Department of Computer Science Drexel University College of Engineering

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Level Set and PDE Methods for Visualization

Notes for IEEE Visualization 2005 Course #6 Minneapolis, MN October 25, 2005

Organizer

David Breen Drexel University

Speakers

Mike Kirby University of Utah

Aaron Lefohn University of California, Davis

Ken Museth Linköping University
Tobias Preusser University of Bremen
Guillermo Sapiro University of Minnesota

Ross Whitaker University of Utah

Course Schedule

Session	1 -	. DDF	and	امریم ا	CΔt	Fund	lamanta	ıle
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- 8:30 Welcome Breen
- 8:40 Introduction to PDEs with Applications Sapiro
- 9:50 Introduction to Level Set Methods Whitaker

10:00 Break

Session 2 – Level Set Software and Numerical Methods

- 10:30 Open Source Level Set Software (ITK) Whitaker
- 10:45 Level Set Applications: Segmentation, Surface Processing and Reconstruction Whitaker
- 11:30 Numerical Methods and Algorithms for PDEs Museth

12:15 Lunch Break

Session 3 – Implementation Details

- 1:45 Level Set Applications: Segmentation and Reconstruction From Sparse Data Museth
- 2:50 Introduction to GPU Computation Lefohn
- 3:20 Level Set Application: Interactive Segmentation Lefohn

3:45 Break

Session 4 - PDE Applications

- 4:15 Level Set Method for 3D Anisotropic Geometric Diffusion Preusser
- 4:45 PDE Methods in Flow Field Post-Processing Preusser
- 5:15 Algorithms for Solving Reaction-Diffusion Equations Kirby
- 5:35 Reaction-Diffusion Models for Vector Visualization Kirby

6:00 Course Ends

Table of Contents

Session 1 - PDE and Level Set Fundamentals

Introduction to PDE's in Image Processing, Visualization, Computer Vision and Computer Graphics (Slides)

G. Sapiro

Isosurfaces, Level Sets and PDEs (Slides)

R. Whitaker

Session 2 - Level Set Software and Numerical Methods

Fundamental Properties of Level Sets (Slides)

K. Museth

Session 3 - Implementation Details

Segmentation with Level Sets (Slides)

K. Museth

Introduction to GPU Computation (Slides)

A. Lefohn

Interactive Level-Set Deformation On The GPU (Slides)

A. Lefohn

Session 4 - PDE Applications

PDE Methods in Flow Field Post-Processing &

Anisotropic Levelset Diffusion (Slides)

T. Preusser

Anisotropic Diffusion in Vector Field Visualization on Euclidean Domains and Surfaces

U. Diewald, T. Preusser, and M. Rumpf

Reaction-Diffusion Models for Vector Visualization: Algorithms and Implementations (Slides)

R.M. Kirby

Display of Vector Fields Using a Reaction-Diffusion Model

A. Sanderson, C. Johnson, and R.M. Kirby

Course Abstract

Level set methods, an important class of partial differential equation (PDE) methods, define dynamic surfaces implicitly as the level set (isosurface) of a sampled, evolving nD function. This course is targeted for researchers interested in learning about level set and other PDE-based methods, and their application to visualization. The course material will be presented by several of the recognized experts in the field, and will include introductory concepts, practical considerations and extensive details on a variety of level set/PDE applications.

The course will begin with preparatory material that introduces the concept of using partial differential equations to solve problems in visualization. This will include the structure and behavior of several different types of differential equations, e.g. the level set, heat and reaction-diffusion equations, as well as a general approach to developing PDE-based applications. The second stage of the course will describe the numerical methods and algorithms needed to implement the mathematics and methods presented in the first stage, including information on implementing the algorithms on GPUs. Throughout the course the technical material will be tied to applications, e.g. image processing, geometric modeling, dataset segmentation, model processing, surface reconstruction, anisotropic geometric diffusion, flow field post-processing and vector visualization.

Prerequisites

Knowledge of calculus, linear algebra, computer graphics, visualization, geometric modeling and computer vision. Some familiarity with differential geometry, differential equations, numerical computing and image processing is strongly recommended, but not required.

Organizer Biography

David Breen is an Assistant Professor in the Computer Science Department at Drexel University. He has held research positions at the Center for Advanced Computing Research and the Computer Graphics Lab at the California Institute of Technology, the European Computer-Industry Research Centre, the Fraunhofer Institute for Computer Graphics, and the Rensselaer Design Research Center. His research interests include level set models for graphics and visualization, medical image analysis and segmentation, geometric modeling and computational biology. He has published over 50 research papers in these and other areas, as well as the book *Cloth Modeling and Animation*. Breen received a B.A. in Physics (Colgate University, 1982), and a Ph.D. in Computer and Systems Engineering (Rensselaer Polytechnic Institute, 1993).

E-mail: david@cs.drexel.edu

Speaker Biographies

Mike Kirby is an Assistant Professor in the School of Computing and is Director of the Computational Engineering and Science (CES) Program at the University of Utah. He is also affiliated with the Institute for Scientific Computing and Imaging. He received ScM degrees in computer science and applied mathematics and the PhD degree in applied mathematics from Brown University. His teaching and research focuses on large-scale scientific computation and visualization, with an emphasis on the scientific cycle of mathematical modeling, high-performance computation and parallelization, visualization, evaluation, and understanding. He has co-authored a textbook on parallel scientific computing, and numerous journal articles and book chapters spanning both theory and applications in computational engineering and science. E-mail: kirby@cs.utah.edu

Aaron Lefohn is a Ph.D. student in computer science at the University of California at Davis and a graphics software engineer at Pixar Animation Studios. His research interests include general-purpose computation on graphics hardware and high-quality interactive rendering. His current research focuses on data structure abstractions for graphics hardware, and his M.S. thesis describes an interactive GPU-based level-set

segmentation and visualization system. Lefohn completed an M.S. in computer science at the University of Utah in 2003, an M.S. in theoretical chemistry from the University of Utah in 2001, and B.A. in chemistry from Whitman College in 1997. He is an NSF graduate fellow in computer science. E-mail: lefohn@cs.ucdavis.edu

Ken Museth is a professor of Computer Graphics at Linköping University and an Adjunct Professor at Aarhus University. He received his Ph.D. in computational quantum dynamics from the University of Copenhagen in 1997. From 1998 to 2003 he was a visiting faculty member in the Chemical Physics Department, then a research scientist in the Computer Science Department at the California Institute of Technology. He has also been a scientific consultant to Digital Domain and NASA's Jet Propulsion Laboratory. His research activities focus on the areas of deforming geometry and level set methods. E-mail: kenmu@itn.liu.se

Tobias Preusser studied Mathematics at the University of Bonn, Germany, and at New York University. In 1999 he finished his Diploma thesis on "Adaptive Methods in Large-Scale Image Processing" at the University of Bonn. During his PhD studies he worked on diffusion methods in scientific visualization and geometric image processing. In 2000 and 2001 he was a visiting researcher at the Max-Planck-Institute for Mathematics in the Sciences in Leipzig, Germany, and at Bologna University, Italy, respectively. Preusser obtained the PhD degree from the Department of Mathematics of the University of Duisburg, Germany. His thesis investigated anisotropic geometric diffusion in image and image-sequence processing. He is now a postdoctoral fellow at the Center for Complex Systems and Visualization at the University of Bremen, Germany, where he is continuing his research on PDE methods in image processing and scientific visualization, as well as image-based computing in the context of medical simulations. E-mail: tp@mevis.de

Guillermo Sapiro is a Distinguished McKnight University Professor with the Department of Electrical and Computer Engineering at the University of Minnesota. He received a B.Sc. (summa cum laude) and Ph.D. in Electrical Engineering from the Technion, Israel Institute of Technology, in 1989 and 1993 respectively. He was a Member of Technical Staff at HP Labs in Palo Alto, CA. He works on differential geometry and geometric partial differential equations, both in theory and applications in computer vision, computer graphics, medical imaging and image analysis. Sapiro

recently co-edited special issues of IEEE Image Processing and the Journal of Visual Communication and Image Representation. He has authored and co-authored numerous papers in this area and has written a book published by Cambridge University Press, January 2001. He has received the Gutwirth Scholarship for Special Excellence in Graduate Studies, the Ollendorff Fellowship for Excellence in Vision and Image Understanding, the Rothschild Fellowship for Post-Doctoral Studies, the ONR Young Investigator Award, the PECASE Award, and the NSF Career Award. E-mail: guille@ece.umn.edu

Ross Whitaker is an Associate Professor in the School of Computing at the University of Utah and is affiliated with the Institute for Scientific Computing and Imaging. He received a B.S. in Electrical Engineering and Computer Science (Princeton University, 1986) and a Ph.D. in Computer Science (University of North Carolina, Chapel Hill, 1993). He was a research staff member at the European Computer-Industry Research Centre, and an Assistant Professor at the University of Tennessee. He teaches and conducts research in computer vision, image processing, medical imaging, and computer graphics/visualization. He has published numerous papers and book chapters on PDE/level set methods for image processing and computer graphics. E-mail: whitaker@cs.utah.edu

Level Set / PDE Web Sites

Sapiro Home Page http://www.ece.umn.edu/users/guille

Whitaker Home Page http://www.cs.utah.edu/~whitaker

VISPack Web Site http://www.cs.utah.edu/~whitaker/vispack

ITK Web Site http://www.itk.org

Museth Home Page http://gg.itn.liu.se

Lefohn Home Page

http://graphics.cs.ucdavis.edu/~lefohn

Preusser Home Page

http://www.mevis.de/~tp

Kirby Home Page

http://www.cs.utah.edu/~kirby

Breen - Geometric Modeling and Deformable Models

http://www.cs.drexel.edu/~david/geom_mod.html

http://www.cs.drexel.edu/~david/deform_mod.html

Osher Home Page

http://www.math.ucla.edu/~sjo

UCLA CAM Technical Reports

http://www.math.ucla.edu/applied/cam

Level Set Systems, Inc.

http://www.levelset.com

Fedkiw Home Page

http://www.graphics.stanford.edu/~fedkiw

Sethian Home Page

http://www.math.berkeley.edu/~sethian

Rumpf Home Page

http://numerik.math.uni-duisburg.de/people/rumpf/rumpf.shtml

Strzodka Home Page

http://numerik.math.uni-duisburg.de/people/strzodka/strzodka.htm

Related Papers

- J. Becker, T. Preusser, and M. Rumpf, "PDE Methods in Flow Simulation Post Processing," *Computing and Visualization in Science*, 2000.
- D. Breen, R. Whitaker, K. Museth and L. Zhukov, "Level Set Segmentation of Biological Volume Datasets," Handbook of Medical Image Analysis, 2005.
- J. Cates, A. Lefohn and R. Whitaker, "GIST: An Interactive, GPU-Based Level-Set Segmentation Tool for 3D Medical Images," *Medical Image Analysis*, 2004.
- U. Diewald, T. Preusser, and M. Rumpf, "Anisotropic Diffusion in Vector Field Visualization on Euclidian Domains and Surfaces," IEEE Transactions on Visualization and Computer Graphics, 2000
- A. Lefohn, J. Kniss, C. Hansen and R. Whitaker, "A Streaming Narrow-Band Algorithm: Interactive Computation and Visualization of Level Sets," *IEEE Transactions on Visualization and Computer Graphics*, 2004.
- A. Lefohn, J. Kniss, R. Strzodka, S. Sengupta and J. Owens, "Glift: An Abstraction for Generic, Efficient GPU Data Structures," ACM Transactions on Graphics, 2005.
- F. Memoli and G. Sapiro, "Fast Computation of Weighted Distance Functions and Geodesics on Implicit Hyper-Surfaces," *Journal of Computational Physics*, 2001.
- K. Museth, D.E. Breen, R.T. Whitaker and A.H. Barr, "Level Set Surface Editing Operators," *Proc. SIGGRAPH 2002 Conference*.
- K. Museth, D.E. Breen, R.T. Whitaker, S. Mauch and D. Johnson, "Algorithms for Interactive Editing of Level Set Model," *Computer Graphics Forum*, 2005.
- M. Nielsen and K. Museth, "Dynamic Tubular Grid: An Efficient Data Structure and Algorithms for High Resolution Level Sets," to be published in *Journal of Scientific Computing*, 2005.
- T. Preusser and M. Rumpf, "A Level Set Method for Anisotropic Geometric Diffusion in 3D Image Processing," SIAM Journal of Applied Mathematics, 2002.
- A. Sanderson, C. Johnson, R.M. Kirby, "Display of Vector Fields Using a Reaction-Diffusion Model," Proc. IEEE Visualization 2004.
- T. Tasdizen, R. Whitaker, P. Burchard and S. Osher, "Geometric Surface Processing via Normal Maps," ACM Transactions on Graphics, 2003.
- R.T. Whitaker, "A Level-Set Approach to 3D Reconstruction From Range Data," *International Journal of Computer Vision*, 1998.
- R.T. Whitaker, "Isosurfaces and Level-Sets," *The Visualization Handbook*, Eds. C. Hansen and C. Johnson, Elsevier, 2005

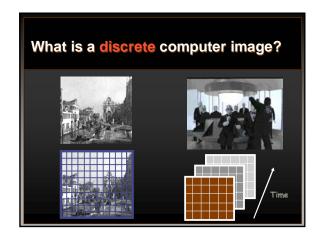
Related Books

- S. Osher and R. Fedkiw, *Level Set Methods and Dynamic Implicit Surfaces*, Springer, New York, 2002.
- S. Osher and N. Paragios (eds.), Geometric Level Set Methods in Imaging, Vision and Graphics, Springer, New York, 2003.
- G. Sapiro, Geometric Partial Differential Equations and Image Analysis, Cambridge University Press, 2001.
- J. Sethian, Level Set Methods and Fast Marching Methods, Cambridge University Press, 1999

Introduction to PDE's in image processing, visualization, computer vision, and computer graphics

Guillermo Sapiro

Electrical and Computer Engineering
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Consequences of discrete image representations

- Classical image processing and computer vision is based on discrete mathematics (most of it)
 - Sums instead of integrals
 - Re-definition of classical continuous operators as a Laplacian, Minkowsky addition, etc

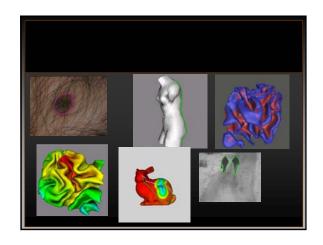
The PDE's approach

- Images are continuous objects
- Image processing is the results of iteration of infinitesimal operations: PDE's
- Differential geometry on images
- Computer image processing is based on numerical analysis

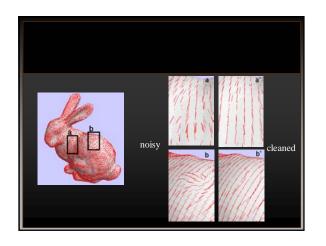
Why? Why Now? Who?

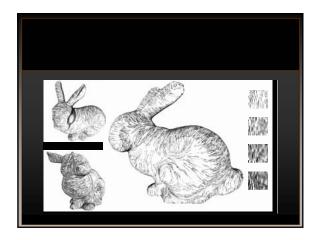
- Why now:
 - Computers!!!
 - People
- Why:
 - New concepts
 - Accuracy
 - Formal analysis (existence, uniqueness, etc)
- Consequences:
 - Many state of the art results

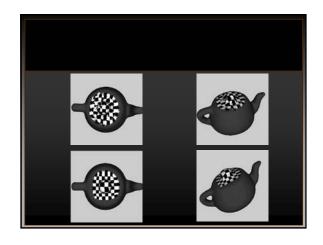
By this afternoon, we will be able to ...













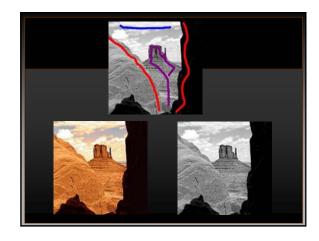


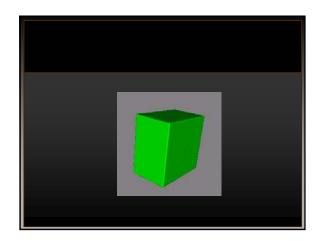




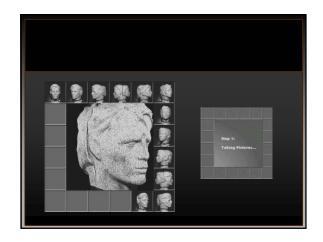


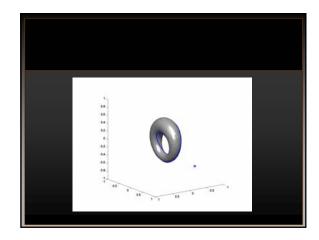


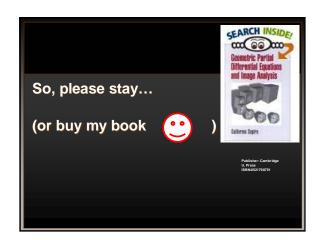


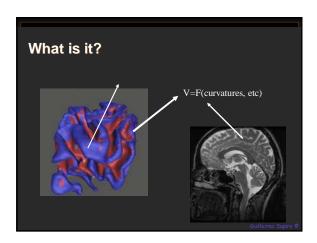






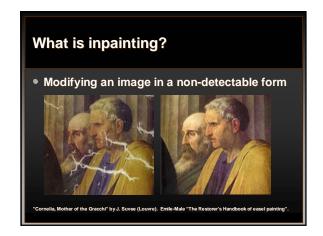








Overview • Goal and background • Art, biology, math, and engineering come together • Related work • Inpainting • Filling-in • Inpainting and image decomposition • 3D surface filling-in





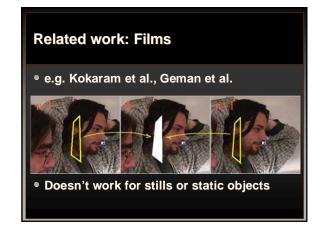


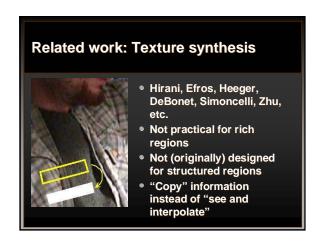






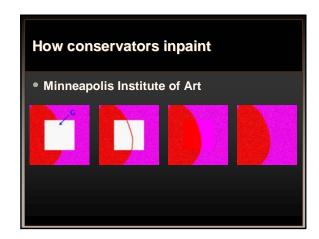


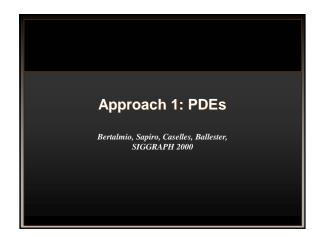


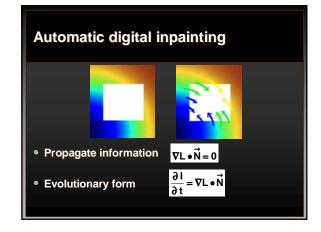


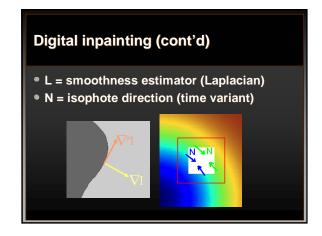


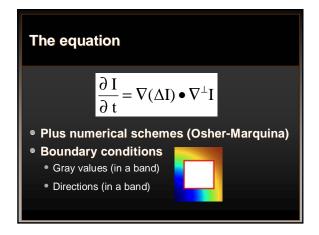


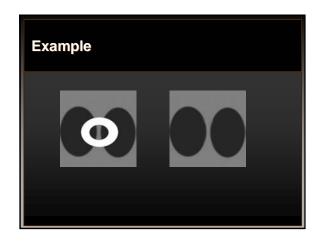




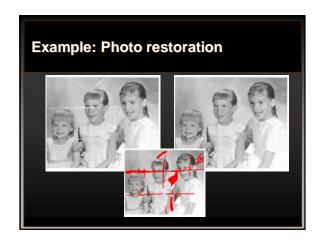




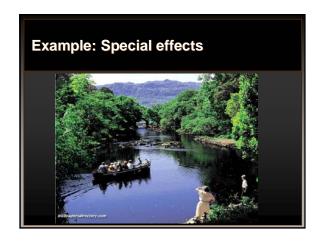


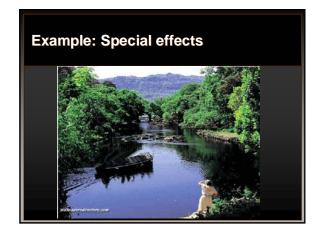


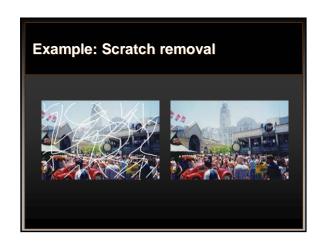


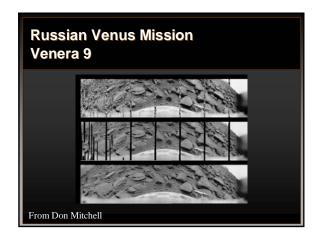


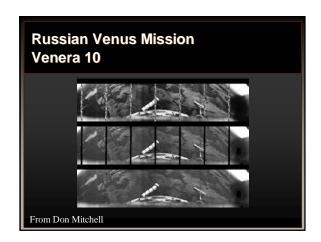


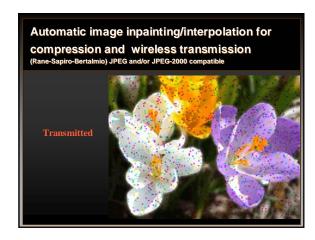








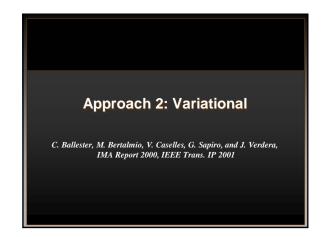


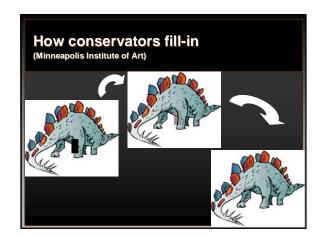


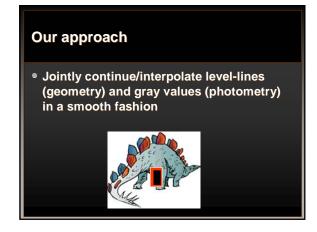


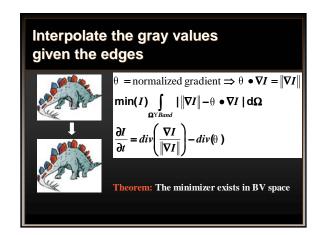


Concluding remarks (cont.) • Connected to fluid dynamics (see Bertalmio-Bertozzi-Sapiro CVPR 2001) $\frac{\partial (\Delta I)}{\partial t} = \nabla (\Delta I) \cdot \nabla I^{\perp}$ • Opens then door to high order PDE's • Extended to a variational formulation: Approach 2...

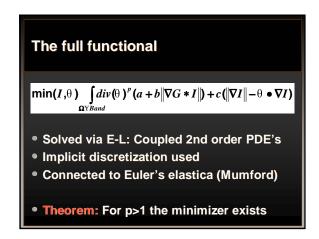


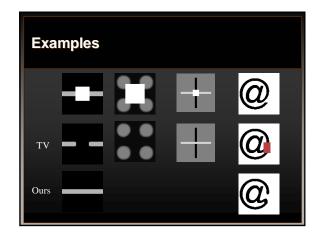




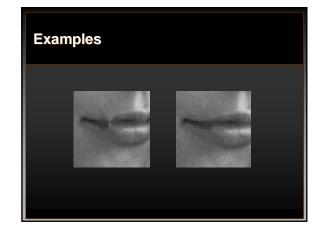


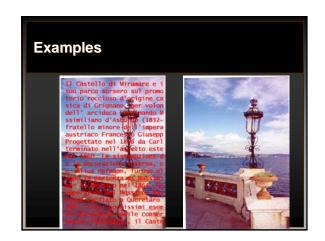




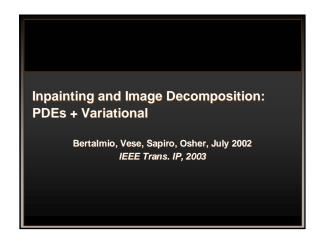


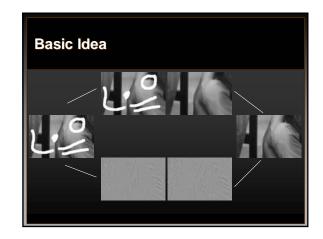


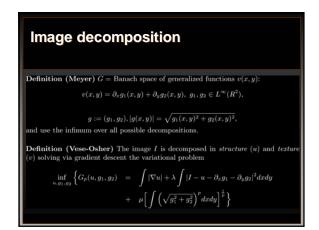


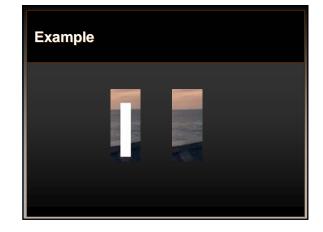




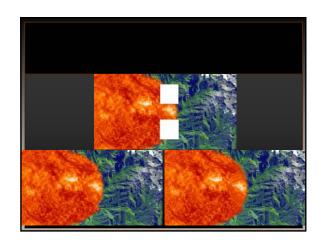




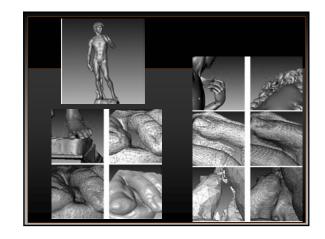








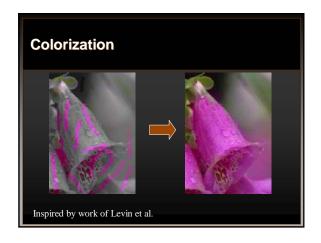


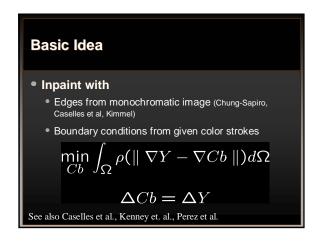


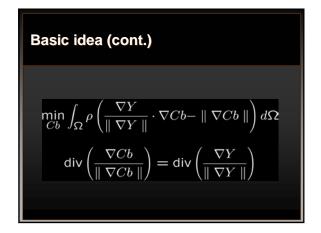




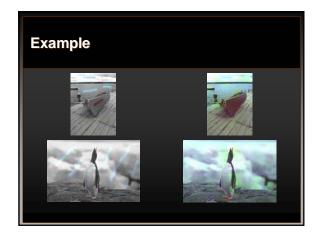


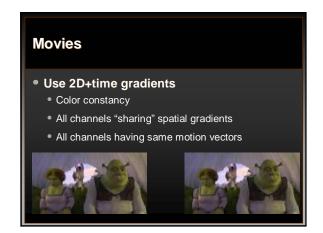


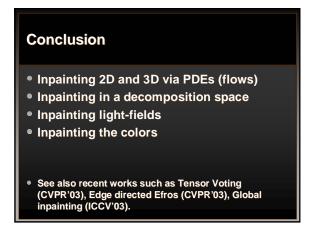












Isosurfaces, Level Sets, and PDEs

Ross Whitaker SCI Institute, School of Computing University of Utah



The Next Several Talks

- 1. Implicit surfaces and level-set geometry
- 2. Level sets, numerical schemes, and software
- 3. Applications of level sets and PDEs to surface/volume processing

SCI

Overview

- · Introduce implicit surfaces/level sets
- · Geometry of level sets
- Application of level-set geometry

SCI

Isosurfaces

· Implicit representation

$$F: \underset{x,\,y,\,z}{U} \mapsto \Re$$

Pomain of volume – where surface lives

$$U \subset \Re^3$$

· Surface S is set of points

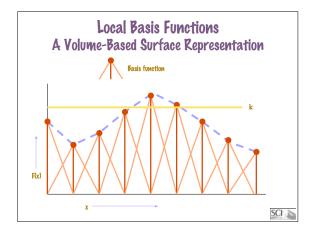
$$S = \{\bar{x} | F(\bar{x}) = k\}$$

SCI

Isosurfaces $\hbox{\bf . Implicit formulation } F(\bar x) = k$ Curve/Surface Model $\hbox{\bf Family of Embedded Contours}$

How Do We Represent F(x)?

- · Linear combination of global basis functions
 - "Blobby" models [Blinn 82]
 - Peformation by modifying size, position, number, etc. [Muraki 91]
- · Linear combination of <u>local</u> basis functions
 - Local deformations defined by neighborhood
 - Many degrees of freedom-arranged on grid
 - Well defined relationship between surface motion and grid values



Local Basis Functions

- · Geometry defined by local operations
- · Continuous mathematics on F(x)
- Grid value (voxel) manipulations determined by well-defined numerical methods
 - Level-set method [Osher & Sethian 88]

SCI

Visualizing Level Sets

- Direct volume rendering methods
 - E.g. cast rays from viewpoint
 - Transfer functions, root finding
- Extract surface primitives
 - E.g. marching cubes [Lorensen & Cline, 87]
 - Others
- · Not subject of this talk

SCI

Geometry of Isosurfaces

- Surface normals
- Curvature
- Goal: express surface geometry in terms of derivatives of F()

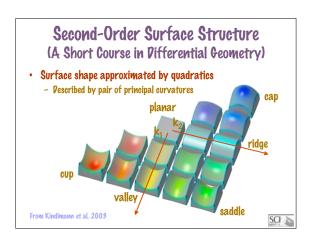
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Surface Normals

· Exist for every point in U

$$\bar{n}(\bar{x}) = \frac{\nabla F(\bar{x})}{|\nabla F(\bar{x})|} \ \text{where} \ \bar{x} \in U.$$

- Gives normal to level-set passing through that point
- · Convention inside or out (be consistent)
- · How to compute? (e.g. central differences)



Surface Curvature

Principle curvatures

 k_1, k_2

- Principle directions
- \bar{e}_1, \bar{e}_2
- Invariants
 - Gaussian curvature

 $K = k_1 \times k_2$

- Mean curvature $H=\frac{k_1+k_2}{2}$ Peviation from flatness $D=\sqrt{k_1^2+k_2^2}$

 - total curvature

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Curvature of Isosurfaces

· Projection operator (tangent plane)

$$P = I - \bar{n} \otimes \bar{n} = \left(\begin{array}{ccc} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{array}\right) - \left(\begin{array}{ccc} n_x n_x & n_x n_y & n_x n_z \\ n_y n_x & n_y n_y & n_y n_z \\ n_z n_x & n_z n_y & n_z n_z \end{array}\right)$$

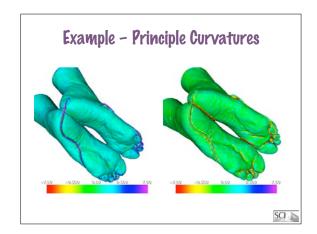
- Hession of F() $D^2F = \left(egin{array}{ccc} F_{xx} & F_{xy} & F_{xz} \\ F_{xy} & F_{yy} & F_{xz} \\ F_{xz} & F_{yz} & F_{zz} \end{array} \right)$
- · Curvature (matrix) given by projected, normalized Hessian of F()

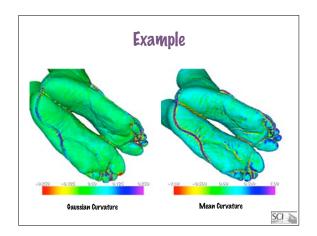
$$W = \frac{P(D^2F)P}{|\nabla F|}$$

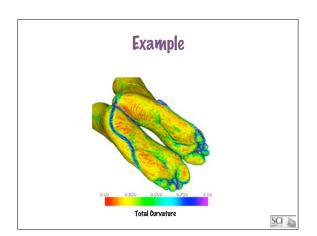
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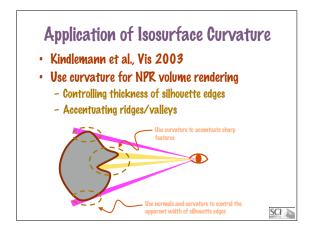
Curvature of Isosurfaces

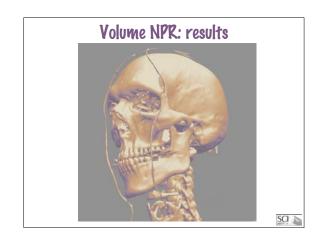
- Eigenvalues of W $\rightarrow k_1, k_2, 0$
- Eigenvectors of W -> e_1 , e_2 , n
- · Trace of W -> 2H
- · Norm of W -> D
- $K = 4H^2 D^2$
- · Note: derivatives must be taken using appropriately smooth basis functions

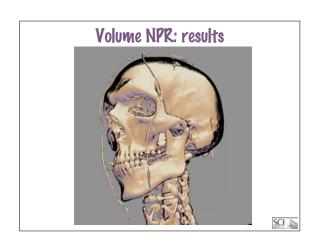


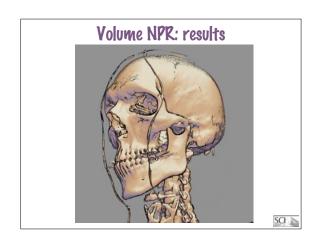


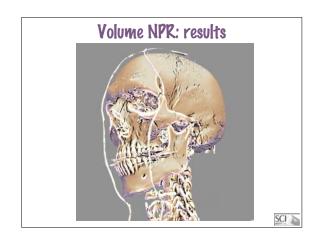


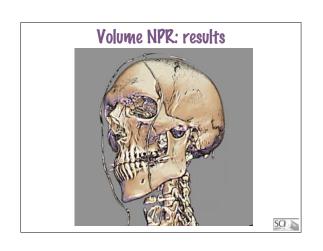












Summary

- Implicit surfaces
- Volumes
- · Geometry of level sets

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The Next Several Talks

- 1. Implicit surfaces and level-set geometry
- 2. Level sets, numerical schemes, and software
- 3. Applications of level sets and PDEs to surface/volume processing

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Level Set Introduction

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Overview

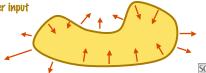
- · Peformable models
- · Level-set equations
- Numerical/computational techniques
- Software

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Freeform Deformable Surfaces

- E.g. "Snakes" [Kass et al. 86]
- Velocity v() at each point x on \$
- Where does v() come from?
 - Pata (e.g. attraction to edges in images)
 - Geometry (e.g. curvature, smoothness)





Level Sets - Moving Isosurfaces

- Osher and Sethian 1988
 - Method for modeling moving wave fronts
 - Formulation and numerical scheme
- Strategy
 - Function F() encodes the motion of the moving interface
 - Allows for a great deal of flexibility of shapes and topologies

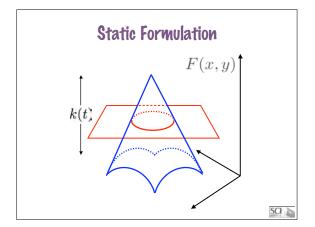
Static vs Dynamic Formulation

 $F(\bar{x}) = k(t)$

 $F(\bar{x},t) = k$

- Static
 - Single F, k varies
 - Motion strictly inward or outward
 - Fast marching method O(N lg M) Sethian 95
- Dynamic
 - Evolving F, k fixed
 - General motions
 - Front tracking schemes





Dynamic Formulation F(x, y; t) $F(x, y; t + \Delta t)$ SCI

Level-Set Equation Perivation

Total derivative is zero

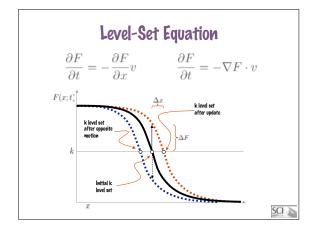
$$\frac{d}{dt}F(\bar{x}(t),t) = \frac{dk}{dt} = 0$$

$$\begin{split} & \frac{\text{Chain rule}}{dt} F(\bar{x}(t),t) = \frac{\partial F(\bar{x},t)}{\partial \bar{x}} \frac{d\bar{x}}{dt} + \frac{\partial F(\bar{x},t)}{\partial t} \end{split}$$

Level set equation (PDE)

$$\frac{\partial F}{\partial t} = -\nabla F \cdot \bar{v}$$

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Numerical Issues

- · Analytical expressions approximated on discrete grid u^n_{ij}
 - Finite forward differences in time

$$u_{ij}^{n+1} = u_{ij}^n + \Delta t \Delta u_{ij}^n$$

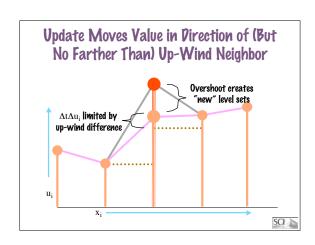
- Spatial derivatives approximated using kernels or stencils

$$\frac{\partial F}{\partial x} \approx \frac{u_{i+1,j}^n - u_{i-1,j}^n}{2} \longrightarrow \text{ 10-1}$$

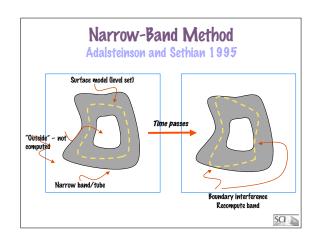
- Special care must be taken in the first-order derivatives in the LS equation (up-wind scheme) $$^{-1/2}$$
 - Must maintain monotonicity -> one-sided derivatives 1/2

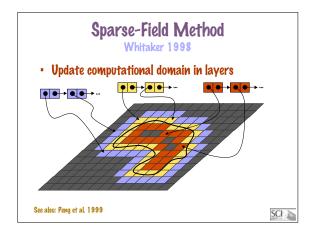


Updates Proportional to 1 st Perivatives • Avoid creation of new level sets Level-set motion Down wind Up wind Up-wind difference Up-wind difference

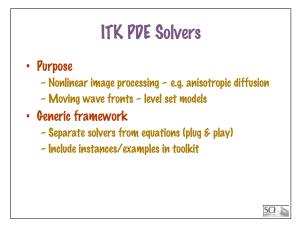


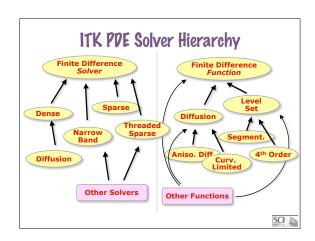
Front Tracking Solution important only in proximity to Is (wf) of interest Maintain computational domain that moves with wf Computational cost grows with surface area not volume

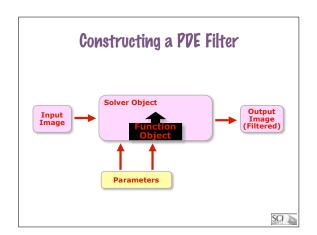


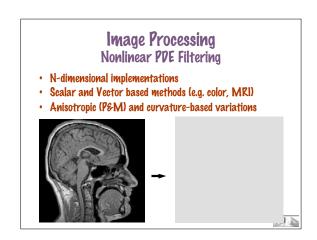


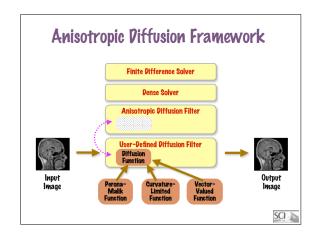
Public Pomain Software Insight Toolkit (ITK) Large toolkit for multidimensional image processing NIH sponsored consortium-industry and academics Includes API (algorithms) and applications www.itk.org Includes a framework for PPE-based image processing in N dimensions Others: Vispack Ian Mitchel http://www.cs.ubc.ca/~mitchell/ToolboxLS

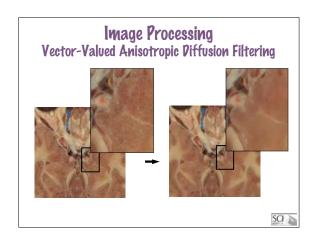


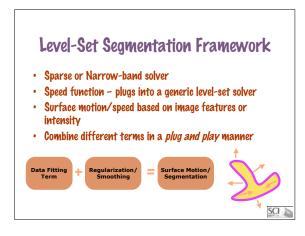


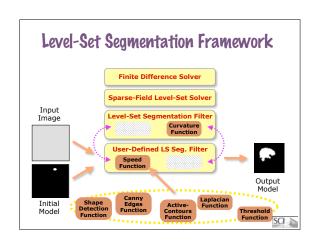


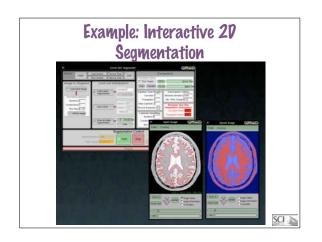


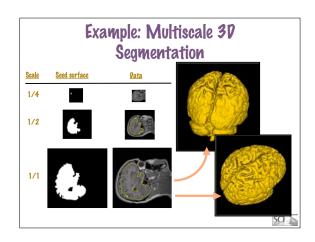




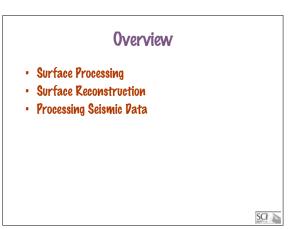








Applications of PDEs



Surface Processing

- Goals
 - Rich set of tools paralleling those available for images
 - Scientific approach that does not depend on content
 - Free from arbitrary decisions (user input)
- Why level sets?
 - Topological changes are part of smoothing
 - Shape based-free from parameterization
 - Good for data-driven applications
 - Not as useful when parameterization is part of model

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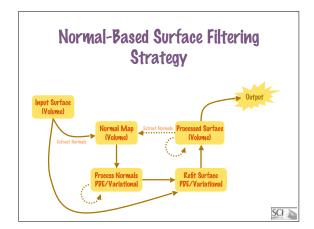
Generalizing IP to Surfaces E.g. Feature-Preserving Flows • Image Processing - Anisotropic smoothing (Perona & Malik) - Markov Random Fields

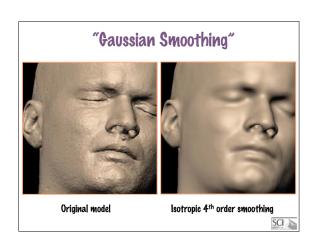
Strategy

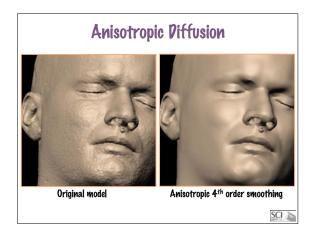
- · Variational formulations of curvature
 - Generalization of grad mag from IP
 - Gradient of normal map
 - Allowance for outliers
- · Decouple the normals from surface
 - Process normals
 - Refit surface to normal map
 - System of 2nd-order equations
- Tasdizen, Whitaker, Burchard, Osher, IEEE Vis 2002, 106 October 2003.

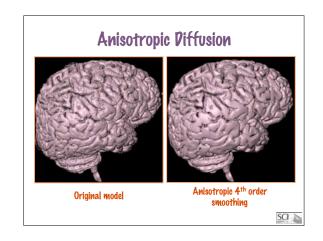
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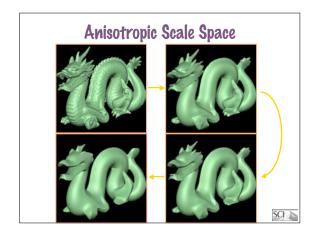
Surface Filtering Strategy Surface & Normals Process Normals Refit Normals



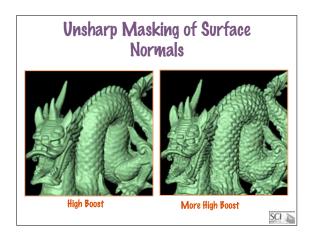














Surface Reconstruction

- · Surface Measurement Technologies
 - LAPAR, structured light, vision
 - Tomography
 - Sonar/ultrasound
 - Radar
- · New computational capabilities
- Applications
 - Visualization
 - Analysis

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Goals

- Foundations
 - Statistical framework
- Implementations
 - Numerical
 - Computational parallelism
- Systems
 - Practical-real data
 - Computation time

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3D Surface Reconstruction

- Noise
- Occlusions
- Geometric Distortion (Calibration)
- Registration

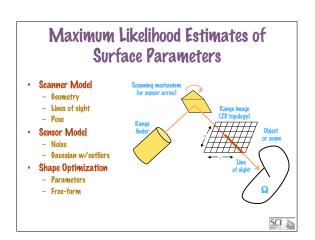




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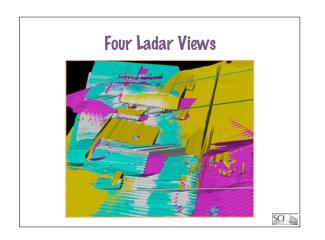
Processing Measured Pata A Systems Approach Noise (Stochastie) Weasured Data Noise (Stochastie) X A(X) B(X) Bayesian reconstruction: maximizing the posterior - likelihood (data) + prior -> MAP

MAP PET Reconstruction (R. Leahy, USC) Linear Reconstruction MAP Reconstruction

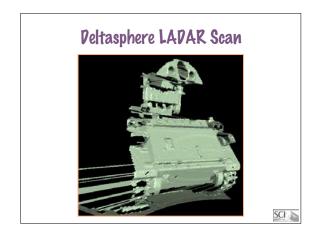




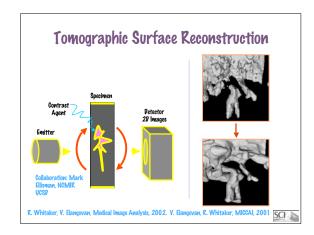












Summary

- · PDEs and geometric signal processing
- · New modeling processing capabilities
- Real applications (real data) running at interactive rates

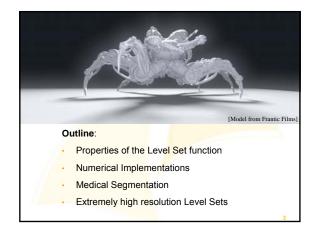
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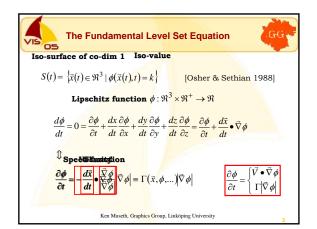
Thanks!

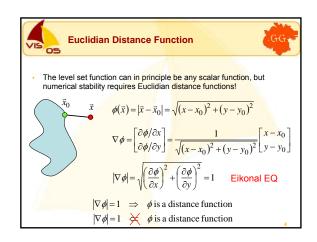
- Sponsors:
 - ONR, NSF, NIH, Exxon-Mobil
- Students and Colleagues:
 - Ernesto Juarez-Valdes, Stan Osher, Tolga Tasdizen, Vidya Elangovan, Paul Burchard, Suyash Awate, Won-Ki Jeong

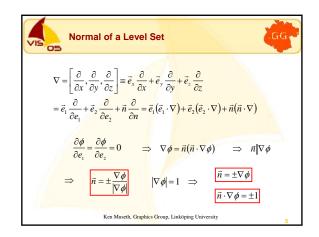
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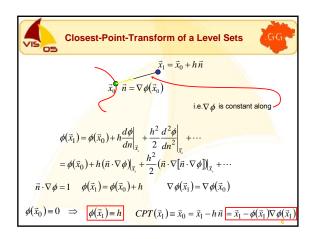


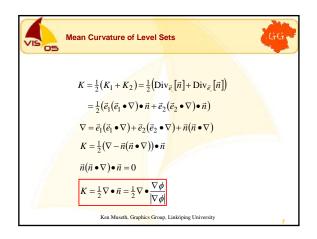


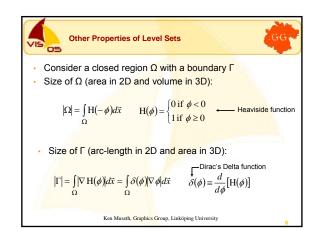


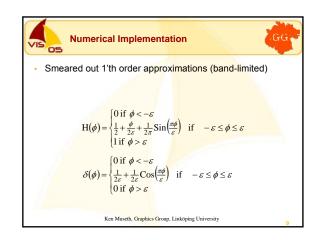


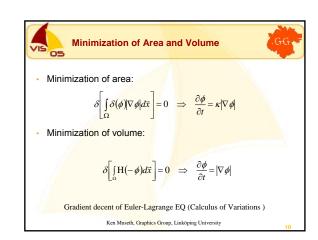


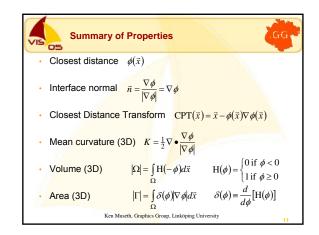


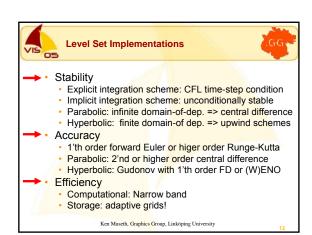


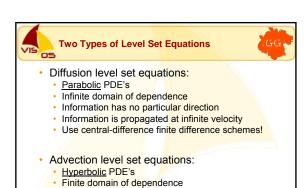






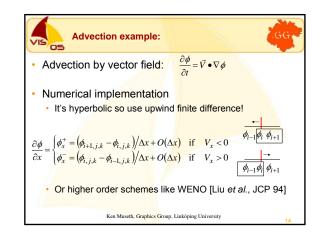


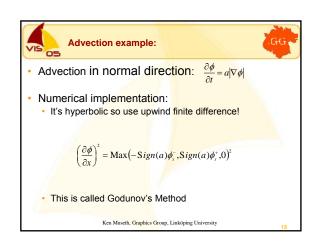


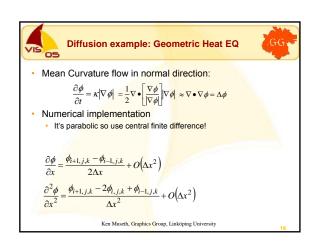


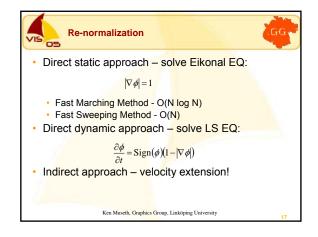
Information is propagated at finite velocity Use upwind finite difference schemes!

Information has a direction

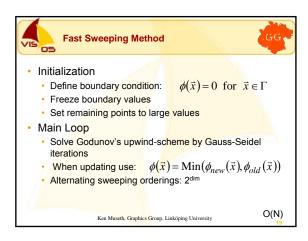


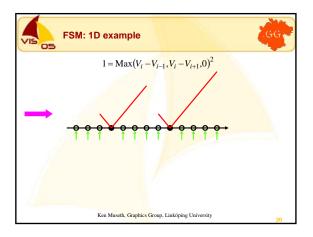


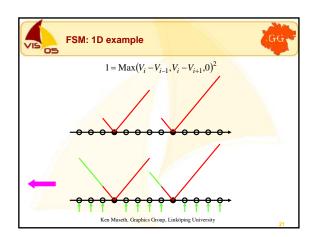


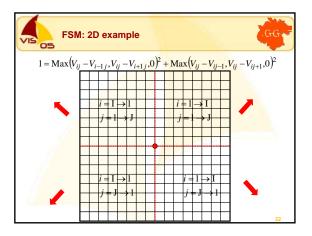


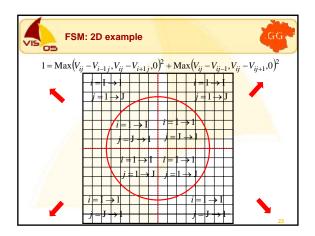


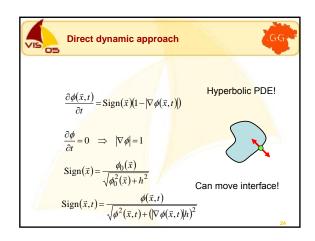


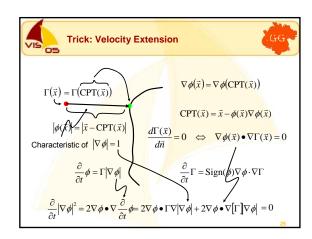


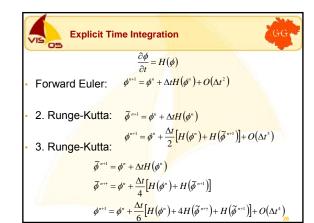


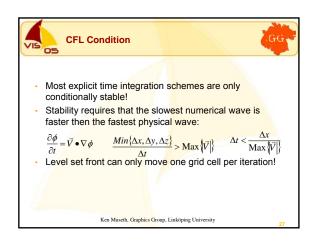


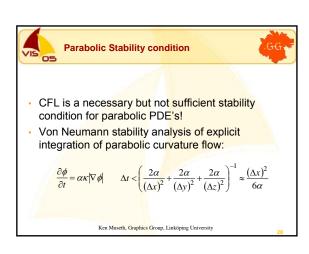


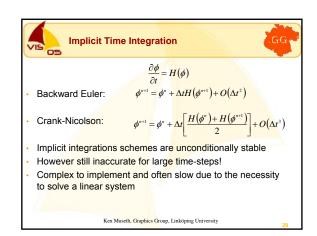


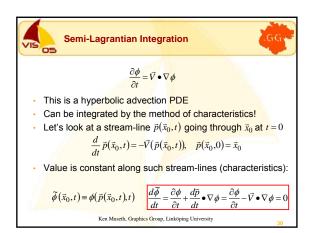


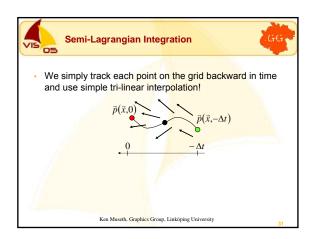


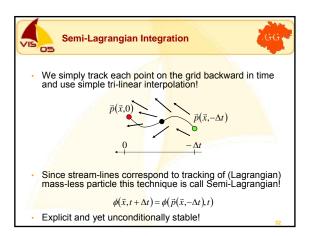


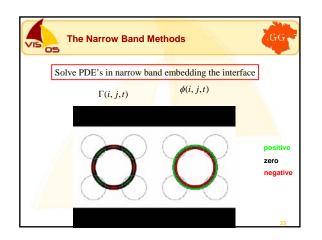


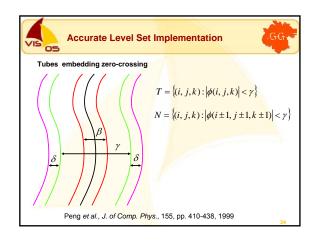


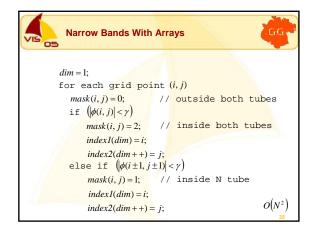


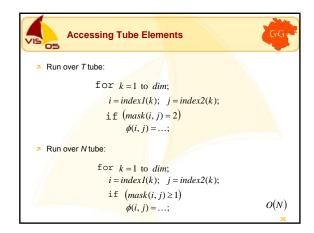


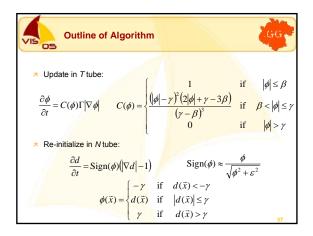


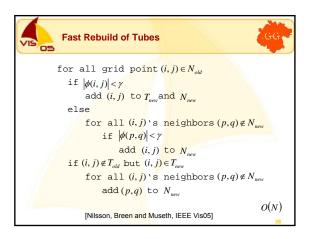


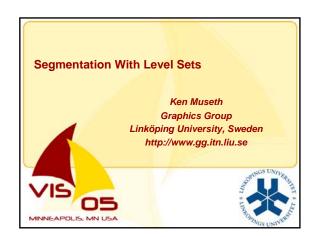


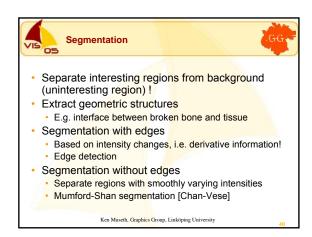


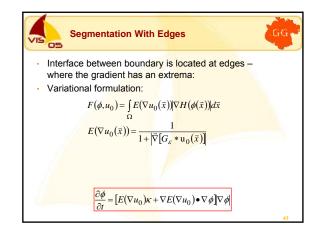


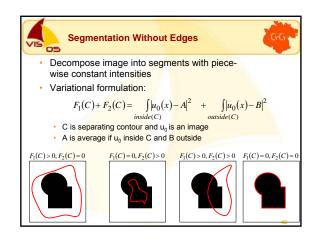


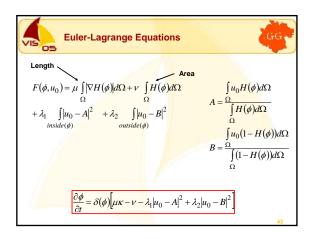


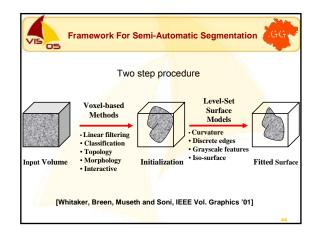


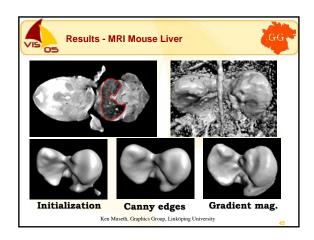


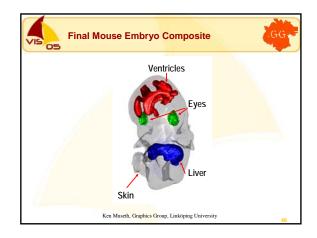


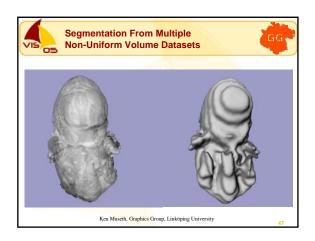


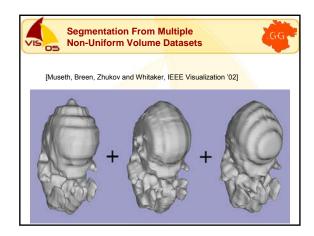


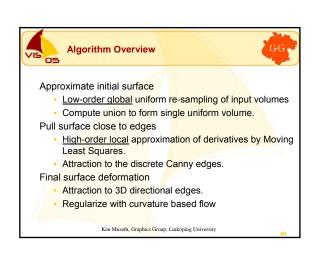


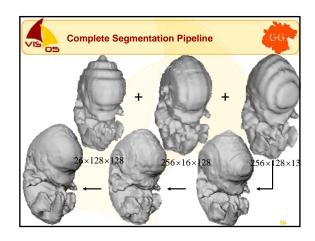


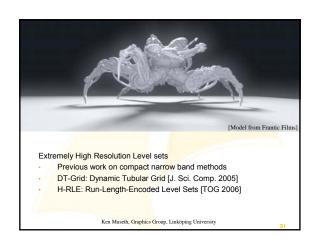


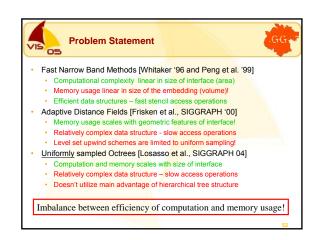


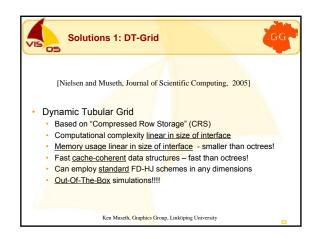


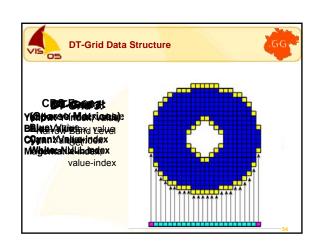


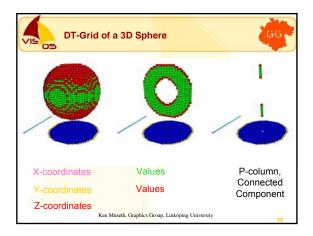


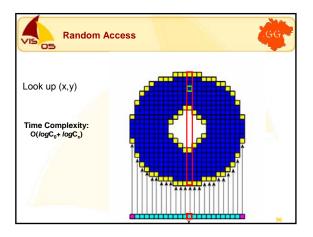


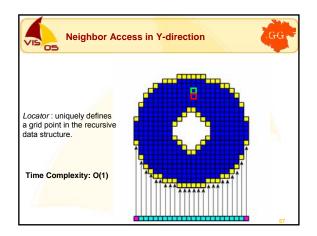


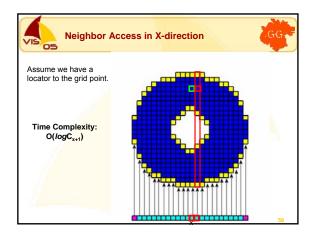


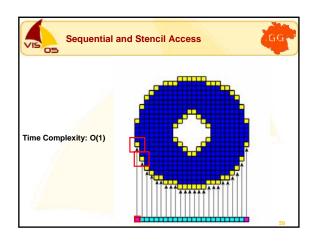


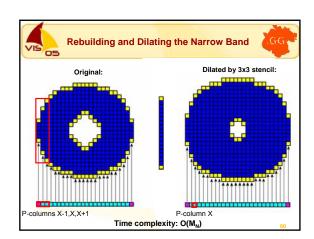


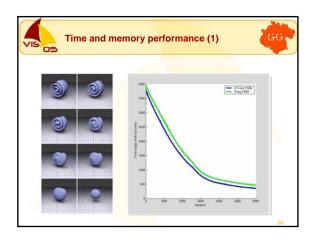


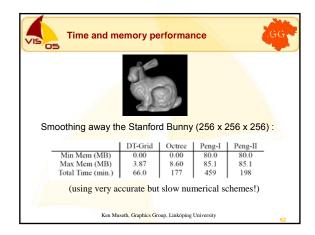






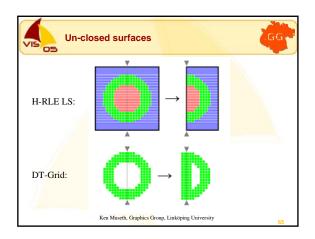


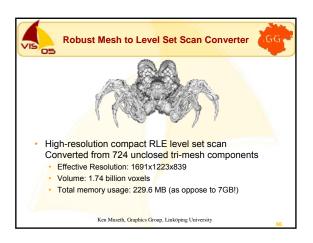




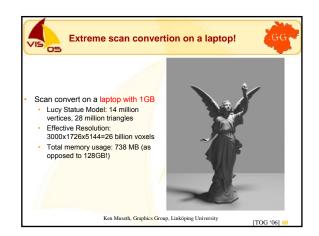




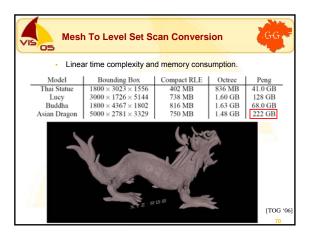


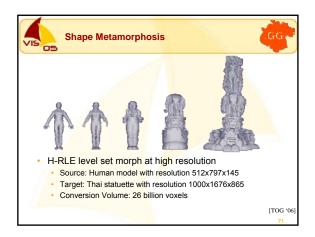


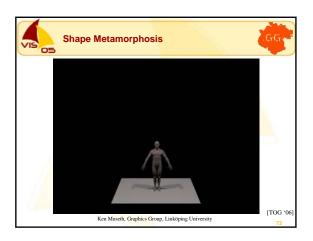


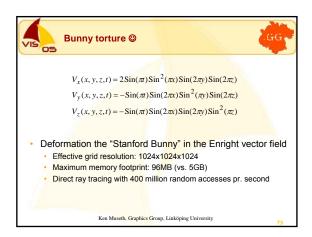






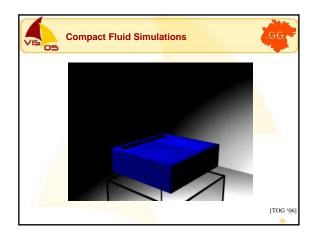


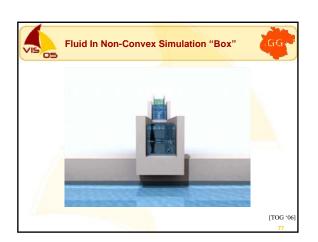








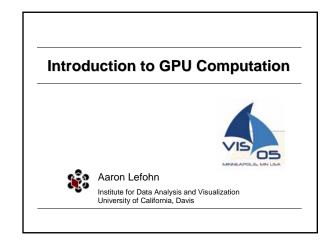


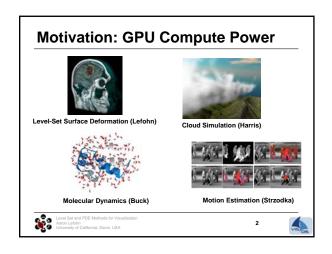


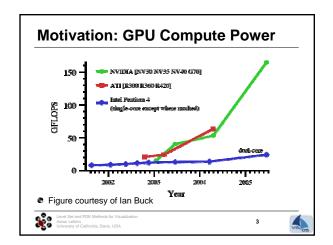


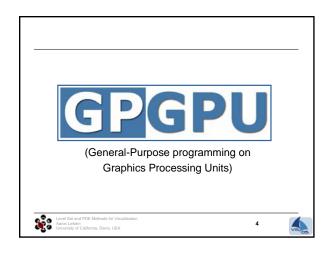


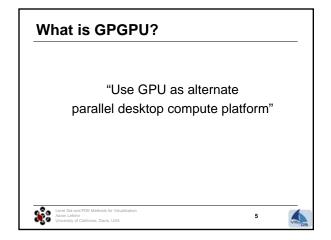


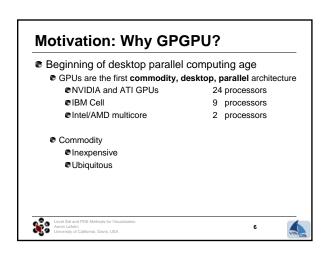


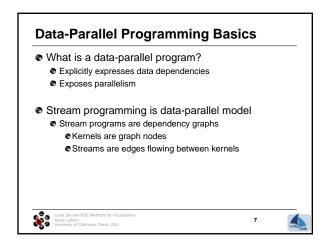


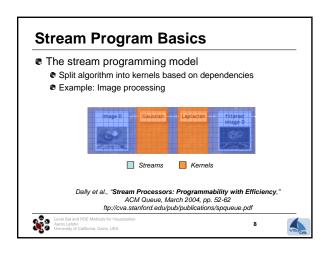


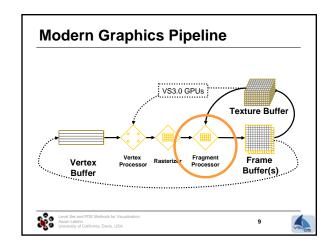


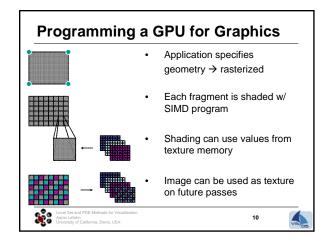


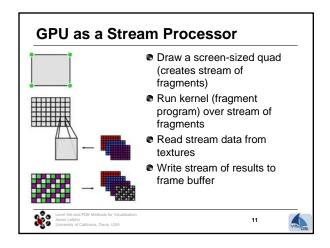


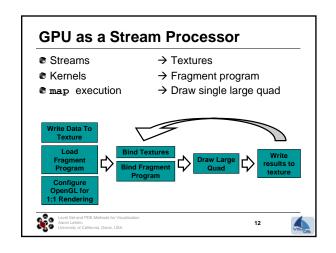












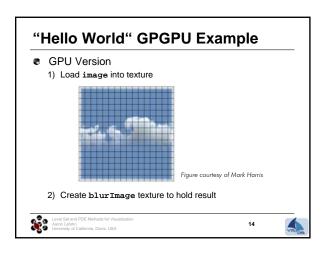
```
"Hello World" GPGPU Example

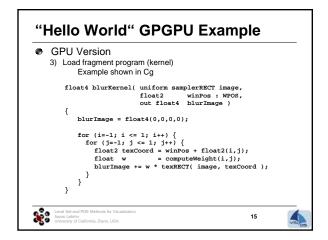
3 x 3 convolution
CPU version

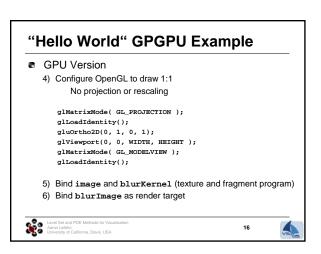
image = loadImage( WIDTH, HEIGHT );
blurImage = allocZeros( WIDTH, HEIGHT );

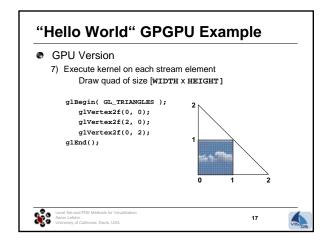
for (x=0; x < WIDTH; x++)
    for (y=0; y < HEIGHT; y++)
    for (j=-1; i <= 1; i++)
        for (j=-1; j <= 1; j++)
        float w = computeWeight(i,j);
        blurImage[x][y] += w * image[x+i, y+j];

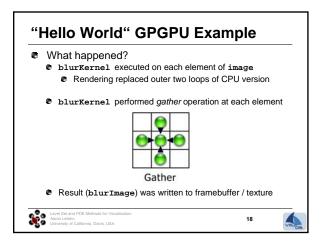
Level Set and PDE Methods for Visualization
Average Labora, Dave, USA
```

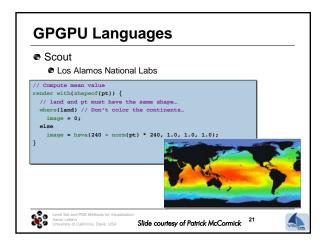


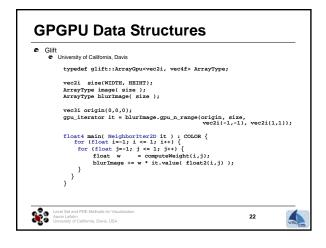










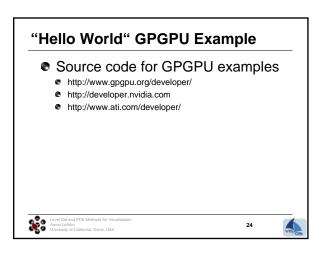


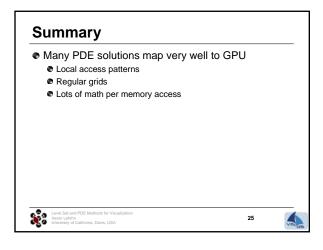
```
For Further GPGPU Information

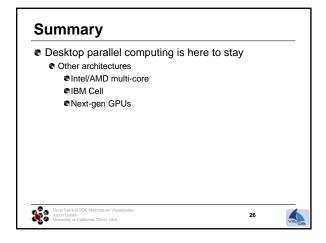
Vivrant research community

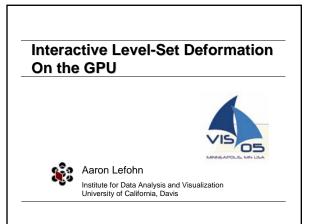
Two major focuses
Application-specific results
New programming models for data-parallel computation

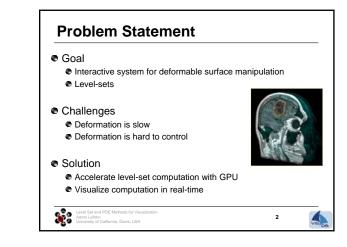
http://www.qpgpu.org/
Paper, forums, source code examples
ACM SIGGRAPH and IEEE Visualization course notes
GPU Gems II
Eurographics 2005 STAR report
```



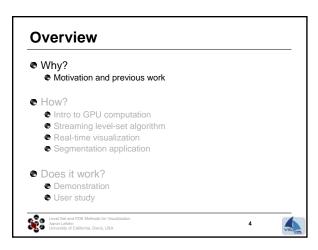


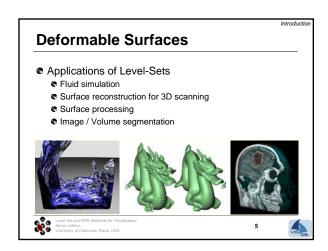


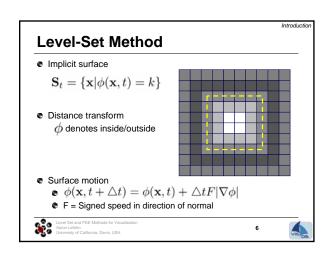


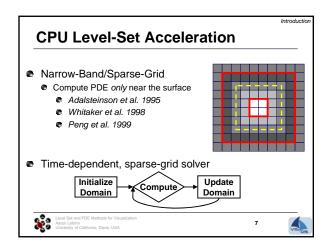


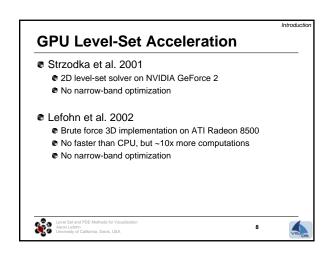


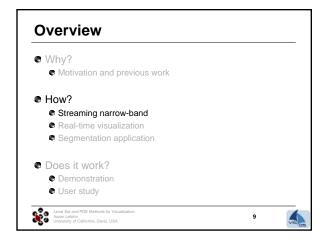


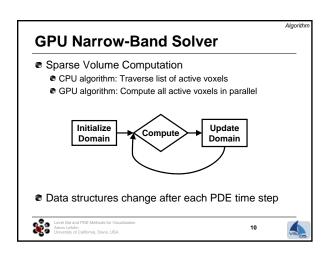


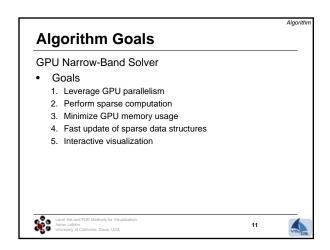


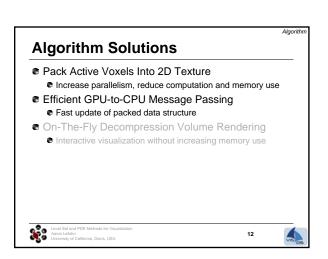


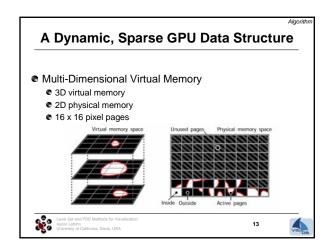


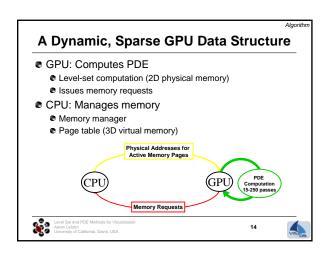


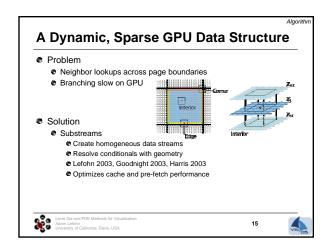


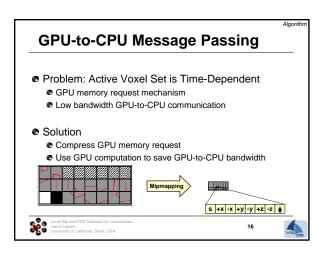




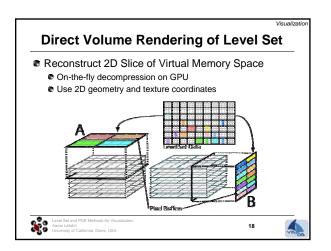


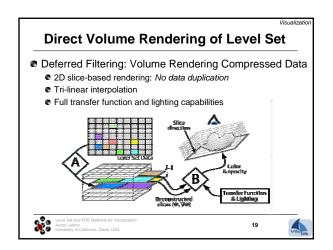


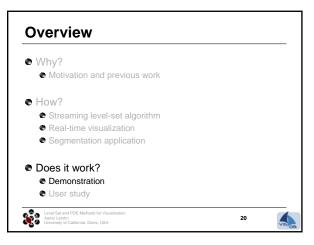


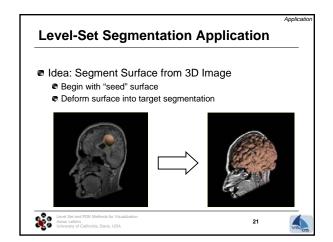


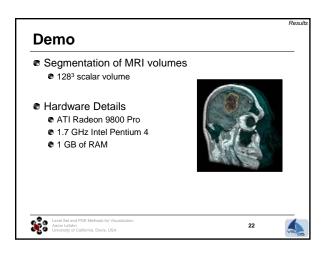


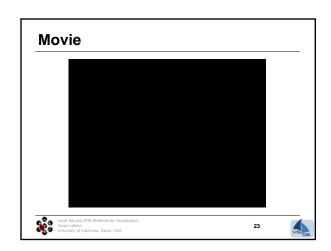


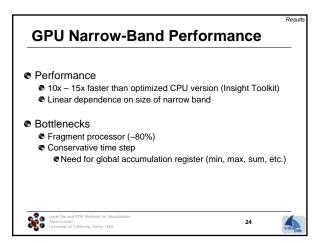


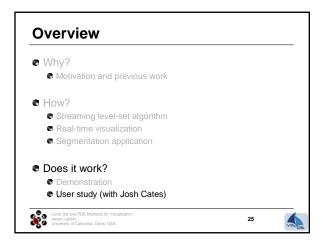


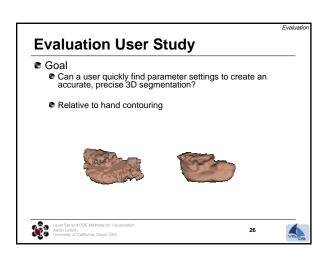


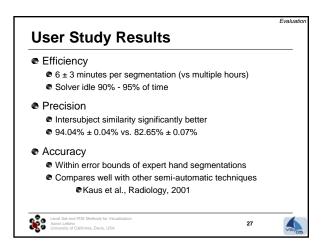


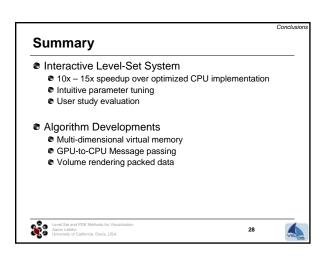


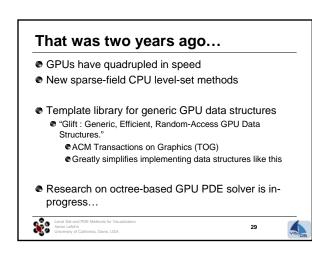


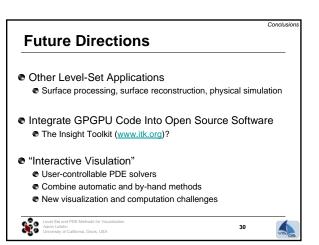












Acknowledgements

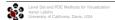
- Volume rendering Joe Kniss Tumor user study
 "Teem" raster-data toolkit Josh Cates

Gordon Kindlmann

- "GLEW" OpenGL extension wrangler

Ross Whitaker, Charles Hansen, Steven Parker and John Owens
 ATI: Evan Hart, Mark Segal, Jeff Royle, and Jason Mitchell
 Brigham and Women's Hospital

National Science Foundation Graduate Fellowship
 Office of Naval Research grant #N000140110033
 National Science Foundation grant #ACl008915 and #CCR0092065



Questions?

For More Information

Google "Lefohn level set" http://graphics.cs.ucdavis.edu/~lefohn/

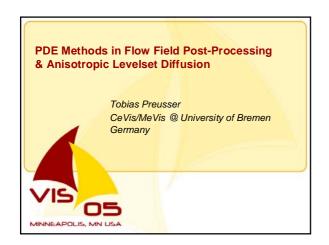
Journal Papers Based on this Work

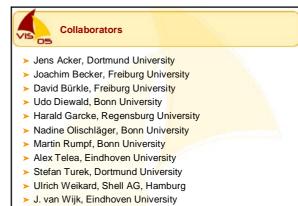
Lefohn, Kniss, Hansen, Whitaker, "A Streaming Narrow Band Algorithm: Interactive Computation and Visualization of Level Sets," IEEE Transactions on Visualization and Computer Graphics, 10 (40), Jul / Aug, pp. 422-433, 2004

Cates, Lefohn, Whitaker, "GIST: An Interactive, GPU-Based Level-Set Segmentation Tool for 3D Medical Images," Medical Image Analysis, to appear 2004









PDE methods in flow field post-processing

Scale space concept of image processing

Anisotropic nonlinear diffusion for flow visualization

Implementation

Phase separation for flow clustering

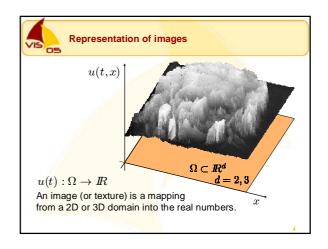
Anisotropic levelset diffusion

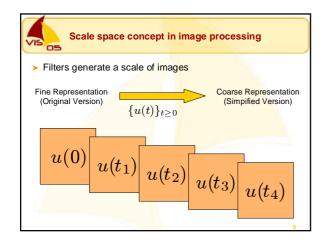
Introduction

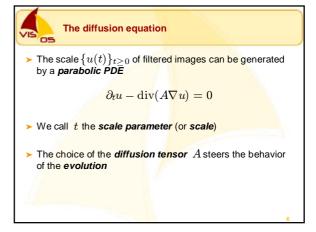
Anisotropic levelset diffusion

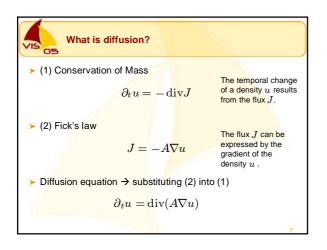
The importance of regularization

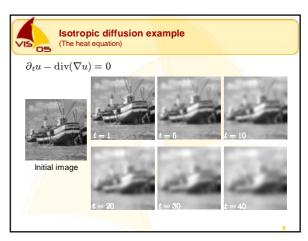
Evaluation of curvature

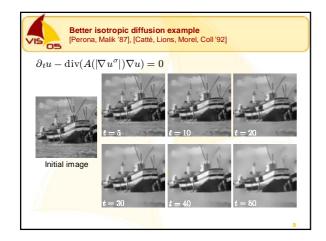


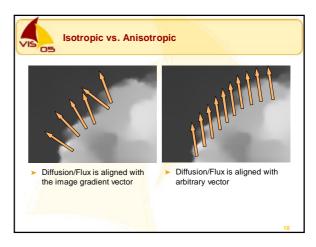


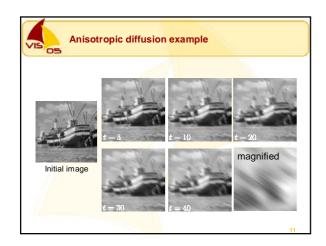


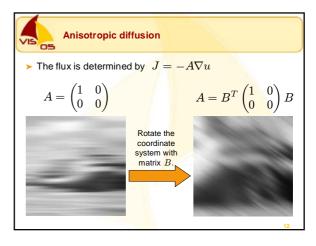


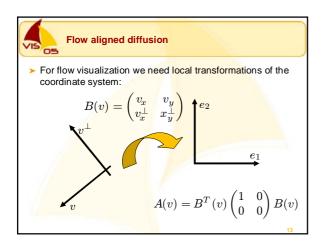


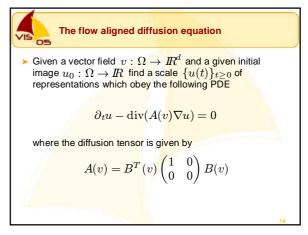


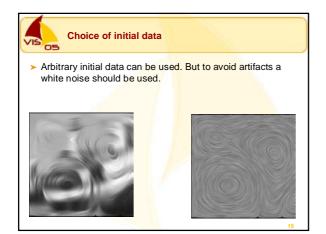


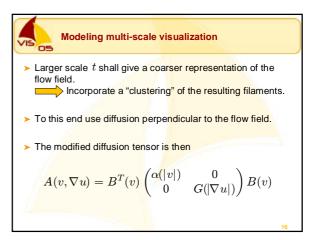


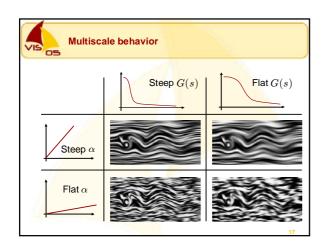


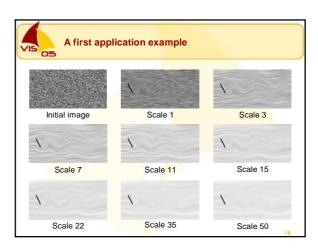


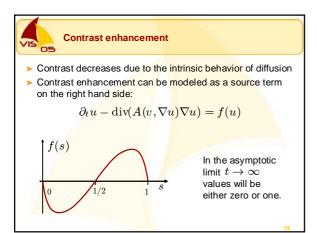


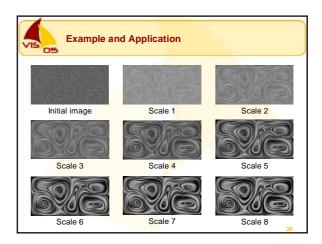














Characterization and comparison of the results

- Generates streamlines
- ➤ Incorporates a continuous scaling possibility
- ➤ Anisotropic diffusion is an asymptotic limit of LIC (Line integral convolution) (cf. [Cabral, Leedom '93], [Stalling, Hege '95] and many others ...)
- Anisotropic diffusion is a parallel version of Spot Noise (cf. [van Wijk '91], [de Leeuw, van Wijk '95] and others ...)



FEM-Discretization I

- Backward difference quotient in scale variable:

$$\underbrace{\frac{u^n - u^{n-1}}{\sqrt{1 - 1}}}_{\approx \partial_t u(n\tau)} - \operatorname{div}(A^{n-1} \nabla u^n) = f(u^{n-1})$$

Spatial discretization: Finite Element Methods (9-pt stencil) can resolve anisotropic diffusion better than Finite Difference Methods (5-pt stencil)



FEM-Discretization II

- Mesh consisting of quadrilaterals or hexahedrals (node of mesh = pixel of output texture)
- ➤ Piecewise bi- or trilinear function space (corresponding to bi- or trilinear interpolation in 2D resp. 3D)

$$u^{n}(x) = \sum_{i=0}^{N} (\vec{u}^{n})_{i} \phi_{i}(x)$$



➤ Computation of one step of the anisotropic diffusion means solving a system of equations for the nodal function values

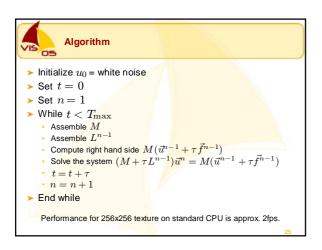
$$(M + \tau L^{n-1})\vec{u}^n = M(\vec{u}^{n-1} + \tau \vec{f}^{n-1})$$

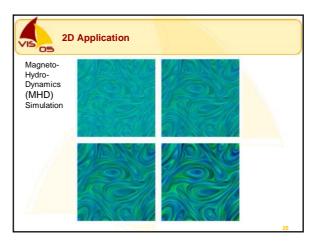
Mass- and stiffness-matrices

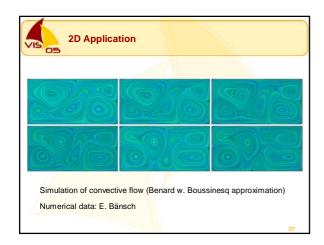
- The shape functions ϕ_i determine the image space and the matrices involved into the system of equations.
- Mass matrix

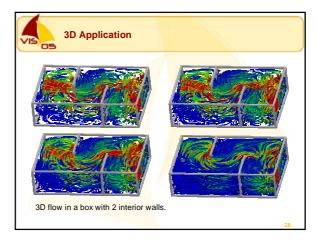
$$M_{ij} = \int_{\Omega} \phi_i(x) \phi_j(x) dx$$

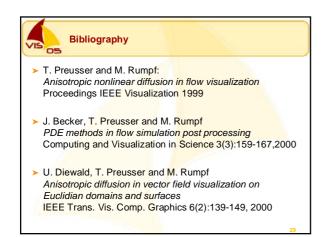
> Stiffness matrix
$$L_{ij}^n = \int_{\Omega} B^T(v) \begin{pmatrix} \alpha & 0 \\ 0 & G^n \end{pmatrix} B(v) \nabla \phi_i(x) \cdot \nabla \phi_j(x) \, dx \\ = \int_{\Omega} \begin{pmatrix} \alpha & 0 \\ 0 & G^n \end{pmatrix} B(v) \nabla \phi_i(x) \cdot B(v) \nabla \phi_j(x) \, dx$$

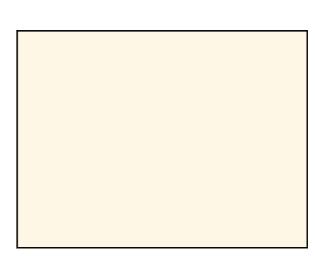


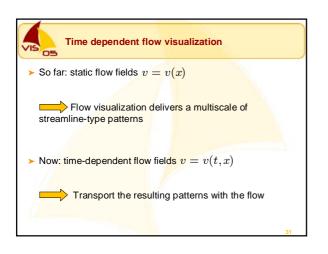


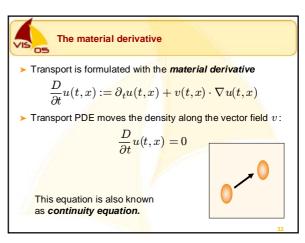


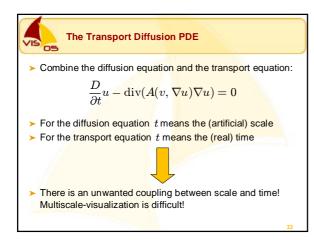


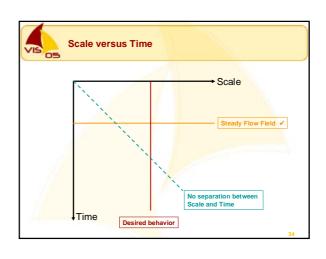


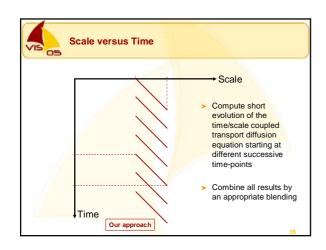


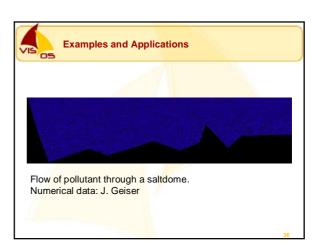


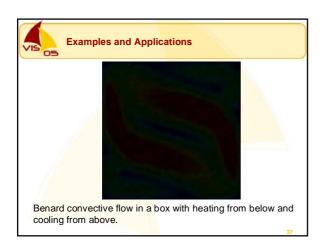


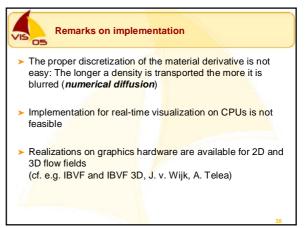


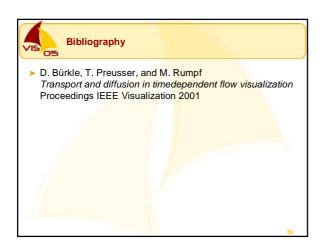


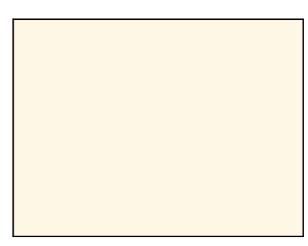












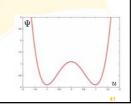


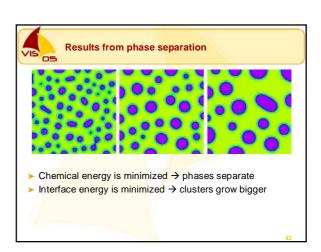
Clustering with Phase Separation (Cahn-Hilliard model)

Phase separation and coarsening in (binary) metal alloys is described by minimizing the Energy

$$E(u) := \int_{\Omega} \left\{ \Psi(u) + \frac{\gamma}{2} |\nabla u|^2 \right\}$$

- Ψ describes the chemical energy (double-well potential)
- $|\nabla u|^2$ describes the interface energy between the phases







Anisotropic interface energy

- ► The interface energy is given by $|\nabla u|^2$
- > All interfaces are penalized equally



ightharpoonup Anisotropic energy penalizing interfaces which are not aligned with a given flow field: $A(v)
abla u \cdot \nabla u$

$$A(v) = B(v)^T \begin{pmatrix} 1 & 0 \\ 0 & \beta \end{pmatrix} B(v) \qquad \beta \ge 1$$



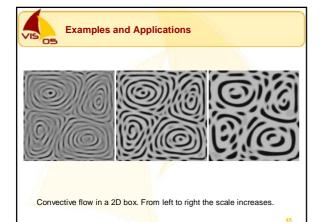
Anisotropic phase separation

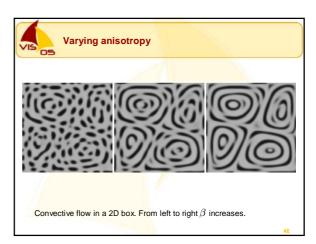
➤ Minimize the anisotropic energy

$$E(u) := \int_{\Omega} \left\{ \Psi(u) + \frac{\gamma}{2} A(v) \nabla u \cdot \nabla u \right\}$$

The first variation of energy and a gradient descent approach leads to the PDE

$$\partial_t u - \Delta(\Psi'(u) - \gamma \operatorname{div}(A(v)\nabla u)) = 0$$





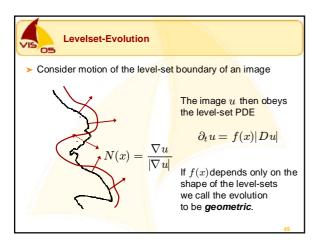


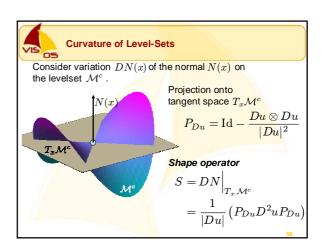
Examples and Applications

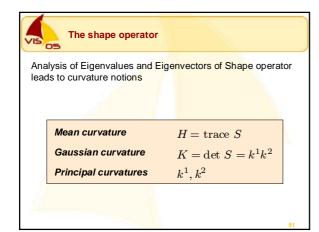
➤ H. Garcke, T. Preusser, M. Rumpf, A. Telea, U. Weikard and J, van Wijk

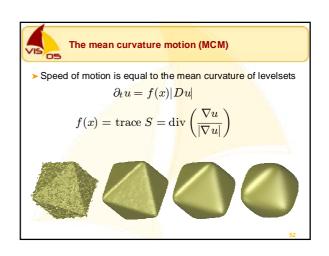
A Phase Field Model for Continuous Clustering on Vector

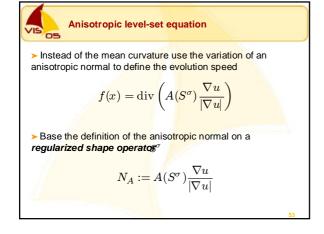
IEEE Trans. Vis. Comp. Graphics, 7, 230-241, 2001.

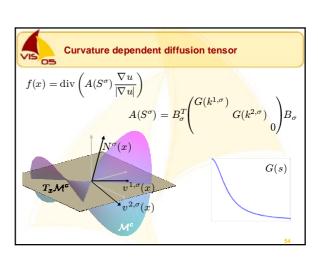


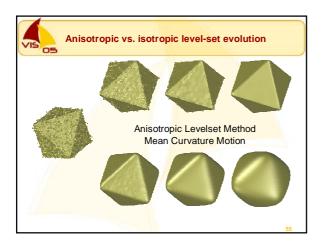


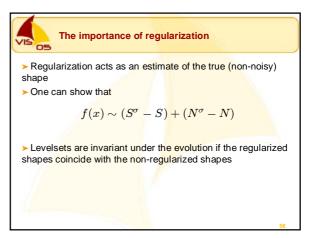


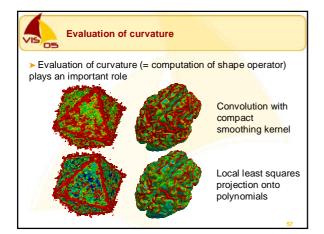


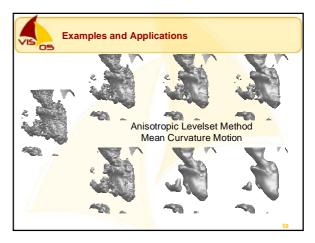


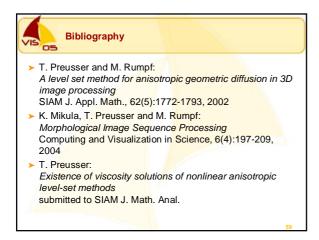












Anisotropic Diffusion in Vector Field Visualization on Euclidean Domains and Surfaces

Udo Diewald, Tobias Preußer, and Martin Rumpf

Abstract—Vector field visualization is an important topic in scientific visualization. Its aim is to graphically represent field data on two and three-dimensional domains and on surfaces in an intuitively understandable way. Here, a new approach based on anisotropic nonlinear diffusion is introduced. It enables an easy perception of vector field data and serves as an appropriate scale space method for the visualization of complicated flow pattern. The approach is closely related to nonlinear diffusion methods in image analysis where images are smoothed while still retaining and enhancing edges. Here, an initial noisy image intensity is smoothed along integral lines, whereas the image is sharpened in the orthogonal direction. The method is based on a continuous model and requires the solution of a parabolic PDE problem. It is discretized only in the final implementational step. Therefore, many important qualitative aspects can already be discussed on a continuous level. Applications are shown for flow fields in 2D and 3D, as well as for principal directions of curvature on general triangulated surfaces. Furthermore, the provisions for flow segmentation are outlined.

Index Terms—Flow visualization, multiscale, nonlinear diffusion, segmentation.

1 Introduction

The visualization of field data, especially of velocity fields from CFD computations, is one of the fundamental tasks in scientific visualization. A variety of different approaches has been presented. The simplest method of drawing vector plots at nodes of some overlaid regular grid in general produces visual clutter because of the typically different local scaling of the field in the spatial domain, which leads to disturbing multiple overlaps in certain regions, whereas, in other areas, small structures such as eddies cannot be resolved adequately. This gets even worse if tangential fields on highly curved surfaces are considered.

The central goal is to come up with intuitively better receptible methods which give an overall, as well as a detailed, view on the flow patterns. Single particle lines only partially enlighten features of a complex flow field. Thus, we want to define a texture which represents the field globally on a 2D or 3D domain and on surfaces, respectively. Here, we confine ourselves to stationary fields. In the Euclidean case, we suppose $v:\Omega\to\mathbb{R}^n$ for some domain $\Omega \subset \mathbb{R}^n$, whereas, in the case of a manifold \mathcal{M} embedded in \mathbb{R}^3 , we consider a tangential vector field v. We ask for a method generating stretched streamline type patterns which are aligned to the vector field v(x). Furthermore, the possibility of successively coarsening this pattern is obviously a desirable property. Methods which are based on such a scale of spaces and enhance certain structures of images are well-known in image processing analysis.

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Actually, nonlinear diffusion allows the smoothing of gray or color images while retaining and enhancing edges [18]. Now, we set up a diffusion problem, with strong smoothing along integral lines and edge enhancement in the orthogonal directions. Applying this to some initial random noise image intensity, we generate a scale of successively coarser patterns which represent the vector field. Finite elements in space and a semi-implicit time stepping are applied to solve this diffusion problem numerically. Furthermore, a suitable modification of the approach allows the identification of topological regions.

Before we explain in detail the method, let us discuss related work on vector field visualization and image processing. Later on we will identify some of the well-known methods as equivalent to special cases or asymptotic limits of the presented new method, respectively.

2 RELATED WORK

The spot noise method proposed by van Wijk [25] introduces spot-like texture splats which are aligned by deformation to the velocity field in 2D or on surfaces in 3D. These splats are plotted in the fluid domain, showing strong alignment patterns in the flow direction. The original first order approximation to the flow was improved by de Leeuw and van Wijk in [6] by using higher order polynomial deformations of the spots in areas of significant vorticity. In an animated sequence, these spots can be moved along streamlines of the flow. Furthermore, in 3D, van Wijk [26] applies the integration to clouds of oriented particles and animates them by drawing similar moving transparent and illuminated splats.

The Line Integral Convolution (LIC) approach of Cabral and Leedom [4] integrates the fundamental ODE describing streamlines forward and backward in time at every

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pixelized point in the domain, convolves a white noise along these particle paths with some Gaussian type filter kernel, and takes the resulting value as an intensity value for the corresponding pixel. According to the strong correlation of this intensity along the streamlines and the lack of any correlation in the orthogonal direction, the resulting texturing of the domain shows dense streamline filaments of varying intensity. Stalling and Hege [21] increased the performance of this method, especially by reusing portions of the convolution integral already computed on points along the streamline. Forssell [10] proposed a similar method on surfaces and Max et al. [17] discussed flow visualization by texturing on contour surfaces. Max and Becker [16] presented a method for visualizing 2D and 3D flows by animating textures.

Shen and Kao [20] applied an LIC type method to unsteady flow fields. Recently, a method [2] has been presented which generates streakline type patterns by numerical calculation of the transport of inlet coordinates and inlet position. Interrante and Grosch [12] generalized line integral convolution to 3D in terms of volume rendering of line filaments.

In [24], Turk and Banks discuss an approach which selects a certain number of streamlines. They are automatically equally distributed all over the computational domain to characterize, in a sketch-type representation, the significant aspects of the flow. An energy minimizing process is used to generate the actual distribution of streamlines.

Especially for 3D velocity fields, particle tracing is a very popular tool. But, a few particle integrations released by the user can hardly scope with the complexity of 3D vector fields. Stalling et al. [22] use pseudorandomly distributed, illuminated, and transparent streamlines to give a denser and more receptible representation, which shows the overall structure and enhances important details.

Van Wijk [27] proposed the implicit stream surface method. For a stationary flow field, the transport equations $v\cdot\nabla\phi=0$ are solved for given v and certain inflow and outflow boundary conditions in a precomputing step. Then, isosurfaces of the resulting function ϕ are streamsurfaces and can be efficiently extracted with interactive frame rates, even for larger data sets.

Most of the methods presented so far have in common, that the generation of a coarser scale requires a recomputation. For instance, if we ask for a finer or coarser scale of the line integral convolution pattern, the computation has to be restarted with a coarser initial image intensity. In the case of spot noise, larger spots have to be selected and their stretching along the field has to be increased. The approach to be presented here will incorporate a successive coarsening as time proceeds in the underlying diffusion problem.

As already mentioned in the introduction, our method of anisotropic nonlinear diffusion to visualize vector fields is derived from well-known image processing methodology. Discrete diffusion type methods have been known for a long time. Perona and Malik [18] introduced a continuous diffusion model which allows the denoising of images together with the enhancing of edges. Alvarez et al. [1] established a rigorous axiomatic theory of diffusive scale

space methods. Kawohl and Kutev [14] investigate a qualitative analysis of the Perona and Malik model. The recovering of lower dimensional structures in images is analyzed by Weickert [28], who introduced an anisotropic nonlinear diffusion method, where the diffusion matrix depends on the so-called structure tensor of the image. A finite element discretization and its convergence properties have been studied by Kacur and Mikula [13].

Concerning the application of diffusion type methods on surfaces, a general introduction to differential calculus on manifolds can be found for instance in the book by do Carmo [7]. Dziuk [8] presented an algorithm for the solution of partial differential equations on surfaces and, in [9], he discussed a numerical method for geometric diffusion applied to the surface itself which coincides with the mean curvature motion.

3 THE NONLINEAR DIFFUSION PROBLEM

Let us now derive our method based on a suitable PDE problem. At first, we confine ourselves to the case of planar domains in 2D and 3D. Here, nonlinear anisotropic diffusion applied to some initial random noisy image will enable an intuitive and scalable visualization of complicated vector fields. Therefore, we pick up the idea of line integral convolution, where a strong correlation in the image intensity along integral lines is achieved by convolution of an initial white noise along these lines. As proposed already by Cabral and Leedom [4], a suitable choice for the convolution kernel is a Gaussian kernel. On the other hand, an appropriately scaled Gaussian kernel is known to be the fundamental solution of the heat equation. Thus, line integral convolution is nothing else than solving the heat equation in 1D on an integral line parameterized with respect to arc length. On pixels which are located on different integral lines, the resulting image intensities are not correlated. Hence, the thickness of the resulting image patterns in line integral convolution is of the size of the random initial patterns, in general, a single pixel. Increasing this size, as has been proposed by Kiu and Banks [15], leads to broader stripes and, unfortunately, less sharp transitions across streamline patterns. As described so far, line integral convolution is a discrete pixel-based method. If we ask for a well-posed continuous diffusion problem with similar properties, we are led to some anisotropic diffusion, now controlled by a suitable diffusion matrix.

To begin with, let us at first introduce a general nonlinear diffusion method from image processing and then discuss the selection of the appropriate diffusion tensor and the righthand side. Here, we consider first the case of an image in Euclidean space either in 2D or 3D. In Section 6, we then generalize this with respect to textures on surfaces. We consider a function $\rho: \mathbb{R}_0^+ \times \Omega \to \mathbb{R}$ which solves the parabolic problem

$$\begin{split} \frac{\partial}{\partial t} \rho - \operatorname{div}(A(\nabla \rho_{\epsilon}) \nabla \rho) &= f(\rho) & \text{ in } \mathbb{R}^{+} \times \Omega, \\ \rho(0, \cdot) &= \rho_{0} & \text{ on } \Omega, \\ \frac{\partial}{\partial \nu} \rho &= 0 & \text{ on } \mathbb{R}^{+} \times \partial \Omega. \end{split}$$

for given initial density $\rho_0: \Omega \to [0,1]$. Here, $\rho_\epsilon = \chi_\epsilon * \rho$ is a mollification of the current density, which will later on turn

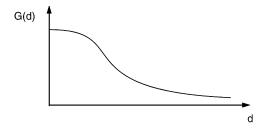


Fig. 1. The shape of $G(\cdot)$ which, applied to the gradient of the mollified image intensity, serves as a diffusion coefficient in image processing.

out to be necessary for the well-posedness of the above parabolic, boundary, and initial value problem. In our setting, we interpret the density as an image intensity, a scalar grayscale or-with a slight extension to the vector valued case—as a vector valued color. Thus, the solution $\rho(\cdot)$ can be regarded as a family of images $\{\rho(t)\}_{t\in\mathbb{R}_{0}^{+}}$, where the time t serves as a scaling parameter. Let us remark that, by the trivial choice A = 1 and $f(\rho) = 0$, we obtain the standard linear heat equation with its isotropic smoothing effect. In image processing, ρ_0 is a given noisy initial image. The diffusion is supposed to be controlled by the gradient of the image intensity. Large gradients mark edges in the image which should be enhanced, whereas small gradients indicate areas of approximately equal intensity. Here, denoising, i.e., intensity diffusion, is considered. For that purpose we prescribe a diffusion coefficient

$$A = G(\|\nabla \rho_{\epsilon}\|),$$

where $G: \mathbb{R}_0^+ \to \mathbb{R}^+$ is a monotone decreasing function with $\lim_{d\to\infty}G(d)=0$ and $G(0)=\beta$, where $\beta\in\mathbb{R}^+$ is constant (cf. Fig. 1), e. g. $G(d)=\frac{\beta}{1+\|d\|^2}$. If we would replace the mollified gradient $\nabla\rho_\epsilon$ as argument of G by the true gradient $\nabla \rho$, which leads to the original Perona Malik model, we would, in general, obtain a backward parabolic problem in areas of high gradients which is no longer wellposed [14]. The invoked mollification avoids this shortcoming and comes along with a desirable presmoothing effect. Nevertheless, the enhancing of steep gradients and, thereby, edges in the image, known from backward diffusion, is retained if we adjust the mollification carefully. A suitable choice [13] for this mollification is a convolution with the heat equation kernel, i.e., we define $\rho_{\epsilon} = \tilde{\rho}(t =$ $\epsilon^2/2$) where $\tilde{\rho}$ is the solution of the heat equation with initial data ρ . Then, ϵ is the width of the corresponding Gaussian filter. Fig. 2 gives an example of such an image smoothing and edge enhancement by nonlinear diffusion. The function $f(\cdot)$ may serve as a penalty which forces the scale of images to stay close to the initial image, e.g., choosing $f(\rho) = \gamma(\rho_0 - \rho)$, where γ is a positive constant.

Now, we incorporate anisotropic diffusion. For a given vector field $v:\Omega\to\mathbb{R}^n$, we consider linear diffusion in the direction of the vector field and a Perona Malik type diffusion orthogonal to the field. Let us suppose that v is continuous and $v\neq 0$ on Ω . Then, there exists a family of continuous orthogonal mappings $B(v):\Omega\to SO(n)$ such that $B(v)v=\|v\|e_0$, where $\{e_i\}_{i=0,\cdots,n-1}$ is the standard base



Fig. 2. The noisy image on the left is successively smoothed by nonlinear diffusion. On the right the resulting smoothed image with enhanced edges is shown.

in \mathbb{R}^n (cf. Fig. 3). We consider a diffusion matrix $A = A(v, \nabla \rho_{\epsilon})$ and define

$$A(v,d) = B(v)^T \begin{pmatrix} \alpha(\|v\|) & \\ G\|d\| \mathrm{Id}_{n-1} \end{pmatrix} B(v),$$

where $\alpha: \mathbb{R}^+ \to \mathbb{R}^+$ controls the linear diffusion in the vector field direction, i.e., along streamlines, and the edge-enhancing diffusion coefficient $G(\cdot)$ introduced above acts in the orthogonal directions. Here, Id_{n-1} is the identity matrix in dimension n-1. We may either choose a linear function α or, in the case of a velocity field which spatially varies over several orders of magnitude, we select a monotone function α (cf. Fig. 4) with

$$\alpha(0) > 0$$
 and $\lim_{s \to \infty} \alpha(s) = \alpha_{\text{max}}.$

In general, it does not make sense to consider a certain initial image. As initial data ρ_0 , we thus choose some random noise of an appropriate frequency range. This can, for instance, be generated by running a linear isotropic diffusion simulation on a discrete white noise for a short time. Hence, patterns will grow upstream and downstream, whereas the edges tangential to these patterns are successively enhanced. Still, there is some diffusion perpendicular to the field which supplies us for evolving time with a scale of progressively coarser representation of the flow field. If we run the evolution for vanishing righthand side f, the image contrast will, unfortunately, decrease due to the diffusion along streamlines. The asymptotic limit would turn out to be an averaged gray value. Therefore, we strengthen the image contrast during the evolution, selecting an appropriate function $f:[0,1] \to \mathbb{R}$ (cf. Fig. 4) with

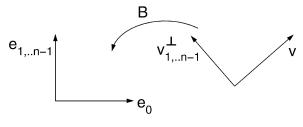


Fig. 3. The coordinate transformation B(v).

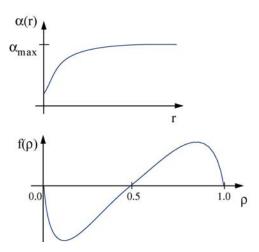


Fig. 4. The graphs of the velocity dependent linear diffusion $\alpha(\cdot)$, respectively, the scalar contrast enhancing right hand side $f(\cdot)$.

$$\begin{split} f(0) &= f(1) = 0 \\ f &> 0 \text{ on } (0.5, 1) \text{ and } f < 0 \text{ on } (0, 0.5). \end{split}$$

If we—at first glance—neglect the diffusive term in the equation, one realizes that perturbations below the average value 0.5 are pushed toward the zero value and, accordingly, values above 0.5 are pushed toward 1. Well-known maximum principles ensure that the interval of gray values [0,1] is not enlarged running the nonlinear diffusion. Here, the first property of f is of great importance. Finally, we end up with the method of nonlinear anisotropic diffusion to visualize complex vector fields. We thereby solve the nonlinear parabolic problem

$$\frac{\partial}{\partial t}\rho - \operatorname{div}(A(v, \nabla \rho_{\epsilon})\nabla \rho) = f(\rho),$$

starting from some random initial image ρ_0 , and obtain a scale of images representing the vector field in an intuitive way (cf. Fig. 5).

The corresponding variational formulation is obviously given by

$$(\partial_t \rho, \theta) + (A(v, \nabla \rho_{\epsilon}) \nabla \rho, \nabla \theta) = (f, \theta),$$

for all $\theta \in C^{\infty}(\Omega)$, where (.,.) denotes the L^2 product on the domain Ω . Our later finite element implementation will be based on this formulation by restriction to finite dimensional function spaces.

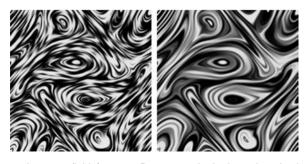


Fig. 5. A vector field from a 2D magneto-hydrodynamics simulation (MHD) is visualized by nonlinear diffusion. A discrete white noise is considered as initial data. We run the evolution on the left for a small and, on the right, for a large constant diffusion coefficient α .

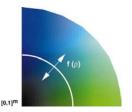


Fig. 6. A sketch of the vector valued contrast enhancing function f which leads to asymptotic states $\rho(\infty,\cdot)\in\{0\}\cup(S^{m-1}\cap[0,1]^m)$. Here, the components of the density are interpreted as blue, respectively green, color values. The arrows indicate the direction of contrast enhancement.

4 COUPLED SYSTEM OF DIFFUSION EQUATIONS

If we ask for pointwise asymptotic limits of the evolution, we expect an almost everywhere convergence to $\rho(\infty,\cdot) \in$ $\{0,1\}$ due to the choice of the contrast enhancing function $f(\cdot)$. Analytically, 0.5 is a third, but unstable, fix point of the dynamics. Thus, numerically, it will not turn out to be locally dominant. The space of asymptotic limits significantly influences the richness of the developing vector field aligned structures. We may ask how to further enrich the pattern which is settled by anisotropic diffusion. This turns out to be possible by increasing the set of asymptotic states. We no longer restrict ourselves to a scalar density ρ_r but consider a vector valued $\rho: \Omega \to [0,1]^m$ for some $m \ge 1$, and a corresponding system of parabolic equations. The coupling is given by the nonlinear diffusion coefficient $G(\cdot)$ which now depends on the norm $\|\nabla \rho\|$ of the Jacobian of the vector valued density $\nabla \rho$ and the righthand side $f(\cdot)$. We define

$$f(\rho) = h(\|\rho\|)\rho$$

with $h(s) = \tilde{f}(s)/s$ for $s \neq 0$, where \tilde{f} is the old righthand side from the scalar case and h(0) = 0. Furthermore, we select an initial density which is now a discrete "white" noise with values in $B_1(0) \cap [0,1]^m$. Thus, the contrast enhancing now pushes the pointwise vector density ρ either to the 0 or to some value on the sphere sector $S^{m-1} \cap [0,1]^m$ in \mathbb{R}^m (cf. Fig. 6). Again, a straightforward application of the maximum principle ensures $\rho(t,x) \in S^{m-1} \cap [0,1]^m$ for all t and t and t and t in t in

Fig. 7 shows an example for the application of the vector valued anisotropic diffusion method applied to a 2D flow field from a MHD simulation convective flow field. Furthermore, Fig. 8 shows results of this method applied to several time steps of a convective flow field. An incompressible Bénard convection is simulated in a rectangular box with heating from below and cooling from above. The formation of convection rolls will lead to an exchange of temperature. We recognize that the presented method is able to nicely depict the global structure of the flow field, including its saddle points, vortices, and stagnation points on the boundary. Fig. 9 shows results for the same data sets obtained by line integral convolution (here, we used the implementation of Stalling and Hege [21]). Finally, Fig. 10 shows a different application to a porous media flow field.

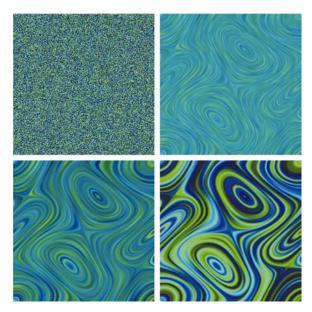


Fig. 7. Different snapshots from the multiscale based on anisotropic diffusion are depicted for a 2D MHD simulation vector field. Here, we consider a two-dimensional diffusion problem and interpret the resulting density as a color in a blue/green color space.

5 APPLICATION IN 3D

The anisotropic nonlinear diffusion problem has been formulated in Section 3 for arbitrary space dimension. It results in a scale of vector field aligned patterns which we then have to visualize. In 2D, this has already been done in a straightforward manner in the above figures. In 3D, we have somehow to break up the volume and open up the view to inner regions. Otherwise, we must confine ourselves with some pattern close to the boundary representing solely the shear flow.

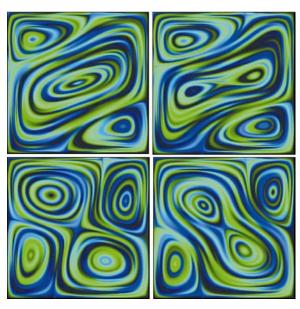


Fig. 8. Convective patterns in a 2D flow field are displayed and emphasized by the method of anisotropic nonlinear diffusion. The images show the velocity field of the flow at different time steps. The resulting alignment is thereby with respect to streamlines of this time dependent flow.



Fig. 9. LIC image generated for one of the data sets that have already been processed in Fig. 8 by nonlinear diffusion (cf. lower left image in Fig. 8).

Here, a further benefit of the vector valued diffusion comes into operation. We know that, for m=2, the asymptotic limits—which differ from 0—are, in mean, equally distributed on $S^1 \cap [0,1]^2$. Hence, we reduce the informational content and focus on a ball-shaped neighborhood $B_{\delta}(\omega)$ of a certain point $\omega \in S^1 \cap [0,1]^2$. Now, we can either look at isosurfaces of the function

$$\sigma(x) = \|\rho(x) - \omega\|^2,$$

where the isolevel δ^2 allows us to depict the boundary of the preimage of $B_{\delta}(\omega)$ with respect to the mapping ρ (cf. Fig. 11 and Fig. 12). Alternatively, we might use volume rendering to visualize this type of subvolumes. A detailed discussion of the latter approach is beyond the scope of this paper.

6 ANISOTROPIC DIFFUSION ON SURFACES

In the above sections, we have discussed anisotropic diffusion in vector field visualization on domains which are subsets of two and three-dimensional Euclidean space.

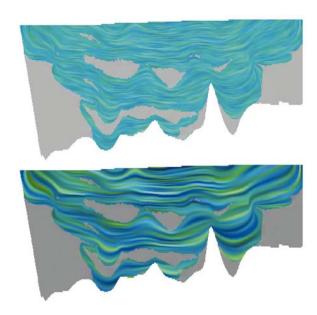


Fig. 10. Field aligned diffusion clearly outlines the principal features of a porous media flow in the vicinity of a salt dome. Lenses of lower permeability force the flow to pass through narrow bridges. We depict two time steps of the diffusion process.

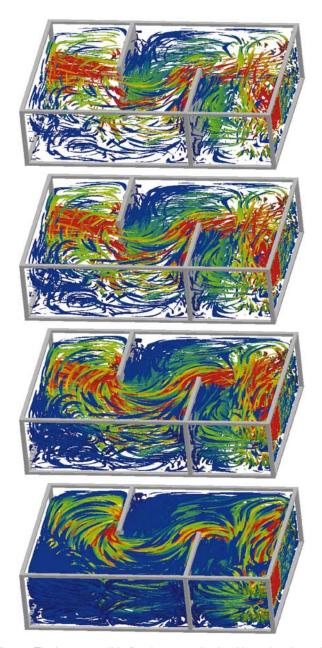


Fig. 11. The incompressible flow in a water basin with two interior walls and an inlet (on the left) and an outlet (on the right) is visualized by the anisotropic nonlinear diffusion method. Isosurfaces show the preimage of $\partial B_{\delta}(\omega)$ under the vector valued mapping ρ for some point ω on the sphere sector. From top to bottom, the radius δ is successively increased. A color ramp blue–green–red indicates an increasing absolute value of the velocity. The diffusion is applied to initial data, which is a relatively coarse grain random noise.

In what follows, we will outline how to carry over this methodology to display tangential vector fields on surfaces. Important examples are results from meteorological simulations, flow fields on streamsurfaces, or vector fields in differential geometry. The applications presented here will focus on the latter case and present multiscale textures on surfaces representing the principal directions of curvature. Based on the well-established intrinsic differential calculus on manifolds [7], we can pick up the same diffusion problems with an appropriate reinterpretation of the operators. Thus, let us first briefly review the basic notation



Fig. 12. Nonlinear anisotropic diffusion applied to the same 3D data set as in Fig. 11, but with a fine grain white noise as initial data.

of manifolds, differential calculus, and geometric diffusion. For a detailed introduction to geometry and differential calculus, we refer to [7] and [5, chapter 1]. For the sake of simplicity, we assume our surfaces to be compact embedded manifolds without boundary. Thus, we consider a smooth manifold \mathcal{M} , which we suppose to be embedded in \mathbb{R}^3 . Let $x:\Omega\to\mathcal{M}$; $\xi\mapsto x(\xi)$ be a coordinate map from an atlas of \mathcal{M} . For each point x on \mathcal{M} , the embedded tangent space $\mathcal{T}_x\mathcal{M}$ is spanned by the basis $\{\frac{\partial x}{\partial \xi_1}, \frac{\partial x}{\partial \xi_2}\}$. By $\mathcal{T}\mathcal{M}$, we denote the tangent bundle. On \mathcal{M} , the metric $g(\cdot,\cdot)$ as a bilinear form on $\mathcal{T}\mathcal{M}\times\mathcal{T}\mathcal{M}$ is prescribed by the metric tensor $g=(g_{ij})_{ij}$ with

$$g_{ij} = \frac{\partial x}{\partial \xi_i} \cdot \frac{\partial x}{\partial \xi_j},$$

where \cdot indicates the scalar product in \mathbb{R}^3 . The inverse of g is denoted by $g^{-1}=(g^{ij})_{ij}$. Based on the metric, we can define the integration of a function f on \mathcal{M} . We split up an integral over \mathcal{M} into separate integrals over subsets which are in the image $x(\Omega)$ of some coordinate map x and define

$$\int_{x(\Omega)} f := \int_{\Omega} f(x(\xi)) \sqrt{\det g} \, d\xi.$$

Integrating either a product of two functions f, g on \mathcal{M} or the product of two vector fields v, w on $\mathcal{T}\mathcal{M}$, we obtain the following scalar products on $C^0(\mathcal{M})$ and $C^0(\mathcal{T}\mathcal{M})$, respectively:

$$(f,g)_{\mathcal{M}} := \int_{\mathcal{M}} fg \, \mathrm{d}x,$$

 $(v,w)_{\mathcal{TM}} := \int_{\mathcal{M}} g(v,w) \, \mathrm{d}x.$

Next, we have to introduce the fundamental intrinsic gradient and divergence operators on \mathcal{M} . The gradient $\nabla_{\mathcal{M}} f$ of f is defined as the representation of df with respect to the metric g. We obtain, in coordinates,

$$\nabla_{\mathcal{M}} f = \sum_{i,j} g^{ij} \frac{\partial (f \circ x)}{\partial \xi_j} \frac{\partial x}{\partial \xi_i}.$$

Furthermore, we define the divergence $\mathrm{div}_{\mathcal{M}}v$ for a vector field $v\in\mathcal{TM}$ as the dual operator of the gradient by

$$\int_{\mathcal{M}} \operatorname{div}_{\mathcal{M}} v \, \phi \, \mathrm{d}x := -\int_{\mathcal{M}} g(v, \nabla_{\mathcal{M}} \phi) \, \, \mathrm{d}x$$



Fig. 13. The principal directions of curvature are visualized by anisotropic diffusion on a minimal surface.

for all $\phi \in C_0^{\infty}(\mathcal{M})$.

Finally, with these differential operators at hand, we can discuss a general and intrinsic diffusion on a manifold in analogy to diffusion in Euclidean space: We ask for a solution $\rho: \mathbb{R}_0^+ \times \mathcal{M} \to \mathbb{R}$ of the parabolic equation

$$\frac{\partial}{\partial t}\rho - \operatorname{div}_{\mathcal{M}}(A\nabla_{\mathcal{M}}\rho) = f(\rho)$$

on $\mathbb{R}_0^+ \times \mathcal{M}$ for given initial data $\rho(0,\cdot) = \rho_0$ on \mathcal{M} . Here, we suppose A to be some positive definite symmetric endomorphism on $T\mathcal{M}$. Testing with any function $\theta \in C^\infty(\mathcal{M}(t))$ and integrating over \mathcal{M} , we obtain the variational formulation

$$(\partial_t \rho, \theta)_{\mathcal{M}} + (A \nabla_{\mathcal{M}} \rho, \nabla_{\mathcal{M}} \theta)_{\mathcal{T} \mathcal{M}} = (f(\rho), \theta)_{\mathcal{M}}.$$

Now, we consider our actual goal, which is the generation of a texture by nonlinear anisotropic diffusion to represent a given vector field $v \in \mathcal{TM}$ on the surface. Thus, we suppose A to depend on the vector field v and the norm of the gradient of a convoluted intensity ρ_{ϵ} :

$$A = A(v, \|\nabla_{\mathcal{M}} \rho_{\epsilon}\|).$$

For no vanishing v, let $w \in \mathcal{T}_x \mathcal{M}$ be some unit vector normal to v, i.e., g(v,w)=0. Hence, $\{\frac{v}{\|v\|},w\}$ is a basis of $\mathcal{T}_x \mathcal{M}$ and, with respect to this basis, we define, as before in the Euclidean case,

$$A(v,d) = \begin{pmatrix} \alpha(\|v\|) & \\ & G\|d\| \end{pmatrix}.$$

As righthand side $f(\cdot)$, we pick up the one already introduced in Section 3 and again assume ρ_0 to be a

random noise, either scalar or vector valued, but now prescribed on the surface \mathcal{M} . Furthermore, we have to give a suitable definition of the regularizing presmoothing to obtain ρ_{ϵ} from the original intensity ρ . Again, we proceed in analogy to the Euclidean case and define ρ_{ϵ} as the result of the above diffusion problem with $A=\operatorname{Id}$ at time $t=\frac{\epsilon^2}{2}$ and for initial data ρ .

Finally, the resulting family $\{\rho(t)\}_{t\geq 0}$ of intensities on \mathcal{M} gives a multiscale of representations of the given vector field v. Fig. 13 and Fig. 14 show results on different surfaces. We consider the principal directions of curvature as tangential vector fields on which we apply the anisotropic diffusion method. On the underlying triangular grids, the shape operator, whose eigenvalues are the principal curvatures, is approximated as follows: Locally, we regard a single triangle T and all the neighboring triangles which have a nonzero intersection with T as a graph over the plane containing T and calculate the L^2 projection of this piecewise linear graph onto the set of quadratic graphs which are tangential to the plane. Then, we evaluate the constant shape operator on this graph. Let us emphasize that the L^2 projection is always defined, although the local graph property of the triangular grid might not hold in certain degenerate cases.

7 DISCRETIZATION IN 2D AND 3D

In what follows, we discuss the discretization and implementation of the field aligned diffusion method. We will first focus on domains in 2D and 3D Euclidean space. For this purpose, a finite element discretization in space and a semi-implicit backward Euler or second order Crank



Fig. 14. For both principal directions of curvature, different timesteps of the anisotropic diffusion are displayed on the surface of a presmoothed Stanford bunny. In addition, the corresponding principle curvature values are color coded.

Nicolsson scheme in time are considered. Here, we have restricted ourselves to regular grids in 2D and 3D generated by recursive subdivision. On these grids, we consider bilinear, respectively, trilinear, finite element spaces. Numerical integration is based on the lumped masses product $(\cdot,\cdot)^h$ [23] for the L^2 product (\cdot,\cdot) in the variational formulation and a one point quadrature rule for the bilinear form $(A\nabla\cdot,\nabla\cdot)$. Semi-implicit means, for the schemes considered here, that the nonlinearity $A(\cdot)$ is evaluated at the old time. Finally, in each step of the discrete evolution, we have to solve a single system of linear equations. We obtain, for a backward Euler discretization,

$$(M^k + \tau L^k(A^k))\bar{\rho}^{k+1} = M^k\bar{\rho}^k + \tau M^k\bar{f}^k.$$

Here, $\bar{\rho}^k = (\bar{\rho}_i^k)_i$ is the vector of nodal intensity values at time $t^k = k\tau$, where τ is the selected time step size. Furthermore, if we denote the "hat shaped" multilinear basis functions by Φ_i and the diffusion tensor with respect to the discrete intensity at time t^k by A^k ,

$$\begin{split} M^k &:= \left((\Phi_i, \Phi_j)^h \right)_{ij} \\ L^k(A^k) &:= \left((A^k \nabla \Phi_i, \nabla \Phi_j) \right)_{ij} \end{split}$$

are the lumped mass matrix and nonlinear stiffness matrix, respectively. Finally, the components of the righthand side \bar{f}^k are evaluated by $(\bar{f}^k)_i = f(\bar{\rho}^k_i)$.

 $ar{f}^k$ are evaluated by $(ar{f}^k)_i = f(ar{
ho}_i^k)$. The global matrices M^k and $L^k(A^k)$ are assembled from local matrices m^E and l^E with respect to a single element. Their entries correspond to all pairings of local basis functions. Due to the applied lumped mass integration, we immediately verify

$$m_{ij}^E = \frac{1}{2^n} \delta_{ij} |E|,$$

where |E| is the volume of the rectangular element E and δ_{ij} the usual Kronecker symbol. For the nonlinear stiffness matrix we obtain

$$\begin{split} l_{ij}^E(A) &= |E| \bigg[\alpha(\|V\|) \bigg(\nabla \Phi^i \cdot \frac{V}{\|V\|} \bigg) \bigg(\nabla \Phi^j \cdot \frac{V}{\|V\|} \bigg) + \\ &G(\|D\|) \bigg(\nabla \Phi^i - \nabla \Phi^i \cdot \frac{V}{\|V\|^2} V \bigg) \\ &\cdot \bigg(\nabla \Phi^j - \nabla \Phi^j \cdot \frac{V}{\|V\|^2} V \bigg) \bigg]. \end{split}$$

where $V = v(c_E)$ for the center of mass c_E of E, D the gradient of the presmoothed discrete intensity at c_E , and $\{\Phi^i\}_i$ the set of local basis functions.

In each time step, the computation of the prefiltered intensity vector $\bar{\rho}^n_{\epsilon}$ is based on a single implicit time step $\epsilon^2/2$ for the corresponding discrete heat equation scheme with respect to initial data $\bar{\rho}^n$.

In our implementation, the regular grids are procedurally interpreted as quadtrees, respectively octtrees [19]. Finally, no matrix is explicitly stored. The necessary matrix multiplications in the applied iterative CG solver are performed in successive tree traversals. Hierarchical BPX type [3] preconditioning is used to accelerate the convergence of the linear solver. The computation of a single time

step on a 257² grid performed on a Silicon Graphics workstation with an R10000 processor requires 1.2 seconds. Computing time in 3D is currently much more expensive. But, there is still a great potential to speed up the algorithm considerably, for instance, by taking into account better ordering strategies for the unknowns which correspond to the anisotropy. This will be exploited in the future. Furthermore, the code is prepared to incorporate spatial grid adaptivity if possible (cf. Fig. 17).

8 DISCRETIZATION ON SURFACES

The discretization of the proposed anisotropic diffusion method on surfaces is completely analogous to the above Euclidean case. We only have to replace the discrete differential operators and bilinear forms by their intrinsic geometric counterparts. We suppose the surface $\mathcal M$ to be approximated with a sufficiently fine triangular grid $\mathcal M_h$ consisting of nondegenerate triangles T with maximal diameter h. Thus, we only focus on the computation of the local mass matrix m^T and the local nonlinear stiffness matrix $l^T(A)$, respectively. We obtain again by lumped mass integration

$$m_{ij}^T = \frac{1}{3}\delta_{ij}|T|,$$

where |T| is the area of the triangle T. Next, let us consider, for every triangle T, the reference triangle $\hat{T} \subset \mathbb{R}^2$ with independent variables ξ_1, ξ_2 and nodes $\xi^0 = (0,0)$, $\xi^1 = (1,0)$, and $\xi^2 = (0,1)$. Then, an affine coordinate mapping X maps \hat{T} onto T and its nodes ξ^i onto the corresponding nodes P^i of T on the discrete surface in \mathbb{R}^3 . Hence, the corresponding metric tensor is as in the continuous case given by $g_{ij} = \frac{\partial X}{\partial \xi_i} \cdot \frac{\partial X}{\partial \xi_j}$, where $\frac{\partial X_k}{\partial \xi_i} = P_k^i - P_k^0$. Hence, we can evaluate gradients of the linear basis functions Φ^l corresponding to the nodes P^l by

$$abla_{\mathcal{M}_h}\Phi^l = \sum_{i,j} g^{ij} rac{\partial \Phi^l}{\partial \xi_j} (P^i - P^0),$$

where the derivatives of Φ^l with respect to the reference coordinates ξ are

$$\begin{pmatrix} \frac{\partial \Phi^l}{\partial \xi_1} \\ \frac{\partial \Phi^l}{\partial \xi_2} \end{pmatrix} = \begin{pmatrix} -1 \\ -1 \end{pmatrix}, \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Finally, we calculate the local nonlinear stiffness matrix

$$\begin{split} l_{ij}^T(A) &= \\ |T| \bigg[\alpha(\|V\|) \bigg(\nabla_{\mathcal{M}_h} \Phi^i \cdot \frac{V}{\|V\|} \bigg) \bigg(\nabla_{\mathcal{M}_h} \Phi^j \cdot \frac{V}{\|V\|} \bigg) + \\ G(\|D\|) \bigg(\nabla_{\mathcal{M}_h} \Phi^i - \nabla_{\mathcal{M}_h} \Phi^i \cdot \frac{V}{\|V\|^2} V \bigg) \\ \cdot \bigg(\nabla_{\mathcal{M}_h} \Phi^j - \nabla_{\mathcal{M}_h} \Phi^j \cdot \frac{V}{\|V\|^2} V \bigg) \bigg]. \end{split}$$

where $V = v(c_T)$ for the center of mass c_T of T, D the geometric gradient of the presmoothed discrete intensity on T, and "·" still indicates the scalar product in \mathbb{R}^3 .

9 COMPARISON TO OTHER METHODS

So far, we have introduced a novel approach which provides us with an intuitive understanding of complex vector fields. We have discussed a variety of important properties and advantages. Let us now rank this method among other visualization methods and compare it with different techniques. Here, we especially pick up the line integral convolution method and the spot noise approach.

For stationary vector fields, we obtain similar results by all methods. Thin field aligned patterns are generated. Line integral convolution leads to comparable results with the essential difference that the PDE-based method carries a nice scale space property, i.e., evolving a longer time in the anisotropic diffusion method, we obtain a successive coarsening of the resulting pattern representing the vector field.

Furthermore, in a restricted sense, line integral convolution (LIC) and spot noise can be regarded as special cases of the anisotropic nonlinear diffusion method. LIC with Gaussian filter kernel can be identified as the asymptotic limit of the latter method for a concentration of the edge enhancing function $G(\cdot)$ at 0. Other filter kernel shapes correspond to different, in general, nonlinear diffusion processes along streamlines. Further on, generating a single deformed spot on the computational domain, as proposed in [6], can be regarded as an early time step in the diffusion starting with initial data, that is, a characteristic function of a circular disk. If we release a bunch of such disks as initial data in such a way that the evolving patterns do not overlap, then the resulting image is comparable to spot noise. Thus, the original spot noise technique can be regarded as a parallel version of short time diffusive vector field visualization.

10 TOWARD FLOW SEGMENTATION

The above applications already show the capacity of the anisotropic nonlinear diffusion method to outline the flow structure not only locally. Indeed, especially for larger evolution times in the diffusion process, the topological skeleton of a vector field becomes clearly visible. We will now investigate a possible flow segmentation by means of the anisotropic diffusion. Let us restrict this to the twodimensional case of an incompressible flow with vanishing velocity v at the domain boundary $\partial\Omega$. Then, topological regions are separated by homoclinic, respectively, heteroclinic, orbits connecting critical points in the interior of the domain and stagnation points on the boundary. Critical points, by definition, points with vanishing velocity v = 0, may either be saddle points or vortices. Furthermore, we assume critical points to be nondegenerate, i.e., ∇v is regular. Saddle points are characterized by two real eigenvalues of ∇v with opposite sign, whereas, at vortices, we obtain complex conjugate eigenvalues with vanishing real part. Stagnation points on $\partial\Omega$ are similar to saddles. For details we refer to [11]. In each topological region, there is a family of periodic orbits close to the heteroclinic, respectively, homoclinic, orbit. This observation gives reason for the following segmentation algorithm. At first, we search for critical points in Ω and stagnation points on $\partial\Omega$. We

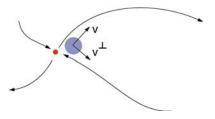


Fig. 15. A sketch of the four sectors at a critical point, the initial spot for the diffusion calculation and the oriented system $\{v,v^{\perp}\}$.

calculate the directions which separate the different topological regions. In the case of saddle points, these are the eigenvectors of ∇v . Next, we successively place an initial spot in each of the sectors and perform an appropriate field aligned anisotropic diffusion. Let us suppose that a single sector is spanned by vectors $\{s_+, s_-\}$, where the sign \pm indicates incoming and outgoing direction. The method presented in Section 3 would lead to a closed pattern along one of the above closed orbits for time t large enough. To fill out the interior region, we modify the diffusion as follows: Up to now, the Perona Malik diffusions enhance edges of the current image in both directions normal to the velocity. Henceforth, we select an orientation for a "one sided" diffusion (cf. Fig. 15), i.e., we

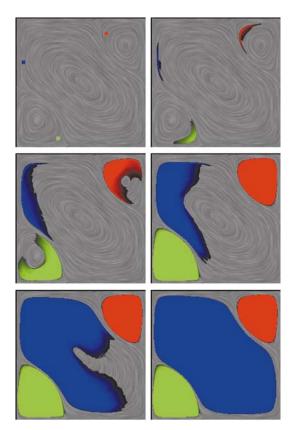


Fig. 16. Nonlinear diffusion segmentation is applied to a velocity field from a Bénard convection. Several time steps are shown starting from initial seed spots in critical point sectors. Here, we have placed these seeds as close as possible in terms of the grid size in the sectors spanned by the eigenvalues of the Jacobian of the velocity. Only to emphasize the evolution process, a single grayscale image from the diffusion calculation (cf. Fig. 8) is underlying the sequence of segmentation time steps.

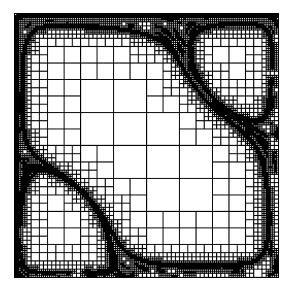


Fig. 17. The adaptive quadtree on which we approximate the segmentation function ρ at a certain time step.

select a unique normal v^{\perp} to v and consider the diffusion matrix

$$A(v, \nabla \rho_{\epsilon}) = B(v)^T \begin{pmatrix} \alpha \\ G((\nabla \rho_{\epsilon} \cdot v^{\perp})_{+}) \end{pmatrix} B(v),$$

where α is a positive constant and $(s)_+ := \max\{s, 0\}$. Furthermore, we consider a nonnegative, concave function $f: \mathbb{R}_0^+ \to \mathbb{R}_0^+$ with f(0), f(1) = 0 as a source term in the diffusion equation. If the orientation of $\{s_+, s_-\}$, coincides with that of $\{v, v^{\perp}\}$, then linear diffusion in the direction toward the interior will fill up the complete topological region. A segmentation of multiple topological regions at the same time is possible if we carefully select the sectors to release initial spots. Fig. 16 shows different time steps of the segmentation applied to a convective incompressible flow. So far, we have seen that anisotropic diffusion has strong provisions for flow segmentation as well. In a certain sense, we thereby identify the complement of what is usually extracted in topology recognition. An outstanding advantage of the new method is its numerical stability and its selfsharpening effect due to the edge enhancing strategy. We pay for this by a higher computational complexity. If we apply a standard implementation on a uniform grid of size n^2 , the segmentation cost is at least $O(n^2)$ compared to an O(n) count of grid cells met by the direct ODE integration to compute the homoclinic and heteroclinic orbits corresponding to the critical points. Fig. 17 shows an adaptive quadtree which allows the same resolution quality for the segmentation function ρ as on a full grid, but now at a much lower cost. We thereby consider a piecewise linear and continuous finite element space on the adaptive quadtree.

11 Conclusions

We have introduced a new method based on the solution of a nonlinear anisotropic diffusion problem for the post processing of vector data. The method can be applied on 2D or 3D domains, as well as on two-dimensional surfaces embedded in \mathbb{R}^3 . From a mathematical point of view, one of

the major advantages is that it is based on a physically intuitive continuous model, i.e., streamline aligned diffusion. Most of the properties can be discussed on this level. Finally, it is discretized in an appropriate way making use of recent and efficient numerical algorithms.

From the authors' point of view, exciting future research directions are further investigations of flow visualization in 3D. The exploiting of adaptive finite element paradigms and ordering strategies for the unknowns will be especially key issues to reducing the computing costs.

Furthermore, a visualization approach based on anisotropic diffusion and applicable for time dependent vector fields is a challenging topic. Finally, the anisotropic diffusion flow segmentation also carries provisions for the identification of interesting flow regions in 3D, such as recirculation zones and vortex cores.

Further results and the algorithm running on an $n \times m$ 2D vector array is available as a source code at http:// www.iam.uni-bonn.de/FktAna_NumMath/Num_Vis/ projekte/flow_visualization/.

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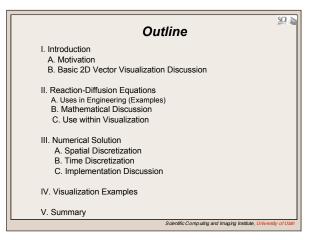


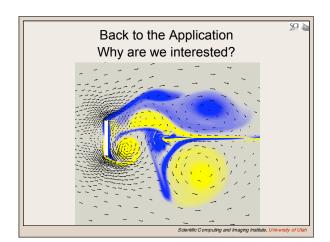
Tobias Preußer studied mathematics at the University of Bonn, Germany, and at the Courant Institute at New York. He received his Diploma degree in 1999. Currently, he is working on his PhD thesis in the field of anisotropic diffusion methods in multiscale feature extraction. While a student, he worked for the National Research Center for Information Technology (GMD). In the spring of 2000, he received a grant from the European Community for his work at the Italian CINECA supercomputing center.



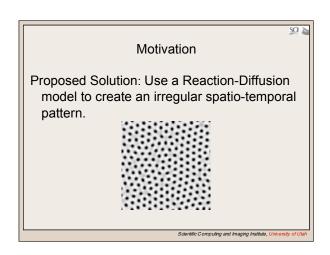
Martin Rumpf received his Diplom degree and his PhD in mathematics from Bonn University in 1989 and 1992, respectively. He has been a professor of applied mathematics at Bonn University since 1996. His research interests are numerical methods for nonlinear partial differential equations, adaptive finite element methods, and computer vision. Furthermore, he is concerned with new flow visualization techniques and efficient multiscale methods in scientific visualization.

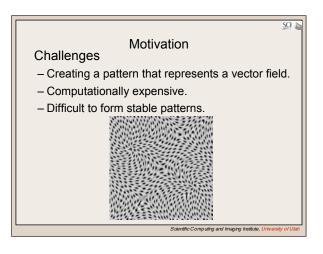
Reaction-Diffusion Models for Vector Visualization: Algorithms and Implementations Speaker: Mike Kirby School of Computing and Magin Institute University of Utan Salt Lake City, UT, USA Work Done In Collaboration With: Allen Sanderson Scientific Computing and Imaging Institute University of Utan Salt Lake City, UT, USA





Motivation Visualization Goals Create a visualization of a vector field that combines the advantages of dense and sparse representations. Dense representations such as Spot Noise and LIC show global orientation but without augmentation lack local magnitude and direction. Sparse representations such as streamlines and glyphs give local magnitude and direction but due to occlusion can miss features. Represent other scalar values such as uncertainty within the visualization without overloading.





Examples of Reaction - Diffusion Models In Engineering

Alan Turing (1952) - The chemical basis of morphogenesis.

$$\frac{\partial a}{\partial t} = s(\alpha - ab) + d_a \nabla^2 a$$

$$\frac{\partial b}{\partial t} = s(ab - b - \beta) + d_b \nabla^2 b$$

Describes the chemical process between two morphigens, a and b.

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Examples of Reaction - Diffusion Models In Engineering

Gray-Scott Equations: Mathematical Biology

$$\frac{\partial u}{\partial t} = -uv^2 + F(1 - u) + d_u \nabla^2 u$$

$$\frac{\partial v}{\partial t} = uv^2 - (F + k)v + d_v \nabla^2 v$$

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Examples of Reaction - Diffusion Models In Engineering

FitzHugh-Nagumo Equations: Electrocardiology

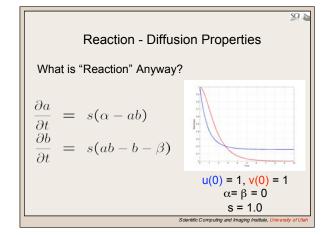
$$C_m \frac{\partial v}{\partial t} = Av(v - \alpha)(1 - v) + d_v \nabla^2 v$$
$$\frac{\partial w}{\partial t} = \epsilon(v - \gamma w) + d_w \nabla^2 w$$

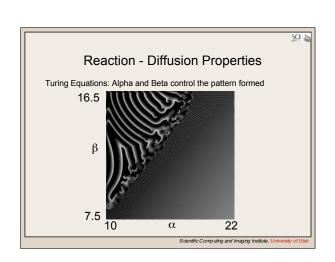
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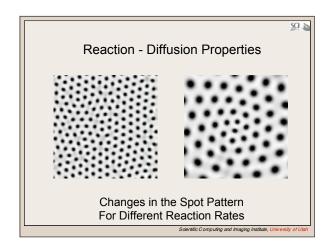
Reaction - Diffusion Previous Research

- Reaction Diffusion for texture generation
 - Turk et al. 1991
- Witkin et al. 1991
- Fowler et al. 1992
- Chambers et al. 1995
- Reaction Diffusion for tensor visualization
 - Kindleman et al. 2000
- Diffusion for vector visualization
 - Preußer et al. 1999
 - Garcke et al. 2001
- Advection for vector visualization
 - Shen et al. 1996
 - Weiskopf et al. 2004

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Reaction - Diffusion Properties

How do we use diffusion?

Embedding the vector field in the diffusion tensor:

$$\frac{\partial u}{\partial t} = H(u, v) + (\nabla \cdot \sigma_u \nabla) u$$

where σ_u is a symmetric positive definite diffusion tensor.

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Reaction - Diffusion Properties

Assume that we are working on a finite domain $[a_1,b_1] \times [a_2,b_2] \in \mathcal{R}^2$, and that we are given a regularly-spaced computational grid of size $N_x \times N_y$. At each point (x_i,y_j) , $i=1,\ldots,N_x,j=1,\ldots,N_y$ suppose that we are given an angle $\theta_{ij} \in [0,2\pi]$.

$$\begin{split} \mathbf{R}_{ij} &= \begin{pmatrix} \cos\theta_{ij} & \sin\theta_{ij} \\ -\sin\theta_{ij} & \cos\theta_{ij} \end{pmatrix} \\ \mathbf{R}_{ij}^{-1} &= \mathbf{R}_{ij}^T = \begin{pmatrix} \cos\theta_{ij} & -\sin\theta_{ij} \\ \sin\theta_{ij} & \cos\theta_{ij} \end{pmatrix} \end{split}$$

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Reaction - Diffusion Properties

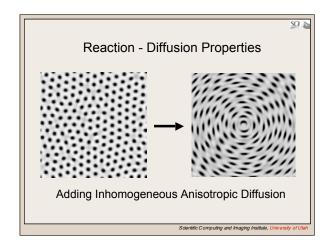
$$\Lambda_{ij} = \begin{pmatrix} (\lambda_1)_{ij} & 0 \\ 0 & (\lambda_2)_{ij} \end{pmatrix}$$

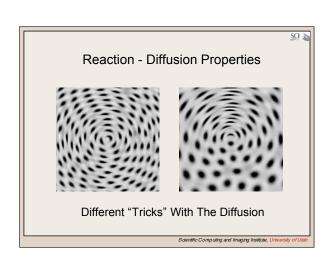
where $(\lambda_1)_{ij}$ is the diffusivity in the first principal direction and $(\lambda_2)_{ij}$ is the diffusivity in the second principal direction.

$$\sigma_{ij} = \mathbf{R}_{ij}^T \wedge_{ij} \mathbf{R}_{ij}.$$

$$\sigma_{ij} = \begin{pmatrix} (\sigma_{11})_{ij} & (\sigma_{12})_{ij} \\ (\sigma_{21})_{ij} & (\sigma_{22})_{ij} \end{pmatrix}.$$

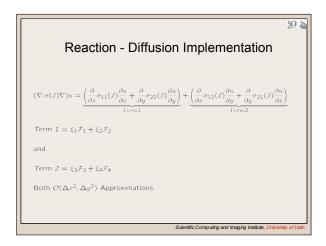
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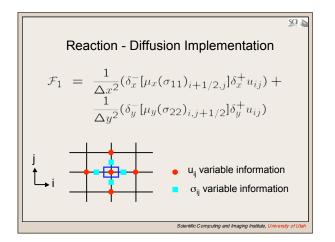


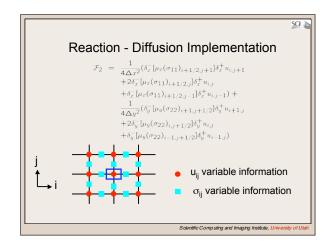


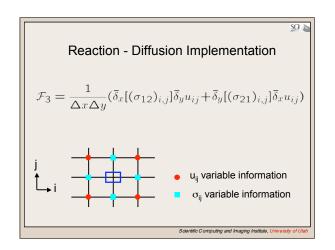
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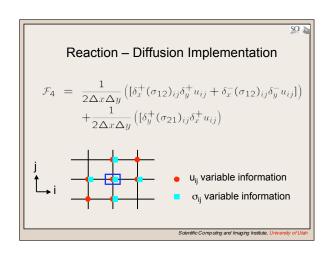
Reaction - Diffusion Implementation $(\nabla \cdot \sigma(\vec{x}) \nabla) u = \frac{\partial}{\partial x} \sigma_{11}(\vec{x}) \frac{\partial u}{\partial x} + \frac{\partial}{\partial x} \sigma_{12}(\vec{x}) \frac{\partial u}{\partial y} + \frac{\partial}{\partial y} \sigma_{21}(\vec{x}) \frac{\partial u}{\partial x} + \frac{\partial}{\partial y} \sigma_{22}(\vec{x}) \frac{\partial u}{\partial y}$ $\delta_x^+ u_{ij} = u_{i+1,j} - u_{i,j}$ $\delta_y^- u_{ij} = u_{i,j} - u_{i-1,j}$ $\delta_y^- u_{ij} = u_{i,j} - u_{i,j-1}$ Basic Finite Difference $\delta_y u_{ij} = u_{i,j+1/2,j} - u_{i-1/2,j}$ $\delta_y u_{ij} = u_{i,j+1/2} - u_{i,j-1/2}$ Operators $\delta_z u_{ij} = \frac{1}{2} (u_{i+1/2,j} - u_{i-1/2,j})$ $\delta_y u_{ij} = \frac{1}{2} (u_{i+1/2,j} + u_{i-1/2,j})$ $\mu_x u_{ij} = \frac{1}{2} (u_{i+1/2,j} + u_{i-1/2,j})$ Soleratic Computing and Imaging Institute, University of Utah.

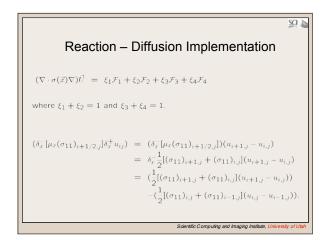


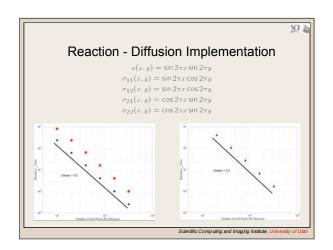


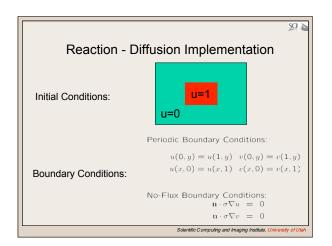


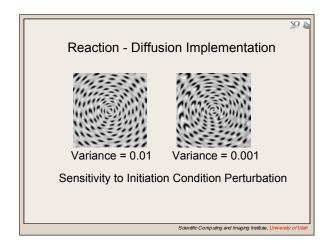


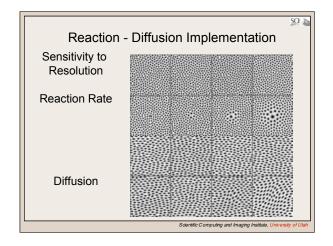


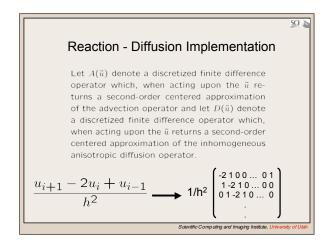


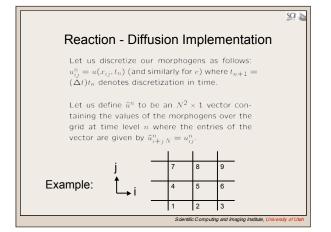


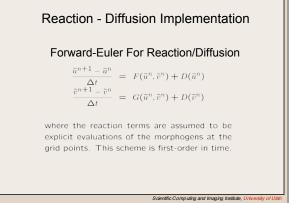








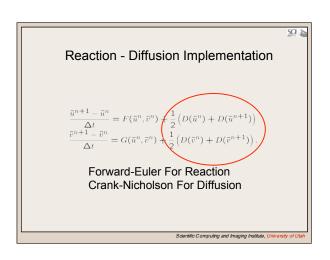




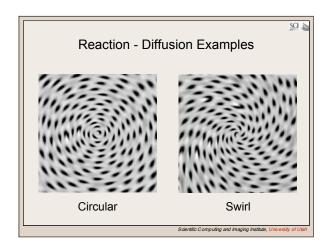
Reaction - Diffusion Implementation $\frac{\tilde{u}^{n+1} - \tilde{u}^n}{\frac{\Delta t}{\tilde{v}^{n+1} - \tilde{v}^n}} = F(\tilde{u}^n, \tilde{v}^n) + D(\tilde{u}^{n+1}) + D(\tilde{v}^{n+1}).$ Forward-Euler for Reaction Backward-Euler For Diffusion

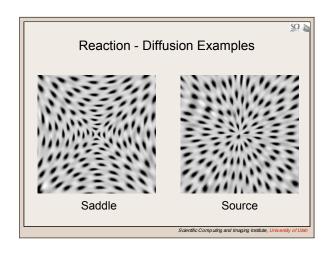
Reaction - Diffusion Implementation Given the linearity of the discretized diffusion operator, we can thus re-write the system to be solved as: $(\mathbf{I} - \Delta t \, \mathbf{D}) \, \tilde{u}^{n+1} \, = \, \tilde{u}^n + (\Delta t) F(\tilde{u}^n, \tilde{v}^n) \\ (\mathbf{I} - \Delta t \, \mathbf{D}) \, \tilde{v}^{n+1} \, = \, \tilde{v}^n + (\Delta t) G(\tilde{u}^n, \tilde{v}^n)$ where \mathbf{I} denotes the $N^2 \times N^2$ identity operator and \mathbf{D} denotes the $N^2 \times N^2$ matrix form of the linear finite difference diffusion operator.

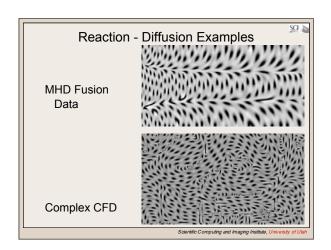
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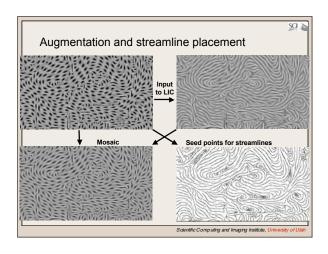


Reaction - Diffusion Implementation Given the linearity of the discretized diffusion operator, we can thus re-write the system to be solved as: $\left(\mathbf{I} - \frac{1}{2}\Delta t\,\mathbf{D}\right)\tilde{u}^{n+1} \ = \ \tilde{u}^n + \Delta t\left(F(\tilde{u}^n,\tilde{v}^n) + \frac{1}{2}\mathbf{D}(\tilde{u}^n)\right) \\ \left(\mathbf{I} - \frac{1}{2}\Delta t\,\mathbf{D}\right)\tilde{v}^{n+1} \ = \ \tilde{v}^n + \Delta t\left(G(\tilde{u}^n,\tilde{v}^n) + \frac{1}{2}\mathbf{D}(\tilde{v}^n)\right)$ where \mathbf{I} denotes the $N^2 \times N^2$ identity operator and \mathbf{D} denotes the $N^2 \times N^2$ matrix form of the linear finite difference diffusion operator.









Summary Reaction-Diffusion provides one means of accomplishing vector visualization Pros: a. Reasonably easy to implement with Finite Differences b. Both CPU and GPU friendly Cons: a. Fickle to things like reaction rates, initial conditions b. Does not provide "controlled" pattern spacing c. Computationally intensive (due to running to steady state) Reaction-Diffusion can be used to augment other procedures

Display of Vector Fields Using a Reaction-Diffusion Model

Allen R. Sanderson, Chris R. Johnson, and Robert M. Kirby*

Scientific Computing and Imaging Institute, University of Utah

ABSTRACT

Effective visualization of vector fields relies on the ability to control the size and density of the underlying mapping to visual cues used to represent the field. In this paper we introduce the use of a reaction-diffusion model, already well known for its ability to form irregular spatio-temporal patters, to control the size, density, and placement of the vector field representation. We demonstrate that it is possible to encode vector field information (orientation and magnitude) into the parameters governing a reaction-diffusion model to form a spot pattern with the correct orientation, size, and density, creating an effective visualization. To encode direction we texture the spots using a light to dark fading texture. We also show that it is possible to use the reaction-diffusion model to visualize an additional scalar value, such as the uncertainty in the orientation of the vector field.

An additional benefit of the reaction-diffusion visualization technique arises from its automatic density distribution. This benefit suggests using the technique to augment other vector visualization techniques. We demonstrate this utility by augmenting a LIC visualization with a reaction-diffusion visualization. Finally, the reaction-diffusion visualization method provides a technique that can be used for streamline and glyph placement.

Keywords: Vector Field Visualization, Flow Visualization, Reaction-Diffusion, Vector Fields

1 Introduction

Visualizing vector fields is important for many computational applications, including fluid dynamics, wind and water currents in climate modeling, bioelectric fields in neuroscience, and magnetic fields in nuclear fusion. To meet the needs arising from this diverse set of applications, many different techniques for visualizing vector fields have been developed [1, 3, 7, 8, 19, 22, 31, 33, 30]. Each technique has strengths and weaknesses in its ability to represent the magnitude, orientation, direction, uncertainty and topological structures of the associated vector field.

For instance, the simplest method for displaying a vector field is to place glyphs representing the vector direction and magnitude at regular intervals. However, because of scaling differences, overlap between the glyphs can occur. This can produce visual clutter and occlusion [25]. The problem is further compounded when data are displayed in three dimensions. Displaying normalized vector values can reduce the clutter but at a loss of information. Even when the visual clutter can be overcome, displaying vector fields using regular intervals may not be appropriate. This is because the grid spacing may not correspond to the underlying vector field.

More complicated techniques, such as streamlines, can provide powerful visual cues [10]. However, enough streamlines must be placed to provide cues without causing visual clutter. Streamlines

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IEEE Visualization 2004 October 10-15, Austin, Texas, USA 0-7803-8788-0/04/\$20.00 ©2004 IEEE may be placed selectively to reduce the clutter, but such selective placement may cause critical areas to be missed [28].

With the exception of a glyph-based method, no technique is singularly able to visualize uncertainty in vector fields. In [17] Pang et al. demonstrated several different glyphs for characterizing uncertainty in vector fields. However, as a glyph based method it can also succumbs to clutter and occlusion.

Given the various shortcomings in many of the current vector field visualization techniques, the main goal of this work was to develop an automated method that uses the vector magnitude, orientation, direction, and uncertainty to control the shape, size, orientation, direction, and density of the objects used to represent the vector field. At the same time, we wanted a method that would be mesh independent and produce a visualization that would be natural and pleasing to the eye. To achieve these goals, we have explored the use of a reaction-diffusion model for vector field visualization.

2 BACKGROUND AND PREVIOUS WORK

Because of its importance to scientific computing applications, creating effective techniques for visualizing vector fields is an active area of research. Given the large number of techniques that have been developed, it is not practical to review each technique; a very complete review can be found in [18]. Instead, we focus on three related areas for visualizing vector fields: the use of random patterns, selective placement, and reaction-diffusion.

The use of random patterns for visualizing a vector field has been explored by van Wijk [29], Cabral and Leedom [3], Shen et al. [23], and others using either spot or white noise to form a dense representation of the vector field by integrating an ODE that represents the basis for a streamline. By dense, we mean that there is value for each grid location. Others, such as Preußer, [20] have formulated the problem as a PDE describing nonlinear anisotropic diffusion. With the ODE and PDE formulations the resulting image has a brush-stroke appearance of variable thickness. While this type of image is useful for showing the vector orientation, it fails to provide information about magnitude and direction. These shortcomings have been addressed in various forms by adding directional [23, 30] and magnitude [5, 13] cues.

More recent work has focused on creating images that are less dense, but still contain useful information about the vector field, e.g. [12, 24, 28]. Turk and Banks [28] explored a method to bundle similar streamlines until an energy function is minimized. Once the function is minimized, the streamlines can be replaced with variable sized curved arrows to show direction and magnitude. Kirby et al. [12], were able to achieve similar results using a random placement of variable sized arrows. Once an arrow is placed, a Poisson distribution disk based on the vector magnitude is used to prevent other arrows from being placed near it. However, because the arrow represents just the value at a single location rather than representing the values in a neighborhood of the sample, it is possible to create the illusion that an area is homogeneous when it is not.

Computer graphics applications have also made use of a reaction-diffusion model to generate texture maps [27, 32]. These types of textures are useful for forming patterns that are natural looking and are typically applied to organic models such as ani-

mals. Turk explored the use of different reaction models to produce a variety of patterns [27]. At the same time, Witkin and Kass [32] used anisotropic diffusion to form different patterns. These patterns can be classified as either spot or stripe patterns.

Rather than forming the texture and then applying it to a model using a traditional texture mapping, Turk exploited the fact that a reaction-diffusion model can be used on an irregular grid. This allowed him to create textures directly on a tessellated surface, avoiding any warping between model space and parameter space. It is possible to make use of this same property to texture isosurfaces, which is a very common visualization tool. In a similar vein, Chambers and Rockwood [4] employed a reaction-diffusion model to generate a solid texture, which is used on a surface and on a volume. Like Witkin and Kass, Chambers and Rockwood also used an anisotropic diffusion technique to form stripe patterns.

Although the reaction-diffusion model was initially mentioned by Cabral and Leedom as a possible model for visualizing data [3], the first researchers to use the method were Kindlmann et al.[11]. They created a solid texture using a reaction-diffusion model based on tensor values from diffusion-weighted magnetic resonance images. This anisotropy formed elliptical "blobs," which were then volume rendered. Our approach is similar in that we also use a tensor but rather than having tensor data supplied we create the tensor from vector data. This gives several additional degrees of freedom that can be used to encode additional data.

Perhaps the most closely related work is that of Garcke et al. [6] were they use the vector field to define the amount anisotropy in the diffusion. Further, they are able to incorporate a shaped particle to show direction as well as orientation. They use both methods to create visualizations of clustered vector fields.

We now present a more detailed look at reaction-diffusion, and describe how it can be used to visualize vector field data.

3 REACTION DIFFUSION

In 1952 Turing [26] proposed a reaction-diffusion model for describing the chemical process between two morphogens within a series of cells. Due to the dynamics of the system, the morphogens both react and diffuse which changes their concentration within each cell. With time, the morphogens may form a stable pattern representing the dynamic equilibrium of the system. The pattern formation is not dependent on the initial state of the system; the dynamics of the system drives the concentrations toward an equilibrium state.

Turing described the reaction-diffusion of a two morphogen model as a set of nonlinear partial differential equations:

$$\frac{\pi a}{\pi t} = F(a,b) + d_a \nabla^2 a \tag{1}$$

$$\frac{\mathfrak{I}b}{\mathfrak{I}t} = G(a,b) + d_b \nabla^2 b \tag{2}$$

where a and b are the morphogen concentrations; F and G are the functions controlling the production rate of a and b, respectively; d_a and d_b are the diffusion rates, and $\nabla^2 a$ and $\nabla^2 b$ are the Laplacians of a and b, respectively. Turing further defined F and G as:

$$F(a,b) = s(a - ab) \tag{3}$$

$$G(a,b) = s(ab - b - b) \tag{4}$$

where a and b again are the morphogen concentrations, a and b are the formation and degradation rates of a and b respectively, and s is the reaction rate.

For the state to change, there must be some perturbation in the initial conditions which is a stable solution. This perturbation can

arise from a non-uniformity in either the initial concentrations, a and b, or the formation and degradation rates, a and b. A nonuniform formation and degradation rate can be interpreted as being the natural variation within each cell.

After the system is put into motion, the morphogen concentrations will change until a dynamic equilibrium is reached and thus a stable pattern is formed. Although the pattern is stable, the morphogen concentration in each cell will continue to change. However, the change is statistically very small.

Turing's equations are just one specific instance of reactiondiffusion phenomena. Other similar variations can be found in the literature, including those in [16]. In this paper, we have focused on the use of Turing's model because it provides results that are indicative of reaction-diffusion in general. We note that the techniques provided in this paper can be similarly applied to other reactiondiffusion systems.

3.1 Mapping the Reaction-Diffusion Kinetics

In order to use a reaction-diffusion model for visualization, a mapping must be established between the vector field and the reaction-diffusion model. There are three possibilities: a mapping between the vector field and the reaction kinetics, a mapping between the vector field and the diffusion kinetics, or a combination of both. Because of the instabilities associated with a reaction-diffusion model, we have focused on finding a mapping for the vector magnitude to the reaction kinetics and the vector orientation to the diffusion kinetics.

3.2 Reaction Kinetics

Within Turing's reaction kinetics there are several free variables. We first mapped the patterns formed as a function of the two reaction values, a and b. The patterns formed can be described as being finger print or spot patterns. However, these formed only inside a very narrow band of values for a and b. On either side of this band a stable pattern did not occur. For this work we chose a and b to be $16\pm1\%$ and $12\pm1\%$ respectively, which produced spot like patterns.

The other free parameter in the reaction kinetics is the reaction rate *s*. Changing the reaction rate changes the size of the pattern formed. This provides an ideal mapping to a scalar value such as the vector magnitude. Thus, it is possible to create patterns of varying size with the size directly relating to the vector magnitude.

It should be noted that similar results can be obtained by varying the diffusion rates d_a and d_b . As such, it is neither the reaction nor the diffusion rates that changes the size, but rather their relative difference. For simplicity and clarity, we vary only the reaction rate for each cell.

3.3 Diffusion Kinetics

The diffusion kinetics as written in Equations (1) and (2) have one free parameter per equation, the diffusion rates d_a and d_b . As previously noted, changing the diffusion rate changes the size of the pattern formed, and for our purposes it is fixed. However, this is not the only free parameter in the diffusion equation. If we relax the isotropic diffusion condition and use anisotropic diffusion we are able to create a broader range of patterns. Assuming anisotropy, the reaction-diffusion equation can be generalized as:

$$\frac{\mathfrak{I}u}{\mathfrak{I}t} = H(u,v) + (\nabla \cdot s_u \nabla)u, \tag{5}$$

where s_u is a symmetric positive definite diffusion tensor into which we will encode the vector field of interest.

Assume that we are working on a finite domain $[a_1,b_1] \times [a_2,b_2] \in \mathcal{R}^2$, and that we are given a regularly-spaced computational grid of size $N_x \times N_y$. At each point (x_i,y_j) , $i=1,\ldots,N_x$, $j=1,\ldots,N_y$ suppose that we are given a vector $\vec{v}_{ij}=(u_{ij},v_{ij})^T$. To embed this vector field into a tensor field, define $q_{ij}=\arctan\frac{v_{ij}}{u_{ij}} \in [0,2_{\mathbb{Z}}]$.

We can now define a rotation matrix and its inverse based upon the angle determined above:

$$\mathbf{R}_{ij} = \begin{pmatrix} \cos q_{ij} & \sin q_{ij} \\ -\sin q_{ij} & \cos q_{ij} \end{pmatrix}$$
 (6)

$$\mathbf{R}_{ij}^{-1} = \mathbf{R}_{ij}^{T} = \begin{pmatrix} \cos q_{ij} & -\sin q_{ij} \\ \sin q_{ij} & \cos q_{ij} \end{pmatrix}$$
 (7)

We now define a principal diffusivity matrix Λ which is a diagonal matrix and gives the diffusivity coefficients along the two principal axes of diffusion:

$$\Lambda_{ij} = \begin{pmatrix} (1_1)_{ij} & 0\\ 0 & (1_2)_{ij} \end{pmatrix}$$
 (8)

where $(1_1)_{ij}$ is the diffusivity in the first principal direction and $(1_2)_{ij}$ is the diffusivity in the second principal direction.

With the definitions above we can define a diffusivity tensor s_{ij} based on our vector field as

$$s_{ii} = \mathbf{R}_{ii}^T \Lambda_{ii} \mathbf{R}_{ii}. \tag{9}$$

We then combine the spatially nonuniform anisotropic diffusion matrix with a discrete finite difference Laplacian as outlined in [9] in a manner that maintains second-order convergence. With this control, we now have our desired mapping between the vector orientation and the diffusion kinetics. Witkin and Kass [32] took a similar approach for creating 2D texture patterns but assumed a spatially uniform anisotropic diffusion matrix.

3.4 Directional Texturing

The final step in the mapping process is to encode directional information in the pattern created. Up to this point we have described the use of a reaction-diffusion model for generating non-specific patterns. The most common pattern formed using a reaction-diffusion model is a spot pattern. The exact formation of the spot pattern will be discussed in the following section. Assuming an oriented elliptical spot pattern, we show the direction by texturing the spots with a contrasting light to dark fading texture. The texturing is local for each spot using the following steps:

Step 1: Assuming dark spots on a light background, normalize the image values to be between 0 and 255.

Step 2: Find the centroid of each spot by first thresholding to remove all pixels with a value greater than 64. Second, thin the remaining pixels to into single pixels [21]. Label the remaining single pixels as the centroid.

Step 3: For each centroid find all connected pixels with a value less than 128.

Step 4: For each connected pixel, calculate the dot product between the vector form by the centroid and the connected pixel and the underlying vector field value at the centroid.

Step 5: Normalize the dot product to be between 0 and 1 based on the minimum and maximum dot product values within the spot.

Step 6. If the normalized dot product is less than .9, interpolate between the minimum and maximum gray scales in the spot. This becomes the new gray scale value that gives the dark to light fading on the spot. Otherwise the gray scale is set to be 255, which produces a contrasting light tip on the spot.

The texturing does not change the size or orientation of the spots; merely highlights the direction using values that maintain their natural appearance in the image. A light to dark fading provides directional cues because the fading has a natural strong to weak association. Wegenkittl et al. [30], took a similar approach to create oriented streamlines. The texturing is fully demonstrated in the following section.

4 IMAGE FORMATION

A reaction-diffusion image is created using a forward Euler integration on the discrete version of Equation (5) for a and b until a dynamic equilibrium state is reached, at which time, a stable pattern will have formed. We have found that using a cell size of 1.0 and a step size of 0.5 provides a balance between numerical stability and the pattern formation.

Figure (1a) shows a spot pattern created with the Turing model. Analysis of the image shows that the spot placement is balanced. That is, there is a uniform density of spots with equal spacing around them. This balancing process can be observed during the integration process when a spot begins to form in an area of lower concentration. Other nearby spots adjust themselves so they are not too close to the newly formed spot. Sometimes this adjustment may come in the form of a change in the position of the spots or when one or more of the spots disappears and its concentration is absorbed by remaining spots. This natural organization is one of the properties of reaction-diffusion equations that makes them very useful for visualization purposes.

Circular spots alone do not show magnitude, orientation, or direction. As discussed previously, to show magnitude we scale the spots using the reaction rate, *s* to reflect the vector magnitude, Figure (1b). To show orientation we compress the spots into elliptical shapes by applying an anisotropic diffusion matrix where the values along the principal axis have a 3:1 ratio. Next, we rotate the diffusion matrix for each vector so that the ellipse's major axis is aligned with the vector field. Once the system comes to a dynamic equilibrium and the spots have formed, the light to dark texture is applied to each spot to show direction, Figure (2a). Figure (2b) shows the magnitude, orientation, and direction combined.

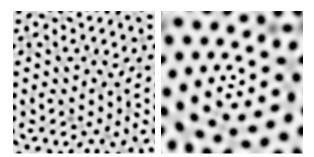


Figure 1: (a) Reaction-Diffusion visualization of circular spots of constant size and (b) variable size.

We now show the application of our reaction-diffusion model to visualize a set of idealized vector fields. Our goal was to see if it was possible to capture the nature of different types of vector fields. These fields include the electric field from a dipole and an electrostatic charge, Figure (3); and a vector field for a saddle and a sink, Figure (4). In each vector field the magnitude and orientation changes smoothly. The properties of vector field can easily be discerned as the ellipsoidal spots have the correct size and orientation. In the case of the circular, saddle, and sink vector fields, a change in the vector magnitude occurs as the vector field moves away from the center and is shown by a corresponding change in the spot size. However, in Figures (2-4), the images have several spots that did

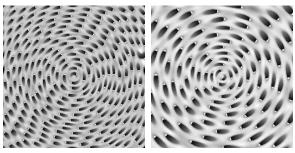


Figure 2: (a) Reaction-Diffusion visualization of a circular vector field showing orientation and direction and (b) magnitude, orientation, and direction.

not form very well, appearing to be smeared together. This is attributable to the variance of the formation and degradation factors, a and b which need to have enough variance to perturb the system but not so much as to cause irregularities in the patterns.

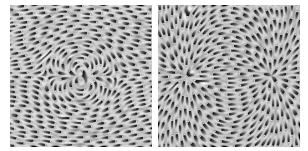


Figure 3: (a) Reaction-Diffusion visualization the electric field from a dipole and (b) an electrostatic charge.

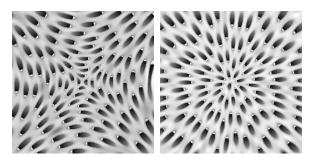


Figure 4: (a) Reaction-Diffusion visualization of a saddle vector field and (b) a sink vector field.

One of the unique features of using a reaction-diffusion model with anisotropic diffusion as we have done is that the spot, although appearing random, naturally align themselves into pseudo streamlines. Further, when the vector field is curved, the spots are not perfectly elliptical, but rather a bean-like shape. This is due to the spatially nonuniform anisotropy influencing the overall spot shape. Another feature of the reaction-diffusion model is that, due to the diffusion, faint streaks emanate from the ends of the spots. These streaks act to connect the spots, further aiding in visualizing the vector field.

This aligning, bending, and streaking all give the observer cues to the underlying vector field. But critical areas may also be of interest. This is another area where the reaction-diffusion model gives visual cues. For instance, at locations where the vector field is diverging, the spots are no longer elliptical but assume odd shapes. If the vector field is diverging equally in all directions the spots will be circular. As such, oddly shaped or circular spots could indicate,

critical areas, or as will be demonstrated in the next section, the location of a large uncertainty in orientation. These are locations the observer may want to inspect further. For instance, in Figures (2-4) the spots are elliptical and are aligned with the vector field throughout the image except at the critical areas.

We have shown that it is possible to view different vector fields using a reaction-diffusion model. However, when implementing the reaction-diffusion method, a question regarding resolution arises: what is the minimum resolution required for individual features to be seen? By its nature, the process of diffusion acts to smooth, lowering high concentrations and raising low concentrations. As such, it is possible to lose individual features that are significantly different than their neighbors.

To determine the minimum resolution at which features can be seen, we oversampled a vector field until it was possible to see the impact of a single vector that was significantly different in both magnitude and orientation than its surrounding, otherwise constant neighbors. This is demonstrated in Figures (5) and (6) for both magnitude and orientation, respectively. It is not until there is an oversampling of eight times the original sample that the magnitude will significantly impact its neighbors to be visually noticeable. Similarly, it takes an oversampling of eight times for the angle to impact its neighbors. Unfortunately, for large vector fields, oversampling is not always practical because it may require significantly more computational time. As such, when visualizing a vector field without oversampling features less than eight nodes in size may be smoothed out.

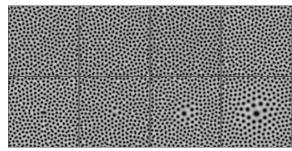


Figure 5: Effect of a single value on the spot size with an oversampling of 0, 1, 2, 4, 8, 16, 32, and 64 times.

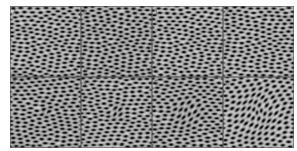


Figure 6: Effect of a single value on the spot orientation with an oversampling of 0, 1, 2, 4, 8, 16, 32, and 64 times.

4.1 Uncertainty Measurements

In the previous examples, we have fixed the amount of anisotropy in the diffusion matrix. However, this is not necessary. If we allow the anisotropy to vary, we can map and visually represent another variable. In this case, we define the amount of anisotropy to be the ratio of the values along the principal axes in the diffusion matrix. When the amount of anisotropy is small, the spot formed is circular. Where when the anisotropy is high, the spot formed is elliptical, at

times, almost to the point of being a thick line. This difference is very well suited to mapping an orientation uncertainty. When the orientation uncertainty is very small the spot maintains its elliptical shape, reflecting a precise orientation. When the uncertainty is very high, the spot is more circular, reflecting the uncertainty in the orientation. In the previous examples, the ratio of the values along the principal axes in the diffusion tensor was fixed at 3:1. We now allow it to vary between 5:3 and 7:1. This is demonstrated in Figure (7) where the uncertainty is a function of the angular position. In a subsequent example, instead of encoding the uncertainty, we encode the vorticity of the vector field,

$$w_{ij} = (\nabla \times u_{ij}), \tag{10}$$

which can be visualized as a scalar in 2D.



Figure 7: Reaction-Diffusion visualization of orientation uncertainty. The orientation uncertainty is a function of the angular position.

4.2 Augmentation and Automatic Streamline/Glyph Placement

Figure (8) shows our reaction-diffusion model for visualizing an idealized vector field that contains three saddles and two vortexes. To augment the LIC image we have taken Kirby et al.'s [12] painter's approach by using a LIC image as an undercoat with the reaction-diffusion as the topcoat. In addition to the trending pattern formed by the spots, the brush stroke appearance from the LIC image enhances the ability of the observer to follow the vector field, while the spots provide magnitude and direction.

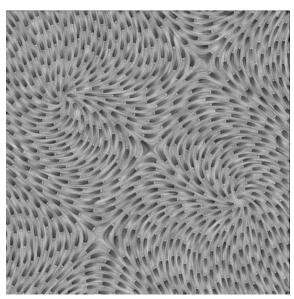


Figure 8: Reaction-Diffusion visualization of multiple vector fields using an underlying LIC image as a base coat.

One of the features of using a reaction-diffusion model is that it also provides a mechanism for automatic streamline or glyph placement. By finding the centroid of each spot using standard image processing techniques [21], it can serve as a seed point for placing a streamline, Figure (9) or a glyph, Figure (10). Although there is one streamline per spot in Figure (9), there appears to be fewer streamlines because of the their alignment and slight overlap. Because the density of the spots is based on the magnitude of the underlying vector field in Figure (10) it was possible to scale the arrow glyphs to reflect the vector magnitude without causing occlusion. Further, because of the density relationship, the number of glyphs also reflects, albeit inversely, the magnitude of the vector field. If a uniform density is desired then all that needs to be done is to run the reaction-diffusion using a constant magnitude term which would then determine the density.

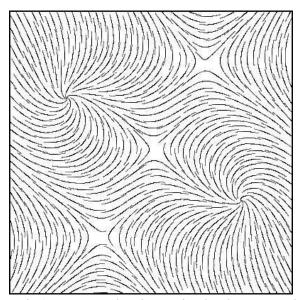


Figure 9: Automatic streamline placement based on the spot centroids in Figure (8).

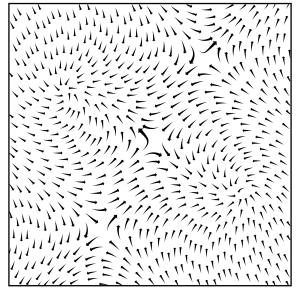


Figure 10: Automatic arrow glyphs placement based on the spot centroids in Figure (8).

5 RESULTS AND DISCUSSION

We now apply our reaction-diffusion model to a numerical simulation of the nonlinear magneto-hydrodynamics (MHD) that occur in the DIII-D tokamak nuclear fusion reactor. The vector field shown in Figures (11-13) is a reoriented two-dimensional slice of the magnetic field in the Tokamak reactor. In Figure (11), just the magnitude and vorticity of the vector field is visualized with no orientation information. The greater the vorticity, the more symmetric the spots become. This gives a good example of how this technique can be used for visualizing two scalar values. Figure (12) is the same vector field showing the orientation and direction. Finally, in Figure (13), the vector field is shown with magnitude, orientation, and direction. In Figure (14), we have encoded the magnitude, orientation, direction, and vorticity of the vector field. Although the spots in Figure (14) appear in different locations than in Figure (13) a comparison of the two images shows a significant difference in the spot shape in those areas with high vorticity.

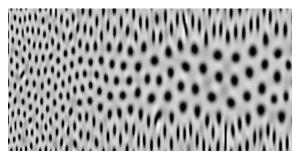


Figure 11: Reaction-Diffusion visualization of a MHD Magnetic vector field. Magnitude and Vorticity are shown.

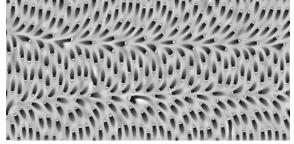


Figure 12: Reaction-Diffusion visualization of a MHD Magnetic vector field. Orientation and direction are shown.

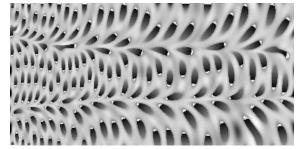


Figure 13: Reaction-Diffusion visualization of a MHD Magnetic vector field. Magnitude, orientation, and direction are shown.

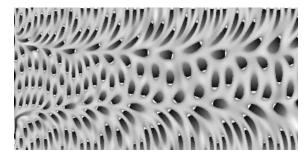


Figure 14: Reaction-Diffusion visualization of a MHD Magnetic vector field. Magnitude, orientation, direction, and vorticity are shown.

5.1 Relationship to Other Vector Field Visualization Techniques

It is also possible to set the reaction rate to zero and use just the diffusion tensor to create other types of visualizations. To produce a LIC-like image, we change the diffusion tensor to be 1D, Figure (15a). To produce convective patterns, such as those proposed by Preußer [20], we use a highly anisotropic 2D diffusion tensor, Figure (15b). This shows that the diffusion model used is consistent with previously published results.

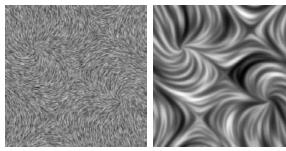


Figure 15: (a) Visualization with a LIC like appearance and (b) with a convective patterns using just the diffusion kinetics.

5.2 Comparison with Other Vector Field Visualization Techniques

We now compare the reaction-diffusion images with four different vector field visualization techniques. Figure (16) shows vector glyphs at regular intervals [25], Figure (17) shows oriented ellipses based on a Poisson distribution [12], Figure (18) shows a line integral convolution [3], and Figure (19) shows image guided streamlines [28]

Placing glyphs at regular intervals is much simpler and quicker than using a reaction-diffusion model, but, as previously discussed, occlusion can be a problem. Using a random Poisson distribution solves the occlusion problem but fails to provide any organization, which is often a key to producing an effective visualization. Using a reaction-diffusion model overcomes the occlusion problem because the spots have a density that is based upon the vector magnitude. Another problem with regular and random intervals is that they may mislead the eye by forming a pattern that may not be part of actual vector field. Conversely, the reaction-diffusion model forms spots in a pattern that follows the underlying structure of the field.

When comparing the reaction-diffusion method to LIC, we can see that both techniques visualize the vector field in a manner that is natural and easy to observe by producing a dense image representation of the vector field field. With reaction-diffusion images, different reaction rates produce spots at different densities. The less dense the spots, the greater the chance that areas of interest may

be missed. However, images with a high density of spots may be difficult to view because of the Moray patterns that can form. As such, the density of the spots is a critical component for an effective reaction-diffusion vector field visualization. Currently, the only way to control the density is by using different reaction rates. However, we are also investigating different reaction-diffusion models, such as the Gray-Scott model.

Unlike traditional LIC images, which do not contain magnitude or direction information, the reaction-diffusion model is able to naturally incorporate this information into the visualization. Including the magnitude and direction greatly enhances the visualization. LIC images, along with other noise-based techniques, can be extended to show the magnitude, but these techniques do so at a loss of vector field detail because of blurring used to emphasize the magnitude [5, 13].

Next, we compare the reaction-diffusion image to a visualization using the image-guided streamline technique developed by Turk and Banks [28]. Both techniques are similar in that both are able to show magnitude, orientation, and direction. However, the reaction-diffusion technique represents magnitude more intuitively than the image guided technique. This is because instead of using length to represent magnitude the reaction-diffusion technique uses a width, which is more intuitive.

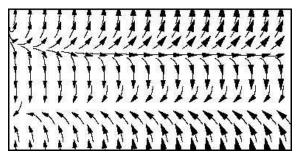


Figure 16: Uniform sampled vector glyph image of the vector field used in Figure (11).

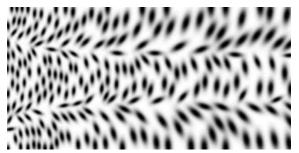


Figure 17: Randomly placed ellipse image of the vector field used in Figure (11).

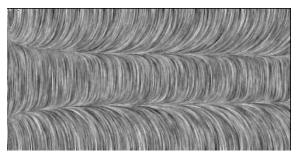


Figure 18: LIC image of the vector field used in Figure (11).

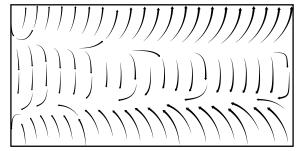


Figure 19: Image guided streamlines generated by Turk and Bank's algorithm of the vector field used in Figure (11).

One of the drawbacks of using a reaction-diffusion model compared to the other techniques is the computational expense. Using an explicit formulation, the patterns take 15-25k iterations to form and become stable. Using a GPU implementation based on Lefohn et al. [15] this takes less than a minute. Whereas it is possible to produce LIC images at interactive rates [2]. We are currently investigating GPU based multi-grid and implicit integration techniques, which should reduce the computational expense.

An additional problem that can occur during pattern formation is that spots can form and fail to separate, as shown in Figure (3a). Where this happens is random and appears to be dependent on the initial conditions. We have observed that it tends to happen more frequently with smaller spots.

One of the greatest benefits of using a reaction-diffusion model is the ability to seamlessly integrate uncertainty measurements in the model. None of the other techniques, with the exception of vector glyphs, are able to show uncertainty as part of their representation [17]. Although vector glyphs can show uncertainty, occlusion remains a problem in their use.

An additional benefit is that the reaction-diffusion technique can be used alone or to augment other techniques. Furthermore, it is possible to use the centroids of the spots to provide a set of seed points for placing streamlines and scaled glyphs.

6 CONCLUSIONS AND FUTURE WORK

We have introduced the use of a reaction-diffusion model that can produce patterns with different shapes, sizes, orientations, and directions for visualizing vector fields. We are able to control the pattern formation by mapping two of the vector field components, orientation and magnitude, to the diffusion and reaction kinetics, respectively. In addition, we also can map an orientation uncertainty to the diffusion kinetics. This mapping produces a spot pattern that is highly representative of the underlying vector field. To show direction we have applied a light to dark fading texture to each spot.

The principle advantage of the reaction-diffusion model over existing vector field visualization techniques is that the pattern size and density that naturally arises from the reaction-diffusion model accurately represents the underlying vector field. Further, the shape of the pattern (e.g. the spots) not only contains information concerning magnitude, orientation, and direction but also may contain other information, such as uncertainty or vorticity.

We have also demonstrated the use of the reaction-diffusion model for the automatic placement of streamlines or glyphs and shown how it can be augment other techniques. Although we have not used color to highlight certain features, one could easily incorporate color to further enhance the visual attributes.

Future work includes extension of the reaction-diffusion algorithm to three dimensions. In such an extension, the reaction kinetics remain the same; only the diffusion kinetics must be extended. The output is a three-dimensional texture that can be volume ren-

dered using various techniques or may be applied to two dimensional surfaces. The image generated would have similar characteristics to those generated by Kindlmann and Weinstein [11] and Chambers and Rockwood [4] and unfortunately suffer from the same visualization problems.

Finally, there are a number of perceptual issues that require further investigation, including a formal user study such as the one performed by Laidlaw et. al. [14] to determine the effectiveness of the reaction-diffusion visualization technique in comparison to other vector field visualization techniques. One area of particular interest is quantifying the effectiveness of the natural patterns that form from using a reaction-diffusion model.

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8 APPENDIX

One of the difficulties in using a reaction-diffusion model is the inherent instability of the system. Below are the parameters used to obtain the stable pattern shown in Figure (1a), which are applied to Eqs. (1-4) using a discrete central difference Laplacian on a uniform grid.

a = 4.0

b = 4.0

 $D_a = 1.0 / 4.0$

 $D_b = 1.0 / 16.0$

 $a = 16.0 \pm 1\%$

 $b = 12.0 \pm 1\%$

s = 1.0 / 64.0

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