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The quest for a consistent signal in ground and GRACE gravity time series

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Abstract

Recent studies show that terrestrial and space based observations of gravity agree over Europe. In this paper, we compare time series of terrestrial gravity (including the contribution due to surface displacement) as measured by superconducting gravimeters (SGs), space based observations from GRACE, and predicted changes in gravity derived from two global hydrological models at 10 SG stations in Central Europe. Despite the fact that all observations and models observe a maximum in the same season due to water storage changes, there is little agreement between the SG time series even when they are separated by distances smaller than the spatial resolution of GRACE. We also demonstrate that GRACE and the SG observations and the water storage models do not display significant correlation at seasonal periods nor at inter-annual periods. These findings are consistent with the fact that the SGs are sensitive primarily to mass changes in the few hundred meters surrounding the station.

Keywords

Time variable gravity, hydrology, GRACE, superconducting gravimeters.

1. Introduction

The Earth is a coupled dynamic system with a climate component composed of the atmosphere, the oceans, the cryosphere and the continental hydrology. The sensitivity of contemporary geodetic techniques to the Earth system makes them a powerful and indispensable tool to monitor its dynamics. Nevertheless, the contribution of geodesy to understanding the Earth relies on the accuracy and quality of the data analysis. In particular, geodetic theory has to be improved to the extent that we can take full advantage of the data precision. For example, estimate the hydrological effects on terrestrial and space gravity measurements remains challenging, as subsurface water dynamics is very difficult to assess, at both local and global scales.

Separation of the couplings can be achieved by benefiting from the combination of multiple geodetic measurements and/or of the climate models. Various studies showed a fair consistency between GNSS, climate models, and GRACE (Gravity Recovery and Climate Experiment) data (Blewitt et al., 2001; Blewitt and Clarke, 2003; van Dam et al., 2007; Tregoning et al., 2009; Tesmer et al., 2011; Valty et al., 2013). Here, we evaluate the insights that can be obtained from a comparison/combination of terrestrial gravity measurements from superconducting gravimeters (SGs) in Central Europe with the equivalent gravity estimated from the GRACE solutions. Previous studies (e.g. Abe et al., 2012; Crossley et al., 2012; Neumeyer et al., 2006, 2008; Weise et al., 2009, 2011) claim they found a common behavior between the times series from the SGs, GRACE and hydrological models.

However, regarding the Newtonian effect of hydrological processes, SGs are sensitive primarily to mass included in the few hundred meters around the station (Creutzfeldt et al., 2008). So, one may have expected larger discrepancies between the SGs and GRACE solutions. To address this problem, we extend the previous study both in time - time series of our study extend up to 2012 - and in the number of SG used, and we test, using a different method, the robustness of the common signal.

2. Data

SG

The SG station locations are shown on the map of Figure 1, and their characteristics are described in Table 1. The time series are corrected for tidal effects using the parameter sets obtained from the tidal analysis of the hourly time series. This analysis was performed with the Eterna 3.4 package (Wenzel, 1996). The atmospheric influence is removed using the 3-D high resolution 3 hourly ECMWF model, assuming an inverted barometer hypothesis, as provided by J.-P. Boy (<u>http://loading.ustrasbg.fr/GGP/</u>) - for a review of the 3-D correction, see Crossley et al., 2013. The centrifugal effect associated with polar motion is also corrected (Wahr, 1985).

Removal of instrumental offsets is a critical step and is probably the most subjective part of the SG processing, as this depends on the operator (Hinderer et al., 2007). For all stations, the offsets are removed either visually, when the gap is not too long (typically, no more than a few hours), or, if the gap is longer, adjusting the SG series using co-located AG measurements when available. For the Pecny (PE), Moxa (MO) and Strasbourg (ST) stations, our processing was found consistent with the residuals provided by the operators; for the other stations the operators provided the data directly. The accumulated impact of remaining differences in the offsets is similar to a random walk process (Hinderer et al., 2007), and is included in the instrumental drift. For all series, after corrections, a second order polynomial was adjusted and subtracted to remove possible non-linear instrumental drift or other very long term geophysical effects, which are out of the scope of this study.

The SG time series used in this study are shown on Figure 2a and on Figure 2b after removing a composite seasonal cycle by means of a stacking technique (Hartmann and Michelsen, 1989). This tool allows removing the mean signal of period *T*. This is done on each SG series separately, by first computing the mean signal for a given phase ϕ by averaging all the value of the time series corresponding to this phase ($t = \phi$, $T + \phi$, $2T + \phi$, ...), then by removing it at every data point of this phase. At Wettzell, a change in the annual signal is observed after 2008, probably caused by major

construction works undertaken in 2009 and by the fact that the SG was moved by 250 m in October 2010.

Global hydrological models

We use hydrological loading effects provided by J.-P Boy (Boy and Hinderer, 2006; http://loading.ustrasbg.fr/GGP/), computed from the continental ground water content provided by the GLDAS/Noah model (Rodell et al., 2004) and ERA interim reanalysis (Uppala et al., 2005). Those datasets will be referred to as GLDAS and ERA, here after. The 6 hourly model based on ERA are interpolated to 3 hourly data to match the SG and GLDAS sampling. The space sampling of GLDAS is 0.25 degree and 0.7 degree for ERA interim (Boy, pers. Comm., 2012). The hydrology grids were decomposed into spherical harmonics, and then converted into ground gravity using the appropriate combination of load Love Number (e.g., Farrell, 1972). The Love numbers were calculated assuming PREM model (Dziewonski & Anderson, 1981) as Earth model.

GRACE

We use GRACE time gravity solutions from seven institutes, as summarized in Table 2:

- The release 5 of the three official solutions, NASA/CSR, NASA/JPL and GFZ groups (noted here as CSR, JPL and GFZ). These solutions are given without filtering but corrected for a dealiasing model for atmosphere and oceans (AOD de-aliasing products).
- Four other independent solutions: ITG monthly solution from Bonn university (Kurtenbach et al., 2009), AIUB monthly solution from Bern university (e.g. Beutler et al., 2010), DTM-1b monthly solution from Delft University of Technology (noted here DTM, Liu et al., 2010), and GRGS 10-days release 2 solution from the CNES French space agency (Bruinsma et al., 2009).

The GRGS and DTM solutions are already regularized using various methods (see above references and websites for more details). In the CSR, JPL, GFZ, ITG and AIUB series, striping noise has to be filtered out prior to investigations. We applied a correlated-error filter and a 500 km Gaussian smoothing based on the Swenson and Wahr (2006) method; this method was shown as the most precise in Valty et al. (2013). We found that AIUB solution presents an anomalously high degree 2 zonal coefficient. Since this coefficient is usually very small in surface gravity time variations (unlike the geoid that presents large J2 time variations), it has been suppressed from the AIUB solution prior to our computation.

To allow comparison between GRACE solutions and ground gravity measurements, which here are not corrected for the non-tidal ocean contribution, we also make a comparison with the three GRACE solutions where the non-tidal ocean contribution has been added back using dealiasing products, when provided by the analysis center (only GRGS, GFZ and ITG). A total of 10 GRACE solutions has consequently been used. As for the hydrology models, GRACE time variable gravity was decomposed into spherical harmonics, and then reconstructed at the SG station location as ground gravity values, using the appropriate combination of load Love numbers. Note that we did not use the classical formulation for gravity perturbation based on the loading gravimetric factor (see Farrel (1972) or Boy et al. (2002)), because it supposes that the load is above the gravimeter. Such assumption is valid for the atmosphere or oceans, but is less adapted for hydrology loading problems, where the load is generally under the sensor. We prefer the formulation used by Crossley et al. (2012), which is the derivative of the gravitational potential perturbation inferred from GRACE measurements plus a free air additional correction due to ground displacements. If we note $(\Delta C_{nm} \Delta S_{nm})$ the Stoke's coefficients of degree n and order m of the gravitational potential perturbation provided by GRACE, h_n the vertical displacement Love number, and k_n the potential perturbation Love number, then the gravimetric signal can be reconstructed as follows:

$$g(\theta,\lambda,t) = \frac{GM}{a^2} \sum_{n=2}^{N} \sum_{m=0}^{+n} P_n^m(\cos\theta) \left(n+1-2 \frac{h_n}{1+k_n}\right) (\Delta C_{nm}(t)\cos(m\lambda) + \Delta S_{nm}(t)\sin(m\lambda)),$$

where P_n^m are the associated Legendre polynomials, GM is Earth's standard gravitational parameter, a the semi-major axis of the ellipsoid, N the maximum degree, and (θ, λ, t) are colatitude, longitude and time.

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The 3-hourly time series of SG and hydrological models are decimated to 5 days; for GRACE, the original sampling rates were kept as provided by the different data centers, except in the case of the EOF analysis, where the GRACE series were linearly interpolated to 5 days to compare directly with the SG data. Scaling the series to the shortest sample interval avoids losing information. Before performing the different analyses, a second degree polynomial was systematically adjusted to all the SG, hydrological and GRACE series, in order to remove any possible bias that may be caused by non-linear slopes caused by SG instrumental drift or by residual long-period geophysical signals which are beyond the scope of this paper (Van Camp and Francis, 2006; Van Camp et al., 2010).

3. Common variability in the SG time series

As GRACE only sees large scale phenomenon, any GRACE/SG agreement would rely on common variability between the SG time series at large scale. A classical method to look for a common variability in time series is correlation study, as done by Neumeyer et al., 2008 and Abe et al., 2012. The correlation coefficients of the series are given in Table 3. However, the interpretation of the correlation coefficient rely on a statistical test which makes no sense when a strong periodic signal is present in the data, as all the data point corresponding to the same phase are not independent (see Von Storch and Zwiers (1999) for more detail on the assumption). As evidenced by Figure 2a, a strong seasonal signal is present in most of the time series. The problem appears clearly when one takes two arbitrary signals that would be pure annual waves:

 $X_1 = \cos(2\pi v t + \varphi) \tag{1}$

 $X_2 = \cos(2\pi\nu t) \tag{2}$

If the time series are properly sampled, the correlation coefficient is a fair approximation of $cos(\phi)$, meaning that even with a 45° phase difference, the correlation amounts to 0.7, which may appear as important. Actually, the correlation analysis cannot be applied when the signal is dominated by the seasonal component. The same problem will appear whatever other comparison method is used, as

the presence of a strong periodic component is only significant if the detection of that period is an interesting result by it-self. For example, discovering the period of the translational motion of the inner core inside the outer core (Slichter mode) (Slichter, 1961) in SG records would be a nice discovery. Conversely, many geodetic time series one could take on Earth would exhibit at least some seasonal signal, and no conclusion can be drawn from such a result. One could argue that the fact that there is an annual signal in both series is significant by it-self, but this is not really instructive. On the contrary, correlation studies can be insightful after removing the seasonal component from the signal. Let us look at the correlation of the time series corrected for the annual component (Table 4 and Figure 3); 10 pairs out of 41 are significantly correlated: BH-MO, BH-PE, BH-ST, BH-WA, BH-WE, MB-VI, MB-WA, MO-PE, PE-WA and VI-WA, which is above the significance level. On the other hand, the fact that only a fourth of the pairs of time series appear significantly correlated when the seasonal cycle is filtered out is not consistent with a dominant coherent signal at the different stations. Note that, in each significant case but one (VI-WA), underground pairs and surface pairs are correlated, while underground-surface pairs are anti-correlated. This is again consistent with the local masses playing the dominant role in SG measurements, as local water would be above the gravimeter for underground station and below it for surface station.

The Empirical Orthogonal Function decomposition is a classical data mining technique, which allows retrieving common signal in a set of time series. Technical information and algorithms can be found in Preisendorfer [1988]. Starting from a set of time series $x_i(t_1)$, i = 1 ... N, l = 1... M, the covariance matrix is computed, and the eigenvectors of the covariance matrix, called principal components or Empirical Orthogonal Functions (EOFs), are used as a new basis in which the time series are written. Then, the time series can be written as:

 $x_i(t_l) = \sum_{k=1}^{N} \alpha_{k,i} \operatorname{T}_k(t_l)$ (3)

Where the $\alpha_{k,i}$ are the EOFs, and the functions $T_k(t_l)$ are their associated time series. Classically, the EOFs are sorted so that the first EOF explains the most variance in the initial set of time series. Most

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of the time, an important part of the variance is explained by only a few EOFs. Starting from a set of *N* time series, the covariance matrix is *NxN*; consequently, there are exactly *N* eigenvector for the matrix.

Let us take the 7 stations as discussed in Crossley et al. (2012): BH, MB, MC, MO, ST, VI and WE, from 2002.6 to 2007.8; note that the annual signal was not filtered out. For the reasons explained in the beginning of this section, it is difficult to interpret the results if the series contains an annual component: the EOF analysis will extract the seasonal signal as the first mode, even with a nonnegligible phase-lag (up to 45°) between the time series. Actually, the EOF analysis then only allows concluding to the presence of a seasonal signal in all the time series. Here, after filtering for the seasonal cycle, we computed the eigenvectors and the associated time series. Then, we computed the variance explained by each of the EOFs for each initial time series. The total variance explained by the first mode over the 7 SG time series is slightly less than 30%. There are three surface SGs (BH, MC, WE) where 78%, 67% and 58% of the variance is explained, the other four stations having less than 10% explained. This result may seem encouraging, but it is important to note that the algorithm focuses on the most significant EOF mode, i.e. the one that explains the most variance. To assess the significance of this result, we compare those results with what would be obtained for random time series. Speaking of climatically induced signal, the hypothesis of a red noise described as a degree one autoregressive process (AR1) is commonly used (Ghil et al., 2002). We estimated the AR1 parameters for each of the SG time series, and then generated a set of 100,000 time series with the same parameters. We then computed the EOF decomposition of each of the 100,000 sets of 7 time series, and computed the variance explained by the first EOF mode. The results are shown on Figure 4, which shows the distribution of the variance explained by the first mode, with a red vertical line at the value obtained with the SG dataset. We observe that the variance explained by the first mode narrows the mode of the distribution obtained with random data; this indicates that the 30% variance explained does not demonstrate that a common source of signal exists, it is simply due to the fact that the algorithm is built to extract the EOF in such a way that most of the variance will be

explained by one time series, whatever the input. This result is consistent with previous studies, where they show a common signal which is mostly annual, and not much beside, although that picture may change when longer series are available.

Nevertheless, the seasonal signal in the SG time series is information that needs to be analyzed. Its amplitude and phase are obtained by a linear least square fit of a sine wave at each station. They are given in Table 5 and represented as phasor diagrams on Figure 5a. Given that the local water content dominates the SG gravity signal, the phasors are also provided on Figure 5b with an opposite sign for the gravity data at the underground stations (CO, MB, MO, ST, VI and WA). This approach was adopted by Boy and Hinderer (2006) and Van Camp et al. (2010). Although the phasors are less dispersed, those diagrams show that the amplitudes and phases do not indicate a common signal, but rather station maxima within a seasonal cycle, as expected. Of course, GRACE does smooth these signals because of its much larger averaging footprint.

The magnitude of the annual signal depends on the local hydrogeological context. Even for homogeneous climate conditions, the topography around the SG stations, as well as the local petrology and the building umbrella effect, result in inhomogeneous ground water storage, as evidenced by several studies (e.g. Creutzfeldt et al., 2008; Deville et al., 2013; Lampitelli and Francis, 2010; Longuevergne et al., 2009; Meurers et al., 2007; Naujoks et al., 2010; Van Camp et al., 2006). Consequently, there is no conclusion to be drawn from either an agreement or a disagreement of amplitude in the seasonal signal. We now focus on the phase, which might be less dependent on the local context and more comparable with large scale information such as GRACE or climate models.

Figure 5b shows that the phases all are within a time interval of about 222 days; if we restrict our analysis to the largest seasonal signal, between MC and WA, the phases are included in a 77 day interval. This simply means that the maximum water load occurs within a season, which is to be expected. In short, the phase distribution does not allow concluding that the seasonal signal is

common for the available set of SG time series, but it is consistent with Central Europe being wettest at the end of the winter.

4. Common variability of SGs, GRACE and hydrological models

With a resolution of 400 km, GRACE barely distinguishes the position of the different stations and is mostly sensitive to the large scale feature of the ground water mass distribution. This would advocate for GRACE being consistent with a common signal in the SG time series, as long as this common signal actually exists, for example resulting from a large scale phenomena, and acts similarly on all terrestrial gravity sensors. In the case of the SG time series, we have shown that there is only little, if any, common signal, both at the annual and interannual time scales. This lack of coherence is at least partially caused by diverse site conditions. Nevertheless, as the subsurface ground water experiences a maximum at the end of the winter; one would expect at least some agreement in phase between the annual component of GRACE, the SG, and the hydrological models. This would not imply, considering their transfer functions, that they agree on the water distribution over Central Europe; it simply means that they more or less agree that winter is wetter than summer.

Figure 6 shows the phasor diagrams for the annual component at the different stations for the SGs, the 10 different GRACE solutions and the GLDAS and ERA hydrological models. As, in most cases, hydrology models predict seasonal cycles larger than the other ones, the corresponding arrows are reduced by a factor of 2 for the sake of clarity.

Globally, we see that, as expected, all GRACE solutions are relatively close in amplitude and in-phase from within 19 days (CO) to 63 days (MB), but not perfectly identical, depending on the location (at MB, WA and to a lesser extent, ST, there are more differences between the GRACE solutions, probably due to the closeness of the ocean). However differences between the solutions are globally smaller than the differences between GRACE solutions and hydrology models or SGs.

At all stations but PE and WE, the hydrological models disagree in amplitude, probably partly due to a simplified treatment of near field effects, and only agree within 4 months in phase (Table 5, Figures 6 and 7). For PE and WE, the amplitudes predicted by the ERA model are comparable to the SG observations, although the possible recent changes in the hydrogeological properties around the WE station may have changed this picture.

For 3 stations located above the ground (BH, MC, PE) and an underground one (MB), there are some phase and/or amplitude agreements between SGs and some of the GRACE solutions, but our sample is too small to draw any real conclusion.

5. Local effects

Obviously, there are to be some common signals within the water mass distribution around stations located within a few hundred kilometers, and these common signals may be emphasized in the GRACE signal. We have shown that this common signal does not dominate the SG series.

The dominant signal in the SG time series comes from the area directly around the instrument, within a few hundred meters, as shown e.g. by Creutzfeldt et al. (2008). A perfect hydrological local model accurately estimates the direct attraction from mass close to the gravimeter but, subtracting it, the corrected gravity signal cannot be consistent with the mass distribution observed by GRACE, as demonstrated in the Appendix. To compare SGs with GRACE, it is necessary to add back $g_{L_{Smooth}}(S)$, the smoothed local effect of the mass distribution, into the corrected SG series (see equations A4, A5 and A6). One could estimate $g_{L_{Smooth}}(S)$ by using GRACE or local models. In the first case, one would create a common signal, even if there is none, which is not appropriate; in the second case, one would have to rely on perfect hydrological models, but then, why using an SG for hydrological investigations?

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Our results, and that from previous studies, show that the agreement with GRACE is worse for underground station; this makes perfect sense considering that the part of the mass closest to the SG is above an underground instrument, which generates a partial cancelation of the signal, as in MO, ST and VI, but not in WA. Obviously, considering those stations as anomalous, as done by Crossley et al. 2012, does improve the coherence of the remaining set. Overall, it shows the limitation of the comparison of very local measurements with regional ones.

6. Conclusion

At first sight, looking for an agreement between SGs and GRACE is a long shot, as numerous studies have shown that most of the gravity effects recorded by SGs are induced by subsurface water dynamics in a radius around the gravimeter smaller than 1000 m. On the other hand, if successful, there would be much to be learned from the intercomparison in terms of validation, calibration, and corrections of geodetic and hydrological measurements.

The analysis of time series from 10 European superconducting gravimeters showed that (1) except for the presence of an annual cycle at most of the stations, as in most geodetic time series, there is no clear common behavior between the different SGs; (2) the consistency between the annual cycles of the different SGs is poor, both in phase and amplitude. Similarly, the annual cycles of the SGs are not consistent with predictions computed from GRACE and hydrological models.

Considering the complexity of the hydrogeological processes governing the conversion between rainfall and water mass distribution, it is easy to justify disagreements both in phase and in amplitude, as observed here. Consequently, our results do not demonstrate that the physical phenomena monitored by the SGs and GRACE are different. On the other hand, a study combining those data sets can only be fruitful if there are at least some degrees of consistency.

Terrestrial gravity measurements can be fruitfully used to perform comprehensive, local hydrogeological investigations, as shown in Wettzell (Creutzfeldt et al., 2010) or in the Larzac karstic area (Jacob et al., 2010); on the other hand GRACE has provided numerous information on large scale hydrological and geodynamic phenomena (Ramillien et al., 2008; Pollitz, 2006). But, this study shows that the feasibility of joined studies is still unclear, in particular because it is impossible to correct SG data for local phenomena to make them comparable with GRACE observations.

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Figure captions

Figure 1. Map showing the location of the SG stations used in this study, see Table 1 for details.

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Figure 6. Phasor diagrams of the annual components at the different SG stations for the 10 different GRACE solutions and the two GLDAS and ERA hydrological models. For clarity the amplitude of the global hydrological models is reduced by a factor 2. The sign of the SG data from the CO, MB, MO, ST, VI and WA underground stations is inverted.

Figure 7. Phase distribution of the annual component in the GRACE solutions, hydrological models and SG time series at the SG stations. The sign of the SG data from the CO, MB, MO, ST, VI and WA underground stations is inverted. For the GRACE solutions the whiskers indicate the upper and lower extreme values, as well as the average and the one sigma confidence interval.

Tables

Acronym	Name	Instrument	Latitude	Longitude	Starting time	Ending time
вн	Bad Homburg	CD030 L	50.2285	8.6113	2002-01-05	2007-04-01
		SG044			2007-04-01	2012-06-01
со	Conrad	C025	47.9288	15.8609	2007-11-14	2012-05-28
МС	Medicina	C023	44.5219	11.6450	2002-01-05	2012-06-01
MB	Membach	C021	50.6092	6.0067	2002-01-05	2012-05-03
мо	Моха	C034 L	50.6447	11.6156	2002-01-05	2011-12-27
PE	Pecny	OSG-050	49.9138	14.7856	2007-05-06	2011-12-15
ST	Strasbourg	C026	48.6217	7.6850	2002-01-05	2010-12-27
VI	Vienna	C025	48.2493	16.3579	2002-01-05	2007-10-23
WA	Walferdange	OSG-040	49.6647	6.1528	2003-12-23	2012-05-28
WE	Wettzell	CD-029 L	49.1440	12.8780	2002-01-05	2010-10-10
		SG-030			2010-10-10	2012-06-30

Table 1. Description of the SG time series used in this study.

1 2
2 3 4 5 6 7 8
6 7
8 9 10
11 12
13 14 15
11 12 13 14 15 16 17 18
10
20 21 22 23 24 25 26 27 28 29 30 31 32 33 34 35 36 37 38
24 25 26
27 28
29 30 31
32 33 34
35 36
39
40 41 42
43 44 45
46 47
48 49 50
51 52 53
54 55
56 57 58
59 60

Solution	Origin	Reference Geoid	Time	Periodicity	Additional	Non-	Note
Name			period		filtering	tidal	
						ocean	
						load	
						added	
CSR	NASA Center	GGM03C		monthly	Destriping	No	Data access from
	for Space		2004-				http://podaac.jpl.nasa.gov/
	Research		2010				
	(USA)						
JPL	NASA Jet	GGM03C		monthly	Destriping	No	Data access from
	Propulsion		2004-				http://podaac.jpl.nasa.gov/
	Laboratory		2010				
	(USA)						
GFZ	GFZ German	EIGEN-6S		monthly	Destriping	No	Dahle et al. (2012)
	Research		2005-				
	Centre for		2010				
	Geosciences		2010				
	(Germany)						
ITG	Bonn	ITG-GRACE2010S	2002-	monthly	Destriping	No	Kurtenbach et al. (2009)
	University		2009				
	(Germany)						
AIUB	Bern	AIUB-GRACE03S	2003-	monthly	Destriping	No	Degree 2 zonal coefficient
	University		2009				corrected. Beutler et al. (2010)
	(Germany)						
GRGS	CNES French	EIGEN-		10-days	None	No	Regularized solution. Bruinsma
	Spatial	GRGS.RL02.MEAN-	2002-				et al. (2009)
	Agency -	FIELD	2012				
	GRGS group						
	(France)						
DTM	Delft	EIGEN-GL04C	2000	monthly	None	No	DTM-1b model. Liu et al. (2010).
	University of		2003-				Wiener filter based solution
	Technology		2010				(Klees et al., 2008).
657.0	(Netherlands).	6 657		6	6	Vee	6 657
GFZ_O	See GFZ	See GFZ		See GFZ	See GFZ	Yes	See GFZ
GRGS_O	See GRGS	See GRGS		See GRGS	See GRGS	Yes	See GRGS
ITG_O	See ITG	See ITG		See ITG	See ITG	Yes	See ITG

Table 2. Summary of the different GRACE solutions used in this study.

	BH	CO	MB	MC	MO	PE	ST	VI	WA
CO	31	-							
MB	-53	3	-						
MC	31	1	-27	-					
MO	-35	-13	-3	17	-				
PE	93	31	-21	58	-11	-			
ST	-54	-16	9	-20	35	15	-		
VI	3	N/A	27	11	9	-49	10	-	
WA	-69	-36	37	-4	14	-70	44	-9	-
WE	78	29	-40	18	-26	55	-39	-25	-49

Table 3. Correlation (in %) between the different time series shown on Figure 2a. Due to the strong annual component, the significance could not be tested.

	BH	CO	MB	MC	MO	PE	ST	VI	WA
CO	19	-							
	(86%)								
MB	-32	0	-						
	(92%)	(51%)							
MC	-15	0	14	-					
	(74%)	(51%)	(79%)						
MO	-52	-27	-3	22	-				
	(99%)	(88%)	(55%)	(87%)					
PE	77	30	-14	20	-52	-			
	(100%)	(91%)	(76%)	(78%)	(96%)				
ST	-52	N/A	-11	-3	40	N/A	-		
	(98%)		(71%)	(56%)	(93%)				
VI	11	N/A	39	1	15	N/A	13	-	
	(70%)		(96%)	(52%)	(71%)		(72%)		
WA	-56	-28	36	20	23	-67	38	-52	-
	(100%)	(92%)	(96%)	(85%)	(84%)	(99%)	(93%)	(100%)	
WE	72	14	-26	-3	-35	17	-26	-22	-1
	(100%)	(80%)	(93%)	(57%)	(94%)	(74%)	(89%)	(82%)	(52%)

Table 4. Correlation coefficients (in %) as shown on Figure 3, between the different time series shown on Figure 2b. The significant correlations are bold-faced. The coefficient is not evaluated when the time series overlapping is shorter than 3.5 years.

60

Station		SG		SG _{Inv}		GLDAS		ERA		GRACE	
	Α	Ph	Α	Ph	А	Ph	А	Ph	А	Ph	
BH	17.2	75.4	17.2	75.4	54.9	61.1	33.5	44.8	18.8	50.8	
CO	8.6	112.5	8.6	-67.5	39.7	61.2	18.8	43.4	25.7	69.9	
MB	9.8	236.4	9.8	56.4	45.8	46.7	16.6	37.2	16.5	41.5	
MC	15.7	34.1	15.7	34.1	57.2	55.1	25.5	41.2	23.0	61.7	
MO	0.9	160.3	0.9	-19.7	55.3	57.5	29.7	47.9	19.5	53.8	
PE	25.5	64.8	25.5	64.8	42.4	60.5	30.3	56.3	22.9	65.2	
ST	7.6	274.5	7.6	94.5	49.7	57.8	29.1	34.5	22.4	56.7	
VI	4.6	335.5	4.6	155.5	44.0	56.6	19.8	47.1	25.7	71.2	
WA	24.2	290.8	24.2	110.8	59.9	60.6	38.9	40.9	19.3	49.6	
WE	29.0	90.4	29.0	90.4	45.1	64.3	25.9	44.7	22.8	62.9	

Table 5. Amplitude (in nm/s²) and phases (in days) as shown in Figure 5 and 6, evaluated using least-squares fit of a purely annual term. SG_{Inv} means that the sign at the CO, MB, MO, ST, VI and WA underground stations is inverted. For GRACE the averages of the different solutions are provided.

Figures

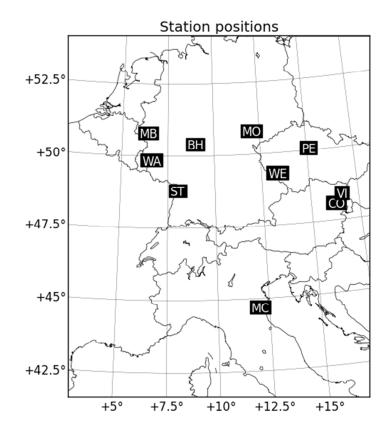


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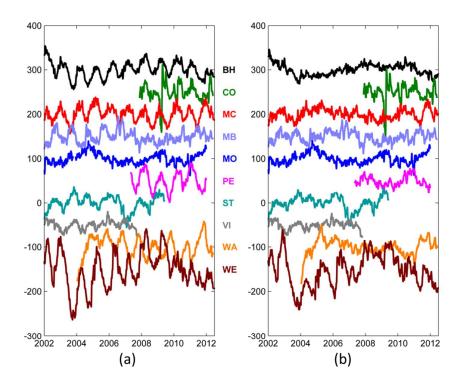


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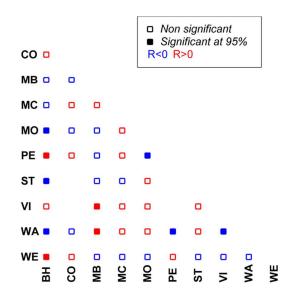


Figure 3. Correlation between the different SG time series as on Figure 2b, after removing a composite annual signal. The squared are filled when the correlation is significant (95% level). The coefficient is not evaluated when the time series overlapping is shorter than 3.5 years.

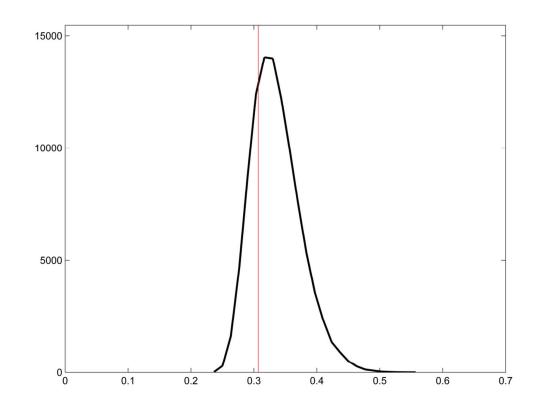


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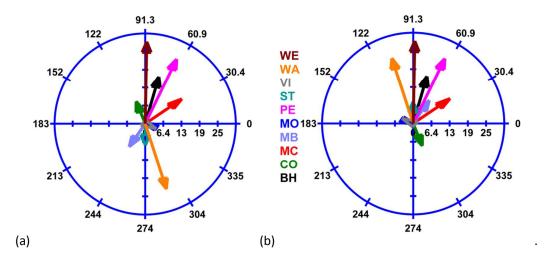


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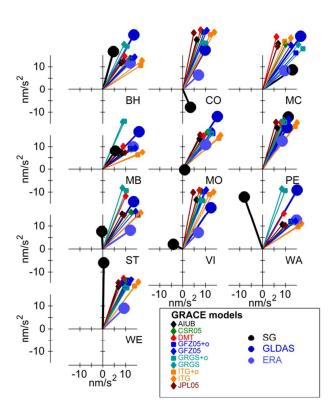


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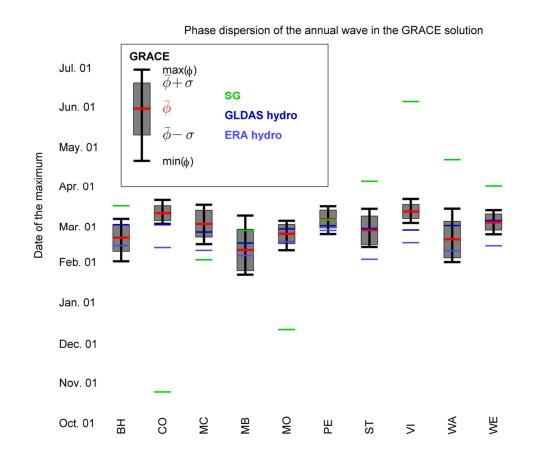


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Appendix

Comparison of the terrestrial and satellite observations of hydrological effects on gravity

Let us consider three hydrologic units, as shown in the figure A1:

- R remote (light blue);
- L local (dark blue);
- *L_{smooth}* local, where the water mass contained in the *L* unit is smoothed by the GRACE
 transfer function (hatched).

The water contained in *L* participates to the mass distribution that is observed by GRACE.

Then, we have:

 $-g_D$: gravity signal caused by the ground displacement due to the loading effect;

- g_R : gravity signal caused by the Newtonian effect from the remote mass in the R unit;

 $-g_L$: gravity signal caused by the Newtonian effect from the local mass in the *L* unit;

These quantities are understood as projection of their corresponding vectors onto the vertical direction.

The gravity measured by the SG at the surface station S reads as:

$$g_{SG}(S) = g_D + g_R(S) + g_L(S)$$
 (A1)

The gravity measured by the SG at the underground station *U* reads as:

$$g_{SG}(U) = g_D + g_R(U) + g_L(U)$$
 (A2)

If the distance *r* to the limit of the global domain *G* is large compared to the distance between points *S* and *U*, we have:

 $h(S) - h(U) \ll r$

and:

1

$$g_R(U) \cong g_R(S)$$

Such that (A2) becomes:

$$g_{SG}(U) \cong g_D + g_R(S) + g_L(U) \tag{A3}$$

The gravity as measured by GRACE, reconstructed at the SG station location *S* as a ground gravity value, reads as:

$$g_{GRACE}(S) = g_D + g_R(S) + g_{L_{Smooth}}(S)$$
(A4)

To make a meaningful comparison between GRACE and SG, we have included the deformation part g_D deduced from the GRACE load models in the estimation of $g_{GRACE}(S)$.

 $g_{L_{Smooth}}(S)$ is the gravity effect from the water contained in the *L* unit, reconstructed at the *S* site. The equivalent water height is of the same order of magnitude as in the *R* unit, although the gravity effect $g_R(S)$ turns out to be small as that water is far from *S*.

If the topography of *L* is flat and the water contained in *L* is homogeneously distributed within *L*, the water mass in *L* and the water mass in L_{smooth} would be similar and $g_L \cong g_{L_{smooth}}$. On the other hand, if the same total water mass is concentrated in a small region within *L*, then the mass in *L* would be far larger than in L_{smooth} and $g_L \gg g_{L_{smooth}}$. Hence, the GRACE signal would be dominated by $g_D + g_R$, and the SG signal by g_L . Yet, in this case, even if GRACE and SG see different sources, the signal maxima would be within 3 months because all water sources present seasonal variations.

Taking (A3) and (A4) into account, the correction for making SG gravity comparable with GRACE requires applying a remove-restore technique (bold):

$$g_{SG,compar}(S) = g_{SG}(S) - g_L(S) + g_{L_{Smooth}}(S)$$
(A5)

$$g_{SG,compar}(U) = g_{SG}(U) - g_L(U) + g_{L_{Smooth}}(S)$$
(A6)

Those equations shows that it is only possible to convert g_{SG} to something that we can compare with GRACE if we know, by some other means, the mass everywhere around the gravimeter.

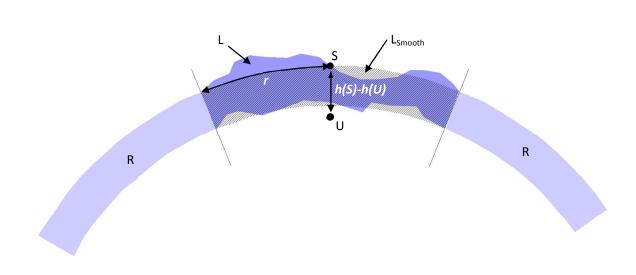


Figure A1: The three different hydrologic units *R*, *L* and L_{smooth} taken into account when investigating hydrological effects on satellite and terrestrial gravity measurements. Note, that the location of *U* shown here is exemplary. *U* can be located below the surface anywhere, within or below the *L* unit"