



Miskolc Mathematical Notes
Vol. 13 (2012), No 2, pp. 429-439

HU e-ISSN 1787-2413
DOI: 10.18514/MMN.2012.520

On the blow-up of solutions for the unstable sixth order parabolic equation

Zhenbang Li and Changchun Liu



ON THE BLOW-UP OF SOLUTIONS FOR THE UNSTABLE SIXTH ORDER PARABOLIC EQUATION

ZHENBANG LI AND CHANGCHUN LIU

Received 3 May, 2012

Abstract. We study the universal blow-up of sixth-order parabolic thin film equation with the initial boundary conditions. We prove that the problem in finite time blow-up will happen, if the initial datum $u_0 \in C^{6+\alpha}(\bar{\Omega})$ with $-\int_{\Omega} (H(u_0) + \frac{1}{2}|\Delta u_0|^2) dx \geq 0$. And then, we get some nondegeneracy results on blow-up for this problem.

2000 *Mathematics Subject Classification:* 35K55; 35K90; 76A20

Keywords: blow-up, nondegeneracy, sixth order parabolic equation

1. INTRODUCTION

In this paper, we consider the following initial boundary problem of sixth-order equation

$$\begin{cases} u_t - \Delta(\Delta^2 u - |u|^{p-1}u) = 0, & \text{in } \Omega \times (0, T), \\ u = \Delta u = \Delta^2 u = 0, & \text{on } \partial\Omega \times [0, T), \\ u = u_0, & \text{in } \Omega \times \{0\}, \end{cases} \quad (1.1)$$

where $\Omega \subset R^N$ is a bounded smooth domain, $p > 1$.

During the past years, only a few works have been devoted to the sixth-order parabolic equation [1, 4, 5, 7].

Recently, Evans, Galaktionov and King [4, 5] considered the sixth-order thin film equation containing an unstable (backward parabolic) second-order term

$$\frac{\partial u}{\partial t} = \operatorname{div} [|u|^n \nabla \Delta^2 u] - \Delta(|u|^{p-1}u), n > 0, p > 1.$$

By a formal matched expansion technique, they show that, for the first critical exponent $p = p_0 = n + 1 + \frac{4}{N}$ for $n \in (0, \frac{5}{4})$, where N is the space dimension, the free-boundary problem with zero-height, zero-contact-angle, zero-moment, and zero-flux conditions at the interface admits a countable set of continuous branches of radially symmetric self-similar blow-up solutions $u_k(x, t) = (T - t)^{-\frac{N}{nN+6}} f_k(y)$, $y = \frac{x}{(T-t)^{\frac{1}{nN+6}}}$, where $T > 0$ is the blow-up time.

In fact, when $n = 0$, the equation (1.1) is obtained. In this paper we study the universal blow-up and some nondegeneracy results on blow-up of the equation (1.1). Our method about universal finite time blow-up is similar to that of Elliott and Zheng [3] which treats the blow-up problem for Cahn-Hilliard equation. We can show that if the initial datum $u_0 \in C^{6+\alpha}(\bar{\Omega})$ with $-\int_{\Omega} (H(u_0) + \frac{1}{2}|\Delta u_0|^2) dx \geq 0$, then the solution to the above problem (1.1) should blow up in finite time.

We also establish some nondegeneracy results on the blow-up of the problem. We mainly follow the purpose of Giga and Kohn [6] and Cheng and Zheng [2]. More accurately, there is a constant $\varepsilon > 0$, depending on n, p and the constant in the estimates of the fundamental solution to $u_t - \Delta^3 u = 0$ (see (3.1) below), such that if u is a solution of the equation

$$u_t - \Delta(\Delta^2 u - |u|^{p-1}u) = 0, \quad \text{on } Q_r = B_r(a) \times [t_1 - r^6, t_1),$$

where $1 < p < 3, a \in \mathbb{R}^n, t_1 \in \mathbb{R}$ and $0 < r \leq 1$, and if

$$|u(x, t)| \leq \varepsilon(t_1 - t)^{-\frac{2}{3(p-1)}} \quad \text{for all } (x, t) \in Q_r, \tag{1.2}$$

then u does not blow up at (a, t_1) .

The following sections include our main results. In Section 2, we establish universal finite time blow-up. Section 3 is devoted to the nondegeneracy results on the blow-up.

2. UNIVERSAL FINITE TIME BLOW-UP

Theorem 1. *Assume $u_0 \in C^{6+\alpha}(\bar{\Omega})$ with $-\int_{\Omega} (H(u_0) + \frac{1}{2}|\Delta u_0|^2) dx \geq 0$. Then the solution of the problem (1.1) must blow up at a finite time, namely, for some $T > 0$*

$$\lim_{t \rightarrow T} \|u(t)\| = +\infty,$$

where $H(u) = -\frac{|u|^{p+1}}{p+1}$.

Proof. Let

$$F(t) = \int_{\Omega} \left(H(u) + \frac{1}{2}|\Delta u|^2 \right) dx,$$

then

$$\begin{aligned} \frac{dF(t)}{dt} &= \int_{\Omega} \left(-|u|^{p-1}u\varphi(u)u_t + \frac{1}{2}\Delta u\Delta u_t \right) dx \\ &= \int_{\Omega} \left(-|u|^{p-1}u + \frac{1}{2}\Delta^2 u \right) u_t dx \\ &= -\int_{\Omega} |\nabla \left(-|u|^{p-1}u + \frac{1}{2}\Delta^2 u \right)|^2 dx \leq 0. \end{aligned}$$

So

$$2 \int_{\Omega} H(u) dx - 2F(0) \leq -\|\Delta u\|^2, \tag{2.1}$$

where

$$F(0) = \int_{\Omega} \left(H(u_0) + \frac{1}{2} |\Delta u_0|^2 \right) dx.$$

Let ϕ be the unique solution to

$$\begin{cases} \Delta \phi = u, & \text{in } \Omega, \\ \nabla \phi = 0, & \text{on } \partial\Omega. \end{cases}$$

It is easy to get that

$$\|\nabla \phi\|^2 \leq C \|\Delta \phi\|_2^2 \leq C \|u\|^2. \tag{2.2}$$

Now multiplying (1.1) by ϕ and integrating with respect x , we obtain

$$\begin{aligned} \frac{d}{dt} \|\nabla \phi\|^2 &= -2 \int_{\Omega} \varphi(u) u dx - 2 \|\Delta u\|^2 dx \\ &\geq 4 \int_{\Omega} H(u) dx - 4F(0) - 2 \int_{\Omega} \varphi(u) u dx \\ &= \int_{\Omega} \left(2 - \frac{4}{p+1} \right) |u|^{p+1} dx - 4F(0) \\ &\geq \frac{2(p-1)}{p+1} \left(\int_{\Omega} u^2 dx \right)^{\frac{p+1}{2}} - 4F(0). \end{aligned} \tag{2.3}$$

Combining (2.2), (2.3) and $-F(0) \geq 0$, we have

$$\frac{d}{dt} \|\nabla \phi\|^2 \geq \frac{2C(p-1)}{p+1} \|\nabla \phi\|^{p+1}. \tag{2.4}$$

Let $y(t) = \|\nabla \phi\|_2^2$ with $t \in [0, T)$, then

$$y'(t) \geq \gamma (y(t))^{\frac{p+1}{2}}, \tag{2.5}$$

where $\gamma = \frac{2C(p-1)}{p+1}$. A direct integration of (2.5) then yields

$$y^{\frac{p-1}{2}}(t) \geq \frac{1}{y^{\frac{1-p}{2}}(0) - \frac{p-1}{2} \gamma t}.$$

It turns out that the solution of the problem (1.1) will blow up in finite time. The proof of this theorem is completed. \square

3. NONDEGENERACY RESULTS ON THE BLOW-UP

Let $\Gamma(x, t)$ be the fundamental solution to $u_t - \Delta^3 u = 0$. According to [8], we have the follow inequalities:

$$|D_t^\mu D_x^\nu \Gamma(x, t)| \leq C t^{-\frac{1}{6}(n+6\mu+\nu)} \exp \left\{ -\omega \frac{|x|^{\frac{6}{5}}}{t^{\frac{1}{5}}} \right\}, \quad t > 0, \tag{3.1}$$

where $C > 0$, $\omega > 0$ are constants, and μ, ν are nonnegative integers.

Our purpose in this section is to have some nondegeneracy results on the blow-up. We state that the solution $u(x, t)$ to blows up at (a, t_1) if it is not locally bounded nearby, i.e., if there is a sequence $\{(x_k, \tau_k)\} \subset \Omega \times [0, t_1]$ with $(x_k, \tau_k) \rightarrow (a, t_1)$ as $k \rightarrow \infty$ such that $|u(x_k, \tau_k)| \rightarrow \infty$.

Theorem 2. *There is a constant $\varepsilon > 0$, depending on n, p and the constant in (3.1), such that if u is a solution of the equation*

$$u_t - \Delta(\Delta^2 u - |u|^{p-1}u) = 0, \quad \text{on } Q_r = B_r(a) \times [t_1 - r^6, t_1),$$

where $1 < p < 3, a \in \mathbb{R}^n, t_1 \in \mathbb{R}$ and $0 < r \leq 1$, and if

$$|u(x, t)| \leq \varepsilon(t_1 - t)^{-\frac{2}{3(p-1)}} \quad \text{for all } (x, t) \in Q_r, \tag{3.2}$$

then u does not blow up at (a, t_1) .

Next, we introduce the two lemma which will be used in the article and whose proofs can be found in [2] and [6].

Lemma 1. *For $0 < a < 1, \theta > 0$, and $0 < h < 1$, the integral*

$$I(h) = \int_h^1 (s - h)^{-a} s^{-\theta} ds,$$

satisfies

- (1) $I(h) \leq \left(\frac{1}{1-a} + \frac{1}{a+\theta-1} \right)$ if $a + \theta > 1$,
- (2) $I(h) \leq \frac{1}{1-a} + |\log h|$ if $a + \theta = 1$,
- (3) $I(h) \leq \frac{1}{1-a-\theta}$ if $a + \theta < 1$.

Lemma 2. *If $y(t), r(t)$ and $q(t)$ are continuous functions defined on $[t_0, t_1]$, such that $y(t) \leq y_0 + \int_{t_0}^t y(s)r(s)ds + \int_{t_0}^t q(s)ds, t_0 \leq t \leq t_1$, and $r(t) \geq 0$ on $[t_0, t_1]$, then*

$$y(t) \leq \exp \left\{ \int_{t_0}^t r(\tau) d\tau \right\} \left[y_0 + \int_{t_0}^t q(\tau) \exp \left\{ - \int_{t_0}^{\tau} r(\sigma) d\sigma \right\} d\tau \right].$$

Then, we began to prove the main Theorem 2.

Proof. Without loss of generality, we may assume $a = 0$ and $t_1 = 0$. By scaling, it is sufficient to consider the case $r = 1$. In the fact, if u satisfies the assumptions of the theorem with $r < 1$, then $u_r(x, t) = r^{\frac{4}{p-1}} u(rx, r^6 t)$ satisfies them with $r = 1$ (using the same ε), and clearly u_r blow up at $(0, 0)$ if u does.

Let ϕ be a smooth function supported on $B_1(0)$ such that $\phi \equiv 1$ on $B_{\frac{1}{2}}(0)$ and $0 \leq \phi \leq 1$. Consider $\omega = \phi u$; then $\omega_t - \Delta^3 \omega = g$ where

$$g = -2\nabla\Delta^2u\nabla\phi - \Delta^2u\Delta\phi - \Delta(u\Delta^2\phi + 4\nabla\Delta u\nabla\phi + 6\Delta u\Delta\phi + 4\nabla u\nabla\Delta\phi) - \phi\Delta(|u|^{p-1}u)$$

The semigroup representation formula for ω gives that

$$\omega(t) = e^{(t+1)\Delta^3}\omega(-1) + \int_{-1}^t e^{(t-s)\Delta^3}g(s)ds \quad \text{for } -1 \leq t < 0, \quad (3.3)$$

where $e^{t\Delta^3}$ is the semigroup associated with the equation $u_t - \Delta^3u = 0$ in \mathbb{R}^n , i.e.,

$$(e^{t\Delta^3}h)(x) = \int_{\mathbb{R}^n} \Gamma(x-y,t)h(y)dy.$$

Notice that $\int_{\mathbb{R}^n} \Gamma(x-y,t)dy = 1$. It follows that

$$\|e^{t\Delta^3}h\| \leq \|h\|_\infty. \quad (3.4)$$

The (3.1) implies that

$$\begin{aligned} |(e^{t\Delta^3}D_ih)(x)| &= \left| \int_{\mathbb{R}^n} \Gamma(x-y,t)D_ih(y)dy \right| \\ &= \left| \int_{\mathbb{R}^n} \frac{\partial}{\partial x_i} \Gamma(x-y,t)h(y)dy \right| \leq Ct^{-\frac{1}{6}}\|h\|_\infty, \quad \forall i = 1, 2, \dots, n, \end{aligned}$$

So, we get that

$$\begin{aligned} \|e^{t\Delta^3}D_ih\|_\infty &\leq Ct^{-\frac{1}{6}}\|h\|_\infty, \quad \|e^{t\Delta^3}D_{ij}h\|_\infty \leq Ct^{-\frac{1}{3}}\|h\|_\infty, \\ \|e^{t\Delta^3}D_{ijk}h\|_\infty &\leq Ct^{-\frac{1}{2}}\|h\|_\infty, \quad \|e^{t\Delta^3}D_{ijkl}h\|_\infty \leq Ct^{-\frac{2}{3}}\|h\|_\infty, \\ \|e^{t\Delta^3}D_{ijklm}h\|_\infty &\leq Ct^{-\frac{5}{6}}\|h\|_\infty, \end{aligned} \quad (3.5)$$

where $i, j, k, m, q \in \{1, 2, \dots, n\}$.

Now let $g = g_1 + g_2$, where $g_2 = -\phi\Delta(|u|^{p-1}u)$. As above, we estimate

$$\begin{aligned} &\left| \int_{-1}^t e^{(t-s)\Delta^3}g_2(s)ds \right| \\ &\leq \int_{-1}^t \left| \int_{\mathbb{R}^n} \Delta(\phi\Gamma(x-y,t-s))(|u|^{p-1}u)(y,s)dy \right| ds \\ &\leq \int_{-1}^t \left| \int_{\mathbb{R}^n} \Delta\Gamma(x-y,t-s)\phi|u|^{p-1}u(y)dy \right| ds \\ &\quad + \int_{-1}^t \left| \int_{\mathbb{R}^n} (\Gamma(x-y,t-s)\Delta\phi + 2\nabla\Gamma(x-y,t-s) \cdot \nabla\phi)|u|^{p-1}u(y)dy \right| ds \\ &\leq C \int_{-1}^t (t-s)^{-\frac{1}{3}}\|\phi u^p\|_\infty(s)ds + C \int_{-1}^t \|\Delta\phi u^p\|_\infty(s)ds \end{aligned}$$

$$\begin{aligned}
& + C \int_{-1}^t (t-s)^{-\frac{1}{6}} \|\nabla \phi u^p\|_{\infty}(s) ds \\
\leq & C \int_{-1}^t (t-s)^{-\frac{1}{3}} \|u\|_{\infty}^{p-1} \|\omega\|_{\infty}(s) ds + C \int_{-1}^t \|u^p\|_{\infty}(s) ds \\
& + C \int_{-1}^t (t-s)^{-\frac{1}{6}} \|u\|_{\infty}^p(s) ds \\
\leq & C \varepsilon^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\frac{2}{3}} \|\omega\|_{\infty}(s) ds + C \varepsilon^p \int_{-1}^t (-s)^{-\frac{2p}{3(p-1)}} ds \\
& + C \varepsilon^p \int_{-1}^t (t-s)^{-\frac{1}{6}} (-s)^{-\frac{2p}{3(p-1)}} ds, \tag{3.6}
\end{aligned}$$

due to our assumption.

On the other hand, it is found similarly that

$$\begin{aligned}
& \left| \int_{-1}^t e^{(t-s)\Delta^3} g_1(s) ds \right| \\
= & \left| \int_{-1}^t \int_{\mathbb{R}^n} \Gamma(x-y, t-s) (-2\nabla \Delta^2 u \nabla \phi - \Delta^2 u \Delta \phi \right. \\
& \left. - \Delta(u \Delta^2 \phi + 4\nabla \Delta u \nabla \phi + 6\Delta u \Delta \phi + 4\nabla u \nabla \Delta \phi)) dy ds \right| \\
\leq & C \int_{-1}^t (t-s)^{-\frac{5}{6}} \|u\|_{\infty}(s) ds \leq C \varepsilon \int_{-1}^t (t-s)^{-\frac{5}{6}} (-s)^{-\frac{2}{3(p-1)}} ds. \tag{3.7}
\end{aligned}$$

By (3.2)-(3.4), (3.6) and (3.7), we get that for $-1 \leq t < 0$,

$$\begin{aligned}
\|\omega(t)\|_{\infty} & \leq \varepsilon + \varepsilon^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\frac{2}{3}} \|\omega\|_{\infty}(s) ds \\
& + C \varepsilon^p \int_{-1}^t (t-s)^{-\frac{1}{6}} (-s)^{-\frac{2p}{3(p-1)}} ds + C \varepsilon \int_{-1}^t (t-s)^{-\frac{5}{6}} (-s)^{-\frac{2}{3(p-1)}} ds \\
& \leq \varepsilon + C \varepsilon^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\frac{2}{3}} \|\omega\|_{\infty}(s) ds + C \varepsilon (-t)^{\frac{1}{6} - \frac{2}{3(p-1)}}, \tag{3.8}
\end{aligned}$$

due to $1 < p < 3$ and Lemma (1).

Let $y(t) = \|\omega(t)\|_{\infty}$; therefore

$$y(t) \leq \varepsilon + C \varepsilon (-t)^{\frac{1}{6} - \frac{2}{3(p-1)}} + C \varepsilon^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\frac{2}{3}} y(s) ds. \tag{3.9}$$

Define $f(t) = \chi_{[-1,0]}(t) y(t)$, $\forall t < 0$. We introduce a special maximal function on $(-\infty, 0)$:

$$(Mf)(t) = \sup_{r>0} \frac{1}{r} \int_{t-r}^t |f(s)| ds, \quad \forall t \in (-\infty, 0).$$

Now $\forall r > 0$,

$$\begin{aligned} \int_{-1}^t (t-s)^{-\frac{1}{3}}(-s)^{-\frac{2}{3}}y(s)ds &= \int_{-\infty}^t (t-s)^{-\frac{1}{3}}(-s)^{-\frac{2}{3}}f(s)ds \\ &= \int_{t-r}^t (t-s)^{-\frac{1}{3}}(-s)^{-\frac{2}{3}}f(s)ds + \int_{-\infty}^{t-r} (t-s)^{-\frac{1}{3}}(-s)^{-\frac{2}{3}}f(s)ds \\ &= I_1 + I_2. \end{aligned}$$

We compute these two integrals, respectively.

$$\begin{aligned} I_1 &\leq (-t)^{-\frac{2}{3}} \int_{t-r}^t (t-s)^{-\frac{1}{3}}f(s)ds \\ &= (-t)^{-\frac{2}{3}} \sum_{k=0}^{\infty} \int_{t-\frac{r}{2^k}}^{t-\frac{r}{2^{k+1}}} (t-s)^{-\frac{1}{3}}f(s)ds \\ &\leq (-t)^{-\frac{2}{3}} \sum_{k=0}^{\infty} \left(\frac{r}{2^{k+1}}\right)^{-\frac{1}{3}} \int_{t-\frac{r}{2^k}}^{t-\frac{r}{2^{k+1}}} f(s)ds \\ &\leq (-t)^{-\frac{2}{3}} \sum_{k=0}^{\infty} \left(\frac{1}{2^{k+1}}\right)^{\frac{2}{3}} r^{\frac{2}{3}}(Mf)(t) \\ &= Cr^{\frac{2}{3}}(-t)^{-\frac{2}{3}}(Mf)(t), \end{aligned}$$

and

$$I_2 \leq r^{-\frac{1}{3}} \int_{-\infty}^{t-r} (-s)^{-\frac{2}{3}}f(s)ds \leq r^{-\frac{1}{3}} \int_{-\infty}^t (-s)^{-\frac{2}{3}}f(s)ds = r^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}}f(s)ds.$$

Then,

$$f(t) \leq \varepsilon + C\varepsilon(-t)^{\frac{1}{6}-\frac{2}{3(p-1)}} + C\varepsilon^{p-1} \left[r^{\frac{2}{3}}(-t)^{-\frac{2}{3}}(Mf)(t) + r^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}}f(s)ds \right],$$

for all $r > 0$ and $t \in (-\infty, 0)$.

Let

$$r = \frac{\int_{-1}^t (-s)^{-\frac{2}{3}}f(s)ds}{(-t)^{-\frac{2}{3}}(Mf)(t)},$$

so we have

$$\begin{aligned} f(t) &\leq \varepsilon + C\varepsilon(-t)^{\frac{1}{6}-\frac{2}{3(p-1)}} + C\varepsilon^{p-1} \left((-t)^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}}f(s)ds \right)^{\frac{2}{3}} ((Mf)(t))^{\frac{1}{3}} \\ &\leq \varepsilon + C\varepsilon(-t)^{\frac{1}{6}-\frac{2}{3(p-1)}} + C\varepsilon^{p-1} (-t)^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}}f(s)ds \\ &\quad + C\varepsilon^{p-1}(Mf)(t). \end{aligned} \tag{3.10}$$

If we define

$$g(t) = (-t)^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}} f(s) ds,$$

then

$$g'(t) = (-t)^{-1} \left[\frac{1}{3} (-t)^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}} f(s) ds + f(t) \right] \geq 0.$$

Hence $g(t)$ is increasing in $(-\infty, 0)$.

Then we get

$$\begin{aligned} \max_{-1 \leq \tau \leq t} f(\tau) &\leq \varepsilon + C\varepsilon(-t)^{\frac{1}{6} - \frac{2}{3(p-1)}} \\ &+ C\varepsilon^{p-1}g(t) + C\varepsilon^{p-1} \max_{-1 \leq \tau \leq t} (Mf)(\tau), \quad \forall t \in [-1, 0), \end{aligned} \quad (3.11)$$

where we have used $\frac{1}{6} - \frac{2}{3(p-1)} < 0$ since $1 < p < 3$.

Clearly, $\max_{-1 \leq \tau \leq t} (Mf)(\tau) \leq \max_{-1 \leq \tau \leq t} f(\tau)$ by our definition of the maximal function. Therefore (3.11) implies that for any $-1 \leq t < 0$,

$$\begin{aligned} \max_{-1 \leq \tau \leq t} f(\tau) &\leq \\ &\frac{1}{1 - C\varepsilon^{p-1}} \left[\varepsilon + C\varepsilon(-t)^{\frac{1}{6} - \frac{2}{3(p-1)}} + C\varepsilon^{p-1}(-t)^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}} f(s) ds \right], \end{aligned}$$

provided that $C\varepsilon^{p-1} < 1$. Especially,

$$\begin{aligned} f(t) &\leq \frac{1}{1 - C\varepsilon^{p-1}} \left[\varepsilon + C\varepsilon(-t)^{\frac{1}{6} - \frac{2}{3(p-1)}} + C\varepsilon^{p-1}(-t)^{-\frac{1}{3}} \int_{-1}^t (-s)^{-\frac{2}{3}} f(s) ds \right] \\ &\forall t \in [-1, 0). \end{aligned}$$

Then for $\varepsilon > 0$ small enough, we obtain

$$\begin{aligned} (-t)^{\frac{1}{3}} f(t) &\leq 2 \left[\varepsilon + C\varepsilon(-t)^{\frac{1}{2} - \frac{2}{3(p-1)}} + C\varepsilon^{p-1} \int_{-1}^t (-s)^{-1} (-s)^{\frac{1}{3}} f(s) ds \right] \\ &\forall t \in [-1, 0). \end{aligned}$$

Define $h(t) = (-t)^{\frac{1}{3}} f(t)$; then

$$h(t) \leq 2\varepsilon + 2C\varepsilon(-t)^{\frac{1}{2} - \frac{2}{3(p-1)}} + 2C\varepsilon^{p-1} \int_{-1}^t (-s)^{-1} h(s) ds, \quad (3.12)$$

Applying Lemma(2), we have

$$\begin{aligned} h(t) &\leq (-t)^{-2C\varepsilon^{p-1}} \left[2\varepsilon + C(p, \varepsilon)\varepsilon(-t)^{\frac{1}{2} - \frac{2}{3(p-1)}} + 2C\varepsilon^{p-1} \right] \\ &\leq 2\varepsilon(-t)^{-2C\varepsilon^{p-1}} + C(p, \varepsilon)\varepsilon(-t)^{\frac{1}{2} - \frac{2}{3(p-1)}}, \quad \forall t \in [-1, 0). \end{aligned}$$

Then $f(t) \leq 2\varepsilon(-t)^{-\frac{1}{3} - 2C\varepsilon^{p-1}} + C(p, \varepsilon)\varepsilon(-t)^{\frac{1}{6} - \frac{2}{3(p-1)}}$, $\forall t \in [-1, 0)$, or

$$y(t) \leq 2\varepsilon(-t)^{-\frac{1}{3} - 2C\varepsilon^{p-1}} + C(p, \varepsilon)\varepsilon(-t)^{\frac{1}{6} - \frac{2}{3(p-1)}}, \quad \forall t \in [-1, 0). \quad (3.13)$$

Choose $\varepsilon > 0$ small enough that $\frac{1}{3} + 2C\varepsilon^{p-1} < \frac{2}{3(p-1)}$ which is possible since $1 < p < 3$. Define $\alpha = \max\{\frac{1}{3} + 2C\varepsilon^{p-1}, \frac{2}{3(p-1)} - \frac{1}{6}\} \leq \frac{2}{3(p-1)}$, it is easy to find that $\alpha > \frac{1}{3}$; then (3.13) implies $y(t) \leq C(p, \varepsilon)\varepsilon(-t)^{-\alpha}, \forall t \in [-1, 0)$. Hence

$$|u(x, t)| \leq C(p, \varepsilon)\varepsilon(-t)^{-\alpha}, \quad \forall (x, t) \in B_{\frac{1}{2}}(0) \times [-1, 0). \quad (3.14)$$

Now let $\tilde{\phi}$ be a function supported on $B_{\frac{1}{2}}(o)$ with $\tilde{\phi} \equiv 1$ on $B_{\frac{1}{4}}(0)$ and $0 \leq \tilde{\phi} \leq 1$, and define $\tilde{\omega} = \tilde{\phi}u$; then we go back to (3.6)-(3.8) and we have that

$$\begin{aligned} \|\tilde{\omega}(t)\|_{\infty} &\leq \varepsilon + C \int_{-1}^t (t-s)^{-\frac{1}{3}} \|u\|_{\infty}^{p-1} \|\tilde{\omega}\|_{\infty} ds + C \int_{-1}^t \|u\|_{\infty}^p ds \\ &\quad + C \int_{-1}^t (t-s)^{-\frac{1}{6}} \|u\|_{\infty}^p ds + C\varepsilon \int_{-1}^t (t-s)^{-\frac{5}{6}} \|u\|_{\infty} ds \\ &\leq \varepsilon + C\varepsilon^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\alpha(p-1)} (-s)^{-\alpha} ds + C\varepsilon^p \int_{-1}^t (-s)^{-\alpha p} ds \\ &\quad + C\varepsilon^p \int_{-1}^t (t-s)^{-\frac{1}{6}} (-s)^{-\alpha p} ds + C\varepsilon \int_{-1}^t (t-s)^{-\frac{5}{6}} (-s)^{-\alpha} ds \\ &\leq \varepsilon + C\varepsilon^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\alpha p} ds + C\varepsilon^p \int_{-1}^t (-s)^{-\alpha p} ds \\ &\quad + C\varepsilon^p \int_{-1}^t (t-s)^{-\frac{1}{6}} (-s)^{-\alpha p} ds + C\varepsilon \int_{-1}^t (t-s)^{-\frac{5}{6}} (-s)^{-\alpha} ds \quad (3.15) \end{aligned}$$

due to (3.14).

Since $\frac{1}{3} < \alpha < \frac{2}{3(p-1)}$, we get

$$\frac{5}{6} - \alpha p > \frac{2}{3} - \alpha p > \frac{1}{6} - \alpha.$$

Hence by Lemma(1), we obtain

$$\|\tilde{\omega}(t)\|_{\infty} \leq \varepsilon + C\varepsilon^{p-1} + C\varepsilon^{p-1}(-t)^{\frac{1}{6}-\alpha} \leq (2 + C\varepsilon^{p-1})(-t)^{\frac{1}{6}-\alpha}, \quad \forall t \in [-1, 0),$$

Which means, for small $\varepsilon > 0$,

$$|u(x, t)| \leq (2 + C\varepsilon^{p-1})(-t)^{\frac{1}{6}-\alpha}, \quad \forall (x, t) \in B_r(0) \times [-1, 0). \quad (3.16)$$

Iterating the argument finitely many times we can get that there is a number $0 < r_0 < \frac{1}{4}$ such that

$$|u(x, t)| \leq K(-t)^{-\frac{1}{6p}}, \quad \forall (x, t) \in B_{r_0}(0) \times [-1, 0), \quad (3.17)$$

where K is constant.

Next, we choose another cut-off function $\hat{\phi}$ supported on B_{r_0} such that $\hat{\phi} \equiv 1$ on $B_{\frac{r_0}{2}}$ and define $\hat{w} = \hat{\phi}u$. Going back to (3.15) and applying Lemma(1), we have

$$\begin{aligned} \|\hat{w}(t)\|_\infty &\leq \varepsilon + C \int_{-1}^t (t-s)^{-\frac{1}{3}} \|u\|_\infty^{p-1} \|\hat{w}\|_\infty ds + C \int_{-1}^t \|u\|_\infty^p ds \\ &\quad + C \int_{-1}^t (t-s)^{-\frac{1}{6}} \|u\|_\infty^p ds + C \varepsilon \int_{-1}^t (t-s)^{-\frac{5}{6}} \|u\|_\infty ds \\ &\leq \varepsilon + CK^{p-1} \int_{-1}^t (t-s)^{-\frac{1}{3}} (-s)^{-\frac{1}{6}} ds + CK^p \int_{-1}^t (-s)^{-\frac{1}{6}} ds \\ &\quad + CK^p \int_{-1}^t (t-s)^{-\frac{1}{6}} (-s)^{-\frac{1}{6p}} ds + CK \int_{-1}^t (t-s)^{-\frac{5}{6}} (-s)^{-\frac{1}{6p}} ds \\ &\leq \varepsilon + CK^{p-1}, \end{aligned} \tag{3.18}$$

which means that $|u(x, t)| \leq C$ in $B_{\frac{r_0}{2}} \times [-1, 0)$. This completes the proof of the theorem. \square

Using the same argument, we can easily draw the following conclusion.

Theorem 3. Suppose $p \geq 3$, then for any $\delta \in (0, \frac{2}{3(p-1)})$, there is a constant $\varepsilon > 0$, depending on n, p and the constant in (3.1), such that if u is a solution of the equation

$$u_t - \Delta(\Delta^2 u - |u|^{p-1}u) = 0, \quad \text{on } Q_r = B_r(a) \times [t_1 - r^6, t_1)$$

where $a \in \mathbb{R}^n, t_1 \in \mathbb{R}$ and $0 < r \leq 1$, and if

$$|u(x, t)| \leq \varepsilon(t_1 - t)^{-\frac{2}{3(p-1)}} \quad \text{for all } (x, t) \in Q_r,$$

then u does not blow up at (a, t_1) .

REFERENCES

- [1] J. W. Barrett, S. Langdon, and R. Nürnberg, "Finite element approximation of a sixth order nonlinear degenerate parabolic equation," *Numer. Math.*, vol. 96, no. 3, pp. 401–434, 2004.
- [2] T. Cheng and G.-F. Zheng, "On the blow-up of solutions for some fourth order parabolic equations," *Nonlinear Anal., Theory Methods Appl.*, vol. 66, no. 11, pp. A, 2500–2511, 2007.
- [3] C. M. Elliott and S. Zheng, "On the Cahn-Hilliard equation," *Arch. Ration. Mech. Anal.*, vol. 96, pp. 339–357, 1986.
- [4] J. D. Evans, V. A. Galaktionov, and J. R. King, "Unstable sixth-order thin film equation. I: Blow-up similarity solutions," *Nonlinearity*, vol. 20, no. 8, pp. 1799–1841, 2007.
- [5] J. D. Evans, V. A. Galaktionov, and J. R. King, "Unstable sixth-order thin film equation. II: Global similarity patterns," *Nonlinearity*, vol. 20, no. 8, pp. 1843–1881, 2007.
- [6] Y. Giga and R. V. Kohn, "Nondegeneracy of blow up for semilinear heat equations," *Commun. Pure Appl. Math.*, vol. 42, no. 6, pp. 845–884, 1989.
- [7] A. Jüngel and J.-P. Milišić, "A sixth-order nonlinear parabolic equation for quantum systems," *SIAM J. Math. Anal.*, vol. 41, no. 4, pp. 1472–1490, 2009.

- [8] V. A. Solonnikov, "On boundary value problems for linear parabolic systems of differential equations of general form," *Proc. Steklov Inst. Math.*, vol. 83, p. 184, 1965.

Authors' addresses

Zhenbang Li

Department of Mathematics, Jilin University, Changchun 130012, China

E-mail address: jamesbom23@yahoo.com.cn

Changchun Liu

Department of Mathematics, Jilin University, Changchun 130012, China

E-mail address: liucc@jlu.edu.cn