

# A predicted astrometric microlensing event by a nearby white dwarf

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## ABSTRACT

We used the Tycho-*Gaia* Astrometric Solution catalogue, part of *Gaia* Data Release 1, to search for candidate astrometric microlensing events expected to occur within the remaining lifetime of the *Gaia* satellite. Our search yielded one promising candidate. We predict that the nearby DQ type white dwarf LAWD 37 (WD 1142-645) will lens a background star and will reach closest approach on 2019 November 11 ( $\pm 4$  d) with impact parameter  $380 \pm 10$  mas. This will produce an apparent maximum deviation of the source position of  $2.8 \pm 0.1$  mas. In the most propitious circumstance, *Gaia* will be able to determine the mass of LAWD 37 to  $\sim 3$  per cent. This mass determination will provide an independent check on atmospheric models of white dwarfs with helium rich atmospheres, as well as tests of white dwarf mass radius relationships and evolutionary theory.

**Key words:** gravitational lensing; micro – astrometry – white dwarfs.

## 1 INTRODUCTION

Einstein’s general theory of relativity predicts that light passing close to a massive object is deflected (Einstein 1916). This later led Einstein to the idea that massive objects can act as gravitational lenses and multiply image background sources (see e.g. Schneider, Ehlers & Falco 1992, for a review). In microlensing, the multiple images are typically separated by milliarcseconds and usually cannot be fully resolved, although the photometric brightening of the source and the astrometric deviation of the light centroid can be in principle detected. Paczynski (1995) noted that microlensing events can be predicted where high proper motion objects (lenses) approach the location of background sources. High proper motion stars are generally nearby and therefore have well-determined distances, which allows the lens mass to be found with high accuracy. The advent of data from the *Gaia* satellite, which is providing parallaxes and proper motions for over a billion stars in the Galaxy, makes it timely to look at Paczynski’s suggestion anew (e.g. Belokurov & Evans 2002; Harding et al. 2018).

This letter is structured as follows. In Section 2, the theory of mass determination via astrometric microlensing is described. Section 3 outlines the methods we used to search for lenses in the Tycho-*Gaia* Astrometric Solution (TGAS) catalogue, part of *Gaia* Data Release 1 (DR1) (Gaia Collaboration 2016a,b; Lindegren et al. 2016). Section 4 gives details of our best candidate event. Finally,

in Section 5, we sum up with an assessment of the feasibility of observing this event with *Gaia* and the *Hubble Space Telescope* (*HST*) and in Section 6 summarize the outlook for, and implications of, a precision measurement of the mass of LAWD 37.

## 2 MASS DETERMINATION BY ASTROMETRIC MICROLENSING

Microlensing occurs when a massive point-like foreground object (lens) focuses the light from a background point-like object (source). In the case of perfect alignment between the lens, source, and observer, a single Einstein ring with angular radius

$$\frac{\Theta_E}{\text{mas}} \approx 90.2 \left( \frac{M}{M_\odot} \right)^{1/2} \left( \frac{\text{pc}}{D_l} \right)^{1/2}, \quad \text{if } D_l \ll D_s \quad (1)$$

is formed. Here,  $M$  is the mass of the lens and  $D_l$  and  $D_s$  are the distances to the lens and source, respectively. We have assumed that the distance to the source is much greater than the distance to the lens (Paczynski 1986). When a lens encounters a source at some non-zero impact parameter ( $\Delta\theta_{\text{min}}$ ), a bright major image and faint minor image of the source are formed. The major image is located outside the Einstein radius and close to the source, whereas the minor image is located inside the Einstein radius and is close to the lens. In order, the major image, source, lens, and minor image always lie along the same line (see Boden, Shao & Van Buren 1998, fig. 2). The position of the images relative to the lens are given as

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(Paczynski 1986)

$$\frac{\theta_{\pm}}{\text{mas}} = \frac{1}{2} \left[ \pm (u^2 + 4)^{1/2} + u \right] \frac{\Theta_E}{\text{mas}}, \quad (2)$$

where we have defined the dimensionless instantaneous angular separation of the source and lens as  $u = \Delta\theta/\Theta_E$ . Here, and in the following equations the positive and negative parts refer to the major and minor images, respectively. At closest approach,  $u = u_{\min} = \Delta\theta_{\min}/\Theta_E$ . The amplifications of the two images are given as (Paczynski 1986)

$$A_{\pm} = \frac{u^2 + 2}{2u(u^2 + 4)^{1/2}} \pm 1. \quad (3)$$

This amplification causes an apparent brightening of the source (photometric microlensing) and an apparent displacement of the image–source light centroid (astrometric microlensing).

In the case of a luminous stellar lens and source, in which the lens, source, and images cannot be resolved, the apparent centroid shift due to both the major and minor images is suppressed. This is due to light contamination from the luminous lens. The centroid shift is reduced by a factor of  $(1 + f_l/f_s)$ , where  $f_l$  and  $f_s$  are the observed fluxes of the lens and source, respectively (Dominik & Sahu 2000). This effect often reduces the astrometric signal by a factor of  $\sim 100$ , making detection difficult (e.g. Proft, Demleitner & Wambsganss 2011). However, for some events, the impact parameter is large enough that the source and luminous lens can be resolved. In this case, we see an apparent shift of the source centroid, caused by the presence of the major image only. The centroid shift is found by taking the difference between the apparent position of the major image and the true position of the source and is given by Sahu et al. (2017) as

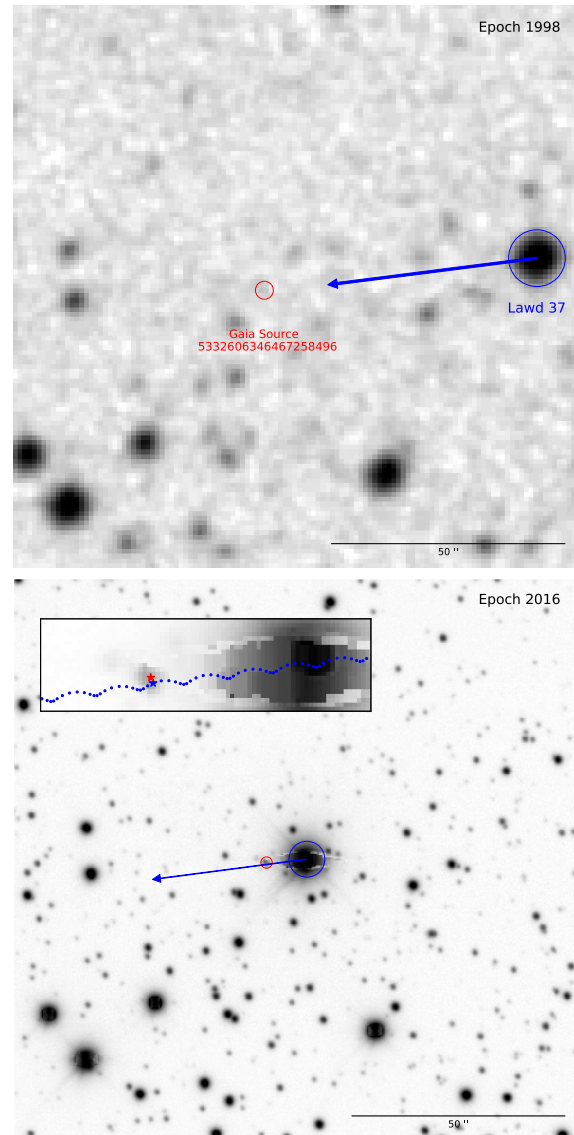
$$\frac{\delta\theta}{\text{mas}} = \theta_+ - \Delta\theta = \frac{1}{2} \left[ (u^2 + 4)^{1/2} - u \right] \frac{\Theta_E}{\text{mas}}. \quad (4)$$

Here, the centroid shift direction is always towards the position of the major image. This is maximal when the lens and source are at closest approach ( $u = u_{\min}$ ). If multi-epoch shifts in the source centroid and lens source separations can be measured for an event, the mass of the lens can be determined using equations (1) and (4), provided that the distance to the lens is known.

### 3 CANDIDATE EVENT PREDICTION

To search for events, a high proper motion ( $>150 \text{ mas yr}^{-1}$ ) sample of 13 206 lens stars from the TGAS catalogue was selected. To narrow our search, the lens sample was cross-matched with the *Gaia* DR1 source catalogue. Each lens was paired with all sources within a search radius of 10 times its proper motion. This produced a catalogue of  $\sim 4000$  lens–source pairs, which we investigated further by calculating time of closest approach and estimated astrometric deflection. The parallax motion of the lens and the proper motion of the source, where available from the ‘Hot Stuff for One Year’ proper motion catalogue (HSOY) (Altmann et al. 2017), was included.

We define a candidate lensing event as a lens–source pair which has a closest approach within the remaining *Gaia* mission time, assumed to be between 2018 and 2022. This left 30 candidate events. Visual inspection of the stellar field around each event removed six suspected erroneous events, which could not be confirmed to be genuine in the images available to us. Of the 24 remaining events, only one had an estimated significant maximum centroid shift in excess of 0.4 mas. It is this event that we report on here.



**Figure 1.** Images of the stellar field around the event. Top: Digitized Sky Survey image at epoch 1998. Bottom: Dark Energy Camera Plane Survey (DECAPS) (Schlafly et al. 2018) image at epoch 2016. On both images the blue circle indicates the position of the lens LAWD 37, the red circle indicates the position of the source, and the blue arrow indicates the proper motion vector of LAWD 37. Bottom inset: Zoom of the DECAPS image. This shows the source and lens position at closest approach marked with red and blue stars, respectively. The blue dashed line indicates the lens trajectory which includes parallax motion.

### 4 THE CANDIDATE

We predict that the known white dwarf LAWD 37 ( $G$  magnitude  $\sim 11$ ) will encounter a background source ( $G$  magnitude  $\sim 18$ ) with a closest approach of  $\Delta\theta_{\min} = 380 \pm 10 \text{ mas}$  ( $u_{\min} = 11.6 \pm 0.5$ ) on 2019 November  $11 \pm 4 \text{ d}$  ( $2019.86 \pm 0.01$  Julian Years). Fig. 1 shows the stellar field around the event and the trajectory of LAWD 37 as it approaches the source. The position and proper motion data for both LAWD 37 (the lens) and background source can be found in Table 1. Errors in the event parameters were calculated using the uncertainties in source and lens position, proper motion and parallax provided by the TGAS and HSOY catalogues.

**Table 1.** Lens LAWD 37 (first row) and source (second row) data. Proper motions of the lens and source are from the TGAS (Gaia Collaboration 2016a,b; Lindegren et al. 2016) and HSOY (Altmann et al. 2017) catalogues, respectively. The coordinates ( $\alpha$ ,  $\delta$ ) are from the *Gaia* DR1 source catalogue, on the ICRF and at epoch 2015.0 Julian Years. Distance to the lens  $D_1$  is obtained by inverting the lens parallax of  $215.8 \pm 0.2$  mas from TGAS.  $G$  is the *Gaia*  $G$ -band magnitude.

<i>Gaia</i> DR1 source id	$\alpha$ (deg $\pm$ mas)	$\delta$ (deg $\pm$ mas)	$\mu_\alpha \cos(\delta)$ (mas yr $^{-1}$ )	$\mu_\delta$ (mas yr $^{-1}$ )	$D_1$ (pc)	$G$ (mag)
5332606518269523072	176.454 9073 $\pm$ 0.2	-64.842 957 14 $\pm$ 0.2	2662.0 $\pm$ 0.2	-345.2 $\pm$ 0.2	4.63 $\pm$ 0.03	11.410 $\pm$ 0.002
5332606346467258496	176.463 604 56 $\pm$ 2	-64.843 297 79 $\pm$ 2	-14 $\pm$ 3	-2 $\pm$ 3	-	18.465 $\pm$ 0.005

At a distance of  $\sim 4.6$  pc, LAWD 37 (also known as WD 1142-645) is the fourth nearest known white dwarf to the Sun (Sion et al. 2009). It is classified as spectral type DQ indicating the presence of carbon in its atmosphere (Koester & Weidemann 1982). By fitting atmospheric models (Dufour, Bergeron & Fontaine 2005) to the photometry of LAWD 37, estimates for its effective temperature ( $T_{\text{eff}} = 7966 \pm 219$  K) and surface gravity ( $\log g = 8.09 \pm 0.02$ ) have been obtained (Giammichele, Bergeron & Dufour 2012). This surface gravity estimate combined with the parallactic distance allows the radius of LAWD 37 to be determined. Assuming LAWD 37 follows the standard evolutionary model for carbon-oxygen (CO) core white dwarfs, Giammichele et al. (2012) estimates that the radius corresponds to a mass of  $0.61 \pm 0.01 M_\odot$ .

Using equation (1), the mass estimate of Giammichele et al. (2012) and the TGAS parallax, we find the Einstein radius for LAWD 37 to be  $\Theta_E = 32.8 \pm 0.3$  mas. We have assumed that the source is sufficiently distant such that  $D_s \gg D_l$ . Fig. 2 shows the estimated astrometric signal and separation of the lens and source during the event. At closest approach, the maximum centroid shift is  $\delta\theta_{\text{max}} = 2.8 \pm 0.1$  mas. *Gaia*'s resolution limit is a function of the orientation of the objects with respect to the focal plane and the magnitude difference of the objects. However, it is potentially  $\sim 100$  mas,<sup>1</sup> as shown on Fig. 2. Due to the event's large impact parameter ( $u_{\text{min}} \gg 1$ ), the photometric signal is estimated to correspond to an apparent maximum brightening of the source of  $\sim 10^{-4}$  mag. Therefore, the photometric signal is unlikely to be detected by *Gaia*, so we consider constraining the mass of LAWD 37 from the astrometric signal only.

## 5 OBSERVATIONAL OUTLOOK

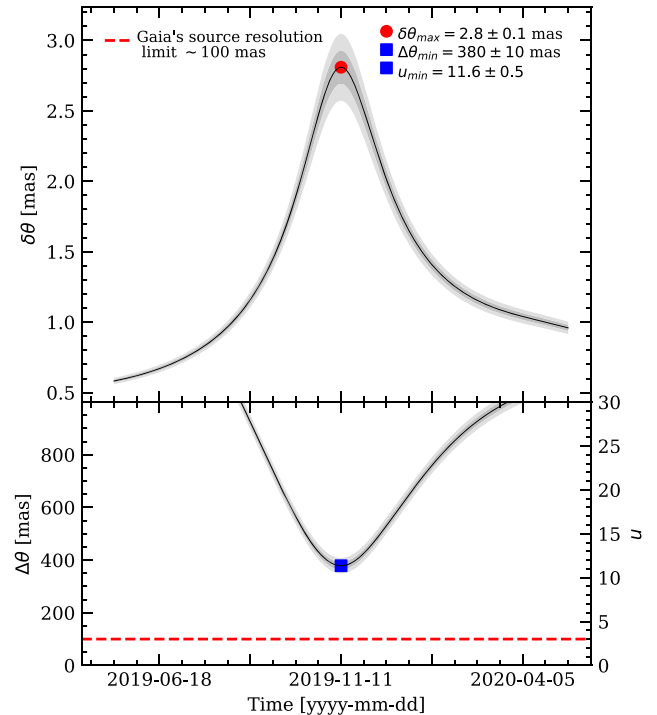
With a closest approach of  $\Delta\theta_{\text{min}} = 380 \pm 10$  mas, a predicted astrometric deflection of  $\delta\theta_{\text{max}} = 2.8 \pm 0.1$  mas and a lens-source magnitude difference of  $\sim 7$ , the viability of successfully observing this event has to be demonstrated.

### 5.1 *Gaia*

Fig. 2 shows that the lens and source should be resolvable by *Gaia* for the duration of the event. In order to assess the feasibility of observing the event with *Gaia*, we use the *Gaia* Observation Schedule Tool (GOST)<sup>2</sup> to predict the dates and scan direction of the expected *Gaia* observations. Due to the scanning law, observations are unevenly spaced and the scan direction, which is an important predictor of the possible centroiding precision, constantly changes. In Fig. 3, we show the propagated positions of the source and lens when *Gaia* is predicted to observe LAWD 37. We have also plotted the scan direction of *Gaia* and the estimated direction of deflection.

<sup>1</sup><https://www.cosmos.esa.int/web/gaia/science-performance>

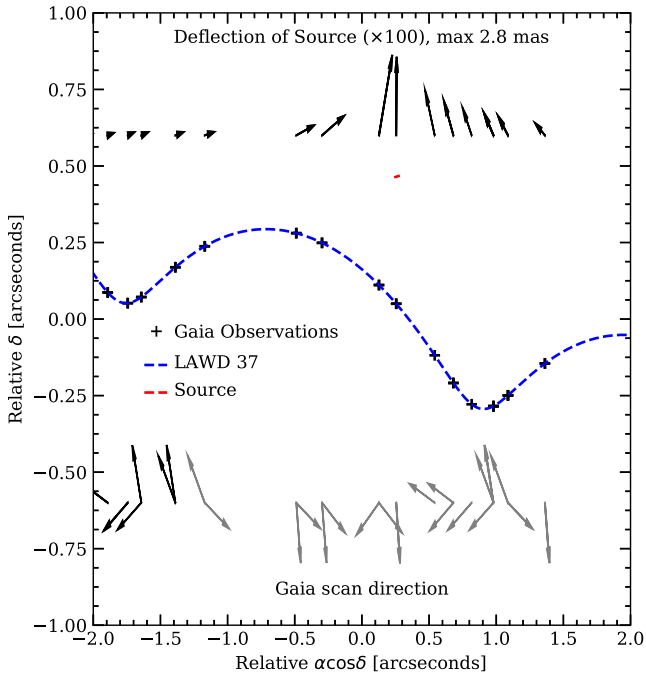
<sup>2</sup><https://gaia.esac.esa.int/gost/>



**Figure 2.** Top panel: Magnitude of the centroid shift around the time of closest approach (2019.86  $\pm$  0.01 Julian Years or 2019 November 11  $\pm$  4 d) for the event. This is calculated using equation (4). Maximum centroid shift (red circle) is  $\delta\theta_{\text{max}} = 2.8 \pm 0.1$  mas. Bottom panel: Lens-source separation around the time of closest approach for the event. The minimum separation (blue square) is  $u_{\text{min}} = 11.6 \pm 0.5$ , corresponding to a minimum angular separation of  $\Delta\theta_{\text{min}} = 380 \pm 10$  mas. Dark and lighter grey bands indicate  $1\sigma$  and  $2\sigma$  errors on both  $\delta\theta$  and  $\Delta\theta$ , respectively. The red dashed line indicates *Gaia*'s resolution limit of  $\sim 100$  mas (see the text).

The final centroiding precision of *Gaia* will be determined by a combination of the scan direction and the relative position of the two objects. Particularly for objects fainter than  $G = 13$ , *Gaia* provides only binned line-spread functions with very precise positions along scan directions, but relatively low precision in the across scan direction. For objects as bright as LAWD 37, *Gaia* will provide a window  $2 \times 1$  arcsec in the across and along scan, respectively (Fabricius et al. 2016). From one CCD transit, it is possible to obtain precisions of 0.06 mas for a  $G = 12$  object (Fabricius et al. 2016). However, because our primary measurement is the distance between the two objects, the floor will be set by the fainter source.

When the objects are observed in the same window, the use of gates to stop LAWD 37 saturating will lead to a significantly reduced signal to noise of the fainter source and a corresponding loss in precision. Even when in the same window, the higher precision along scan will remain because the pixels are rectangular and ap-



**Figure 3.** Blue and red dashed lines indicates LAWD 37’s and the source’s trajectory around the time of closest approach. Crosses mark the time of *Gaia*’s predicted observations. The vectors at the top indicate the predicted source deflection direction, the largest deflection is 2.8 mas. The arrows at the bottom indicate *Gaia*’s scan directions (grey arrows indicate provisional scan directions after 2019 June). When the deflection arrow and *Gaia*’s scan direction are aligned, the measurement is along scan and they will be the most precise.

proximately three times larger in the across compared to the along scan direction. For the best-case scenario, with both objects in the window and aligned along scan, the error on the apparent separation could be lower than 0.2 mas while in the worst-case scenario, with the orientation across scan, the error could be as high as 1 mas. This precision will be improved by a factor of 3 as we have nine independent estimates, one for each column in the focal plane. We simulated a uniform distribution of scan angles and assumed the along and across scan errors above and that the nine observations provide independent measurements. From this, we find the per epoch median error for the apparent lens source separation is  $\sigma_{ls} = 0.24$  mas. Current GOST results from around the event maximum indicate that there will be approximately 30 scans in which the astrometric deflection will be  $>2\sigma_{ls}$ .

Assuming that  $\Theta_E = 32.8$  for LAWD 37, we may estimate the precision at which *Gaia* could determine its mass. At each *Gaia* transit with an expected astrometric deflection  $>2\sigma_{ls}$ , we draw  $10^6$  samples from a Gaussian centred at the expected deflection and with variance  $\sigma_{ls}^2$ . We have assumed that the error on the true lens source separation is small compared to the error on the apparent lens source separation, so that  $\sigma_{ls} \approx \sigma_{\text{deflection}}$  since the apparent lens–source separation is the sum of the true lens–source separation and the deflection. Using these samples and inverting equation (4) for the mass of the lens, we calculate  $10^6$  simulated measurements for the mass of LAWD 37 at each transit. By taking the mean and variance of the mass measurement distributions for each transit and then calculating the inverse variance weighted average across all transits, we produce a final mass measurement and error. We

estimate in the best case that *Gaia* should be able to determine the mass of LAWD 37 to  $\sim 3$  per cent precision.

## 5.2 Hubble Space Telescope

Single measurement accuracies of  $\sim 0.2$  mas through pointed imaging by the Advanced Camera for Surveys and the Wide Field Camera 3 (WFC3) have been achieved with *HST* (see e.g. Bellini, Anderson & Bedin 2011). Although spatial scanning modes with WFC3 have enabled astrometric measurements with a precision 20–40  $\mu\text{as}$  Casertano et al. (2016), the magnitude difference in our event precludes this technique (Casertano, private communication). However, for our event the maximum centroid shift is estimated to be  $\sim 2.8$  mas, which is well within *HST*’s capabilities. Large-scale observing campaigns with *HST* to constrain masses of single objects via astrometric microlensing are already underway (Kains et al. 2017). Recently, the mass of white dwarf Stein 2015 B was determined with an accuracy  $\sim 7$  per cent via astrometric microlensing (Sahu et al. 2017). This event had a lens–source closest approach  $\sim 100$  mas. At the point that it was still resolvable by *HST* (separation  $\sim 500$  mas), this produced a deflection of the background source position of the order of  $\sim 2$  mas. The Stein 2015 B event is a similar brightness and contrast ratio to the LAWD 37 event and still the deflection was successfully measured by *HST*, providing an optimistic outlook for our event.

## 6 DISCUSSION AND CONCLUSIONS

White dwarfs are comprised mainly of degenerate matter. They are expected to obey a theoretical mass radius relationship (MRR) as they evolve and cool. Observational confirmation of the MRR is problematic, mainly due to the difficulty of determining the mass of white dwarfs. In a small number of cases when a white dwarf is found in an eclipsing or astrometric binary system (see e.g. Liebert et al. 2013; Parsons et al. 2016), or with a binary main-sequence companion in wide orbit whose radial velocity can be measured independently (Falcon et al. 2010), its mass can be calculated. However, for the majority of white dwarfs, the mass has to be determined indirectly using parameters ( $T_{\text{eff}}$ ,  $\log g$ ) derived from atmospheric models. These models are fitted using spectroscopy or broad-band photometry and require assumptions about the interior structure of white dwarfs. Specifically, the thickness of the non-degenerate hydrogen layers usually has to be prescribed, leading to poor constraints on MRRs. For example, Tremblay et al. (2017) mentions that MRRs derived from atmospheric models can vary between 1 and 15 per cent depending on whether a thin or thick hydrogen layer is assumed. Additionally, white dwarfs found in eclipsing binaries are post-common envelope, meaning they have interacted with their companion and potentially evolved differently from isolated white dwarfs.

LAWD 37 is a DQ white dwarf, so it has a helium-rich atmosphere. This often means thin hydrogen layers are prescribed in the atmospheric models (Tremblay et al. 2017). A mass determination of LAWD 37 by astrometric microlensing is completely independent of atmospheric models. Therefore, in addition to providing an independent check of model assumptions for white dwarfs with helium-rich atmospheres, it will provide an important comparison point between theoretical and observed MRRs, and white dwarf evolutionary theory.

In conclusion, we have predicted that the white dwarf LAWD 37 will lens the light from a background source, causing an apparent deflection of the source position. Maximally, this deflection

will be  $2.8 \pm 0.1$  mas on 2019 November  $11 \pm 4$  d. If LAWD 37 and the source are read out in the same window by *Gaia*, a mass determination to  $\sim 3$  per cent precision should be achieved. Recent observations with *HST* of a comparable astrometric microlensing event have allowed the successful determination of the mass of white dwarf Stein 2015 B with  $\sim 7$  per cent accuracy. This provides an optimistic outlook for a precision mass determination of LAWD 37 from our event with *HST*.

*Gaia*'s second data release (DR2) is set for 2018 April 25. In addition to providing a refined prediction of the event presented in this *letter*, DR2 will likely provide us with the ability to predict a large number of astrometric microlensing events, and hence precise mass measurements of a rich variety of stars.

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