

# **Tiago Daniel Teixeira Gonçalves**

Bachelor in Micro and Nanotechnologies Engineering

# Implementation of X-ray reflectivity on the characterization of ultra-thin films for memory devices

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Adviser:	Joana Vaz Pinto, Assistant Professor,
	Faculty of Sciences and Technology,
	NOVA University of Lisbon
Co-adviser:	Hugo Manuel Brito Águas, Associate Professor,
	Faculty of Sciences and Technology,

NOVA University of Lisbon

**Examination Committee** 

Chairperson: Rui Alberto Garção Barreira do Nascimento Igreja Raporteurs: Luís Miguel Nunes Pereira Joana Vaz Pinto



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# Implementation of X-ray reflectivity on the characterization of ultra-thin films for memory devices

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Marco Polo describes a bridge, stone by stone. "But which is the stone that supports the bridge?" Kublai Khan asks. "The bridge is not supported by one stone or another," Marco answers, "but by the line of the arch that they form." Kublai Khan remains silent, reflecting. Then he adds: "Why do you speak to me of the stones? It is only the arch that matters to me." Polo answers: "Without stones there is no arch." — Italo Calvino, Invisible Cities

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In this work X-ray reflectivity (XRR) was implemented as a characterization technique to study very thin films produced at CENIMAT. Three groups of samples were analysed: thin films of  $Ta_2O_5$ , thin films  $Ta_2O_5$  mixed with SiO<sub>2</sub> (TaSiO) and thin films of zinc-tin oxide (ZTO). Spectroscopic ellipsometry was used as a complementary technique and all results were compared with the results of profilometry and Rutherford backscattering spectrometry (RBS).

A good agreement was found between the thicknesses obtained by all techniques. XRR revealed to be a better technique than SE to determine the thickness of the native oxide of silicon. In the same way the determination of the roughness was more accurate with XRR.

The objective of the characterization of the TaSiO and  $Ta_2O_5$  samples was to study the influence of the radio frequency (rf) power in the  $Ta_2O_5$  target during the sputtering process. In the case of the TaSiO samples an increase of the rf power resulted in an increase of the percentage of  $Ta_2O_5$  and Ar in the films, as well as an increase of the density and dielectric constant. On the other side, the short-range order decreased. With the increase of the rf power the  $Ta_2O_5$  samples only revealed a lower short-range order.

The aim of the characterization of the ZTO was to study of the effect of rf power in the ZTO target during the sputtering process, as well as the influence of the  $H_2$  and  $O_2$  flows. The increase of the rf power resulted in denser films with lower short-range order. In a less clear way a change in the  $H_2$  influenced the films thickness. The increase of the  $O_2$  flow had as a result films less dense and less rough.

Keywords: X-ray reflectivity, spectroscopic ellipsometry, thin films, Ta<sub>2</sub>O<sub>5</sub>, ZTO

Neste trabalho foi implementada a técnica de refletometria de raios-X (XRR) para a caracterização de filmes muito finos produzidos no CENIMAT. Três grupos de amostras foram analisadas: filmes finos de  $Ta_2O_5$ , filmes finos de  $Ta_2O_5$  misturado com SiO<sub>2</sub> (TaSiO) e filmes de óxido de zinco e estanho (ZTO). Elipsometria espetroscópica (SE) foi utilizada como técnica complementar e os resultados foram comparados com os resultados de perfilometria e espectroscopia de retrodispersão de Rutherford (RBS).

Houve uma boa concordância entre as espessuras obtidas para todas as técnicas. XRR revelou ser uma melhor técnica do que SE para a determinação da espessura do óxido nativo das bolachas de silício. Da mesma forma, a determinação da rugosidade foi mais precisa com XRR.

O objetivo da caracterização das amostras de TaSiO e Ta<sub>2</sub>O<sub>5</sub> era estudar a influência da potência de rádio frequência (rf) no alvo de Ta<sub>2</sub>O<sub>5</sub> durante o processo de pulverização catódica. No caso das amostras de TaSiO o aumento da potência de rf resultou num aumento da percentagem de Ta<sub>2</sub>O<sub>5</sub> e Ar nos filmes, assim como num aumento da densidade e da constante dielétrica. Por outro lado, a organização a curto alcance diminuiu. Com o aumento da potência de rf as amostras de Ta<sub>2</sub>O<sub>5</sub> apenas revelaram uma diminuição da organização a curto alcance.

A caracterização das amostras de ZTO tinha como objetivo estudar o efeito da potência de rf no alvo de ZTO durante o processo de pulverização catódica, bem como a influência do fluxo de  $H_2$  e de  $O_2$ . O aumento da potência de rf resultou em filmes mais densos e com menor organização a curto alcance. De forma menos clara a mudança do fluxo de  $H_2$  influenciou as espessuras dos filmes. O aumento do fluxo de  $O_2$  teve como resultado filmes menos densos e menos rugosos.

**Palavras-chave:** Refletometria de raios-X, elipsometria espectroscópica, filmes finos, Ta<sub>2</sub>O<sub>5</sub>, ZTO.

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- $\Delta$  Change in the polarization of the light after the interaction with the material.
- $\psi$  Change in the amplitude of the components of the electric field after the interaction with the material.
- $\lambda$  Wavelength.
- $\theta$  Angle.
- $\chi^2$  Error function.
- $\varepsilon_{\infty}$  High frequency dielectric constant.
- $\varepsilon_1$  Real part of the dielectric function.
- $\varepsilon_2$  Imaginary part of the dielectric function.
- $\alpha_c$  Critical angle.
- $\sigma_{n,n+1}$  Roughness or diffusenes of a layer.
- *A* Transition matrix element.
- *C* Broadening term.
- *E* Photon energy.
- $E_0$  Peak transition energy.
- $E_g$  Band gap energy.
- *Q* Momentum transfer.
- $d_K$  Distance between two peaks of the Kiessing fringes.
- $k_z$  Wave factor.
- $r_p$  Amplitude reflection coefficients of the incident parallel polarized wave.
- $r_s$  Amplitude reflection coefficients of the incident perpendicularly polarized wave.

AFM	Atomic Force Microscopy.
CENIMAT	Materials Research Center.
DRAM	Dynamic Random Access Memory.
GIZO	Indium-Gallium-Zinc Oxide.
MIS	Metal-Insulator-Semiconductor.
RBS	Rutherford Backscattering Spectrometry.
rf	radio frequency.
RRAM	Resistive Random Access Memory.
SE	Spectroscopic Ellipsometry.
SLD	Scattering Length Density.
TaSiO	mixture of $Ta_2O_5$ and $SiO_2$ .
TFT	Thin-Film Transistor.
XRR	X-ray Reflectivity.
ZTO	Zinc-Tin Oxide.

Nowadays the memory devices are based on thin films that are becoming even thinner. However the development of thinner films leads to a problem: the characterization of those films. With thicknesses that can be lower than 10 nm, the implementation of techniques capable of measure such thicknesses are crucial in order to continue the scientific research in the field of memory devices.

Taking that in mind the main objective of this work was the implementation of X-ray Reflectivity (XRR) to characterize thin and ultra-thin films at Materials Research Center (CENIMAT). Some of the XRR experiments were performed before the beginning of the thesis but due to the lack of a data analysis method, the determination of the properties of the films was not possible. The implementation of such method turned possible the use of XRR at CENIMAT. By using spectroscopic ellipsometry (SE) one aimed to confirm the success of the XRR data analysis method as well as study the limitations and information one can get from both techniques.

To test the analysis method, some sets of samples of the research group were analysed. Films of  $Ta_2O_5$  mixed with SiO<sub>2</sub>,  $Ta_2O_5$  and zinc-tin oxide (ZTO) were used. The variation of the used radio frequency power in the target during the sputtering process was studied in all samples. In the ZTO samples the variation of the H<sub>2</sub> and O<sub>2</sub> was also studied.

#### Motivation

One must do everything that is possible with the available resources. That was the spirit behind this work. CENIMAT has a diffractometer that enables XRR characterization. The implementation of a protocol to analyse the XRR measured data would turn possible the use of the diffractometer for another characterization technique (besides X-ray Diffraction) with no additional costs.

Besides that, until the beginning of this thesis the characterization of some thin films produced at CENIMAT had been a challenge. Semiconductor films for thin-film transistors with thicknesses between 20 and 40 nm, dielectric and semiconductor films produced by solution process and storage layers in memory devices below 10 nm are some examples where other characterization techniques can not be successfully used.

In order to fill the need of using a technique that could be useful in the characterization of such films and by taking advantage of the possibility that XRR data can be obtained at CENIMAT, the implementation of a method to analyse XRR data was crucial. Using SE that already was a very well established technique at CENIMAT, it was possible to complement the XRR characterization with SE. So, a more complete characterization of the samples was achieved.

## **1.1 X-ray reflectivity**

Nanostructured devices are becoming more used in many fields like microelectronics, optoelectronics and sensors. These devices use thin films that are becoming thinner with the evolution of fabrication processes. With this reduction of thickness and with the use of multilayers films, some physical properties of the films and their interfaces become crucial for the correct function of the devices [1].

X-ray Reflectivity (XRR) is a non-destructive method of characterization that can analyse crystalline and amorphous single and multilayer thin films. This technique takes advantage of the fact that matter has a refractive index for X-rays smaller than one. As a result, for angles below a specific critical angle, the incident X-rays are reflected. In other words, occurs a total external refraction. This physical property makes it possible to get information about the material density. Besides that, with this technique it is possible to obtain the film thickness (precision below 1 nm) and the roughness of the surface and interfaces [1, 2].

The critical angle  $(\alpha_c)$  can easily be obtained from the angle from which the measured intensity drops rapidly. After being obtained, the critical angle can be used for a density analysis. The thickness of a film can also be easily obtained and this parameter is related to the angular distance between oscillation waves in the intensity of reflected light -Kiessing fringes  $(d_K)$  [1]. As an example, a XRR curve is plotted in Figure 1.1 illustrating the determination of the  $\alpha_c$  and  $d_K$ .



Figure 1.1: Data obtained by XRR of a tantalum pentoxide film on a silicon wafer.

The Kiessing fringes are due to the constructive and destructive interactions between the reflected waves. An explanation of this interactions is represented in Figure 1.2. In this figure one takes as an example a sample with a single film. Part of the incident beam is reflected and other the part penetrates the film. The part of the beam that penetrates the film will be divided again: part will be reflected in the interface between medium 2 and 3 and part will penetrate the medium 3 and so on. All the components of the initial incident beam that will go out to medium 1 will interact with each other [3]. Here it is considered just the specular part of the beam (when the angle of the incident beam is equal to the angle of the reflected beam) [1].



Figure 1.2: Representation of the light interactions in a single layer system that leads to the creation of the Kiessing fringes.

It is usual to use mathematical formalism to analyse the measured data. In this work the analyses of XRR data were performed in the *MOTOFIT* package [4]. This package uses the Abelès matrix formalism [5]. In Appendix A is possible to find more information this formalism.

There are a few set-ups that can be used to perform XRR analyses. In the set-up used in this work both the X-ray source and the detector moved and the sample is kept fixed. The other possibility in rotate the sample at the same time as the source moves. In this case the detector is fixed. Nevertheless, in these two set-ups the angle between the source-sample and sample-detector is the same, as one only uses the specular part of the reflected beam [1, 2].

Despite same differences, in all set-ups a high intense X-ray source and a low noise detector is necessary. This is due the fact that the range of orders of magnitude of the measure intensity can be higher than 5. In the example shown in Figure 1.1 the range of the measured intensity has almost 6 orders of magnitude. Not only that but the alignment of the sample is crucial and should be performed carefully [1, 2].

In this thesis XRR was used only to know thicknesses, roughnesses and Scattering Length Density (SLD) (a parameter proportional to the mass density) in a dry medium. Despite that, XRR has been used in more applications. For example, XRR can be used to study layers in a liquid medium [6] as well as liquid-liquid and air-liquid interfaces [7], to do image reconstruction of the samples [8], to determine the element depth profiles in low contrast multilayers systems [9], to determine the electron density profile of a electrolyte solution inside nanotubes [10] or even to calibrate ellipsometers [11].

# 1.2 Spectroscopic ellipsometry

#### 1.2.1 Theory

Ellipsometry is a non-destructive optical characterization technique that characterizes light that is reflected from a sample. It can also characterize transmitted light, but this work focuses only on the reflection mode. It measures the change that occurs in the polarization of light with the interaction light-material. This technique is named «ellipsometry» because, with this interaction, linearly polarized light usually becomes «elliptical» [12].

Light is an electromagnetic wave composed by an electric field and a magnetic induction which are perpendicular to each other. The electric field is used to describe the polarization state of light. For that, when a light beam is travelling along the z axis, its associated electric field can be decomposed in two components: one parallel to the x axis and another parallel to the y axis (Figure 1.3a). These components are described by the propagation number (that depends on the wavelength), their initial magnitude and initial phase [12].



(a) Linear polarization

(b) Elliptical polarization

Figure 1.3: Representation of the components of the electric field of a light beam [12].

The polarization state is directly correlated with the difference between the two components, in the amplitude and in phase. When there is no difference in the initial phases the result is an electromagnetic wave with planar polarization state. When this difference is not zero, the synthesized vector rotates with the propagation of light, which results in an elliptical polarization (Figure 1.3b). The particular case when the difference is equal to  $\pi/2$  the polarization state is circular [12].

With the interaction with matter, the phase of the electric field can change and this change depends on the refraction index. Because of that, one of the values measured in ellipsometry to characterize the sample is the change in the polarization of the light, witch is represented by  $\Delta$  [12].

Besides the change in phase, the amplitude of the components of the electric field can also change with the reflection. The change is represented by  $\psi$ , that is defined by:

$$tg(\psi) = \frac{r_p}{r_s} \tag{1.1}$$

where  $r_p$  and  $r_s$  are the amplitude reflection coefficients of the incident parallel polarized wave  $(r_p)$  and the perpendicularly polarized wave  $(r_s)$  to the plane of incidence. These coefficients are given by Fresnel equations [12].

These two values ( $\psi$ ,  $\Delta$ ) express the change in phase and amplitude that appends in light with the reflection on the sample. So, in ellipsometry, the change in the emitted light is measured as the change in the polarization state and this is represented by the fundamental equation of ellipsometry [12]:

$$\rho = tg(\psi)exp(i\Delta) \tag{1.2}$$

If one takes as an example a simple structure,  $\psi$  characterize the refractive index of the sample and  $\Delta$  describes the extinction coefficient. Therefore, these two values measured by Spectroscopic Ellipsometry (SE) describe directly two properties of the sample. This example represents the basic principle of elipsometry [12].

Until here, it has been described the ellipsometry technique, however in this work it was used SE that has the same fundaments of the ellipsometry but uses a light spectra instead of a monochromatic light. Using SE today it is possible to determine many physical properties of the sample with the two values measured ( $\psi$ ,  $\Delta$ ) as function of the wavelength or energy, namely the complex refractive index, complex dielectric function, absorption coefficient, film thickness, band gap and free-carrier absorption. To obtain these properties, it is necessary to create optical models to analyse the SE data. For the creation of those models, it is crucial to have as much information as possible of the sample, like approximate thicknesses and materials [12].

A spectroscopic ellipsometer is required to do SE analyses. In the system used in this thesis, the light ray comes from a polychromatic light source, passes through a linear polarizer and then the sample is subjected to this incident light. After the reflection, light (now elliptical polarized) passes through a periodic modulator and an analyser. Finally, light reaches the detector that measures the intensity of light [12]. A scheme of this system can be found in the figure below.



Figure 1.4: Scheme of the spectroscopic ellipsometer used in this work (adapted from [12]).

#### 1.2.2 Tauc-Lorentz model

In order to analyse the experimental data, one needs a mathematical function. One of the most well-known is the Tauc-Lorentz dispersion formula [13]. This formula is the junction of the Tauc joint density states [14] and the Lorentz calculation of the imaginary part of the dielectric function ( $\varepsilon_2$ ) [15]. This is given by [13]:

$$\varepsilon_{2}(E) = \begin{cases} \frac{AE_{0}C(E-E_{0})^{2}}{\left(E^{2}-E_{0}^{2}\right)+C^{2}E^{2}}\frac{1}{E}, & E > E_{g} \\ 0, & E \le E_{g} \end{cases}$$
(1.3)

where A is the transition matrix element, C is the broadening term, E is the photon energy,  $E_0$  is the peak transition energy and  $E_g$  is band gap. All these parameters are in energy units and, except E, are fitting parameters. E is the variable and the fitting is performed as function of E.

Doing a Kramers–Kronig integration of the equation 1.3 it is possible to obtain the real part of the dielectric function. With this integration a new fitting parameter appears,  $\varepsilon_{\infty}$ , which represents the high frequency dielectric constant [13].

# **1.3** $Ta_2O_5$ films

Given the decrease of devices dimensions, it is important to find materials with high dielectric constant to be an alternative to  $SiO_2$  [16–18]. One of the most studied materials is the  $Ta_2O_5$  that has a high dielectric constant, high refractive index [19, 20] and has been studied using both XRR [21–23] and SE [17, 24–27]. Due to its properties,  $Ta_2O_5$  has been used in several applications such as Dynamic Random Access Memory (DRAM) [19, 28] and redox-based Resistive Random Access Memory (RRAM) [29].

Taking into account the number and diversity of applications, it is crucial to understand the relation between the fabrication methods and the film's properties. For example, a change in the dielectric constant of  $Ta_2O_5$  with variations in the deposition process has already been reported [30, 31]. Furthermore, for the  $Ta_2O_5$  application as redox-based RRAM has been shown that the change in the film density can affect the device behaviour [32].

In order to improve the properties of  $Ta_2O_5$  films to be used as dielectric layer in transistors, a mixture of  $Ta_2O_5$  and SiO<sub>2</sub> (TaSiO) can be used. TaSiO film presents a lower leakage current, a lower hysteresis magnitude and a faster recovery of the threshold voltage, compared to a  $Ta_2O_5$  film. Besides that, the use of a substract bias can even lower the leakage current. So, using a mixture with SiO<sub>2</sub> and substract bias, one can have thinner films and with better properties [24, 33].

In this work TaSiO films were produced by radio frequency (rf) co-sputtering with subtract bias. The objective was to study the influence of the rf power in the  $Ta_2O_5$  target in the thickness, density and roughness of the films. After that, a study of the same

nature was performed in  $Ta_2O_5$  films, also produced by rf sputtering with subtract bias. The fabrication conditions were the same for both studies to be possible to compare the results.

## 1.4 ZTO films

Thin-Film Transistor (TFT) is a device with many applications such as displays and solar cells. In order to have a better performance of TFTs for this application, amorphous oxide semiconductor has been studied to substitute technologies like a-Si:H TFTs. This is due to the excellent uniformity in large areas, high transparency and good electric performance [34–36].

Indium-Gallium-Zinc Oxide (GIZO) is one of the amorphous oxide semiconductors used. This material presents a good electrical performance in displays and solar cells and can be fabricated at room temperature [35]. However, according to a report published in January 2018 by the European Commission [37], both indium and gallium are in the list of critical raw materials and have an end-of-life recycling input rate close to 0%. So, indium- and gallium-free TFTs have been studied.

Zinc-Tin Oxide (ZTO) has been widely studied as an alternative to GIZO due to its performance and the fact that it does not use critical raw materials. It presents many advantages such as being cheap, non-toxic, having good physical robustness, film smoothness, chemical stability and good transparency [38–42]. Recently was shown that ZTO can be used to fabricate flexible TFTs at low temperature. Not only that but it was also demonstrated that the electric performance of the TFTs is improved with the incorporation of hydrogen in the ZTO layer [34].

In the present work ZTO films were produced using different parameters. With this approach one aimed to study the influence of parameters like  $H_2$  and  $O_2$  flow during the rf sputtering fabrication in the density, roughness and thickness of the films.

#### Methods

#### 2.1 Characterization techniques

XRR was performed using the PANalytical's X'Pert PRO MRD X-ray diffractometer with a monochromatic source of CuK $\alpha$  ( $\lambda = 1.540598$ Å). To perform a XRR analysis both the X-ray source and detector move, changing the angle ( $\theta$ ) relatively to the sample. The  $2\theta$  scan range used in this work was from 0.1000° to 4.0000° with a step size of 0.0050°. With this analysis one can take the thickness, roughness and SLD for each film. The data analysis was done in the *MOTOFIT 4.1* package [4]. A tutorial of the data analysis can be found in Appendix B and all error functions of the performed fits in the Appendix C.

However, the files obtained by XRR cannot be directly used in *MOTOFIT 4.1*. As an output, the diffractometer gives the angle  $(2\theta)$  and the respective measured intensity. As an input, *MOTOFIT 4.1* needs the photon momentum as well as the respective normalized intensity. To create files with the data properly converted, a function was made in *Python 3.6.3*. The code is in Appendix D.

To obtain the real and imaginary part of the dielectric function a HORIBA-Jobin Yvon Uvisel spectroscopic ellipsometry system was used with an incident angle of 70° and in a spectral range between 1.5 and 6.5 eV. The acquired data was analysed with *DeltaPsi* 2 software (HORIBA) and all error functions of the performed fits can be found in the Appendix C.

The models used to analyse the SE data were composed by the following layers: crystalline silicon as a substract, SiO<sub>2</sub> (silicon native oxide), the material to be analysed and a roughness layer (50 % of void and 50 % of the material) (Figure E.1). For the crystalline silicon it was used a reference model of the software. The model of the SiO<sub>2</sub> was fitted using the Tauc-Lorentz dispersion formula [13] at the same time as its thickness (with an initial guess of 20 Å) from a silicon wafer. Then, that model of the SiO<sub>2</sub> was used in all samples without fitting. For the films to be analysed both the parameters of the Tauc-Lorentz dispersion formula [13] and thickness were fitted. The initial guesses for the thicknesses were the thicknesses obtained by the other techniques or the values one expected from the fabrication methods. In the samples where the thickness of the roughness layer was close to 0.1 Å the layer was not used because that is the lower result of a thickness fitting, so it is meaningless.

In order to have a more direct measurement of the roughness of the film, Atomic Force Microscopy (AFM) analyses were performed using Asylum MFP-3D system in tapping mode. The AFM images analysis were done in the *Gwyddion* software. Also, to have a more direct measure of some film thicknesses, a XP-200 profilometer from Ambios Technology were used. The compositions of the films were obtained by Rutherford Backscattering Spectrometry (RBS) using a 2.5 MV van de Graaf Accelerator with a 2 MeV He<sup>+</sup> beam. The RBS spectra were analysed with the IBA DataFurnace NDF software [43]. All error functions of the performed fits are in Appendix C.

The characterization of the Metal-Insulator-Semiconductor (MIS) devices were performed in Keysight B1500A Semiconductor Device Analysis. A scanning range from 5 to -11 V with a frequency of 100 kHz were used. The metal contact was made of molybdenum and was the grounded contact (Figure E.2).

#### 2.2 Materials fabrication

#### 2.2.1 Ta<sub>2</sub>O<sub>5</sub> samples deposition process

The first samples to be analysed were  $Ta_2O_5$  films. With this samples one aimed to study the influence of rf power in the  $Ta_2O_5$  target in the density and roughness of the films and to assess how XRR can be used to extract information. In order to achieve that, six samples were prepared with a range of 0.31 to 1.85 W cm<sup>-2</sup> rf power in the  $Ta_2O_5$  target.

All films were deposited in pieces of a 4" p-type silicon wafer with a resistivity from 1 to 2  $\Omega$  cm by rf sputtering from a 2" Ta<sub>2</sub>O<sub>5</sub> target. The deposition pressure was 0.3 Pa and a flow ratio of Ar/O<sub>2</sub> was 14/1. It was also used a 15 W sample bias.

A piece of the same wafer was analysed to obtain the  $SiO_2$  thickness. The goal was to be sure of the most amount of properties possible so one has less fitting parameters in the SE and XRR data analysis.

#### 2.2.2 Ta<sub>2</sub>O<sub>5</sub>/SiO<sub>2</sub> samples deposition process

Two sets of samples were produced with two different objectives. The main goal was to produce films to be used as the dielectric material in transistors and in MIS capacitors. Here a thickness of 250 nm were desired for the dielectric layer. In order to achieve that, a first set of samples were fabricated to calculate the growth rate of the films. Knowing the growth rate, a second set of samples were produced with the desired thickness of 250 nm.

The deposition process of these samples was similar to the  $Ta_2O_5$  samples. Nevertheless these films were deposited by rf co-sputtering from a 2"  $Ta_2O_5$  and a 2"  $SiO_2$  AJA systems targets. Besides that, six samples for each set were produced with a rf power of  $1.23 \,\mathrm{W \, cm^{-2}}$  in the  $SiO_2$  target and in a range of 0.31 to  $1.85 \,\mathrm{W \, cm^{-2}}$  in the  $Ta_2O_5$  target. A sample with only  $Ta_2O_5$  was also produced with a rf power of  $1.23 \,\mathrm{W \, cm^{-2}}$ .

#### 2.2.3 ZTO samples deposition process

In order to study the influence of hydrogen in ZTO films to be used as the dielectric layer in TFTs, several films were produced at different fabrication conditions.

All films were fabricated in a 4" p-type silicon wafer with a resistivity from 1 to  $2\Omega$  cm by rf sputtering from a 2" ZTO AJA systems target. The time of pre-sputtering used was 15 min, the deposition pressure of 0.3 Pa with 50 sccm of Ar flow. The conditions that changed between samples were: the rf power in the target (combined with a change of the deposition time) and the flow of both O<sub>2</sub> and H<sub>2</sub>.

## **Results and discussion**

## **3.1** $Ta_2O_5$ films

The first set of samples considered here consisted of  $Ta_2O_5$  films grown by rf sputtering on silicon. These films are usually used as the dielectric layer in devices such as TFTs at Materials Research Center (CENIMAT). This set of samples was produced with different rf powers (from 0.3 to 1.85 W cm<sup>-2</sup>) in the  $Ta_2O_5$  target during the sputtering process and using a substrate bias of 15 W. With this it was aimed to study the influence of the rf power in the fabricated films.

In the following table it is possible to found the nomenclature used in the section. The nomenclature was chosen in function of the rf power.

Sample	Rf power (W cm <sup>-2</sup> )
Ta1	0.30
Ta2	0.62
Ta3	0.93
Ta4	1.23
Ta5	1.54
Ta6	1.85

Table 3.1: Nomenclature of the Ta<sub>2</sub>O<sub>5</sub> samples.

Before the characterization of the  $Ta_2O_5$  films, a piece of the same silicon wafer was characterized. The objective was to lower the number of the fitting parameters of the data analysis. The analysis of the silicon wafer can be found in Appendix F. With that analysis one concluded that it is possible to use a native oxide thickness of 20 Å in the SE data analyses.

An important remark about the XRR is that it was used the theoretical SLD of Si and SiO<sub>2</sub>:  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup>, respectively (more details in Appendix F). The XRR data was analysed using the procedure presented in Appendix B.

In Figure 3.1 it is possible to see the measured data of two samples with the respective fits: the sample with a low rf power (Ta2) and the sample with a high rf power (Ta6). Just with this figure it is noticed that the sample Ta1 should have a similar density to the sample Ta6 because it has a similar critical angle ( $\alpha_c$ ). However, it has a bigger distance between two peaks of the Kiessing fringes ( $d_K$ ) so it is clearly thinner. The results of the fits will be shown and discussed below in more detail but only with the measured data one can already conclude that the increase of rf power seems to originate thicker films.



Figure 3.1: Measured data by XRR and respective fits of Ta1 and Ta6 samples.

The first parameter that will be discussed is the thickness of the  $Ta_2O_5$  films. The obtained thicknesses are shown in Table 3.2. As one can see both techniques are in excellent agreement and the results for the thickness have a very low uncertainty. This is the first sign of the success of the implementation of XRR. Nevertheless, it is worth to point out a discrepancy in the case of the thickness of the Ta1 sample. This will be discussed later.

Sample	Deposition	Ta <sub>2</sub> O <sub>5</sub> thickness (Å)		Ta <sub>2</sub> O <sub>5</sub> growth rate (Å/min)		SiO <sub>2</sub> thickness (Å)
×	time (min)	XRR	SE	XRR	SE	XRR
Ta1	30	$19.1 \pm 0.4$	$28.9\pm0.3$	0.64	0.96	$17 \pm 8$
Ta2	30	$133.3\pm0.9$	$136 \pm 2$	4.4	4.5	$18 \pm 3$
Ta3	26	$316 \pm 2$	$319 \pm 2$	12.2	12.3	$18 \pm 11$
Ta4	20	$238.2\pm0.6$	$232 \pm 2$	11.9 <sup>1</sup>	11.6 <sup>1</sup>	$17 \pm 3$
Ta5	20	$573 \pm 1$	$572 \pm 3$	28.7	28.6	$18 \pm 5$
Ta6	20	$718 \pm 2$	$718 \pm 5$	35.9	35.9	$19 \pm 6$

Table 3.2: Thicknesses obtained from  $Ta_2O_5$  samples by XRR and SE using 20.0 Å as thickness in the SE data analysis.

<sup>1</sup> The flow ratio of  $Ar/O_2$  used to grow this sample was 14/3 instead of 14/1 and this may be the explanation for a smaller growth ration.

To have a better view of the differences between the thicknesses obtained by SE and XRR, a representation of the thicknesses obtained by both techniques is presented in Figure 3.2. A linear fit was performed in order to compare these differences with the results reported in the literature (witout the thickness of Ta1 sample). The result of this fit revealed to be similar to the ones described in the literature [11, 44–46]: there is a clear

linear tendency and the slope of the linear fit is near 1. Besides that the interception is equal to 0.



Figure 3.2: Comparison between the thicknesses obtained from XRR and SE.

Besides the thickness of the  $Ta_2O_5$ , the thickness of the native oxide was also fitted in the XRR data analysis. The results (see Table 3.2) are very consistent between samples when one takes in consideration the associated uncertainty. Not only that, but the results are also similar to the 20 Å used as the thickness of the native oxide in SE data analysis.

The thicknesses of both the  $Ta_2O_5$  and the native oxide were not the only fitted parameters. The SLD and roughness of the samples were also fitted. Previously, with only the representation of the XRR data of the samples Ta2 and Ta6 (Figure 3.1) one concluded that the density should not be changing with the increase of the rf power in the target.

In Table 3.3 the results of the fits are presented and the same conclusion can be taken: no evidence of a change in the density of the samples was found for different values of rf power. This is true not only for the density, but also for the roughness of the films that were obtained by XRR and AFM (here SE did not give viable results).

Samula	SLD $(10^{-6} \text{ Å}^{-2})$	Ta <sub>2</sub> O roug	ghness (Å)	
Sample	XRR	XRR	AFM	
Ta1	$48 \pm 9$	$5.0 \pm 0.5$	2.8	
Ta2	$49 \pm 3$	$3.4 \pm 0.2$	4.2	
Ta3	$48 \pm 3$	$4 \pm 2$	4.2	
Ta4	$48 \pm 2$	$5.1 \pm 0.5$	3.7	
Ta5	$50 \pm 4$	$5 \pm 2$	3.3	
Ta6	$48 \pm 4$	$5 \pm 2$	3.5	

Table 3.3: The roughness obtained by XRR and AFM and the SLD obtained by XRR using Si and SiO<sub>2</sub> SLD values of  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup> respectively.

#### CHAPTER 3. RESULTS AND DISCUSSION

From the AFM images used to get the roughness of the samples, one can see how smooth and regular the surfaces of the films are. As an example, two of those images are show below. The small dots in the surfaces may be impurities acquired by the exposition of the samples to the air.



Figure 3.3: Images of the surfaces of two samples obtained by AFM.

As previously said, with the XRR analysis one concluded that the SLD of the samples did not change with the variation of the rf power. In other words, the density remained unchanged. The same conclusion can be made with the analysis of the Tauc-Lorentz models obtained to describe the samples. This is shown in Figure 3.4 where the real and imaginary part of the dielectric function ( $\varepsilon_1$  and  $\varepsilon_2$ , respectively) are represented as functions of the photon energy. The models for these two properties are very similar for all samples, except for the Ta1 sample(Table H.1). This may be the reason for the difference in the thickness determined by SE and XRR of this sample.



Figure 3.4: Real and imaginary part of the dielectric function obtained by SE.
For the other samples a change in density is not found, as the maximum of  $\varepsilon_2$  does not clearly change [47]. Besides that, it is possible to notice that in all samples the Ta<sub>2</sub>O<sub>5</sub> is amorphous, because the increase of the  $\varepsilon_2$  with the photon energy is not abrupt [26].

Furthermore, from the parameters of the Tauc–Lorentz model (Table H.1) it is possible to conclude that the increase of the rf power in the  $Ta_2O_5$  target leads to a small decrease of the band gap of the films and a decrease of short-range order (increase in *C*) [17, 47]. Again, no difference in the parameter *A* was found, so the density of the films did not change with different values of the rf power in the  $Ta_2O_5$  target [17, 47], which is consistent with the results of the XRR data analysis.

To summarise, with this set of  $Ta_2O_5$  samples grown by rf sputtering it was possible to conclude that: no difference in the density and roughness of the films was found with the variation of rf power in the  $Ta_2O_5$  target, the increase of the rf power decreased the band gap and short-range order of the films and all samples were amorphous.

# 3.2 TaSiO samples

After the successful implementation of XRR and the use of the procedure presented in Appendix B to characterize  $Ta_2O_5$  films, one aimed to use this technique to characterize more complex samples. So, the second and third sets of samples that were characterized consisted in a mixture of  $Ta_2O_2$  and  $SiO_2$  compound films (TaSiO) for a dielectric study performed in CENIMAT. The films were grown by co-sputtering and different rf powers were used on the  $Ta_2O_2$  target to achieve samples with different compositions.

Two sets of samples were analysed: films used to determine the growth rate of the specific target powers (named gr-TaSiO) and films used for the dielectric layer in TFTs and MIS devices, where the thickness of all layers were set to be approximately 250 nm (named d-TaSiO). In the following table the nomenclature and the details of the deposition process of each sample is presented. The used nomenclature was set based on the rf power and composition of each sample.

Samp	le	Rf power	$(W  cm^{-2})$
growth rate determination	dielectric layer	Ta <sub>2</sub> O <sub>5</sub> target	SiO <sub>2</sub> target
gr-Ta1Si	-	0.30	1.23
gr-Ta2Si	d-Ta2Si	0.62	1.23
gr-Ta3Si	d-Ta3Si	0.93	1.23
gr-Ta4Si	d-Ta4Si	1.23	1.23
gr-Ta5Si	d-Ta5Si	1.54	1.23
gr-Ta6Si	d-Ta6Si	1.85	1.23
gr-Ta4	d-Ta4	1.23	0

Table 3.4: Nomenclature of both gr-TaSiO and d-TaSiO samples.

#### 3.2.1 Growth rate determination (gr-TaSiO samples)

Since these samples were produced to determine the growth rate, the thickness is the most important parameter to be measured. For this set of samples four techniques were used to do so: profilometry, XRR, SE and RBS.

As in the previous section, an important remark about the XRR needs to be done. Since the fit of the SLD of the native oxide and silicon were not possible to perform, the theoretical SLD of Si and SiO<sub>2</sub> used were  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup> respectively (more details in Appendix F). Despite that it is possible to see in Figure 3.5 de quality of the fits. Just with the information found in the figure below, it is possible to predict that the increase of rf power in the Ta<sub>2</sub>O<sub>5</sub> leads to thicker samples because the sample gr-Ta6Si has a bigger distance between two peaks of the Kiessing fringes ( $d_K$ ) and leads to denser samples since the critical angle ( $\alpha_c$ ) is higher in the sample gr-Ta6Si.



Figure 3.5: Measured data and fit of the samples gr-Ta1Si and gr-Ta6Si.

Taking this remark in mind for the XRR data analysis the thickness in length units is a parameter that can be fitted. The same happens in the SE data analysis. For the other techniques the thickness is obtained in a different way. On one hand, by profilometry the thickness is obtained directly with the measure in length units. On the other hand, with RBS the thickness is a fitting parameter, but it is expressed in atoms cm<sup>-2</sup> and to be converted to length units one must to consider the film's density.

The RBS spectra with the obtained fits with the energy barriers indicated for each element is shown in Figure 3.6. It was needed to assume a fixed stoichiometry for the tantalum and silicon oxides in order to extract a relative concentration of each oxide. In other words, the fit was not done for the elements but for the molecules (it was assumed that the both components formed a solid solution). Looking at the figure bellow one can see that the number of counts for the peak of Ta is increasing with the increase of the rf power in the target. This means that the quantity of Ta is also increasing.



Figure 3.6: RBS spectra from all samples with the respective fits in red.

The same result can be seen in Table 3.5 where the results of the fits are presented. The molecular % of  $Ta_2O_5$  is increasing relatively to the molecular % of  $SiO_2$ . The molecular % of Ar remains between 3 and 5 % for almost all samples.

Sample	Thickness (atoms cm <sup>-2</sup> )	Ta <sub>2</sub> O <sub>5</sub> (%)	SiO <sub>2</sub> (%)	Ar (%)
gr-Ta1Si	38.0	25.5	66.4	8.1
gr-Ta2Si	251.0	31.4	65.8	2.8
gr-Ta3Si	374.0	40.8	40.8	3.8
gr-Ta4Si	500.0	67.3	29.1	5.0
gr-Ta5Si	377.5	70.1	25.5	3.3
gr-Ta6Si	376.4	72.0	24.5	3.5
gr-Ta4	390.0	95.2	-	4.8

Table 3.5: Results of the RBS fits.

However, as previously said, since one wants to compare the thickness obtained by this technique with the others it is fundamental to convert the thickness from  $atoms cm^{-2}$  to length units, in this case Å. To do so the procedure in Appendix G was followed. One needs to take in mind that for the conversion it was necessary to use the mass densities of the components 2.2 and 8.2 g cm<sup>-3</sup> for the Ta<sub>2</sub>O<sub>5</sub> and SiO<sub>2</sub>, respectively [48].

With this conversion it is possible to show the thicknesses for all techniques in Table 3.6. The first conclusion is that all techniques are in general agreement. The profilometry revealed to be the less accurate technique to measure these films, as expected. It is also noticeable that for the gr-Ta1Si and gr-Ta2Si samples it was not possible to get a good model to describe the measured SE data and the results for these two samples are not accurate. This will be discussed in more detail below. Furthermore, for the thinner sample (gr-Ta1Si) the fit of the RBS data does not agree with the other techniques. That can be

explained by the difficult fitting of the measured data of films with such a low thickness.

C 1	Film thickness (Å)					
Sample	profilometry	XRR	SE	RBS		
gr-Ta1Si	$(24 \pm 3) \times 10$	$212 \pm 4$	$0.20 \pm 0.05$	55 ± 3		
gr-Ta2Si	$(32 \pm 4) \times 10$	$330.8\pm0.5$	$266 \pm 4$	$358 \pm 18$		
gr-Ta3Si	$(57 \pm 3) \times 10$	$497 \pm 4$	$507 \pm 3$	$511 \pm 26$		
gr-Ta4Si	$(73 \pm 5) \times 10$	$672 \pm 5$	$678 \pm 5$	$670 \pm 34$		
gr-Ta5Si	$(52 \pm 4) \times 10$	$451 \pm 4$	$462 \pm 2$	$503 \pm 26$		
gr-Ta6Si	$(46 \pm 3) \times 10$	$445 \pm 4$	$458 \pm 2$	$501 \pm 26$		
gr-Ta4	-	$505 \pm 2$	$507 \pm 4$	-		

Table 3.6: Thicknesses obtained by profilometry, XRR, SE and RBS.

Despite the results being similar, the thickness from RBS are slightly different in some cases relatively to the thicknesses obtained by XRR and SE. Three main reasons can be pointed out. The first one is that the fits were from the molecules and not the elements and the stoichiometry can be different from the stoichiometry assumed. The second one is that the mass density assumed may not be the right one. Finally, the fits are not perfect and one sign of that is that the molecular % of Ar is not consistently increasing with the increase of the rf power. Nevertheless, the difference between techniques is very small.

With the thicknesses determined it is possible to calculate the growth rates as one can see in Table 3.7. At the time the d-TaSiO samples were fabricated, only the thicknesses obtained by profilometry were known. This does not represent a problem because the growth rates calculated based on the four techniques are very close to each other.

C 1 .	Deposition	Growth rate (Å/min)			
Sample	time (min)	profilometry	XRR	SE	RBS
gr-Ta1Si	30	8.0	7.1	-	1.8
gr-Ta2Si	30	10.7	11.0	-	11.9
gr-Ta3Si	30	19.0	16.6	16.9	17.0
gr-Ta4Si	30	24.3	22.4	22.6	22.3
gr-Ta5Si	15	34.7	30.1	30.8	33.5
gr-Ta6Si	12	38.3	37.1	38.2	41.8
gr-Ta4	30	-	16.8	16.9	-

Table 3.7: Calculated growth rates based on the obtained thicknesses.

Although the main objective has already been achieved, more analyses can be done with these samples. One relation that can be study is the relation between the thicknesses obtained by XRR and SE data analysis presented in Figure 3.7. The result is similar to the results reported in literature [11, 44–46]: there is a clear linear tendency, the slope of the

linear fit is near 1 and the interception is different than 0. Interceptions different than 0 have been reported in different studies but this fact remains without a good explanation [45].



Figure 3.7: Comparison between thicknesses obtained from XRR and SE.

Continuing the comparison between XRR and SE results, an agreement is not found in the roughness of the films (Table 3.8). The roughnesses obtained by XRR seems to be very similar for all samples. The same does not happen with the SE data analysis. Not only that, but by SE it was not possible to determine the roughnesses of some films because those values tended to 0.1 Å, which is the limit of the used software, so those values were meaningless. In those cases, no roughness layer was used in the data analysis. Here XRR revealed to be a more useful technique to measure the film roughness.

Sample	Fi. roughr	lm 1ess (Å)	SLD	thic	SiO <sub>2</sub> kness (Å)	SiO <sub>2</sub> roughness (Å)
	XRR	SE	XRR	XRR	SE	XRR
gr-Ta1Si	$4.1 \pm 0.3$	-	$20.0 \pm 0.9$	$12 \pm 3$	$249.9 \pm 0.4$	13 ± 2
gr-Ta2Si	$4.4\pm0.6$	-	$30 \pm 2$	$27 \pm 3$	$70 \pm 3$	$15.8 \pm 0.9$
gr-Ta3Si	$4 \pm 1$	-	$37 \pm 2$	$25 \pm 4$	$35 \pm 2$	$13 \pm 4$
gr-Ta4Si	$4.9\pm0.9$	$6 \pm 2$	$39 \pm 2$	$23 \pm 5$	$22 \pm 4$	$16 \pm 8$
gr-Ta5Si	$3 \pm 2$	$2.9\pm0.7$	$41 \pm 2$	$21 \pm 5$	$30 \pm 2$	$12 \pm 6$
gr-Ta6Si	6 ± 2	$3.9 \pm 0.8$	$41 \pm 2$	$21 \pm 6$	$32 \pm 2$	$14 \pm 6$
gr-Ta4	$4.0\pm0.9$	$7 \pm 1$	$50 \pm 2$	$21 \pm 7$	$62 \pm 3$	-

Table 3.8: Results from XRR and SE data analysis of the Ta<sub>2</sub>O<sub>5</sub> and of the native oxide using Si and SiO<sub>2</sub> SLD values of  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup> respectively.

For the determination of the native oxide thickness, XRR revealed again to be a more

accurate technique. It is possible to see that the thicknesses of the native oxide were not in agreement for all samples for both techniques. In the case of the XRR, the results are close to the expected and, given the uncertainty, the results are close to each other. However, the results of the native oxide thickness obtained by the SE data analysis are not so reliable because of the incoherence between samples and in some cases, like in the sample gr-Ta4, is not even close to the expected values.

This incoherence highlights the importance of analysing a piece of the substrate in the first place. This can give the opportunity to reduce the number of the fitting parameters in the data analysis, like the native oxide thickness and roughness.

For the last conclusion of the XRR results, given the fact that  $Ta_2O_5$  is a denser material than SiO<sub>2</sub> [48] it was predictable that, with the increase of  $Ta_2O_5$  in the film, the film becomes denser. This is proven correct by the fact that both the SLD and the molecular % of  $Ta_2O_5$  are increasing as the rf power in the  $Ta_2O_5$  target increases (Table 3.5 and 3.8).

Previously in Tables 3.6 and 3.8 was evident that the models for the gr-Ta1Si and gr-Ta2Si do not give physically reliable results. The problem with the models of gr-Ta1Si and gr-Ta2Si is shown in Figure 3.8 and in Table H.2. There is a huge discrepancy between this two models and the other models. Despite the attempts, better models were not achieved.



Figure 3.8: Real and imaginary part of the dielectric function of the samples.

With the representation of  $\varepsilon_1$  and  $\varepsilon_2$  it is possible to see that with higher concentration of Ta<sub>2</sub>O<sub>5</sub> the peaks become more intense. This is also perceptible in Table H.2 where the parameter A1 increases. This means that the films become denser [17, 47]. The parameter C1 also increases with the concentration of Ta<sub>2</sub>O<sub>5</sub> which reveals that the films have a decrease of short-range order with the increase of the rf power in the Ta<sub>2</sub>O<sub>5</sub> target [17, 47] which contradicts the reported results [24]. However, the decrease of  $E_g$  with the increase of the Ta<sub>2</sub>O<sub>5</sub> concentration is in agreement with the literature [24]. It is also possible to conclude that all Ta<sub>2</sub>O<sub>5</sub> films are amorphous, because the increase of the  $\varepsilon_2$ with the photon energy is not abrupt [26]. In summary, with this set of samples it was possible to conclude that the density,  $Ta_2O_5$  composition and growth rate increase with the rf power in the  $Ta_2O_5$  target at the same time as the short-range order decreases. Once again, the success of the implementation of XRR was demonstrated.

#### 3.2.2 Dielectric layer (d-TaSiO samples)

After the characterization of the gr-TaSiO samples to calculate the growth rate, the characterization of the d-TaSiO samples will be presented. The first part of this characterization is the analysis of the thickness, roughness and composition of the film and of the native oxide. The second part to be shown is the electric characterization of the films as a dielectric layer in a MIS structure.

In this section the thicknesses were obtained by two techniques: SE and RBS. The results are presented in Table 3.9. The deposition time used to fabricate these samples was based on the growth rate calculated using the thicknesses obtained by profilometry (Table 3.7). With the exception of the sample d-Ta3Si, all samples have a thickness close to the desired 250 nm. Besides that, as have been noticed, the results for both techniques are very close to each other. This agreement can lead us to two other conclusions.

Sample	Film tl	Film thickness $(atoms  cm^{-2})$	
	SE	RBS	RBS
d-Ta2Si	$2395 \pm 14$	$(236 \pm 12) \times 10$	1668
d-Ta3Si	$2049\pm7$	$(204 \pm 11) \times 10$	1495
d-Ta4Si	$2425\pm8$	$(241 \pm 13) \times 10$	1796
d-Ta5Si	$2311 \pm 10$	$(228 \pm 12) \times 10$	1712
d-Ta6Si	$2550 \pm 11$	$(256 \pm 13) \times 10$	1923
d-Ta4	$2520\pm13$	$(237 \pm 12) \times 10$	1855

Table 3.9: Thicknesses obtained by SE and RBS.

First it is important to take in mind that in order to calculate the thicknesses from the RBS results it was assumed the mass density of the  $SiO_2$  and  $Ta_2O_5$  of 2.2 and  $8.2 \text{ g cm}^{-3}$ , respectively. Not only that, but it is fundamental to understand that the fits were done assuming that the components of the films were  $Ta_2O_5$ ,  $SiO_2$  and Ar (although Ar is not used to calculate the thicknesses). This means that the results presented are the results of the fit with the molecules and not the fit with the elements (Appendix G).

Therefore, if the results of both techniques are so similar with these assumptions, it means that the assumed stoichiometry may be close to the real stoichiometry of the components of the films and that the components have a density close to the bulk densities.

To calculate the thickness presented in Table 3.9 obtained by RBS, the thickness in  $atoms cm^{-2}$  and the composition of the films were used. The composition obtained by

RBS and SE is shown in Table 3.10. The results are the expected ones: with the increase
of the rf power in the $Ta_2O_5$ target the concentrations of $Ta_2O_5$ and Ar increase.

0 1	RBS			Normalized RBS	SE
Sample	Ta <sub>2</sub> O <sub>5</sub> (%)	SiO <sub>2</sub> (%)	Ar (%)	Ta <sub>2</sub> O <sub>5</sub> (%)	Ta <sub>2</sub> O <sub>5</sub> (%)
d-Ta2Si	37.5	59.1	3.4	38.8	40.6
d-Ta3Si	57.6	38.6	3.8	59.9	56.4
d-Ta4Si	65.7	30.0	4.3	68.7	60.5
d-Ta5Si	70.6	25.0	4.5	73.8	61.5
d-Ta6Si	71.7	24.0	4.3	74.9	66.3
d-Ta4	94.5	-	5.5	100	-

Table 3.10: Composition in molecular % obtained by RBS and SE data analysis.

The results were the expected because on one hand in the RBS data (Figure 3.9) the peak of the Ta also increases with the increase of the rf power. This means that the quantity of Ta is also increasing. On the other hand, the trapping of rare gases in sputtering is proportional to the atomic mass of the targets. Since  $Ta_2O_5$  has a higher atomic mass than SiO<sub>2</sub>, with the increase of the molecular % of  $Ta_2O_5$ , more Ar atoms are trapped in the deposited film [49].



Figure 3.9: RBS spectra from all samples with the respective fits in red.

Also in Table 3.10 it is compared the molecular % of the  $Ta_2O_5$  obtained by RBS and SE. In this comparison it was used a normalised concentration in the RBS data analysis. The normalised concentration only takes into account the concentration of  $Ta_2O_5$  and SiO<sub>2</sub>. This was used because not only it is not possible to calculate the thickness with films that have gases in its constituents, but also because in the SE data analysis presented here, Ar were not considered as well.

The molecular % of the  $Ta_2O_5$  presented for the SE data analysis was calculated too. Using a different model (Figure E.3), the volumetric concentration of the  $Ta_2O_5$  was fitted. The mass densities were assumed to calculate the molecular %. Knowing that RBS gives more viable results, one can conclude that it is possible to obtain a concentration with SE, but the result can only be used to give an idea of the concentrations and cannot be assumed as a certain result. This is the case because in these analyses the concentration of  $Ta_2O_5$  and the Tauc-Lorentz model of  $Ta_2O_5$  were fitted at the same time.

Besides the thickness and composition, SE give more information, namely the thickness of the silicon oxide and the roughness of the  $Ta_2O_5$  films. However, as one can see in the table below (Table 3.11), the models obtained by SE revealed to be not very accurate for these two parameters. On one hand, the values of the native oxide thickness are not similar between samples. On the other hand, the roughness could not be obtained for two samples and is not consistent in the others. Nevertheless, it seems that the thickness of the  $Ta_2O_5$  was not affected by these problems. Nevertheless, the parameters of the Tauc-Lorentz models can give more information (Figure 3.10 and Table H.3)

Table 3.11: Thicknesses of the native oxide and roughnesses of the Ta<sub>2</sub>O<sub>5</sub>.

Sample	SiO <sub>2</sub> thickness (Å)	Ta <sub>2</sub> O <sub>5</sub> roughness (Å)
d-Ta2Si	$14 \pm 8$	NS
d-Ta3Si	$22 \pm 3$	NS
d-Ta4Si	$9\pm3$	$5 \pm 2$
d-Ta5Si	$16 \pm 4$	$8 \pm 2$
d-Ta6Si	$18 \pm 4$	$4 \pm 2$
d-Ta4	21 ± 5	$11 \pm 2$



Figure 3.10: Real and imaginary part of the dielectric function of the samples.

To conclude the first part of the characterization of the d-TaSiO samples one may look to the results of the Tauc-Lorentz models. Those results are in line with the results of the gr-TaSiO samples: the films become denser and have a decrease of short-range order with the increase of the rf power in the  $Ta_2O_5$  target and all films are amorphous. This is shown in both Figure 3.10 and Table H.3.

The second part of this characterization consisted in the electrical characterization of MIS structures using the d-TaSiO films as dielectric layers. Therefore, the capacitance was measured as a function of the frequency (capacitance characteristic) and voltage.

The objective of the study of these films was to use them as the dielectric layer in transistors. So, the possibility of measuring these films in high frequency was desired. However, as one can see in Figure 3.11 for frequencies higher than 100 kHz the capacitance starts to increase. This effect is due to the high impedance of the cables. Because of that, the remain electrical tests were performed using a frequency of 100 kHz.



Figure 3.11: Electrical measurements of the d-TaSiO samples.

To know the dielectric constant of the films the capacitance was measured in function of the applied potential, as shown in the figure above. Knowing that the contacts' size was  $494 \,\mu\text{m} \times 494 \,\mu\text{m}$ , using the thicknesses obtained by SE (Table 3.9) and with the maximum capacitance of the films (Figure 3.11), the dielectric constant was calculated.

To get a better understanding of the calculated dielectric constants, the results are presented in Figure 3.12. As it is possible to see there, the dielectric function increases with the increase of molecular % of  $Ta_2O_5$  in the films (caused by the increase of the rf power in the  $Ta_2O_5$  target). The results are in agreement with the expected ones. The dielectric constants of SiO<sub>2</sub> and amorphous  $Ta_2O_5$  are 3.9 [50] and 22-25 [51, 52], respectively. So, with the increase of the molecular % of  $Ta_2O_5$ , the dielectric constant of the films gets closer to the dielectric constant of the amorphous  $Ta_2O_5$ .



Figure 3.12: Dielectric constants obtained from the electrical measurements.

In conclusion, with these sets of samples of films with different mixtures of  $Ta_2O_5$  and  $SiO_2$  was possible to see that the films become denser and have a decrease of short-range order with the increase of the rf power in the  $Ta_2O_5$ . Is was also possible to conclude that with the increase of the rf power the molecular % of  $Ta_2O_5$  also increased which lead to an increase of the dielectric constant.

# 3.3 ZTO samples

The final set of samples characterized was a set of ZTO films grown by rf sputtering. These samples were fabricated at the same time as the samples used by C. Fernandes *et al.* as the dielectric layer in TFTs for flexible electronics [34]. Here the objective was to understand the role of some fabrication parameters, namely the rf power in the ZTO target and the flow of  $O_2$  and  $H_2$ . The nomenclature used is presented in Table 3.12.

Sample	Rf power (W cm <sup>-2</sup> )	H <sub>2</sub> flow (sccm)	O <sub>2</sub> flow (sccm)
P1-H1-O2	1	0	5.36
P1-H2-O2	1	0.5	5.36
P1-H3-O2	1	1	5.36
P2-H1-O1	2	0	1.78
P2-H1-O2	2	0	5.36
P2-H2-O2	2	0.5	5.36
P2-H3-O2	2	1	5.36
P2-H1-O3	2	0	10.72
P2-H2-O3	2	0.5	10.72
P2-H3-O3	2	1	10.72

Table 3.12: Samples nomenclature used in function of the fabrication parameters.

### CHAPTER 3. RESULTS AND DISCUSSION

In line with the other sections of this work, the XRR data was analysed using the procedure presented in Appendix B and once again the fit of the SLD of the native oxide and of the silicon were not viable. To solve this problem the assumed values of the SLD used for the Si and SiO<sub>2</sub> were  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup>, respectively. These values are the theoretical ones. A more detailed explanation of these problem can be found in Appendix F.

Nevertheless, and as in the previous sections, very good fits of the measures XRR data was obtained. In Figure 3.13 the real data and the fit of a sample is presented. This is a representative example because for the remain samples the quality of the fits was as good as the fit of this sample.



Figure 3.13: Real data and XRR fit of the sample P1-O2-H1.

To study the correlation of the thicknesses determined by both techniques, a comparison was made and it is presented in Figure 3.14. The result obtained with this comparison is similar to other results presented in this work as well as to the results reported in the literature [11, 44–46]. In other words, the result of the linear fit was a slope close to 1 and the interception is different than 0. As in the previous section, an explanation for the interception being different than 0 was not found.



Figure 3.14: Comparison between thicknesses obtained from XRR and SE.

## 3.3.1 Influence of the rf power

The first parameter studied was the rf power in the ZTO target. Two different rf powers were tested:  $1 \text{ W cm}^{-2}$  (P1) and  $2 \text{ W cm}^{-2}$  (P2). Besides the change in the rf powers, the deposition time was also changed. In the samples where  $1 \text{ W cm}^{-2}$  was used, the deposition time was 20 min. The deposition time of the samples where  $2 \text{ W cm}^{-2}$  was used was 5 min. This means that using twice the rf power, the time used was a quarter. With this, half of the thickness was expected for the P2-samples compared to the P1-samples and that was the result, as one can see in Table 3.13. Besides that, it is noticeable that both techniques are in very good agreement.

Sample	Film thickness (Å)		Film SLD $(10^{-6} \text{ Å}^{-2})$	Film roughness (Å)
	SE	XRR	XRR	XRR
P1-O2-H1 P2-O2-H1	299 ± 2 157 ± 2	$305 \pm 5$ 161.1 ± 0.6	$\begin{array}{c} 44\pm3\\ 47\pm3\end{array}$	$8 \pm 2$ $6.4 \pm 0.4$
P1-O2-H2 P2-O2-H2	$311 \pm 11$ $172 \pm 2$	$330 \pm 3$ 161.5 ± 0.4	$\begin{array}{c} 45\pm3\\ 47\pm3\end{array}$	$\begin{array}{c} 8\pm2\\ 6.0\pm0.4\end{array}$
P1-O2-H3 P2-O2-H3	$306 \pm 2$ $165 \pm 2$	$309 \pm 3$ 156.0 ± 0.5	$\begin{array}{c} 44\pm3\\ 49\pm3\end{array}$	$9 \pm 2$ $6.4 \pm 0.4$

Table 3.13: Results from XRR and SE data analyses of the samples with different target rf power using Si and SiO<sub>2</sub> SLD values of  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup>, respectively.

The differences between P1- and P2-samples are not just in the thicknesses. P2samples also have a lower roughness, but this could be due to the fact that these samples are thinner. To test this hypothesis, samples with similar thicknesses should be fabricated using different rf powers. Besides the roughness, the SLD also changes. The SLD increases with the increase of the rf power. Although the change is slight for the given uncertainties, it is possible to conclude that the density of the films increases with the increase of the rf power (supported by the increase of A1 with the rf power shown in Table H.4).

The increase of the density is also confirmed by SE because as shown in Table H.4 the parameter A1 is also higher for P2-samples [17, 47]. Looking at the Tauc-Lorentz parameters for this sample is also possible to see an increase of the parameters *E*1 and C1. So, despite being denser, the sample also has a lower short-range order [17, 47]. With SE one can also conclude that the samples are amorphous, because as shown in Figure I.1 the  $\varepsilon_2$  does not present an abrupt change with the photon energy [26].

#### **3.3.2** Influence of the H<sub>2</sub> flow

The second parameter studied was the  $H_2$  flow during the rf sputtering. To test the influence of this parameter, samples with three different flows were prepared. Some samples were fabricated without  $H_2$  flow (H1), some with a flow of 0.5 sccm (H2) and others with a flow of 1 sccm (H3).

Starting with the thickness, one can see that the thicknesses are higher in the H2samples. This may indicate that the use of  $H_2$  flow in the fabrication of the samples can lead to higher thicknesses. Nevertheless, a critical point of  $H_2$  flow may exist as the thicknesses are lower in the H3-samples than the H2-samples. More tests with more values of  $H_2$  flow should be performed to obtain better conclusions. Besides that, it seems that the SLD increases with the  $H_2$  flow but when one takes in consideration the associated uncertainty, that conclusion is not very clear.

Sample	F thick	Film ness (Å)	Film SLD $(10^{-6} \text{ Å}^{-2})$	Film roughness (Å)
	SE	XRR	XRR	XRR
P1-O2-H1	299 ± 2	$305 \pm 5$	$44 \pm 3$	8 ± 2
P1-O2-H2	$311 \pm 11$	$330 \pm 3$	$45 \pm 3$	8 ± 2
Р1-О2-Н3	$306 \pm 2$	$309 \pm 3$	$44 \pm 3$	9 ± 2
P2-O2-H1	157 ± 2	$161.1 \pm 0.6$	$47 \pm 3$	$6.4 \pm 0.4$
P2-O2-H2	$172 \pm 2$	$161.5\pm0.4$	$47 \pm 3$	$6.0 \pm 0.4$
Р2-О2-Н3	$165 \pm 2$	$156.0\pm0.5$	$49 \pm 3$	$6.4 \pm 0.4$
P2-O3-H1	$165 \pm 2$	$152.2 \pm 0.9$	$43 \pm 3$	$6.3 \pm 0.9$
P2-O3-H2	$152 \pm 1$	$150 \pm 1$	$44 \pm 3$	$6 \pm 1$
Р2-О3-Н3	$143 \pm 2$	$139.5\pm0.8$	$46 \pm 4$	$6.0 \pm 0.7$

Table 3.14: Results from XRR and SE data analyses of the samples with different H<sub>2</sub> flow using Si and SiO<sub>2</sub> SLD values of  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup>, respectively.

Through the Tauc-Lorentz parameters (see Table H.5) no conclusions can be taken because no pattern was found in any parameter. The same happens looking at the representation of both parts of the complex dielectric function in Figure I.2, in which one can only conclude that all samples are amorphous [26].

## **3.3.3** Influence of the O<sub>2</sub> flow

Finally, the last parameter to be study was the influence of the  $O_2$  flow during the fabrication process. Three different flows were tested to understand the role of the  $O_2$  flow in the parameters of the samples: 1.78 (O1), 5.36 (O2) and 10.72 sccm (O3).

The results of the characterization of these samples can be seen in Table 3.15. A difference between the thicknesses was found. This difference may indicate that the increase of the  $O_2$  flow results in a decrease of the thickness. The sample P2-O2-H1 is an exception for this hypothesis. To make a clear conclusion, more samples are needed to test this hypothesis. Besides the decrease of the thickness, the roughness and the SLD also present lower values for a higher  $O_2$  flow. So, the increase of  $O_2$  flow may lead to less dense and smoother films.

Table 3.15: Results from XRR and SE data analyses of the samples with different O<sub>2</sub> flow using Si and SiO<sub>2</sub> SLD values of  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6}$  Å<sup>-2</sup>, respectively.

Sample	Film thickness (Å)		Film SLD $(10^{-6} \text{ Å}^{-2})$	Film roughness (Å)
	SE	XRR	XRR	XRR
P2-O1-H1 P2-O2-H1 P2-O3-H1	$176 \pm 2$ $157 \pm 2$ $165 \pm 2$	$175 \pm 2$ $161.1 \pm 0.6$ $152.2 \pm 0.9$	$47 \pm 3$ $47 \pm 3$ $43 \pm 3$	$6.7 \pm 0.9$ $6.4 \pm 0.4$ $6.3 \pm 0.9$
Р2-О2-Н2 Р2-О3-Н2	$172 \pm 2$ $152 \pm 1$	$161.5 \pm 0.4$ $150 \pm 1$	$\begin{array}{c} 47\pm3\\ 44\pm3\end{array}$	$6.0 \pm 0.4$ $6 \pm 1$
Р2-О2-Н3 Р2-О3-Н3	$\begin{array}{c} 165 \pm 2 \\ 143 \pm 2 \end{array}$	$156.0 \pm 0.5$ $139.5 \pm 0.8$	$\begin{array}{c} 49\pm3\\ 46\pm4\end{array}$	$6.4 \pm 0.4$ $6.0 \pm 0.7$

Looking at the Tauc-Lorentz parameters in Table H.6 and the dielectric function in Figure I.3 no tendency can be found. As a result of that, one can only conclude that all samples were amorphous [26].

## Conclusions and future perspectives

In this work the implementation of XRR at CENIMAT as a characterization technique of thin and ultra-thin films was the main objective. Through the use of the package *MOTOFIT 4.1* to analyse the measured data obtained from the diffractometer present at CENIMAT the objective was successfully achieved. Now it is possible to use this technique to measure the thickness, roughness and the SLD of thin films and ultra-thin films at CENIMAT.

The use of the SE as a complementary technique was also one of the objectives of this work. Using the ellipsometer of the CENIMAT and using Tauc-Lorentz dispersion formula this objective was also successfully completed. A very good agreement in the thicknesses measured by both techniques was found. In all groups of samples analysed in this work, a linear relation between the thicknesses obtained by XRR and SE data analyses was found. The slope of this relation for all set of samples was close to 1 and the interception was different than 0. In none of the group of samples analysed a good explanation for the slope being different than 0 was found.

RBS and profilometer analyses were also performed. An agreement was also found with the thicknesses obtained by these techniques. Besides the agreement between these four techniques, RBS and profilometer present some disadvantages. In order to calculate the thickness of a film with the RBS data the composition of the film must be well known, which does not always happen. The profilometer can give a direct measure of the thickness but it is less accurate and it is not suitable to measure ultra-thin films. So, XRR and SE can be more appropriate when RBS and profilometer cannot be carried out.

The use of XRR and SE as complementary techniques lead to very good results, being possible to say that these techniques work very well as complementary techniques. On one hand, throughout this work XRR was shown to be a better technique to determine the roughness of the films and to determine the thickness of the native oxide. On the other hand, SE can give more details about the films. Using the Tauc-Lorentz dispersion formula it was possible to compare the short-range order between films, to obtain the band gap of the films and to determine whether the films were amorphous or not.

All these conclusions resulted from the analysis of four groups of samples. The first group of those samples was  $Ta_2O_5$  films. The second group of samples was prepared to know the growth rate of films with  $Ta_2O_5$  and  $SiO_2$  (TaSiO) in order to produce the third group of samples, in which 250 nm of thickness was desired. The final group of samples was ZTO films. All samples were produced by rf sputtering.

In the case of the  $Ta_2O_5$  one aimed to study the influence of the rf power in the  $Ta_2O_5$  target. One concluded that the films were amorphous and that the increase of the rf power in the  $Ta_2O_5$  target lead to a decrease of the short-range order. No differences in the roughness or density were found.

From the TaSiO it was also studied the influence of the rf power in the Ta<sub>2</sub>O<sub>5</sub> target.

The results showed that the films got a higher concentration of  $Ta_2O_5$  and Ar, as predictable. The films, that revealed to be amorphous, got denser and the dielectric constant also increased with the increase of the rf power. On the other hand the short-range order decreased.

Samples with ZTO were produced in order to study the influence of the rf power in the ZTO target, the  $H_2$  flow and the  $O_2$  flow during the fabrication process. It can be concluded that all samples were amorphous according to the SE results.

The ZTO samples produced with higher rf power had a lower deposition time, so they were thinner. These samples revealed to have a lower roughness, but it was not clear if it was due to the use of higher rf power or due to the fact that these samples are thinner. Samples with similar thickness should be prepared to have a better understanding of the role of the rf power in the roughness of the films. It was also noticeable that the increase of the rf power lead to denser films and with lower short-range order.

The study of the use of  $H_2$  flow indicates that it may have a role in the thickness of the film. For the intermediate value of flow the thickness were higher than for the lower and higher values of tested flows. This can mean that it may exist a critical point where the thickness has a maximum value. To be sure more samples with more values of  $H_2$  flows should be characterized.

An increase of the  $O_2$  flow lead to films less dense and with higher roughness. It seems that also lead to thinner films, but this conclusion is not clear. Here more tests are also desired to have a better understanding of the influence of the  $O_2$  flow.

In the future, besides all the tests suggested above, other methods to analyse the XRR could be tested. The package *MOTOFIT 4.1* presents many advantages, but it showed some limitation for the analyses of thinner samples (below 30 Å). Besides that, *MOTOFIT 4.1* works within *IGOR*, which is a paid program. Once there are free alternatives, those alternatives must be studied. In this way, any CENIMAT's worker and student could analyse XRR data in their personal computer.

To test the use of the current analysis method more diverse samples should be analysed. Not only single layer but also multilayer systems. Samples with subtracts different from silicon and fabricated with different processes should also be characterized.

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# Abelès matrix formalism

In order to extract the properties of the films out of the data obtained by XRR, the analyses were performed in the *MOTOFIT 4.1* package [4]. This package works within *IGOR* (Wavematrics) and uses the Abelès matrix formalism. In this formalism, the reflectivity is calculated as a function of the momentum transfer (Q):

$$Q = \frac{4\pi}{\lambda} sin(\theta) = 2k_z \tag{A.1}$$

where  $\lambda$  is the wavelength of the incident radiation,  $\theta$  is the angle of the incident radiation and  $k_z$  the wave factor. The SLD profile is what makes the measured reflectivity vary and it is normally changing as a depth function. Despite that, a good approximation can be done if one considers the samples as a group of layers with determined thickness, roughness and SLD (which is a function that depends on the mass density and the composition of the layers).

In the Abelès matrix formalism the obtained reflectivity is calculated from the multiplication of the characteristic matrix of each layer. To calculate this matrix, three variables are needed: the wave factor (that is a function of the SLD), the layer thickness and the Fresnel reflection coefficient. The Fresnel reflection coefficient between two layers is given by:

$$r_{n,n+1} = \frac{k_n - k_{n+1}}{k_n + k_{n+1}} \tag{A.2}$$

This does not takes into account the roughness/diffuseness of the layer. So, it is used a Fresnel reflection coefficient with an error function [53]:

$$r_{n,n+1} = \frac{k_n - k_{n+1}}{k_n + k_{n+1}} (-2k_n^2 + \sigma_{n,n+1}^2)$$
(A.3)

where  $\sigma_{n,n+1}$  is the roughness/diffuseness [4].

# Tutorial of XRR data analysis in MOTOFIT

This tutorial assumes that the data analysis will be performed in the computer used at CENIMAT to analyse XRR and X-ray diffraction data. On that computer *IGOR*, *MOTOFIT* and *Python* with the respective packages have already been installed. This tutorial is based in the experience acquired during the present work and in the *MOTOFIT* manual that can be found at http://motofit.sourceforge.net/manual/motofit/motofitmanual.pdf that was written by Andrew Nelson. In the manual you can find more detailed information about *MOTOFIT*.

In the beginning you must have a *.csv* file generated by the PANalytical's X'Pert PRO MRD X-ray diffractometer in the CENIMAT. These files contain the information of the characterization as well as the measured data (detected intensity as a function of the angle of incidence). In order to use the *MOTOFIT* package the data must be converted to obtained the normalised intensity as a function of the photon momentum.

To convert the data you must use the *Python* function *data\_converter\_Motofit.py*. After opening the function a file selection dialog box appears. There you must select all files you want to convert. After the conversion new files will appear in the same folder in which the original files are stored. If the original file name is *example.csv*, the converted file name will be *example\_motofit.csv*.

In order to open the *MOTOFIT* package, you have to open *IGOR* and go to *Motofit* > *Fit Reflectivity data* > *Continue*. In the *Reflectivity panel* it is possible to find four menu tabs and in the *Fit* tab it is possible to find four menu boxes (Figure B.1).

Reflectivity Panel
Fit Constraints thickness estimation plot control
Joad data datasetnone V Save fits
Plotting       type     logR vs Q       vg     resolution dq/q % 5       type     Add cursor
Pring           Do fit           Genetic
$\begin{array}{c c c c c c c c c c c c c c c c c c c $
layer         thickness         SLD         solvent         roughness           fronting         INF         0         1         1         25         3.47         0         4         1           backing         INF         2.07         4         1

Figure B.1: Reflectivity panel of *MOTOFIT* with the menu taps in the red boxes and the menu boxes in the green boxes.

In the *Dataset* menu box, by clicking in *load data* you can choose the file to be analysed. Beware of two points: choose the *\_motofit.csv* files (the converted files) and confirm that you are working in the right file. After loading one file, if you load a second one the program will continue to work in the first file by default. You can confirm and choose the right file to work in *data set*.

In the *Model* menu box you can load a file with the coefficients if you already have it. If not you can define all the parameters. All the parameters that you are certain about you should set them as fixed. An example is presented in Figure B.2. In this case it is a sample of Ta<sub>2</sub>O<sub>5</sub> film in a silicon wafer. For a silicon wafer the *backing* has a fixed value for the SLD as  $20.124 \times 10^{-6}$  Å<sup>-2</sup>. The *layer 2* is the native silicon oxide, so the SLD value was fixed in  $18.886 \times 10^{-6}$  Å<sup>-2</sup>. The third layer is the Ta<sub>2</sub>O<sub>5</sub>. These SLD values were obtained using the *SLD calculator* in the *Motofit* tab.



Figure B.2: An example of a model with fixed parameters.

In the *Fitting* menu box you can perform the fit by clicking in *Do fit*. A first fit should be done with the *Genetic* option selected. This is the best way to have a better idea of the parameters. Then the *Gencurvefit limits* windows will appear. Here you must define the limits for the parameters' fitting. The number of iterations and the population size of the fit can be higher if necessary. After clicking in *Do fit* the results will be shown. It may be necessary to adjust the *resolution dq/q* % in the *Plotting* menu box.

Once you already have a good idea of the parameters you should perform a fitting with the *Genetic* + *LM* option selected. This performs a genetic optimization and uses the Levenberg-Maquardt algorithm to have more accurate results and the respective uncertainties. You must be aware that if the defined limits of the parameters' fitting are too vague, the fit may not present viable results. This is the reason why the genetic fitting must be done in the first place. The final results will be the SLD profile in the *Scattering lengh density* window (Figure B.3) and the parameters with the respective uncertainty in the window below (Figure B.4). To export the fitted graphic, you can use the *Save fits* in the *Dataset* box menu. You must choose the option with the same name as the file.



Figure B.3: Results of the SLD profile of the sample.



Figure B.4: Results of the parameters of the sample with the respective uncertainties.

If you are sure that some parameters are correlated, you can use that information in the *Constraints* tab menu. As an example, if that information is «the SLD of the layer 1 is equal do the SLD of layer 2», the input will be K7=K15. The explanation of the number of the parameters is shown in Figure B.5.



Figure B.5: Results of the parameters of the sample with the respective uncertainties.

The *plot control* tab menu can be used to define which lines can or cannot be shown. By clicking twice in the lines of the graphic you can modify them.

One final remark must be done: every time you use *MOTOFIT* to get information to include in a publication you have to give credits to the author citing the paper:

A. Nelson. "Co-refinement of multiple-contrast neutron/X-ray reflectivity data using MOTOFIT." In: *Journal of Applied Crystallography* 39.2 (2006), pp. 273–276.

Sample	XRR	SE
Si wafer	0.0178582	0.373754
Ta1	0.00613824	0.175486
Ta2	0.0150852	4.190142
Ta3	0.00780566	1.683170
Ta4	0.0208817	3.329095
Ta5	0.0284124	3.497609
Ta6	0.0267322	5.451688

Table C.1:  $\chi^2$  of the fits performed in the Ta<sub>2</sub>O<sub>5</sub> samples (section 3.1).

Error functions of the performed fits

Table C.2:  $\chi^2$  of the fits performed in the TaSiO samples (section 3.2).

C	gr-TaSi			d-TaSi		Table 3.10
Sample	XRR	SE	RBS	SE	RBS	SE
Ta1Si	0.0158015	0.195114	60.67	-		-
Ta2Si	0.055409	1.524357	30.9	7.187333		7.353949
Ta3Si	0.017355	0.973171	8.2	4.365729		4.574820
Ta4Si	0.012244	1.868353	7.2	4.143032		4.395866
Ta5Si	0.0233148	1.222181	8.0	5.474410		5.877838
Ta6Si	0.0148769	1.136910	11.9	6.183555		7.297378
Ta4	0.0249664	2.261642	8.525	7.525571		7.513732

Table C.3:  $\chi^2$  of the fits performed in the ZTO samples (section 3.3).

Sample	XRR	SE
P1-O2-H1	0.0364218	1.288003
P1-O2-H2	0.0605636	1.142378
P1-O2-H3	0.0490492	1.177760
P2-O1-H1	0.0428692	0.545678
P2-O2-H1	0.0188471	0.510383
P2-O2-H2	0.0211289	0.732384
Р2-О2-Н3	0.0237625	0.579150
P2-O3-H1	0.0637746	0.567632
P2-O3-H2	0.0546619	0.413325
P2-O3-H3	0.0481622	0.393520

# Python fuction to convert data to be used in MOTOFIT

```
# -*- coding: utf-8 -*-
1
   .....
2
3
   Created on Mon Apr 30 09:16:58 2018
4
   @author: Tiago Gonçalves
5
   .....
6
   #This program converts files from XRR analysis from CENIMAT's diffractometer
7
8
   #into data that can be analysed by Motofit
9
   import numpy as np
10
   from tkinter import filedialog
11
12
   from tkinter import *
13
   #To select the files:
14
   root = Tk()
15
   root.withdraw()
16
   files_path = filedialog.askopenfilenames(filetypes = (("CSV_files","*.csv"),
17
                                                            ("all_files","*.*")))
18
   number_files = len(files_path)
19
   z = 0
20
21
   #Runs through all files
22
   while z < number_files:</pre>
23
        data = np.genfromtxt(files_path[z], delimiter = ',', skip_header = 29)
24
        length = len(data)
25
       max_intensity = np.max(data)
26
27
        i = 0
28
        #Converts the values
       while i < length:</pre>
29
            data[i , 0] = 4 * np.pi * np.sin(np.radians(data[i , 0])) / 1.5405980
30
31
            data[i , 1] /= max_intensity
            i += 1
32
        #Save the new files
33
       np.savetxt(files_path[z][:-4] + '_motofit.csv', data,
34
                   fmt = '%.9f,%.25f')
35
        z += 1
36
```

# **Complementary schemes**



Figure E.1: Model used to perform the fit of the data obtained by SE.



Figure E.2: Scheme of the MIS structure.



Figure E.3: Model used to perform the fit of the data obtained by SE to get the composition of the films.
### Analysis of a silicon wafer

For the characterization of the Ta<sub>2</sub>O<sub>5</sub> films a piece of the silicon wafer where all films were grown was previously characterized. Here the objective was to lower the number of the fitting parameters of the XRR and the SE data analysis. More specifically one aimed to get the thickness and roughness of the native oxide and the SLD of both the native oxide and the Si, so these parameters could be used as fixed values for all samples.

Despite the attempts, one had to use the theoretical values of the SLD to get reliable fits in the XRR data analysis. Those values were obtained in the SLD calculator of MOTOFIT. The reason for this problem is the similarity of Si and SiO<sub>2</sub> electron densities [45].

K. Hasche et al. [44] were able to find a good contrast between the SLD of Si and SiO<sub>2</sub>, which lead them to achieve thicknesses with low uncertainties. This was due to the photon energy of 1841 eV used because, at this photon energy, the Si extinction coefficient has a local maximum while the SiO<sub>2</sub> extinction coefficient remains low.

In this work, the fitting of those parameters revealed to be impracticable. When fitted the obtained results were physically incorrect or both values of SLD were identical when one considered the uncertainty of those values. So, the SLD of Si and SiO<sub>2</sub> used were  $18.886 \times 10^{-6}$  and  $20.124 \times 10^{-6} \text{ Å}^{-2}$  respectively.

Besides the problem with the SLD, the roughness of the native oxide could not be fitted in the SE data analysis. Here the roughness is represented by the thickness of a layer where the components are 50 % of void and 50 % of the native oxide. When the thickness of this layer was fitted, the result was 0.1 Å which is the limit of the software, so it is meaningless.

Nevertheless, as one can see in Table F.1, the SiO<sub>2</sub> thickness obtained by SE has a very low uncertainty. Because of that, for the SE data analysis of the Ta<sub>2</sub>O<sub>5</sub> samples, 20.0 Å was the fixed value of the  $SiO_2$  thickness used. However, the same did not happen with the XRR results. The uncertainty associated with SiO<sub>2</sub> thickness is about 11 % of its value so the SiO<sub>2</sub> thickness remained a fitting parameter in the XRR analysis. Despite that it is important to see that, given the associated uncertainties, both techniques are in agreement.

Table F.1: Results of XRR and SE analyses of a silicon	ı wafer usi	ing Si a	and SiO <sub>2</sub>	SLD	values	of
$18.886 \times 10^{-6}$ and $20.124 \times 10^{-6} \text{ Å}^{-2}$ , respectively.						

Talata	Si	O <sub>2</sub>
lechnique	Thickness (Å)	Roughness (Å)
SE	$20.0 \pm 0.3$	-
XRR	$18 \pm 2$	$5.2 \pm 0.3$

One important remark about the silicon wafer XRR analysis, as it is possible to see in Figure F.1, is the fact that the fitting does not present the  $SiO_2$  Kiessing fringes as it should. Nevertheless, the results obtained above (Table F.1) were the expected ones. This can be caused by a possible limitation of the software used.



Figure F.1: Real data and XRR fit of the silicon wafer.

### **RBS** analysis

In this Appendix the analysis of the RBS measured data of a single sample will be presented as an example. The data analysis procedure was the same for all samples.

In RBS a beam of ions is directed onto the sample to be measured. A He<sup>+</sup> beam was used in the present work. The number and the energy of the backscattered ions are measured. On one hand the number of backscattered ions with a specific energy is proportional to the concentration of the element of that respective energy. On the other hand the energy of each backscattered ion depends on two factors: the mass of the atom in which it was scattered and the depth of that atom in the sample.

In Figure G.1 the energy barriers are indicated for each element present in the sample. The Ar was also included since its presence is predicted due to its usage in the rf sputtering process. To perform the fit one assumed that the sample was made of  $Ta_2O_5$  and  $SiO_2$  besides Ar. The fit could have been done with the individual elements instead of the molecules but the quantification of O is very difficult and uncertain.



Figure G.1: Fit and treated data measured by RBS from the gr-Ta5Si sample.

From the fitting it is possible to get the information presented in Table G.1. The units of the obtained thickness are atoms  $cm^{-2}$ , but to compare with the other techniques it is necessary to convert to Å.

Table G.1: Results from the performed fit of the gr-Ta5Si sample.

Thickness $(10^{15}  ext{ atoms cm}^{-2})$	Ta <sub>2</sub> O <sub>5</sub>	SiO <sub>2</sub>	Ar
	(molecular%)	(molecular%)	(molecular%)
377.5	70.1	25.5	3.3

To do the conversion it is necessary to use the composition of the film, the mass density and molecular mass of each component. The problem is that Ar is a gas which means that it does not have a mass density value. To solve this problem an approximation was done: the Ar was not considered and the normalized composition was used. In other words, the assumed composition was 73.3 molecular % of  $Ta_2O_5$  and 26.7 molecular % of  $SiO_2$ . The mass densities used for the  $Ta_2O_5$  and  $SiO_2$  were 8.2 and 2.2 g cm<sup>-3</sup>, respectively [48].

The first step is to get a density in  $atoms cm^{-3}$  for each component. Here one uses the mass density, molecular mass and the Avogadro constant. Then, dividing the RBS thickness by the weighted density (taking into account the normalized composition) one gets the thickness in a length unit (in this case cm).

# Parameters of the Tauc-Lorentz models obtained by spectroscopic ellipsometry

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Sample	$E_g$ (eV)	ε <sup>ε</sup>	A~(eV)	E (eV)	C (eV)
Ta1	$3.6 \pm 0.2$	$3.2 \pm 0.2$	81 ± 17	$5.3 \pm 0.2$	$2.6 \pm 0.2$
Ta 2	$4.31\pm0.03$	$2.5 \pm 0.2$	$276 \pm 20$	$5.08\pm0.04$	$1.77\pm0.04$
Ta3	$4.31\pm0.02$	$2.12\pm0.05$	$323 \pm 10$	$4.97\pm0.03$	$1.84\pm0.03$
Ta4	$4.28\pm0.02$	$2.42\pm0.08$	$278 \pm 11$	$5.06 \pm 0.03$	$1.81\pm0.03$
Ta5	$4.20\pm0.02$	$2.14\pm0.05$	$264 \pm 15$	$5.09\pm0.04$	$2.02\pm0.03$
Ta6	$4.23 \pm 0.02$	$2.21 \pm 0.06$	$264 \pm 12$	$5.10\pm0.04$	$1.92 \pm 0.03$

Table H.2: Parameters from a double Tauc-Lorentz oscillator of TaSiO samples used to determine the grow rate (section 3.2.1).

Sample	Eg (eV)	ε <sup>∞</sup>	A1 (eV)	E1 (eV)	C1 (eV)	A2 (eV)	E2 (eV)	C2 (eV)
gr-Ta1Si	$-1.8 \pm 0.5$	$-46 \pm 10$	$494 \pm 75$	$7.38 \pm 0.04$	$1.84\pm0.06$	$112 \pm 17$	$4.37 \pm 0.02$	$0.55 \pm 0.4$
gr-Ta2Si	$4.773 \pm 0.007$	$0.8 \pm 0.3$	$42 \pm 10$	$24 \pm 4$	$1.0 \pm 0.4$	$353 \pm 10$	$4.40\pm0.03$	$2.46\pm0.03$
gr-Ta3Si	$4.47\pm0.02$	$0.4 \pm 0.2$	$69 \pm 5$	$10.4 \pm 0.3$	$0.02 \pm 0.097$	$90 \pm 4$	$5.58\pm0.02$	$1.78\pm0.03$
gr-Ta4Si	$4.30 \pm 0.02$	$1.9 \pm 0.1$	$24 \pm 2$	$8.0 \pm 0.4$	$2.6 \pm 0.4$	$132 \pm 6$	$5.44\pm0.02$	$1.89\pm0.03$
gr-Ta5Si	$4.20 \pm 0.02$	$1.0 \pm 0.2$	$64 \pm 7$	$10.2 \pm 0.6$	$7.8 \pm 0.6$	$116 \pm 4$	$5.48\pm0.02$	$1.70 \pm 0.03$
gr-Ta6Si	$4.18\pm0.02$	$1.48\pm0.07$	$77 \pm 9$	$7.1 \pm 0.3$	$5.6 \pm 0.4$	$87 \pm 5$	$5.45\pm0.02$	$1.49\pm0.05$
gr-Ta4	$4.09\pm0.02$	$1.4 \pm 0.2$	$83 \pm 15$	$7.9 \pm 0.9$	$9 \pm 2$	$143 \pm 9$	$5.27 \pm 8.3$	$1.57 \pm 0.04$

Sample	$E_g$ (eV)	ε <sub>∞</sub>	A1 (eV)	E1 (eV)	C1 (eV)	A2 (eV)	E2 (eV)	C2 (eV)
d-Ta2Si	$4.53 \pm 0.03$	$1.74 \pm 0.06$	$54 \pm 9$	$7.3 \pm 0.4$	$4.0 \pm 0.3$	$40 \pm 5$	$5.65 \pm 0.07$	$1.57 \pm 0.08$
d-Ta3Si	$4.25 \pm 0.02$	$1.96 \pm 0.03$	$59 \pm 4$	$6.71 \pm 0.08$	$2.9 \pm 0.2$	$40 \pm 2$	$5.52 \pm 0.02$	$1.11 \pm 0.03$
d-Ta4Si	$4.237 \pm 0.007$	$1.78 \pm 0.03$	$65 \pm 4$	$7.0 \pm 0.2$	$4.1 \pm 0.2$	72 ± 3	$5.48\pm0.02$	$1.40 \pm 0.03$
d-Ta5Si	$4.254 \pm 0.008$	$1.79 \pm 0.04$	$81 \pm 7$	$6.5 \pm 0.2$	$4.1 \pm 0.2$	$77 \pm 4$	$5.42\pm0.02$	$1.42 \pm 0.03$
d-Ta6Si	$4.19 \pm 0.02$	$1.80\pm0.04$	$96 \pm 8$	$6.3 \pm 0.2$	$4.0 \pm 0.2$	$66 \pm 4$	$5.39 \pm 0.02$	$1.19 \pm 0.03$
d-Ta4	$4.172 \pm 0.008$	$1.8 \pm 0.3$	$102 \pm 10$	$6.6 \pm 0.4$	$9.5 \pm 0.6$	$161 \pm 5$	$5.24 \pm 0.02$	$1.59 \pm 0.02$

Table H.4: Parameters from a double Tauc-Lorentz oscillator of ZTO samples in the study of the influence of the rf power (section 3.3.1).

Sample	$E_g$ (eV)	$\varepsilon_\infty$	A1 (eV)	E1~(eV)	C1 (eV)	A2 (eV)	E2 (eV)	C2 (eV)
P1-H1-O2	$2.85 \pm 0.05$	$1.2 \pm 0.5$	$94 \pm 4$	$10.5 \pm 0.6$	$15.6 \pm 0.8$	12 ± 2	$4.49 \pm 0.04$	$2.2 \pm 0.2$
P2-H1-O2	$3.0 \pm 0.1$	$1.0 \pm 0.7$	$136 \pm 13$	$13 \pm 1$	22 ± 2	$19 \pm 5$	$4.29\pm0.08$	$1.7 \pm 0.1$
P1-H2-O2	$3.00 \pm 0.04$	$1.5 \pm 0.4$	$91 \pm 4$	$9.5 \pm 0.5$	$14.3 \pm 0.6$	$12 \pm 2$	$4.51 \pm 0.03$	$1.9 \pm 0.1$
P2-H2-O2	$3.0 \pm 0.1$	$0.9 \pm 0.7$	$166 \pm 21$	$16 \pm 2$	$29 \pm 3$	$24 \pm 7$	$4.2 \pm 0.1$	$2.0 \pm 0.1$
P1-H3-O2	$2.96 \pm 0.04$	$1.1 \pm 0.5$	$93 \pm 4$	$10.3 \pm 0.6$	$15.1 \pm 0.6$	$12 \pm 2$	$4.54 \pm 0.04$	$2.0 \pm 0.2$
P2-H3-O2	$3.0 \pm 0.1$	$0.4 \pm 0.6$	$174 \pm 17$	$16 \pm 2$	$29 \pm 3$	$21 \pm 6$	$4.22 \pm 0.09$	$1.9 \pm 0.1$

## APPENDIX H. PARAMETERS OF THE TAUC-LORENTZ MODELS OBTAINED BY SPECTROSCOPIC ELLIPSOMETRY

Table H.5: Parameters from a double Tauc-Lorentz oscillator of ZTO samples in the study of the influence of the H<sub>2</sub> flow (section 3.3.2).

Sample	$E_g$ (eV)	ε <sup>ε</sup>	A1 (eV)	E1~(eV)	C1 (eV)	A2 (eV)	E2 (eV)	C2 (eV)
P1-H1-O2	$2.85 \pm 0.05$	$1.2 \pm 0.5$	$94 \pm 4$	$10.5 \pm 0.6$	$15.6 \pm 0.8$	12 ± 2	$4.49 \pm 0.04$	$2.2 \pm 0.2$
P1-H2-O2	$3.00 \pm 0.04$	$1.5 \pm 0.4$	$91 \pm 4$	$9.5 \pm 0.5$	$14.3 \pm 0.6$	$12 \pm 2$	$4.51 \pm 0.03$	$1.9 \pm 0.1$
P1-H3-O2	$2.96 \pm 0.04$	$1.1 \pm 0.5$	$93 \pm 4$	$10.3 \pm 0.6$	$15.1 \pm 0.6$	$12 \pm 2$	$4.54\pm0.04$	$2.0 \pm 0.2$
P2-H1-O2	$3.0 \pm 0.1$	$1.0 \pm 0.7$	$136 \pm 13$	$13 \pm 1$	22 ± 2	$19 \pm 5$	$4.29\pm0.08$	$1.7 \pm 0.1$
P2-H2-O2	$3.0 \pm 0.1$	$0.9 \pm 0.7$	$166 \pm 21$	$16 \pm 2$	$29 \pm 3$	$24 \pm 7$	$4.2 \pm 0.1$	$2.0 \pm 0.1$
P2-H3-O2	$3.0 \pm 0.1$	$0.4 \pm 0.6$	$174 \pm 17$	$16 \pm 2$	$29 \pm 3$	$21 \pm 6$	$4.22 \pm 0.09$	$1.9 \pm 0.1$
P2-H1-O3	$3.0 \pm 0.1$	$0.5 \pm 0.6$	$182 \pm 22$	$16 \pm 2$	$31 \pm 4$	$16 \pm 2$	$4.20\pm0.09$	$1.92 \pm 0.08$
P2-H2-O3	$3.01 \pm 0.09$	$-0.5 \pm 0.6$	$186 \pm 17$	$16 \pm 1$	$27 \pm 3$	$25 \pm 6$	$4.24\pm0.08$	$1.98\pm0.09$
P2-H3-O3	$3.0 \pm 0.1$	$-0.1 \pm 0.6$	$162 \pm 16$	$15 \pm 1$	$24 \pm 3$	22 ± 6	$4.34\pm0.08$	$1.9 \pm 0.1$

Table H.6: Parameters from a double Tauc-Lorentz oscillator of ZTO samples in the study of the influence of the O<sub>2</sub> flow (section 3.3.3).

Samla	F (aV)	с С	41 (AV)	E1 (aV)	C1 (aV)	4 7 (A)	F 7 (aV)	
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P2-H1-O1	$2.94 \pm 0.08$	$1.6 \pm 0.5$	$135 \pm 14$	$14 \pm 1$	25 ± 2	29 ± 6	$3.94 \pm 0.09$	$2.1 \pm 0.1$
P2-H1-O2	$3.0 \pm 0.1$	$1.0 \pm 0.7$	$136 \pm 13$	$13 \pm 1$	22 ± 2	$19 \pm 5$	$4.29\pm0.08$	$1.7 \pm 0.1$
P2-H1-O3	$3.0 \pm 0.1$	$0.5 \pm 0.6$	$182 \pm 22$	$16 \pm 2$	$31 \pm 4$	$16 \pm 2$	$4.20 \pm 0.09$	$1.92 \pm 0.08$
P2-H2-O2	$3.0 \pm 0.1$	$0.9 \pm 0.7$	$166 \pm 21$	$16 \pm 2$	29 ± 3	24 ± 7	$4.2 \pm 0.1$	$2.0 \pm 0.1$
P2-H2-O3	$3.01 \pm 0.09$	$-0.5 \pm 0.6$	$186 \pm 17$	$16 \pm 1$	27 ± 3	$25 \pm 6$	$4.24\pm0.08$	$1.98\pm0.09$
P2-H3-O2	$3.0 \pm 0.1$	$0.4 \pm 0.6$	$174 \pm 17$	$16 \pm 2$	29 ± 3	21 ± 6	$4.22 \pm 0.09$	$1.9 \pm 0.1$
P2-H3-O3	$3.0 \pm 0.1$	$-0.1 \pm 0.6$	$162 \pm 16$	$15 \pm 1$	$24 \pm 3$	22 ± 6	$4.34 \pm 0.08$	$1.9 \pm 0.1$

# The real and imaginary part of the dielectric function of the ZTO samples



Figure I.1: Real and imaginary part of the dielectric function of the samples characterized in section 3.3.1.

# APPENDIX I. THE REAL AND IMAGINARY PART OF THE DIELECTRIC FUNCTION OF THE ZTO SAMPLES



Figure I.2: Real and imaginary part of the dielectric function of the samples characterized in section 3.3.2.



Figure I.3: Real and imaginary part of the dielectric function of the samples characterized in section 3.3.3.