{pmatrix},$$

satisfy symplecticity condition (6) requires rather tedious computations. Consider two successive terms in (6)

$$(S^{[n-j+1]})^* \mathcal{J}S^{[j-1]} + (S^{[n-j]})^* \mathcal{J}S^{[j]}$$

and let us compute the block matrix in the left upper corner. By a direct substitution, we obtain the expression

$$-(A^{*})^{n-j-1}CA^{j} + (A^{*})^{n-j}CA^{j-1} - (A^{*})^{n-j}CA^{j-1} + (A^{*})^{n-j+1}CA^{j-2}.$$
(15)

Hence, we immediately see that the second and third term cancel. Moreover, the first one cancels against a term coming from the expression $(S^{[n-j+2]})^* \mathcal{J}S^{[j-2]}$, and similarly, the last term in (15) also disappears. By the same reasoning, we also show that the remaining three blocks in (6) equal zero. Also, it is shown in [7] that

$$S^{*}(\lambda)\mathcal{J}S(\lambda) = \mathcal{J} + (\bar{\lambda} - \lambda)D^{*}(\lambda) \begin{pmatrix} -C & A^{*} \\ A & B \end{pmatrix} D(\lambda),$$
(16)

where

$$D(\lambda) = \begin{pmatrix} (I - \lambda A)^{-1} & \lambda (I - \lambda A)^{-1} B \\ 0 & I \end{pmatrix}.$$
 (17)

Therefore, we see that Hamiltonian system (14) is a particular case of our symplectic system (2) satisfying (7).

(iii) Exponential system.

By an 'exponential' system we mean the system (2), where

$$S_k(\lambda) = I + \sum_{j=1}^{\infty} R_k^j \frac{\lambda^j}{j!} = \exp\{\lambda R_k\}$$
(18)

with the matrices $R_k \in \mathbb{C}^{2n \times 2n}$ satisfying

$$R_k^* \mathcal{J} + \mathcal{J} R_k = 0, \qquad R_{k+N} = R_k.$$

When *S* is a symplectic matrix 'close' to the identity matrix (||I - S|| < 1), then there is a well-defined matrix log *S*, and we have

$$S^{\lambda} = \exp\{\lambda \log S\} = \sum_{j=0}^{\infty} (\log S)^j \frac{\lambda^j}{j!},$$

and it is known (see, *e.g.*, [12, p.8]) that the logarithm of a symplectic matrix is the Hamiltonian matrix, *i.e.*, $(\log S)^* \mathcal{J} + \mathcal{J} \log S = 0$.

The symplectic system where $S_k(\lambda)$ depend on λ in the form (18) was investigated in [13]. Here we have

$$S_k^{[j]} = \frac{R_k^j}{j!},$$

so identity (6) holds as can be verified by a short direct computation. Moreover, it was shown in [13] that the matrix $\mathcal{A}(\lambda, \overline{\lambda})$ in (7) is of the form

$$\mathcal{A}(\lambda,\bar{\lambda}) = \sum_{j=0}^{\infty} (-1)^{j} \frac{(\bar{\lambda}-\lambda)^{2j}}{(2j+1)!} (R^{*})^{j} (-\mathcal{J}R) R^{j} + \sum_{j=1}^{\infty} (-1)^{j} \frac{(\bar{\lambda}-\lambda)^{2j-1}}{(2j)!} (R^{*})^{j} \mathcal{J}R^{j}.$$
(19)

Also, [13] shows, using the time scales calculus, that the dependence of $S(\lambda)$ on λ given by (18) is, in a certain sense, a natural discrete counterpart of the multiplicative dependence on λ in (3).

(iv) Symplectic systems with nonlinear dependence on a spectral parameter.

In [14], oscillatory and spectral properties of (2) with real-valued $S_k(\lambda)$ were investigated. The crucial assumption in that paper is that the matrices (with $= \frac{d}{d\lambda}$)

$$\Psi_k(\lambda) := \hat{\mathcal{S}}_k(\lambda) \mathcal{J} \mathcal{S}_k^T(\lambda) \quad \text{are symmetric}$$
(20)

(and nonpositive definite). We show that our assumptions (4), (5), and (6) imply (20). First of all, observe that the symplecticity of a matrix S is equivalent to the symplecticity of its

transpose S^T . Hence, identities (4), (5), (6) are equivalent (for real $S^{[j]}$) to

$$S^{[1]}\mathcal{J} + \mathcal{J}(S^{[1]})^{T} = 0,$$

$$S^{[2]}\mathcal{J} + \mathcal{J}(S^{[2]})^{T} + S^{[1]}\mathcal{J}(S^{[1]})^{T} = 0,$$

$$\vdots$$

$$S^{[n]}\mathcal{J} + \dots + S^{[n-j]}\mathcal{J}(S^{[j]})^{T} + \dots + \mathcal{J}(S^{[n]})^{T} = 0.$$
(21)

We have (suppressing the index *k*)

$$\begin{split} \dot{S}(\lambda)\mathcal{J}S^{T}(\lambda) \\ &= \left(\sum_{j=1}^{\infty} j\lambda^{j-1}S^{[j]}\right)\mathcal{J}\left(\sum_{j=0}^{\infty} \lambda^{j}S^{[j]}\right)^{T} \\ &= S^{[1]}\mathcal{J} + \lambda \left[S^{[1]}\mathcal{J}(S^{[1]})^{T} + 2S^{[2]}\mathcal{J}\right] + \lambda^{2} \left[S^{[1]}\mathcal{J}(S^{[2]})^{T} + 2S^{[2]}\mathcal{J}(S^{[1]})^{T} + 3S^{[3]}\mathcal{J}\right] + \cdots \\ &+ \lambda^{n} \left[S^{[1]}\mathcal{J}(S^{[n]})^{T} + 2S^{[2]}\mathcal{J}(S^{[n-1]})^{T} + \cdots + jS^{[j]}\mathcal{J}(S^{[n+1-j]})^{T} + \cdots \\ &+ (n+1)S^{[n+1]}\mathcal{J}\right] + \cdots . \end{split}$$

We will show that the matrices by all powers of λ are symmetric. Suppose that n = 2j - 1 is odd; if *n* is even, the situation is similar. We arrange the summands by λ^n in brackets in the previous formula into the triangular scheme as follows:

$$S^{[1]}\mathcal{J}(S^{[n]})^{T} + \dots + S^{[j]}\mathcal{J}(S^{[j]})^{T} + S^{[j+1]}\mathcal{J}(S^{[j-1]})^{T} + \dots + S^{[n+1]}\mathcal{J}$$

$$\vdots$$

$$+ S^{[j]}\mathcal{J}(S^{[j]})^{T} + S^{[j+1]}\mathcal{J}(S^{[j-1]})^{T} + \dots + S^{[n+1]}\mathcal{J}$$

$$+ S^{[j+1]}\mathcal{J}(S^{[j-1]})^{T} + \dots + S^{[n+1]}\mathcal{J}$$

$$\vdots$$

$$+ S^{[n+1]}\mathcal{J}.$$
(22)

Substituting for the first *j* lines in this expression from (21) (with n + 1 instead of *n*), we can write (22) as (we arrange again the first *j* lines into a triangular-type scheme)

$$-\mathcal{J}(S^{[n+1]})^{T} - \mathcal{J}(S^{[n+1]})^{T} - S^{[1]}\mathcal{J}(S^{[n]})^{T}$$

$$\vdots -\mathcal{J}(S^{[n+1]})^{T} - S^{[1]}\mathcal{J}(S^{[n]})^{T} - \dots - S^{[j-1]}\mathcal{J}(S^{[j+1]})^{T} + S^{[j+1]}\mathcal{J}(S^{[j-1]})^{T} + \dots + S^{[n]}\mathcal{J}(S^{[1]})^{T} + S^{[n+1]}\mathcal{J}$$

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$$+ S^{[n]} \mathcal{J} (S^{[1]})^T + S^{[n+1]} \mathcal{J}$$
$$+ S^{[n+1]} \mathcal{J}.$$

Now we see that the sum of the terms on the first and the last line is a symmetric matrix (since $\mathcal{J}^T = -\mathcal{J}$). Similarly, the sum of the second and last but one line terms is a symmetric matrix, *etc.* Altogether, the term by λ^n , $n \in \mathbb{N}$, is a symmetric matrix, which means that the matrix Ψ is really symmetric.

3 Central stability zone

We rewrite the matrices $S(\lambda)$ in system (2) into the form

$$S_k(\lambda) = I + \lambda S_k^{[1]} + o(\lambda) \quad \text{as } \lambda \to 0$$

with $S^{[1]}$ satisfying (4). The monodromy matrix of (2) is

$$\mathcal{U}_N(\lambda) := \mathcal{S}_{N-1}(\lambda) \cdots \mathcal{S}_0(\lambda) = I + \lambda \sum_{k=0}^{N-1} S_k^{[1]} + o(\lambda)$$

as $\lambda \rightarrow 0$. Denote

$$S^{[1]} = \sum_{k=0}^{N-1} S_k^{[1]}.$$

Then, for $\lambda \in \mathbb{R}$ with $|\lambda|$ small, we have $\log \mathcal{U}_N(\lambda) = \lambda \mathcal{S}^{[1]} + o(\lambda)$. If we denote by $\gamma_j(\lambda)$ the eigenvalues of the matrix $\log \mathcal{U}_N(\lambda)$, we have

$$\gamma_j(\lambda) = \lambda s_j + o(\lambda),$$

where s_i are the eigenvalues of $\mathcal{S}^{[1]}$.

The proof of the following statement (which shows the non-emptiness of the central stability zone under an additional assumption that the eigenvalues of a certain matrix are distinct) is essentially the same as that of [2, Theorem 2.2].

Theorem 1 Suppose that the matrix $\mathcal{JS}^{[1]}$ is negative definite and the eigenvalues s_j of $\mathcal{S}^{[1]}$ are distinct. Then there exists l > 0 such that the solutions of (2) are bounded on \mathbb{Z} for real λ satisfying $|\lambda| < l$.

Proof The proof is based on the statement that if *A* is a Hermitean positive definite matrix, then the equation det($A + i\mu \mathcal{J}$) = 0 has only real roots; see [15]. So, let $A = -\mathcal{J}\mathcal{S}^{[1]}$ and look for the eigenvalues of $\mathcal{S}^{[1]}$ in the form $i\sigma$. Then

$$0 = \det(\mathcal{S}^{[1]} - i\sigma I) = \det \mathcal{J}(-\mathcal{J}\mathcal{S}^{[1]} + i\sigma \mathcal{J}) = \det(A + i\sigma \mathcal{J})$$

(since det $\mathcal{J} = 1$, see [1]), and hence $\sigma \in \mathbb{R}$ and $s_j = i\sigma_j$. Now, for $\lambda \in \mathbb{R}$, $|\lambda|$ small, the numbers $\gamma_j(\lambda) = i\sigma_j\lambda + o(\lambda)$ are still different. However, at the same time, for $\lambda \in \mathbb{R}$, the matrix $\mathcal{U}_N(\lambda)$ is symplectic, *i.e.*, $\mathcal{U}_N^*(\lambda)\mathcal{J}\mathcal{U}_N(\lambda) = \mathcal{J}$, and hence its eigenvalues $\rho_j(\lambda) = \exp{\{\gamma_j(\lambda)\}}$

are symmetric with respect to the unit circle, and hence $\gamma_j(\lambda)$ are symmetric with respect to the imaginary axis. But this implies that $\gamma_j(\lambda)$ themselves are on the imaginary axis, *i.e.*, $|\rho_j(\lambda)| = 1$. This means, in view of the later given formula (25), that solutions of (2) are bounded on \mathbb{R} for $\lambda \in \mathbb{R}$, $|\lambda|$ small.

4 Krein's traffic rules

So far we did not need any assumption on the matrices $S(\lambda)$ for $\lambda \in \mathbb{C}$. Complex values of λ are needed to get an extension of the previous theorem to the case when the eigenvalues of $S^{[1]}$ are *not necessarily distinct*. In this situation, the stability behavior of solutions of (2) is heavily based on Krein's traffic rules for the eigenvalues of the monodromy matrix $U_N(\lambda)$.

The crucial concept is the concept of a *positive-type system*. In the Hamiltonian setting (14), this means

$$H_k \ge 0$$
, $k = 0, ..., N - 1$ and $\sum_{k=0}^{N-1} H_k > 0$;

see [7], where

$$H_k = \begin{pmatrix} -C_k & A_k^* \\ A_k & B_k \end{pmatrix},\tag{23}$$

here > (\geq , <, \leq) means positive (nonnegative, negative, nonpositive) definiteness of the matrix indicated. Under this assumption, we have for *solutions* $z = \begin{pmatrix} x \\ \mu \end{pmatrix}$ of system (14)

$$\Delta(z_k^*\mathcal{J}z_k) = (\bar{\lambda} - \lambda)\binom{x_{k+1}}{u_k}^*H_k\binom{x_{k+1}}{u_k} \geq 0;$$

see [7]. Consequently, one can speak about \mathcal{J} -monotonicity of the monodromy matrix $\mathcal{U}_N(\lambda)$ depending on whether Im $\lambda > 0$ or Im $\lambda < 0$, in the same way as in [2].

In our symplectic setting, we have

$$S_k^*(\lambda)\mathcal{J}S_k(\lambda) - \mathcal{J} = (\bar{\lambda} - \lambda)[-\mathcal{J}S_k^{[1]} + G_k(\lambda, \bar{\lambda})],$$

where G_k is the matrix containing powers of λ and $\overline{\lambda}$ such that

$$G_k(\lambda, \overline{\lambda}) = \mathcal{J}S_k^{[1]} + \mathcal{A}_k(\lambda, \overline{\lambda});$$

see (7). Actually, as we will see later (proof of Theorem 2), the eigenvalues of the monodromy matrix (*i.e.*, multipliers) follow Krein's traffic rules if we suppose that there exists $\varepsilon > 0$ such that

$$\mathcal{A}_k(\lambda,\bar{\lambda}) \ge 0, \quad k = 0, \dots, N-1 \quad \text{and} \quad \sum_{k=0}^{N-1} \mathcal{A}_k(\lambda,\bar{\lambda}) > 0 \quad \text{for } |\operatorname{Im} \lambda| < \varepsilon.$$
 (24)

A closer examination of the proofs in [2] reveals that when investigating the behavior of eigenvalues of the monodromy matrix $U_N(\lambda)$ in dependence on the parameter λ , one

can forget about the original system from which this monodromy matrix originates (differential system (3) or difference system (2)). The only important fact is that this matrix is symplectic for $\lambda \in \mathbb{R}$ and \mathcal{J} -monotonic for $\lambda \in \mathbb{C} \setminus \mathbb{R}$. Under assumption (24), the traffic rules of eigenvalues of $\mathcal{U}_N(\lambda)$ are the same as in [2, 6].

Recall now some important facts and concepts. Let $\xi \in \mathbb{C}^{2n}$ be an eigenvector of $\mathcal{U}_N(\lambda)$ corresponding to the eigenvalue ρ , which is called the *multiplier* of $\mathcal{U}_N(\lambda)$. The solution z of (2) given by the initial condition $z_0 = \xi$ satisfies

$$z_N = \mathcal{U}_N(\lambda)\xi = \rho\xi = \rho z_0,\tag{25}$$

and generally $z_{k+N} = \rho z_k$. Consequently, solutions of (2) are bounded for $k \in \mathbb{Z}$ if and only if all eigenvalues of $\mathcal{U}_N(\lambda)$ have modulus one and are of simple type, *i.e.*, multiple eigenvalues have the same algebraic and geometric multiplicity.

Clearly, $\lambda = 0$ is the stability point of (2) since all solutions are constant in this case. Theorem 1 states that if the eigenvalues of $S^{[1]}$ are distinct, we have a central stability zone for (2). The restriction on distinctness of eigenvalues is removed in [2] (see also [6, 7] in the discrete case) using a perturbation principle, then Theorem 1 remains true when Krein's traffic rules for eigenvalues are preserved. For these rules, the crucial concept is the concept of a *multiplier of definite type* which is defined as follows.

Definition 1 Let $|\rho_0| = 1$ be an eigenvalue of the monodromy matrix \mathcal{U}_N , and let \mathcal{L}_{ρ_0} be its eigenspace. If $iu^* \mathcal{J}u > 0$ for every $u \in \mathcal{L}_{\rho_0}$, then the eigenvalue ρ_0 is said to be the multiplier of the *first kind* and it is called of *second kind* if $iu^* \mathcal{J}u < 0$ for $u \in \mathcal{L}_{\rho_0}$. If there exists $0 \neq u \in \mathcal{L}_{\rho_0}$ such that $u^* \mathcal{J}u = 0$, the multiplier ρ_0 is of *indefinite (mixed) type*. If $|\rho_0| \neq 1$ is an eigenvalue of \mathcal{U}_N , then it is called of the first kind if $|\rho_0| > 1$ and of the second kind if $|\rho_0| < 1$.

The following statement is proved similarly as [6, Theorem 4.6].

Theorem 2 Suppose that (24) holds and that for $\lambda_0 \in \mathbb{R}$ the monodromy matrix $\mathcal{U}_N(\lambda_0)$ has a multiplier ρ_0 , with $|\rho_0| = 1$, of definite kind with the multiplicity m > 0. Then in a δ -neighborhood of λ_0 , m eigenvalues $\rho_1(\lambda), \ldots, \rho_m(\lambda)$ of $\mathcal{U}_N(\lambda)$, for which $\rho_j(\lambda) \to \rho_0$ as $\lambda \to \lambda_0$, $j = 1, \ldots, m$, are situated. Moreover, if $\lambda \in \mathcal{O}_{\delta}(\lambda_0)$ is real, then $\rho_j(\lambda), j = 1, \ldots, m$, stay on the unit circle. The coefficients α_j in the expansion

$$\rho_j(\lambda) = \rho_0 \Big[1 + \alpha_j (\lambda - \lambda_0) + o(\lambda - \lambda_0) \Big]$$

are given by

$$\alpha_{j} = -\frac{1}{u_{j}^{*}\mathcal{J}u_{j}}\sum_{k=0}^{N-1} (z_{k}^{[j]})^{*} [-\mathcal{J}S_{k}^{[1]} + G_{k}(\lambda,\bar{\lambda})]z_{k}^{[j]},$$

where u_j are properly taken eigenvectors of $\mathcal{U}_N(\lambda_0)$ corresponding to ρ_0 and $z_k^{[j]} = \mathcal{U}_k(\lambda_0)u_j$ are the solutions of (2) with $\lambda = \lambda_0$ given by the initial condition $z_0^{[j]} = u_j$. *Proof* The only difference with respect to the proof of Theorem [6, Theorem 4.6] is that for $\lambda \in \mathcal{O}_{\delta}(\lambda_0)$ and solutions $z^{[j]}$, j = 1, ..., m, we have

$$(z_N^{[j]})^* \mathcal{J} z_N^{[j]} - u_j^* \mathcal{J} u_j = \sum_{k=0}^{N-1} (z_k^{[j]})^* [S_k^*(\lambda) \mathcal{J} S_k(\lambda) - \mathcal{J}] z_k^{[j]}$$

= $-2i \operatorname{Im} \lambda \sum_{k=0}^{N-1} (z_k^{[j]})^* [-\mathcal{J} S_k^{[1]} + G_k(\lambda, \bar{\lambda})] z_k^{[j]}.$

Hence

$$(|\rho_{j}(\lambda)|^{2}-1)u_{j}^{*}\mathcal{J}u_{j}=-2i\operatorname{Im}\lambda\sum_{k=0}^{N-1}(z_{k}^{[j]})^{*}[-\mathcal{J}S_{k}^{[1]}+G_{k}(\lambda,\bar{\lambda})]z_{k}^{[j]}$$

and for $\lambda = \lambda_0 + ih$, $h \in \mathbb{R}$, similarly as in [2, 6],

$$\lim_{h\to 0}\frac{1}{2ih}\left(1-\left|\rho_j(\lambda_0+ih)\right|^2\right)=-\frac{\rho_j'(\lambda_0)}{\rho_j(\lambda_0)}=-\frac{\rho_j'(\lambda_0)}{\rho_0}.$$

The previous parts of our paper show that if (2) is of definite type, *i.e.*, (24) holds, then we have

$$-i\frac{\rho_j'(\lambda)}{\rho_j(\lambda)}\Big|_{\lambda=\lambda_0} = \frac{d}{d\lambda}\arg\rho(\lambda)|_{\lambda=\lambda_0} < 0$$

for the first kind multipliers and $\frac{d}{d\lambda} \arg \rho(\lambda) > 0$ for the second kind ones. This means that the multipliers resulting from splitting of first kind multipliers move clockwise and second kind multipliers move counterclockwise with increasing $\lambda \in \mathbb{R}$, both remain on the unit circle in some neighborhood of λ_0 .

Now, let us take $\lambda_0 = 0$, which is surely the stability point. Then, under assumptions of Theorem 1, system (2) is stable in some neighborhood of $\lambda_0 = 0$. The multipliers move around the unit circle as described above and may exit the unit circle only when multipliers of different kind meet on the unit circle. However, this may happen only when one of the multipliers comes through the point [-1,0] on the unit circle, *i.e.*, $\rho(\lambda) = -1$ is the eigenvalue of $U_N(\lambda)$, which is the same as that the boundary value problem

$$z_{k+1} = S_k(\lambda) z_k, \qquad z_N + z_0 = 0$$

has a solution.

Here we have, similarly to [7], the different situation in comparison with [6], where the symplectic system corresponding to the second-order system (12) is considered; see the part (i) of Section 2. Since in our case $U_N(\lambda)$ is no longer a polynomial, generally it can happen that the equation $\det(U_N(\lambda) + I) = 0$ has no root, and hence (2) is stable for every λ .

Let us comment shortly on assumption (24). In the Hamiltonian setting treated in the part (ii) of Section 2, we have

$$\mathcal{A}_k(\bar{\lambda},\lambda) = D_k^*(\lambda)H_kD_k(\lambda),$$

where *H* is given by (23) and $D(\lambda)$ by (17). Hence, (24) is satisfied for all $\lambda \in \mathbb{C}$, not only for those with $|\operatorname{Im} \lambda|$ sufficiently small. Concerning the exponential case of Section 2(iii), we have (compare (19))

$$\mathcal{A}_k(\bar{\lambda}, \lambda) = -\mathcal{J}R_k + o(\operatorname{Im}\lambda) \text{ as } \operatorname{Im}\lambda \to 0.$$

Hence, (24) is satisfied if

$$-\mathcal{J}R_k \ge 0$$
, $k = 0, \dots, N-1$ and $-\sum_{k=0}^{N-1} \mathcal{J}R_k > 0$.

This condition is in agreement with assumptions on the matrices H_k in the Hamiltonian setting. In the general setting, (24) is satisfied if

$$-\mathcal{J}S_k^{[1]} \ge 0, \quad k = 0, \dots, N-1, \qquad -\sum_{k=0}^{N-1} \mathcal{J}S_k^{[1]} > 0$$

and

$$G_k(\bar{\lambda}, \lambda) = o(\operatorname{Im} \lambda)$$
 as $\operatorname{Im} \lambda \to 0, k = 0, \dots, N-1$.

Finally, note that we have made only the first step towards the elaboration of a consistent stability theory of periodic symplectic difference systems in our paper. For example, we have not been able yet to establish a 'symplectic version' of the perturbation principle which enables to prove the statement of Theorem 1 without the additional assumption that the eigenvalues of the matrix $S^{[1]}$ are distinct. This is one of the problems which are subject of the present investigation, and we hope to solve it, together with other open problems, in subsequent papers.

Competing interests

The author declares that they have no competing interests.

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