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## Quantifying Performance in Fading Channels Using the Sampling Property of a Delta Function

### Won Mee Jang

Abstract—We apply the sampling property of a delta function to obtain the probability of error in fading channels. Our approach reduces the integration to a sampling. The sampling point is obtained in terms of fading parameters and the average signal-to-noise ratio (SNR) to provide the closed form solution of the performance.

*Index Terms*—Fading channel, performance, delta function, sampling.

#### I. INTRODUCTION

**O** UR sampling method is remarkably simple compared to other performance analysis methods for wireless communication in fading channels discussed in [1-8]. The bit error rate (BER) obtained from the proposed method, called delta approximation denoted by ' $\doteq$ ', is graphically or numerically displayed and compared to the exact theoretical result for various fading channels. We extend the result to the integer power of the *Q*-function to exhibit the average symbol error probability (ASEP) of the differentially encoded quadrature phase-shift keying (DE-QPSK).

#### **II. DERIVATION OF DELTA APPROXIMATION**

We can express the probability density function (PDF) of the fading channel as

$$p(\beta) = K \exp\{-b\beta\}\beta^{c-1}f(\beta) \tag{1}$$

where  $\beta$  is the magnitude square of the fading gain. K is a constant, and  $f(\beta)$  is an auxiliary function that depends on fading characteristics. For example, K = 1, b = 1, c = 1 and  $f(\beta) = 1$  for Rayleigh fading channels. Then the integer power probability of error in fading channels can be obtained as

$$P_b^p(\bar{\gamma}) = \int_0^\infty Q^p(\sqrt{\bar{\gamma}\beta})p(\beta)d\beta \tag{2}$$

$$= K \int_0^\infty Q^p(\sqrt{\bar{\gamma}\beta}) \exp\{-b\beta\} \beta^{c-1} f(\beta) d\beta$$
(3)

where  $Q(\alpha) = \int_{\alpha}^{\infty} (2\pi)^{-1/2} e^{-y^2/2} dy$ .  $\bar{\gamma}$  is the average received SNR. For binary phase-shift keying (BPSK),  $\bar{\gamma} = 2\bar{\gamma}_b$  where  $\bar{\gamma}_b = E_b/N_o$ .  $E_b$  is the average bit energy and  $N_o$  is the one-sided noise power spectral density (PSD).

*Coherent detection:* To find the sampling point, we introduce the *Q*-function approximation,

$$Q_1(\alpha) \approx \frac{1}{\sqrt{2\pi\alpha}} \exp\{-\alpha^2/2\} - \frac{1}{\sqrt{2\pi}(\alpha+1)} \exp\{-(\alpha+1)^2/2\}$$
(4)

Manuscript received October 15, 2010. The associate editor coordinating the review of this letter and approving it for publication was Y. Chen.

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Digital Object Identifier 10.1109/LCOMM.2011.011011.101964

where the first term (we call it  $Q_o(\alpha)$ ) is the well-known upper bound of the Q-function. We show  $Q_o(\alpha) \ge Q_1(\alpha) \ge Q(\alpha)$  in Appendix A. The integer power of the Q-function approximation can be found using the binomial expansion:

$$Q_{1}^{p}(\alpha) \approx \sum_{k=0}^{p} {p \choose k} (-1)^{k} \left(\frac{1}{\sqrt{2\pi\alpha}} \exp\{-\alpha^{2}/2\}\right)^{p-k} \\ \left(\frac{1}{\sqrt{2\pi}(\alpha+1)} \exp\{-(\alpha+1)^{2}/2\}\right)^{k}.$$
 (5)

With a change of variables  $(\bar{\gamma}\beta = x)$  and  $Q_1^p(\alpha)$  into Eq. (3),

$$P_b^p(\bar{\gamma}) \approx \frac{K}{\bar{\gamma}^c} \int_0^\infty Q_1^p(\sqrt{x}) \exp\{-b(x/\bar{\gamma})\} x^{c-1} f(x/\bar{\gamma}) dx.$$
(6)

Applying Eq. (5) to Eq. (6) and with another change of variables  $(x = y^N)$ ,

$$P_b^p(\bar{\gamma}) \approx \left(\frac{1}{\sqrt{2\pi}}\right)^p \frac{K}{\bar{\gamma}^c} \sum_{k=0}^p \binom{p}{k} (-1)^k \exp\{-k/2\}$$
$$\int_0^\infty f(y^N/\bar{\gamma})(\sqrt{y^N}+1)^{-k}(\sqrt{y^N})^{-(p-k)} \exp\{-k\sqrt{y^N}\}$$
$$\exp\{-(p/2+b/\bar{\gamma})y^N\}y^{N(c-1)}Ny^{N-1}dy.$$
(7)

With  $a = p/2 + b/\bar{\gamma}$ , let us define

$$g(y^{N}) := \exp\{-ay^{N}\}Ny^{cN-1}.$$
(8)

Then, we find

$$\int_0^\infty g(y^N) dy = \frac{\Gamma(c)}{a^c} \tag{9}$$

with the Gamma function defined as  $\Gamma(\alpha) := \int_0^\infty y^{\alpha-1} e^{-y} dy$  for  $\alpha > 0$ . We also show in Appendix B that

$$\lim_{N \to \infty} g(y^N) = \frac{\Gamma(c)}{a^c} \delta\left(y^N - \frac{c}{a}\right).$$
(10)

Now, applying the sampling property of a delta function to Eq. (7),

$$P_b^p(\bar{\gamma}) \doteq \left(\frac{1}{\sqrt{2\pi}}\right)^p \frac{K}{\bar{\gamma}^c} \frac{\Gamma(c)}{a^c} \sum_{k=0}^p \binom{p}{k} (-1)^k \exp\{-k/2\}$$

$$\int_0^\infty f(y^N/\bar{\gamma})(\sqrt{y^N}+1)^{-k} (\sqrt{y^N})^{-(p-k)} \exp\{-k\sqrt{y^N}\}$$

$$\delta\left(y^N - \frac{c}{a}\right) dy = \left(\frac{1}{\sqrt{2\pi}}\right)^p \frac{K\Gamma(c)}{(\bar{\gamma}a)^c} f\left(\frac{c}{a\bar{\gamma}}\right)$$

$$\left[\sqrt{\frac{a}{c}} - \left(\sqrt{\frac{c}{a}}+1\right)^{-1} \exp\left\{-\left(\sqrt{\frac{c}{a}}+\frac{1}{2}\right)\right\}\right]^p.$$
(11)

TABLE I Parameters for Fading Channels

fading	K	a	c	$f\left(\frac{c}{a\bar{\gamma}}\right)$
Nakagami-m	$m^m/\Gamma(m)$	$p/2 + m/ar{\gamma}$	m	1
Nakagami-n	$(1+n^2)\exp\{-n^2\}$	$p/2 + 1/\bar{\gamma}$	1	$\exp\{-n^2/(a\bar{\gamma})\}I_o\left(2n\sqrt{(1+n^2)/(a\bar{\gamma})}\right)$
Nakagami- $q$	$(1+q^2)/(2q)$	$p/2 + (1+q^2)^2/(4q^2\bar{\gamma})$	1	$I_o\left((1-q^4)/\{(4q^2)(a\bar{\gamma})\}\right)$

)



Fig. 1. Nakagami-m fading, BPSK, m=0.5, 1, 1.5, 2, 2.5, 3, 3.5, 4, 5, 6.

Noncoherent detection: The BER in a noncoherent additive white Gaussian noise (AWGN) channel can be expressed as  $P_e(\bar{\gamma}) = 2^{-1}e^{-\bar{\gamma}/2}$  where  $\bar{\gamma} = 2\bar{\gamma}_b$  for differential PSK (DPSK) [3, Eq. (14.3-9)]. Then, the integer power of the BER in fading channels can be expressed using Eq. (1) as

$$P_b^p(\bar{\gamma}) = \frac{K}{2^p} \int_0^\infty \exp\{-p\bar{\gamma}\beta/2\} \exp\{-b\beta\}\beta^{c-1} f(\beta)d\beta.$$
(12)

With a change of variables  $(\bar{\gamma}\beta = x)$ ,

$$P_b^p(\bar{\gamma}) = \frac{K}{2^p} \frac{1}{\bar{\gamma}^c} \int_0^\infty \exp\{-px/2\} \exp\{-bx/\bar{\gamma}\} x^{c-1} f(x/\bar{\gamma}) dx.$$
(13)

With another change of variables  $(x = y^N)$ ,

$$P_b^p(\bar{\gamma}) = \frac{K}{2^p} \frac{1}{\bar{\gamma}^c} \int_0^\infty \exp\{-(p/2 + b/\bar{\gamma})y^N\} \\ y^{N(c-1)} f(y^N/\bar{\gamma})Ny^{N-1} dy.$$
(14)

Applying the same process in coherent detection, the delta approximation can be obtained as

$$P_b^p(\bar{\gamma}) \doteq \frac{K}{2^p} \frac{\Gamma(c)}{(\bar{\gamma}a)^c} f(c/a\bar{\gamma}) \,. \tag{15}$$

#### **III. NUMERICAL RESULTS**

*Nakagami-m fading:* The PDF of the fading is [1]

$$p(\beta) = \frac{m^m \beta^{m-1}}{\Gamma(m)} \exp\{-m\beta\}, \quad \beta \ge 0.$$
 (16)

Making reference to Eq. (1), we choose  $K = m^m/\Gamma(m)$ , c = m,  $a = p/2 + m/\bar{\gamma}$  and  $f(\frac{c}{a\bar{\gamma}})=1$ . We summarize the parameter selection in different fadings in Table I. The result is compared to the exact performance of BPSK [2, Eq. (5.80)]

$$P_e = \frac{1}{\pi} \int_0^{\pi/2} \left( 1 + \frac{\bar{\gamma}_b}{m \sin^2 \phi} \right)^{-m} d\phi.$$
 (17)



Fig. 2. Nakagami-n fading, BPSK, n=1, 2, 2.5, 3.

In Fig. 1 we plot the exact BER and the delta approximation for m=0.5 to 6. We can see that the delta approximation and the exact performance agree well.

*Nakagami-n fading:* For Nakagami-n fading channels, the fading PDF is shown as [1]

$$p(\beta) = (1+n^2) \exp\{-n^2\} \exp\{-(1+n^2)\beta\}$$
$$I_o(2n\sqrt{(1+n^2)\beta}), \quad \beta \ge 0$$
(18)

and we choose  $K = (1 + n^2) \exp\{-n^2\}$ , c = 1,  $a = p/2 + 1/\bar{\gamma}$  and  $f(\frac{c}{a\bar{\gamma}}) = \exp\{-n^2/(a\bar{\gamma})\}$   $I_o(2n\sqrt{(1 + n^2)/(a\bar{\gamma})})$ . It is important to choose an auxiliary function  $f(\cdot)$  as flat as possible with respect to (w.r.t.) the SNR to be robust against the sampling point error that may be introduced by the Q-function approximation. The modified Bessel function  $I_o(\cdot)$  is a decreasing function w.r.t. the SNR while  $\exp\{-n^2/(a\bar{\gamma})\}$  is an increasing function. As a result, the product of the two is rather flat w.r.t. the SNR and thus less sensitive to the sampling point error. The BER of Nakagami-n with BPSK is shown in Fig. 2 for n=1, 2, 2.5 and 3. The delta approximation agrees well with the exact BER at a moderate and high SNR. The exact performance is numerically obtained from Eq. (2) with p = 1 using a double integration with the modified Bessel function.

*Nakagami-q fading:* The PDF of Nakagami-*q* fading channels is presented as [1]

$$p(\beta) = \frac{1+q^2}{2q} \exp\left\{-\frac{(1+q^2)^2\beta}{4q^2}\right\} I_o\left(\frac{(1-q^4)\beta}{4q^2}\right), \ \beta \ge 0$$
(19)

and we choose  $K = (1 + q^2)/(2q)$ , c = 1,  $a = p/2 + (1 + q^2)^2/(4q^2\bar{\gamma})$ , and  $f(\frac{c}{a\bar{\gamma}}) = I_o((1 - q^4)/\{(4q^2)(a\bar{\gamma})\})$ . Since Nakagami-q fading with q = 1 corresponds to the Rayleigh fading channel, we compare the delta approximation to the exact performance of BPSK [3, Eq. (14.3-7)]

$$P_e = (1/2) \left( 1 - \sqrt{\bar{\gamma}_b / (1 + \bar{\gamma}_b)} \right).$$
 (20)



Fig. 3. Nakagami-q fading, q=0.6 (DPSK), q=1 (BPSK).

The delta approximation of DPSK is obtained from Eq. (15) with p = 1 and then compared to the exact performance [2, Eqs. (2.12) & (8.201)]

$$P_e = (1/2) \left\{ 1 + 2\bar{\gamma}_b + (2\bar{\gamma}_b q)^2 / (1+q^2)^2 \right\}^{-1/2}.$$
 (21)

The delta approximation agrees well with the exact performance as shown in Fig. 3.

*Differentially Encoded QPSK:* The ASEP of the DE-QPSK in the AWGN channel is given by [4]

$$P_e = 4Q(\sqrt{\overline{\gamma}}) - 8Q^2(\sqrt{\overline{\gamma}}) + 8Q^3(\sqrt{\overline{\gamma}}) - 4Q^4(\sqrt{\overline{\gamma}}).$$
 (22)

Therefore, the ASEP of the DE-QPSK in fading channels can be obtained as

$$P(\bar{\gamma}) \doteq 4P_b^1(\bar{\gamma}) - 8P_b^2(\bar{\gamma}) + 8P_b^3(\bar{\gamma}) - 4P_b^4(\bar{\gamma})$$
(23)

with  $P_b^p(\bar{\gamma})$  defined in Eq. (11). The ASEP of the DE-QPSK in Nakagami-*m* fading channels is shown in Tables II and III for *m*=2.5 and 3.5, respectively. The result is compared to approximations in Chiani [5, Eq. (14)], Börjesson [6, Eq. (9)], Karagiannidis [7, Eq. (12)] and Isukapalli [8, Eq. (11)]. The delta approximation provides a tighter approximation than Chiani and Börjesson for  $\bar{\gamma} \geq 10$  dB for both *m*=2.5 and 3.5. Our method also provides a tighter approximation than Karagiannidis and Isukapalli for  $\bar{\gamma} \geq 20$  dB for both *m*=2.5 and 3.5. Karagiannidis ASEP is obtained with parameter values specified in [7]. Isukapalli ASEP is obtained with the number of terms in Taylor series expansion,  $N_a$ =14. The order of its computational complexity is  $O(\exp\{N_a \ln(N_a)\})$  that can make Isukapalli method infeasible for a large value of *m* or SNR.

#### APPENDIX A

 $Q_o(\alpha) \ge Q_1(\alpha) \ge Q(\alpha)$ : Let  $e(\alpha) = Q_1(\alpha) - Q(\alpha)$ . Then  $e(0) = \infty$  and  $e(\infty) = 0$ . If we assume  $de(\alpha)/d\alpha < 0, \forall \alpha \ge 0$ ,

$$de(\alpha)/d\alpha = (2\pi)^{-1/2} [-\alpha^{-2} \exp\{-\alpha^2/2\}$$
(24)  
+(\alpha+1)^{-2} \exp\{-(\alpha+1)^2/2\} + \exp\{-(\alpha+1)^2/2\}] < 0, \text{ or } (25)

$$\{\alpha^2/(\alpha+1)^2\}\exp\{-(\alpha+1/2)\} + \alpha^2\exp\{-(\alpha+1/2)\} < 1.$$
(26)

TABLE II ASEP OF DE-QPSK (*m*=2.5)

$\bar{\gamma}$	Exact	Chiani	Börj.	Karag.	Isukap.	Eq. (23)
0	4.814e-1	5.061e-1	5.000e-1	5.259e-1	4.841e-1	5.532e-1
5	2.133e-1	2.436e-1	2.250e-1	2.292e-1	2.130e-1	2.298e-1
10	4.326e-2	5.122e-2	4.570e-2	4.528e-2	4.270e-2	4.430e-2
15	4.410e-3	5.257e-3	4.653e-3	4.546e-3	4.318e-3	4.412e-3
20	3.098e-4	3.694e-4	3.266e-4	3.174e-4	3.022e-4	3.073e-4
25	1.877e-5	2.239e-5	1.978e-5	1.919e-5	1.828e-5	1.857e-5

TABLE III ASEP OF DE-QPSK (*m*=3.5)

$\bar{\gamma}$	Exact	Chiani	Börj.	Karag.	Isukap.	Eq. (23)
0	4.763e-1	5.045e-1	4.964e-1	5.217e-1	4.793e-1	5.557e-1
5	1.950e-1	2.269e-1	2.063e-1	2.085e-1	1.944e-1	2.128e-1
10	2.889e-2	3.480e-2	3.053e-2	2.977e-2	2.832e-2	2.986e-2
15	1.502e-3	1.802e-3	1.580e-3	1.512e-3	1.452e-3	1.512e-3
20	4.081e-5	4.869e-5	4.285e-5	4.068e-5	3.909e-5	4.069e-5
25	8.398e-7	9.995e-7	8.811e-7	8.342e-7	8.009e-7	8.345e-7

Let  $y_1(\alpha)$  and  $y_2(\alpha)$  be the first and second term of Eq. (26). Then,

$$y(\alpha) = y_1(\alpha) + y_2(\alpha) < 1.$$
 (27)

 $y_1(\alpha)$  and  $y_2(\alpha)$  are unimodal functions for  $\alpha \ge 0$  with their peaks at  $\alpha=1$  and  $\alpha=2$ , respectively. Thus,  $y(\alpha) < y_1(1) + y_2(2) < 1$ . Indeed,  $de(\alpha)/d\alpha < 0$  for  $\alpha \ge 0$ . In result,  $Q_o(\alpha) \ge Q_1(\alpha) \ge Q(\alpha)$ .

#### APPENDIX B

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A limit impulse: To obtain the critical point  $y_*^N$  of  $g(y^N)$ :

$$\frac{dg(y^N)}{dy} = -aNy^{N-1}\exp\{-ay^N\}Ny^{cN-1} + \exp\{-ay^N\}N(cN-1)y^{cN-2} = 0,$$
(28)

or  $y_*^N = \lim_{N \to \infty} y^N = \frac{c}{a}$ . Applying the above result to Eq. (8), we find that

$$\lim_{N \to \infty} g(y_*^N) = \lim_{N \to \infty} \exp\{-c\} N(\frac{c}{a})^c = \infty.$$
(29)

From Eq. (28), we can see that  $dg(y^N)/dy > 0$  for  $0 < y^N < y_*^N$ , and  $dg(y^N)/dy < 0$  for  $y^N > y_*^N$ . Together with Eqs. (9) and (29), we obtain Eq. (10).

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