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# PREDICTION OF COLLAPSED LOAD OF STEEL COLUMNS USING FINITE STRIP METHOD

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Covenant University Ota, Lagos. Nigeria.\*Corresponding Author : [olatokunbo.ofuyatan@covenantuniversity.edu.ng](mailto:olatokunbo.ofuyatan@covenantuniversity.edu.ng)**ABSTRACT**

*The solution of inelastic bifurcation of column is the eigenvalue and eigenvectors which can be solved by many methods such as Power Method and William-Wittrick algorithm which has been proven to be efficient. This study utilised a new method of solution algorithm that utilizes the decomposition of the net assembled structural global matrix into the lower and upper triangles leading to the calculation of the determinant of the decomposed matrix. The product of the squared leading diagonals is the determinant which is the error of the present calculation. This method has been demonstrated to be about 0.02% error. The conventional Finite Strip Method is programmed with the new solution algorithm using an object oriented programming concept (Java). The method is compared with the published work and the agreement*

*between the current work and the reported experimental and numerical work is excellent. Residual stress is successfully implemented with the new solution algorithm.*

**Key words:** Finite strip method, column, numerical, collapse, residual stress.

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## 1. INTRODUCTION

Predicting the collapsed load of steel columns requires numerical computation because of material and geometric nonlinearities. The ability to reliably predict the collapsed load is fundamental to assessing the effects of local buckling, initial geometric imperfections in the section, residual stresses produced by manufacturing processes and material yielding and many other phenomena as noted by Ofuyatan *et al.* One of such methods is Finite Strip Method which is a versatile tool for the prediction of local, distortional and overall buckling of rolled and cold formed sections. The extent of the application of the method to the current work is limited to the material non-linearity. Although, the initial imperfection could be easily incorporated as demonstrated by Olawale, the emphasis now is the demonstration of the new solution algorithm. However, residual stress synonymous to rolled section is investigated. Li and Schafer have explored the solutions and provided design recommendations for two practical issues that develop when integrating computational member analysis with the conventional finite strip method (FSM) into cold-formed steel member design utilizing the direct strength method (DSM). First, FSM often fails to uniquely identify the relevant local and distortional member buckling modes. However, advancement in the field of FSM such as constrained finite strip method (cFMS) resolves some issues such as distortional buckling. cFMS also has some limitations such as inability to account for rounded corners in the model of cross-section. G. Eccher, K.J.R. Rasmussen, R. Zandonini presented the application of the isoparametric spline finite strip method to the elastic buckling analysis of perforated folded plate structures. The general theory of the isoparametric spline finite strip method was introduced. The kinematics assumptions, strain– displacement and constitutive relations of the Mindlin plate theory described and applied to the spline finite strip method. The corresponding matrix formulation was utilized in the equilibrium and stability equations to derive the stiffness and stability matrices. A number of numerical examples of flat and folded perforated plate structures illustrate the applicability and accuracy of the proposed.

## 2. FINITE STRIP METHOD

The general concept of Finite Strip Method has been well established as an efficient means of predicting buckling of plate and plated structures. In most cases, elastic approach has been very popular even in the age of powerful computational power. The extension of the elastic finite strip method to the inelastic range is accomplished by using the deformation theory of plasticity applied to thin plates. Based on this theory of plasticity, the nonlinear material properties are given by the following equations:

$$\{\sigma\} = E_{20}[F]\{\varepsilon\} \quad (1)$$

where the elasto-plastic modular matrix [F] is given by

$$[F] = \begin{bmatrix} f_{11} & f_{12} & 0 \\ f_{12} & f_{22} & 0 \\ 0 & 0 & f_{33} \end{bmatrix} \quad (2)$$

and the coefficients of [F] are given by

$$f_{11} = \frac{(1 + 3\tau)}{\mu}$$

$$f_{12} = \frac{2(1 - (1 - 2\nu)\tau)}{\mu}$$

$$f_{22} = \frac{4}{\mu}$$

$$f_{33} = \frac{1}{2 + 2\nu + 3e}$$

$$e = E_{20}/E_s$$

$$\tau = E_T/E_{20}$$

$$\mu = 5 - 4\nu + 3e - (1 - 2\nu)^2\tau \quad (3)$$

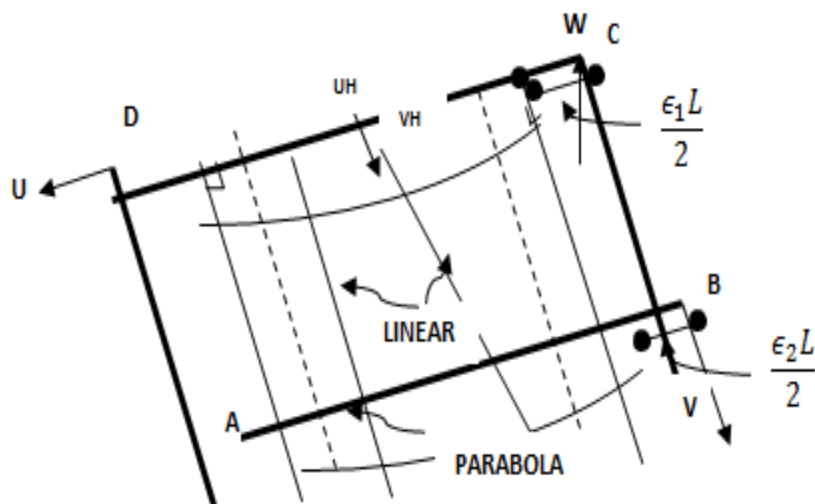
where

$E_t$  is the tangent modulus

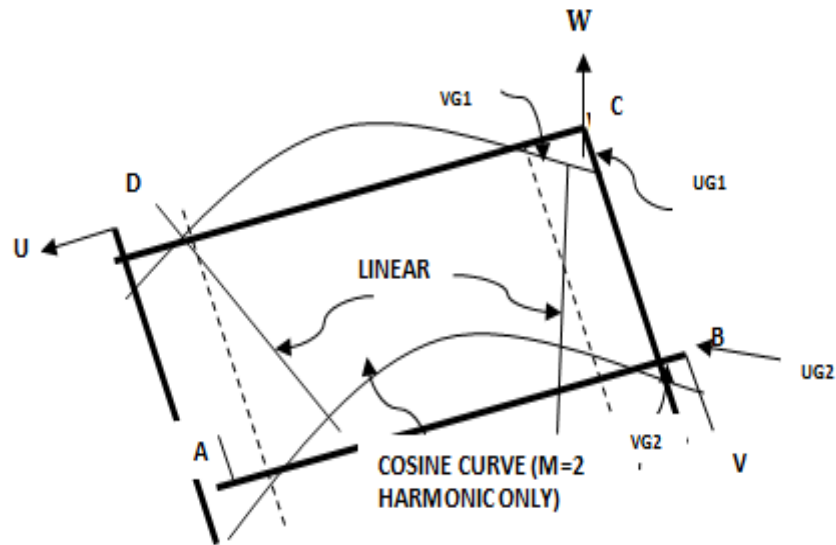
$E_s$  is the secant modulus

$E_{20}$  is Young's modulus of elasticity at ambient temperature.

These moduli are obtained from the uniaxial stress-strain relation.



**Figure 1a.** A basic plate showing the Hookean and membrane displacements of a strip



**Figure 1b** A basic plate showing the Hookean and membrane displacements of a strip  
The moment-curvature relationship is given by

$$\{M\} = \frac{E_{20}t^3}{12}[F]\{\chi\} \tag{4}$$

and the curvature is given by

$$\{\chi\} = \left\{ -\frac{\partial^2 W}{\partial x^2} - \frac{\partial^2 W}{\partial y^2} 2\frac{\partial^2 W}{\partial x \partial y} \right\}$$

and  $t$  is the plate thickness

To obtain the stiffness and stability matrices, the out-of-plane and in-plane effects are considered individually.

**a) Out-of-Plane:** The combination of the resulting equations enables all buckling modes to be considered. The out-of-plane displacement function is given by

$$f_o = \{W\} \tag{5}$$

$$f_o = \{Z\}^T \sin \pi \xi \{\delta_o\} \tag{6}$$

where

$$\{Z\}^T = \begin{Bmatrix} C_3 \\ C_4 \\ C_5 \\ C_6 \end{Bmatrix} \tag{7}$$

And  $C_3 = b(\eta - 2\eta^2 + \eta^3)$

$$C_4 = 1 - 3\eta^2 + \eta^3$$

$$C_5 = b(-\eta^2 + \eta^3)$$

$$C_6 = 3\eta^2 - 2\eta^3$$

$$\xi = \frac{x}{\lambda}$$

$$\eta = \frac{y}{b} \tag{8}$$

$$\{\delta_o\} = \begin{Bmatrix} \theta_1 \\ w_1 \\ \theta_2 \\ w_2 \end{Bmatrix} \tag{10}$$

The out-of-plane strain vector is related to the

curvature as given below

$$\{\varepsilon_o\} = z\{\chi\} \tag{11}$$

$$\{\varepsilon_o\} = [B_o]\{\delta_o\} \tag{12}$$

and  $[B_o]$  is given as

$$[B_o] = \begin{bmatrix} \frac{z\pi^2}{\lambda^2} \{Z\}^T \sin\pi\xi \\ -\frac{z}{b^2} \{Z'_{\eta\eta}\}^T \sin\pi\xi \\ \frac{2\pi z}{b\lambda} \{Z'_{\eta}\}^T \cos\pi\xi \end{bmatrix} \tag{13}$$

Applying the principle of virtual work, the change in the internal virtual work for the out-of-plane behaviour can easily be established leading to the out-of-plane stiffness matrix as documented in many literature.

$$[K_o] = \int_{vol} [B_o]^T [F] [B_o] dvol \tag{14}$$

Similarly, the out-of-plane stability matrix is obtained considering the out-of-plane displacement functions shown in eqn. 15 and membrane bending strain vector in eqn. 16 respectively.

$$\{f_o\} = [N_o]\{\delta_o\} \tag{15}$$

$$\{\varepsilon_{bo}\} = \frac{1}{2} \begin{Bmatrix} \left(\frac{\partial W}{\partial x}\right)^2 \\ \left(\frac{\partial W}{\partial y}\right)^2 \\ \frac{\partial W}{\partial x} \frac{\partial W}{\partial y} \end{Bmatrix} \tag{16}$$

$$\{\varepsilon_{bo}\} = \frac{1}{2} \begin{bmatrix} \frac{\partial W}{\partial x} & 0 \\ 0 & \frac{\partial W}{\partial y} \\ \frac{\partial W}{\partial y} & \frac{\partial W}{\partial x} \end{bmatrix} \begin{Bmatrix} \frac{\partial W}{\partial x} \\ \frac{\partial W}{\partial y} \end{Bmatrix} \tag{17}$$

$$\{\varepsilon_{bo}\} = \frac{1}{2} \{\delta_o\}^T \begin{bmatrix} \frac{1}{\lambda^2} [N'_{o,\xi}]^T [N'_{o,\xi}] \\ \frac{1}{b^2} [N'_{o,\eta}]^T [N'_{o,\eta}] \\ \frac{1}{b\lambda} \{ [N'_{o,\eta}]^T [N'_{o,\xi}] + [N'_{o,\xi}]^T [N'_{o,\eta}] \} \end{bmatrix} \{\delta_o\} \tag{18}$$

The out-of-plane stability matrix is obtained as shown in eqn. 19. Detail of its development is documented in Olawale <sup>[2]</sup>.

$$[S_o] = \int_{vol} [B_{so}] \{\sigma\}^T dvol \quad (19)$$

**b) In-Plane:**

The in-plane displacement function is given by

The in-plane displacement function is given by

$$\{f_i\} = \{U \ V\} = [N_i] \{\delta_i\} \quad (20)$$

$$f_i = \begin{bmatrix} \{X\}^T \cos \pi \xi \\ \{Y\}^T \sin \pi \xi \end{bmatrix} \{\delta_i\} \quad (21)$$

where

$$\begin{aligned} \{X\}^T &= \{0 \ C_1 \ 0 \ C_2\} \\ \{Y\}^T &= \{C_1 \ 0 \ C_2 \ 0\} \end{aligned} \quad (22)$$

and

$$\begin{aligned} C_1 &= 1 - \eta \\ C_2 &= \eta \end{aligned} \quad (23)$$

$$\{\delta_i\} = \{v_1 \ u_1 \ v_2 \ u_2\} \quad (24)$$

The in-plane strain vector is given by

$$\{\varepsilon_i\} = \begin{Bmatrix} \frac{\partial U}{\partial x} \\ \frac{\partial V}{\partial y} \\ \frac{\partial U}{\partial y} + \frac{\partial V}{\partial x} \end{Bmatrix} \quad (25)$$

$$\{\varepsilon_i\} = [B_i] \{\delta_i\} \quad (26)$$

and

$$[B_i] = \begin{bmatrix} -\frac{\pi}{\lambda} \{X\}^T \sin \pi \xi \\ \frac{1}{b} \{Y'_{\eta}\}^T \sin \pi \xi \\ (\frac{1}{b} \{X'_{\eta}\}^T + \frac{\pi}{\lambda} \{Y\}^T) \cos \pi \xi \end{bmatrix} \quad (27)$$

The in-plane stiffness matrix is obtained following the same principle of increment in virtual work and is given by eqn. 28

$$[K_i] = \int_{vol} [B_i]^T [F] [B_i] dvol \quad (28)$$

In a similar vein as the out-of-plane stability matrix, the in-plane stability matrix  $[S_i]$  is obtained using in-plane bending strain given in eqn. 29.

$$\{\varepsilon_{bi}\} = \frac{1}{2} \left\{ \begin{array}{l} \left(\frac{\partial U}{\partial x}\right)^2 + \left(\frac{\partial V}{\partial x}\right)^2 \\ \left(\frac{\partial U}{\partial y}\right)^2 + \left(\frac{\partial V}{\partial y}\right)^2 \\ 2\left(\frac{\partial U}{\partial x} \frac{\partial U}{\partial y} + \frac{\partial V}{\partial x} \frac{\partial V}{\partial y}\right) \end{array} \right\} \quad (29)$$

This can be expressed as

$$\{\varepsilon_{bi}\} = \frac{1}{2} \begin{bmatrix} \frac{\partial U}{\partial x} & \frac{\partial V}{\partial x} & 0 & 0 \\ 0 & 0 & \frac{\partial U}{\partial y} & \frac{\partial V}{\partial y} \\ \frac{\partial U}{\partial y} & \frac{\partial V}{\partial y} & \frac{\partial U}{\partial x} & \frac{\partial V}{\partial x} \end{bmatrix} \left\{ \begin{array}{l} \frac{\partial U}{\partial x} \\ \frac{\partial V}{\partial x} \\ \frac{\partial U}{\partial y} \\ \frac{\partial V}{\partial y} \end{array} \right\} \quad (30)$$

The in-plane stability matrix is thus obtained following the same principle as out-plane stability matrix as shown in eqn. 31

$$[S_i] = \frac{bt\pi^2}{\lambda} \int_0^1 \sigma_x [\{X\} \{X\}^T + \{Y\} \{Y\}^T] d\eta \quad (31)$$

The resulting out-of-plane and in-plane eigenvalue problems are solved using the new algorithm reported by Olawale and Ogunbiyi<sup>[1]</sup>.

$$[[K_o] - [S_o]]\{\sigma_o\} = 0 \quad (32)$$

Also, in-plane equilibrium is given as

$$[[K_i] - [S_i]]\{\sigma_i\} = 0 \quad (33)$$

The new solution algorithm used for the computation of the buckling load of column has been reported in Olawale and Ogunbiyi<sup>[1]</sup> and briefly described here for completeness.

The assembled global matrix is given as  $[A_G]$ . This can be decomposed into lower and upper triangular matrices such that  $[L] [L]^T = [A_G]$ , where

$$[L] = \begin{bmatrix} L_{11} & 0 & 0 & \dots & \dots & \dots & 0 \\ L_{21} & L_{22} & 0 & \dots & \dots & \dots & 0 \\ \vdots & \vdots & \vdots & & & & \vdots \\ \vdots & \vdots & \vdots & & & & \vdots \\ L_{n1} & L_{n1} & \dots & \dots & \dots & \dots & L_{nn} \end{bmatrix} \quad (34)$$

and  $[L]^T$  is the transpose of the  $[L]$  matrix. The multiplication of the lower and upper triangular matrices  $[L]$  and  $[L]^T$  yields the global matrix which can be assembled in the form below

$$[A_G] = \begin{bmatrix} a_{11} & \dots & \dots & \dots & \dots & a_{1n} \\ \vdots & & & & & \vdots \\ \vdots & & & & & \vdots \\ \vdots & & & & & \vdots \\ a_{n1} & \dots & \dots & \dots & \dots & a_{nn} \end{bmatrix}$$

where

$$L_{11} = (a_{11})^{1/2}$$

$$L_{1j} = a_{1j} / L_{11}$$

$$L_{ii} = \left( a_{ii} - \sum_{k=1}^{i-1} L_{ki}^2 \right)^{1/2}$$

$$L_{ij} = \left( a_{ij} - \sum_{k=1}^{i-1} L_{ki}L_{kj} \right) / L_{ii}$$

$$i = 2, 3, \dots, n$$

$$j = i + 1, i + 2, \dots, n$$

$$i > j$$

The determinant of the matrix  $[A_G]$  is given by

$$[A_G] = |[L][L]^T| = |[L]||[L]^T|$$

$$= L_{11}^2 L_{22}^2 L_{33}^2 \dots L_{nn}^2$$

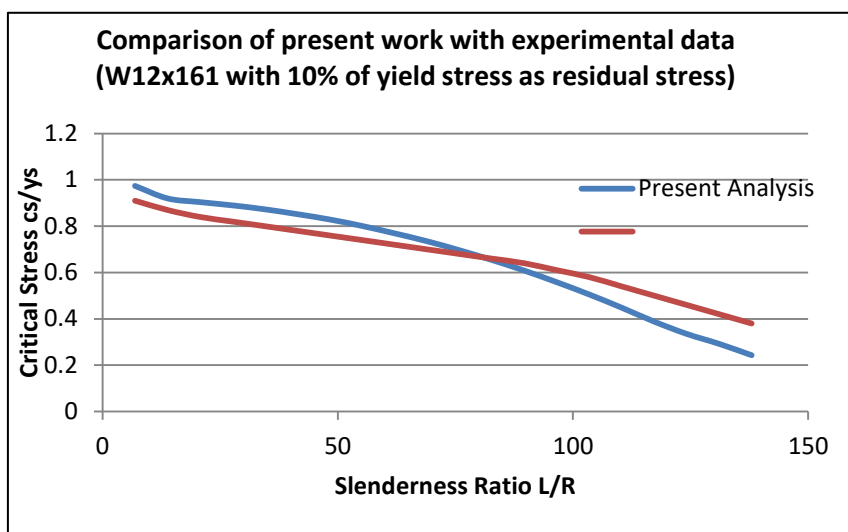
$$= \prod_{L=1}^n L_{11}^2 = \prod_{L=1}^n \left( a_{ii} - \sum_{k=1}^{i-1} L_{ki}^2 \right) i = 2, 3, \dots, n$$

$$\lambda = \frac{\sigma}{L^*}$$

$$\sigma_{building} = \lambda L^*$$

One of the advantages of the new solution method is that the need for the introduction of sub strips with a strip to achieve accuracy is totally eliminated. This, in effect reduces the bandwidth requirement and the computational effort is greatly reduced.

### 3. RESULTS AND DISCUSSION



**Figure 2** Comparison of the current analysis to European Convection Curve with Residual Stress of 10% Yield Stress.



The development is programmed in Java which is an object oriented programming language and the result of the current work is compared with European Convention curve B3-24 with 10% of yield stress for the residual stress level that is multi-linear distribution. The two computational results compare very closely and it can be said that current work is accurate to an appreciable extent for further parametric study. The plot of this comparison is shown in figure 2.

The behaviour of various rolled steel shapes is studied at various residual stress levels ranging from 10% to 50% yield stress. The residual stress distribution is a multi-linear relationship as shown in figure 3. The variation of the critical stress with the slenderness ratio for different residual stress levels indicates that the critical stress is adversely affected by the presence of residual stress that can be either locked in during rolling process or welding procedure. This effect is clearly shown in figure 4 & 5 and it can be seen that the bigger the shape the more adversely the critical stress is affected. The depreciation in the critical stress for small shape can be seen to be gradual throughout the range of slenderness ratios considered. Another interesting observation is that residual does not affect very stocky columns with slenderness ratio of about 20. It can be seen that only the intermediate and slender columns are adversely affected by the presence of residual stress.

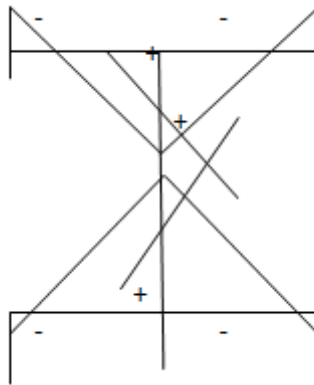


Figure 3 Residual stress distribution

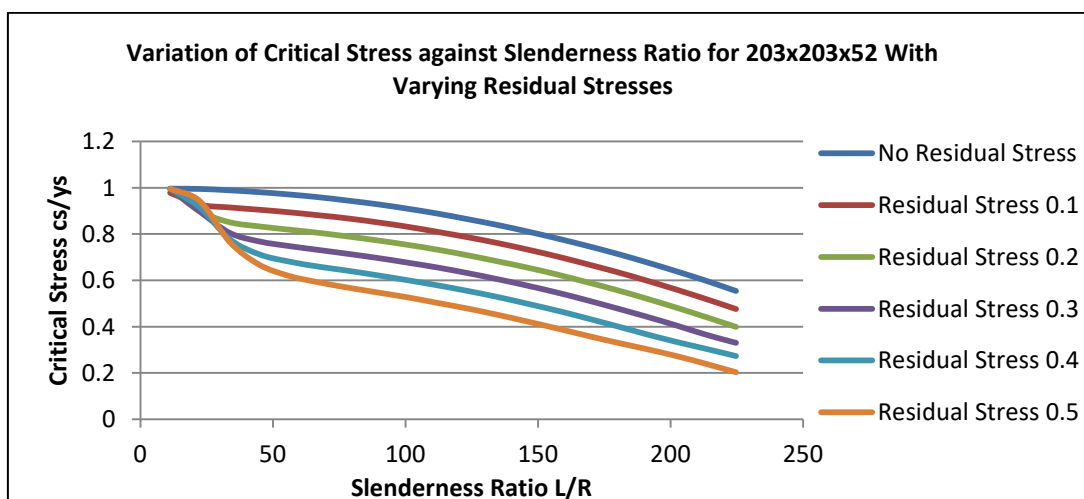
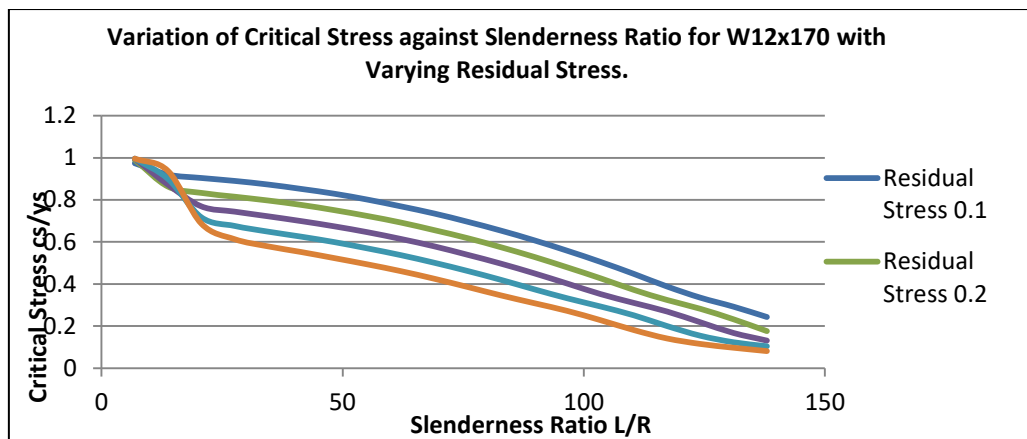
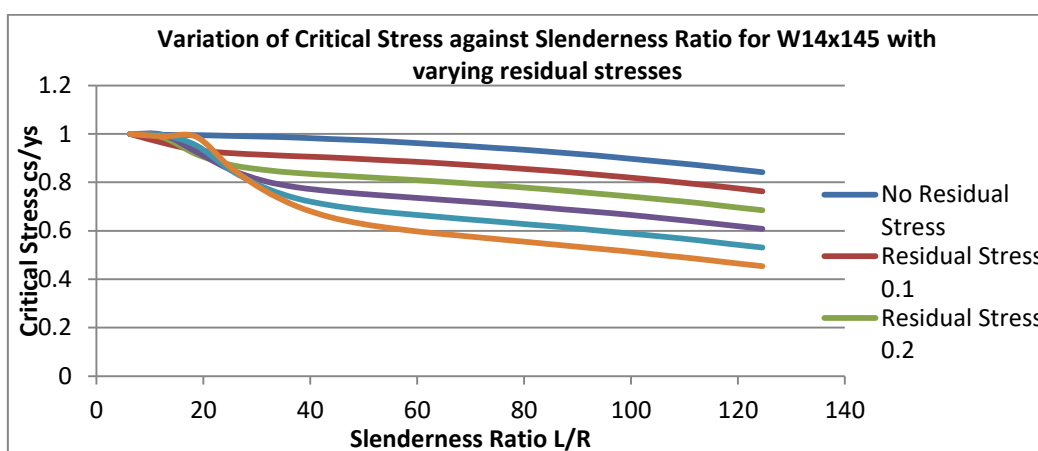


Figure 4 Critical Stress variation with slenderness ratio at various level of residual stress for normal rolled section.



**Figure 5** Critical Stress variation with slenderness ratio at various level of residual stress for wide flange rolled section



**Figure 6** Critical Stress variation with slenderness ratio at various level of residual stress for wide flange rolled section

#### 4. CONCLUSIONS

From this study, the power of new solution algorithm has been demonstrated in predicting the collapsed load of steel columns with or without residual stresses. It is also observed that residual stresses do not affect short columns but the intermediate and slender columns are susceptible to the presence of residual stress. The residual stress can have the capacity of the column if the magnitude of the residual stress is so high if not controlled at the time of production of the steel shapes. Also, the bigger the shape, the more the effect of residual stress on the capacity of columns is pronounced.

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