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# Analysis of a tritium enhanced water spectrum between 7200 and $7245 \mathrm{~cm}^{-1}$ using new variational calculations 

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#### Abstract

A tritium enhanced water absorption spectrum previously recorded in the $7200-7245 \mathrm{~cm}^{-1}$ region is analysed. Variational calculations for HTO predict absorption to be dominated by the $2 v_{3}$ vibrational band in this region. New assignment are made for HTO based on this line list with a band origin measured to be at $7236.03 \mathrm{~cm}^{-1}$. A calculated $\mathrm{T}_{2} \mathrm{O}$ line list predicts absorption in this region to be below the experimental detection limit despite the large quantity of tritium present. From 170 lines observed 37 known $\mathrm{H}_{2}{ }^{16} \mathrm{O}$ lines are identified and 111 new HTO assignments are made.


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## 1. Introduction

The heavy radioactive tritium $\left({ }^{3} \mathrm{H}\right)$ isotope of hydrogen has a half life of approximately 12.32 years [1]. It forms several isotopologues of water the most common being the singly substituted $\mathrm{HT}^{16} \mathrm{O}$ and the doubly substituted $\mathrm{T}_{2}{ }^{16} \mathrm{O}$ which have trace natural abundances (see Table 1 below). Whilst sharing similar properties to $\mathrm{H}_{2}{ }^{16} \mathrm{O}$, both form corrosive liquids due to self radiolysis and are highly toxic.

As radioactive water isotopologues, these species have been widely used as tracers in life science water transport studies including Refs. [2-5] to name but a few. Their radioactivity has also been put to use in the dating of water based liquids including vintage wines [6]. Rare isotopologue abundances have also often been used to trace atmospheric processes since isotopic variations can be caused by specific atmospheric drivers such as cometary and meteorite deposits at the top of the atmosphere. The spectra of these species are therefore of some importance to these communities, as they provide an opportunity for isotopologue specific detection.

Understanding and controlling tritium and the molecules it forms is also key in nuclear physics. In particular the harmful nature of tritiated water make its detection indispensable. High resolution spectroscopy provides a means of detecting tritiated water species with several benefits, including the potential for in situ real time observation and the ability to make observations without direct sampling.

[^0]Furthermore, documenting the spectra of these species is of value to the theoretical community [7], in their attempt to understand the breakdown of the Born Oppenheimer (BO) approximation. Recorded spectra of these species provide an important tool for evaluating $a b$ initio non BO calculations, designed to model the breakdown of the BO approximation.

Previous high resolution studies of these isotopologues are limited to a handful of works [8-15]. Unlike stable water isotopologues only the fundamentals of these tritiated species have been previously observed at high resolution. Prior to these works Staats et al. [16] recorded near-infrared spectra for these species at significantly lower resolution. Rovibrational assignments are documented by Refs. [12,13,11,10], for the fundamentals and by Refs. [8,9] for pure rotational spectra. As an aid to new assignments new line lists for HTO and $\mathrm{T}_{2} \mathrm{O}$ were hence computed as part of this work.

In this work we analyse a water spectrum recorded by Kobayashi et al. [17], expected to contain a mixture of $\mathrm{H}_{2} \mathrm{O}$, HTO and $\mathrm{T}_{2} \mathrm{O}$. The $7200-7245 \mathrm{~cm}^{-1}$ region studied has not previously been used to probe HTO and $\mathrm{T}_{2} \mathrm{O}$. We present new rotational assignments for HTO, dominated by the $2 v_{3}$ vibrational band, and the new line lists for HTO and $\mathrm{T}_{2} \mathrm{O}$ which we hope to contribute to future analyses.

The paper is structured as follows. The experimental data is described in Section 2. Section 3 presents a theoretical prediction of the observed spectrum according to approximate concentrations, including the new calculations. The techniques used in the analysis are described briefly in Section 4 and the results of the analysis and our new assignments are discussed in Section 5. Data

Table 1
Experimental (Expt.) and natural abundances for the three dominant water isotopologues present in the gas mixture of Kobayashi et al. [20]. The experimental abundances are estimated statistically based on the relative intensities of $\mathrm{H}_{2}{ }^{16} \mathrm{O}$ and $\mathrm{HT}^{16} \mathrm{O}$ assignments.

| Isotopologue | Expt. | Natural |
| :--- | :--- | :--- |
| $\mathrm{H}_{2}{ }^{16} \mathrm{O}$ | 0.01 | 0.997317 |
| $\mathrm{HT}^{16} \mathrm{O}$ | 0.16 | $4.987 \times 10^{-17}$ |
| $\mathrm{~T}_{2}{ }^{16} \mathrm{O}$ | 0.83 | $6.235 \times 10^{-34}$ |

sets arising from this work have been placed in the Supplementary Data.

## 2. The observed spectrum

Kobayashi et al. [17] recorded a frequency modulation nearinfrared spectrum at 296 K scanning the $7200-7245 \mathrm{~cm}^{-1}$ spectral region containing some 170 line positions, with a resolution of $0.02 \mathrm{~cm}^{-1}$ and estimated precision of $0.003 \mathrm{~cm}^{-1}$. The relative intensity limit of detectable lines is thought to be only around $10^{-2}$ times the strongest line, which places some limitation on the analysis. It should be noted that the spectral resolution does not limit the precision of line positions except in the case of blended lines. Although strong lines are well determined, having line widths (FWHM) of the order $0.02 \mathrm{~cm}^{-1}$ the presence of strong air moisture absorption lines reduces the reliability of intensity measurements and limits the sensitivity in those regions.

The gas chamber was filled with tritiated water produced using the technique described in the previous work of Marr et al. [18], and had an unknown ratio of $[\mathrm{H}] /[\mathrm{T}]$. Since deuterium and the rare oxygen isotopologues occur at natural abundances we expect the presence of $\mathrm{HT}^{16} \mathrm{O}, \mathrm{T}_{2}{ }^{16} \mathrm{O}, \mathrm{H}_{2}{ }^{16} \mathrm{O}$ to dominate the spectrum. Table 1 gives the abundance of the different isotopomers in the sample.

## 3. Predicted spectrum

Initial analysis made assignments using the HITRAN line list for $\mathrm{H}_{2}{ }^{16} \mathrm{O}$ and an HTO line list available on-line at the Tomsk web

Table 2
Comparison of observed and calculated band origins for the fundamental vibrational modes of HTO and $\mathrm{T}_{2} \mathrm{O}$. Both observed (Obs.) and calculated (Calc.) positions and their residuals $\Delta$ are in units of $\mathrm{cm}^{-1}$.

| Band | HTO |  |  | $\mathrm{T}_{2} \mathrm{O}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Obs. | Calc. | $\Delta$ | Obs. | Calc. | $\Delta$ |
| $v_{1}$ | 2299.8 | 2299.525 | 0.275 | 2237.2 | 2237.055 | 0.145 |
| $v_{2}$ | 1332.5 | 1332.431 | 0.069 | 995.4 | 995.360 | 0.040 |
| $v_{3}$ | 3716.6 | 3716.791 | -0.191 | 2366.6 | 2366.568 | 0.032 |

address http://spectra.iao.ru/, based on the Partridge and Schwenke (PS) potential energy surface [19]. These line lists have the benefit that they are fully labelled with the usual vibrational normal mode and asymmetric top quantum numbers for water. Indeed, these assignments indicate the spectrum to be dominated by the $2 v_{3}$ band in this region as expected. However the nearby $v_{1}+v_{2}+v_{3}$ and $2 v_{1}+2 v_{2}$ bands also contribute. No assignments of $\mathrm{T}_{2}{ }^{16} \mathrm{O}$ using an equivalent line list were made.

These preliminary assignments indicate a ratio [HTO]/ $\left.\left[\mathrm{H}_{2} \mathrm{O}\right)\right] \approx 20$, based on comparison $I_{\text {obs }} / I_{\text {calc }}$ values for assignments to the two species. Based on purely statistical arguments and ignoring fractionation effects or any radiation induced chemistry this translates to the approximate concentrations for the three species given in Table 1 . The corresponding ratio $[\mathrm{T}] /[\mathrm{H}] \approx 10$.

Due to the lack of documented line lists for the two tritiated species our own variational calculations were undertaken based on the DVR3D program suite [21]. For each of these isotopologues calculations were undertaken up to $J_{\mathrm{MAX}}=15$ using Radau coordinates and an atomic mass for tritium taken from NIST [22]. Full 296 K line lists in the $0-10000 \mathrm{~cm}^{-1}$ range are included in the Supplementary Material. In both cases intensities were calculated using the LTP2011 DMS of Lodi et al. [23].

For HTO the PES of Voronin et al. was used [24], with parameters based on their calculations for HDO [25], which were fully converged up to $20000 \mathrm{~cm}^{-1}$. For $\mathrm{T}_{2} \mathrm{O}$ the PES of Shirin et al. [26] was used and parameters were based on their work for the $\mathrm{D}_{2}{ }^{16} \mathrm{O}$ isotopologue [26] for which convergence in this region was fully tested. Although mass specific non Born-Oppenheimer contributions to the PES were originally introduced, the fundamental band origins were not improved, and the original PES was used without alteration. Fig. 1 presents an overview of our new line lists.

Fundamental band origins were in general agreement with previously observed values [7] for both of these line lists, as seen in Table 2. Partition functions were calculated based on the computed energy levels and were in general agreement with those available at Tomsk on line, once allowance is made for different conventions for the treatment of nuclear spin, with $Q_{\text {нто }}=775.4$ and $Q_{\mathrm{T}_{2} \mathrm{O}}=785.1$ at 296 K , including all nuclear spin factors.

The line lists were used to produce a predicted spectrum where intensities are scaled by the estimated concentrations. This is shown in Fig. 2. Absorption is thus expected to be dominated by HTO and $\mathrm{H}_{2} \mathrm{O}$, with the $2 v_{2}$ band of HTO centred on this region. Although the gas mixture contains a high concentration of $\mathrm{T}_{2} \mathrm{O}$, we do not expect any absorption for this region as it falls between vibrational bands. The strongest $\mathrm{T}_{2} \mathrm{O}$ lines can be seen to be over $10^{4}$ weaker than other strong lines in this region, and this is beyond the sensitivity of the experiment. This is confirmed by failure to assign $\mathrm{T}_{2} \mathrm{O}$ lines to the spectrum using either our own calculations or the Tomsk PS line list (see Table 3).


Fig. 1. Calculated line lists for HTO and $\mathrm{T}_{2} \mathrm{O}$ in the $0-10000 \mathrm{~cm}^{-1}$ range.


Fig. 2. Predicted line positions and intensities for the three isotopologues of water in the $7000-7500 \mathrm{~cm}^{-1}$ region. Intensities are scaled by the estimated experimental abundances given in Table 1. The horizontal arrow indicates the region in question whilst the vertical arrow indicates the difference between maximum HTO and $\mathrm{T}_{2} \mathrm{O}$ intensities in that region. The horizontal dashed line indicates the intensity cut off used in the analysis.

## 4. Method

Our approach for assigning the previously unobserved lines in the spectrum is similar to previous work on a deuterated water spectrum [27], aiming first to identify and remove previously observed $\mathrm{H}_{2}{ }^{16} \mathrm{O}$ lines from the experimental line list. This was done by comparison with the HITRAN database $[28,29]$ which quotes position accuracy of $0.01-0.1 \mathrm{~cm}^{-1}$ in this region. However, as previously mentioned, there is no published data in this region for HTO.

While traditional methods of spectral analysis employ combination differences (CD's) as the major method for assigning spectra, new assignments were made by comparison with the line lists described in Section 3. There are several reasons for this approach. Primarily, the small spectral window available is unlikely to yield CD matches since P, Q and R branch transitions are unlikely to fall close together. Additionally, the lack of published experimental ground state energy levels, as required to calculate accurate CD's reduces the quality of this approach. Furthermore large line densities, the presence of multiple species in the spectrum and limited precision of the experimental data all increase the likelihood of false positive

Table 3
HTO assignments in the $2 v_{3}$ band made using variational calculations performed as part of this work. Positions and residuals are also shown for equivalent assignments made using the calculations available on-line at the Tomsk web address http://spectra.iao.ru/, for the sake of comparison. $v_{\text {obs }}, v_{\text {calc }}, v_{\text {Tomsk }}, \Delta, \Delta_{\text {Tomsk }}$ and $\mathrm{E}^{\prime \prime}$ are given in units of $\mathrm{cm}^{-1}$, whilst $I_{\text {calc }}$ are in units of $\mathrm{cm} /$ molecule. The units of observed intensity are arbitrary. Here the residuals are $\Delta=v_{\text {obs }}-v_{\text {calc }}$ and $\Delta_{\text {Tomsk }}=v_{\text {obs }}-v_{\text {Tomsk }}$. Lines marked with an asterisk are part of blended features.

| $v_{\text {obs }}$ | $v_{\text {calc }}$ | $v_{\text {Tomsk }}$ | $\triangle$ | $\Delta_{\text {Tomsk }}$ | $E^{\prime \prime}$ | $I_{\text {obs }}$ | $I_{\text {calc }}$ | $I_{\text {obs }} / I_{\text {calc }}$ | $J^{\prime}$ | $K_{a}^{\prime}$ | $K_{c}^{\prime}$ | $J^{\prime \prime}$ | $K_{a}^{\prime \prime}$ | $K_{c}^{\prime \prime}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 7200.680 | 7200.771 | 7200.689 | -0.091 | -0.009 | 114.576 | 372.5 | $8.173 \mathrm{E}-22$ | $4.558 \mathrm{E}+23$ | 3 | 1 | 3 | 4 | 0 | 4 |  |
| 7200.887 | 7200.993 | 7200.977 | -0.106 | -0.090 | 591.369 | 64.9 | $2.481 \mathrm{E}-23$ | $2.616 \mathrm{E}+24$ | 8 | 3 | 5 | 7 | 4 | 4 |  |
| 7201.231 | 7201.309 | 7201.239 | -0.078 | -0.008 | 69.188 | 484.7 | $1.640 \mathrm{E}-21$ | $2.956 \mathrm{E}+23$ | 2 | 0 | 2 | 3 | 0 | 3 |  |
| 7201.703 | 7201.787 | 7201.716 | -0.084 | -0.013 | 81.695 | 524.9 | $1.371 \mathrm{E}-21$ | $3.830 \mathrm{E}+23$ | 2 | 1 | 2 | 3 | 1 | 3 |  |
| 7201.904 | 7201.988 | 7201.912 | -0.084 | -0.008 | 202.009 | 306.2 | $9.447 \mathrm{E}-22$ | $3.241 \mathrm{E}+23$ | 5 | 0 | 5 | 5 | 1 | 4 |  |
| 7202.494 | 7202.573 | 7202.531 | -0.079 | -0.037 | 266.677 | 35.3 | $6.853 \mathrm{E}-23$ | $5.147 \mathrm{E}+23$ | 5 | 2 | 4 | 4 | 3 | 1 |  |
| 7203.562 | 7203.701 | 7203.945 | -0.139 | -0.383 | 934.939 | 14.1 | $2.505 \mathrm{E}-23$ | $5.633 \mathrm{E}+23$ | 9 | 5 | 5 | 9 | 5 | 4 | * |
| 7203.562 | 7203.710 | 7203.954 | -0.149 | -0.392 | 934.933 | 14.1 | $2.509 \mathrm{E}-23$ | $5.625 \mathrm{E}+23$ | 9 | 5 | 4 | 9 | 5 | 5 | * |
| 7206.521 | 7206.648 | 7206.537 | -0.127 | -0.016 | 499.924 | 16.2 | 7.070E-23 | $2.295 \mathrm{E}+23$ | 8 | 2 | 7 | 8 | 2 | 6 |  |
| 7208.253 | 7208.348 | 7208.264 | -0.094 | -0.010 | 202.009 | 31.8 | $1.185 \mathrm{E}-22$ | $2.680 \mathrm{E}+23$ | 5 | 1 | 5 | 5 | 1 | 4 |  |
| 7208.438 | 7208.516 | 7208.447 | -0.078 | -0.010 | 140.563 | 460.7 | $1.213 \mathrm{E}-21$ | $3.798 \mathrm{E}+23$ | 4 | 0 | 4 | 4 | 1 | 3 | * |
| 7208.438 | 7208.516 | 7208.457 | -0.078 | -0.020 | 27.619 | 460.7 | $6.160 \mathrm{E}-22$ | $7.479 \mathrm{E}+23$ | 0 | 0 | 0 | 1 | 1 | 1 | * |
| 7209.636 | 7209.776 | 7209.716 | -0.140 | -0.080 | 684.490 | 37.4 | $1.125 \mathrm{E}-22$ | $3.325 \mathrm{E}+23$ | 8 | 4 | 5 | 8 | 4 | 4 |  |
| 7209.855 | 7209.904 | 7209.848 | -0.050 | 0.007 | 684.406 | 141.1 | $1.134 \mathrm{E}-22$ | $1.244 \mathrm{E}+24$ | 8 | 4 | 4 | 8 | 4 | 5 |  |
| 7209.965 | 7210.208 | 7210.104 | -0.244 | -0.140 | 202.009 | 74.1 | 2.178E-22 | $3.401 \mathrm{E}+23$ | 4 | 2 | 3 | 5 | 1 | 4 |  |
| 7210.721 | 7210.879 | 7210.743 | -0.159 | -0.023 | 591.397 | 38.8 | $1.750 \mathrm{E}-22$ | $2.218 \mathrm{E}+23$ | 7 | 4 | 4 | 7 | 4 | 3 |  |
| 7210.791 | 7210.882 | 7210.836 | -0.091 | -0.045 | 591.369 | 26.8 | $1.948 \mathrm{E}-22$ | $1.377 \mathrm{E}+23$ | 7 | 4 | 3 | 7 | 4 | 4 |  |
| 7211.693 | 7211.805 | 7211.771 | -0.113 | -0.078 | 510.028 | 251.9 | 3.220E-22 | $7.822 \mathrm{E}+23$ | 6 | 4 | 3 | 6 | 4 | 2 | * |
| 7211.693 | 7211.815 | 7211.780 | -0.122 | -0.088 | 510.020 | 251.9 | $3.224 \mathrm{E}-22$ | $7.813 \mathrm{E}+23$ | 6 | 4 | 2 | 6 | 4 | 3 | * |
| 7212.317 | 7212.396 | 7212.353 | -0.079 | -0.036 | 325.018 | 12.7 | $8.274 \mathrm{E}-23$ | $1.535 \mathrm{E}+23$ | 6 | 2 | 5 | 5 | 3 | 2 |  |
| 7212.554 | 7212.662 | 7212.638 | -0.108 | -0.084 | 440.342 | 297.0 | 5.183E-22 | $5.731 \mathrm{E}+23$ | 5 | 4 | 2 | 5 | 4 | 1 | * |
| 7212.554 | 7212.664 | 7212.640 | -0.110 | -0.086 | 440.340 | 297.0 | 5.184E-22 | $5.730 \mathrm{E}+23$ | 5 | 4 | 1 | 5 | 4 | 2 | * |
| 7212.795 | 7212.871 | 7212.805 | -0.076 | -0.010 | 34.765 | 141.1 | $1.294 \mathrm{E}-21$ | $1.091 \mathrm{E}+23$ | 1 | 0 | 1 | 2 | 0 | 2 |  |
| 7212.820 | 7212.901 | 7212.833 | -0.081 | -0.014 | 49.279 | 121.4 | $9.038 \mathrm{E}-22$ | $1.343 \mathrm{E}+23$ | 1 | 1 | 1 | 2 | 1 | 2 |  |
| 7213.176 | 7213.254 | 7213.191 | -0.078 | -0.016 | 91.185 | 356.3 | $1.355 \mathrm{E}-21$ | $2.629 \mathrm{E}+23$ | 3 | 0 | 3 | 3 | 1 | 2 |  |
| 7213.191 | 7213.220 | 7213.306 | -0.029 | -0.116 | 394.820 | 3.5 | $1.018 \mathrm{E}-23$ | $3.465 \mathrm{E}+23$ | 3 | 1 | 3 | 2 | 2 | 0 |  |
| 7213.297 | 7213.401 | 7213.387 | -0.104 | -0.089 | 382.311 | 601.1 | $8.191 \mathrm{E}-22$ | $7.339 \mathrm{E}+23$ | 4 | 4 | 1 | 4 | 4 | 0 |  |
| 7213.297 | 7213.401 | 7213.387 | -0.104 | -0.089 | 382.310 | 601.1 | $8.191 \mathrm{E}-22$ | $7.339 \mathrm{E}+23$ | 4 | 4 | 0 | 4 | 4 | 1 |  |
| 7213.565 | 7213.728 | 7213.595 | -0.163 | -0.031 | 402.088 | 28.9 | 7.303E-23 | $3.961 \mathrm{E}+23$ | 6 | 3 | 4 | 7 | 2 | 5 |  |
| 7214.008 | 7214.162 | 7214.051 | -0.154 | -0.043 | 678.615 | 67.0 | $5.777 \mathrm{E}-23$ | $1.160 \mathrm{E}+24$ | 9 | 3 | 7 | 9 | 3 | 6 |  |
| 7214.202 | 7214.294 | 7214.218 | -0.093 | -0.016 | 69.188 | 261.1 | $6.172 \mathrm{E}-22$ | $4.230 \mathrm{E}+23$ | 2 | 1 | 2 | 3 | 0 | 3 |  |
| 7214.371 | 7214.458 | 7214.442 | -0.087 | -0.071 | 684.406 | 34.6 | 2.004E-23 | $1.725 \mathrm{E}+24$ | 9 | 3 | 6 | 8 | 4 | 5 | * |
| 7214.371 | 7214.486 | 7214.389 | -0.116 | -0.018 | 402.088 | 34.6 | $1.267 \mathrm{E}-22$ | $2.729 \mathrm{E}+23$ | 7 | 2 | 6 | 7 | 2 | 5 | * |
| 7215.795 | 7215.651 | 7215.530 | 0.144 | 0.265 | 678.615 | 16.9 | $2.450 \mathrm{E}-23$ | $6.912 \mathrm{E}+23$ | 8 | 4 | 5 | 9 | 3 | 6 | * |
| 7216.391 | 7216.467 | 7216.408 | -0.076 | -0.017 | 54.029 | 558.8 | $1.267 \mathrm{E}-21$ | $4.412 \mathrm{E}+23$ | 2 | 0 | 2 | 2 | 1 | 1 |  |
| 7217.022 | 7217.113 | 7217.037 | -0.092 | -0.015 | 140.563 | 43.7 | $1.881 \mathrm{E}-22$ | $2.326 \mathrm{E}+23$ | 4 | 1 | 4 | 4 | 1 | 3 |  |
| 7218.274 | 7218.433 | 7218.376 | -0.159 | -0.103 | 29.203 | 214.5 | $9.066 \mathrm{E}-22$ | $2.366 \mathrm{E}+23$ | 1 | 0 | 1 | 1 | 1 | 0 |  |
| 7218.342 | 7218.508 | 7218.452 | -0.166 | -0.110 | 789.360 | 217.3 | $1.662 \mathrm{E}-23$ | $1.307 \mathrm{E}+25$ | 10 | 3 | 8 | 9 | 4 | 5 |  |
| 7219.091 | 7219.223 | 7219.132 | -0.132 | -0.041 | 477.420 | 77.6 | $1.928 \mathrm{E}-22$ | $4.025 \mathrm{E}+23$ | 7 | 3 | 5 | 7 | 3 | 4 |  |
| 7219.502 | 7219.728 | 7219.619 | -0.226 | -0.117 | 674.583 | 75.5 | $2.601 \mathrm{E}-23$ | $2.902 \mathrm{E}+24$ | 8 | 4 | 4 | 9 | 3 | 7 |  |
| 7220.343 | 7220.418 | 7220.368 | -0.075 | -0.026 | 137.258 | 71.3 | $1.059 \mathrm{E}-22$ | $6.727 \mathrm{E}+23$ | 4 | 1 | 4 | 3 | 2 | 1 |  |
| 7220.384 | 7220.487 | 7220.401 | -0.104 | -0.018 | 316.927 | 137.6 | $2.178 \mathrm{E}-22$ | $6.316 \mathrm{E}+23$ | 6 | 2 | 5 | 6 | 2 | 4 |  |

Table 3 (continued)

| $v_{\text {obs }}$ | $v_{\text {calc }}$ | $v_{\text {Tomsk }}$ | $\Delta$ | $\Delta_{\text {Tomsk }}$ | $E^{\prime \prime}$ | $I_{\text {obs }}$ | $I_{\text {calc }}$ | $I_{\text {obs }} / I_{\text {calc }}$ | $J^{\prime}$ | $K_{a}^{\prime}$ | $K_{c}^{\prime}$ | $J^{\prime \prime}$ | $K_{a}^{\prime \prime}$ | $K_{c}^{\prime \prime}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 7220.468 | 7220.598 | 7220.514 | -0.130 | -0.047 | 395.218 | 156.6 | $3.289 \mathrm{E}-22$ | $4.762 \mathrm{E}+23$ | 6 | 3 | 4 | 6 | 3 | 3 |  |
| 7221.277 | 7221.356 | 7221.308 | -0.079 | -0.031 | 395.218 | 21.2 | 7.858E-23 | $2.693 \mathrm{E}+23$ | 7 | 2 | 6 | 6 | 3 | 3 |  |
| 7221.413 | 7221.545 | 7221.467 | -0.132 | -0.054 | 325.018 | 207.4 | $5.415 \mathrm{E}-22$ | $3.831 \mathrm{E}+23$ | 5 | 3 | 3 | 5 | 3 | 2 | * |
| 7221.515 | 7221.650 | 7221.573 | -0.135 | -0.058 | 394.820 | 121.4 | $3.215 \mathrm{E}-22$ | $3.775 \mathrm{E}+23$ | 6 | 3 | 3 | 6 | 3 | 4 |  |
| 7221.675 | 7221.815 | 7221.742 | -0.140 | -0.067 | 476.443 | 50.1 | $1.803 \mathrm{E}-22$ | $2.779 \mathrm{E}+23$ | 7 | 3 | 4 | 7 | 3 | 5 |  |
| 7221.762 | 7221.897 | 7221.821 | -0.135 | -0.060 | 324.884 | 201.8 | $5.382 \mathrm{E}-22$ | $3.750 \mathrm{E}+23$ | 5 | 3 | 2 | 5 | 3 | 3 |  |
| 7222.103 | 7222.224 | 7222.151 | -0.121 | -0.048 | 266.677 | 318.9 | $8.711 \mathrm{E}-22$ | $3.661 \mathrm{E}+23$ | 4 | 3 | 2 | 4 | 3 | 1 |  |
| 7222.198 | 7222.312 | 7222.240 | -0.115 | -0.042 | 266.643 | 266.6 | $8.700 \mathrm{E}-22$ | $3.065 \mathrm{E}+23$ | 4 | 3 | 1 | 4 | 3 | 2 |  |
| 7222.509 | 7222.644 | 7222.579 | -0.136 | -0.070 | 569.717 | 19.8 | $9.533 \mathrm{E}-23$ | $2.072 \mathrm{E}+23$ | 8 | 3 | 5 | 8 | 3 | 6 |  |
| 7222.614 | 7222.729 | 7222.660 | -0.114 | -0.045 | 220.093 | 588.4 | $1.399 \mathrm{E}-21$ | $4.207 \mathrm{E}+23$ | 3 | 3 | 1 | 3 | 3 | 0 | * |
| 7222.614 | 7222.741 | 7222.672 | -0.127 | -0.058 | 220.089 | 588.4 | $1.398 \mathrm{E}-21$ | $4.208 \mathrm{E}+23$ | 3 | 3 | 0 | 3 | 3 | 1 | * |
| 7223.527 | 7223.594 | 7223.553 | -0.067 | -0.026 | 101.881 | 14.8 | 8.626E-23 | $1.718 \mathrm{E}+23$ | 3 | 1 | 2 | 2 | 2 | 1 |  |
| 7224.076 | 7224.162 | 7224.092 | -0.086 | -0.015 | 91.185 | 152.4 | $3.001 \mathrm{E}-22$ | $5.079 \mathrm{E}+23$ | 3 | 1 | 3 | 3 | 1 | 2 |  |
| 7224.453 | 7224.508 | 7224.446 | -0.055 | 0.007 | 11.627 | 382.4 | 7.236E-22 | $5.285 \mathrm{E}+23$ | 0 | 0 | 0 | 1 | 0 | 1 |  |
| 7224.674 | 7224.764 | 7224.685 | -0.090 | -0.011 | 244.486 | 132.6 | $3.615 \mathrm{E}-22$ | $3.669 \mathrm{E}+23$ | 5 | 2 | 4 | 5 | 2 | 3 |  |
| 7225.615 | 7225.689 | 7225.635 | -0.074 | -0.020 | 184.668 | 38.1 | $1.097 \mathrm{E}-22$ | $3.473 \mathrm{E}+23$ | 5 | 1 | 5 | 4 | 2 | 2 |  |
| 7225.759 | 7225.874 | 7225.777 | -0.116 | -0.018 | 140.563 | 34.6 | $1.670 \mathrm{E}-22$ | $2.070 \mathrm{E}+23$ | 3 | 2 | 2 | 4 | 1 | 3 |  |
| 7227.331 | 7227.415 | 7227.343 | -0.084 | -0.012 | 34.765 | 83.3 | $3.331 \mathrm{E}-22$ | $2.500 \mathrm{E}+23$ | 1 | 1 | 1 | 2 | 0 | 2 |  |
| 7227.453 | 7227.549 | 7227.476 | -0.097 | -0.023 | 184.668 | 173.6 | $5.848 \mathrm{E}-22$ | $2.968 \mathrm{E}+23$ | 4 | 2 | 3 | 4 | 2 | 2 |  |
| 7228.703 | 7228.784 | 7228.724 | -0.081 | -0.021 | 244.486 | 33.2 | $8.711 \mathrm{E}-23$ | $3.807 \mathrm{E}+23$ | 6 | 1 | 6 | 5 | 2 | 3 |  |
| 7229.201 | 7229.342 | 7229.436 | -0.141 | -0.235 | 317.105 | 39.5 | $3.825 \mathrm{E}-23$ | $1.033 \mathrm{E}+24$ | 6 | 2 | 4 | 7 | 1 | 7 |  |
| 7229.373 | 7229.453 | 7229.387 | -0.080 | -0.015 | 54.029 | 145.3 | $5.021 \mathrm{E}-22$ | $2.895 \mathrm{E}+23$ | 2 | 1 | 2 | 2 | 1 | 1 |  |
| 7229.484 | 7229.636 | 7229.515 | -0.152 | -0.031 | 316.927 | 16.2 | 7.499E-23 | $2.164 \mathrm{E}+23$ | 5 | 3 | 3 | 6 | 2 | 4 |  |
| 7229.922 | 7230.017 | 7229.949 | -0.096 | -0.027 | 101.992 | 605.4 | $1.574 \mathrm{E}-21$ | $3.845 \mathrm{E}+23$ | 2 | 2 | 1 | 2 | 2 | 0 |  |
| 7230.168 | 7230.266 | 7230.198 | -0.098 | -0.030 | 101.881 | 396.5 | $1.574 \mathrm{E}-21$ | $2.520 \mathrm{E}+23$ | 2 | 2 | 0 | 2 | 2 | 1 |  |
| 7231.072 | 7231.163 | 7231.097 | -0.091 | -0.025 | 183.039 | 103.0 | $5.832 \mathrm{E}-22$ | $1.766 \mathrm{E}+23$ | 4 | 2 | 2 | 4 | 2 | 3 |  |
| 7231.664 | 7231.776 | 7231.669 | -0.113 | -0.006 | 495.022 | 12.0 | 8.894E-23 | $1.349 \mathrm{E}+23$ | 11 | 3 | 8 | 11 | 3 | 9 | * |
| 7232.772 | 7232.832 | 7232.789 | -0.060 | -0.017 | 240.794 | 87.5 | $3.631 \mathrm{E}-22$ | $2.409 \mathrm{E}+23$ | 5 | 2 | 3 | 5 | 2 | 4 |  |
| 7232.895 | 7232.976 | 7232.914 | -0.081 | -0.019 | 29.203 | 342.9 | $1.005 \mathrm{E}-21$ | $3.411 \mathrm{E}+23$ | 1 | 1 | 1 | 1 | 1 | 0 |  |
| 7235.643 | 7235.746 | 7235.662 | -0.103 | -0.019 | 178.456 | 16.9 | $9.661 \mathrm{E}-23$ | $1.753 \mathrm{E}+23$ | 4 | 2 | 2 | 5 | 1 | 5 |  |
| 7236.071 | 7236.113 | 7236.108 | -0.042 | -0.036 | 394.820 | 27.5 | $6.728 \mathrm{E}-23$ | $4.090 \mathrm{E}+23$ | 7 | 2 | 5 | 6 | 3 | 4 |  |
| 7236.181 | 7236.252 | 7236.192 | -0.071 | -0.011 | 27.619 | 421.2 | $1.003 \mathrm{E}-21$ | $4.201 \mathrm{E}+23$ | 1 | 1 | 0 | 1 | 1 | 1 |  |
| 7236.462 | 7236.583 | 7236.693 | -0.121 | -0.232 | 309.865 | 151.7 | $1.151 \mathrm{E}-22$ | $1.318 \mathrm{E}+24$ | 6 | 2 | 4 | 6 | 2 | 5 | * |
| 7236.546 | 7236.643 | 7236.568 | -0.097 | -0.023 | 1034.761 | 20.5 | $6.455 \mathrm{E}-24$ | $3.170 \mathrm{E}+24$ | 12 | 3 | 10 | 11 | 4 | 7 |  |
| 7236.770 | 7236.916 | 7236.806 | -0.146 | -0.036 | 309.865 | 50.8 | $6.308 \mathrm{E}-23$ | $8.053 \mathrm{E}+23$ | 5 | 3 | 2 | 6 | 2 | 5 | * |
| 7237.737 | 7237.795 | 7237.757 | -0.058 | -0.020 | 136.706 | 54.3 | $1.744 \mathrm{E}-22$ | $3.115 \mathrm{E}+23$ | 4 | 1 | 3 | 3 | 2 | 2 |  |
| 7239.198 | 7239.275 | 7239.216 | -0.077 | -0.018 | 49.279 | 299.2 | 4.980E-22 | $6.007 \mathrm{E}+23$ | 2 | 1 | 1 | 2 | 1 | 2 |  |
| 7239.586 | 7239.746 | 7239.666 | -0.160 | -0.080 | 678.615 | 37.4 | $2.575 \mathrm{E}-23$ | $1.452 \mathrm{E}+24$ | 10 | 2 | 9 | 9 | 3 | 6 | * |
| 7240.713 | 7240.806 | 7240.764 | -0.093 | -0.052 | 390.128 | 14.1 | $1.030 \mathrm{E}-22$ | $1.370 \mathrm{E}+23$ | 7 | 2 | 5 | 7 | 2 | 6 |  |
| 7240.738 | 7240.825 | 7240.734 | -0.087 | 0.004 | 91.185 | 24.0 | $8.683 \mathrm{E}-23$ | $2.763 \mathrm{E}+23$ | 2 | 2 | 1 | 3 | 1 | 2 |  |
| 7242.810 | 7242.878 | 7242.828 | -0.067 | -0.018 | 27.619 | 28.0 | $3.449 \mathrm{E}-22$ | $8.122 \mathrm{E}+22$ | 2 | 0 | 2 | 1 | 1 | 1 |  |

CD's. Lack of resolution in general exacerbates this problem, since estimated precisions are lower for blended lines. Conversely, the accuracy of computed line lists allows assignments to be made confidently by comparison with calculations.

Assignments are based upon a direct matching routine which compares positions of observed peaks and known or calculated line positions within a frequency interval, with the possibility of a constant frequency shift being applied. The chance that lines match coincidentally is deemed to be small for the line density of this spectrum and the small frequency intervals employed. The presence of a large number of matches found within a small frequency interval was used to determine an appropriate shift for each line list and the matching frequency range was then trimmed correspondingly to minimise the chance of coincidental matches. All assignments have also been verified using the Tomsk PS based line list described above, which provide the usual vibrational and rotational labels for water.

Assignments were also required to have consistent intensity agreement. Computed $I_{\text {obs }} / I_{\text {calc }}$ ratios were used not only to determine approximate concentrations as discussed above, but also for validation of assignments, for which intensities should show consistent behaviour. This was limited by the uncertainties in the experimental relative intensity data.

Assignments have been checked using CD's, based on calculated energy levels, where possible. However, for the small numbers of lines involved few energy levels are associated to more than one observed line. As a result only 25 energy levels can be confidently derived. These are tabulated in Table 7.

## 5. Results and discussion

Initially, $37 \mathrm{H}_{2}{ }^{16} \mathrm{O}$ features were identified using HITRAN out of 68 HITRAN lines in the region above the intensity threshold of $1 \times 10^{24} \mathrm{~cm} /$ molecule. Assignments were made within a $0.01 \mathrm{~cm}^{-1}$ window of the observed position except for two outliers, with errors less than $0.025 \mathrm{~cm}^{-1}$. The small average shift of $0.0005 \mathrm{~cm}^{-1}$ observed is deemed to be below the experimental accuracy of this data. These lines were removed from the analysis, and are present in the Supplementary Material but will not be discussed further here.

In total 111 new features were assigned to HTO using our calculations, including 38 blended features, from the 183 lines with calculated intensity above the threshold. These belonged to three vibrational bands. Both blended features and outliers in terms of position and intensity agreement must be considered less reliable.

The new assignments were dominated by the strongest band, the $2 v_{3}$ band, with the 83 assignments made including 20 blended features. An average band shift of $-0.103 \mathrm{~cm}^{-1}$ was observed between the calculated and observed positions, and lines were assigned within $0.1 \mathrm{~cm}^{-1}$ of this value, except for one outlier with a $0.23 \mathrm{~cm}^{-1}$ residual. Table 3 shows these assigned lines; blended lines are clearly indicated.

A further 19 assignments were made to the next strongest band, the $v_{1}+v_{2}+v_{3}$ band, and are tabulated in Table 4. These show a large spread around an average shift of $0.043 \mathrm{~cm}^{-1}$, with the standard deviation of residuals at $0.124 \mathrm{~cm}^{-1}$ and as such must be considered less reliable. This band also contains 7 blended features.

Table 4
HTO assignments in the $v_{1}+v_{2}+v_{3}$ band made using variational calculations performed as part of this work. Positions and residuals are also shown for equivalent assignments made using the calculations available on-line at the Tomsk web address http://spectra.iao.ru/, for the sake of comparison. $v_{\text {obss }}, v_{\text {calc }}, v_{\text {Tomsk }}, \Delta, \Delta_{\text {Tomsk }}$ and $E^{\prime \prime}$ are given in units of $\mathrm{cm}^{-1}$, whilst $I_{\text {calc }}$ are in units of $\mathrm{cm} /$ molecule. The units of observed intensity are arbitrary. Here the residuals are $\Delta=v_{\text {obs }}-v_{\text {calc }}$ and $\Delta_{\text {Tomsk }}=v_{\text {obs }}-v_{\text {Tomsk }}$. Lines marked with an asterisk are part of blended features.

| $v_{\text {obs }}$ | $v_{\text {calc }}$ | $v_{\text {Tomsk }}$ | $\Delta$ | $\Delta_{\text {Tomsk }}$ | $E^{\prime \prime}$ | $I_{o b s}$ | $I_{\text {calc }}$ | $I_{\text {obs }} / I_{\text {calc }}$ | $J^{\prime}$ | $K_{a}^{\prime}$ | $K_{c}^{\prime}$ | $J^{\prime \prime}$ | $K_{a}^{\prime \prime}$ | $K_{c}^{\prime \prime}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 7201.107 | 7201.072 | 7201.176 | 0.035 | -0.069 | 820.175 | 21.2 | $1.161 \mathrm{E}-23$ | $1.824 \mathrm{E}+24$ | 10 | 2 | 9 | 11 | 2 | 10 |  |
| 7202.705 | 7202.728 | 7202.842 | -0.023 | -0.137 | 716.932 | 522.1 | $3.289 \mathrm{E}-23$ | $1.588 \mathrm{E}+25$ | 10 | 0 | 10 | 11 | 1 | 11 |  |
| 7203.562 | 7203.451 | 7203.564 | 0.111 | -0.002 | 716.209 | 14.1 | $4.319 \mathrm{E}-23$ | $3.267 \mathrm{E}+23$ | 10 | 0 | 10 | 11 | 0 | 11 | * |
| 7206.162 | 7206.215 | 7206.439 | -0.053 | -0.277 | 266.677 | 28.9 | $6.209 \mathrm{E}-23$ | $4.659 \mathrm{E}+23$ | 3 | 2 | 2 | 4 | 3 | 1 |  |
| 7209.925 | 7209.980 | 7210.128 | -0.056 | -0.203 | 481.445 | 45.9 | $5.253 \mathrm{E}-23$ | $8.731 \mathrm{E}+23$ | 7 | 1 | 6 | 8 | 2 | 7 |  |
| 7215.795 | 7215.637 | 7215.832 | 0.158 | -0.037 | 905.592 | 16.9 | 8.342E-24 | $2.030 \mathrm{E}+24$ | 9 | 4 | 6 | 10 | 4 | 7 | * |
| 7218.645 | 7218.524 | 7218.722 | 0.121 | -0.077 | 600.711 | 19.1 | $4.930 \mathrm{E}-23$ | $3.864 \mathrm{E}+23$ | 9 | 1 | 9 | 10 | 0 | 10 |  |
| 7218.730 | 7218.648 | 7218.880 | 0.082 | -0.150 | 220.093 | 62.8 | $6.951 \mathrm{E}-23$ | $9.034 \mathrm{E}+23$ | 2 | 2 | 1 | 3 | 3 | 0 |  |
| 7218.813 | 7218.781 | 7219.013 | 0.032 | -0.200 | 220.089 | 37.4 | $6.956 \mathrm{E}-23$ | $5.376 \mathrm{E}+23$ | 2 | 2 | 0 | 3 | 3 | 1 |  |
| 7220.498 | 7220.388 | 7220.646 | 0.110 | -0.149 | 1050.746 | 19.1 | $2.422 \mathrm{E}-24$ | $7.867 \mathrm{E}+24$ | 9 | 5 | 4 | 10 | 5 | 5 | * |
| 7220.498 | 7220.399 | 7220.657 | 0.099 | -0.160 | 1050.730 | 19.1 | $2.427 \mathrm{E}-24$ | $7.849 \mathrm{E}+24$ | 9 | 5 | 5 | 10 | 5 | 6 | * |
| 7221.413 | 7221.290 | 7221.444 | 0.123 | -0.031 | 610.265 | 207.4 | $4.938 \mathrm{E}-23$ | $4.201 \mathrm{E}+24$ | 8 | 2 | 6 | 9 | 2 | 7 | * |
| 7227.037 | 7227.147 | 7227.260 | -0.110 | -0.223 | 496.822 | 17.6 | $9.153 \mathrm{E}-23$ | $1.927 \mathrm{E}+23$ | 8 | 0 | 8 | 9 | 1 | 9 |  |
| 7231.664 | 7231.511 | 7231.697 | 0.153 | -0.034 | 184.668 | 12.0 | $5.824 \mathrm{E}-23$ | $2.059 \mathrm{E}+23$ | 3 | 1 | 3 | 4 | 2 | 2 | * |
| 7234.865 | 7234.494 | 7234.680 | 0.371 | 0.185 | 240.794 | 64.9 | $7.263 \mathrm{E}-23$ | $8.938 \mathrm{E}+23$ | 4 | 1 | 3 | 5 | 2 | 4 |  |
| 7239.133 | 7238.870 | 7238.975 | 0.263 | 0.158 | 680.789 | 60.7 | 8.084E-24 | $7.506 \mathrm{E}+24$ | 10 | 0 | 10 | 10 | 1 | 9 |  |
| 7239.586 | 7239.674 | 7239.844 | -0.088 | -0.258 | 481.445 | 37.4 | $7.499 \mathrm{E}-23$ | $4.987 \mathrm{E}+23$ | 7 | 2 | 6 | 8 | 2 | 7 | * |
| 7241.878 | 7241.990 | 7242.133 | -0.112 | -0.255 | 399.130 | 28.9 | $1.476 \mathrm{E}-22$ | $1.960 \mathrm{E}+23$ | 7 | 0 | 7 | 8 | 0 | 8 |  |

Table 5
HTO assignments in the $2 v_{1}+2 v_{2}$ band made using variational calculations performed as part of this work. Positions and residuals are also shown for equivalent assignments made using the calculations available on-line at the Tomsk web address http://spectra.iao.ru/, for the sake of comparison. $v_{\text {obss }}, v_{\text {calc, }}, v_{\text {Tomsk }}, \Delta, \Delta_{\text {Tomsk }}$ and $E^{\prime \prime}$ are given in units of $\mathrm{cm}^{-1}$, whilst $I_{\text {calc }}$ are in units of $\mathrm{cm} /$ molecule. The units of observed intensity are arbitrary. Here the residuals are $\Delta=v_{\text {obs }}-v_{\text {calc }}$ and $\Delta_{\text {Tomsk }}=v_{\text {obs }}-v_{\text {Tomsk. }}$. Lines marked with an asterisk are part of blended features.

| $v_{\text {obs }}$ | $v_{\text {calc }}$ | $v_{\text {Tomsk }}$ | $\Delta$ | $\Lambda_{\text {Tomsk }}$ | $E^{\prime \prime}$ | $I_{\text {obs }}$ | $I_{\text {calc }}$ | $I_{\text {obs }} / I_{\text {calc }}$ | $J^{\prime}$ | $K_{a}^{\prime}$ | $K_{c}^{\prime}$ | $J^{\prime \prime}$ | $K_{a}^{\prime \prime}$ | $K_{c}^{\prime \prime}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 7207.452 | 7207.249 | 7207.557 | 0.203 | -0.105 | 738.177 | 106.5 | $7.466 \mathrm{E}-23$ | $1.427 \mathrm{E}+24$ | 7 | 5 | 3 | 7 | 5 | 2 |
| 7207.452 | 7207.249 | 7207.557 | 0.202 | -0.105 | 738.176 | 106.5 | $7.467 \mathrm{E}-23$ | $1.427 \mathrm{E}+24$ | 7 | 5 | 2 | 7 | 5 | 3 |
| 7209.023 | 7208.794 | 7209.128 | 0.228 | -0.105 | 657.196 | 52.2 | $1.221 \mathrm{E}-22$ | $4.274 \mathrm{E}+23$ | 6 | 5 | 2 | 6 | 5 | 1 |
| 7209.023 | 7208.794 | 7209.128 | 0.228 | -0.105 | 657.196 | 52.2 | $1.222 \mathrm{E}-22$ | $4.274 \mathrm{E}+23$ | 6 | 5 | 1 | 6 | 5 | 2 |
| 7210.434 | 7210.156 | 7210.511 | 0.278 | -0.077 | 587.803 | 36.0 | $1.947 \mathrm{E}-22$ | $1.848 \mathrm{E}+23$ | 5 | 5 | 0 | 5 | 5 | 1 |
| 7210.434 | 7210.156 | 7210.511 | 0.278 | -0.077 | 587.803 | 36.0 | $1.947 \mathrm{E}-22$ | $1.848 \mathrm{E}+23$ | 5 | 5 | 1 | 5 | 5 | 0 |
| 7221.206 | 7220.982 | 7221.523 | 0.224 | -0.317 | 360.209 | 7.8 | $1.061 \mathrm{E}-23$ | $7.315 \mathrm{E}+23$ | 8 | 1 | 7 | 7 | 1 | 6 |
| 7236.770 | 7236.681 | 7237.237 | 0.089 | -0.466 | 440.340 | 50.8 | $2.859 \mathrm{E}-23$ | $1.777 \mathrm{E}+24$ | 6 | 4 | 3 | 5 | 4 | 2 |
| 7236.770 | 7236.689 | 7237.244 | 0.082 | -0.474 | 440.342 | 50.8 | $2.859 \mathrm{E}-23$ | $1.777 \mathrm{E}+24$ | 6 | 4 | 2 | 5 | 4 | 1 |

Table 6
Estimated HTO band origins for $2 v_{3}, 2 v_{1}+2 v_{2}$, and $v_{1}+v_{2}+v_{3}$ derived from this work, using calculated band origins and the average observed band shifts $\bar{\Delta}$, alongside the number and standard deviation $\sigma$ of assignments in each band, all in units of $\mathrm{cm}^{-1}$.

| Band | Origin |  | $\sigma$ | No. assigned | No. blended |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $2 v_{3}$ | 7236.07 | -0.103 | 0.045 | 83 | 20 |
| $2 v_{1}+2 v_{2}$ | 7129.67 | 0.201 | 0.059 | 9 | 8 |
| $v_{1}+v_{2}+v_{3}$ | 7335.85 | 0.043 | 0.124 | 19 | 7 |

Tentative assignments to the $2 v_{1}+2 v_{2}$ band were also made and are shown in Table 5. These assignments showed an average shift of $0.201 \mathrm{~cm}^{-1}$ with a relatively small standard deviation of $0.059 \mathrm{~cm}^{-1}$. Nearly all of these assignments belong to blended features and as such must again be considered less reliable.

The remaining 41 unassigned features are all weak with observed relative intensity less than 0.1 . Such lines become hard to assign when low experimental accuracies combine with large band shifts for the calculated positions. While it is possible higher values of J in the computed line list could account for some of these features, it is thought to be unlikely for $J>15$. Neither is it thought, as explained above, that any features belong to the $\mathrm{T}_{2} \mathrm{O}$ species, for which similar analysis was unsuccessfully attempted. Any further assignments to this spectrum using the presently available data must at best be considered tentative.

Table 7
Energy levels and standard deviations $\sigma$ derived from experimental positions and calculated lower energies using CD's. $N$ denotes the number of transitions sharing the same upper state in each case. Please note for $N=2, \sigma$ corresponds to half the difference between the two derived energy values.

| $v_{1}$ | $v_{2}$ | $v_{3}$ | $J$ | $K_{a}$ | $K_{c}$ | $E\left(\mathrm{~cm}^{-1}\right)$ | $\sigma\left(\mathrm{cm}^{-1}\right)$ | $N$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0 | 0 | 2 | 0 | 0 | 0 | 7236.068 | 0.011 | 2 |
| 0 | 0 | 2 | 1 | 0 | 1 | 7247.518 | 0.042 | 2 |
| 0 | 0 | 2 | 1 | 1 | 1 | 7262.097 | 0.002 | 3 |
| 0 | 0 | 2 | 1 | 1 | 0 | 7263.793 | 0.007 | 2 |
| 0 | 0 | 2 | 2 | 0 | 2 | 7270.423 | 0.005 | 3 |
| 0 | 0 | 2 | 2 | 1 | 2 | 7283.396 | 0.005 | 3 |
| 0 | 0 | 2 | 3 | 1 | 3 | 7315.259 | 0.003 | 2 |
| 0 | 0 | 2 | 2 | 2 | 1 | 7331.918 | 0.005 | 2 |
| 0 | 0 | 2 | 4 | 1 | 4 | 7357.593 | 0.008 | 2 |
| 0 | 0 | 2 | 5 | 1 | 5 | 7410.273 | 0.010 | 2 |
| 0 | 0 | 2 | 4 | 2 | 3 | 7412.047 | 0.073 | 2 |
| 0 | 0 | 2 | 4 | 2 | 2 | 7414.105 | 0.006 | 2 |
| 0 | 0 | 2 | 5 | 2 | 4 | 7469.165 | 0.005 | 2 |
| 0 | 0 | 2 | 6 | 2 | 5 | 7537.323 | 0.012 | 2 |
| 0 | 0 | 2 | 6 | 2 | 4 | 7546.316 | 0.010 | 2 |
| 0 | 0 | 2 | 5 | 3 | 3 | 7546.421 | 0.010 | 2 |
| 0 | 0 | 2 | 5 | 3 | 2 | 7546.640 | 0.005 | 2 |
| 0 | 0 | 2 | 6 | 3 | 4 | 7615.670 | 0.017 | 2 |
| 0 | 0 | 2 | 7 | 2 | 6 | 7616.477 | 0.018 | 2 |
| 0 | 0 | 2 | 7 | 2 | 5 | 7630.866 | 0.025 | 2 |
| 1 | 1 | 1 | 7 | 0 | 7 | 7641.009 | 0.001 | 2 |
| 0 | 0 | 2 | 8 | 3 | 5 | 7792.241 | 0.015 | 2 |
| 0 | 0 | 2 | 8 | 4 | 4 | 7894.171 | 0.090 | 2 |
| 0 | 0 | 2 | 8 | 4 | 5 | 7894.268 | 0.142 | 2 |
| 1 | 1 | 1 | 10 | 0 | 10 | 7919.703 | 0.067 | 2 |

The assignments were used to produce estimates of experimental band origins for the three observed bands. These values shown in Table 6 were derived from the calculated band origins and the observed band shifts described above. Based on the number of assignments used only the $2 v_{3}$ band origin should be considered reliable.

CD checks proved possible for 53 of the new assignments, comprising 25 upper energy levels. Agreement was found to be better than a standard deviation of $0.1 \mathrm{~cm}^{-1}$ in all cases, except two blended transitions. The large CD errors can in general be expected, as a result of using calculated lower energy levels and the low accuracy of the experiment, particularly for blended lines. The energy levels observed are presented in Table 7 and are included in the Supplementary Material.

The available PS based Tomsk line list provide more accurate positions in many, but not all cases, compared with the calculations made as part of this work. In particular positions calculated in this work were in general more accurate for the $v_{1}+v_{2}+v_{3}$ band. Intensities from the two sets of computations are of comparable quality.

## 6. Conclusions

Our analysis gives 111 new assignments to three different vibrational bands of HTO in the $7200-7245 \mathrm{~cm}^{-1}$ spectral region. Although there are many blended lines and the experimental accuracy is limiting, we consider the vast majority of the 83 presented $2 v_{3}$ assignments to be thoroughly reliable and an important development in this region for HTO.

The sensitivity of the experiment did not allow detection of any $\mathrm{T}_{2} \mathrm{O}$ features. Our line lists predicts this was due to the low $\mathrm{T}_{2} \mathrm{O}$ intensities in the small spectral region employed, in a region dominated by HTO. Future experiments will need to exploit greater sensitivity and/or greater tritium concentrations in order to assign this species in this region.

The derived band origins presented here are the first inferred experimentally for HTO beyond the fundamentals. As such they should be of some value to the theoretical community.

We present the analysed spectrum as Supplementary Material, alongside the newly observed energy levels. The line lists computed for HTO and $\mathrm{T}_{2} \mathrm{O}$ at 296 K , in the $0-10000 \mathrm{~cm}^{-1}$ range are also included. These include all $J \leqslant 15$ transitions, and relative intensities greater than $10^{-12}$ times the strongest line. Considering the undocumented nature of the Tomsk line lists available on-line, we hope these line lists will prove a valuable resource in future analysis of tritiated water spectra.

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## Appendix A. Supplementary material

Supplementary data for this article are available on ScienceDirect (www.sciencedirect.com) and as part of the Ohio State University Molecular Spectroscopy Archives (http://library.osu.edu/sites/ msa/jmsa_hp.htm). Supplementary data associated with this article can be found, in the online version, at http://dx.doi.org/ 10.1016/j.jms.2013.05.016.

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