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Ultrathin BaTiO₃ templates for multiferroic nanostructures

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Abstract. The structural, electronic and dielectric properties of high-quality ultrathin BaTiO₃ films were investigated. The films, which were grown by ozone-assisted molecular beam epitaxy on Nb-doped SrTiO₃(001) substrates and have thicknesses as low as 8 unit cells (u.c.) (3.2 nm), are unreconstructed and atomically smooth with large crystalline terraces. A strain-driven transition to three-dimensional (3D) island formation is observed for films of 13 u.c. thickness (5.2 nm). The high structural quality of the surfaces, together with dielectric properties similar to bulk BaTiO₃ and dominantly TiO₂ surface termination, makes these films suitable templates for the synthesis of high-quality metal-oxide multiferroic heterostructures for the fundamental study and exploitation of magneto-electric effects, such as a recently proposed interface effect in Fe/BaTiO₃ heterostructures based on Fe–Ti interface bonds.

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1. Introduction

The coexistence of ferroelectricity and ferromagnetism in two-phase multiferroic systems comprising metals and oxides can result in interesting and useful phenomena, such as magneto-electric effects [1]. One model two-phase system exhibiting magneto-electric behavior, that is, electric field dependence of magnetization, is epitaxial Fe on BaTiO₃ substrates. The magnetic anisotropy of Fe, observable as the magnitude of the coercive field, depends strongly on BaTiO₃'s electric polarization, as has been shown by Sahoo *et al* [2]. The underlying mechanism here is a lateral strain exerted by the ferroelectric on the ferromagnet and the associated change in the magneto-elastic contribution to the total magnetic anisotropy. A second mechanism, different in nature but also leading to magneto-electric behavior, has been predicted recently for Fe/BaTiO₃ heterostructures [3, 4]. Characteristic of this new effect are modulations of the Fe–Ti bonds at the interface by the piezo-electric distortion of the BaTiO₃, resulting in changes in the effective magnetic moments of Fe and Ti atoms and the anisotropy. A second and related example is magnetic tunnel junctions, which are also based on metal-oxide interfaces [5]. Experimental investigation of any metal-oxide heterostructures depends critically on the quality of the oxide layer, since their performance is extremely sensitive to the chemical and structural properties of the interfaces. The realization of model structures or even devices exhibiting and exploiting the above effects is complicated by the experimental difficulty of growing metal-oxide interfaces of sufficiently high interface quality.

Atomically smooth BaTiO₃ substrates are a requirement for the study of magneto-electric effects in all BaTiO₃-based heterostructures. They can be obtained by annealing bulk samples under ultrahigh vacuum or in hydrogen atmosphere to temperatures around 1000 K [6–8], but both the atomic structure and the surface termination depend critically on the preparation conditions [8–10]. As an alternative to bulk BaTiO₃ substrates, ultrathin films of BaTiO₃ were prepared and studied recently [11–27]. Their structural characterization has so far been limited to reflection high-energy electron diffraction (RHEED) during growth [17, 19, 20, 25, 28], and low-energy electron diffraction (LEED) [12]. It is known from these studies that the films grow layer by layer only below a critical film thickness, which is of the order of 12 unit cells (u.c.) [18]. Characterization by piezo-response force microscopy (PFM) [29] has demonstrated robust ferroelectricity in films with thickness as low as 1 nm (under compressive strain) [30], and layers with thickness as low as 1 u.c. can be ferroelectric in BaTiO₃/SrTiO₃ superlattices (1 u.c. = 0.399 nm). However, a comprehensive *in situ* characterization with surface-sensitive

methods to address the atomistic structure, surface termination and dielectric properties is currently lacking, but is urgently required for the optimization of synthesis strategies and to achieve a significant performance increase in magneto-electric structures.

This paper aims to fill this gap. We present a comprehensive study of the structural, dielectric and electronic properties of BaTiO₃ films, which are only a few u.c. thick and grown on Nb-doped SrTiO₃ substrates, with a combination of local probe methods and electron diffraction and spectroscopy. We will demonstrate that such films are superior to bulk BaTiO₃ substrates regarding the structural quality at the surface, while still showing the dielectric properties of bulk BaTiO₃. We propose that BaTiO₃ thin films are suitable for fundamental research and applications, such as for the study of magneto-electric effects and magneto-tunnel junctions.

2. Experimental procedure

The BaTiO₃ films were grown by molecular beam epitaxy (MBE) by co-evaporation of Ba and Ti from Knudsen cells and using pure ozone as the oxidizing agent. A steady flow of ozone gas was delivered to the growth chamber with the pressure maintained at 2×10^{-6} torr. Prior to film growth, the Nb : SrTiO₃ substrates (0.2% Nb doping) were prepared with a buffered HF dip to obtain a TiO₂-terminated surface. The substrate was heated in ozone to a growth temperature of 650 °C and then cooled in ozone after film deposition. Each BaTiO₃ unit cell was deposited in ~ 50 s followed by 30 s annealing (with Ba and Ti shutters closed). The deposition was monitored using RHEED.

After growth, the samples were transferred through air into separate ultrahigh vacuum systems for further studies with scanning tunneling microscopy (STM), LEED, photoemission spectroscopy (UPS), etc. While immediately after transfer no LEED pattern was observable, extremely sharp (1×1) diffraction patterns could be established by thermal annealing at approximately 650 K under O₂ pressure of 3.75×10^{-7} torr, using an oxygen doser that faces the sample surface at a distance of ~ 5 cm. STM images were taken at 45 K using an Omicron variable temperature STM. Photoelectron spectroscopy (PES) was carried out at a soft x-ray beam line (3A1) at the Pohang Light Source, Korea. The UPS spectra presented here were taken with the sample kept at room temperature using a hemispherical electron energy analyzer with multichannel detector (Scienta). The energy resolution is 0.3 eV, as determined from the measured shape of Fermi edge of a Cu reference sample. The binding energy is measured with respect to the Fermi edge of Cu. PFM measurements were carried out in air [31, 32, 39]. An external ac bias voltage is applied to the PFM tip, and the local piezo-electric response from the ferroelectric layer was measured. By scanning the sample surface a two-dimensional (2D) map of the piezo-response amplitude and phase signals could be generated, providing information on the polarization magnitude and direction. For local hysteresis loop measurement, a dc offset voltage of controlled magnitude is applied to the tip. Dielectric properties, such as the coercive field and remanent polarization, were deduced from local hysteresis loops.

3. Results and discussion

3.1. Surface morphology studies with scanning tunneling microscopy

Our studies focus on BaTiO₃ films of 8 and 13 u.c. thickness. RHEED images taken along the [100] axis of the pristine Nb : SrTiO₃ substrate surface prior to film growth and of the

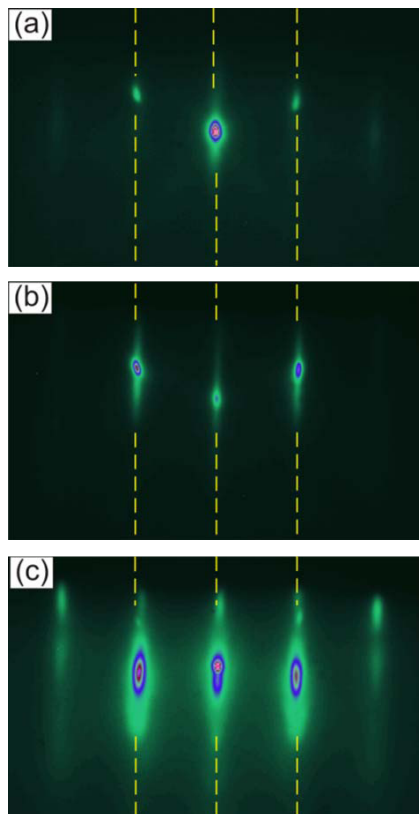


Figure 1. RHEED images of BaTiO₃ grown by MBE, taken at different stages of growth. (a) Pristine Nb-SrTiO₃(100) substrate; (b) after growth of 8 u.c. of BaTiO₃; and (c) after growth of 13 u.c. of BaTiO₃.

8 and 13 u.c. BaTiO₃ films are shown in figures 1(a)–(c), respectively. In particular, for the 8 u.c. film, spots rather than streaks were observed (figure 1(b)), which is consistent with large, crystalline terraces at the surface of the films. Those spots broaden into streaks as the film thickness is increased to 13 u.c., indicating a gradual increase in the surface roughness (figure 1(c)). The films were transferred through air into a separate STM UHV chamber after growth, where they were annealed in oxygen to recover ordered LEED diffraction patterns. A weak LEED pattern became observable at an annealing temperature of $T_a = 150^\circ\text{C}$. With increasing temperature the peak intensity and sharpness improved and the background intensity decreased until high-quality (1×1) diffraction patterns were observed at $T_a \sim 650\text{--}700^\circ\text{C}$ (figure 2). Further annealing was found to be detrimental to the diffraction image quality. STM images were taken on both films immediately after the annealing, at a sample temperature of 45 K. Flat terraces of approximately 100 nm width were found on the 8 u.c. film and the terrace step height corresponds to a single unit cell of BaTiO₃ (figure 2). Despite the atomically smooth terraces, the films are not expected to be globally smooth because of the potential existence of twin boundaries, even though they were not observed in this study. The 13 u.c. films also exhibit atomically flat terraces; however, the terraces are now covered with small islands of average diameter ~ 10 nm. A height histogram analysis reveals that there are on average four open layers (not shown here). This is also consistent with the observation of streaks in the RHEED images and is attributed to the relaxation of epitaxial film strain. It is concluded that crystalline films

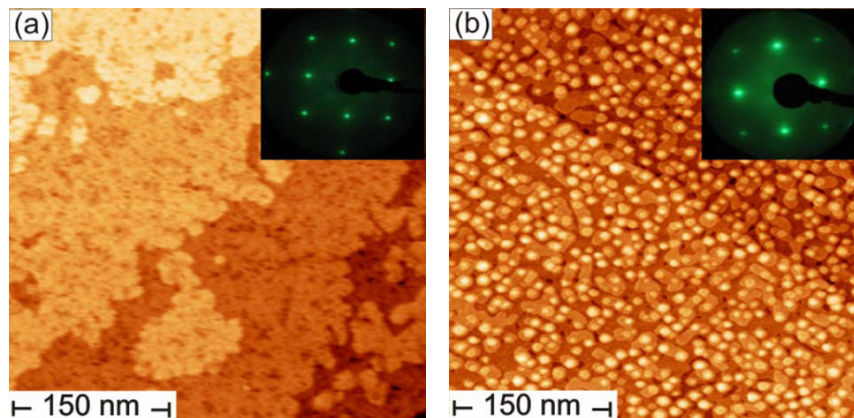


Figure 2. STM images and LEED patterns of BaTiO₃ films on Nb-SrTiO₃. The thickness of BaTiO₃ films is 8 u.c. (a) and 13 u.c. (b). The LEED images were taken at 85 eV.

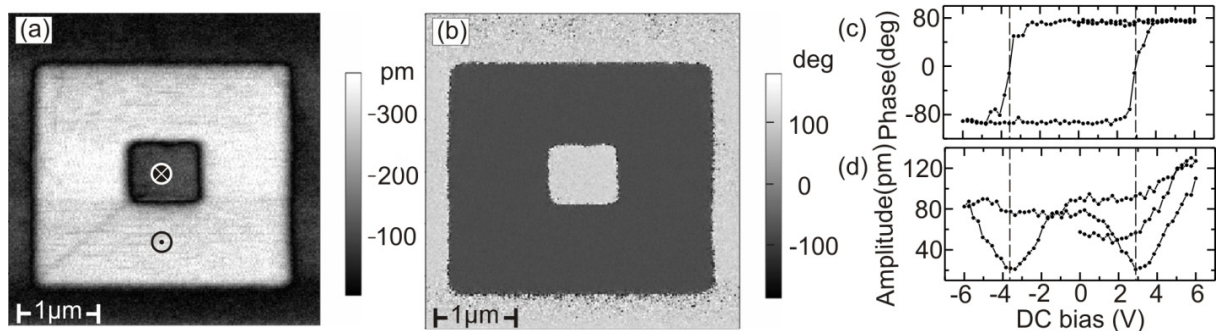


Figure 3. Maps of the remanent PFM amplitude signal (a) and phase signal (b) of a BaTiO₃(8 u.c.)/Nb-SrTiO₃ film, showing areas of opposite ferroelectric polarization, produced by scanning with the tip under ± 4 V dc bias. Hysteresis loops of the PFM in-field phase (c) and amplitude (d) acquired from the same film.

showing no detectable surface reconstruction and exhibiting terraces as wide as 100 nm in the case of the 8 u.c. film can be recovered after sample transfer through air.

3.2. Piezo-response force microscopy and dielectric properties

Resonance-enhanced PFM measurements [31, 32] have been carried out at room temperature on the BaTiO₃ films in air after extraction of the samples from the growth chamber, without further sample treatment. Prior to PFM characterization, bi-domain square-in-square polarization patterns have been generated in the BaTiO₃ films by scanning the surface with an applied dc bias voltage of ± 4 V. The PFM phase and amplitude maps of the resulting patterns, together with local polarization hysteresis loops, are shown in figure 3 for 8 u.c.-thick films. Identical results were obtained from 13 u.c.-thick films and are not shown here, for brevity. The amplitude signal in figure 3(a) is a measure of the polarization magnitude, whereas the phase signal in figure 3(b) shows the polarization direction. The hysteresis loops of the phase and the amplitude signals in panels (c) and (d) confirm the existence of non-zero remanent polarization, which can be

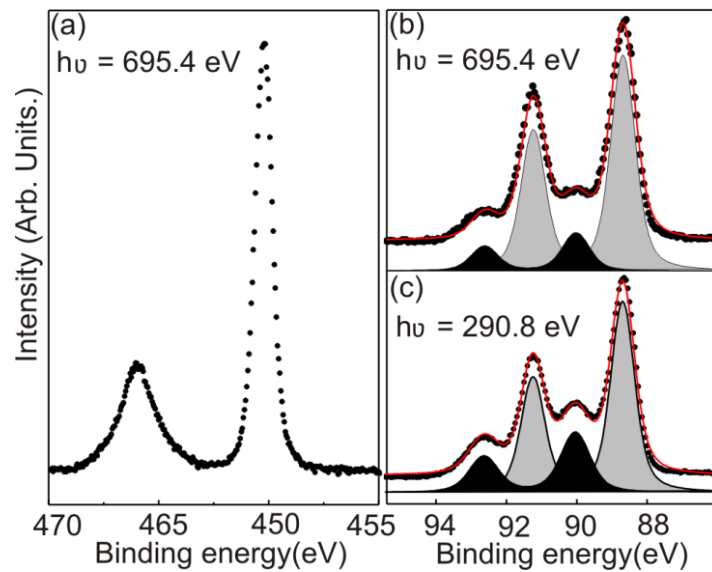


Figure 4. Photoelectron spectra of BaTiO₃(001). Ti 2p spectrum (a) and Ba 4d spectra taken at 695.4 eV (b) and 290.8 eV (c), respectively, are shown. The binding energy is in reference to Fermi energy. Black- and gray-shaded areas in (b, c) represent fits to the bulk and surface peaks of Ba 4d_{5/2} and 4d_{3/2}, respectively.

reversed by applying electric fields larger than those corresponding to the measured coercive bias voltage of approximately 3.5 V. The visible asymmetry in the amplitude signal hysteresis loop for opposite polarization is due to differences in the interfaces on both sides of the BaTiO₃. Film boundaries and surface terminations are known to result in the accumulation of surface charges, which influence the symmetry of the observed PFM loops.

3.3. Photoelectron spectroscopy and surface termination

PES data have been collected on the 13 u.c. BaTiO₃ film, to determine the surface termination and to learn about the electronic structure of such films. Even though the growth of the films was stopped after deposition of the TiO₂ layer, the actual surface termination after transfer through air and annealing in oxygen needs verification. The Ba 4d and the Ti 2p spectra are summarized in figure 4. The Ti 2p spectrum shows two characteristic peaks identified as the Ti 2p_{3/2} and Ti 2p_{1/2} peaks. In these spectra, no surface core level shift (SCLS) is observable. Such a shift is often observed at surfaces due to the reduced coordination of the atoms and the resulting change in charge state there. However, to the best of our knowledge, no SCLS has ever been reported for TiO₂-terminated surfaces. The absence of SCLS for the Ti 2p peaks can be expected, because the effective charge of Ti ions in the TiO₂-terminated BaTiO₃ surface is similar to that of Ti ions in the bulk BaTiO₃, as has been predicted in first-principles calculations [10]. Thus, the absence of SCLS of Ti does not exclude TiO₂ surface termination.

By contrast, the effective charge of Ba ions in a BaO-terminated surface is less than half that of Ba ions in the bulk, thereby producing substantial SCLS, as has already been reported for bulk BaTiO₃(001) [33]. For the present films, we found that the 4d peaks of Ba split into the well-known 4d_{5/2} and 4d_{3/2} peaks, and two additional peaks shifted by ~ 1.5 eV towards higher binding energies with respect to those 4d peaks (figures 4(b) and (c)). Similarly,

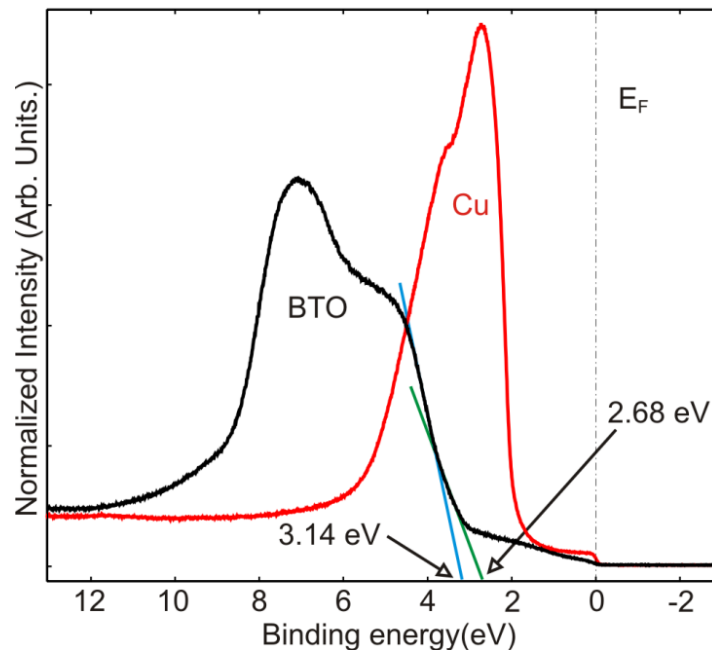


Figure 5. Photoelectron spectra near the Fermi edge of clean BaTiO₃(001) and Cu. The photon energy was 111.04 eV. Extrapolating the band edges gives the band gap (blue and green lines). The non-zero signal in the band gap is due to the existence of mid-gap states (see text).

shifted peaks have been reported for BaO surface layers and are ascribed to SCLS [34]. Ba 4d spectra have been taken with two different photon energies, 695.4 and 290.8 eV. We observe an increase in the ratio of surface to bulk peaks at lower photon energies, where the photoelectrons have smaller kinetic energy and are thus more surface-sensitive. Both observations, the SCLS and the energy dependence of the surface to bulk peak ratio of the Ba 4d states, suggest the existence of BaO on the surface of the films; however, to perform a quantitative analysis of the fractions of BaO- and TiO₂-terminated surface areas requires a detailed analysis of the photoemission peak intensities.

The standard approach for the peak intensity analysis in PES [37] has been adapted here for surfaces containing two atomic species. A detailed discussion of the spectral areas of the Ba 4d bulk and SCLS peak pairs is given in the supplementary material (stacks.iop.org/NJP/13/083037/mmedia). We conclude from this analysis that the BaTiO₃ layer is dominantly TiO₂-terminated, specifically we find the fractions of the TiO₂- and BaO-terminated surface areas to be ~ 70 and $\sim 30\%$, respectively. Uncertainties in our analysis result from estimates of the electron mean free path. We performed further complementary experiments with angular-dependent photoemission spectroscopy on several BaTiO₃ thin films on SrTiO₃ and LaSrMO₃ substrates, fabricated by different groups. Those samples have been prepared at varied annealing temperatures (200–650 °C) and oxygen partial pressures under UHV. All measurements consistently find surfaces dominantly terminated by TiO₂. We conclude that the TiO₂ termination of BaTiO₃(001) films is energetically favored and robust during *in situ* sample preparation.

The valence band spectra of the BaTiO₃ samples are shown in figure 5, together with the spectra of a Cu reference sample for comparison. Most prominently, the valence band, which is

mostly of O 2p character, spans from ~ 3 to 10 eV binding energy. The width of the band gap is determined by extrapolating the band edge. This is shown in figure 5 for two extreme choices of linear extrapolations, resulting in upper and lower limits for the band gap of 2.7–3.2 eV, respectively. In-gap states are observable between the valence band edge and the Fermi edge. Such states are most commonly found to originate from oxygen defects [33, 35, 36] and Ti 3+ ions located in the vicinity of oxygen defects [36]. While first-principles calculations predict that the valence band reaches the Fermi edge [35], we observe instead a band shift to higher binding energy. This shift is similar to the optical band gap of BaTiO₃ of 3.22 eV. It is reasonable to attribute this shift to the pinning of the Fermi level at the bottom of the conduction band to the effective n-type doping, which is a result of the reduction of the sample during the annealing. This indicates that the film has a band gap analogous to bulk BaTiO₃, with in-gap states due to local defects. The latter finding is not surprising for the 13 u.c. film given the observed transition to a 3D island structure.

4. Conclusions

The nanometer-thick films of BaTiO₃ studied in this paper have the following advantages over bulk samples: the formation of large, atomically smooth surface terraces of 1 u.c. height and the absence of surface reconstruction, while still being insulating and exhibiting remanent electric polarization. Our samples show distinctively improved dielectric properties compared with the epitaxial BTO films published so far. Ultrathin BaTiO₃ films are thus expected to become substrates for the fabrication of various model systems for the study of magneto-electric effects and also to facilitate theoretical analysis. An important result of this paper is the demonstration that high-quality surfaces of BaTiO₃ can be recovered after sample transfer through air. This will simplify the sample exchange in most laboratory settings where oxide growth, surface analytics and device fabrication are done in separate UHV systems. We have shown that BaO and TiO₂ surface terminations can coexist and are thus similar in energy, thereby providing experimental evidence to support the mostly theory-based discussion in the literature. Due to the prevalent TiO₂ termination of the surface, these films potentially enable one to synthesize high-quality Fe/BaTiO₃ interfaces that might exhibit the recently proposed magneto-electric effect due to Fe–Ti-bond modulation.

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