

Manuscript version: Published Version

The version presented in WRAP is the published version (Version of Record).

Persistent WRAP URL:

http://wrap.warwick.ac.uk/114476

How to cite:

The repository item page linked to above, will contain details on accessing citation guidance from the publisher.

Copyright and reuse:

The Warwick Research Archive Portal (WRAP) makes this work by researchers of the University of Warwick available open access under the following conditions.

Copyright © and all moral rights to the version of the paper presented here belong to the individual author(s) and/or other copyright owners. To the extent reasonable and practicable the material made available in WRAP has been checked for eligibility before being made available.

Copies of full items can be used for personal research or study, educational, or not-for-profit purposes without prior permission or charge. Provided that the authors, title and full bibliographic details are credited, a hyperlink and/or URL is given for the original metadata page and the content is not changed in any way.

Publisher's statement:

Please refer to the repository item page, publisher's statement section, for further information.

For more information, please contact the WRAP Team at: wrap@warwick.ac.uk



High-resolution Transit Spectroscopy of Warm Saturns

Emily K. Deibert 1 , Ernst J. W. de Mooij 2 , Ray Jayawardhana 3, Jonathan J. Fortney 4 , Matteo Brogi 5 , Zafar Rustamkulov 4 , and Motohide Tamura 6,7,8

Department of Astronomy and Astrophysics, University of Toronto, 50 St. George Street, Toronto, ON M5S 3H4, Canada; deibert@astro.utoronto.ca

² School of Physical Sciences and Centre for Astrophysics and Relativity, Dublin City University, Glasnevin, Dublin, Ireland

³ Department of Astronomy, Cornell University, Ithaca, New York 14853, USA

⁴ Department of Astronomy and Astrophysics, University of California Santa Cruz, Santa Cruz, CA 95064, USA

⁵ Department of Physics, University of Warwick, Coventry CV4 7AL, UK

⁶ Department of Astronomy, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-0033, Japan

⁷ National Astronomical Observatory of Japan, NINS, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan

⁸ Astrobiology Center, NINS, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan

Received 2018 September 24; revised 2018 November 30; accepted 2018 November 30; published 2019 January 18

Abstract

We present high-resolution optical transmission spectroscopy of two sub-Saturn mass transiting exoplanets, HATP-12b and WASP-69b. With relatively low densities and high atmospheric scale heights, these planets are particularly well-suited to characterization through transit spectroscopy, and serve as ideal candidates for extending previously tested methods to lower planetary masses. Using a single transit for each planet, we take advantage of the Doppler cross-correlation technique to search for sodium, potassium, and water absorption features. Our analysis reveals a likely (3.2σ) detection of sodium absorption features in the atmosphere of HAT-P-12b, and enables us to place constraints on the presence of alkaline and molecular species in the atmospheres of both planets. With our results, we highlight the efficacy of ground-based campaigns for characterizing exoplanetary atmospheres and pave the way for future analyses of low-mass planets.

Key words: planets and satellites: atmospheres – planets and satellites: gaseous planets – techniques: spectroscopic

1. Introduction

While thousands of transiting exoplanets have been discovered, relatively little is known about their atmospheric properties and compositions. The main limitation in characterizing a planet's atmosphere lies in the extreme difference in brightness between the planet and its host star, typically preventing direct atmospheric emission from being detected. However, recent progress in the field has made use of transmission spectroscopy during transits, when the light from the host star passes through the exoplanet's atmosphere and allows for the detection of atomic or molecular species.

Charbonneau et al. (2002) used transmission spectroscopy to find evidence of sodium in the atmosphere of the hot Jupiter HD 209458b, marking the first detection of an atmosphere around a planet outside our solar system; this observation was later confirmed by Snellen et al. (2008). Other notable contributions include detections of hydrogen, oxygen, and carbon in an evaporating planetary exosphere (Vidal-Madjar et al. 2003, 2004), as well as detections of carbon monoxide (Snellen et al. 2010) and potassium (Sing et al. 2011). Analyses of transmission spectra have also led to a number of discoveries of molecular compounds in exoplanetary atmospheres, including in particular various water vapor detections (e.g., Deming et al. 2013 and Mandell et al. 2013, among others). More recent work has focused on the detection of atmospheric clouds and hazes (e.g., Pont et al. 2013; Nikolov et al. 2015, and Sing et al. 2016, among others). Despite a great deal of progress in the study of exoplanetary atmospheres, however, transmission spectroscopy has to date almost exclusively targeted high-mass hot Jupiters at relatively low spectral resolutions, with a few exceptions such as the warm Neptune GJ 436b (e.g., Lothringer et al. 2018), and the super-Earths GJ 1214b and 55Cnc e (de Mooij et al. 2013, 2014, and Kreidberg et al. 2014, among

others). In the case of 55Cnc e, Esteves et al. (2017) also conducted a high-spectral-resolution search for water vapor.

The present paper focuses on characterizing the atmospheres of HAT-P-12b (Hartman et al. 2009) and WASP-69b (Anderson et al. 2014), two sub-Saturn mass transiting exoplanets with equilibrium temperatures of ~1000 K and strongly inflated radii of $\sim 0.96R_I$ and $\sim 1.06R_I$ respectively, resulting in large atmospheric scale heights corresponding to \sim 1% of their radii. The atmospheres of both exoplanets have previously been observed, resulting in findings that will guide our analysis. In the case of HAT-P-12b, Line et al. (2013) made use of transmission spectra obtained using the Wide Field Camera 3 on board the *Hubble Space Telescope* to show a lack of expected water absorption features in its atmosphere, suggesting the presence of high-altitude clouds. Sing et al. (2016) followed up with the Space Telescope Imaging Spectrograph to observe HAT-P-12b in the full optical range, and found evidence of cloud or haze aerosols as well as potassium. In the case of WASP-69b, Casasayas-Barris et al. (2017) analyzed two transits observed with the High Accuracy Radial velocity Planet Searcher (HARPS-North) spectrograph and reported a 5σ detection of atmospheric sodium absorption in the D₂ line of the sodium doublet, but not in the D₁ line.

Here we present high-resolution spectroscopy of HAT-P-12b and WASP-69b, focusing on sodium and potassium lines as well as water absorption features. Our paper is structured as follows. In Section 2 we describe the observations obtained for each planet, and in Section 3 we describe data reduction techniques used to correct for various systematic effects. We present our analysis in Section 4 and comment on our results in Section 5. Our concluding remarks are provided in Section 6, and an appendix containing further details on our techniques is included.

Table 1Summary of Observations

Planet	Date (UT)	Instrument/Telescope	Duration (hr)	No. of Frames	Exp. Time (s)	Avg. S/N per Spectral Pixel
HAT-P-12b	2017 May 5	HDS/Subaru	8.4	84	300	0.0043
WASP-69b	2017 Jun 10	GRACES/Gemini	5.2	107	140	0.0014

2. Observations

We observed HAT-P-12b using the Echelle High-Dispersion Spectrograph (HDS; Noguchi et al. 2002) on the Subaru telescope. The observations were taken during a period of 8.4 hr on 2017 May 5th UT, and had a typical spectral resolution of 80,000. The wavelength coverage of the observations was 538–799 nm, with a gap from 658 to 680 nm between the two CCDs of the spectrograph. The observations were taken with the StdRb observing mode⁹ and the #2 image slicer, which allowed for an extremely high resolution but prevented us from removing sky emission lines directly from the spectra (for further details on our sky emission reduction, see Section 3).

Each observation was made with an exposure time of 300 s. The total 8.4 hr observation period represents one transit of HAT-P-12b, with 24 frames during the exoplanet's ingress, transit, and egress.

WASP-69b was observed using the Gemini Remote Access to the Canada-France-Hawaii Telescope (CFHT) ESPaDOnS Spectrograph (GRACES; Chené et al. 2014), which uses the Gemini North telescope in tandem with the Echelle Spectro-Polarimetric Device for the Observation of Stars (ESPaDOns; Donati 2003) at the CFHT. The observations were taken during a period of approximately 5.2 hr on 2017 June 10th UT. The wavelength coverage of the data was 395-1044 nm, but due to low fiber throughput the data were only useful at wavelengths of 420-1010 nm, with maximum sensitivity between 490 and 950 nm. We therefore focused on these wavelengths, and achieved a typical spectral resolution of ~60,000. Our observations made use of the Star-Only (4 slice) mode of the spectrograph, which uses a single fiber to yield extremely high spectral resolution but, again, prevented us from removing sky background emission. The reduction process is detailed in Section 3.

Each observation was made with an exposure time of 140 s. The full observation represents one transit of the exoplanet, with 58 out of 107 frames during the ingress, transit, and egress.

Our observations are summarized in Table 1.

3. Data Reduction

The data were initially reduced using the HDSQL¹⁰ pipeline available at the Subaru observatory (Noguchi et al. 2002) for HAT-P-12b, and the Open source Pipeline for ESPaDOnS Reduction and Analysis (Martioli et al. 2012) for WASP-69b. Following this step, we interpolated all spectra from each set of observations to a common wavelength grid in the telluric rest frame using a cubic spline in order to facilitate our analysis.

Contaminating cosmic rays were removed through median filtering. The data were binned and a threshold of 5 median absolute deviations was applied, with points falling above this threshold flagged as cosmic rays and discarded from further analysis.

We then applied a correction to adjust for the blaze function, which is a grating-dependent variation in brightness at different orders in the data introduced by the use of an Echelle spectrograph. This effect was corrected for on a per-order basis for both sets of observations. We accomplished this by first dividing out a low-order polynomial fit to the ratio of the science spectrum with respect to a reference spectrum—in our case, the first spectrum of the data set—from each individual frame. The divided spectra were then binned and fit with a polynomial, which was evaluated at the wavelength range of the original, unbinned spectra, and divided out of the original spectra in order to remove the effects of the blaze function. The results can be seen in Figures 8 and 9 in Appendix A, where the wavelength-dependent effects of the blaze response in the first panels have been corrected for in the subsequent panels.

3.1. Removal of Telluric Emission Lines

For both planets, strong telluric sodium emission features were visible throughout the observations. As mentioned in Section 2, the image slicers used in our observations prevented us from obtaining separate sky observations, and we thus had to manually remove sky emission lines from our data.

We note that the strength of the sodium emission features in our data seemed to vary in accordance with the change in airmass during the observations. Furthermore, contaminating emission features were *not* present in the spectral absorption lines of potassium, which is the other element targeted in this analysis.

To correct for these emission features, we fit them with the sum of a Gaussian and polynomial function, and subtract off the Gaussian in order to preserve the noise in the data. We found that a first-order polynomial was sufficient in characterizing the shape of the emission features, as they were present only in a very small region of the sodium doublet that could be approximated as linear (see Figure 1). The emission features are removed in this way from both lines of the sodium doublet for every spectrum of both the HAT-P-12b and WASP-69b observations.

Figure 1 shows a comparison of one frame of our HAT-P-12b observations before and after the telluric sodium emission has been corrected for. Similar features, though with varying strengths, were observed throughout both sets of observations.

3.2. Correction of Systematic Effects

After initial corrections, the spectra are largely dominated by contaminating telluric and stellar absorption lines (see Figure 2, and Figures 8 and 9 in Appendix A). To correct for these, we make use of the SYSREM algorithm (Tamuz et al. 2005), which was originally designed to correct systematic effects in a large set of photometric light curves. Following the example of previous atmospheric characterization work, however (e.g., Birkby et al. 2013; Esteves et al. 2017), we take advantage of the fact that SYSREM is well-suited to removing systematic effects that appear in many data sets (in our case, spectra) of the same sample. Due to the large, rapidly changing radial velocities of the exoplanets during their transits, SYSREM can

⁹ https://www.naoj.org/cgi-bin/hds_efs.cgi

¹⁰ https://www.naoj.org/Observing/Instruments/HDS/hdsql-e.html

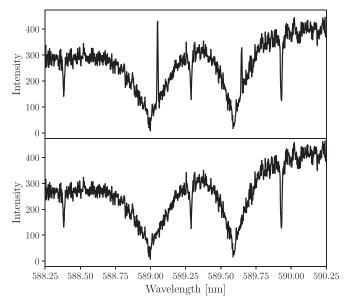


Figure 1. Correction for telluric sodium emission as applied to one HAT-P-12b observation. The top panel shows the data before the correction, with emission features present in both lines of the sodium doublet, while the bottom panel shows the same observation after the features have been removed.

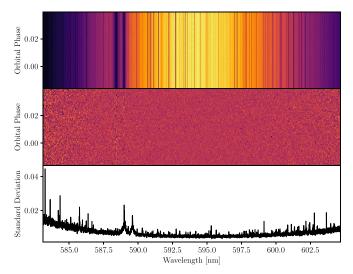


Figure 2. Typical data reduction process used in our analysis, for a WASP-69b spectrum. The top panel shows the data in one order after the blaze correction is applied and cosmic rays are flagged. The middle panel shows the results of the SYSREM algorithm, and the bottom panel shows the standard deviation of each wavelength band after SYSREM.

be used to remove the telluric and stellar absorption features, which are stable in time, without affecting the signal from the exoplanets' atmospheres. We invite the reader to consult Tamuz et al. (2005) for full details of the SYSREM algorithm.

We took the average airmass measured during each exposure as a first approximation to the systematic trends to be removed by the SYSREM algorithm. To account for possible variations in the data between separate orders, SYSREM was applied to each order individually. We tested various numbers of iterations of the algorithm and eventually chose to apply six iterations to each order, after finding that this was sufficient for removing the majority of systematic effects present in the data, minimizing the rms per order, and preventing overfitting. Figures 8 and 9 in Appendix A show the results of this process as applied to all orders, while Figure 2 shows a specific

 Table 2

 Stellar and Planetary Parameters for HAT-P-12b Used in This Analysis

Parameter	Value	References	
Spectral Type	K4	Hartman et al. (2009)	
$R_* (R_{\odot})$	0.701	Hartman et al. (2009)	
$K_p^{\rm a} ({\rm m \ s}^{-1})$	35.8	Hartman et al. (2009)	
$T_{14}^{\ b}$ (days)	0.0974	Hartman et al. (2009)	
$m_p (M_{\text{Jup.}})$	0.201	Mancini et al. (2018)	
$R_p (R_{\text{Jup.}})$	0.919	Mancini et al. (2018)	
$V_{\rm sys}~({\rm km~s^{-1}})$	-40.4589	Mancini et al. (2018)	
Orbital Period (days)	3.21305992	Mancini et al. (2018)	
Mid-transit (BJD)	2455328.49068	Mancini et al. (2018)	
Semimajor Axis (au)	0.0384	Hartman et al. (2009)	
Inclination (degree)	89.10	Mancini et al. (2018)	
R_p/R_*	0.13898	Mancini et al. (2018)	
a/R_*	11.77	Hartman et al. (2009)	
μ_1	0.6508	Claret (2004)	
μ_2	0.1184	Claret (2004)	

Notes. R_* , stellar radius; m_p , planetary mass; R_p , planetary radius.

example of the reduction process as applied to the order containing the sodium doublet in our WASP-69b observations.

We note that the algorithm performs poorly around strong lines. To account for this, we followed Snellen et al. (2010) and Esteves et al. (2017) and weighted each pixel by its standard deviation. Regions in which the algorithm performs poorly are thus suppressed in our analysis. This can be seen in the final panels of Figures 8 and 9 in Appendix A.

4. Analysis

Our initial analysis of the data involved generating light curves as in Snellen et al. (2008). Following their methodology, we integrated the flux within narrow (0.075 nm), medium (0.15 nm), and wide (0.3 nm) passbands around absorption lines of interest to create light curves and look for a signal during the in-transit portion of the data. This was done for both exoplanets in order to search for absorption due to sodium (using the narrow, medium, and wide passband sizes noted above centered on the D₁ and D₂ lines of the sodium doublet, located at 589.592 nm and 588.995 nm respectively) and potassium (again using the narrow, medium, and wide passbands noted above centered on the absorption feature at 769.896 nm). We compared the generated light curves with model light curves made using the methods described in Mandel & Agol (2002), as well as the parameters described in Tables 2 and 3. In both bases, we do not detect any signal; an example is shown in Figure 3 where we present our analysis of the D₁ absorption line in the WASP-69b data.

However, we note that although our light curves were generated using relatively narrow bins, features may have been obscured or diluted by clouds in the exoplanets' atmospheres. To this end, we chose to subsequently employ the Doppler cross-correlation method that was first successfully used by Snellen et al. (2010) to detect carbon monoxide in the atmosphere of HD 209458b, and later used in numerous other works; for example, to detect water in Brogi et al. (2014, 2016), Brogi et al. (2018), and Birkby et al. (2013, 2017), among others. The method involves cross-correlating absorption models with our spectra at a range of Doppler shifts, and then phase-

^a K_p : radial velocity semiamplitude.

 $^{^{\}rm b}$ T_{14} : total transit duration, time between first and last contact.

 Table 3

 Stellar and Planetary Parameters for WASP-69b Used in This Analysis

Parameter	Value	Reference	
Spectral Type	K5	Anderson et al. (2014)	
$R_* (R_{\odot})$	0.813	Anderson et al. (2014)	
$K_p^{\ a} \ ({\rm m \ s}^{-1})$	38.1	Anderson et al. (2014)	
$T_{14}^{\mathbf{b}}$ (days)	0.0929	Anderson et al. (2014)	
$m_p (M_{\text{Jup.}})$	0.260	Anderson et al. (2014)	
$R_p (R_{\text{Jup.}})$	1.057	Anderson et al. (2014)	
$V_{\rm sys}~({\rm km~s}^{-1})$	-9.62826	Anderson et al. (2014)	
Orbital Period (days)	3.8681382	Anderson et al. (2014)	
Mid-transit (BJD)	2455748.83344	Anderson et al. (2014)	
Semimajor Axis (au)	0.04525	Anderson et al. (2014)	
Inclination (degree)	86.71	Anderson et al. (2014)	
R_p/R_*	0.134	Anderson et al. (2014)	
a/R_*	11.97	Anderson et al. (2014)	
μ_1	0.5503	Claret (2004)	
μ_2	0.2012	Claret (2004)	

Notes. R_* , stellar radius; m_p , planetary mass; R_p , planetary radius.

folding the correlation from the in-transit frames and summing over each velocity in order to obtain a map of correlation strength as a function of systemic and Keplerian velocities.

4.1. Atmospheric Models

4.1.1. Alkali Metals + Water Models

Model atmospheres for planets WASP-69b and HAT-P-12b were generated using the methods described in Fortney et al. (2005, 2008). Briefly, under the assumption of chemical equilibrium for solar abundances, a model atmosphere was generated, with the atmospheric temperature structure and chemical abundances arrived at iteratively, given models for the incident fluxes of each planet's parent star, and the assumption of radiative-convective equilibrium. The models represent planetwide average conditions. The model's atmospheric opacities, updated since Fortney et al. (2008), include Barber et al. (2006) for water and Allard et al. (2016) for the alkali metals.

The high-resolution transmission spectra for the planets were generated using the line-by-line code described in Morley et al. (2017). The transmission spectrum code uses the atmospheric temperature profile, atmospheric abundances, and opacities as used in the generation of the model atmosphere. Spectra were calculated at a resolution of R = 500,000.

These models are presented in Figure 10.

4.1.2. Water-only Models

Transmission spectra of WASP-69b and HAT-P-12b were computed with the line-by-line, plane-parallel radiative transfer code utilized for past work on CRIRES data (e.g., de Kok et al. 2014). The code takes as input a prescribed *T-p* profile from Parmentier & Guillot (2014) and corresponding equilibrium abundances from Kempton et al. (2017), as calculated through the NASA Planet Spectrum Generator (Villanueva et al. 2018). The radiative transfer calculations account for H₂–H₂ collision-induced absorption (Borysow et al. 2001; Borysow 2002) and molecular opacity from 157,000 water vapor lines extracted from HITEMP 2010 (Rothman et al. 2010) and accounting for all the main lines contributing to the transmission spectrum,

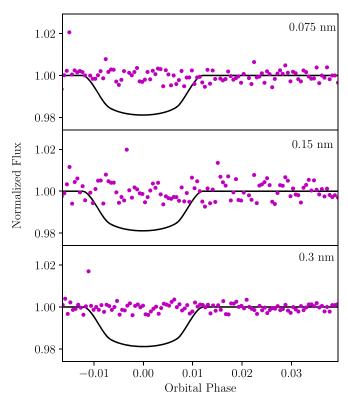


Figure 3. An example of using the Snellen et al. (2008) method to analyze the sodium D_1 line in the WASP-69b data set. The panels represent (from top to bottom) passbands of 0.075, 0.15, and 0.3 nm, and the magenta points correspond to the data while the black line corresponds to a model light curve generated using the methods of Mandel & Agol (2002). We do not detect any signal in any of the three passbands.

plus additional weak lines significantly contributing to the pseudo-continuum. The choice of HITEMP 2010 has led to several detections of water vapor in high-resolution observations of exoplanet atmospheres and is therefore adopted as a benchmark against which other molecular databases can be compared. The radiative transfer is computed across 50 layers of the planet's atmosphere equally spaced in $\log(p)$, and it accounts for the slanted geometry during transit. Additional input parameters of the models are the stellar and planet radii, and planet surface gravity.

These models are presented in Figure 11.

4.2. Doppler Cross-correlation for Alkali Metal Absorption Features

In this section, we describe our search for sodium and potassium absorption features in the atmospheres of HAT-P-12b and WASP-69b. The methodology is similar for both elements in both atmospheres.

We correlated each spectrum around a particular absorption line (either potassium or sodium) with models that had been Doppler-shifted to a range of radial velocities between $-100\,{\rm km\,s^{-1}}$ and $+100\,{\rm km\,s^{-1}}$, with a $1\,{\rm km\,s^{-1}}$ velocity step between each model. At this point, a particularly strong planetary absorption feature would be present in the correlation map as a diagonal line with a slope corresponding to the changing radial velocity of the planet through its transit. On their own, however, the signals from HAT-P-12b and WASP-69b were not strong enough in a single transit to be detected in the correlation map.

 K_p : radial velocity semiamplitude.

 $^{^{\}rm b}$ T_{14} : total transit duration, time between first and last contact.

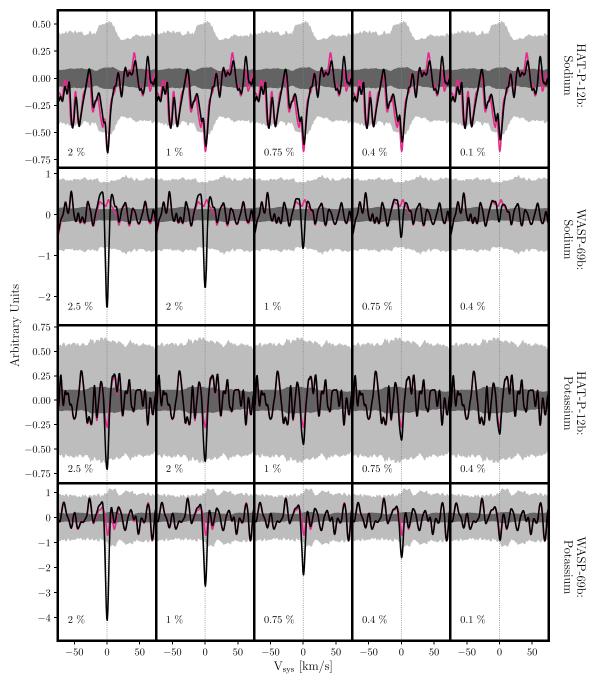


Figure 4. Results of injecting models of various strengths into our data, and repeating the Doppler cross-correlation process. In all panels, the magenta line represents the original data and the black line represents the data with a model injected, where the strength of the injected model is indicated in the bottom left corner of the panel. We note that this percentage refers to a percentage relative to the normalized flux, i.e., a percentage of the stellar surface area. The dark gray contour represents a 1σ confidence level, while the lighter gray contour represents a 3σ confidence level. The relevant planet and element are indicated on the right-hand side of each row. For HAT-P-12b we phase-fold the data to the planet's orbital radial velocity $K_p = 130 \text{ km s}^{-1}$, and for WASP-69b we phase-fold the data to the planet's orbital radial velocity $K_p = 127 \text{ km s}^{-1}$. We also take into account the orbital systemic velocity of each system; see Tables 2 and 3.

To search for a signal from the exoplanet, we therefore phase-folded the correlation signal from the in-transit spectra of the observation. In-transit spectra were determined through the creation of a model light curve generated with the occultquad package by Mandel & Agol (2002) and using limb-darkening parameters obtained from Claret (2004). The parameters used in these calculations are summarized in Tables 2 and 3.

Next we interpolated the correlation signal to the rest frame of the planet and summed over time in order to increase the strength of the signal. Since the planetary radial velocity semiamplitude is not known a priori, we interpolated to values ranging from 1 to $200 \,\mathrm{km \, s^{-1}}$, with a $0.5 \,\mathrm{km \, s^{-1}}$ step between each value. The end result is a correlation map over a range of systemic velocities (V_{sys}) and planetary radial velocity semiamplitudes (K_p) . If a feature is present in the planet's atmosphere, we expect to see it as a strong correlation signal at the planetary radial velocity semiamplitude K_p and a systemic velocity of $V_{\mathrm{sys}} = 0 \,\mathrm{km \, s^{-1}}$.

The results of this process are presented in Section 4 and Figure 4.

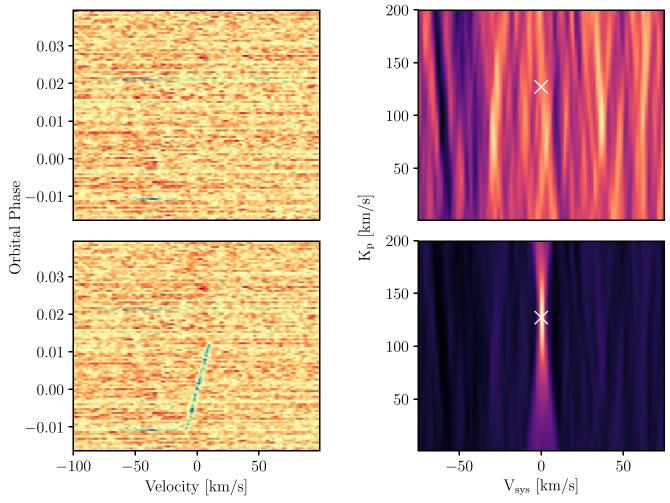


Figure 5. Example correlation maps (left panels) and phase-folded correlations (right panels) obtained from cross-correlating the thousands of water lines across each spectrum of our WASP-69b observations with models at Doppler shifts of -100 km s^{-1} to $+100 \text{ km s}^{-1}$. The color represents the correlation strength. The top panels show the analysis process for our data, and the bottom panels show the same process repeated on the data with a water model injected (see Figure 11). The signal from water is visible in the bottom left panel as a diagonal line across the in-transit portion of the map, corresponding to roughly -0.01 to +0.01 orbital phase. Likewise, the signal is visible at the Keplerian velocity of the exoplanet in the bottom-right panel, at the expected $V_{\rm sys}$ and K_p of the planet: 0 km s^{-1} and $\sim 127 \text{ km s}^{-1}$ respectively. The cross indicates this position in both panels. The fact that no similar signal is visible in the original analysis of our data indicates that water is not present in the atmosphere of WASP-69b at the strength of the injected model.

4.3. Doppler Cross-correlation for Water Absorption Features

Whereas alkali metals produce only a few absorption lines, there are thousands of absorption lines due to water spanning multiple orders in the data. The methodology used to search for water is similar to that described above in Section 4.2; however, rather than correlating over a small window including the relevant absorption lines, we correlate each order of the data with the corresponding wavelength range in the model, and sum the resulting correlation maps. We avoid the orders containing sodium and potassium (orders 3 and 11 for HAT-P-12b and 7 and 16 for WASP-69b), as these are known to be strong absorption features and would dominate the correlation.

A detailed example of the process is presented in Figure 5, and the results of our analysis are presented in Section 4 and Figure 6.

4.4. Model Injection/Recovery Tests

To determine the significance of our results and detection limits of our observations, we performed model injection/ recovery tests for several absorption features at varying strengths for each planet. This was done by multiplying the in-transit frames of the raw data (i.e., the data processed through the reduction piplines available at the telescopes but before any further reduction) by an atmospheric model in the planet's reference frame. We used the atmospheric models described in Section 4.1, but varied the model strength throughout the injection/recovery tests. In this way, we are able to determine the minimum strength at which the model would be recovered by our analysis, and can use this to constrain the atmospheric properties of the planets.

After injecting the models, we followed the data reduction processes described in Section 3 and carried out the Doppler cross-correlation technique described in Section 4.2.

4.5. Detection Significance

If a candidate signal does not surpass 3σ , we do consider it to be significant. If a candidate signal *does* surpass 3σ , we consider the candidate as warranting further investigation. As will be discussed further in Section 5.1, we caution that noise in

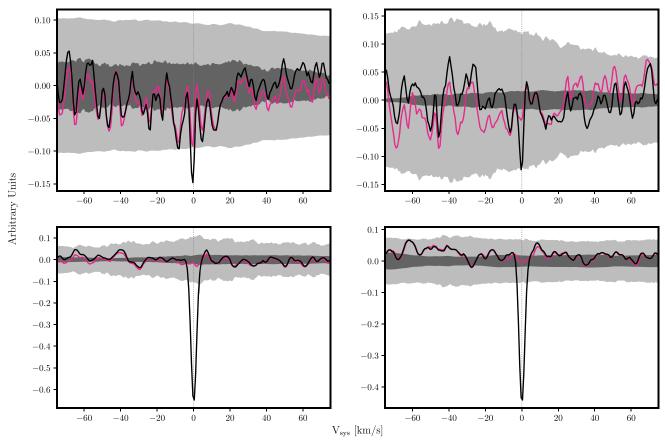


Figure 6. Results of injecting atmospheric absorption models into our data and repeating the Doppler cross-correlation process. In all panels, the magenta line represents the original data and the black line represents the data with a model injected. This process is repeated for each order of the data (aside from those containing strong absorption features due to sodium and potassium). The top row shows our HAT-P-12b observations and the bottom our WASP-69b observations; the left column shows injections with the models containing alkali metals and water, while the right shows injections with the models containing only water.

the data can lead to features approaching the 3σ level, and thus any potential signals must be treated carefully.

In order to determine the significance of our results, we followed the methods described in Esteves et al. (2017). This involved creating 1σ and 3σ confidence levels for each feature in each planet by randomly selecting 24 frames in the case of HAT-P-12b and 58 frames in the case of WASP-69b (corresponding to the total number of in-transit frames for each data set), assigning each an in-transit phase, and carrying out the correlation and phase-folding as described in Section 4.2. The process was repeated 10,000 times in order to determine 1σ and 3σ confidence levels.

5. Results and Discussion

In this section, we present the results of applying the Doppler cross-correlation method as described in Sections 4.2 and 4.3 to our observations. In particular, we are looking for features caused by atmospheric absorption due to sodium, potassium, and water.

5.1. HAT-P-12b

Our results are summarized in Figure 4. As can be seen in the top row of Figure 4, there is a peak at 3.2σ in the data at a systemic velocity of 0 km s^{-1} and a planetary orbital velocity of $\sim 130 \text{ km s}^{-1}$, consistent with a signal due to atmospheric

sodium absorption. This feature has not been observed in any previous analyses of HAT-P-12b (see, e.g., Line et al. 2013 and Sing et al. 2016).

In the case of potassium, we observe a peak in the phase-folded correlations at $V_{\rm sys}=0$ outside of the 1σ contour, but note that it does not exceed the 3σ confidence level. We also note, however, that there are many additional peaks at $>1\sigma$ that are likely due to noise in the data. A similar phenomenon was observed by Esteves et al. (2017) in their analysis of 55 Cnc e's atmosphere, and they found that repeating the Doppler cross-correlation with a pure white noise spectrum resulted in similar features in the phase-folded correlations.

Based on the model injections, we can rule out potassium in the atmosphere of HAT-P-12b down to an amplitude of 2% relative to the normalized flux. Any models injected lower than this are not detected beyond the 3σ level in our analysis.

In the case of water, our results are presented in Figure 6. As in Figure 4, 1σ and 3σ confidence levels are shown, as well as phase-folded correlations for the data alone and the data with an atmospheric model injected. We make use of two separate models (see Section 4.1): one that includes a full atmospheric treatment of each planet, and one that only includes the signal due to water in the exoplanetary atmosphere. The results for HAT-P-12b are in the top row of Figure 6. An injection of the first model yields a detection outside the 3σ level, indicating

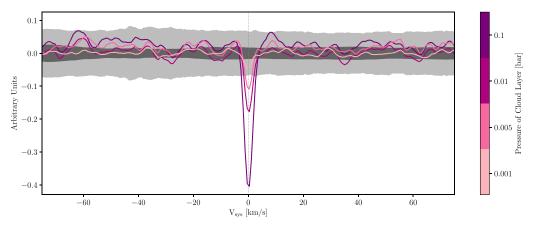


Figure 7. Results of injecting model spectra with cloud layers at various pressures into our WASP-69 data. The pressures corresponding with each layer are indicated in the colorbar. This analysis was performed using the water-only models. We see that at a pressure of 1 mbar, the injected signal is no longer detectable in our data.

that water is not present in the atmosphere of HAT-P-12b at this model strength.

Previous analyses of HAT-P-12b's atmosphere have revealed a cloudy spectrum with very little detectable atmospheric absorption (Line et al. 2013; Sing et al. 2016; Alexoudi et al. 2018), which is consistent with our observations. Sing et al. (2016) and Alexoudi et al. (2018) did detect evidence of weak potassium absorption in the planet's atmosphere, and although we did not confidently observe the same, we note that this could be due to the noise in our data or the fact that only one planetary transit was used in our analysis, resulting in a lower signal-to-noise ratio than would be required to confidently detect atmospheric potassium.

5.2. WASP-69b

As in Section 5.1, the results of cross-correlating our original and injected data with model atmospheric spectra for various absorption features can be seen in Figure 4. In the case of both sodium and potassium, there are no significant (i.e., $>3\sigma$) peaks in the WASP-69b data at a systemic velocity of 0 km s⁻¹. For the sodium correlations (row 2 in Figure 4), we do not observe any features at $V_{\rm sys}=0$, and the majority of features present at other systemic velocities are contained within the 1σ contour. The data with a model injected yields a 3σ detection down to 1% of the flux with respect to the normalized flux, meaning that these model strengths can confidently be ruled out from the planet's atmosphere.

For the potassium correlations (row 4 in Figure 4), we note that there is a peak in the data at $V_{\rm sys}=0$ well outside the 1σ level; however, it does not surpass the 3σ level. Furthermore, injected models at strengths down to $\sim\!1\%$ relative to the normalized flux result in $>\!3\sigma$ peaks in the phase-folded correlations. If potassium is present in the planet's atmosphere, it is therefore at a relatively low amplitude and would require more observations or a greater signal-to-noise ratio to observe.

We also note the presence of a feature slightly exceeding 3σ at $V_{\rm sys} \sim +70\,{\rm km\,s^{-1}}$ in the potassium correlations. We believe that this feature, similar to the many features at $>1\sigma$ discussed in Section 5.1, is caused by remaining noise/systematics in the data. This reinforces the fact that potential candidate signals (i.e., those surpassing 3σ , such as that

detected at $V_{\rm sys} = 0 \, {\rm km \, s}^{-1}$ for sodium absorption in HAT-P-12b as discussed in Section 5.1) must be treated with caution.

Our analysis of water absorption features is presented in Figure 6. In both cases, the analysis of the data with an injected model results in a strong signal well outside the 3σ range, whereas the signal from the data itself does not surpass this level. Thus we conclude that water is not present in the atmosphere of WASP-69b at the strength of either model. Furthermore, the strong signal that is detected when water models are injected into the WASP-69b data allows us to investigate the presence of cloud decks in the exoplanet's atmosphere. Following the analysis presented in Pino et al. (2018), we inject models with a simulated cloud layer into the data, allowing us to place a limit on the altitude at which clouds might exist in the atmosphere of WASP-69b. We choose various cloud pressures and corresponding transit depths in our models, and repeat the cross-correlation after cutting off the model below this depth. In Figure 7, we show injections of models with simulated cloud layers at pressures of 0.1 bar, 0.01 bar, 5 mbar, and 1 mbar. Down to a pressure of 5 mbar the signal is still detected outside the 3σ confidence level; however, a cloud layer at a pressure of 1 mbar does not surpass this level and therefore cannot be ruled out in the atmosphere of WASP-69b.

The atmosphere of WASP-69b has previously been studied by Casasayas-Barris et al. (2017) with the medium-resolution HARPS-North spectrograph. In contrast with our analysis, their observations led to a 5σ detection of atmospheric sodium absorption in just the D_2 line of the sodium doublet, at 588.995 nm (Casasayas-Barris et al. 2017). In order to better compare our analysis with theirs we repeated the Doppler cross-correlation technique as well as the model injection/recovery tests with each line of the sodium doublet separately; however, this did not significantly change our results. In particular, we did not observe a strong correlation signal when carrying out our analysis in just the D_2 line of the sodium doublet.

5.3. The Rossiter–McLaughlin (RM) Effect and Center-to-Limb (CLV) Variations

Here we consider the possible consequences of the RM effect and CLV on our analysis.

The consequences of the RM effect/CLV have been addressed in previous work in the field. In particular, centerto-limb variations of stellar lines may affect transmission spectra (Yan et al. 2017), while the RM effect can introduce misalignments or spurious detections/nondetections into final results (Louden & Wheatley 2015; Barnes et al. 2016). However, we do not think that the RM effect and CLV will have had a noticeable impact on our results. In particular, we note that HAT-P-12 is an extremely slow rotator, with $\nu \sin i_* = 0.5 \pm 0.5 \, \mathrm{km \, s^{-1}}$, as determined in Mancini et al. (2018) using HARPS-N measurements. Furthermore, measurements of the sky-projected orbital obliquity angle λ are unconstrained, but may be consistent with spin-orbital alignment ($\lambda = -54^{+41}_{-13}$).

WASP-69 is a slightly faster rotator, with $v \sin i_* = 2.2 \pm 0.4 \, \mathrm{km \, s^{-1}}$ (Anderson et al. 2014). In their analysis using the HARPS-N spectrograph, Casasayas-Barris et al. (2017) model the RM effect of WASP-69b. They derive a value of $\lambda = 0.4^{+2.0}_{-1.9}$ for the sky-projected orbital obliquity angle. It is found that the RM effect does not significantly impact their results.

6. Conclusion

We have presented our analysis of high-resolution optical transmission spectroscopy of two sub-Saturn mass transiting exoplanets, HAT-P-12b and WASP-69b. While the majority of high-resolution, ground-based efforts to study exoplanetary atmospheres have targeted massive planets, HAT-P-12b and WASP-69b represent an unexplored area of parameter space, with lower masses and cooler temperatures than any hot Jupiters that have previously had their atmospheres characterized via ground-based transmission spectroscopy.

In this paper, we presented the results of the Doppler cross-correlation technique, which takes advantage of the high resolution of our observations and the large radial velocities of our exoplanets in order to correlate sophisticated model spectra with our data in the search for atmospheric absorption features. In particular, we targeted sodium and potassium, two alkaline absorption features that should extend high into the exoplanets' atmospheres. We also used similar techniques to search for water, which is made up of thousands of absorption lines in the wavelength range of our observations.

Our analysis revealed a 3.2σ detection of atmospheric sodium in HAT-P-12b, which has not been detected in any previous analyses of this exoplanet's atmosphere.

E.K.D. is supported in part by an Ontario Graduate Scholarship and an NSERC CGS-M. M.T. is supported by MEXT/JSPS KAKENHI grant No. 18H05442, 15H02063, and 22000005. We thank Abhinav Jindal, Ryan Cloutier, and Lisa Esteves for helpful discussions. The authors acknowledge the role that Roxana Lupu, Richard Freedman, Didier Saumon, and Caroline Morley played in curating the opacity database used in some of these model spectrum calculations.

Appendix A Data Reduction

Figures 8 and 9 show the full data reduction process as applied to each order of the HAT-P-12b and WASP-69b observations, respectively. An explanation of this process is provided in Section 3.

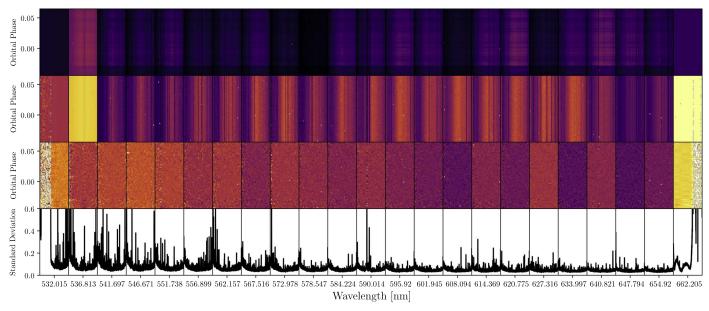


Figure 8. Full data reduction process as applied to the HAT-P-12b observations from the blue CCD. The top panel shows the raw data as they are received after the initial reduction pipeline at the telescope (see Section 3). The second panel shows the results of the blaze correction and cosmic-ray filtering (see Section 3). The third panel shows the data after five iterations of SYSREM have been applied, and the bottom panel shows the standard deviation along each wavelength band of the third panel.

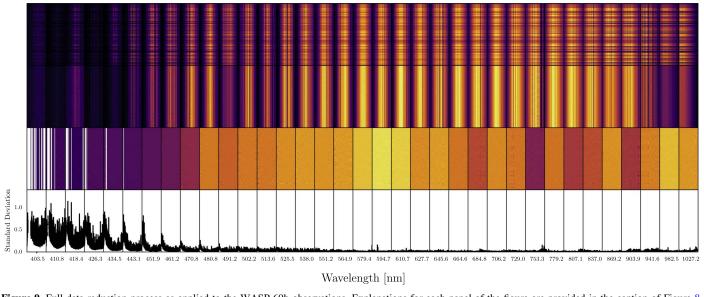


Figure 9. Full data reduction process as applied to the WASP-69b observations. Explanations for each panel of the figure are provided in the caption of Figure 8.

Appendix B Models

The model spectra described in Section 4.1.1 are presented in Figure 10. The top panel shows the model generated for HAT-P-12b, while the bottom panel shows the model generated for WASP-69b. The model spectra

described in Section 4.1.2 are presented in Figure 11. As in Figure 10, the top panel shows the model generated for HAT-P-12b and the bottom panel shows the model generated for WASP-69b.

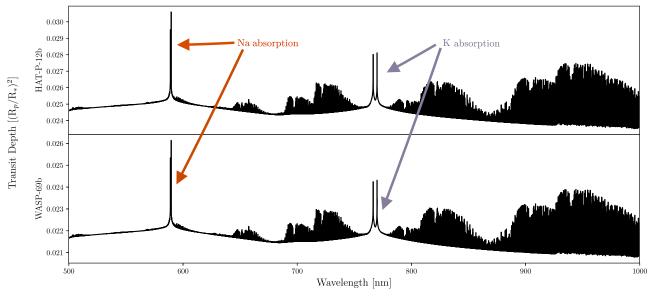


Figure 10. Model spectra used in our analysis. The top panel shows the spectrum calculated for HAT-P-12b, and the bottom panel shows the spectrum calculated for WASP-69b. Water bands can be identified from Figure 11, but the locations of other important species are marked.

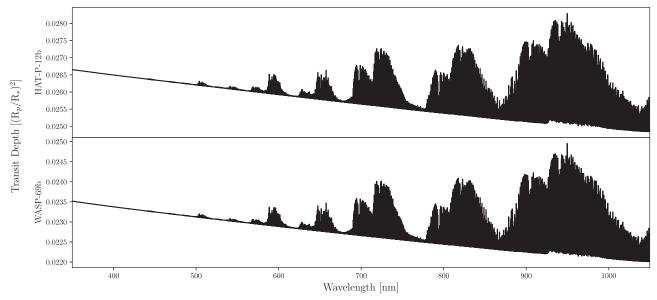


Figure 11. Model spectra used in our analysis. The top panel shows the spectrum calculated for HAT-P-12b, and the bottom panel shows the spectrum calculated for WASP-69b. These models only include lines from water.

ORCID iDs

Emily K. Deibert https://orcid.org/0000-0001-9796-2158
Ernst J. W. de Mooij https://orcid.org/0000-0001-6391-9266

Jonathan J. Fortney https://orcid.org/0000-0002-9843-4354 Matteo Brogi https://orcid.org/0000-0002-7704-0153 Zafar Rustamkulov https://orcid.org/0000-0003-4408-0463 Motohide Tamura https://orcid.org/0000-0002-6510-0681

References

Alexoudi, X., Mallonn, M., von Essen, C., et al. 2018, A&A, 620, A142
Allard, N. F., Spiegelman, F., & Kielkopf, J. F. 2016, A&A, 589, 21
Anderson, D. A., Collier Cameron, A., Delrez, L., et al. 2014, MNRAS, 445, 1114

Barber, R. J., Tennyson, J., Harris, G. J., & Tolchenov, R. J. 2006, MNRAS, 368, 1087 Barnes, J. R., Haswell, C. A., Staab, D., & Anglada-Escudé, G. 2016, MNRAS, 462, 1012 Birkby, J. L., de Kok, R. J., Brogi, M., et al. 2013, MNRAS, 436, L35 Birkby, J. L., de Kok, R. J., Brogi, M., et al. 2017, AJ, 153, 138 Borysow, A. 2002, A&A, 390, 779 Borysow, A., Jorgensen, U. G., & Fu, Y. 2001, JQSRT, 68, 235 Brogi, M., de Kok, R. J., Albrecht, S., et al. 2016, ApJ, 817, 106 Brogi, M., de Kok, R. J., Birkby, J. L., et al. 2014, A&A, 565, 124 Brogi, M., Giacobbe, P., Guilluy, G., et al. 2018, A&A, 615, 16 Casasayas-Barris, N., Palle, E., Nowak, G., et al. 2017, A&A, 608, A135 Charbonneau, D., Brown, T. M., Noyes, R. W., & Gilliland, R. L. 2002, ApJ, 568, 377 Chené, A.-N., Padzer, J., Barrick, G., et al. 2014, Proc. SPIE, 9151, 47 Claret, A. 2004, A&A, 428, 1001 de Kok, R. J., Birkby, J., Brogi, M., et al. 2014, A&A, 561, 150 de Mooij, E. J. W., Brogi, M., de Kok, R. J., et al. 2013, ApJ, 771, 109

```
de Mooij, E. J. W., López-Morales, M., Karjalainen, R., et al. 2014, ApJ, 797, 21
```

```
Deming, D., Wilkins, A., McCullough, P., et al. 2013, ApJ, 774, 95
Donati, J.-F. 2003, in ASP Conf. Ser. 307, Solar Polarization, ed.
  J. Trujillo-Bueno & J. Sanchez Almeida (San Francisco, CA: ASP),41
Esteves, L. J., de Mooij, E. J. W., Jayawardhana, R., et al. 2017, AJ, 153, 268
Fortney, J. J., Marley, M. S., Lodders, K., et al. 2005, ApJ, 627, 69
Fortney, J. J., Marley, M. S., Saumon, D., & Lodders, K. 2008, ApJ, 683, 1104
Hartman, J. D., Bakos, G. Á, Torres, G., et al. 2009, ApJ, 706, 785
Kempton, E. M.-R., Lupu, R., Owusu-Asare, A., et al. 2017, PASP, 129, 974
Kreidberg, L., Bean, J. L., Désert, J.-M., et al. 2014, Natur, 505, 69
Line, M. R., Knutson, H., Deming, D., et al. 2013, ApJ, 778, 183
Lothringer, J. D., Barman, T., & Koskinen, T. 2018, ApJ, 866, 27
Louden, T., & Wheatley, P. J. 2015, ApJL, 814, L24
Mancini, L., Esposito, M., Covino, E., et al. 2018, A&A, 613, 41
Mandel, K., & Agol, E. 2002, ApJL, 580, L171
Mandell, A. M., Haynes, K., Sinukoff, E., et al. 2013, ApJ, 779, 128
Martioli, E., Teeple, D., Manset, N., et al. 2012, Proc. SPIE, 8451, 84512B
Morley, C. V., Knutson, H., Line, M., et al. 2017, AJ, 153, 86
```

```
Nikolov, N., Sing, D. K., Burrows, A. S., et al. 2015, MNRAS, 447, 463
Noguchi, K., Aoki, W., Kawanomoto, S., et al. 2002, PASJ, 54, 855
Parmentier, V., & Guillot, T. 2014, A&A, 562, 133
Pino, L., Ehrenreich, D., Allart, R., et al. 2018, A&A, 619, A3
Pont, F., Sing, D. K., Gibson, N. P., et al. 2013, MNRAS, 432, 2917
Rothman, L. S., Gordon, I. E., Barber, R. J., et al. 2010, JQSRT, 111, 2139
Sing, D. K., Fortney, J. J., Nikolov, N., et al. 2016, Natur, 529, 59
Sing, D. K., Pont, F., Aigrain, S., et al. 2011, MNRAS, 416, 1443
Snellen, I. A. G., Albrecht, S., de Mooij, E. J. W., & Le Poole, R. S. 2008, A&A, 487, 357
Snellen, I. A. G., de Kok, R. I. de Mooij, E. J. W. & Albrecht, S. 2010, Natur, Specifical Control of the Control of th
```

Snellen, I. A. G., de Kok, R. J., de Mooij, E. J. W., & Albrecht, S. 2010, Natur, 465, 1049

Tamuz, O., Mazeh, T., & Zucker, S. 2005, MNRAS, 356, 1466
 Vidal-Madjar, A., Désert, J.-M., Lecavelier des Etangs, A., et al. 2004, ApJ, 604, 69

Vidal-Madjar, A., Lecavelier des Etangs, A., Désert, J.- M., et al. 2003, Natur, 422, 143

Villanueva, G. L., Smith, M. D., Protopapa, S., et al. 2018, JQSRT, 217, 86 Yan, F., Pallé, E., Fosbury, R. A. E., et al. 2017, A&A, 603, A73