

Predictions on the transverse momentum spectra for charged particle production at LHC-energies from a two component model

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Abstract Transverse momentum spectra, $d^2\sigma/(d\eta dp_T^2)$, of charged hadron production in pp -collisions are considered in terms of a recently introduced two component model. The shapes of the particle distributions vary as a function of the c.m.s. energy in the collision and the measured pseudorapidity interval. As a result the pseudorapidity of a secondary hadron in the moving proton rest frame is shown to be a universal parameter describing the shape of the spectra in pp -collisions. In order to extract predictions on the double-differential cross sections $d^2\sigma/(d\eta dp_T^2)$ of hadron production for future LHC-measurements the different sets of available experimental data have been used in this study.

1 Introduction

Since the past decades a large volume of experimental data on charged hadron production spectra has been collected starting from the very first experiments performed on the Intersecting storage rings (ISR) at CERN till contemporary high statistics measurements carried out at the large hadron collider (LHC). It was shown [1] that the spectra of the produced hadrons are characterized by the exponential behavior for the low- p_T part of the spectra, while the high- p_T tail observes the power-like distribution typical for pQCD independently on the energy or type of the collision. However, these phenomena still lack theoretical explanation and can be described by various phenomenological models only. The most widely used model nowadays that combines exponential and power-like shapes is the Tsallis parameterization [2],

$$\frac{d^2\sigma}{d\eta dp_T^2} = \frac{A}{(1 + \frac{E_T}{T \cdot N})^N}. \quad (1)$$

Indeed, in the low $p_T \rightarrow 0$ limit this function becomes similar to the Boltzmann exponent, while for high- p_T it changes to a power-like shape. Recently a new qualitative model considering two sources of hadroproduction has been introduced [3,4], parameterizing the charged particle spectra by the sum of an exponential (Boltzmann-like) and a power-law p_T distribution:

$$\frac{d^2\sigma}{d\eta dp_T^2} = A_e \exp(-E_{Tkin}/T_e) + \frac{A}{(1 + \frac{p_T^2}{T^2 \cdot N})^N}, \quad (2)$$

where $E_{Tkin} = \sqrt{p_T^2 + M^2} - M$ with M equal to the produced hadron mass. A_e , A , T_e , T , N are free parameters, to be determined by a fit to the data. The detailed arguments for this particular choice are given in [3,4]. The exponential term in this model is associated with thermalized production of hadrons by valence quarks and a quark–gluon cloud coupled to them. The power-law term is related to the mini-jet fragmentation of the virtual partons (pomeron in pQCD) exchanged between two colliding partonic systems.

A typical charged particle spectrum as a function of the transverse momentum fitted with this function (2) is shown in Fig. 1. As one can see, the exponential term dominates the particle spectrum at low p_T values.

In [3,4] it was also shown that this model allows one to get a much better description of the available experimental data than Tsallis. The comparison of these two parameterizations is also shown in Fig. 1. As one can see from the data/fit ratio, the Tsallis parameterization does not describe the details of the spectra resulting in “a roughly log-periodic behavior on top of the q -exponential one” [6]. Moreover, it was shown that the relative contributions of the two components of (2) to the spectra vary with the type of the collisions [3,4] (eq , pp , γ

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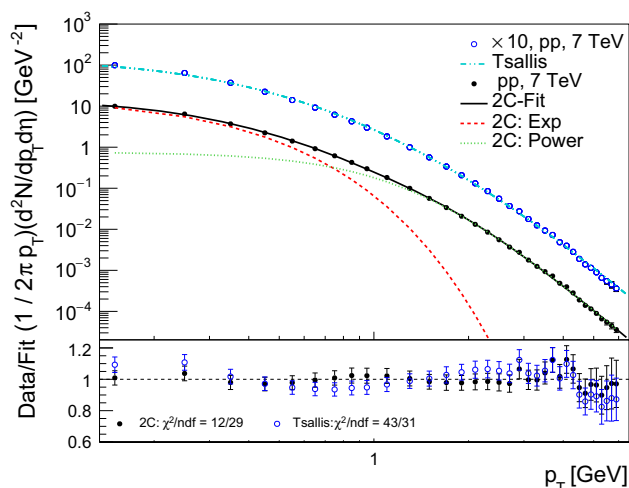


Fig. 1 Charged particle differential cross section $(1/2\pi p_T)d^2\sigma/d\eta dp_T$ [5] fitted to the two component (2C-fit) function (2) and Tsallis (1). The red (dashed) line shows the exponential term of the 2C-fit and the green (dotted) line shows the power-law term, the blue (dash-dot) line shows the Tsallis fit

$p, \gamma\gamma, AA$), the type of the produced hadron [7], the charged multiplicity [8] and the measured pseudorapidity region [9], supporting the idea that the functional form (2) is not just a better parameterization due to the larger number of free parameters, but rather reflecting the underlying dynamics of hadroproduction.

Separating “soft” and “hard” contributions with this model allowed one to make the predictions on the mean $\langle p_T \rangle$ values as a function of multiplicity in the collision [8] and pseudorapidity distributions of charged particles [9]. However, the major interest of many studies in QCD is the transverse momentum spectrum itself. Therefore, in this article it is discussed how its shape varies in different experiments under various conditions. In [3,4] it was shown that the parameters of the fit (2) show a strong dependence on the collision energy. Unfortunately, due to the fact that different collaborations measure charged particle production in their own phase space and under various experimental configurations, the dependences observed in [3,4] were smeared and did not allow one to make strong predictions for further measurements. Thus, an approach to correct the measurements in order to allow for an accurate combination of different experimental data is proposed here.

2 Parameter variations

In [9] it was shown that two sources of hadroproduction described above contribute to different pseudorapidity regions: while the power-law term of (2) prevails in the mid-rapidity region ($\eta \sim 0$), the exponential term dominates at

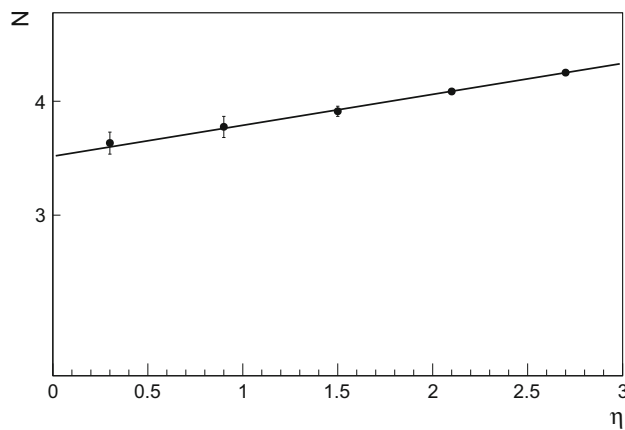


Fig. 2 Variation of the N parameter of (2) obtained from the fits to the experimental data [11] as a function of pseudorapidity. The line shows a power-law fit of this variation

high values of η . Since each collaboration presents measurements on transverse momentum spectra in various pseudorapidity intervals, these variations might explain the smearing of the dependences in [3,4]. The idea to study parameter variations as a function of both collision energy and pseudorapidity region has already been successfully tested in [10].

To further study the variations of the spectra shapes as a function of pseudorapidity we use the data published by the UA1 experiment [11], which present charged particle spectra in five pseudorapidity bins, covering the total rapidity interval $|\eta| < 3.0$. Figure 2¹ shows how the parameter N varies with the pseudorapidity together with a power-law fit of this variation. Note that the parameter shows a growth with pseudorapidity [9], which is explained by the higher thermalization of the spectra, as found in [9].

This variation can be parametrized in the following way:

$$N = N^0 \cdot \left(1 + 0.06 \cdot |\eta|^{1.52}\right) \tag{3}$$

where N^0 denotes the parameter value at $\eta \sim 0$ and η might be taken as a mean $\langle \eta \rangle$ of the measured pseudorapidity interval.

Since the variations of the parameter as a function of the pseudorapidity have been found, it is desirable to exclude its influence when studying the dependences of N on the c.m.s. energy in a collision. This is possible if one combines only those data that have been measured in more or less the same pseudorapidity intervals. Hence, it is prudent to look at the combined data taken by the ISR [13], PHENIX [14] and ALICE [15] that were measured in the $|\eta| < 0.8$ pseudorapidity region.

¹ The error bars on this and further plots are obtained from the fits (2) to the experimental data (with statistical and systematic errors added in quadrature) performed in the ROOT framework [12].

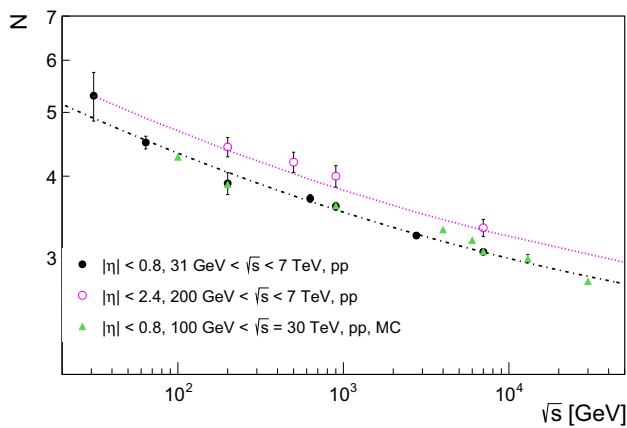


Fig. 3 Variation of the N parameter of (2) obtained from the fits to the experimental data [13–15] and to the Monte Carlo generated data (full points) as a function of the c.m.s. energy \sqrt{s} in a collision. The dash-dot line shows a fit of this variation. In addition, the open points show the parameters for the data measured in another pseudorapidity interval [5, 16] and the dotted line shows the predictions calculated according to (3)

Figure 3 shows the parameter N variation as a function of the c.m.s. energy in a collision. One can notice that it can be characterized by the falling N -value. It is related to the fact that the probability to produce a high- p_T mini-jet should grow with \sqrt{s} . Notably, this behavior correlates with the fact that N decreases when the rapidity interval between the secondary hadron and the initial proton increases.

One can extract the following parametrization from the fit shown in Fig. 3²:

$$N = 2.01 + 5.46 \cdot s^{-0.093} \tag{4}$$

Remarkably, in the $s \rightarrow \infty$ limit $N \rightarrow 2$, which corresponds at large p_T to $d\sigma/dp_T^2 \propto 1/p_T^4$, such a behavior is expected in pQCD where the elementary parton-parton cross section $d\hat{\sigma}/dp_T^2 \propto \alpha_s^2/p_T^4$, that is, up to the logarithms $d\sigma/dp_T^2 \propto 1/p_T^4$. However, in a real collision one should take the initial conditions and the kinematic restrictions into account, resulting in a higher value of N for lower \sqrt{s} . Obviously, such initial conditions should become negligible in the $s \rightarrow \infty$ limit, which is confirmed by the observed behavior (4).

In addition, Fig. 3 shows UA1 [16] and CMS [5] data measured under different experimental conditions. In these measurements the pseudorapidity interval was much wider ($|\eta| < 2.5$) than in [13–15]. Therefore, one can compare the parameter values obtained from the fit of these data (open points in Fig. 3) to the values extrapolated from (3) with N^0 calculated according to (4) and $|\eta| = 1.25$ (pointed lines) and see a rather good agreement. Notably, this behavior correlates with the fact that N decreases when the rapidity inter-

² In the (4) and below \sqrt{s} is given in GeV.

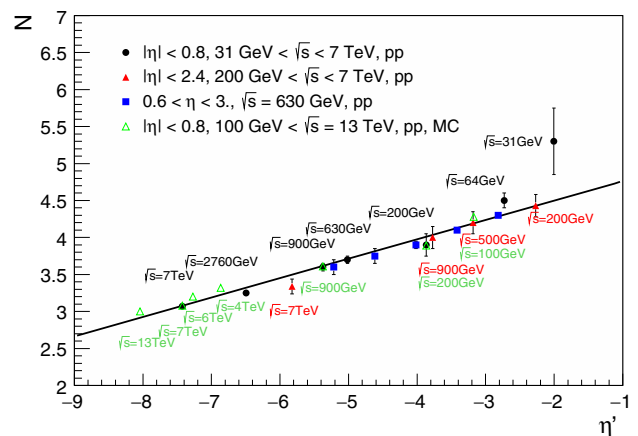


Fig. 4 The dependence of the N parameter on the pseudorapidity of the secondary hadron in the moving proton rest frame. The data points from different experiments [5, 11, 13–16] are shown together with the generated data (MC). The line shows the linear fit as a function of η' expressed by (6)

val between the secondary hadron and the same side beam proton increases.

Let us now check this correlation explicitly and calculate the rapidity interval in the moving proton rest frame according to a simple formula:

$$\eta' = |\eta| - \log(\sqrt{s}/2m_p), \tag{5}$$

where m_p is the mass of the incoming proton. The results of this procedure are shown in Fig. 4. Surprisingly, all the points came to a single line in this interpretation. To understand the origin of this universality one might use Monte Carlo (MC) generators: hard processes at large p_T are known to be described by MC generators pretty well, thus it is expected to get the value of the N -parameter from the fits of the MC-generated spectra rather close to the real data, but with a higher accuracy and in a wider collision energy range. To check this universality, we have produced the Monte Carlo samples for proton-proton collisions at different energies for inelastic (INEL) events with the PYTHIA 8.2 generator [17]. Indeed, the values of the parameter N extracted from the fits to the MC-generated spectra are nicely placed at the same line. Thus, a universal parameter describing the shape of the transverse momentum spectra in pp -collisions has been found.

$$N = 5.04 + 0.27\eta' \tag{6}$$

A further check of the observed phenomena can be made if one plots N as a function of the logarithm of the maximal kinematically allowed transverse momentum $p_{T,max}$ of the secondary hadron at the specific pseudorapidity η .

$$p_{T,max} = \frac{\sqrt{s}}{2} \sin[2 \cdot \tan^{-1}(\exp(-|\eta|))] \tag{7}$$

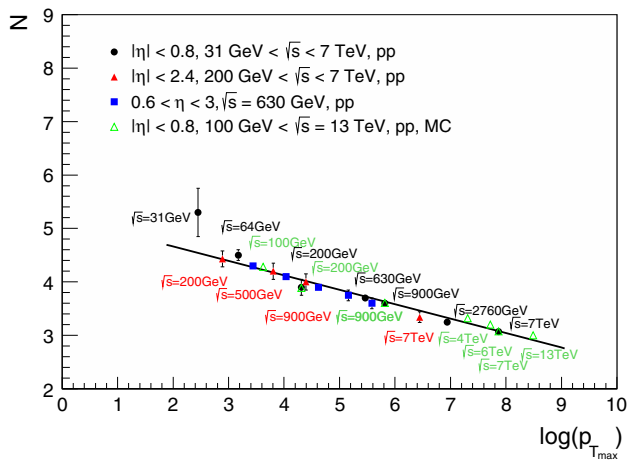


Fig. 5 The dependence of the N parameter on the maximal kinematically allowed transverse momentum of the secondary hadron. The data points from different experiments [5, 11, 13–16] are shown together with the generated data (MC). The line shows the linear fit as a function of p_{Tmax}

This dependence is shown in Fig. 5. One can see that larger N corresponds to smaller p_{Tmax} . That should correspond to the $x \rightarrow 1$ limit of PDFs, where the decrease of the perturbative cross section is modified by the fall-off of the parton distribution functions (PDFs).

Remarkably, similarly to N , the T and T_e also show dependences as a function of both the collision energy \sqrt{s} and the measured pseudorapidity interval η . The variations of the T and T_e parameters were studied in [10]. In [10] the possible theoretical explanation of the thermalized particle production was presented and the following proportionalities were established:

$$T = 409 \cdot (\sqrt{s})^{0.06} \cdot \exp(0.06|\eta|) \text{ MeV}, \tag{8}$$

$$T_e = 98 \cdot (\sqrt{s})^{0.06} \cdot \exp(0.06|\eta|) \text{ MeV}. \tag{9}$$

The parametrizations for T and T_e differ only by a constant factor. However, both the T and the T_e parameters reflect the thermalization which is stronger at higher energies and when closer to the valence quarks. Therefore, the (8) and (9) parametrizations which are functions of the center of mass energy and rapidity interval can be rewritten in a form with only one universal parameter. This universal parameter is the rapidity distance η'' from the farther incoming proton.

$$\eta'' = |\eta| + \log(\sqrt{s}/2m_p) \tag{10}$$

Using (8, 9) we get the universal dependence³:

$$T = 409 \cdot \exp(0.06\eta'') \cdot (2m_p)^{0.06} \text{ MeV}, \tag{11}$$

$$T = 98 \cdot \exp(0.06\eta'') \cdot (2m_p)^{0.06} \text{ MeV}. \tag{12}$$

The dependences are shown on Fig. 6.

³ In (11), (12), and (23) m_p is given in units of 1 GeV/c².

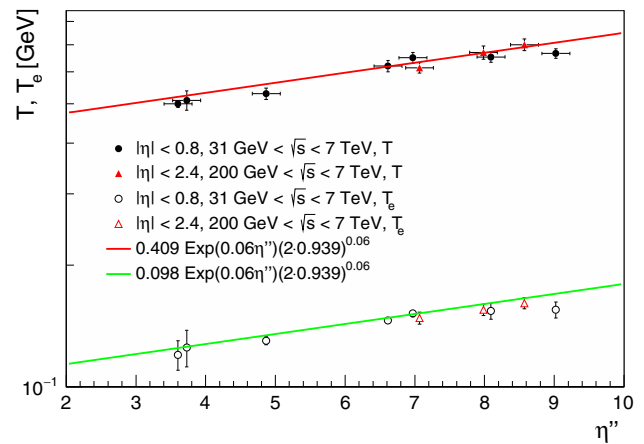


Fig. 6 The dependence of the T and T_e parameters on the pseudo-rapidity of the secondary hadron in the moving opposite side proton rest frame. The data points from different experiments [5, 11, 13–16] are shown. The lines show the dependences (11)–(12), respectively

3 Prediction for further measurements

Though MC was shown to provide a nice description of the high-pt part of the spectra, the nature of the soft particle production still remains ambiguous and varies for different MC generators. Therefore, the next step in understanding the underlying dynamics of high energy hadronic processes was done in a recent analysis [10].

In [9] it was shown that the introduced approach is able to give predictions on the pseudorapidity distributions in high energy collisions for non-single diffractive events (NSD). Using the parameterizations from [9] in addition with (3)–(9) one can provide a formula that describes the shapes of charged particle spectra, being a function only of the center of mass energy and a measured pseudorapidity region. Let us now summarize all the equations to obtain the final result.⁴

$$\left(\frac{d\sigma}{d\eta}\right)_{\text{power}} = A_{\text{power}} \cdot \exp\left(-\frac{\eta^2}{2\sigma_{\text{power}}^2}\right), \tag{13}$$

$$\begin{aligned} \left(\frac{d\sigma}{d\eta}\right)_{\text{exp}} &= A_{\text{exp}} \cdot \exp\left(-\frac{(\eta - \eta_{\text{exp}})^2}{2\sigma_{\text{exp}}^2}\right) \\ &+ A_{\text{exp}} \cdot \exp\left(-\frac{(\eta + \eta_{\text{exp}})^2}{2\sigma_{\text{exp}}^2}\right), \end{aligned} \tag{14}$$

$$\sigma_{\text{power}} = 0.217 + 0.235 \cdot \ln \sqrt{s}, \tag{15}$$

$$\eta_{\text{exp}} = 0.692 + 0.293 \cdot \ln \sqrt{s}, \tag{16}$$

$$\sigma_{\text{exp}} = 0.896 + 0.136 \cdot \ln \sqrt{s}, \tag{17}$$

$$A_{\text{power}} = 0.13 \cdot s^{0.175}, \tag{18}$$

$$A_{\text{exp}} = 0.76 \cdot s^{0.106}. \tag{19}$$

⁴ In (13) and below, A and A_e is measured in mb/GeV², A_{pow} and A_{exp} in mb.

Now, with the knowledge of the variations of the T , N , and T_e parameters and the exponential and power-law contributions one can calculate the normalization parameters A and A_e in (2) in the following way:

$$\frac{d\sigma}{d\eta} = \left(\frac{d\sigma}{d\eta}\right)_{\text{exp}} + \left(\frac{d\sigma}{d\eta}\right)_{\text{pow}} = \int_0^\infty \frac{d^2\sigma}{d\eta dp_T^2} dp_T^2, \quad (20)$$

$$\left(\frac{d\sigma}{d\eta}\right)_{\text{pow}} = \int_0^\infty \frac{A}{\left(1 + \frac{p_T^2}{T^2 \cdot N}\right)^N} dp_T^2 = \frac{AN T^2}{N-1}, \quad (21)$$

$$\begin{aligned} \left(\frac{d\sigma}{d\eta}\right)_{\text{exp}} &= \int_0^\infty A_e \exp(-E_{Tkin}/T_e) dp_T^2 \\ &= 2A_e T_e (m + T_e). \end{aligned} \quad (22)$$

Thus, we get the set of equations allowing us to make a prediction for a double-differential cross section, using the formula (2):

$$\begin{cases} |\eta| = (|\eta|_{\text{max}} + |\eta|_{\text{min}})/2, \\ N = 5.04 + 0.27\eta' \\ T = 409 \cdot \exp(0.06\eta'') \cdot (2m_p)^{0.06}, \\ T = 98 \cdot \exp(0.06\eta'') \cdot (2m_p)^{0.06}, \\ A_e = \frac{1}{2T_e(m+T_e)} (d\sigma/d\eta)_{\text{exp}}, \\ A = \frac{(N-1)}{NT^2} (d\sigma/d\eta)_{\text{pow}}, \end{cases} \quad (23)$$

where η' and η'' can be calculated using (5) and (10) correspondingly.

Now, one can calculate double-differential cross sections $d^2\sigma/(d\eta dp_T^2)$ of charged particle production in high energy collisions at different energies for NSD events. These predictions are shown in Fig. 7 for $|\eta| < 0.8$ and the $|\eta| < 2.4$ pseudorapidity intervals together with the experimental data measured by ALICE [18] and CMS [5]. A good agreement of the prediction with the data can be observed. Thus, the

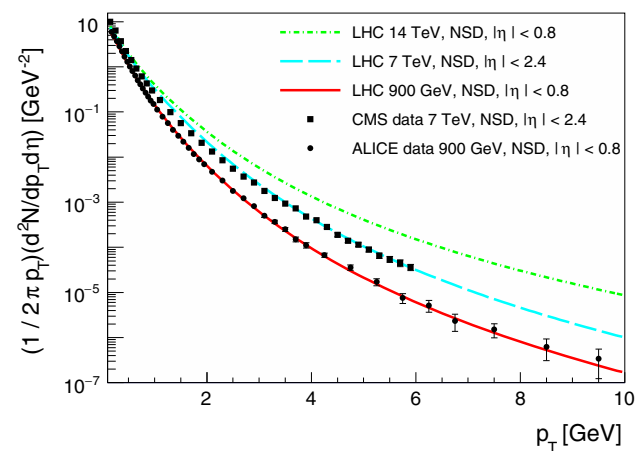


Fig. 7 Predictions of the yield of charged particles $(1/2\pi p_T)d^2N/(d\eta dp_T)$ in high energy collisions in NSD events together with data points from the ALICE [18] and CMS [5] experiments

results (23) give us a powerful tool for predicting the spectral shapes in NSD events.

4 Conclusion

In conclusion, transverse momentum spectra in pp -collisions have been considered using a two component model. Variations of the parameters obtained from the fit have been studied as a function of the pseudorapidity η and the c.m.s. energy \sqrt{s} in the collision. A universal parameter describing a shape of the spectra in pp -collisions was found to be the pseudorapidity of a secondary hadron in the moving proton rest frame. Finally, the observed dependences, together with previous investigations allowed one to make predictions on double-differential spectra $d^2\sigma/(dp_T^2 d\eta)$ at higher energies, successfully tested on the available experimental data.

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