

ACCEPTED MANUSCRIPT

Phase-contrast breast CT: the effect of propagation distance

To cite this article before publication: Luca Brombal *et al* 2018 *Phys. Med. Biol.* in press <https://doi.org/10.1088/1361-6560/aaf2e1>

Manuscript version: Accepted Manuscript

Accepted Manuscript is “the version of the article accepted for publication including all changes made as a result of the peer review process, and which may also include the addition to the article by IOP Publishing of a header, an article ID, a cover sheet and/or an ‘Accepted Manuscript’ watermark, but excluding any other editing, typesetting or other changes made by IOP Publishing and/or its licensors”

This Accepted Manuscript is © 2018 Institute of Physics and Engineering in Medicine.

During the embargo period (the 12 month period from the publication of the Version of Record of this article), the Accepted Manuscript is fully protected by copyright and cannot be reused or reposted elsewhere.

As the Version of Record of this article is going to be / has been published on a subscription basis, this Accepted Manuscript is available for reuse under a CC BY-NC-ND 3.0 licence after the 12 month embargo period.

After the embargo period, everyone is permitted to use copy and redistribute this article for non-commercial purposes only, provided that they adhere to all the terms of the licence <https://creativecommons.org/licenses/by-nc-nd/3.0>

Although reasonable endeavours have been taken to obtain all necessary permissions from third parties to include their copyrighted content within this article, their full citation and copyright line may not be present in this Accepted Manuscript version. Before using any content from this article, please refer to the Version of Record on IOPscience once published for full citation and copyright details, as permissions will likely be required. All third party content is fully copyright protected, unless specifically stated otherwise in the figure caption in the Version of Record.

View the [article online](#) for updates and enhancements.

Note

Phase-contrast breast CT: the effect of propagation distance

Luca Brombal^{1,2}, Sandro Donato^{1,2}, Diego Dreossi³, Fulvia Arfelli^{1,2}, Deborah Bonazza⁴, Adriano Contillo^{5,6}, Pasquale Delogu^{7,8}, Vittorio Di Trapani^{7,8}, Bruno Golosio^{9,10}, Giovanni Mettivier^{11,12}, Piernicola Oliva^{13,10}, Luigi Rigon^{1,2}, Angelo Taibi^{5,6} and Renata Longo^{1,2}

¹Department of Physics, University of Trieste, Via A. Valerio 2, 34127 Trieste, Italy

²INFN Division of Trieste, Via A. Valerio 2, 34127 Trieste, Italy

³Elettra-Sincrotrone Trieste S.C.p.A, 34149 Basovizza Trieste, Italy

⁴Department of Medical Science, Cattinara Hospital, University of Trieste, Strada di Fiume 447, 34149 Trieste, Italy

⁵Department of Physics and Earth Science, University of Ferrara, Via Savonarola 9, 44121 Ferrara, Italy

⁶INFN Division of Ferrara, via Saragat 1, 44122 Ferrara, Italy

⁷Department of Physical sciences, Earth and Environment, University of Siena, Strada Laterina 8, 53100 Siena, Italy

⁸INFN Division of Pisa, Largo Bruno Pontecorvo 3, 56127 Pisa, Italy

⁹Department of Physics, University of Cagliari, Monserrato-Sestu, I-09042 Monserrato, Italy

¹⁰INFN Division of Cagliari, Monserrato-Sestu, I-09042 Monserrato, Italy

¹¹Department of Physics, University of Napoli Federico II, Via Cinthia, 80126 Fuorigrotta, Napoli, Italy

¹²INFN Division of Napoli, Via Cinthia, 80126 Fuorigrotta, Napoli, Italy

¹³Department of Chemistry and Pharmacy, University of Sassari, Via Vienna, 07100 Sassari SS, Italy

E-mail: luca.brombal@ts.infn.it

July 2018

Abstract. X-ray phase imaging has the potential to dramatically improve soft tissue contrast sensitivity, which is a crucial requirement in many diagnostic applications such as breast imaging. In this context, a program devoted to perform *in-vivo* phase-contrast synchrotron radiation breast computed tomography is ongoing at the Elettra facility (Trieste, Italy). The used phase-contrast technique is the propagation-based configuration, which requires a spatially coherent source and a sufficient object-to-detector distance. In this work the effect of this distance on image quality is quantitatively investigated scanning a large breast surgical specimen at 3 object-to-detector distances (1.6, 3, 9 m) and comparing the images both before and after applying the phase-retrieval procedure. The sample is imaged at 30 keV with a 60 μm pixel pitch CdTe single-photon-counting detector, positioned at a fixed distance of 31.6 m from the source. The detector fluence is kept constant for all acquisitions. The study shows that, at the largest distance, a 20-fold SNR increase can be obtained by

1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41
42
43
44
45
46
47

applying the phase-retrieval procedure. Moreover, it is shown that, for phase-retrieved images, changing the object-to-detector distance does not affect spatial resolution while boosting SNR (4-fold increase going from the shortest to the largest distance). The experimental results are supported by a theoretical model proposed by other authors, whose salient results are presented in this paper.

48 Submitted to: *Phys. Med. Biol.*

49 1. Introduction

50 X-ray breast imaging is an extremely demanding task since high contrast sensitivity,
51 high spatial resolution and low delivered dose are required. In this context, x-
52 ray phase-contrast-imaging is a powerful tool to dramatically enhance soft tissues
53 contrast sensitivity without increasing dose. The advantage of phase-contrast imaging
54 over the conventional absorption imaging is based on the fact that, considering soft
55 tissues and energies in the range 10 – 100 keV, the decrement from unity (δ) of the
56 refraction index (n), responsible for phase effects, is about 3 orders of magnitude
57 higher than the absorption term (β), used in the conventional radiology (Rigon 2014).
58 Several approaches exist to transform the object-induced phase shift into intensity
59 modulations on the detector: interferometric (e.g., gratings), analyzer-based, edge-
60 illumination and free-space-propagation techniques are in use with synchrotron and,
61 in some cases, conventional sources (Bravin et al. 2012, Rigon 2014). From the
62 experimental point of view, the single-shot free-space propagation-based technique
63 is the easiest to implement since it only requires to increase the object-to-detector
64 distance without using optical elements or multi-exposure acquisition. On the contrary,
65 propagation-based imaging has more stringent requirements on the x-ray source spatial
66 coherence and detector spatial resolution with respect to other techniques (e.g., gratings,
67 edge illumination) (Pfeiffer et al. 2006, Olivo & Speller 2007). Images acquired with
68 the propagation-based technique show an enhanced contrast in the tissue interfaces
69 (i.e., edge-enhancement), which in the ray-optical approximation is proportional to the
70 Laplacian of the object-induced phase-shift (Peterzol et al. 2005). The edge-enhanced
71 images can be processed by applying a phase-retrieval (PhR) algorithm which allows,
72 under certain approximations, to recover the induced phase-shift (Burvall et al. 2011). In
73 this work the PhR algorithm based on the homogeneous transport of intensity equation
74 proposed by Paganin and co-workers in 2002 is used (Paganin et al. 2002). In fact,
75 the combined effect of free-space propagation and PhR is to increase the image signal-
76 to-noise ratio (SNR) preserving spatial resolution and, far from sharp interfaces where
77 edge-enhancement is present, contrast (Gureyev et al. 2017).

78 Along with phase effects, breast imaging can also take advantage of 3D techniques,
79 such as breast tomosynthesis and breast CT (BCT), which overcome the superposition of
80 the structures inherent in planar imaging potentially hindering the detection of massive

1
2
3
4
5 81 lesions. At present, the development of BCT systems is a hot topic for several research
6 82 groups and companies, the main challenge being the trade-off between spatial resolution
7 83 and delivered dose (Sechopoulos 2013, Sarno et al. 2015, Rößler et al. 2017, Kalender
8 84 et al. 2017).

9 85 In this context, the SYRMA-3D (synchrotron radiation for mammography)
10 86 collaboration aims to set-up the first clinical study of phase-contrast synchrotron
11 87 radiation BCT at the Elettra synchrotron facility (Trieste, Italy) and promising
12 88 results on breast specimens have been recently obtained (Longo et al. 2016, Brombal
13 89 et al. 2018a, Donato et al. 2018).

14 90 In this work, the effect of the propagation distance on the image quality, based on
15 91 scans of total mastectomy specimen acquired at 3 propagation distances, is discussed.
16 92 Specifically, both the effects of propagation distance and PhR on image metrics
17 93 as signal-to-noise-ratio, contrast and spatial resolution are reported and compared
18 94 with a theoretical model proposed by Gureyev, Nesterets and collaborators (Gureyev
19 95 et al. 2017, Nesterets & Gureyev 2014), which is briefly described in the next section.
20 96 A major improvement in signal-to-noise ratio at longer propagation distances and at a
21 97 constant spatial resolution is experimentally demonstrated.

28 29 98 **2. Materials and methods**

30 31 99 *2.1. Theoretical model*

32 100 Let us consider an object positioned at a distance R_1 from a monochromatic point x-
33 101 ray source and at a distance R_2 from a 1D detector (the extension to a 2D detector is
34 102 straightforward). We further suppose that the incident scalar electromagnetic wave
35 103 obeys to the homogeneous transport of intensity equation (TIE-hom), so that the
36 104 intensity at the detector plane ($I_{R_2}(x)$, with x the pixel coordinate) is:

$$37 105 \quad I_{R_2}(x) \simeq \left[1 - \sigma^2 \nabla_x^2\right] I_0(x) \quad , \quad (1)$$

38 106 $I_0(x)$ being the transmitted intensity in the object plane while $\sigma^2 = \gamma R' \lambda / (4\pi)$ accounts
39 107 for the (effective) propagation distance $R' = (R_1 R_2) / (R_1 + R_2)$, for the x-ray wavelength
40 108 λ and for the proportionality factor between the refraction and absorption properties
41 109 of an interface between 2 materials $\gamma = (\delta_2 - \delta_1) / (\beta_2 - \beta_1)$ (Gureyev et al. 2017). It is
42 110 worth noticing that, along with its validity conditions, equation 1 implies that the image
43 111 recorded at a given distance from the object will be similar to the (absorption) contact
44 112 plane image (i.e., at a null propagation distance) apart from the object's interfaces,
45 113 where the Laplacian of the intensity is expected to be significantly different from zero.
46 114 Therefore, within uniform regions of the images (i.e., far from sharp details), neither the
47 115 detected signal nor the noise are expected to change significantly upon the propagation
48 116 process. On the contrary, the spatial resolution improves in the free-space propagation.
49 117 This can be qualitatively understood considering that the (phase) contrast is increased
50 118 close to sharp interfaces (i.e., where the Laplacian is not negligible), hence the high
51 119 spatial frequencies are boosted. The quantitative demonstration of the spatial resolution

improvement associated to free-space propagation imaging can be found in (Gureyev et al. 2017). Once the propagation image has been collected, the phase-retrieved image is obtained by inverting the equation 1. In practice, this is accomplished by convolving the image with a low-pass filter which, in the spatial frequency domain (u), can be described as (Brombal et al. 2018a):

$$H(u) = [1 + 4\pi^2\sigma^2u^2]^{-1} \quad . \quad (2)$$

From the noise reduction perspective, the phase-retrieval filter has nothing special since any low-pass filter would reduce noise enhancing the SNR. Anyway, the peculiarity of the phase-retrieval procedure, applied along with the free-space propagation technique, is that it restores the resolution that would have been observed in the contact plane image while improving the SNR. This means that, once the phase retrieval has been applied, the spatial resolution of the image is the same at all propagation distances, except for magnification effects. In addition, considering flat portions of the image (i.e., far from sharp interfaces), phase retrieval does not modify the image contrast. This can be understood considering that, in practice, phase retrieval acts as a low pass filter, thus not altering the large area contrast of the image (Gureyev et al. 2017, Kitchen et al. 2017). Moreover, the fact that large area contrast is not affected by the phase-retrieval procedure can also be understood from a physical perspective. In fact, in the analytical derivation of the phase-retrieval formula it is assumed that absorption and phase properties are proportional throughout the object, thus, far from interfaces, large area contrast is not expected to vary (Paganin et al. 2002, Burvall et al. 2011).

If both phase-retrieval and tomographic reconstruction process are considered, it can be demonstrated that the SNR gain associated to the application of phase retrieval in the tomographic image is expected to be (Nesterets & Gureyev 2014):

$$SNR_{gain}(A) = \left[(8/3\pi) \frac{A^2}{\ln(A) - 1} \right]^{1/2} \quad , \quad (3)$$

being $A = \sigma^2/(16h'^2)$ a dimensionless parameter accounting for the object composition, irradiation geometry, beam energy (all described by σ) and the detector effective pixel size $h' = h/M$, where h is the physical pixel size and $M = (R_1 + R_2)/R_2$ the magnification factor. To obtain this result the detector is assumed to be ideal, i.e., with $MTF = 1$ up to the Nyquist frequency. The equation 3 is the central result of the model and, as a first approximation, it implies that the SNR gain increases almost linearly with the propagation distance and with the inverse of the square of the effective pixel size. Considering realistic parameters in terms of energy (tens of keV), propagation distance (meters) and pixel size (less than 100 μm), the expected SNR gain is between 1 and 2 orders of magnitude with respect to conventional imaging (Kitchen et al. 2017).

2.2. Experimental setup and sample

The images are acquired at the SYRMEP beamline at Elettra (Tromba et al. 2010). The x-ray beam is produced by one storage ring bending magnet and the energy is

1
2
3
4
5 selected in the range 8.5 - 40 keV by means of a Si(111) double-crystal monochromator,
6 providing an energy resolution of 0.1%. The beam's cross section at the detector
7 is 220 (horizontal) \times 4 mm² (vertical, Gaussian shape, FWHM) while the source-to-
8 detector distance is kept at 31.6 m for all measurements. Images are collected at
9 30 keV positioning the sample at 3 different object-to-detector distances, 1.6, 3 and
10 9 m, respectively. The laminar shape of the beam, along with long object-to-detector
11 distances, allows to work in a scatter-free geometry without the need of anti-scattering
12 grids (Brombal et al. 2018a). To be consistent with the notation of the model presented
13 in the previous section, the propagation distance (R') is defined as the object-to-detector
14 distance scaled by the magnification factor. Given that the magnifications at the 3
15 sample positions are 1.05, 1.1 and 1.4, the propagation distances will be 1.5, 2.7 and
16 6.4 m, respectively. It is worth noticing that, especially at high magnifications, the
17 actual finite dimension of the source should be taken into account since it contributes
18 to the overall image blurring, thus reducing the spatial resolution (Gureyev et al. 2008).
19 Anyway, in this work small magnifications (up to 1.4) are used and the detector spatial
20 resolution is similar to the source size ($\sim 100 \mu\text{m}$) (SYRMEP 2016), therefore making
21 the source size contribution to the image blurring (as a first approximation) negligible.

22 Each scan is performed in 40 seconds, collecting the projections over 180 deg with
23 a rotation speed of 4.5 deg sec⁻¹. The dose, expressed as mean glandular dose (MGD),
24 is evaluated by multiplying the air kerma at the patient position (i.e., 1.6 m object-to-
25 detector distance) by a conversion factor accounting for breast size and glandularity,
26 derived from an ad-hoc developed Monte Carlo simulation based on a GEANT4
27 code (Mettivier et al. 2015, Fedon et al. 2015). In this study, the delivered MGD at
28 the shortest propagation distance was 25 mGy. At larger distances, since the fluence on
29 the detector was kept roughly constant, the delivered dose was slightly increased ($\sim 5\%$
30 higher at 3 m and $\sim 30\%$ higher at 9 m) due to x-ray attenuation in air. In *in-vivo*
31 applications, this issue can easily be overcome by positioning a vacuum pipe between
32 the object and the detector, thus avoiding air attenuation. In addition, as it will be
33 clear in the next section, it can be argued that air attenuation is largely compensated
34 by the SNR increase at larger distances, leaving room for the possibility of a major dose
35 reduction.

36 The sample is a total breast mastectomy containing an epithelial and stromal
37 sarcomatoid carcinoma. After the formalin fixation and sealing in a vacuum bag,
38 the sample diameter is of about 12 cm. The Directive 2004/23/EC of the European
39 Parliament and of the Council of 31 March 2004 on setting standards of quality and
40 safety for the donation, procurement, testing, processing, preservation, storage and
41 distribution of human tissues were followed.

42 The images are collected with a CdTe single-photon-counting detector with a 60 μm
43 pixel pitch (Pixirad-8), comprising an array of 8 modules tiling a total surface of
44 246 \times 25 mm², operated in dead-time-free mode at a frame rate of 30 Hz (Bellazzini
45 et al. 2013, Delogu et al. 2017). Each scan is constituted by 1200 projections which first
46 undergo an ad-hoc pre-processing procedure (Brombal et al. 2018b) and subsequently are
47

1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41
42
43
44
45
46
47
48
49
50
51
52
53
54
55
56
57
58
59
60

200 phase-retrieved ($\gamma = 795$) and reconstructed via a GPU-based filtered back projection
201 with a Shepp-Logan filtering (Brun et al. 2015). The value of the retrieval parameter
202 γ has been extracted from a publicly available database (Taylor 2015) considering a
203 glandular/adipose interface.

204 2.3. Image analysis

205 As a first step the SNR of the images prior to the phase retrieval is measured selecting
206 circular ROIs (4000 pixels each) embedded within tumoral tissue, avoiding sharp edges.
207 According to the TIE model (equation 1) SNR should not change significantly varying
208 the propagation distance if no phase-retrieval is applied, being equal to the SNR that
209 would be observed in the contact plane. To compensate for the beam's magnification,
210 SNR is normalized to square root of the effective pixel size $h' = h/M$, where $h = 60 \mu\text{m}$
211 is the physical pixel pitch and M is the magnification. Moreover, to make up for small
212 fluence variations in different acquisitions, SNR is also normalized to the square root of
213 the average number of counts in the detector N , and defined to be:

$$214 \quad SNR = \frac{\langle I \rangle}{s(I)} \sqrt{\frac{h'_0}{h'}} \sqrt{\frac{N_0}{N}} \quad , \quad (4)$$

215 where $\langle I \rangle$ is the mean pixel value, $s(I)$ the standard deviation in the ROI, h'_0 and
216 N_0 are the reference pixel size and number of counts corresponding to the 1.5 m
217 propagation distance acquisition, respectively. The error associated to the SNR is given
218 by the standard deviation of 5 SNR measurements performed in non-overlapping ROIs.
219 SNR is measured also after the application of the phase-retrieval algorithm and a gain
220 factor is defined as:

$$221 \quad SNR_{gain} = \frac{SNR_{PhR}}{SNR_{noPhR}} \quad . \quad (5)$$

222 Subsequently, the image contrast is measured selecting ROIs both within tumor
223 (subscript 1) and adipose (subscript 2) regions:

$$224 \quad C = \frac{\langle I_1 \rangle - \langle I_2 \rangle}{\langle I_2 \rangle} \times 100 \quad . \quad (6)$$

225 Since phase retrieval is affecting only image noise while free space propagation is
226 affecting spatial resolution, the contrast should not change neither with the application
227 of the phase retrieval, nor varying the propagation distance. As for the SNR, the
228 error associated to the contrast is given by the standard deviation of 5 contrast values
229 measured in non-overlapping ROI pairs.

230 The spatial resolution is measured in the phase-retrieved images selecting, for each
231 distance, 3 line profiles across a sharp fat/tumor interface produced by a surgical cut.
232 The line profiles are fitted with an *erf* and the FWHM of its derivative is measured.
233 The spatial resolution is evaluated as the mean value of the 3 FWHMs and the error
234 is estimated to be the maximum fluctuation around the mean value. According to the
235 theory, excluding the effect of the magnification, the spatial resolution after the PhR
236 should not vary by changing the propagation distance since, for each distance, the PhR is

237 expected to produce the same resolution that would have been measured in the contact
238 plane image. In order to consider only the intrinsic system's spatial resolution, the
239 FWHM is measured in number of pixels instead of an absolute length.

240 3. Results and discussion

241 In figure 1 the reconstructed slices at different propagation distances (1.5, 2.7, 6.4 m)
242 without (a-c) and with (d-f) the phase retrieval are reported. With the aim of a
243 visualization allowing a straightforward comparison between images with and without
244 phase retrieval, all the images have been scaled by a normalization factor such that
245 the average value of fibroglandular tissues far from interfaces is 1 while air is 0. Since
246 tissue relaxation occurred and sample repositioning was needed, some morphological
247 changings (e.g., different position of air gaps within the tissue) are observed at different
248 propagation distances. Care was taken to ensure the best match at all distances in
249 the region enclosed by the dashed line of figure 1 (a), where all the measurements are
250 performed. From the images it can be qualitatively noted that, if no PhR is applied,
251 no major variation in signal and noise is observed by varying the propagation distance,
252 except for the sharp interfaces between adipose (dark grey) and tumor or fibroglandular
253 (bright gray) tissue. On the contrary, increasing the propagation distance, the phase-
254 retrieved reconstructions are smoother and no spatial resolution degradation is observed.

255 The same effect is reported in a finer detail in figure 2, where a zoom on a sharp
256 adipose/tumor interface produced by a surgical cut is displayed. Considering the non-
257 phase retrieved images (a-c) it is clear that the edge-enhancement effect at the interfaces
258 between the two different tissues is amplified at increasing propagation distances, i.e.,
259 the high-spatial frequencies are boosted. This can be better visualized in panels (g-
260 i) reporting line profiles (see dashed line in panels (a-c)) of the non-phase-retrieved
261 images at increasing propagation distances. Besides the edge-enhancement effect, clearly
262 visible in panel (i), the profiles show a high level of noise, possibly hampering tissue
263 differentiation. On the other hand, when the PhR is applied (d-f) the edge appearance
264 does not change by varying the propagation distance and the edge-contrast is not longer
265 present. Considering the respective line profiles reported in panels (j-l), a similar edge
266 sharpness is observed at all distances and, when compared with the non-phase-retrieved
267 images profiles, the noise level is significantly lower.

268 The quantitative results are reported in table 1. As predicted by the theory
269 the SNR, calculated according to equation 4, does not vary significantly with the
270 propagation distance prior to the PhR, while its increase associated with the phase
271 retrieval is greater than a factor of 20 when considering 6.4 m of propagation distance.
272 In addition, it must be noted that only little contrast variations (below 3%) are observed
273 when changing the distance while, at a given position, no significant contrast alterations
274 are associated to the PhR algorithm whose action is limited to image noise. Furthermore,
275 considering phase-retrieved images, the FWHM measured in pixel units does not vary
276 significantly with the propagation distances and, in all cases, was found to be slightly

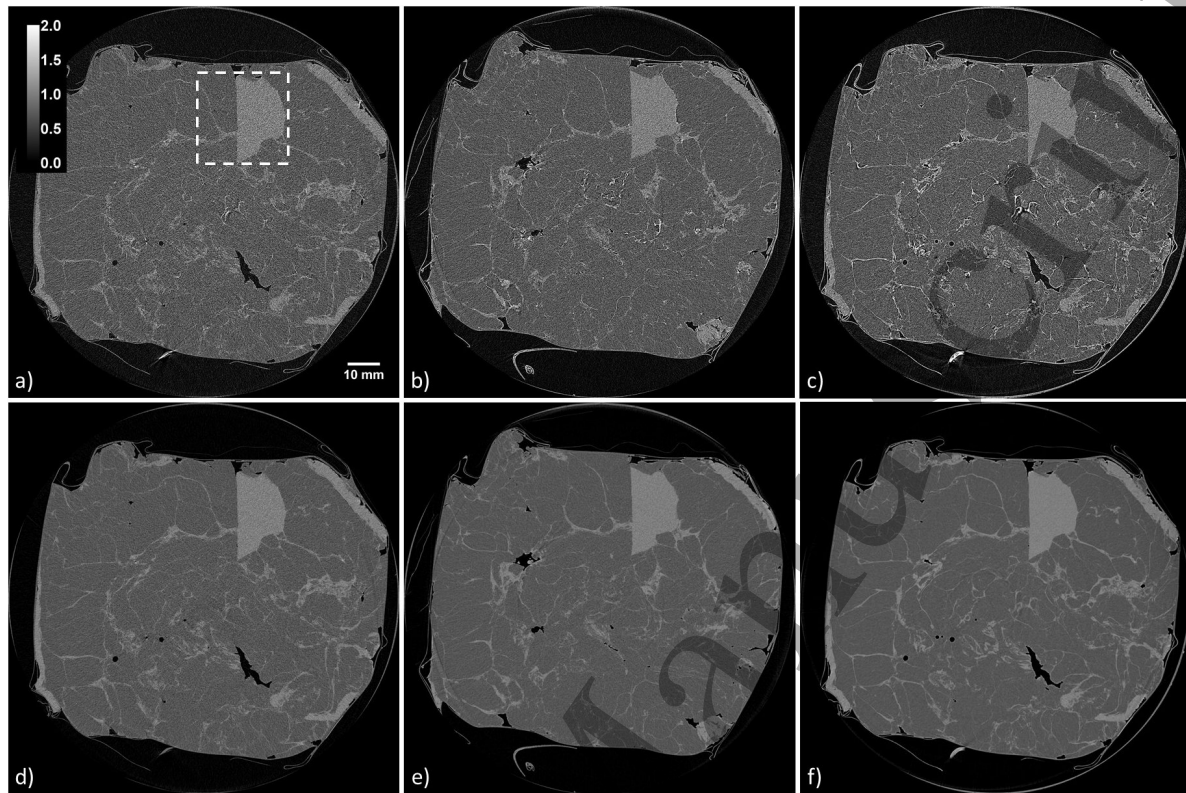


Figure 1. Reconstructed slice acquired at propagation distances 1.5 m (a, d), 2.7 m (b, e) and 6.4 m (c, f). Images in the first row (a-c) are reconstructed without PhR, images in the second row (d-f) with PhR. The dashed square in (a) is the zoom region reported in figure 2. After the normalization described in text, images are displayed in a gray scale window ranging from 0 to 2, where 0 is a typical value of air and 1 a typical value of fibroglandular tissue. Morphological variations at different distances are due to sample repositioning and tissue relaxation within the sample holder.

277 higher than 2 pixels ($\sim 120\mu\text{m}$ on the detector plane). This implies that, taking into
 278 account the magnification, the actual spatial resolution is improved at longer distances
 279 (FWHM $\sim 100\mu\text{m}$) at the expense of a smaller field of view.

Table 1. Quantitative results. The uncertainty associated to each measure is enclosed between round brackets.

	PhR	Distance R'		
		1.52 m	2.72 m	6.44 m
SNR	No	1.63 (0.02)	1.63 (0.03)	1.62 (0.01)
	Yes	8.45 (0.13)	13.3 (0.3)	33.8 (0.7)
Contrast	No	32.8 (0.4)	30.6 (0.3)	33.3 (0.2)
	Yes	32.7 (0.2)	30.7 (0.1)	32.9 (< 0.1)
FWHM (px)	Yes	2.1 (0.5)	2.3 (0.3)	2.4 (0.2)

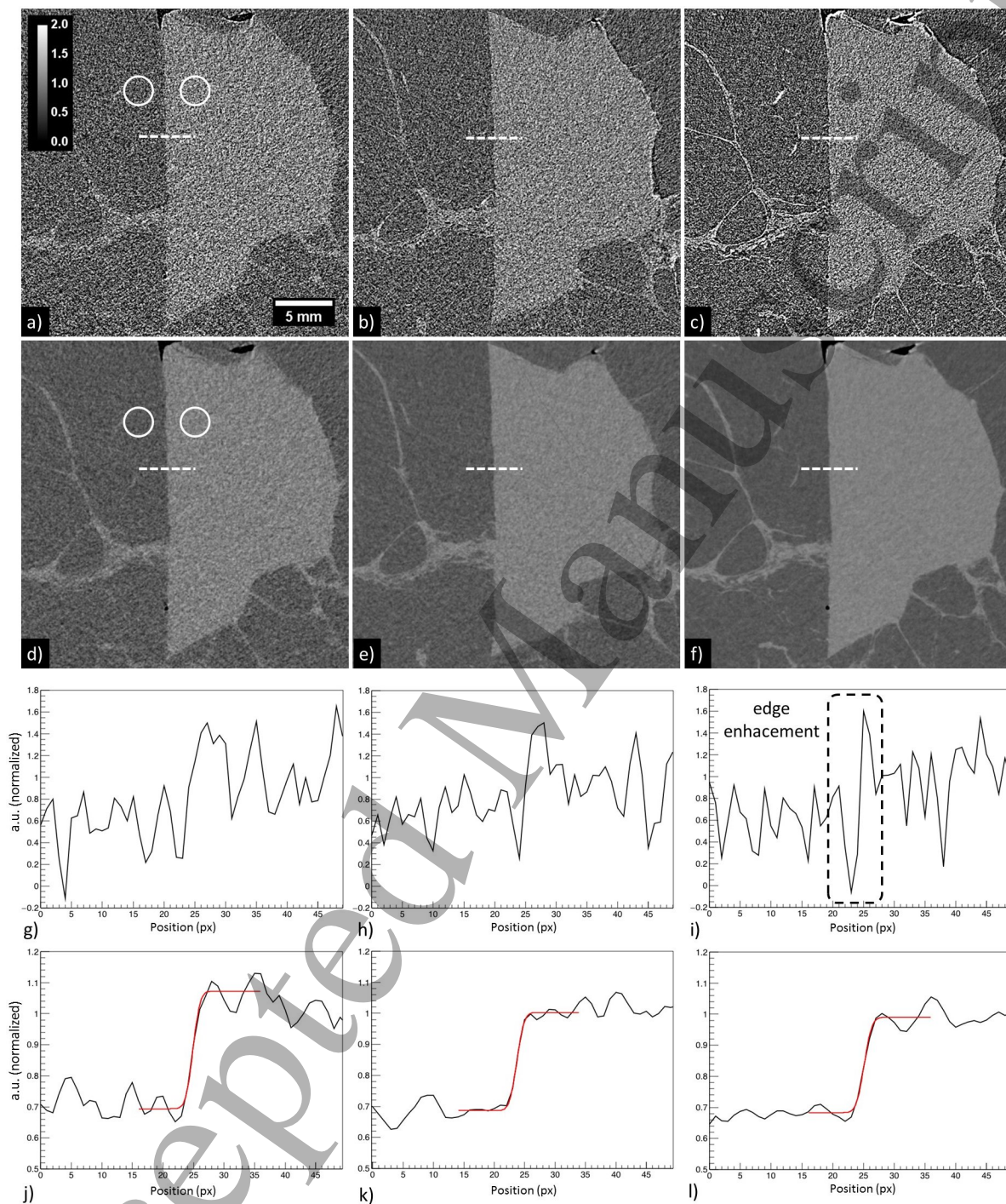


Figure 2. Zoomed detail of figure 1 without (a-c) and with (d-f) phase retrieval at increasing propagation distances (from left to right). In panels (g-i) profiles obtained from the dashed lines in (a-c) are reported. In panels (j-l) profiles obtained from the dashed lines in (d-f) are reported along with the *erf* fit (red curve). In (a) and (d) one of the five pairs of circular ROIs used to determine contrast and SNR are displayed as an example.

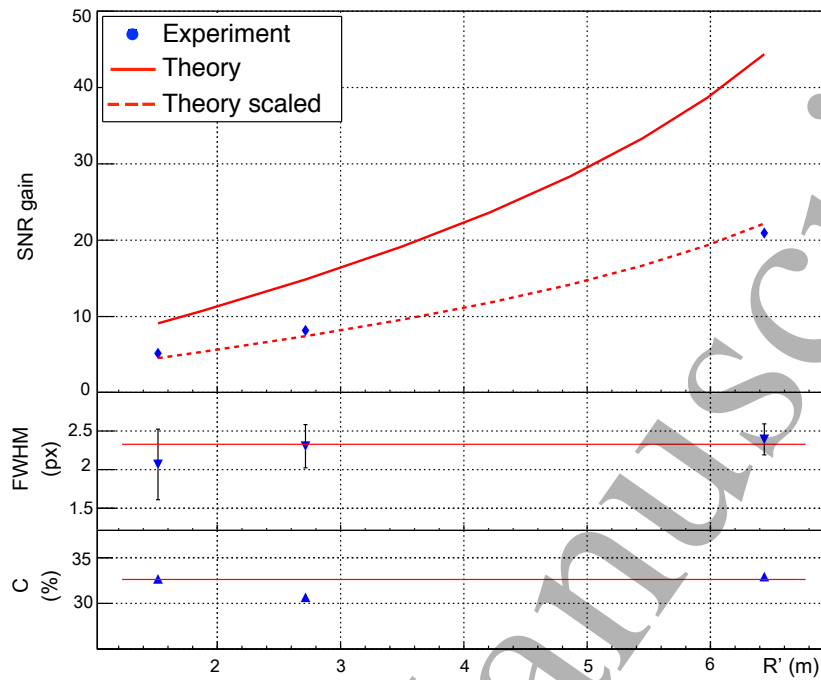


Figure 3. Comparison between experimental results (blue points) and theoretical predictions (solid red lines) as a function of the propagation distance. In the top panel the theoretical prediction scaled by a factor of 2 (dashed red line) is also reported. Some error bars are smaller than points.

With the aim of a better data visualization, the measured SNR gain, contrast and spatial resolution concerning the phase-retrieved images (blue points) and the theoretical predictions (solid red lines) are plotted as a function of the propagation distance in figure 3. From the top panel it can be seen that the measured SNR gain is lower than the predicted value at all propagation distances by roughly a factor of 2. This can be simply explained considering that the model assumes an ideal detector with a constant MTF up to the Nyquist frequency, thus constituting in practice an upper limit for the SNR gain when a real detector is considered. Once the theoretical curve is scaled (dashed red line), the experimental points match the theoretical trend. Moreover, comparing phase-retrieved images, a 4-fold increase in SNR can be obtained at 6.4 m with respect to the shortest propagation distance (1.5 m). At the same time, as predicted by the model, the spatial resolution is kept constant at all the distances (central panel) while only little contrast variations are observed (bottom panel).

4. conclusions

This study on a surgical breast specimen indicates that, combining the free-space propagation phase-contrast technique and the phase-retrieval algorithm, it is possible to obtain a major SNR improvement with respect to conventional imaging, at a constant spatial resolution. Specifically, at a fixed detector fluence, the longer propagation

distances provide higher SNR while leaving spatial resolution unaltered. The maximum observed SNR gain associated with the phase-retrieval algorithm is found to be 20 at 6.4 m while, at all propagation distances, the gain is about a factor of 2 smaller than the one predicted by the presented theoretical model which considers an ideal detector with a constant MTF. This means that the trend of the experimental points is consistent with the theory while the quantitative discrepancies should be attributed to the realistic (non-ideal) detector MTF. For the phase retrieved images, the spatial resolution measured across a sharp adipose/tumoral interface, is slightly higher than $100\mu\text{m}$ at all the propagation distances. In addition, it has been shown that, with the described experimental setup, major contrast variations are not observed neither changing propagation distance nor applying the phase-retrieval. This fact is of great importance in sight of the clinical application of this technique, since the image appearance will look "familiar" to the clinician's eye, who will not need a specific training to read the images. The presented work, where one sample was scanned at a limited number of propagation distances, will be expanded using different samples, propagation distances and including the detector's MTF in the theoretical model. Moreover, the SYRMEP beamline is being re-designed to accommodate larger patient-to-detector distances (1.6 m in the present configuration) to better exploit the advantages of the free-space propagation technique in breast CT clinical applications.

Acknowledgments

The authors gratefully thank Timur Gureyev and Yakov Nesterets for the fruitful discussion on the theoretical model and Giuliana Tromba for her support in the experimental activities. The authors gratefully acknowledge all the members of the SYRMA-3D project, Elettra-Sincrotrone Trieste SCpA and the Cattinara University Hospital for their support. Sandro Donato is partially supported by Consorzio per la Fisica Trieste.

References

- Bellazzini, R., Spandre, G., Brez, A., Minuti, M., Pinchera, M. & Mozzo, P. (2013). Chromatic x-ray imaging with a fine pitch cdtc sensor coupled to a large area photon counting pixel asic, *Journal of Instrumentation* **8**(02): C02028.
- Bravin, A., Coan, P. & Suortti, P. (2012). X-ray phase-contrast imaging: from pre-clinical applications towards clinics, *Physics in Medicine & Biology* **58**(1): R1.
- Brombal, L., Donato, S., Brun, F., Delogu, P., Fanti, V., Oliva, P., Rigon, L., Di Trapani, V., Longo, R. & Golosio, B. (2018b). Large-area single-photon-counting cdtc detector for synchrotron radiation computed tomography: a dedicated pre-processing procedure, *Journal of Synchrotron Radiation* **25**(4).
- Brombal, L., Golosio, B., Arfelli, F., Bonazza, D., Contillo, A., Delogu, P., Donato, S., Mettievier, G., Oliva, P., Rigon, L. et al. (2018a). Monochromatic breast ct: absorption and phase-retrieved images, *Medical Imaging 2018: Physics of Medical Imaging*, Vol. 10573, International Society for Optics and Photonics, p. 1057320.

- 338 Brun, F., Pacilè, S., Accardo, A., Kourousias, G., Dreossi, D., Mancini, L., Tromba, G. & Pugliese,
339 R. (2015). Enhanced and flexible software tools for x-ray computed tomography at the italian
340 synchrotron radiation facility elettra, *Fundamenta Informaticae* **141**(2-3): 233–243.
- 341 Burvall, A., Lundström, U., Takman, P. A., Larsson, D. H. & Hertz, H. M. (2011). Phase retrieval in
342 x-ray phase-contrast imaging suitable for tomography, *Optics express* **19**(11): 10359–10376.
- 343 Delogu, P., Golosio, B., Fedon, C., Arfelli, F., Bellazzini, R., Brez, A., Brun, F., Di Lillo, F., Dreossi,
344 D., Mettievier, G. et al. (2017). Imaging study of a phase-sensitive breast-ct system in continuous
345 acquisition mode, *Journal of Instrumentation* **12**(01): C01016.
- 346 Donato, S., Brombal, L., Tromba, G., Longo, R. et al. (2018). Phase-contrast breast-ct: Optimization
347 of experimental parameters and reconstruction algorithms, *World Congress on Medical Physics
348 and Biomedical Engineering 2018, IFMBE Proceedings*, Vol. 61, Springer, pp. 109–115.
- 349 Fedon, C., Longo, F., Mettievier, G. & Longo, R. (2015). Geant4 for breast dosimetry: parameters
350 optimization study, *Physics in Medicine & Biology* **60**(16): N311.
- 351 Gureyev, T. E., Nesterets, Y. I., Kozlov, A., Paganin, D. M. & Quiney, H. M. (2017). On the
352 unreasonable effectiveness of transport of intensity imaging and optical deconvolution, *JOSA A*
353 **34**(12): 2251–2260.
- 354 Gureyev, T. E., Nesterets, Y. I., Stevenson, A. W., Miller, P. R., Pogany, A. & Wilkins, S. W. (2008).
355 Some simple rules for contrast, signal-to-noise and resolution in in-line x-ray phase-contrast
356 imaging, *Optics express* **16**(5): 3223–3241.
- 357 Kalender, W. A., Kolditz, D., Steiding, C., Ruth, V., Lück, F., Rößler, A.-C. & Wenkel, E.
358 (2017). Technical feasibility proof for high-resolution low-dose photon-counting ct of the breast,
359 *European radiology* **27**(3): 1081–1086.
- 360 Kitchen, M. J., Buckley, G. A., Gureyev, T. E., Wallace, M. J., Andres-Thio, N., Uesugi, K., Yagi, N.
361 & Hooper, S. B. (2017). Ct dose reduction factors in the thousands using x-ray phase contrast,
362 *Scientific reports* **7**(1): 15953.
- 363 Longo, R., Arfelli, F., Bellazzini, R., Bottigli, U., Brez, A., Brun, F., Brunetti, A., Delogu, P., Di Lillo,
364 F., Dreossi, D. et al. (2016). Towards breast tomography with synchrotron radiation at elettra:
365 first images, *Physics in Medicine & Biology* **61**(4): 1634.
- 366 Mettievier, G., Fedon, C., Di Lillo, F., Longo, R., Sarno, A., Tromba, G. & Russo, P. (2015). Glandular
367 dose in breast computed tomography with synchrotron radiation, *Physics in Medicine & Biology*
368 **61**(2): 569.
- 369 Nesterets, Y. I. & Gureyev, T. E. (2014). Noise propagation in x-ray phase-contrast imaging and
370 computed tomography, *Journal of Physics D: Applied Physics* **47**(10): 105402.
- 371 Olivo, A. & Speller, R. (2007). A coded-aperture technique allowing x-ray phase contrast imaging with
372 conventional sources, *Applied Physics Letters* **91**(7): 074106.
- 373 Paganin, D., Mayo, S., Gureyev, T. E., Miller, P. R. & Wilkins, S. W. (2002). Simultaneous phase
374 and amplitude extraction from a single defocused image of a homogeneous object, *Journal of
375 microscopy* **206**(1): 33–40.
- 376 Peterzol, A., Olivo, A., Rigon, L., Pani, S. & Dreossi, D. (2005). The effects of the imaging system
377 on the validity limits of the ray-optical approach to phase contrast imaging, *Medical physics*
378 **32**(12): 3617–3627.
- 379 Pfeiffer, F., Weitkamp, T., Bunk, O. & David, C. (2006). Phase retrieval and differential phase-contrast
380 imaging with low-brilliance x-ray sources, *Nature physics* **2**(4): 258.
- 381 Rigon, L. (2014). x-ray imaging with coherent sources, *Comprehensive Biomedical Physics* **2**: 193–216.
- 382 Rößler, A.-C., Kalender, W., Kolditz, D., Steiding, C., Ruth, V., Preuss, C., Peter, S. C., Brehm,
383 B., Hammon, M., Schulz-Wendtland, R. et al. (2017). Performance of photon-counting breast
384 computed tomography, digital mammography, and digital breast tomosynthesis in evaluating
385 breast specimens, *Academic radiology* **24**(2): 184–190.
- 386 Sarno, A., Mettievier, G. & Russo, P. (2015). Dedicated breast computed tomography: basic aspects,
387 *Medical physics* **42**(6Part1): 2786–2804.
- 388 Sechopoulos, I. (2013). A review of breast tomosynthesis. part i. the image acquisition process, *Medical*

1
2
3
4
5 389 *physics* **40**(1).

6 390 SYRMEP (2016). SYRMEP specifications.

7 391 **URL:** www.elettra.trieste.it/lightsources/elettra/elettra-beamlines/symep/specification.html

8 392 Taylor, J. A. (2015). CSIRO TS imaging.

9 393 **URL:** www.ts-imaging.net/Services/Simple/ICUtilXdata.aspx

10 394 Tromba, G., Longo, R., Abrami, A., Arfelli, F., Astolfo, A., Bregant, P., Brun, F., Casarin, K., Chenda,
11 395 V., Dreossi, D. et al. (2010). The symep beamline of elettra: Clinical mammography and bio-
12 396 medical applications, *AIP Conference Proceedings*, Vol. 1266, AIP, pp. 18–23.

Accepted Manuscript