

The Gaia-ESO Survey: Structural and dynamical properties of the young cluster Chamaeleon I $\star \star \star$

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ABSTRACT

Investigating the physical mechanisms driving the dynamical evolution of young star clusters is fundamental to our understanding of the star formation process and the properties of the Galactic field stars. The young (~ 2 Myr) and partially embedded cluster Chamaeleon I is one of the closest laboratories to study the early stages of star cluster dynamics in a low-density environment. The aim of this work is to study the structural and kinematical properties of this cluster combining parameters from the high-resolution spectroscopic observations of the Gaia-ESO Survey with data from the literature. Our main result is the evidence of a large discrepancy between the velocity dispersion ($\sigma_{\text{stars}} = 1.14 \pm 0.35 \text{ km s}^{-1}$) of the stellar population and the dispersion of the pre-stellar cores ($\sim 0.3 \text{ km s}^{-1}$) derived from submillimeter observations. The origin of this discrepancy, which has been observed in other young star clusters is not clear. It has been suggested that it may be due to either the effect of the magnetic field on the protostars and the filaments, or to the dynamical evolution of stars driven by two-body interactions. Furthermore, the analysis of the kinematic properties of the stellar population put in evidence a significant velocity shift ($\sim 1 \text{ km s}^{-1}$) between the two sub-clusters located around the North and South main clouds of the cluster. This result further supports a scenario, where clusters form from the evolution of multiple substructures rather than from a monolithic collapse.

Using three independent spectroscopic indicators (the gravity indicator γ , the equivalent width of the Li line at 6708 \AA , and the $H\alpha$ 10% width), we performed a new membership selection. We found six new cluster members all located in the outer region of the cluster, proving that Chamaeleon I is probably more extended than previously thought. Starting from the positions and masses of the cluster members, we derived the level of substructure Q , the surface density Σ and the level of mass segregation Λ_{MSR} of the cluster. The comparison between these structural properties and the results of N-body simulations suggests that the cluster formed in a low density environment, in virial equilibrium or supervirial, and highly substructured.

Key words. Stars: kinematics and dynamics – Stars: pre-main sequence – open clusters and associations: individual: Chamaeleon I – Techniques: spectroscopic

1. Introduction

The majority of stars do not form in isolation, but in clusters following the fragmentation and collapse of giant molecular clouds (Lada & Lada 2003; McKee & Ostriker 2007). Studying the formation and evolution of young clusters is fundamental to understand the star formation process and the properties of stars and planetary systems observed in the Galactic field, since they may depend on the formation environment (e.g., Johnstone et al. 1998; Parker & Goodwin 2009; Rosotti et al. 2014).

Despite the large number of multi-wavelength observations of nearby star forming regions carried out during the last two decades (e.g., Carpenter 2000; Getman et al. 2005; Güdel et al. 2007; Gutermuth et al. 2009; Feigelson et al. 2013), the scien-

tific debate on the initial conditions (i.e., stellar density, level of substructure, level of mass segregation) of star clusters and on the mechanisms driving the dissolution of most of them within 10 Myr is still open. In particular, it is not clear if the majority of stars form in very dense ($\gtrsim 10^4 \text{ stars pc}^{-3}$) and mass segregated clusters (e.g., Kroupa et al. 2001; Banerjee & Kroupa 2014), or in a hierarchically structured environment spanning a large density range (e.g., Elmegreen 2008; Bressert et al. 2010). Furthermore, the cluster dispersion may be triggered by gas expulsion due to the feedback of high-mass stars (e.g., Goodwin & Bastian 2006; Baumgardt & Kroupa 2007), or the dynamical evolution of star clusters could be driven by two-body interactions and the effect of the feedback may not be relevant (e.g., Parker & Dale 2013; Wright et al. 2014).

From the observational point of view, the main requirements to solve this debate are: a) an unbiased census of young stellar populations in star-forming regions spanning a large range of properties (i.e., density, age, total mass); b) the determination of the structural properties of young clusters (density, level of substructure and mass segregation) based on robust statistical

* This work is one of the last ones carried out with the help and support of our friend and colleague Francesco Palla, who passed away on the 26/01/2016

** Based on observations made with the ESO/VLT, at Paranal Observatory, under program 188.B-3002 (The Gaia-ESO Public Spectroscopic Survey)

methods that can be used for comparison with models; c) precise measurements of stellar velocities that allow us to resolve the internal dynamics of clusters and derive their dynamical status (e.g., virial ratio, velocity gradients).

During the last few years, progress has been achieved thanks to the theoretical efforts dedicated to better define the structural properties of clusters, (e.g., Cartwright & Whitworth 2004; Allison et al. 2009; Parker & Meyer 2012), to the comparison between models and observations (Sánchez & Alfaro 2009; Parker et al. 2011, 2012; Wright et al. 2014; Da Rio et al. 2014; Mapelli et al. 2015), and to dedicated observational studies of the dynamical properties of young clusters based on accurate radial velocities (RVs) (Fűrész et al. 2006, 2008; Cottaar et al. 2012a; Jeffries et al. 2014; Foster et al. 2015; Sacco et al. 2015; Tobin et al. 2015; Rigliaco et al. 2016; Stutz & Gould 2016). However, it is essential to extend these studies to a large number of clusters to cover the full space of relevant physical parameters (e.g., number of stars, stellar density, and age).

The young cluster Chamaeleon I (Cha I) is located around one of the dark clouds of the Chamaeleon star forming complex (see Luhman 2008 for an exhaustive review). It is the ideal laboratory to study the formation and early evolution of a low-mass cluster due to its proximity (distance= 160 ± 15 pc; Whittet et al. 1997), the presence of a molecular cloud actively forming stars, and a stellar population composed of ~ 240 members (Luhman 2008; Tsitali et al. 2015) distributed over an area of a few square parsecs.

The Cha I molecular cloud has been studied by Cambresy et al. (1997), who obtained an extinction map up to $A_V \sim 10$ mag, and by radio surveys of $C^{18}O$ or ^{12}CO emission (e.g., Boulanger et al. 1998; Mizuno et al. 1999, 2001; Haikala et al. 2005). In particular, Mizuno et al. (2001) used the ^{12}CO to estimate the total mass of the cloud ($\sim 1000 M_\odot$), while Haikala et al. (2005) mapped the structure of the filaments by observing the $C^{18}O$ emission. The filaments follow the structure of the cloud that is elongated in the NW-SE direction and perpendicular to the magnetic field. Protostellar cores within the cloud have been identified by Belloche et al. (2011) observing the continuum emission at $870 \mu m$ and, more recently, studied by Tsitali et al. (2015) in several molecular transitions. Tsitali et al. (2015) measured the velocity dispersion of the cores ($\sim 0.3 \text{ km s}^{-1}$) and conclude that their dynamical evolution is not affected by interaction and competitive accretion, since the collisional timescale is much longer than the core lifetime.

Several multi-wavelength studies have been carried out to identify the stellar and brown dwarf population of Cha I (e.g. Carpenter et al. 2002; Luhman 2004b; Luhman et al. 2008; Luhman 2008; López Martí et al. 2013; Lopez Martí et al. 2013, and references therein). Luhman (2008) compiled a list of 237 members using many membership indicators such as: the position in the HR diagram, high optical extinction, intermediate gravity between giants and main sequence stars, the presence of the Li absorption line at 6708 \AA , infrared excess emission, the presence of emission lines, proper motions and RVs. From the position in the HR diagram, Luhman (2007) derived a median age of 2 Myr and suggested that Cha I is divided in two sub-clusters, one concentrated in the northern part of the cloud ($\delta > -77^\circ$) and one in the south ($\delta < -77^\circ$). The former started to form stars 6 Myr ago, while the latter started later (4 Myr ago) and retains a larger amount of gas mass. Furthermore, they calculated an upper limit to the star formation efficiency of $\sim 10\%$. Several studies have been dedicated to measure the RVs and study the kinematics of the stellar and substellar populations of Cha I (Dubath et al. 1996;

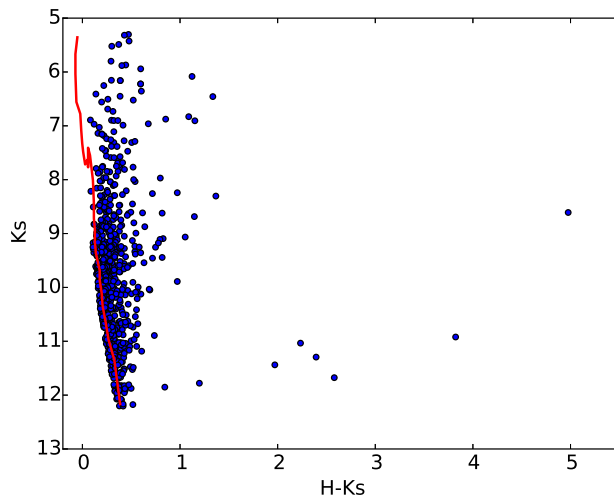


Fig. 1. Color magnitude diagram of the observed targets in the Cha I region based on photometry from the Two Micron All Sky Survey (Skrutskie et al. 2006), with overplotted the 10 Myr isochrone from the Siess et al. (2000) models.

Covino et al. 1997; Joergens & Guenther 2001; Joergens 2006; Guenther et al. 2007). In particular, Joergens (2006) measured a RV dispersion of 1.2 km s^{-1} , but this result is based on 25 stars only.

Cha I has been one of the first young clusters observed by the Gaia-ESO Survey (GES). GES is a large public spectroscopic survey carried out with the multi-object instrument FLAMES at the VLT, which feeds the medium- and high-resolution spectrographs GIRAFFE and UVES. The main goal of the survey is to derive RVs, stellar parameters (i.e., effective temperature, gravity, metallicity) and chemical abundances of 10^5 Milky Way stars in the field and in clusters (Gilmore et al. 2012; Randich & Gilmore 2013). A study of stellar activity, rotation and accretion based on the GES observations of Cha I is reported in Frasca et al. (2015), while the iron abundances of a selected sample of stars observed at high resolution are reported in Spina et al. (2014). Here, we investigate the structural and dynamical properties of Cha I combining the new GES results with data available in the literature. The paper is organized as follows: in Sect. 2, we present the method used for selecting the targets, the observations, and the data retrieved from the GES archive; in Sect. 3 we describe how we select the cluster members; in Sect. 4, we derive its structural properties; in Sect. 5, we derive the dynamical properties of the cluster; in Sect. 6, we discuss our results on the basis of the current models describing the dynamical evolution of low-mass star clusters, and in Section 7, we summarize the results and draw our conclusions.

2. Target selection, observations and data

The target selection and the fiber allocation procedure have been carried out independently for each cluster observed during the survey, however in order to maintain the homogeneity within the GES dataset we followed common guidelines described in Bragaglia et al. (2017, in preparation).

The selection of the targets for the observations of Cha I have been mostly based on the infrared photometry from the 2 micron all sky survey (2MASS, Skrutskie et al. 2006), since optical photometric catalogues available in the literature are incomplete and

Table 1. Members of the cluster observed by the Gaia-ESO Survey

Cname	RA (J2000)	DEC (J2000)	RV ^a (km s ⁻¹)	T_{eff} (K)	γ^b	EW(Li) (mÅ)	H α 10% ^c (km s ⁻¹)	Inst. ^d	memb ^e
10550964-7730540	163.79017	-77.51500	16.83±0.90	-	-	725±23	128±5	G	Y
10555973-7724399	163.99887	-77.41108	-	3640±300	-	120±58	441±17	U	Y
10561638-7630530	164.06825	-76.51472	12.56±2.01	-	-	-	215±5	G	Y
10563044-7711393	164.12683	-77.19425	15.69±0.27	4351±505	0.968±0.005	300±17	390±8	G	Y
10563146-7618334 ^f	164.13108	-76.30928	15.41±0.54	3319±48	-	597±11	142±4	G	N
10574219-7659356	164.42579	-76.99322	16.14±0.26	3452±183	0.895±0.004	571±16	272±6	G	Y
10575376-7724495	164.47400	-77.41375	15.71±0.32	3426±88	0.857±0.013	632±10	138±8	G	N
10590108-7722407	164.75450	-77.37797	15.64±0.40	4135±125	-	375±19	375±8	U	Y
10590699-7701404	164.77912	-77.02789	17.93±0.26	4981±260	-	390±1	440±9	G	Y
11004022-7619280	165.16758	-76.32444	15.73±0.42	-	-	584±18	230±5	G	Y

Notes. A full version of the table is available at the CDS.

(^a) For spectra with a signal-to-noise ratio lower than three we did not report any velocity.

(^b) Empirical gravity indicator defined by Damiani et al. (2014).

(^c) Width at 10% of the peak of the H α line.

(^d) The letters "G" and "U" indicate GIRAFFE and UVES, respectively.

(^e) The letters "Y" and "N" indicate the star is a known member or not, respectively

(^f) This star is likely a member of the ϵ Cha association.

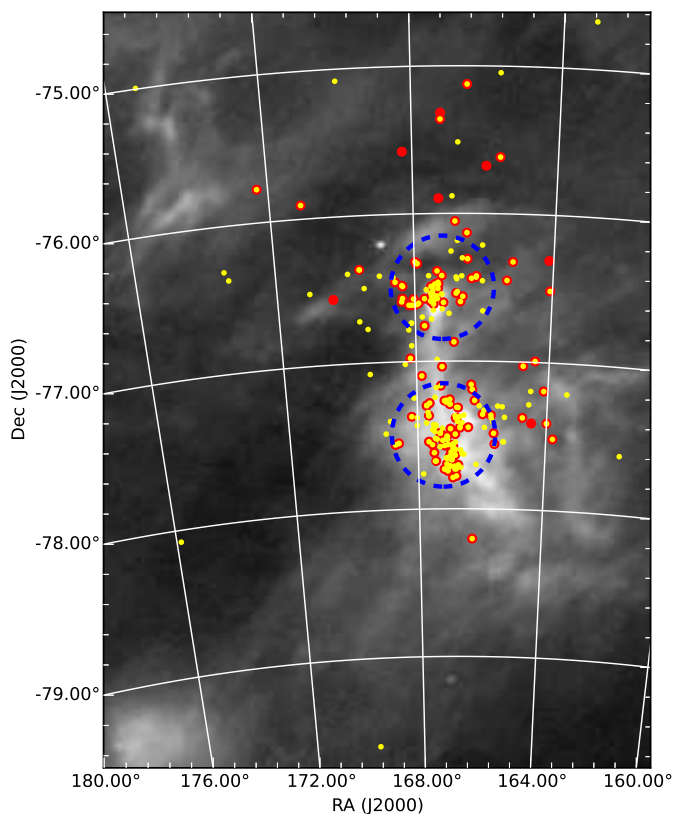


Fig. 2. Far infrared (140 μ m) map of the region around the young cluster Cha I from the *AKARI* all-sky survey (Doi et al. 2015). Yellow dots indicate the positions of all the known members from the literature, while the red bigger dots indicate the positions of all the members selected by the GES observations according to the criteria discussed in Sect. 3. The dashed blue circles (centers $RA_1 = 167.2^\circ$, $Dec_1 = -76.5^\circ$, $RA_2 = 167.2^\circ$, $Dec_2 = -77.5^\circ$ and radius 0.35°) delimit the north and south sub-clusters (see Sect. 4.2 and Sect. 5.3).

not homogeneous. The target selection and the fiber allocation process can be divided in two steps: we first compiled a list of

candidate members in the region of the sky around the cloud ($10:45 \leq RA \leq 11:30$ and $-79:00 \leq DEC \leq -75:00$), then we defined the position of the FLAMES fields of view (FOVs, diameter $25'$) and we allocated the largest possible number of fibers on candidate members.

For compiling the list of candidate members, we collected all the 2MASS sources with an optical counterpart from the Tycho 2 or USNO-B1 (Høg et al. 2000; Monet et al. 2003) catalogues brighter than $R=17$, which corresponds to the magnitude limit of the survey ($V=19$) for very low-mass stars. Then, from this list we selected only the sources that in a K vs. H-K color-magnitude diagram are located above the 10 Myr isochrone retrieved from the Siess et al. (2000) evolutionary models (see Fig. 1). Using this method, we compiled a list of 1933 candidate members. On the basis of this list and the positions of the known members, we chose 25 FOVs. Many of them were located along the main cloud, where most of the known members of the cluster are distributed, while a few FOVs were located on the outer regions, with the aim of looking for new members in regions which are poorly studied (the structure of the cloud and the positions of the known members are shown in Fig. 2). FOVs in the outer regions have been chosen in order to cover each latitude and longitude around the main cloud, focusing on the regions with higher spatial density of sources.

We observed a total of 674 stars with GIRAFFE and 49 with UVES (3 in common between the two spectrographs), of which 113 are known members of the cluster on the basis of catalogues reported in the literature (Luhman 2004a, 2007; Luhman et al. 2008). Most of the known members have been excluded because too faint to be observed with FLAMES. Several fibres have been allocated to the sky in order to allow a good background subtraction. Observations have been carried out during three different runs between March and May 2012¹, using the HR15N setup ($R \sim 17,000$, $\Delta\lambda = 647-679$ nm) for GIRAFFE and the 580 nm ($R \sim 47,000$, $\Delta\lambda = 480-680$ nm) setup for UVES. The median signal-to-noise ratio of the final spectra is 58 and 62 for GIRAFFE and UVES spectra, respectively.

¹ Technical details on the fibre allocation procedure and observations are discussed in Bragaglia et al. (2017, in preparation)

All GES data are reduced and analysed using common methodologies and software to produce an uniform set of spectra and stellar parameters, which is periodically released to all the members of the consortium via a science archive². In this paper, we use only data from the third internal data releases (GESviDR3) of February 2015, with the exception of errors on the GIRAFFE RVs, which are calculated on the basis of the empirical formulae provided by Jackson et al. (2015).

The methodologies used for the data reduction and the derivation of RVs are described in section 2.2 and 2.3 of Jeffries et al. (2014) for the GIRAFFE data, and in Sacco et al. (2014) for the UVES data. Lanzafame et al. (2015) describes in details the procedures used to derive the stellar parameters (i.e., effective temperature and gravities), accretion indicators (e.g., $H\alpha$ width at 10% of the peak – $H\alpha 10\%$), and the equivalent width of the Li line at 6708 Å (hereafter EW(Li)).

3. Membership selection

A detailed selection of members among the stars observed with UVES has been carried out by Spina et al. (2014), who confirmed all the known members from the literature and did not find any new members, therefore we will focus only on stars observed with GIRAFFE.

Since all the stars formed in the same region have very similar velocities, spectroscopic measurements of the RVs are often considered one of the most robust tools to select the members of a cluster. However, one of the main goals of this work is to study the dynamical properties of Cha I (e.g. the RV dispersion, the presence of multiple populations) therefore, we will use the RVs only to discard the obvious non-members, namely the stars outside of the range $0 < RV < 30 \text{ km s}^{-1}$, and for a more accurate selection of the cluster members, we will use three independent spectroscopic parameters included in the GES database: the gravity index γ , the EW(Li) and the $H\alpha 10\%$. The major source of contamination within a sample of candidate members of a nearby young cluster selected on the basis of photometric data are background giants. The giants can be identified using the surface gravity index γ , defined by Damiani et al. (2014) with the specific goal of measuring gravities using the GES GIRAFFE spectra observed with the HR15N setup. The upper panel of Fig. 3 shows γ as a function of the effective temperature for the GES targets observed in Cha I. The locus of the giant stars is clearly visible in the upper part of the plot and well separated from the main sequence and pre-main sequence stars. Similarly to previous works (Prisinzano et al. 2016; Damiani et al. 2014), we classified as giants all stars with an effective temperature lower than 5600 K and $\gamma > 1$.

After the giants stars have been excluded, we need to exclude stars older than Cha I located in the foreground. The most powerful tool to perform this selection is the EW(Li), since late-type stars rapidly deplete their photospheric lithium after 5-30 Myr (e.g., Soderblom 2010). At constant age, the EW(Li) depends on the effective temperature, therefore we cannot define a single threshold for the whole sample. The lower panel of Fig. 3 shows the EW(Li) as a function of T_{eff} for the observed stars in Cha I and for the stars of the 30-50 Myr open cluster IC 2602 observed by Randich et al. (1997, 2001). We select as cluster members all the stars with EW(Li) above the dashed line in the lower panel of Fig. 3, which represents the upper envelope of the

the EW(Li) measured for the stars belonging to IC 2602. In a few cases, when the EW(Li) but not the effective temperature have been derived from the GES spectra (see Lanzafame et al. 2015 for details), we assume the highest threshold (EW(Li)=300 Å). However, the EW(Li) can be underestimated in stars with a very strong mass accretion rate due to the continuum emission in excess with respect to the photospheric one produced by the accretion shock (e.g., Palla et al. 2005; Sacco et al. 2007), therefore, we include in the sample of members all the stars that can be classified as accretors according to the criterium based on the width at 10% of the peak of the $H\alpha$ line ($H\alpha 10\% > 270 \text{ km s}^{-1}$) defined by White & Basri (2003). In the bottom panel of Fig. 3 the cluster members are indicated with filled blue circles. The members below the dashed line have been included because of the $H\alpha 10\%$ while the non-members above the line have been excluded because of the gravity index γ .

Using these criteria, we selected as members of Cha I 89 stars observed with GIRAFFE. This sample includes 7 new members and 82 known members from the literature. Fourteen known members do not meet the membership criteria. Seven of them have been excluded because they are out of the range of RVs ($0 < RV < 30 \text{ km s}^{-1}$). However, all these stars are strong accretors ($H\alpha 10\% > 300 \text{ km s}^{-1}$), therefore the RV derived by the GES pipeline could be wrong due to the presence of strong emission lines produced by material moving at a different velocity with respect to the photosphere. We will use these stars for the analysis of the structural properties of the cluster (see Sect. 4), but we will exclude them from the analysis of the dynamical properties, which is based on the RVs. Six known members have been excluded because the SNR < 10 is too low to derive EW(Li) and $H\alpha 10\%$ (see Lanzafame et al. 2015), but there is no evidence suggesting that these stars are not members, therefore, we will include them in the final catalogue. Finally, the star HD 97300 is too hot (SpT= B9) to exhibit the Li absorption feature at 6708 Å, but it is surrounded by a ring of dust due to a bubble blown by the star (Kóspál et al. 2012). This proves that it belongs to the Cha I star forming region and can be included in the catalogue of members used for the analysis discussed in the following sections. To summarize, the final catalogue of members includes 103 stars (96 already known and 7 new) observed with GIRAFFE and 17 stars observed with UVES (discussed in Spina et al. 2014) for a total of 120 members. We note that our analysis proves that the 7 new members are young stars, but does not demonstrate that belong to Cha I, since they could be members of the two young stellar associations ϵ Cha and η Cha, which are slightly older (4-9 Myr, Torres et al. 2008), have a similar radial velocities and are located in the foreground of Cha I. As shown by Lopez Martí et al. (2013) (see Fig. 1 of that paper), the two associations can be easily separated from Cha I using tangential motions, so we retrieved the proper motions of these stars from the UCAC 4 (Zacharias et al. 2013) and, when not available, from the PPMXL (Roeser et al. 2010) catalogues. We found that all the new members have tangential motions consistent with the Cha I cluster except one star (GES ID. 10563146-7618334), which is closer to the ϵ Cha association.

The list of cluster members and the data used for the membership selection are reported in Table 1, while their positions are plotted on the map in Fig. 2 (red dots) together with all the known members of the clusters compiled from the literature (yellow dots). The number of new members does not significantly increase the population of Cha I. However, they belong to the sparse population located in the outer region surrounding the

² The GES science archive is run by the Royal Observatory of Edinburgh. More information on the archive are available at the website ges.roe.ac.uk

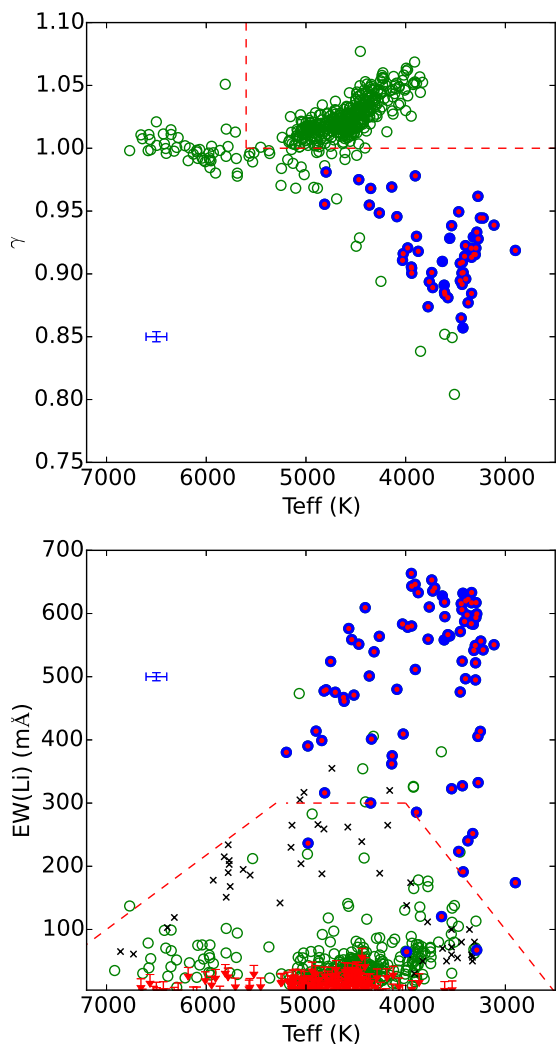


Fig. 3. The upper and the lower panels show the gravity index γ and EW(Li) as a function of the stellar effective temperature, respectively. The index γ and the EW(Li) measured from the GES spectra are indicated with circles: the empty green circles are the non-members, the filled blue circles are the cluster members, and the red dots represent the members already known from previous studies. Upper limits are reported with red downwards arrows. Median error bars are reported on the left for both panels. The black crosses in the bottom panel are the EW(Li) measured for stars in the 30-50 Myr open cluster IC 2602 by Randich et al. (1997, 2001). In both panels, the dashed red lines indicate the threshold used to separate members and non-members. In the bottom panel, a few stars are reported as non-members despite being above the dashed red line because they have been excluded according to the criterium based on the gravity index, and other stars are classified as members despite being below the dashed line, because they are strong accretors, as discussed in Sect. 3.

main cloud. So our study proves that this outer population is richer than previously thought.

4. Structural properties

As proved by several studies (e.g., Scally & Clarke 2002; Schmeja et al. 2008; Allison et al. 2010; Moeckel & Bate 2010; Malmberg et al. 2011; Kruijssen et al. 2012; Parker et al. 2014), a knowledge of the structural properties of open clusters is fundamental to understand their origin, their dynamical evolution and the effects of the star formation environment on the properties of stars and planetary systems. In this work we will focus on three structural properties: the level of substructure, the stellar density and the mass segregation.

4.1. Sample and stellar masses

The sample of stars used for the structural analysis includes all the previously known members (observed or not by GES) and the new members discovered by GES. We exclude stars with $A_J > 1.2$, because catalogues available in the literature are not complete for higher extinction (Luhman 2007).

For the analysis of the mass segregation, we derived homogeneous estimates of stellar masses from the positions of stars in the HR diagram plotted in Fig. 4, using the pre-main sequence evolutionary models developed by Tognelli et al. (2011, 2012). We decide to use this set of models instead of that provided by Siess et al. (2000) and recommended by the Gaia-ESO guidelines for the target selection process, because they have been more recently updated. No masses have been estimated for stars cooler than 3000 K, younger than 0.5 Myr and older than 20 Myr (i.e., above/below the upper/lower isochrones plotted in the HR diagram). Very cool and very young stars have been excluded because mass estimation based on pre-main sequence evolutionary models could be very uncertain. Stars located below the 20 Myr isochrone have been excluded because, as suggested by Luhman (2007); Luhman & Muench (2008), their luminosity is underestimated due to the presence of a circumstellar disk seen edge-on, which absorbs most of the photospheric emission.

To build the HR diagram we use 2MASS photometry and the GES parameters when available, otherwise we use the parameters from Luhman (2004a) and Luhman (2007). Specifically, the effective temperatures have been directly measured from the GIRAFFE and UVES spectra (Lanzafame et al. 2015) in the GES sample and from low resolution spectra for the other stars (Luhman 2004a, 2007). A comparison of the effective temperatures derived by the GES spectra and those obtained from low resolution spectroscopy is reported in Fig. 5, which shows that the latter are slightly lower in the range between 3600 and 4600 K. We believe that this systematic discrepancy can be due to either the scale used by Luhman (2004a, 2007) to convert spectral types into temperature (GES temperature are derived directly from the spectra) or to the different spectroscopic resolution of the spectra used for the analysis.

Luminosities have been derived from the 2MASS J magnitude corrected for absorption, using the bolometric correction reported in Pecaut & Mamajek (2013) and assuming a distance of 160 pc. For the GES sample, we calculate the absorption A_J from $E(J - H) = (J - H) - (J - H)_0$, assuming $A_J/E(J - H) = 0.38$ as in Luhman (2004a). The intrinsic color $(J - H)_0$ was calculated from the effective temperature using the color-temperature transformations from Pecaut & Mamajek (2013). For stars that do not belong to GES catalogues, the infrared absorption has been retrieved from Luhman (2007, 2004b). The masses of the stars used for the structural analysis are reported in table 2.

To understand if the use of different catalogues for the estimation of masses can affect our results, we performed all our

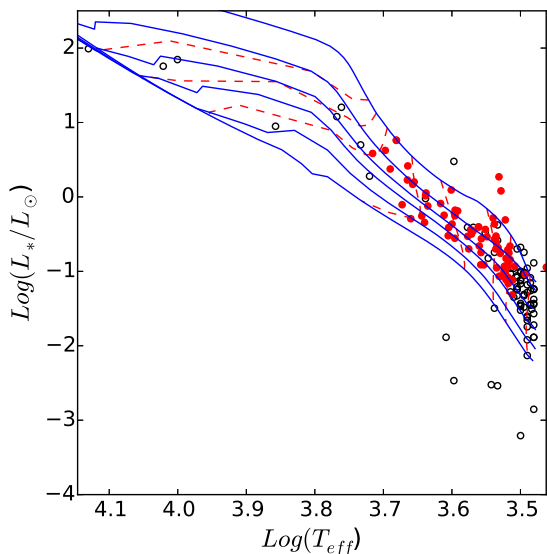


Fig. 4. HR diagram of members in Cha I selected from the literature and the GES data with $A_J < 1.2$ and $T_{eff} > 3000$ K. Temperatures and luminosities have been derived from the GES spectra and the 2MASS photometry for the red dots, and from the literature for the other stars. Isochrones (at 0.5, 1.5, 3.0, 5.0, 10.0, and 20.0 Myr) and tracks (at 0.1, 0.2, 0.3, 0.6, 1.0, 1.8, 2.4, 3.0 solar masses) from an improved version of the Tognelli et al. (2011, 2012) pre-main sequence evolutionary models are reported with continuous blue and dashed red lines, respectively.

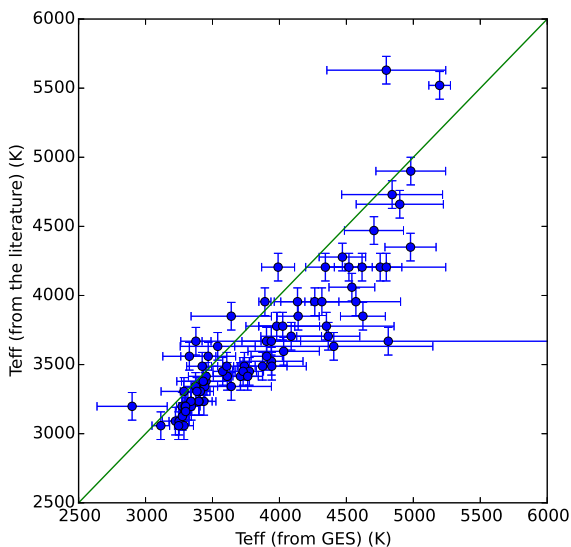


Fig. 5. Comparison between the effective temperature retrieved from the literature (y-axis) and measured from the GES spectra (x-axis).

analysis using only data from the literature for all the known members finding no relevant differences. We stress that in the context of this paper the stellar masses are only used for studying the level of the mass segregation (see Sect. 4.3 and 4.4). For this scope, we do not need the specific values of the masses but only to put them in order from the most to the less massive.

Table 2. Known cluster members from the literature and new members used to study the structural properties of the cluster

RA (J2000)	DEC (J2000)	T_{eff} (K)	$\text{Log}\left(\frac{L_{Bol}}{L_{\odot}}\right)$	prov ^a	Mass (M_{\odot})
161.65812	-77.60097	7200	0.95	L	1.64
163.15392	-74.67464	3161	-1.00	L	0.14
163.41575	-77.20939	3451	-1.49	L	0.30
163.79017	-77.51500	3198	-1.08	L	0.15
163.99887	-77.41108	3640	-0.36	G	0.37
163.99887	-77.41108	3640	-0.36	G	0.37
164.06825	-76.51472	3044	-1.51	L	0.09
164.12683	-77.19425	4350	0.06	G	0.85
164.42579	-76.99322	3451	-0.33	G	0.28

Notes. A full version of the table is available at the CDS.

^(a) The letters G and L indicate that the data used for deriving the stellar mass are retrieved from the literature and the GES archive, respectively.

4.2. Level of substructure

To measure the level of substructure, we use the Q -parameter introduced by Cartwright & Whitworth (2004). This parameter is defined as the ratio:

$$Q = \frac{\bar{m}}{\bar{s}} \quad (1)$$

where \bar{m} is the mean length of the edges of the minimum spanning tree (MST) connecting the stars, normalized by the factor $(NA)^{1/2}/(N-1)$ (N is the total number of stars and A is the area of the cluster), and \bar{s} is the mean separation between the stars divided by the cluster radius³. Several simulations demonstrate that clusters with $Q > 0.8$ are characterized by a smooth radially concentrated structure, which is probably the result of dynamical evolution occurring after the cluster formation, while clusters with $Q < 0.8$ are characterized by a high level of substructure and closely resemble the filamentary structure of the molecular clouds where they formed (e.g., Schmeja & Klessen 2006; Parker & Meyer 2012). To calculate the best estimator \hat{Q} and the standard error $\sigma(\hat{Q})$ of the Q -parameter, we use the *Jackknife* method (Quenouille 1949; Tukey 1958). This is a resampling method that consists in calculating the value of Q_i for N different samples composed of all the stars except for the i -th star (with i from 1 to N). The best estimator and the standard error are equal to:

$$\hat{Q} = \frac{1}{N} \sum_{i=1}^N Q_i \quad \sigma(\hat{Q}) = \sqrt{\frac{N-1}{N} \sum_{i=1}^N (Q_i - \hat{Q})^2} \quad (2)$$

The relation between Q and the level of substructure has been calculated by Cartwright & Whitworth (2004) through simulations considering isotropic clusters. However, Cartwright & Whitworth (2009) pointed out that in elongated clusters Q could be biased towards lower values and estimated a correction factor which depends on the cluster aspect ratio. Since Cha I is characterized by a slightly elongated structure, we applied these corrections on our results.

³ The radius and the area of the cluster are calculated as described in Cartwright & Whitworth (2004). Specifically the former is the distance between the center of the cluster and the most distant stars, and the latter is the area of a circular surface with the same radius.

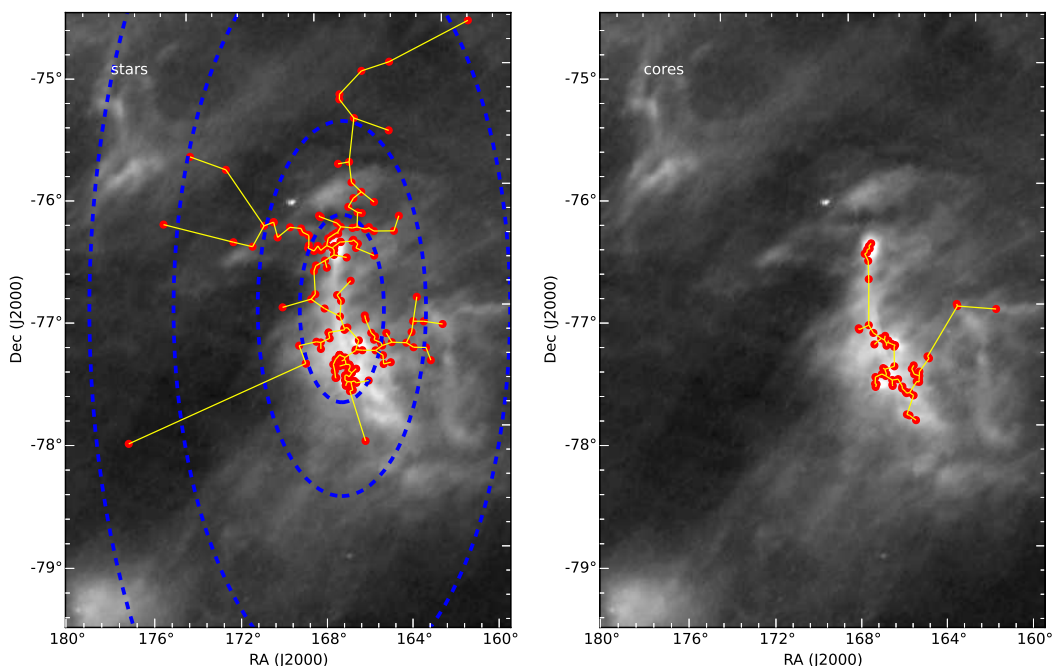


Fig. 6. Spatial distribution of the stars (left panel) and starless cores (right panel) used to calculate the Q parameter. The minimum spanning tree is plotted with a yellow line. The blue dashed lines in the left panel show the elliptic boundaries of the regions including the stars used to calculate the values of the Q parameter reported in Fig. 7.

For our sample, the resulting $\hat{Q} = 0.80 \pm 0.08$ is higher than what found by previous studies carried out by Cartwright & Whitworth (2004) and Schmeja & Klessen (2006) ($Q = 0.68$ and 0.67 , respectively). Cartwright & Whitworth (2004) did not consider the elongation in their calculation, but this does not explain such a large discrepancy, since the elongation factor for our sample is just 1.03. Therefore, this discrepancy is most likely due to the different sample of members used for the calculations. In fact, Cartwright & Whitworth (2004) used the sample of members selected by Lawson et al. (1996) and Ghez et al. (1997), while Schmeja & Klessen (2006) results are based on a sample of members retrieved from Cambresy et al. (1998). Both samples are less complete and cover a smaller area of the sky than ours. As discussed in the previous section, Cha I is composed of an inner denser region characterized by the presence of a molecular cloud still forming stars and an outer sparser region with no gas. Furthermore, our discovery of new members only in the outer part of the cluster proves that the level of completeness of the star catalogue in the outer region is lower than in the inner one. To understand how this can affect the Q -parameter, we calculated Q for four different samples composed of stars enclosed within the area of the sky delimited by ellipses with the same center ($RA = 167.2^\circ$, $DEC = -77.1^\circ$) and eccentricities $e = 0.89$, but different semi-major axes (see table 3 and the left panel of Fig. 6). In particular, the smallest ellipse contains only the stars in the inner embedded region, the largest includes all the stars and the two intermediate ones have a semi-major axes twice and four times the smallest.

The value of Q as a function of the semi-major axes of the ellipses is shown in Fig. 7 and reported in table 3. The Q -parameter gradually increases from $Q \sim 0.6$ – when we consider only the stars in the inner and denser region of the cluster – to $Q \sim 0.8$ for the full sample. This clear correlation between the Q -parameter and the area of the cluster considered to perform the calculation may be due to one or more of the following reasons. a) The stars in the inner region are younger

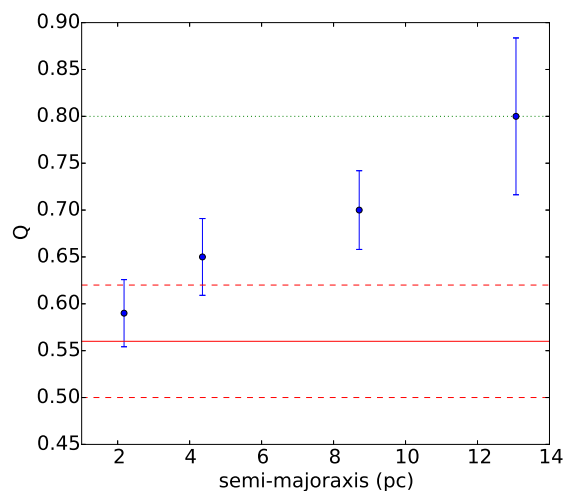


Fig. 7. The Q -parameter as function of the semi-major axes of the ellipses, shown in Fig. 6, which delimit the area enclosing all the stars used for the calculation. The dotted green line indicates the value of Q expected for a sample of stars randomly distributed, while the continuous and the dashed red lines indicate the the Q -parameter with error bars calculated from the positions of the pre-stellar cores.

and located very close to where they formed, so Q is similar to what expected at the initial stage of the cluster formation (e.g., Parker et al. 2014), when the distribution of stars resembles the distribution of gas in filaments, while the stars in the outer regions migrated from their formation site, so their spatial distribution has been randomized; b) Cha I is composed of multiple populations with different structural properties. In the different subsets defined by the ellipses, these populations have different weights, so the value of Q changes according to which population is weighed the most; c) The presence of patchy extinction in the inner region produces substructures, which will not

Table 3. Properties of the ellipses used for investigating the relation between Q and the completeness of the member sample.

Semi-major axis (degree)	N_{star}	Q
0.78	98	0.59 ± 0.04
1.56	143	0.65 ± 0.04
3.12	157	0.70 ± 0.04
4.68	160	0.80 ± 0.08

be present if the full sample of stars would be visible. d) The member selection in the outer regions is not complete. Missing some of the members, we can miss some of the substructures, so Q increases, when we include the outer part of the cluster.

Luhman (2007) suggested that Cha I is composed of two subclusters with different star formation histories, that may have followed independent dynamical evolutions. Starting from this assumption we estimated Q for two independent samples including the stars within the two red circles represented in Fig. 2, which are likely to include only stars belonging to one of the two subclusters. For both subclusters we found a value of $Q = 0.76 \pm 0.06$, consistent with a randomly distributed sample. This result suggests that if the two subclusters are independent their dynamical evolution already erased primordial substructures. However, simulations suggest that the Q -parameter is statistically robust only for clusters with a number of stars larger than 100 (e.g., Cartwright & Whitworth 2004; Parker & Dale 2015), while the regions delimited in Fig. 2 include only 48 and 63 stars.

Furthermore, we derived the Q -parameter ($Q = 0.56 \pm 0.06$) from the positions of 60 pre-stellar cores found by Belloche et al. (2011) with a submillimeter survey. The agreement between the value of Q measured for the cores and for the stars in the smaller ellipse supports the explanation at point a). However, as discussed above the number of objects is too low to consider this measurement statistically robust.

4.3. Stellar density

The stellar density is a key parameter to derive the effects of the the environment on the evolution of the star-disc systems and the dynamical status of the clusters. To derive the surface density Σ we used the same definition as Bressert et al. (2010):

$$\Sigma = \frac{N - 1}{\pi D_N^2} \quad (3)$$

where N is the N th nearest neighbour and D_N is the projected distance to that neighbour. For our calculation we set $N = 7$ as in Bressert et al. (2010). The top panel in Fig. 8 shows the density distribution of the stars used to calculate the Q parameter with the best fit of the distribution with a log-normal function (peak ~ 8 stars pc^{-2} and dispersion $\sigma_{\log_{10}\Sigma} = 0.67$). The profile of the distribution is very similar to what observed by Bressert et al. (2010) (see Fig. 1 of the paper) and is well described by a log-normal function in the low-density tail, while at high density the observed distribution decreases faster than a log-normal function. The reason for this deviation from the log-normal model is not clear. However a similar deviation is also observed in the much larger sample analysed by Bressert et al. (2010). The peak of the distribution is located at lower densities and the dispersion is lower with respect to the results found by Bressert et al. (2010), but this is not surprising since $\sim 70\%$ of the young stellar

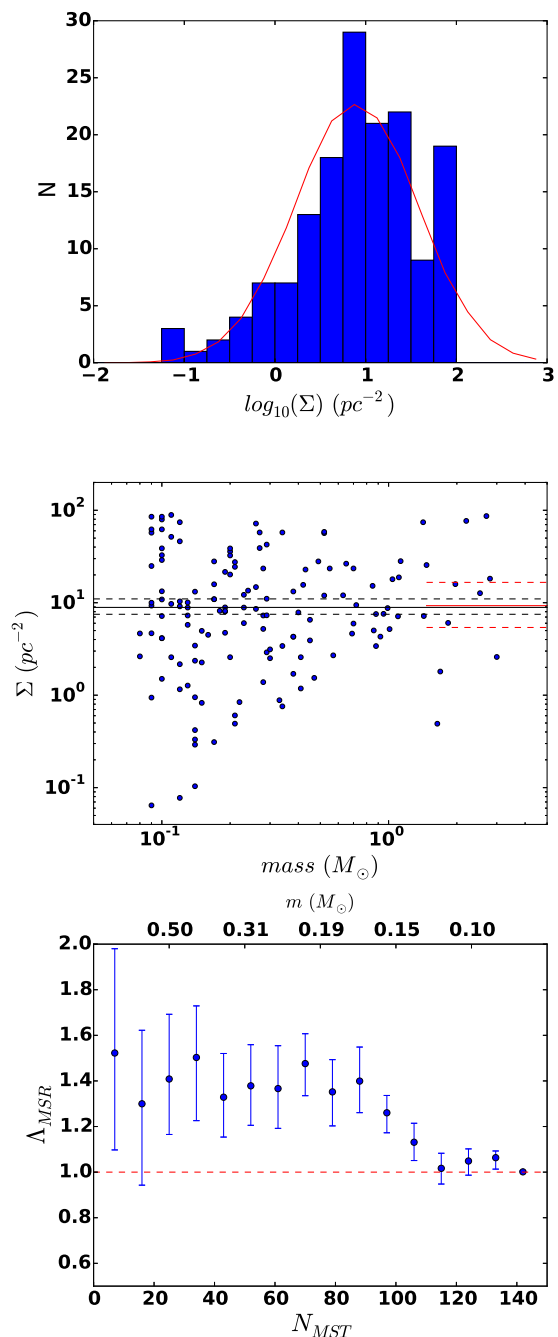


Fig. 8. The top panel reports the distribution of the surface density Σ and the best-fit model with a lognormal function (continuous red line). The middle panel shows the stellar density as defined in Eq. (3) as a function of the star mass. The black and the red lines indicate the median density of all the stellar sample and the ten most massive stars, respectively. The continuous lines represent the best value, while the dashed lines represent their error bars. The bottom panel shows the evolution of the mass segregation ratio, Λ_{MSR} , for the N_{MST} most massive stars. The top x-axis indicates the lowest mass star m_L within N_{MST} and the dashed line corresponds to $\Lambda_{MSR} = 1$, i.e., no mass segregation.

objects used for their analysis belong to the Orion star forming region, which is more massive and much denser than Cha I.

Simulations describing the dynamical evolution of young star clusters suggest that the stellar density may depend on stellar mass, i.e. the density of stars near massive objects can be higher because massive stars act as a potential well and trap low

mass stars (e.g., Parker et al. 2014). The relation between density and mass is plotted in the middle panel of Fig. 8, which shows that the density around the most massive stars (median density $\Sigma = 9.3_{-3.9}^{+7.3} \text{ pc}^{-2}$) is consistent within errors with the surface density for the rest of the stars (median density $\Sigma = 8.9_{-1.4}^{+2.1} \text{ pc}^{-2}$)⁴. This result proves that either the cluster did not go through a sufficient dynamical evolution to determine an increase of density around the most massive stars or that these stars are not massive enough to act as a potential well and attract low-mass stars.

4.4. Mass segregation

The last structural property to analyse is the amount of mass segregation. In mass-segregated clusters, the more massive stars are concentrated in a smaller volume (or projected area on the line of sight) than lower mass stars. To estimate the level of mass segregation we used the method introduced by Allison et al. (2009) and based on the mass segregation ratio Λ_{MSR} :

$$\Lambda_{MSR} = \frac{\langle l_{average} \rangle^{+\sigma_{5/6}/l_{subset}}}{l_{subset}^{-\sigma_{5/6}/l_{subset}}} \quad (4)$$

where l_{subset} is the length of the MST of a subset of stars composed of the number N_{MST} most massive stars of the cluster, and $l_{average}$ is the average of the lengths of the MSTs of 50 different subsets composed of a number N_{MST} random stars. If $\Lambda_{MSR} > 1$ the MST of the more massive stars is smaller than the MST of a random sample so the cluster is mass-segregated, otherwise if $\Lambda_{MSR} < 1$ it is inversely mass segregated (i. e., the most massive stars are spread over a larger area than other stars). For estimating the uncertainties on this ratio, we used the same method as in Parker et al. (2012), namely we considered as lower (upper) error the length of the MST, which lies at $1/6(5/6)$ of an ordered list including all the MSTs of the random subsets used to calculate $l_{average}$. In the bottom panel of Fig. 8 we show the evolution of Λ_{MSR} as function of N_{MST} . The upper x-axis shows the smallest mass within the sample of N_{MST} stars. The plot shows only a marginal evidence of mass segregation, which is not significant since the value of Λ_{MSR} at higher masses is consistent with $\Lambda_{MSR}=1$, and Λ_{MSR} for intermediate mass stars is above 1 by less than 2-3 error bars, which as estimated by simulations performed by Parker & Goodwin (2015) means a significance lower than 95%.

5. Kinematical properties

The precision of the RVs derived from the GES spectra (Jackson et al. 2015) allows us to study the kinematical properties of the cluster. We will use the RVs to determine its global RV dispersion σ_c , to investigate the presence of a RV gradient and to understand, if the two different populations identified by Luhman et al. (2008) have different kinematical properties. For this analysis, we will use only members of the cluster observed by GES and reported in Table 1, since only for these stars we have precise measurements of the RV with a proper evaluation of the errors.

⁴ The best values and the error bars of the densities have been calculated by generating 2000 bootstrap resamples. Namely, the best value is the median of the bootstrap distribution, while the lower and upper values defined by error bars correspond to the 15th and 85th percentiles.

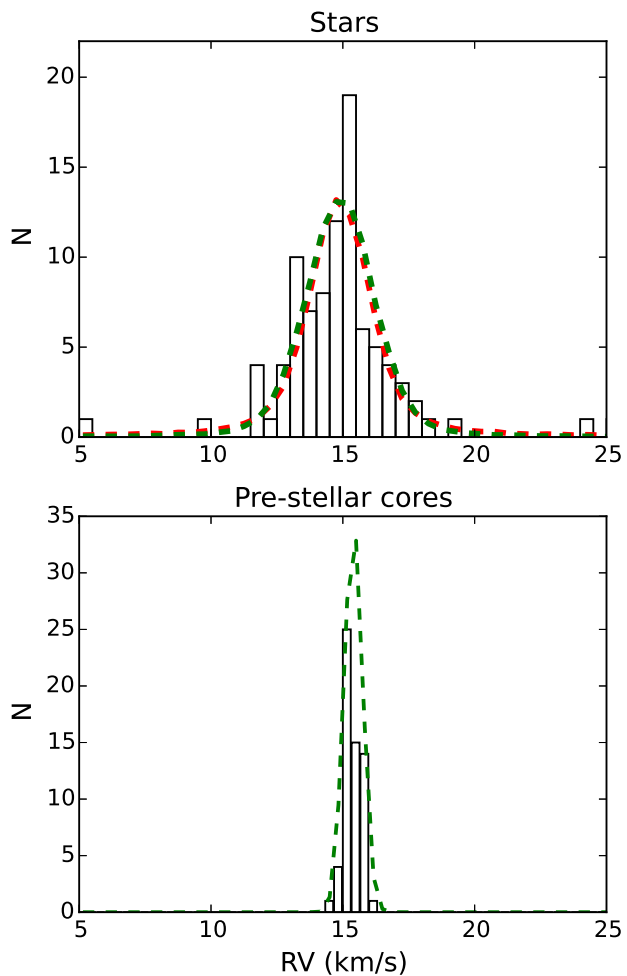


Fig. 9. The top panel shows the RV distribution of the full sample of cluster members observed by GES. The red and green dashed lines describe the best fit models with a gaussian broadened by the measurement errors and the velocity offsets due to binaries, assuming a fixed binary fraction ($f_{bin} = 0.5$) and letting the binary fraction free to vary, respectively. The bottom panel shows the distribution of the RVs of the pre-stellar cores measured by Tsitali et al. (2015) from the C^{18}O (2 – 1) molecular transition, with overlying the best fit models with a gaussian function.

5.1. Radial Velocity dispersion

The RV distribution of the cluster members is shown in the top panel of Fig. 9. We modeled the distribution using a maximum-likelihood method developed by Cottaar et al. (2012b) and already used in several works (e.g., Jeffries et al. 2014; Foster et al. 2015; Sacco et al. 2015), which allows us to properly take into account the errors on each star and the presence of binaries. Specifically, we assume that the stellar RVs have an intrinsic Gaussian distribution (with mean v_c and standard deviation σ_c) broadened by the measurement uncertainties and the velocity offsets due to binary orbital motion. The distribution of the offsets is calculated numerically by a code⁵ developed by Cottaar et al. (2012b), considering three different assumptions: a) binary periods follow a log-normal distribu-

⁵ Available online at <https://github.com/MichielCottaar/velbin>.

Table 4. Parameters obtained from the fits of the RV distributions with 1σ errors .

	v_c (km s ⁻¹)	σ_c (km s ⁻¹)	f_{bin}	α (km s ⁻¹ deg ⁻¹)	β (km s ⁻¹ deg ⁻¹)	$\ln(L_{max})$
Cha I (fit 1)	14.88±0.15	0.94±0.15	0.5	0	0	-183.68
Cha I (fit 2)	14.90±0.15	1.11±0.11	0.18±0.11	0	0	-181.52
Cha I (fit 3)	14.87±0.15	1.08±0.14	0.17±0.11	-0.21±0.13	0.11±0.24	-180.16
Cha I N	15.29±0.22	0.95±0.18	0.18	0	0	-51.94
Cha I S	14.36±0.22	0.87±0.24	0.18	0	0	-78.06
Cha I Outer	14.98±0.30	1.17±0.28	0.18	0	0	-47.05

tion with mean period 5.03 and dispersion 2.28 in \log_{10} days (Raghavan et al. 2010); b) the secondary to primary ratio (q) follows a power-law $\frac{dN}{dq} \sim q^{-0.5}$ for $0.1 < q < 1$ (Reggiani & Meyer 2011); c) the distribution of eccentricities is flat between 0 and a maximum value which depends on the period according to Eqn. 6 from Parker & Goodwin (2009).

We performed two fits. In the first one (fit 1 in Table 4), we kept the fraction of binaries fixed at $f_{bin} = 0.5$, while in the second one (fit 2 in Table 4), it was left free to vary. In both fits, we consider only stars in the range $0 < RV < 30$ km s⁻¹, since stars outside this range are either binaries or stars with a miscalculated RV due to the presence of strong emission lines. The parameters derived by the two fits are reported in the first two rows of Table 4 and the best fit functions are plotted in Fig. 9. Since the two models are nested, to evaluate the parameters of which fit to adopt we can perform a likelihood-ratio test. This gives a probability $P(L_{fit1}/L_{fit2}) = 3.8\%$, which indicates, with a marginal level of significance, that the fit with a fixed binary fraction can be rejected and the parameters from the second fit can be adopted.

Our results are in agreement with previous estimates of the central cluster velocity and of the velocity dispersion from Joergens (2006) ($v_c = 14.7$ km s⁻¹ and $\sigma_c = 1.3$ km s⁻¹).

5.2. Radial velocity gradient

To investigate the presence of a RV gradient in the stellar population, we fitted the RV distribution with the same function discussed in the previous section, but instead of considering the mean cluster velocity v_c as a single free parameter, we assumed that the velocity $v_c = v_{c0} + \alpha\Delta RA + \beta\Delta DEC$, where ΔRA and ΔDEC are the RA and DEC shifts of each star with respect to a fixed position calculated as the median of the star positions, and v_{c0} , α and β are free parameters of the fit together with σ_c , which is assumed to be constant over the whole region. The result of this fit is reported in Table 4 (fit 3). The parameters v_{c0} and σ_c are in agreement with the results found with the previous fits and the components of the RV gradient α and β are consistent within two standard deviations with zero. Since the function used for fit2 is the same as fit 3, when we fix the parameters α and β to zero, we can use the likelihood-ratio test to compare the model with and without a gradient. The probability $P(L_{max}(fit2)/L_{max}(fit3)) = 26\%$, so we conclude that there is no evidence of the presence of a RV gradient in the cluster.

5.3. Kinematical properties of the sub-clusters

Luhman (2007) suggested that Cha I is composed of two sub-clusters with different star formation histories. To understand if these two populations have different kinematical properties we divided our sample in three groups: 1) one composed of 29 stars located within the upper circle drawn in Fig.2 with a blue dashed

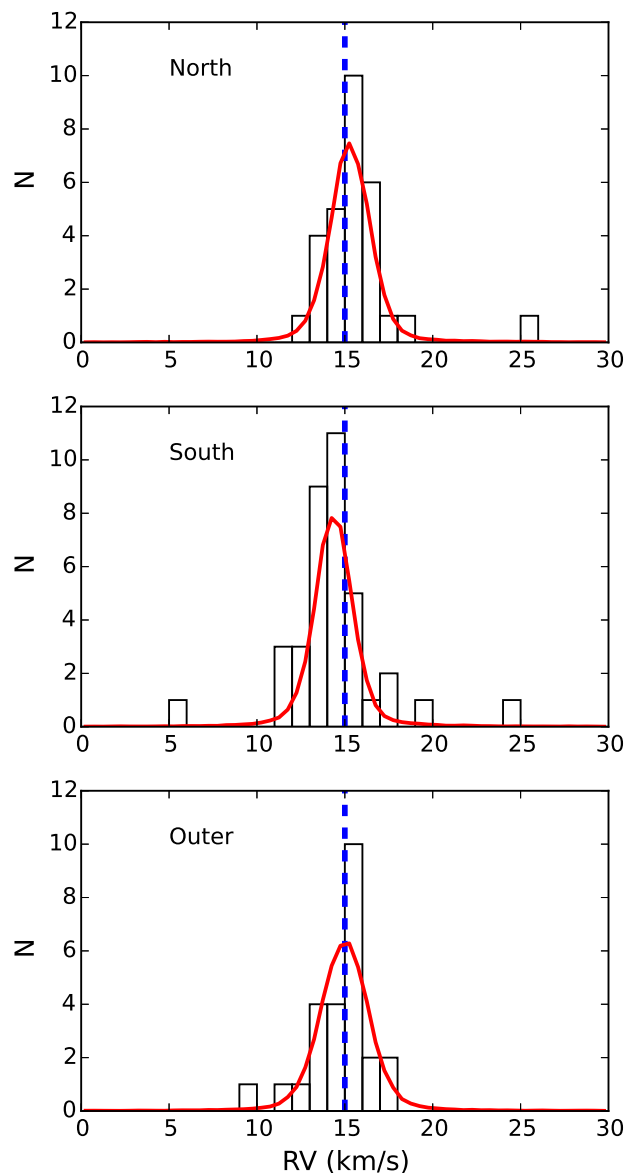


Fig. 10. The three panels from the top to the bottom show the RV distributions of the North, the South subclusters, and the stars dispersed in the outer regions, respectively. The dashed blue line marks the velocity of 15 km s⁻¹ and the red line is the best-fit the distribution with the same model used for the full sample.

line, which approximately defines the boundary of the northern part of the cloud; 2) one composed of 37 stars within the lower circle, which defines the boundary of the southern cloud; 3) one composed of 25 stars located in the outer regions..

The RV distributions of the three samples are shown in Fig. 10 and the results from the fits of the distributions are reported in the last rows of Table 4. The central RVs of the two clusters concentrated around the clouds differ by $\sim 1 \text{ km s}^{-1}$ at 2σ level of significance and on the basis of a Kolmogorov-Smirnov test the probability that the two distributions are part of the same population is $<1\%$. The kinematical properties of the stars located in the outer regions are closer to what found for the northern stars, suggesting that the majority of the outer stars belong to the northern cluster. This is consistent with the hypothesis suggested by Luhman (2007), that the northern cluster started to form earlier and, therefore, it is going through a more advanced stage of its evolution. Lopez Martí et al. (2013) already tried to kinematically separate the two subclusters using proper motions, finding no evidence of different velocities. They do not report any upper limit on the velocity separation between subclusters, so it is difficult to compare their data with our result. However, the precision of proper motions used for their work is lower than RVs from the Gaia-ESO Survey.

6. Discussion

The main goal of this work is to study the physical processes leading to the formation and the dynamical evolution of small star clusters. In the next sections, we will compare the structural and dynamical properties of the stellar populations in Cha I with the properties of pre-stellar cores and with some numerical models describing the early stages of the star cluster evolution.

6.1. Structural properties

Using N-body simulations, Parker et al. (2014) studied the evolution of the level of substructure and the mass segregation in young star clusters. In Fig. 11 we compare the results of the simulations performed by Parker et al. (2014) for clusters with high ($n_{\text{stars}} \sim 1000 \text{ stars pc}^{-1}$) and low ($n_{\text{stars}} \sim 100 \text{ stars pc}^{-1}$) stellar density with the structural properties of Cha I derived in Sect. 4. The simulations differ for the initial virial ratio (α_{vir}) and the initial fractal dimension D , which indicated the level of substructure ($D = 1.6$ is a highly substructured cluster and $D = 3.0$ is a roughly uniform sphere). The figure shows the initial conditions of the simulated clusters ($t = 0 \text{ Myr}$) and their status after 2 Myr.

None of the simulated clusters with high stellar density reproduce the structural properties of Cha I, with the exception of the case of a supervirial cluster with no substructure. However, even in this case, the properties of the simulated clusters at the initial conditions are not consistent with the properties of the pre-stellar cores. This results is not surprising, since the stellar density of the simulated clusters is much higher than the observed density in Cha I of both stars and pre-stellar cores, and further supports the hypothesis that Cha I did not form in a high density environment, in contrast to the hypothesis advanced by Marks & Kroupa (2012).

The properties of the low-density simulated clusters are much closer to the properties of Cha I. In particular, virial and supervirial simulated clusters are consistent with the overall properties of the cluster after 2 Myr of dynamical evolution. Furthermore, the simulations with a high level of substructure at $t = 0 \text{ Myr}$ are consistent with the properties of embedded stars and

pre-stellar cores, if we assume that these represent the properties of the cluster at its formation. Otherwise, according to the simulations, for a cluster which is initially sub-virial we should observe a level of mass segregation after 2 Myr, which we do not observe in Cha I.

To summarize, according to this analysis, Cha I formed in a low-density environment with a virial ratio $\alpha_{\text{vir}} \geq 0.5$ and a high level of substructures. It has erased substructure due to dynamical interactions and will likely disperse in the Galactic field. A similar scenario has been proposed for the more evolved cluster Gamma Velorum (Jeffries et al. 2014; Mapelli et al. 2015; Sacco et al. 2015). It would be interesting to perform a direct measurement of the virial ratio in Cha I. However, due to the highly asymmetric structure of both the stellar and the gas component of the cluster it is difficult to estimate the virial ratio without any information about its structure along the line of sight. This information will be provided by the astrometric mission *Gaia* for most of the optically visible stars.

It is worth mentioning a few caveats concerning the comparisons between the simulations performed by Parker et al. (2014) and our results: a) as proven by the large area of the parameter space covered by the simulations in each panels of Fig. 11, N-body simulations, especially of low-density clusters, are partially degenerate namely, the same initial conditions may lead to clusters with very different properties in the Q vs. Σ_{LDR} and Q vs. Λ_{MSR} plots. In particular, the comparison between our results and models give strong constrains about the initial density of the cluster, but less stronger constrains about other critical properties, like the initial virial ratio. The analysis of other young star clusters similar to Cha I and the definition of new diagnostics of the dynamical status of star clusters using also kinematic data can help to overcome this limitation; b) N-body simulations do not include the presence of gas, which in the case of Cha I is the main component of the potential energy of the systems. The effects of the gas in the evolution of star clusters is a very debated topic. Some authors (e.g., Kruijssen et al. 2012) suggest that the influence of the gas in the cluster dynamics is negligible, but a direct comparison between simulations describing in a consistent way the evolution of gas and stars is required to provide final answers to this issue.; c) the simulations discussed in this paper assume that all the stars are coeval, while the age spread in Cha I is larger than its median age. The origin of the age spread in young clusters is not clear and is not reproduced by any of the state-of-the-art simulations. More sophisticated simulations and precise and complete data are required to fully understand the star formation history of clusters like Cha I.

6.2. Kinematical properties

As shown in Fig. 9, the most clear result of the kinematical analysis is the large discrepancy between the velocity dispersion of the stars ($\sigma_{\text{stars}} = 1.10 \pm 0.15 \text{ km s}^{-1}$) and that of pre-stellar cores ($\sigma_{\text{cores}} \sim 0.3 \text{ km s}^{-1}$) derived from submillimeter observations of molecular transitions by Tsitali et al. (2015). As noted in Tab. 4, the velocity dispersion of the stellar component does not depend on the sample of stars used for the fit. In fact, the two sub-clusters around the molecular cloud and the sample of stars located in the outer region have similar velocity dispersions, which are in all cases much higher than the dispersion measured for the pre-stellar cores. A similar discrepancy between the pre-stellar cores and the stars has been observed in the ρ Oph star forming region, in the young cluster NGC 1333, and in Orion. The velocity dispersion of the stellar component in ρ Oph ($\sigma_{\text{stars}} = 1.14 \pm 0.35 \text{ km s}^{-1}$) was derived from the

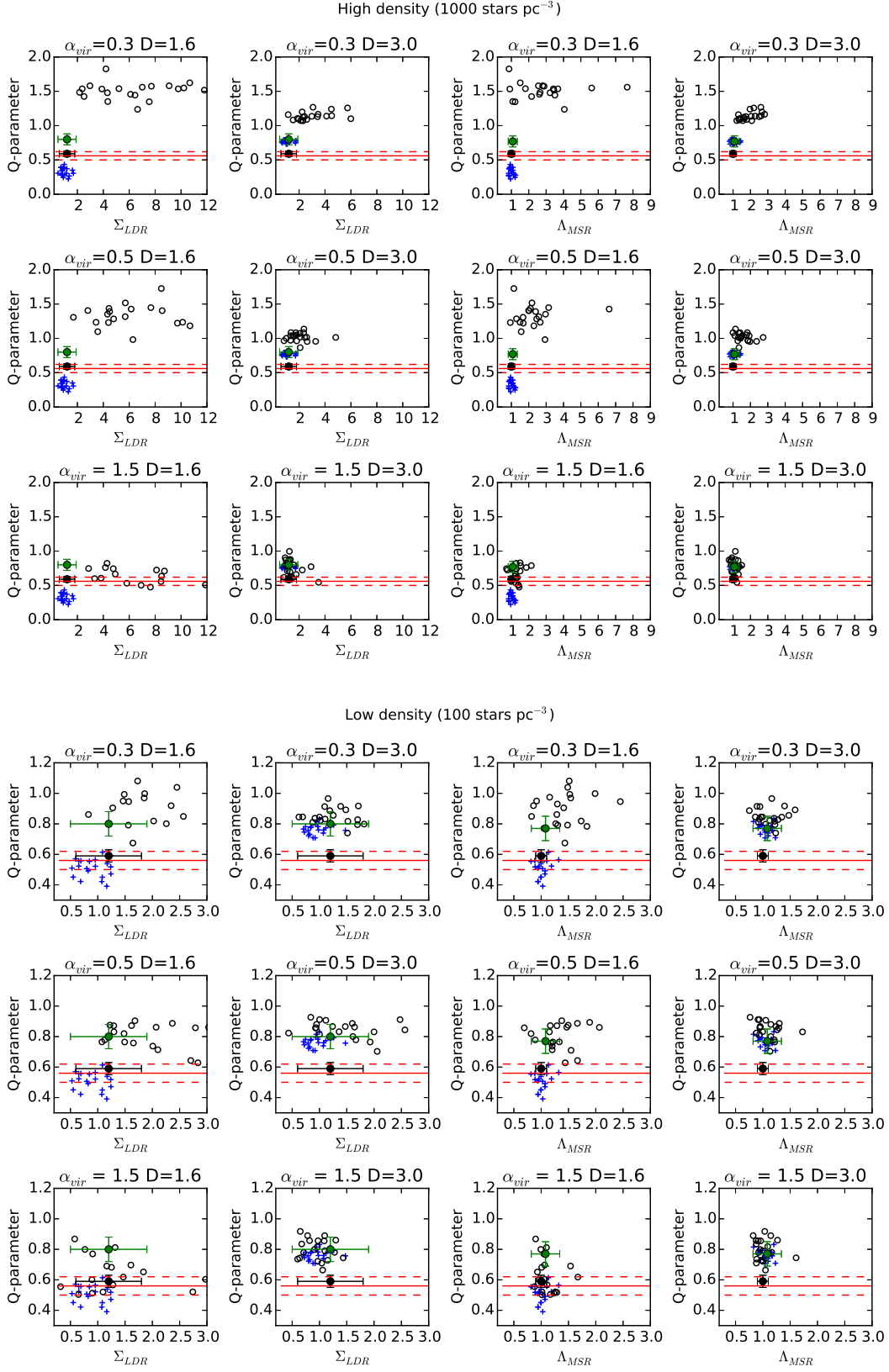


Fig. 11. Comparison between the observed structural properties of Cha I and simulated clusters with high ($n_{\text{stars}} = 1000 \text{ stars pc}^{-3}$, top panels) and low ($n_{\text{stars}} = 100 \text{ stars pc}^{-3}$, bottom panels) initial stellar density from Parker et al. (2014). The panels on the left and on the right show the Q -parameter as function of Σ_{LDR} (i.e., the ratio between the median superficial density of the most massive stars and the rest of the sample), and Λ_{MSR} (for the ten most massive stars), respectively. The simulations differ for the initial virial ratio α_{vir} and the initial level of substructure ($D = 1.6$ is a highly substructured cluster and $D = 3.0$ is a roughly uniform sphere). Blue crosses and black circles represent the simulated clusters at the initial conditions and after 2 Myr evolution, respectively. Green and black dots represent the properties of Cha I for the full sample of members and only for the stars in the embedded region within the smallest ellipse in Fig. 6, respectively. The red lines trace the Q -parameter estimated for the pre-stellar cores with errors.

Gaia-ESO observations of the optically visible stars around the main cloud L1688 by Rigliaco et al. (2016), who suggested that the cluster is bound and in virial equilibrium, while the velocity dispersion of the cores ($\sigma_{\text{cores}} \sim 0.4 \text{ km s}^{-1}$) was estimated by André et al. (2007), who suggested that the cores are sub-virial. The kinematical properties of both the cores and the stars of NGC 1333 have been analysed by Foster et al. (2015), who also found that the stars are virial ($\sigma_{\text{stars}} = 0.92 \pm 0.12 \text{ km s}^{-1}$), while the cores ($\sigma_{\text{cores}} \sim 0.5 \text{ km s}^{-1}$) are sub-virial. They suggested that the discrepancy between stars and cores can be due either to the magnetic field having a strong influence on the cores and/or to the global collapse of the cluster after the protostellar phase. A similar conclusion has been obtained by N-body simulations carried out by Parker & Wright (2016), who found that clusters starting as subvirial undergo cool collapse, so the dynamical interaction among stars quickly inflates the distribution. However, for small low-density clusters Parker & Wright (2016) found a lower velocity dispersion ($\sigma \sim 0.5 \text{ km s}^{-1}$) than observed. This discrepancy could be associated to the lack of gas in the N-body simulations, since the presence of a significant amount of gas reduces the virial ratio and leads to the collapse of clusters with higher velocity dispersion than in the case without gas (Leigh et al. 2014, Mapelli et al., in prep.). The morphology and the kinematic of gas, protostars and pre-main sequence stars has been studied in Orion A by Stutz & Gould (2016). They propose that protostars are ejected from the filaments due to magnetically induced transverse waves. This slingshot-like mechanism is responsible of the velocity discrepancy between young stars and protostars still within the filaments.

The second result of our kinematical analysis is the discrepancy ($\sim 1 \text{ km s}^{-1}$) between the central velocities of the two sub-clusters located around the northern and southern clouds. This is not surprising, since Luhman (2007) suggested that Cha I is composed of two components with different star formation histories. Furthermore, recent studies show that multiple populations (e.g., Jeffries et al. 2014; Sacco et al. 2015) and RV gradients (e.g., Tobin et al. 2015) are common in young clusters and star forming regions. According to the submillimeter observations, the mean velocities of the cores in the north and the south clusters also differ by $\sim 0.3 \text{ km s}^{-1}$. However, the discrepancy is in the opposite direction with respect to what we found for the stars, namely the cores in the south have a higher redshift than in the north. The reason of this anti-correlation between stars and cores is not clear, but this result supports a scenario where the dynamics of the cores is independent from the dynamics of the stellar populations.

7. Conclusions

In this work we present a new analysis of the spectroscopic parameters derived from the Gaia-ESO Survey observations of the young cluster Cha I aimed at investigating the structural and dynamical properties of the cluster. We obtained the following main results.

1. An evident discrepancy between the velocity dispersion of the stellar component ($\sigma_{\text{star}} = 1.10 \pm 0.15 \text{ km s}^{-1}$) derived by Gaia-ESO spectra and the dispersion ($\sigma_{\text{cores}} = 0.3 \text{ km s}^{-1}$) of pre-stellar cores derived by submillimeter observations. A similar discrepancy has been observed in the young embedded clusters ρ Oph, NGC 1333 and in Orion. The origin of such a large discrepancy is not clear. It could be related to the effect of the magnetic field on the protostars or the filaments where they form, or to two-body stellar dynamical interactions following the cluster formation. We will further investigate this issue in a forthcoming paper (Mapelli et al. in prep.).
2. Analysing independently the RV distributions of the two sub-clusters located around the two main molecular clouds, we found that the central RVs differ by $\sim 1 \text{ km s}^{-1}$. This result supports the evidence found by Luhman (2007) that Cha I is composed of two sub-clusters with different star formation histories and a scenario where young clusters do not form as monolithic systems but from the merging of smaller subsystems.
3. A new membership analysis based on three independent spectroscopic criteria led to the confirmation of all the previously known members, for which new astrophysical parameters from the Gaia-ESO Survey are available, and to the discovery of six new members in Cha I and one new member of the ϵ Cha association, which are all located in the outer part of the cluster.
4. The level of substructure of the cluster measured using the Q -parameter defined by Cartwright & Whitworth (2004) depends on the sample used for the calculation. If we consider only the stars in the inner region the value of Q indicates that the cluster is highly substructured, while if we take into account the full sample of members the spatial distribution of the cluster is consistent with a random sample. It is not clear if this trend has a physical origin or if it is the result of a bias due to differential extinction in the inner region of the cluster or incomplete target selection in the outer region.
5. As observed in other low-mass young star clusters, Cha I is not mass-segregated and its superficial density follows a log-normal distribution, with the exception of its high mass end, which follows a steeper trend.
6. The comparison between the observed structural properties of Cha I and the results of N-body simulations performed by Parker et al. (2014) suggests that the cluster formed as highly substructured, and virial or supervirial. However, discrepancies between the simulated clusters and Cha I (e.g., the lack of gas in the simulated clusters) may affect this comparison.

Acknowledgements. These data products have been processed by the Cambridge Astronomy Survey Unit (CASU) at the Institute of Astronomy, University of Cambridge, and by the FLAMES/UVES reduction team at INAF/Osservatorio Astrofisico di Arcetri. These data have been obtained from the Gaia-ESO Survey Data Archive, prepared and hosted by the Wide Field Astronomy Unit, Institute for Astronomy, University of Edinburgh, which is funded by the UK Science and Technology Facilities Council. This work was partly supported by the European Union FP7 programme through ERC grant number 320360 and by the Leverhulme Trust through grant RPG-2012-541. We acknowledge the support from INAF and Ministero dell' Istruzione, dell' Università e della Ricerca (MIUR) in the form of the grant "Premiale VLT 2012". This work has been partially supported by PRIN-INAF-2014. The results presented here benefit from discussions held during the Gaia-ESO workshops and conferences supported by the ESF (European Science Foundation) through the GREAT Research Network Programme. MM acknowledges financial support from the Italian Ministry of Education, University and Research (MIUR) through grant FIRB 2012 RBF12P2M1F, from INAF through grant PRIN-2014-14, and from the MERAC Foundation.

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Table 1. Members of the cluster observed by the Gaia-ESO Survey

Cname	RA (J2000)	DEC (J2000)	RV ^a (km s ⁻¹)	T_{eff} (K)	γ^b	EW(Li) (mÅ)	$H\alpha 10\%^c$ (km s ⁻¹)	Inst. ^d	memb ^e
10550964-7730540	163.79017	-77.51500	16.83±0.90	-	-	725±23	128±5	G	Y
10555973-7724399	163.99887	-77.41108	-	3640±300	-	120±58	441±17	U	Y
10561638-7630530	164.06825	-76.51472	12.56±2.01	-	-	-	215±5	G	Y
10563044-7711393	164.12683	-77.19425	15.69±0.27	4351±505	0.968±0.005	300±17	390±8	G	Y
10563146-7618334 ^f	164.13108	-76.30928	15.41±0.54	3319±48	-	597±11	142±4	G	N
10574219-7659356	164.42579	-76.99322	16.14±0.26	3452±183	0.895±0.004	571±16	272±6	G	Y
10575376-7724495	164.47400	-77.41375	15.71±0.32	3426±88	0.857±0.013	632±10	138±8	G	N
10590108-7722407	164.75450	-77.37797	15.64±0.40	4135±125	-	375±19	375±8	U	Y
10590699-7701404	164.77912	-77.02789	17.93±0.26	4981±260	-	390±1	440±9	G	Y
11004022-7619280	165.16758	-76.32444	15.73±0.42	-	-	584±18	230±5	G	Y
11011875-7627025	165.32812	-76.45069	13.27±0.26	3940±236	0.905±0.004	580±21	175±11	G	Y
11021927-7536576	165.58029	-75.61600	15.31±0.90	-	-	566±33	94±4	G	Y
11022491-7733357	165.60379	-77.55992	14.33±0.40	4519±174	-	471±12	382±8	U	Y
11023265-7729129	165.63604	-77.48692	15.35±0.31	3610±104	0.885±0.008	618±13	148±10	G	Y
11025504-7721508	165.72933	-77.36411	14.97±0.54	2900±262	0.919±0.013	174±11	322±10	G	Y
11035144-7540335	165.96433	-75.67597	9.66±1.45	-	-	443±73	144±6	G	N
11035682-7721329	165.98675	-77.35914	5.32±0.32	3443±95	0.865±0.011	616±25	135±2	G	Y
11045100-7625240	166.21250	-76.42333	13.77±0.40	4571±334	-	576±17	175±7	U	Y
11045285-7625514	166.22021	-76.43094	13.91±0.26	4031±267	0.911±0.003	583±18	120±3	G	Y
11045701-7715569	166.23754	-77.26581	-67.23±5.83	-	-	-	372±48	G	Y
11050752-7812063	166.28133	-78.20175	15.61±1.66	-	-	399±15	181±87	G	Y
11051467-7711290	166.31113	-77.19139	17.02±0.38	3433±200	0.892±0.010	525±10	241±13	G	Y
11052272-7709290	166.34467	-77.15806	13.34±1.57	-	-	-	126±5	G	Y
11052472-7626209	166.35300	-76.43914	15.23±0.28	3576±73	0.881±0.007	567±17	104±3	G	Y
11054300-7726517	166.42917	-77.44769	15.29±0.38	3223±98	0.944±0.007	542±8	138±5	G	Y
11055261-7618255	166.46921	-76.30708	15.00±0.55	4406±740	-	609±13	202±9	G	Y
11055780-7607489	166.49083	-76.13025	-	-	-	-	-	U	Y
11060011-7507252	166.50046	-75.12367	13.88±1.20	-	-	504±17	142±6	G	Y
11062555-7633418	166.60646	-76.56161	164.11±1.61	-	-	-	305±25	G	Y
11064180-7635489	166.67417	-76.59692	14.96±0.51	3273±16	0.928±0.016	332±14	416±10	G	Y
11064510-7727023	166.68792	-77.45064	14.69±0.40	4343±147	-	401±51	503±38	U	Y
11065733-7742106	166.73888	-77.70294	12.88±0.28	3436±62	0.901±0.007	606±20	155±7	G	Y
11065906-7718535	166.74608	-77.31486	-606.45±2.79	-	-	-	432±9	G	Y
11070919-7723049	166.78829	-77.38469	-	-	-	-	-	G	Y
11071148-7746394	166.79783	-77.77761	17.24±0.54	3709±133	-	640±8	288±7	G	Y
11071206-7632232	166.80025	-76.53978	16.59±0.25	3980±135	0.921±0.003	578±9	523±10	G	Y
11071330-7743498	166.80542	-77.73050	14.46±0.76	-	-	590±50	119±5	G	Y
11071622-7723068	166.81758	-77.38522	-	-	-	-	-	G	Y
11071915-7603048	166.82979	-76.05133	15.10±0.28	3607±114	0.884±0.008	595±32	371±8	G	Y
11072022-7738111	166.83425	-77.63642	14.91±0.50	3400±126	0.923±0.011	497±12	297±5	G	Y
11072040-7729403	166.83500	-77.49453	14.67±0.27	3296±64	0.921±0.004	594±4	115±5	G	Y
11072825-7652119	166.86771	-76.86997	15.12±0.27	3453±132	0.908±0.004	476±4	330±2	G	Y
11073519-7734493	166.89663	-77.58036	14.39±0.37	3397±20	0.896±0.012	620±17	110±4	G	Y
11073832-7747168	166.90967	-77.78800	13.82±0.42	3316±14	0.915±0.013	542±20	119±5	G	Y
11074245-7733593	166.92687	-77.56647	12.40±1.68	-	-	508±18	332±10	G	Y
11074366-7739411	166.93192	-77.66142	15.02±0.28	4140±157	0.969±0.005	362±14	351±7	G	Y
11075225-7736569	166.96771	-77.61581	11.98±1.29	-	-	564±65	128±5	G	Y
11075588-7727257	166.98283	-77.45714	17.36±0.40	4752±162	-	524±14	116±22	U	Y
11075792-7738449	166.99133	-77.64581	24.26±2.80	-	-	218±21	612±12	G	Y
11075809-7742413	166.99204	-77.71147	19.03±0.41	3423±116	0.901±0.012	191±34	501±10	G	Y
11075993-7715317	166.99971	-77.25881	11.66±3.07	-	-	-	290±17	G	Y
11080002-7717304	167.00008	-77.29178	15.62±0.70	-	-	574±18	112±4	G	Y
11080148-7742288	167.00617	-77.70800	10.65±0.40	-	-	-	-	U	Y
11080297-7738425	167.01237	-77.64514	17.69±0.50	4813±1391	0.955±0.009	316±9	432±9	G	Y
11081509-7733531	167.06287	-77.56475	14.13±0.28	4798±444	0.981±0.004	479±4	384±2	G	Y
11081648-7744371	167.06867	-77.74364	15.33±0.26	3409±72	0.914±0.005	587±4	114±3	G	Y
11081703-7744118	167.07096	-77.73661	14.64±0.58	3114±66	0.939±0.014	551±20	112±4	G	Y
11082237-7730277	167.09321	-77.50769	11.72±1.53	-	-	390±18	413±9	G	Y
11082410-7741473	167.10042	-77.69647	13.17±0.80	-	-	584±44	123±5	G	Y

Table 1. continued.

Cname	RA (J2000)	DEC (J2000)	RV ^a (km s ⁻¹)	T_{eff} (K)	γ^b	EW(Li) (mÅ)	$H\alpha 10\%^c$ (km s ⁻¹)	Inst. ^d	memb ^e
11083905-7716042	167.16271	-77.26783	13.10±0.27	4264±303	0.948±0.003	564±3	474±10	G	Y
11083952-7734166	167.16467	-77.57128	13.18±0.84	-	-	417±16	318±8	G	Y
11084069-7636078	167.16954	-76.60217	13.06±0.26	3874±183	0.918±0.003	633±6	202±14	G	Y
11084296-7743500	167.17900	-77.73056	-	-	-	-	-	G	Y
11085090-7625135	167.21208	-76.42042	14.16±0.44	3251±75	0.945±0.011	413±23	308±7	G	Y
11085242-7519027	167.21842	-75.31742	14.09±0.28	3558±161	0.928±0.004	564±6	275±11	G	N
11085367-7521359	167.22363	-75.35997	14.79±0.27	3539±280	0.938±0.004	323±4	376±0	G	Y
11085422-7732115	167.22592	-77.53653	13.43±0.56	-	-	560±21	112±4	G	Y
11085464-7702129	167.22767	-77.03692	15.41±0.29	4025±275	0.916±0.005	409±47	466±20	G	Y
11090512-7709580	167.27133	-77.16611	13.08±4.23	-	-	647±44	238±82	G	Y
11090915-7553477	167.28813	-75.89658	14.09±0.32	3614±100	0.891±0.009	558±8	120±3	G	N
11091172-7729124	167.29883	-77.48678	14.76±0.40	3905±117	-	646±17	119±5	U	Y
11091297-7729115	167.30404	-77.48653	13.26±0.26	3763±110	0.894±0.004	610±4	131±4	G	Y
11091380-7628396	167.30750	-76.47767	14.92±0.95	3294±45	-	617±24	281±20	G	Y
11091769-7627578	167.32371	-76.46606	15.06±0.40	4541±171	-	559±10	133±4	U	Y
11091812-7630292	167.32550	-76.50811	16.15±0.46	3903±41	0.978±0.016	511±37	303±7	G	Y
11092378-7623207	167.34908	-76.38908	-	3990±123	-	65±8	566±20	U	Y
11092855-7633281	167.36896	-76.55781	-	-	-	-	-	G	Y
11094006-7628392	167.41692	-76.47756	15.20±0.26	3939±219	0.901±0.005	643±7	154±11	G	Y
11094525-7740332	167.43854	-77.67589	12.67±1.43	-	-	414±12	200±106	G	Y
11094621-7634463	167.44254	-76.57953	14.34±0.58	3431±35	0.909±0.015	327±14	375±9	G	Y
11095003-7636476	167.45846	-76.61322	12.79±0.83	-	-	14±5	-	G	Y
11095340-7634255	167.47250	-76.57375	-536.10±0.73	-	-	269±8	584±12	G	Y
11095407-7629253	167.47529	-76.49036	15.16±0.70	-	-	419±19	321±14	G	Y
11095873-7737088	167.49471	-77.61911	-635.00±2.41	3376±113	-	240±16	465±1	G	Y
11100010-7634578	167.50042	-76.58272	25.00±0.30	4979±191	-	236±8	554±11	G	Y
11100369-7633291	167.51538	-76.55808	15.48±0.42	3891±45	0.930±0.012	285±11	440±12	G	Y
11100469-7635452	167.51954	-76.59589	15.65±0.26	4088±227	0.946±0.003	480±9	437±9	G	Y
11100704-7629377	167.52933	-76.49381	14.04±0.40	4624±168	-	467±12	448±9	U	Y
11101141-7635292	167.54754	-76.59144	17.14±0.35	4470±174	0.975±0.006	551±6	438±13	G	Y
11101153-7733522	167.54804	-77.56450	13.17±0.41	3302±46	0.915±0.008	549±24	122±4	G	Y
11102852-7716596	167.61883	-77.28322	14.97±0.28	3287±71	0.933±0.007	600±20	107±13	G	Y
11103481-7722053	167.64504	-77.36814	14.76±1.87	-	-	572±22	-	G	Y
11103801-7732399	167.65838	-77.54442	13.70±0.36	4842±377	-	399±4	701±14	G	Y
11104959-7717517	167.70663	-77.29769	-435.00±2.29	3327±198	-	252±19	498±10	G	Y
11105076-7718032	167.71150	-77.30089	14.83±0.38	3339±16	0.884±0.013	617±19	119±5	G	Y
11105333-7634319	167.72221	-76.57553	49.84±0.97	3287±171	-	67±4	425±9	G	Y
11105597-7645325	167.73321	-76.75903	18.02±1.12	-	-	321±10	257±8	G	Y
11112260-7705538	167.84417	-77.09828	13.75±0.31	3342±66	0.920±0.008	584±7	124±11	G	Y
11113474-7636211	167.89475	-76.60586	14.91±0.25	3943±258	0.905±0.002	663±2	150±18	G	Y
11113965-7620152	167.91521	-76.33756	16.17±0.25	3468±205	0.949±0.002	223±33	309±7	G	Y
11114632-7620092	167.94300	-76.33589	16.38±0.40	4617±207	-	461±10	373±8	U	Y
11115400-7619311	167.97500	-76.32531	15.50±0.25	3738±80	0.901±0.002	653±23	225±4	G	Y
11120327-7637034	168.01362	-76.61761	13.94±0.85	3247±15	-	556±109	164±7	G	Y
11120984-7634366	168.04100	-76.57683	15.20±0.30	3275±51	0.962±0.005	406±5	338±7	G	Y
11122441-7637064	168.10171	-76.61844	15.16±0.29	4899±326	-	414±5	389±8	G	Y
11124210-7658400	168.17542	-76.97778	15.04±0.26	4365±235	0.955±0.003	501±24	180±7	G	Y
11124268-7722230	168.17783	-77.37306	14.18±0.40	5196±81	-	380±7	-	U	Y
11124299-7637049	168.17912	-76.61803	14.46±0.40	4706±221	-	475±9	107±6	U	Y
11130450-7534369	168.26875	-75.57692	15.15±0.50	3321±50	0.929±0.011	583±31	291±9	G	N
11132446-7629227	168.35192	-76.48964	15.25±0.27	3374±90	0.877±0.006	622±9	233±10	G	Y
11132737-7634165	168.36404	-76.57125	13.87±0.25	3775±106	0.874±0.002	559±11	151±16	G	Y
11132971-7629012	168.37379	-76.48367	16.40±0.27	3339±67	0.913±0.006	633±6	106±3	G	Y
11133356-7635374	168.38983	-76.59372	16.91±0.63	3299±47	-	522±67	254±14	G	Y
11141565-7627364	168.56521	-76.46011	16.92±0.42	3384±124	-	598±7	248±5	G	Y
11142454-7733062	168.60225	-77.55172	16.01±1.08	-	-	342±29	396±15	G	Y
11145031-7733390	168.70962	-77.56083	14.44±0.25	3729±100	0.889±0.003	636±3	137±4	G	Y
11182024-7621576	169.58433	-76.36600	13.82±0.40	4317±127	-	539±18	116±12	U	Y
11213079-7633351	170.37829	-76.55975	14.99±0.30	3632±102	0.910±0.005	628±5	147±16	G	N
11242981-7554237	171.12421	-75.90658	11.56±0.50	3300±46	-	495±10	163±5	G	Y

Table 1. continued.

Cname	RA (J2000)	DEC (J2000)	RV ^a (km s ⁻¹)	T _{eff} (K)	γ ^b	EW(Li) (mÅ)	Hα10% ^c (km s ⁻¹)	Inst. ^d	memb ^e
11291261-7546263	172.30254	-75.77397	15.20±0.40	4818±96	-	477±12	130±5	U	Y

^(a)

For spectra with a signal-to-noise ratio lower than three we did not report any velocity.

^(b) Empirical gravity indicator defined by Damiani et al. (2014).

^(c) Width at 10% of the peak of the Hα line.

^(d) The letters "G" and "U" indicate GIRAFFE and UVES, respectively.

^(e) The letters "Y" and "N" indicate the star is a known member or not, respectively

^(f) This star is likely a member of the ε Cha association.

Table 2. Known cluster members from the literature and new member used to study the structural properties of the cluster.

RA (J2000)	DEC (J2000)	T_{eff} (K)	$Log(L_{Bol}/L_{\odot})$	prov ^a	Mass M_{\odot}
161.65812	-77.60097	7200	0.95	L	1.64
163.15392	-74.67464	3161	-1.00	L	0.14
163.41575	-77.20939	3451	-1.49	L	0.30
163.79017	-77.51500	3198	-1.08	L	0.15
163.99887	-77.41108	3640	-0.36	G	0.37
163.99887	-77.41108	3640	-0.36	G	0.37
164.06825	-76.51472	3044	-1.51	L	0.09
164.12683	-77.19425	4350	0.06	G	0.85
164.42579	-76.99322	3451	-0.33	G	0.28
164.47400	-77.41375	3426	-0.94	G	0.28
164.52487	-77.19725	3091	-1.66	L	0.10
164.56987	-77.28806	5250	0.28	L	1.43
164.75450	-77.37797	4135	-0.04	G	0.68
164.77912	-77.02789	4980	0.66	G	1.71
165.16758	-76.32444	3306	-1.11	L	0.20
165.30708	-77.37742	3091	-1.60	L	0.10
165.32812	-76.45069	3940	-0.24	G	0.56
165.58029	-75.61600	3198	-1.21	L	0.15
165.60379	-77.55992	4519	0.22	G	1.01
165.60875	-75.04464	3161	-1.12	L	0.14
165.63604	-77.48692	3610	-0.73	G	0.40
165.67429	-77.40681	3125	-1.31	L	0.12
165.92442	-77.44778	3058	-1.44	L	0.10
165.94850	-77.33231	3125	-0.74	L	0.13
165.98675	-77.35914	3442	-0.56	G	0.29
166.01771	-76.65911	3234	-1.35	L	0.16
166.03787	-76.45536	4350	-0.02	L	0.88
166.09479	-77.30222	4060	-1.89	L	-
166.17742	-77.69919	3270	-1.03	L	0.19
166.21250	-76.42333	4571	-0.26	G	1.01
166.22021	-76.43094	4031	-0.51	G	0.72
166.28133	-78.20175	3161	-1.11	L	0.14
166.31113	-77.19139	3433	-0.49	G	0.28
166.34467	-77.15806	3161	-1.46	L	0.13
166.35300	-76.43914	3576	-0.90	G	0.38
166.42917	-77.44769	3223	-0.93	G	0.17
166.46921	-76.30708	4405	-0.23	G	0.96
166.49083	-76.13022	13500	1.99	L	3.00
166.50042	-75.12367	3198	-1.33	L	0.14
166.56417	-77.36575	5770	1.20	L	2.53
166.60642	-76.56161	3091	-1.28	L	0.11
166.66437	-77.60144	3091	-2.13	L	0.10
166.67417	-76.59692	3273	-0.81	G	0.19
166.68108	-77.44286	3415	-0.38	L	0.26
166.68792	-77.45064	4343	0.12	G	0.83
166.73888	-77.70294	3435	-0.65	G	0.29
166.74608	-77.31486	3234	-0.96	L	0.17
166.74746	-75.51553	3091	-1.92	L	0.10
166.76667	-76.52917	3198	-1.24	L	0.14
166.79783	-77.77761	3708	-0.40	G	0.42
166.79921	-76.43058	3091	-1.74	L	0.10
166.80025	-76.53978	3979	-0.31	G	0.62
166.82979	-76.05133	3606	-0.53	G	0.38
166.83425	-77.63642	3399	0.28	G	-
166.83500	-77.49453	3295	-0.88	G	0.20
166.83642	-77.63536	5860	1.08	L	2.20
166.85179	-77.73025	3024	-1.44	L	0.09
166.86771	-76.86997	3452	-0.56	G	0.29
166.89663	-77.58036	3397	-1.03	G	0.26

Table 2. continued.

RA (J2000)	DEC (J2000)	T_{eff} (K)	$\text{Log}(L_{Bol}/L_{\odot})$	prov ^a	Mass M_{\odot}
166.90967	-77.78800	3316	-1.08	G	0.21
166.91000	-75.88108	3161	-1.47	L	0.13
166.92687	-77.56647	3091	-0.96	L	0.11
166.94208	-77.66914	3024	-1.43	L	0.09
166.94400	-76.25483	3024	-1.89	L	0.08
166.96771	-77.61581	3058	-1.15	L	0.10
166.99971	-77.25881	3024	-1.24	L	0.09
167.00617	-77.70800	3955	0.48	L	-
167.01371	-77.65483	10010	1.85	L	2.70
167.06287	-77.56475	4798	0.83	G	-
167.06867	-77.74364	3409	-0.65	G	0.27
167.07096	-77.73661	3113	-1.02	G	0.12
167.07900	-77.65472	3058	-1.26	L	0.10
167.08075	-77.53117	3161	-3.21	L	-
167.09325	-77.50769	3125	-1.06	L	0.12
167.10042	-77.69647	3058	-1.09	L	0.10
167.11042	-77.26528	3024	-1.72	L	0.09
167.16467	-77.57128	3024	-1.24	L	0.09
167.21208	-76.42042	3250	-1.30	G	0.17
167.21842	-75.31742	3558	-0.41	G	0.34
167.22363	-75.35997	3538	-0.50	G	0.33
167.22587	-77.53653	3091	-1.18	L	0.11
167.22767	-77.03692	4024	-0.35	G	0.67
167.22904	-76.54472	3058	-1.46	L	0.10
167.27133	-77.16611	3161	-1.17	L	0.13
167.28813	-75.89658	3614	-0.90	G	0.41
167.29883	-77.48678	3905	-0.17	G	0.52
167.30125	-77.48667	3560	-0.48	L	0.34
167.30125	-77.48667	3415	-0.58	L	0.27
167.30404	-77.48653	3762	-0.69	G	0.51
167.30750	-76.47767	3294	-1.16	G	0.19
167.32371	-76.46606	4541	0.24	G	1.03
167.34908	-76.38908	3990	0.19	G	-
167.37138	-76.98833	3091	-1.62	L	0.10
167.41692	-76.47756	3939	-0.15	G	0.54
167.43854	-77.67589	3024	-1.24	L	0.09
167.44254	-76.57953	3430	-0.52	G	0.28
167.45492	-77.52214	3058	-1.44	L	0.10
167.45846	-76.61322	10500	1.76	L	2.80
167.46925	-77.67633	3415	-2.54	L	-
167.47887	-76.58614	3024	-2.85	L	-
167.49471	-77.61911	3376	0.15	G	-
167.51954	-76.59589	4087	-0.20	G	0.68
167.52933	-76.49381	4623	0.30	G	1.12
167.54804	-77.56450	3301	-0.88	G	0.21
167.61883	-77.28322	3286	-1.12	G	0.19
167.70663	-77.29769	3327	-0.22	G	-
167.71150	-77.30089	3338	-0.88	G	0.23
167.72221	-76.57553	3287	-0.72	G	0.20
167.73321	-76.75903	3024	-0.89	L	-
167.79513	-76.69928	3488	-2.52	L	-
167.84417	-77.09828	3341	-0.98	G	0.23
167.89475	-76.60586	3942	-0.54	G	0.65
167.91521	-76.33756	3467	-0.23	G	0.28
167.94300	-76.33589	4617	0.47	G	1.10
167.97500	-76.32531	3738	-0.44	G	0.45
168.01362	-76.61761	3246	-1.11	G	0.17
168.01462	-77.43358	3161	-1.52	L	0.12
168.04100	-76.57683	3275	-0.64	G	0.20
168.10171	-76.61844	4898	0.42	G	1.49

Table 2. continued.

RA (J2000)	DEC (J2000)	T_{eff} (K)	$Log(L_{Bol}/L_{\odot})$	prov ^a	Mass M_{\odot}
168.11550	-76.73953	5410	0.70	L	1.97
168.12883	-76.74003	3705	-0.41	L	0.42
168.17542	-76.97778	4364	-0.33	G	0.93
168.17783	-77.37306	5196	0.62	G	1.88
168.17912	-76.61803	4706	-0.08	G	1.13
168.20254	-76.78517	3270	-1.28	L	0.18
168.26875	-75.57692	3321	-1.10	G	0.21
168.33383	-77.01789	3451	-0.74	L	0.30
168.35192	-76.48964	3373	-1.05	G	0.25
168.36404	-76.57125	3775	-0.52	G	0.49
168.37379	-76.48367	3338	-0.97	G	0.23
168.38983	-76.59372	3299	-1.08	G	0.20
168.56521	-76.46011	3383	-0.73	G	0.26
168.60225	-77.55172	3270	-0.96	L	0.19
168.60879	-77.55117	3024	-1.38	L	0.09
168.62108	-76.42775	3161	-1.43	L	0.13
168.70962	-77.56083	3728	-0.49	G	0.45
168.84083	-77.40117	3161	-1.02	L	0.14
168.99279	-77.48461	3198	-1.24	L	0.14
169.01196	-76.41481	3955	-2.47	L	-
169.40417	-77.07725	3778	-0.41	L	0.47
169.40800	-76.77203	3024	-1.89	L	0.08
169.46712	-76.49422	3198	-1.06	L	0.15
169.58154	-76.36703	3524	-0.82	L	0.34
169.58433	-76.36600	4317	-0.11	G	0.88
169.64079	-76.71781	3125	-1.01	L	0.12
169.64883	-79.59856	3161	-0.68	L	-
169.92558	-76.39239	3125	-1.48	L	0.11
170.37829	-76.55975	3632	-0.65	G	0.40
171.04942	-76.51181	3125	-1.33	L	0.12
171.12421	-75.90658	3299	-0.70	G	0.21
173.34696	-76.36922	3198	-0.70	L	0.17
173.45525	-76.31108	3198	-1.24	L	0.14
175.20696	-74.99428	3024	-1.57	L	0.09
175.86121	-78.07928	3125	-1.02	L	0.12

Notes. ^(a) The letters G and L indicate that the data used for deriving the stellar mass are retrieved from the literature and the GES archive, respectively.