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# Research Article

# **Eigenvalues of Vectorial Sturm-Liouville Problems with Parameter Dependent Boundary Conditions**

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We generalize the *regularized sampling method* introduced in 2005 by the author to compute the eigenvalues of scalar Sturm-Liouville problems (SLPs) to the case of vectorial SLP with parameter dependent boundary conditions. A few problems are worked out to illustrate the effectiveness of the method and show by the same token that we have indeed a general method capable of handling with ease very broad classes of SLPs, whether scalar or vectorial.

#### 1. Introduction

In [1] we introduced the regularized sampling method, a method to compute the eigenvalues of scalar Sturm-Liouville problems (SLPs) with parameter dependent boundary conditions. We subsequently used this method to compute the eigenvalues of singular and non-self-adjoint Sturm-Liouville problems. The scope of the method was further extended to include the computation of the eigenvalues of discontinuous/impulsive, nonlocal ([2] and the references therein), and two-parameter SLPs [3]. Continuing our effort we will tackle in this paper vectorial SLP with parameter dependent non-separated boundary conditions. Vectorial Sturm-Liouville problems have been considered in [4–13] and the references therein while corresponding inverse problems appeared in [14–17].

#### 2. The Characteristic Function

Consider the vectorial Sturm-Liouville problem,

$$-y'' + Q(x) y = \mu^{2} y, \quad 0 < x < 1$$

$$A(\mu) \begin{pmatrix} y(0, \mu) \\ y'(0, \mu) \end{pmatrix} + B(\mu) \begin{pmatrix} y(1, \mu) \\ y'(1, \mu) \end{pmatrix} = 0,$$
(1)

where *Q* is an  $n \times n$  matrix function, *A* and *B* are real  $2n \times 2n$  matrix functions of the parameter  $\mu$  such that the matrix  $[A(\mu) \mid B(\mu)]$  has full rank.

Let  $Y_c$ ,  $Y_s$  be the solutions of the Sturm-Liouville matrix equation  $-Y'' + Q(x)Y = \mu^2 Y$  subject to the initial conditions  $Y(0, \mu) = I$ ,  $Y'(0, \mu) = 0$  and  $Y(0, \mu) = 0$ ,  $Y'(0, \mu) = I$ , respectively, I being the  $n \times n$  identity matrix and 0 being the  $n \times n$  zero matrix.

The general solution of the Sturm-Liouville equation

$$-Y'' + O(x)Y = u^{2}Y$$
 (2)

is given by  $Y = Y_c a + Y_s b$  with arbitrary constant vectors a and b. Replacing in the boundary conditions, we get

$$A(\mu) \begin{pmatrix} a \\ b \end{pmatrix} + B(\mu) \begin{pmatrix} Y_c(1,\mu) a + Y_s(1,\mu) b \\ Y'_c(1,\mu) a + Y'_s(1,\mu) b \end{pmatrix} = 0,$$

$$\left\{ A(\mu) + B(\mu) \begin{pmatrix} Y_c(1,\mu) & Y_s(1,\mu) \\ Y'_c(1,\mu) & Y'_s(1,\mu) \end{pmatrix} \right\} \begin{pmatrix} a \\ b \end{pmatrix} = 0.$$
(3)

To have a nontrivial solution  $\binom{a}{b}$  a necessary and sufficient condition is that  $F(\mu) = 0$  where the characteristic function is

$$F(\mu) = \det \{ F \operatorname{mat}(\mu) \}, \tag{4}$$

where

$$F \operatorname{mat}(\mu) = A(\mu) + B(\mu) \begin{pmatrix} Y_c(1,\mu) & Y_s(1,\mu) \\ Y'_c(1,\mu) & Y'_s(1,\mu) \end{pmatrix}.$$
 (5)

The eigenvalues of (1) are the square of the zeroes of F. It is well known that the multiplicities of these eigenvalues are at most n.

## 3. Main Results

Let  $PW_{\sigma}$  be the Paley-Wiener space

$$PW_{\sigma} = \left\{ f \text{ entire, } \left| f\left(\mu\right) \right| \le Ce^{\sigma |\text{Im }\mu|}, \int_{-\infty}^{\infty} \left| f\left(\mu\right) \right|^2 d\mu < \infty \right\}, \tag{6}$$

and recall the celebrated Whittaker-Shannon-Kotel'nikov theorem [18].

**Theorem 1.** Let  $f \in PW_{\sigma}$ ; then

$$f(\mu) = \sum_{k=-\infty}^{\infty} f\left(\frac{k\pi}{\sigma}\right) \frac{\sin\sigma(\mu - k)}{\sigma(\mu - k)},\tag{7}$$

where the series converges uniformly on compact subset of C and in  $L^2(R)$ .

It is known that, in the case of scalar Sturm-Liouville problems,  $y(x,\mu)$  is an entire function of  $\mu$  for each fixed  $x \in (0,1]$ .  $y(x,\mu)$  is in a Paley-Wiener space as a function of  $\mu$  for each x only in the Dirichlet case. So, we had to subtract some terms from  $y(x,\mu)$  to make the difference fall in an appropriate PW $_\sigma$  space. We had even to subtract terms involving multiple integrals to get sharper results when it comes to computing of the eigenvalues. The regularized sampling method has been introduced recently [1] to overcome this problem; we do not have to subtract any term involving any (multiple) integration. In fact we multiplied  $y(x,\mu) - \phi(x,\mu)$  and  $y'(x,\mu) - \psi(x,\mu)$  by an appropriate function of  $\mu$  and got the eigenvalues with much greater precision at a reduced cost. Here  $\phi$  and  $\psi$  are known simple functions.

For the vectorial Sturm-Liouville problem at hand, we will use the regularized sampling method to recover the matrices  $Y_c(1,\mu), Y_c'(1,\mu), Y_s(1,\mu)$ , and  $Y_s'(1,\mu)$  from which we obtain  $F(\mu)$ , the characteristic function whose zeroes are the square roots of the sought eigenvalues of the problem.

Consider the compatible vector and matrix norms given by

$$||Y|| = \max_{i=1,\dots,n} |Y_i|, \qquad ||P|| = \max_{i=1,\dots,n} \sum_{j=1}^{n} |P_{ij}|, \qquad (8)$$

where  $Y \in \mathbb{R}^n$ ,  $P \in \mathbb{R}^{n \times n}$ . In the following we will make use of the standard estimate.

Lemma 2. Consider

$$\left|\cos u\right| \le e^{\left|\operatorname{Im} u\right|}, \qquad \left|\frac{\sin u}{u}\right| \le \frac{\gamma_0}{1+\left|u\right|}e^{\left|\operatorname{Im} u\right|}, \qquad (9)$$

where  $\gamma_0$  is some constant (we may take  $\gamma_0 = 1.72$ ).

To cover both cases  $(Y(0, \mu) = I, Y'(0, \mu) = 0$  and  $Y(0, \mu) = 0, Y'(0, \mu) = I)$  we will consider the following initial value problem:

$$Y'' + \mu^{2}Y = Q(x)Y,$$
 (10)  
 $Y(0,\mu) = E_{1}, \qquad Y'(0,\mu) = E_{2},$ 

where  $E_1$  and  $E_2$  are  $n \times n$  matrices or n-vector. We have

$$Y(x,\mu) = E_1 \cos \mu x + E_2 \frac{\sin \mu x}{\mu} + \int_0^x \frac{\sin \mu (x-t)}{\mu} Q(t) Y(t,\mu) dt.$$
 (11)

Our first result is the following theorem.

**Theorem 3.**  $Y(x, \mu)$  is an entire matrix function of  $\mu$  for each fixed  $x \in (0, 1]$  and satisfies the growth conditions

$$||Y(x,\mu)|| \le \left(\left\{||E_{1}|| + \frac{\gamma_{0}}{1+|\mu|}||E_{2}||\right\} e^{\gamma_{0} \int_{0}^{1}||Q(t)||dt}\right) e^{x|\operatorname{Im}\mu|}$$

$$\le \gamma_{1}e^{x|\operatorname{Im}\mu|},$$

$$||Y(x,\mu) - \left\{E_{1}\cos\mu x + E_{2}\frac{\sin\mu x}{\mu}\right\}||$$

$$\le \frac{\gamma_{2}}{1+|\mu|}e^{x|\operatorname{Im}\mu|} \le \gamma_{2}e^{x|\operatorname{Im}\mu|},$$

$$||Y'(x,\mu) - \left\{-\mu E_{1}\sin\mu x + E_{2}\cos\mu x\right\}|| \le \gamma_{3}e^{x|\operatorname{Im}\mu|},$$

$$||Y'(x,\mu) + \mu E_{1}\sin\mu x - E_{2}\cos\mu x$$

$$-\int_{0}^{x}\cos\mu(x-t)Q(t)\left(E_{1}\cos\mu t + E_{2}\frac{\sin\mu t}{\mu}\right)dt||$$

$$\le \frac{\gamma_{4}}{1+|\mu|}e^{x|\operatorname{Im}\mu|},$$
(12)

for some positive constants  $\gamma_1$ ,  $\gamma_2$ ,  $\gamma_3$ , and  $\gamma_4$ .

*Proof.* From (11) and using standard arguments, we conclude that  $Y(x, \mu)$  is an entire matrix function of  $\mu$  for each x in (0, 1]. Its derivative with respect to x,

$$Y'(x,\mu) = -\mu E_1 \sin \mu x + E_2 \cos \mu x$$

$$+ \int_0^x \cos \mu (x-t) Q(t) Y(t,\mu) dt,$$
(13)

is also an entire matrix function of  $\mu$  for each x in (0, 1]. Going back to (11) we get at once

$$\|Y(x,\mu)\| \leq \|E_{1}\cos\mu x + E_{2}\frac{\sin\mu x}{\mu}\|$$

$$+ \int_{0}^{x} \left|\frac{\sin\mu (x-t)}{\mu (x-t)}\right|$$

$$\cdot (x-t) \|Q(t)\| \cdot \|Y(t,\mu)\| dt$$

$$\leq e^{x|\operatorname{Im}\mu|} \left\{ \|E_{1}\| + \frac{\gamma_{0}x}{1+x|\mu|} \|E_{2}\| \right\}$$

$$+ \int_{0}^{x} \gamma_{0} (x-t) e^{(x-t)|\operatorname{Im}\mu|} \|Q(t)\| \cdot \|Y(t,\mu)\| dt$$

$$\leq e^{x|\operatorname{Im}\mu|} \left\{ \|E_{1}\| + \frac{\gamma_{0}}{1+|\mu|} \|E_{2}\| \right\}$$

$$+ e^{x|\operatorname{Im}\mu|} \int_{0}^{x} \gamma_{0} \|Q(t)\| \cdot e^{t|\operatorname{Im}\mu|} \|Y(t,\mu)\| dt.$$
(14)

Multiplying by  $e^{-x|\operatorname{Im}\mu|}$ , using Gronwall's lemma, and multiplying back by  $e^{x|\operatorname{Im}\mu|}$  we get

$$||Y(x,\mu)|| \le \left( \left\{ ||E_1|| + \frac{\gamma_0}{1+|\mu|} ||E_2|| \right\} e^{\gamma_0 \int_0^1 ||Q(t)|| dt} \right) e^{x|\operatorname{Im} \mu|}$$

$$\le \gamma_1 e^{x|\operatorname{Im} \mu|},$$
(15)

where  $\gamma_1=\{\|E_1\|+\gamma_0\|E_2\|\}\exp(\gamma_0\int_0^1\|Q(t)\|dt)$ . Now, using the above estimate in (10), we get

$$\left\| Y(x,\mu) - \left\{ E_{1} \cos \mu x + E_{2} \frac{\sin \mu x}{\mu} \right\} \right\| \\
\leq \int_{0}^{x} \left| \frac{\sin \mu (x-t)}{\mu (x-t)} \right| \cdot (x-t) \|Q(t)\| \cdot \|Y(t,\mu)\| dt \\
\leq \int_{0}^{x} \frac{\gamma_{0} e^{(x-t)|\operatorname{Im}\mu|}}{1 + |\mu| (x-t)} \cdot (x-t) \|Q(t)\| \gamma_{1} e^{t|\operatorname{Im}\mu|} dt \qquad (16) \\
\leq e^{x|\operatorname{Im}\mu|} \frac{\gamma_{0} \gamma_{1}}{1 + |\mu|} \int_{0}^{1} \|Q(t)\| dt \\
\leq \frac{\gamma_{2}}{1 + |\mu|} e^{x|\operatorname{Im}\mu|} \leq \gamma_{2} e^{x|\operatorname{Im}\mu|},$$

where  $\gamma_2 = \gamma_0 \gamma_1 \int_0^1 ||Q(t)|| dt$ . Likewise we have

$$\|Y'(x,\mu) - \{-\mu E_1 \sin \mu x + E_2 \cos \mu x\}\|$$

$$\leq \int_0^x |\cos \mu (x-t)| \cdot \|Q(t)\| \cdot \|Y(t,\mu)\| dt$$

$$\leq \int_{0}^{x} e^{(x-t)|\operatorname{Im}\mu|} \|Q(t)\| \gamma_{1} e^{t|\operatorname{Im}\mu|} dt$$

$$= e^{x|\operatorname{Im}\mu|} \gamma_{1} \int_{0}^{1} \|Q(t)\| dt = \gamma_{3} e^{x|\operatorname{Im}\mu|},$$
(17)

where  $\gamma_3 = \gamma_1 \int_0^1 \|Q(t)\| dt$ . As in the scalar case,  $Y(x, \mu)$  is in a Paley-Wiener space only in the Dirichlet case; however,  $Y(x, \mu) - \{E_1 \cos \mu x + E_2 (\sin \mu x/\mu)\}$  is. As for  $Y'(x, \mu)$ , it is not; nor is  $Y'(x, \mu) - \{-\mu E_1 \sin \mu x + E_2 \cos \mu x\}$  since they are not square integrable over the reals for fixed x in (0, 1]. Also,

$$\|Y'(x,\mu) + \mu E_{1} \sin \mu x - E_{2} \cos \mu x$$

$$- \int_{0}^{x} \cos \mu (x-t) Q(t) \left( E_{1} \cos \mu t + E_{2} \frac{\sin \mu t}{\mu} \right) dt \|$$

$$\leq \int_{0}^{x} |\cos \mu (x-t)| \cdot \|Q(t)\|$$

$$\cdot \|Y(t,\mu) - \left\{ E_{1} \cos \mu t + E_{2} \frac{\sin \mu t}{\mu} \right\} \|dt$$

$$\leq \int_{0}^{x} e^{(x-t)|\operatorname{Im} \mu|} \|Q(t)\| \frac{\gamma_{2}}{1+|\mu|} e^{t|\operatorname{Im} \mu|} dt$$

$$= e^{x|\operatorname{Im} \mu|} \frac{\gamma_{2}}{1+|\mu|} \int_{0}^{1} \|Q(t)\| dt = \frac{\gamma_{4}}{1+|\mu|} e^{x|\operatorname{Im} \mu|},$$
where  $\gamma_{4} = \gamma_{2} \int_{0}^{1} \|Q(t)\| dt$ .

We get at once the following corollaries.

**Corollary 4.**  $Y_c(x, \mu)$ ,  $Y'_c(x, \mu)$ ,  $Y_s(x, \mu)$ ,  $Y'_s(x, \mu)$ ,  $Y_0(x, \mu)$ , and  $Y'_0(x, \mu)$  are entire matrix functions of  $\mu$  for each fixed  $x \in (0, 1]$  and satisfy the growth conditions

$$\begin{aligned} \|Y_{c}(x,\mu)\| &\leq \left(e^{\gamma_{0} \int_{0}^{1} \|Q(t)\| dt}\right) e^{x|\operatorname{Im} \mu|}, \\ \|Y_{c}(x,\mu) - I \cos \mu x\| &\leq \frac{\gamma_{2}}{1 + |\mu|} e^{x|\operatorname{Im} \mu|} \leq \gamma_{2} e^{x|\operatorname{Im} \mu|}, \\ \|Y'_{c}(x,\mu) + \mu I \sin \mu x\| &\leq \gamma_{3} e^{x|\operatorname{Im} \mu|}, \\ \|Y'_{c}(x,\mu) - \left\{-\mu I \sin \mu x + \int_{0}^{x} \cos \mu (x - t) Q(t) \cos \mu t \, dt\right\} \| \\ &\leq \frac{\gamma_{4}}{1 + |\mu|} e^{x|\operatorname{Im} \mu|}, \\ \|Y_{s}(x,\mu)\| &\leq \left(\frac{\gamma_{0}}{1 + |\mu|} e^{\gamma_{0} \int_{0}^{1} \|Q(t)\| dt}\right) e^{x|\operatorname{Im} \mu|} \leq \gamma_{1} e^{x|\operatorname{Im} \mu|}, \\ \|Y_{s}(x,\mu) - I \frac{\sin \mu x}{\mu} \| &\leq \frac{\gamma_{2}}{1 + |\mu|} e^{x|\operatorname{Im} \mu|} \leq \gamma_{2} e^{x|\operatorname{Im} \mu|}, \\ \|Y'_{s}(x,\mu) - I \cos \mu x\| &\leq \gamma_{3} e^{x|\operatorname{Im} \mu|}, \end{aligned}$$

$$\left\| Y_{s}'(x,\mu) - \left\{ I \cos \mu x + \int_{0}^{x} \cos \mu (x-t) Q(t) \frac{\sin \mu t}{\mu} dt \right\} \right\| \\
\leq \frac{\gamma_{4}}{1+|\mu|} e^{x|\operatorname{Im}\mu|}, \\
\left\| Y_{0}(x,\mu) \right\| \leq \left( \left\{ \|D_{1}\| + \frac{\gamma_{0}}{1+|\mu|} \|D_{2}\| \right\} e^{\gamma_{0} \int_{0}^{1} \|Q(t)\| dt} \right) e^{x|\operatorname{Im}\mu|} \\
\leq \gamma_{1} e^{x|\operatorname{Im}\mu|}, \\
\left\| Y_{0}(x,\mu) - \left\{ D_{1} \cos \mu x + D_{2} \frac{\sin \mu x}{\mu} \right\} \right\| \\
\leq \frac{\gamma_{2}}{1+|\mu|} e^{x|\operatorname{Im}\mu|} \leq \gamma_{2} e^{x|\operatorname{Im}\mu|}, \\
\left\| Y_{0}'(x,\mu) - \left\{ -\mu D_{1} \sin \mu x + D_{2} \cos \mu x \right\} \right\| \leq \gamma_{3} e^{x|\operatorname{Im}\mu|}, \\
\left\| Y_{0}'(x,\mu) + \mu D_{1} \sin \mu x - D_{2} \cos \mu x - \int_{0}^{x} \cos \mu (x-t) Q(t) \left( D_{1} \cos \mu t + D_{2} \frac{\sin \mu t}{\mu} \right) dt \right\| \\
\leq \frac{\gamma_{4}}{1+|\mu|} e^{x|\operatorname{Im}\mu|}, \tag{19}$$

for some generic positive constants  $\gamma_1$ ,  $\gamma_2$ ,  $\gamma_3$ , and  $\gamma_4$ .

Corollary 5. The functions,

$$Y_{c}(1,\mu) - I\cos\mu,$$

$$Y'_{c}(1,\mu) - \left\{-\mu I\sin\mu + \int_{0}^{1}\cos\mu (1-t)Q(t)\cos\mu t dt\right\},$$

$$Y_{s}(1,\mu) - I\frac{\sin\mu}{\mu},$$

$$Y'_{s}(1,\mu) - \left\{I\cos\mu + \int_{0}^{1}\cos\mu (1-t)Q(t)\frac{\sin\mu t}{\mu} dt\right\},$$

$$Y'_{0}(1,\mu) - \left\{D_{1}\cos\mu + D_{2}\frac{\sin\mu}{\mu}\right\},$$

$$Y''_{0}(1,\mu)$$

$$- \left\{-\mu D_{1}\sin\mu + D_{2}\cos\mu + D_{2}\frac{\sin\mu t}{\mu}\right\},$$

$$(20)$$

belong to the Paley-Wiener space  $P\dot{W}_1$  as functions of  $\mu$  and thus can be recovered from their samples at  $\mu_k = k\pi$ ,  $k \in Z$  using the WSK series.

**Theorem 6.** Let  $\theta$  be positive real number and  $m \ge 2$  a positive integer. Consider

$$U_{1}(x,\mu)$$

$$= \left(Y(x,\mu) - \left\{E_{1}\cos\mu x + E_{2}\frac{\sin\mu x}{\mu}\right\}\right) \left(\frac{\sin(\theta\mu)}{\theta\mu}\right)^{m},$$

$$U_{2}(x,\mu)$$

$$= \left(Y'(x,\mu) - \left\{-\mu E_{1}\sin\mu x + E_{2}\cos\mu x\right\}\right) \left(\frac{\sin(\theta\mu)}{\theta\mu}\right)^{m}$$
(21)

belong to the Paley-Wiener space  $PW_{\sigma}$  where  $\sigma = x + m\theta$  as functions of  $\mu$  for each fixed  $x \in (0,1]$  for  $m \geq 1$  and satisfy the growth condition  $\|U_1(x,\mu)\|$ ,  $\|U_2(x,\mu)\| \leq \gamma_5 e^{\sigma|\text{Im }\mu|}/(1+\theta|\mu|)^m$  where  $\gamma_5$  is some positive constant  $(\gamma_5 = \gamma_0^m \max(\gamma_2, \gamma_3))$ .

*Proof.* It is enough to note that  $\sin(\theta \mu)/\theta \mu$  is an entire function of  $\mu$  and satisfies the estimate in the above Lemma and the fact that  $Z(x, \mu)$  is the product of two entire functions thus entire.

*Remark 7.* To avoid the first singularity of  $(\sin(\theta \mu)/\theta \mu)^{-m}$  we will take  $\theta < (N-m)^{-1}$ .

The use of the WSK theorem allows us to recover  $U_1(1,\mu)$  and  $U_2(1,\mu)$  as

$$U_{1}(1,\mu) = \sum_{k \in \mathbb{Z}} \alpha_{k} \frac{\sin \sigma (\mu - \mu_{k})}{\sigma (\mu - \mu_{k})},$$

$$U_{2}(1,\mu) = \sum_{k \in \mathbb{Z}} \beta_{k} \frac{\sin \sigma (\mu - \mu_{k})}{\sigma (\mu - \mu_{k})},$$
(22)

where  $\alpha_k = U_1(1, \mu_k)$ ,  $\beta_k = U_2(1, \mu_k)$ ,  $\mu_k = k\pi/\sigma$ , and  $\sigma = 1 + m\theta$ .

Hence,  $Y(1, \mu)$  or  $Y'(1, \mu)$  can be recovered as

$$Y(1,\mu) = E_1 \cos \mu + E_2 \frac{\sin \mu}{\mu} + \left(\frac{\sin(\theta\mu)}{\theta\mu}\right)^{-m} \sum_{k \in \mathbb{Z}} \alpha_k \frac{\sin \sigma(\mu - \mu_k)}{\sigma(\mu - \mu_k)},$$

$$Y'(1,\mu) = -\mu E_1 \sin \mu + E_2 \cos \mu + \left(\frac{\sin(\theta\mu)}{\theta\mu}\right)^{-m} \sum_{k \in \mathbb{Z}} \alpha_k \frac{\sin \sigma(\mu - \mu_k)}{\sigma(\mu - \mu_k)}.$$
(23)

In practice, we take  $|k| \le N$  for some positive integer N, large enough, so that  $F(\mu)$  can be reconstructed whose zeros are the square roots of the sought eigenvalues.

Since  $\mu^{m-1}U_1(1,\mu)$  and  $\mu^{m-1}U_2(1,\mu)$  are in  $L^2(-\infty,\infty)$ , Jagerman's result [18] is applicable and yields the following better estimate.

**Lemma 8** (truncation error). Let  $U_j^N(1,\mu) = \sum_{k=-N}^N U_j(1,\mu_k)(\sin\sigma(\mu-\mu_k)/\sigma(\mu-\mu_k))$  denote the truncation of  $U_j(1,\mu)$ , j=1,2. Then, for  $|\mu| < N\pi/\sigma$ ,

$$\left| U_{j} \left( 1, \mu \right) - U_{j}^{N} \left( 1, \mu \right) \right| \\
\leq \frac{\left| \sin \gamma \mu \right| \gamma_{5,j}}{\pi \left( \pi / \sigma \right)^{m-1} \sqrt{1 - 4^{-m+1}}} \\
\cdot \left[ \frac{1}{\sqrt{(N\pi/\sigma) - \mu}} + \frac{1}{\sqrt{(N\pi/\sigma) + \mu}} \right] \frac{1}{(N+1)^{m-1}}, \tag{24}$$

where  $\gamma_{5,j} = \|\mu^{m-1}U_j(1,\mu)\|_2$ .

**Lemma 9.** Consider  $|\mu| < N\pi/\sigma$ ,

$$\left|Y'(1,\mu) - Y_N(1,\mu)\right|,$$

$$\left|Y'(1,\mu) - Y_N'(1,\mu)\right|$$

$$\leq \left|\frac{\sin(\theta\mu)}{\theta\mu}\right|^{-m} \times \frac{\left|\sin\gamma\mu\right|\gamma_5}{\pi\left(\pi/\sigma\right)^{m-1}\sqrt{1-4^{-m+1}}}$$

$$\cdot \left[\frac{1}{\sqrt{(N\pi/\sigma)-\mu}} + \frac{1}{\sqrt{(N\pi/\sigma)+\mu}}\right] \frac{1}{(N+1)^{m-1}},$$
(25)

where  $\gamma_5 = \max\{\|\mu^{m-1}Y(1,\mu)\|_2, \|\mu^{m-1}Y'(1,\mu)\|_2\}.$ 

The approximation of  $Y(1,\mu)$  and  $Y'(1,\mu)$  by  $Y_N(1,\mu)$  and  $Y'_N(1,\mu)$ , respectively, induces an approximation of the characteristic function F by  $F_N$ , whose zeros are the square root of the eigenvalues of the problem.

Let  $\overline{\mu}^2$  denote an eigenvalue of the problem; then independent eigenfunctions associated can be obtained using basis vectors of the null space of the matrix Fmat( $\overline{\mu}$ ) as initial conditions to the differential equation  $-y'' + Q(x)y = \overline{\mu}^2 y$ , 0 < x < 1.

#### 4. Numerical Examples

In this section we will illustrate the power of the regularized sampling method as applied to vectorial Sturm-Liouville problems with parameter dependent boundary conditions. We will take m=6, N=40 and a precision of  $10^{-20}$  for the first three examples involving two dimensional SLPs. We will also work out two three-dimensional SLPs one of them involving parameter dependent boundary conditions. In these last two examples we take different values of N, namely, N=20,40,60,80, and 100, and take m=4 and a precision of  $10^{-20}$ . The reported multiplicities of the eigenvalues  $\overline{\mu}^2$  are just the dimensions of the null space of the corresponding matrices Fmat( $\overline{\mu}$ ).

Example 1 (Chanane [1], 1D-version taken from fom Binding and Browne [19]). Consider

$$-y_1''(x) = \lambda y_1(x), \quad -y_2''(x) = \lambda y_2(x), \quad 0 \le x \le 1$$

$$y_1(0) + (\lambda + d) y_1'(0) = 0, \quad y_2(0) + (\lambda + d) y_2'(0) = 0$$

$$y_1(1) - \lambda y_1'(1) = 0, \quad y_2(1) - \lambda y_2'(1) = 0,$$
(26)

where  $d=-4\pi^2$ . The first three eigenvalues were obtained as 9.730886578213082033, 88.76331625258976337, and 157.88411043863472059 putting them at about  $10^{-18}$  from the exact eigenvalues. All these are double eigenvalues.

Example 2. Consider

$$-z_1'' + xz_1 = \mu^2 z_1, \quad -z_2'' + x^2 z_2 = \mu^2 z_2, \quad 0 < x < 1$$

$$z_1(0) = 0, \quad z_2(0) = 0$$

$$z_1(1) = 0, \quad z_2(1) = 0.$$
(27)

The first four eigenvalues were obtained as 10.149980317596192645, 39.426774741845613693, 88.33456043776637171, and 157.20995003768636950. Their multiplicity is two. Figure 1 illustrates the graph of the characteristic function.

In the next example we change the boundary conditions in Example 2 to a parameter dependent one.

Example 3. Consider

$$-z_{1}'' + xz_{1} = \mu^{2}z_{1}, \quad -z_{2}'' + x^{2}z_{2} = \mu^{2}z_{2}, \quad 0 < x < 1$$

$$z_{1}(0) = 0, \quad z_{2}(0) = 0$$

$$z_{1}(1) + \mu z_{2}(1) + \mu^{2}z_{1}'(1) = 0$$

$$\mu^{2}z_{1}(1) + z_{2}(1) + z_{1}'(1) + z_{2}'(1) = 0.$$
(28)

The first ten eigenvalues were obtained as 1.8774711215942920040, 5.743710132061151193, 19.343498090220217814, 27.524136648884737570, 57.53122978141141088, 68.17260384634582088, 115.66805223522017586, 128.14561531120321356, 193.70672667853536264, and 207.68609472970380538. All these are simple eigenvalues. Figure 2 illustrates the graph of the characteristic function.

Next we consider three-dimensional vectorial SLPs, with different boundary conditions.

Example 4. Consider

$$-z_{1}'' + x^{2}z_{1} = \mu^{2}z_{1}, \qquad -z_{2}'' + \frac{3x}{2}z_{2} - \frac{x}{2}z_{3} = \mu^{2}z_{2},$$

$$-z_{3}'' - \frac{x}{2}z_{2} + \frac{3x}{2}z_{3} = \mu^{2}z_{3}, \quad 0 < x < 1$$

$$z_{1}(0) = 0, \quad z_{2}(0) = 0, \quad z_{3}(0) = 0$$

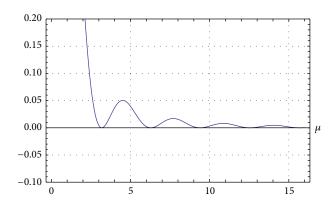
$$z_{1}(1) = 0, \quad z_{2}(1) = 0, \quad z_{3}(1) = 0.$$
(29)

6.32072912835397407977

$\mu_1$	$\mu_2$	$\mu_3$	$\mu_4$
3.18255001139908084139	3.24667169525908938767	6.30074330869793149011	6.32071003783954269867
3.18255272015329488134	3.24667965370940216187	6.30075208373683858194	6.32073018962762953754
3.18255257457282246273	3.24667922487683254112	6.30075161400550151819	6.32072910805919989533
3.18255257724071613616	3.24667923278762738065	6.30075162260634961554	6.32072912799411873146

Table 1:  $\mu$  as a function of N for Example 4.

3.24667923293045271115



3.18255257728954451935

FIGURE 1: *F* for Example 2.

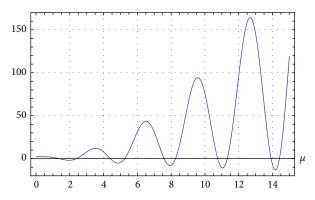
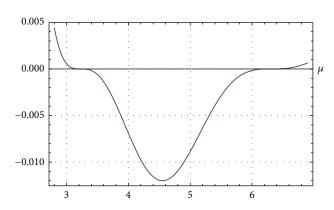


FIGURE 2: *F* for Example 3.

Here, we will take m=4, N=20, 40, 60, 80, and 100, and a precision of  $10^{-20}$ . Figure 3 illustrates the characteristic function  $F_N$  over the range [2.8, 6.9], while Figures 4 and 5 zoom into the regions containing the eigenvalues. Note that, in the range of interest [2.8, 6.9], the graphs of  $F_N$ , N=20, 40, 60, 80, and 100, are on the top of each other. A  $10^{-5}$  precision on  $\mu$  can be obtained with just N=40. It appears clearly that in this example we have a simple eigenvalue  $\lambda_1=\mu_1^2$  and a double eigenvalue  $\lambda_2=\mu_2^2$ , followed by a simple eigenvalue  $\lambda_3=\mu_3^2$  and a double eigenvalue  $\lambda_4=\mu_4^2$  (Figures 7, 8, 9, and 10). To obtain the double eigenvalues we look for the roots  $\overline{\mu}$  of  $F'(\mu)$  and then evaluate  $F(\overline{\mu})$  which happened to be in each case of the order of  $10^{-20}$ . Table 1 illustrates these  $\mu$  as function of N, the number of sampling points.



6.30075162276373835846

FIGURE 3: *F* for Example 4.

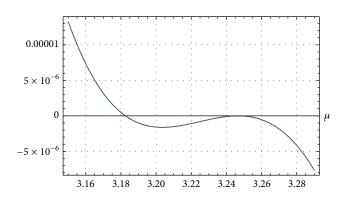


FIGURE 4: *F* around the first cluster of zeroes for Example 4.

The first few *alpha* coefficients in the cardinal series expansion of  $U_1(1, \mu)$  are given as follows:

$$\alpha_0 = \begin{pmatrix} 0.0507 & 0 & 0 \\ 0 & 0.130 & -0.0447 \\ 0 & -0.0447 & 0.130 \end{pmatrix},$$

$$\alpha_1 = -\alpha_{-1} = \begin{pmatrix} 0.0165 & 0 & 0 \\ 0 & 0.0451 & -0.0157 \\ 0 & -0.0157 & 0.0451 \end{pmatrix},$$

$$\alpha_2 = -\alpha_{-2} = \begin{pmatrix} -0.00455 & 0 & 0 \\ 0 & -0.0105 & 0.00352 \\ 0 & 0.00352 & -0.0105 \end{pmatrix},$$

$$\alpha_3 = -\alpha_{-3} = \begin{pmatrix} 0.00197 & 0 & 0 \\ 0 & 0.00442 & -0.00147 \\ 0 & -0.00147 & 0.00442 \end{pmatrix}.$$
(30)

The above data have been reported with only a few digits.

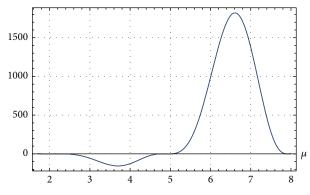


FIGURE 5: *F* for Example 5.

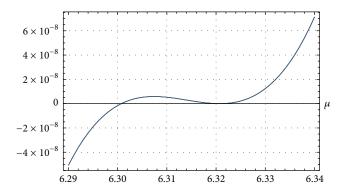


FIGURE 6: *F* around the second cluster of zeroes for Example 4.

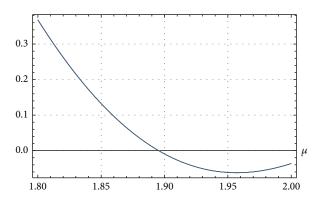
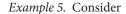


FIGURE 7: F around  $\mu_1$  in Example 5.



$$-z'' + Qz = \mu^2 z, \quad 0 < x < 1$$

$$z(0) = 0 \quad z(1) + Bz'(1) = 0,$$
(31)

where

$$Q = \begin{pmatrix} x^2 & 0 & 0 \\ 0 & \frac{3x}{2} & -\frac{x}{2} \\ 0 & -\frac{x}{2} & \frac{3x}{2} \end{pmatrix}, \qquad B = \begin{pmatrix} \mu^2 & \mu & 1 \\ \mu & \mu^2 & 1 \\ 1 & 1 & 1 \end{pmatrix}. \tag{32}$$

Here, we will take m = 4, N = 20, 40, 60, 80, and 100, and a precision of  $10^{-20}$ . Figure 6 illustrates the characteristic

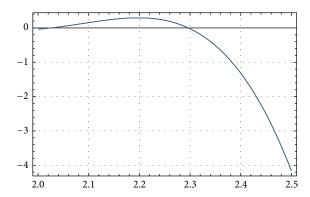


FIGURE 8: F around  $\mu_2$  and  $\mu_3$  in Example 5.

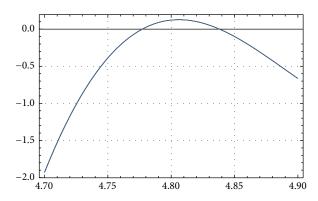


FIGURE 9: F around  $\mu_4$  and  $\mu_5$  in Example 5.

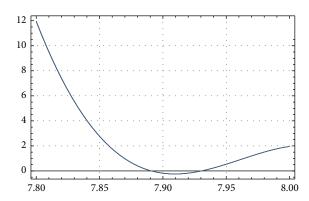


Figure 10: F around  $\mu_6$  and  $\mu_7$  in Example 5.

function  $F_N$  over the range [0,8]. In this range, the graphs of  $F_N$ , N=20,40,60,80, and 100, are on the top of each other. In this example the first seven (07) eigenvalues  $\lambda_k=\mu_k^2, k=1,\ldots,7$ , are all simple. Tables 2(a) and 2(b) illustrate these  $\mu$  as function of N, the number of sampling points.

#### 5. Conclusion

In this paper we have extended the domain of application of the regularized sampling method to the case of vectorial

Table 2:  $\mu$  as a function of N for Example 5.

(a)

N	$\mu_1$	$\mu_2$	$\mu_3$	$\mu_4$		
20	1.89539417301926842876	2.02558758968702183901	2.29660628213819048323	4.77718462334715123614		
40	1.89532715021073954015	2.02545891470303642172	2.29650089361152297043	4.77709199050137280895		
60	1.89531503182513465256	2.02543500687372306198	2.29648223596348958117	4.77707495489416773295		
80	1.89531699030377280216	2.02543884958899241306	2.29648526430437415553	4.77707769693838662656		
100	1.89531696279196052243	2.02543879569567539948	2.29648522170989172904	4.77707765846484979002		
(b)						
N	$\mu_5$		$\mu_6$	$\mu_7$		
20	4.83834197825672	007101 7.8909	9828962247999966	7.93058654329498644499		
40	4.83814507749939	589802 7.8909	0804120709347972	7.93039290552171137187		
60	4.83810790006947131246		9138476342607463	7.93035620848368882791		
80	4.83811385300866362529		9406045343050839	7.93036207289652547074		
100	4.83811376961267	7.8908	9402293178420005	7.93036199078631128916		

Sturm-Liouville problems with parameter dependent boundary conditions. We have presented the theoretical foundation of the method and worked out a few examples to illustrate the method and shown by the same token that we have indeed a general method capable of handling with ease very broad classes of SLPs, whether scalar or vectorial, and providing the results at a reduced cost.

#### **Conflict of Interests**

The author declares that there is no conflict of interests regarding the publication of this paper.

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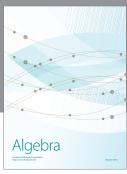
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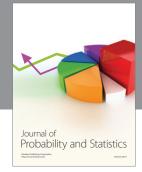
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