Quantum Size Effects And Transport Phenomena In PbSe Quantum Wells And PbSe/EuS Superlattices

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Abstract. It is established that the room-temperature dependences of transport properties on the total thickness of PbSe layers d in PbSe/EuS superlattices exhibit an oscillatory behavior. It is shown that the oscillation period Δd practically coincides with the period of the thickness oscillations observed earlier in single PbSe/EuS quantum well. The non-monotonic character of these dependences is attributed to quantum size effects. The theoretically estimated and experimentally determined Δd values are in good agreement.

Keywords: PbSe, quantum well, superlattice, transport phenomena, quantum size effects. **PACS:** 73.21.-b; 73.50.Jt; 73.50.Lw.

INTRODUCTION

IV-VI semiconductors are known as promising materials for optoelectronic, thermoelectric, and other applications [1]. These compounds are also convenient objects for studying quantum size effects (QSEs) by measuring transport properties due to a small effective mass and a high charge carrier mobility. In [2], we reported the observation of oscillations in the thickness dependences of the transport properties in PbSe quantum wells (QWs) and attributed them to the QSEs. Later we observed similar oscillations in other IV-VI QWs (see e.g., [3,4]).

The goal of this work is to compare the character of the thickness dependences of the transport properties in single PbSe QWs and PbSe/EuS superlattices (SLs) grown on (001) KCl substrates.

EXPERIMENTAL

In [2,5] we described in detail the preparation technique and the growth mechanism for (001) KCl/ PbSe/EuS SLs. In the present work the thickness of PbSe layers (d_0) in SLs was varied in the range of $d_0 = 0.5 - 15$ nm, whereas the thickness of a EuS layers was ~10 nm. The upper (covering) EuS layer was 30-40 nm thick. The number of SLs periods was varied from 5 to 25. All measurements were conducted on freshly prepared samples. The electrical conductivity σ and the Hall coefficient R_H were measured using dc method, with an error not exceeding 5 %. The Hall concentration n and mobility μ of carriers were calculated as $n = 1/(R_H \cdot e)$ and $\mu = R_H \cdot \sigma$, respectively. The Seebeck coefficient S was measured relative to Cu with an accuracy of \pm 3 %. All SLs exhibited *n*-type conductivity.

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RESULTS AND DISCUSSION

In Fig. 1, the dependences of σ , n, μ , and S on the total thickness d of all PbSe layers ($d = N \cdot d_0$) for PbSe/EuS SLs are presented. All the dependences exhibit an oscillatory behavior with the mean value of the oscillation period $\Delta d = 22 \pm 3$ nm. The Δd value was determined using the value of d corresponding to the position of the first extremum in the dependences of all kinetic coefficients as well as using the mean distance between the minima and maxima.

In the inserts to Fig. 1, the thickness dependences of the kinetic coefficients for the single (001)/PbSe/EuS QWs obtained in [2] are presented. The oscillation period measured as was indicated above for SLs ($\Delta d = 20 \pm 5$ nm) is very close to Δd for the SLs. The experimental value of Δd for single QWs ($\Delta d = 35 \pm 3$ nm) reported in [2] is an overestimate because Δd was determined only from the distance between extrema, which is less accurate in case when the number of studied samples with different thicknesses is limited.

One can assume that the observed oscillatory behavior of the *d*-dependences of the transport properties both in PbSe/EuS SLs and in single PbSe/EuS QWs can be interpreted in terms of quantum-mechanical concepts about a gradual filling of the lateral subbands with increasing *d*. The difference between QWs and SLs consists in the fact

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FIGURE 1. Dependences of the electrical conductivity σ , Seebeck coefficient *S*, electron mobility μ , and concentration *n* on total PbSe layers thickness *d* in PbSe/EuS superlattices. In the inserts: the dependences of σ , μ , *n*, and S on thickness for single PbSe/EuS quantum wells reported in [2].

that in SLs one should take into account the total thickness of all PbSe layers $(d = N \cdot d_0)$, which can be

changed either by changing the PbSe layer thickness (d_0) or by changing the number of periods (N). It is taken into account that the conductivity in the studied SLs occurs in the interface planes and that the barrier width is sufficiently small (10 nm) for the realization of the tunneling processes. Every time *d* increases by $\lambda_F/2$, where λ_F is the de Broglie wavelength at the Fermi level ε_{F} , an additional subband drops below ε_{F} , and contributes to electrical conduction. One can estimate the oscillation period Δd (see [2-5]) using the following equation:

$$\Delta d = \frac{\lambda_{\rm F}}{2} = \frac{h}{\sqrt{8m_z^*\varepsilon_{\rm F}}} \tag{1}$$

where m_z^* is the effective mass for motion perpendicular to the QWs. The critical thickness d_1 , at which the first subband passes through the Fermi level $(E_I = \varepsilon_F)$, can be estimated as $d_I = h/\sqrt{8}m_z^* \varepsilon_F$. = Δd . Knowing effective masses of the electrons in *n*-PbSe $(m_t^* = 0.04 \ m_0, \ m_1 = 0.07 \ m_0 \ [1])$ and estimating ε_F from the average electron concentration $(n \approx 2 \cdot 10^{18} \ \text{cm}^{-3})$, one can calculate Δd using Eq. 1, thus obtaining $\Delta d = 20 \pm 5 \ \text{nm}$, which is in good agreement with the experimentally determined values of Δd for QWs (Δd = 20 $\pm 5 \ \text{nm} \ [2]$) and for SLs ($\Delta d = 22 \pm 3 \ \text{nm}$).

Thus, the experimentally determined periods of the quantum thickness oscillations of the transport properties in PbSe/EuS single QWs and PbSe/EuS SLs practically coincide and correspond to the results of theoretical calculations.

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