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A geometrical 1 % distance to the short-period binary Cepheid V1334 Cygni

A. GALLENNE,¹ P. KERVILLA,² N. R. EVANS,³ C.R. PROFFITT,⁴ J. D. MONNIER,⁵ A. MÉRAND,⁶ E. NELAN,⁴ E. WINSTON,³
G. PIETRZYŃSKI,⁷ G. SCHAEFER,⁸ W. GIEREN,^{9,10} R. I. ANDERSON,⁶ S. BORGNIET,¹¹ S. KRAUS,¹²
R. M. ROETTENBACHER,^{5,13} F. BARON,^{5,14} B. PILECKI,⁷ M. TAORMINA,⁷ D. GRACZYK,¹⁵ N. MOWLAVI,¹⁶ AND L. EYER¹⁶

¹European Southern Observatory, Alonso de Córdova 3107, Casilla 19001, Santiago, Chile

²LESIA, Observatoire de Paris, Université PSL, CNRS, Sorbonne Université,

Univ. Paris Diderot, Sorbonne Paris Cité, 5 place Jules Janssen, 92195 Meudon, France

³Smithsonian Astrophysical Observatory, MS 4, 60 Garden Street, Cambridge, MA 02138, USA

⁴Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

⁵Astronomy Department, University of Michigan, 941 Dennison Bldg, Ann Arbor, MI 48109-1090, USA

⁶European Southern Observatory, Karl-Schwarzschild-Str. 2, 85748 Garching, Germany

⁷Centrum Astronomiczne im. Mikołaja Kopernika, PAN, Bartycka 18, 00-716 Warsaw, Poland

⁸The CHARA Array of Georgia State University, Mount Wilson CA 91023, USA

⁹Universidad de Concepción, Departamento de Astronomía, Casilla 160-C, Concepción, Chile

¹⁰Millennium Institute of Astrophysics (MAS), Chile

¹¹LESIA (UMR 8109), Observatoire de Paris, PSL, CNRS, UPMC, Univ. Paris-Diderot

5 place Jules Janssen, 92195 Meudon, France

¹²University of Exeter, School of Physics and Astronomy, Stocker Road, Exeter, EX4 4QL

¹³Department of Astronomy, Stockholm University, SE-106 91 Stockholm, Sweden

¹⁴Department of Physics and Astronomy, Georgia State University, Atlanta, GA 30303, USA

¹⁵Centrum Astronomiczne im. Mikołaja Kopernika, PAN, Rabiańska 8, 87100 Toruń, Poland

¹⁶Département d'Astronomie, Université de Genève, 51 Ch. des Maillettes, 1290 Sauverny, Switzerland

ABSTRACT

Cepheid stars play a considerable role as extragalactic distances indicators, thanks to the simple empirical relation between their pulsation period and their luminosity. They overlap with that of secondary distance indicators, such as Type Ia supernovae, whose distance scale is tied to Cepheid luminosities. However, the Period–Luminosity (P–L) relation still lacks a calibration to better than 5%. Using an original combination of interferometric astrometry with optical and ultraviolet spectroscopy, we measured the geometrical distance $d = 720.35 \pm 7.84$ pc of the 3.33 d period Cepheid V1334 Cyg with an unprecedented accuracy of $\pm 1\%$, providing the most accurate distance for a Cepheid. Placing this star in the P–L diagram provides an independent test of existing period–luminosity relations. We show that the secondary star has a significant impact on the integrated magnitude, particularly at visible wavelengths. Binarity in future high precision calibrations of the P–L relations is not negligible, at least in the short-period regime. Subtracting the companion flux leaves V1334 Cyg in marginal agreement with existing photometric-based P–L relations, indicating either an overall calibration bias or a significant intrinsic dispersion at a few percent level. Our work also enabled us to determine the dynamical masses of both components, $M_1 = 4.288 \pm 0.133 M_\odot$ (Cepheid) and $M_2 = 4.040 \pm 0.048 M_\odot$ (companion), providing the most accurate masses for a Galactic binary Cepheid system.

Keywords: instrumentation: high angular resolution – techniques: interferometric, radial velocities – astrometry – binaries: close, spectroscopic – stars: variables: Cepheids

1. INTRODUCTION

Cepheids play a particularly important role as extragalactic distance indicators thanks to the empirical relation between their pulsation period and their luminosity: the Period–Luminosity (P–L) relation or Leavitt law (Leavitt & Pickering 1912). The ability to predict their

Corresponding author: Alexandre Gallenne
agallenn@eso.org

intrinsic luminosity from their variation period makes them precious standard candles, as its combination with their measured apparent magnitude gives the distance modulus in a simple, empirical manner. Cepheids are radially pulsating, young and bright supergiant stars that are easily identified in distant galaxies ($\lesssim 25$ Mpc). Their distance scale overlaps with secondary distance indicators, such as supernovae Ia (SNIa), whose distance scale is securely tied to Cepheid luminosities. Therefore, an accurate and precise calibration of the Cepheid luminosities is necessary to ensure a reliable calibration for the secondary distance indicators. However, despite considerable efforts over more than one century, the calibration of the zero point (ZP) of the P–L relation has not yet reached a 1% accuracy level. A better control of the systematic uncertainties is needed to establish the P–L relation calibration at this level, and secondary effects on the photometric measurements such as a metallicity-dependence or interstellar reddening must be understood and taken into consideration.

It is now clear that the intrinsic scatter of the relation is smaller at longer wavelengths (see e.g. [Freedman & Madore 2010](#); [Madore & Freedman 2012](#)) than it is for V band ($\sim 0.5 \mu\text{m}$) measurements. In optical bands, the line-of-sight extinction is stronger than in the infrared, and flux contamination from any main-sequence companions is likely to be larger – from an evolutionary time-scale point of view, most of the companions should be stars close to the main-sequence (see e.g. [Bohm-Vitense 1985](#); [Evans 1992](#); [Szabados & Klagyivik 2012](#)). As they are relatively massive stars, the binary fraction of Cepheids is known to be at least 50% ([Szabados 2003](#)), and the presence of companions may bias the calibration of the P–L relation if their photometric contribution is not removed ([Gaia Collaboration et al. 2017](#); [Szabados & Klagyivik 2012](#)). Furthermore, this bias is likely not constant as a function of the pulsation period as the bright long-period Cepheids (hence massive) have a higher contrast with their companions (reducing the potential photometric bias) compared to the shortest periods (less massive), but they are also more likely to be members of multiple stars systems with more massive companions than are short period Cepheids (see e.g. [Sana & Evans 2011](#), increasing the bias). In the near- and mid-IR, Cepheids usually strongly overshadow their companions, and such effects are reduced. On the other hand, the presence of circumstellar envelopes increases the measured flux beyond that of the Cepheid itself ($\lesssim 5\%$ and $1 - 30\%$ of IR excess in the near- and mid-IR, respectively), and so systematically affects the distance estimates ([Gallenne et al. 2017, 2013a](#); [Kervella et al. 2006](#); [Mérand et al. 2006](#)). The optimal band

seems to be the K band ($\sim 2.2 \mu\text{m}$) where the contributions of both the companions and the envelopes are mitigated, but the Cepheid distance scale must nevertheless be tested at the percent level. Such test seems to be even more critical with the growing evidence of a $\sim 3.7\sigma$ discrepancy between the value of the Hubble constant H_0 predicted from Cosmic Microwave Background anisotropies ([Planck Collaboration et al. 2018](#)) and the recent empirical calibration using Cepheids and SNIa ([Riess et al. 2016](#)). Before invoking new physics in the standard cosmological model, new high accuracy data are necessary.

For now, the calibration of the P–L relation is partially anchored to 17 Galactic Cepheids having trigonometric parallax measurements, obtained from the Hubble Space Telescope (HST) Fine Guidance Sensor ([Benedict et al. 2007, 2002](#), 10 Cepheids) and Wide Field Camera 3 spatial scanning ([Riess et al. 2018](#), 7 Cepheids). Unfortunately, their average accuracy is limited to 8%, preventing a reliable 1% calibration of their luminosity. While the final Gaia data release expected in 2022 will provide hundreds of Cepheid parallaxes to 3% or less, our knowledge of the interstellar extinction will still be an important limitation to derive highly accurate Cepheid luminosities. In addition, close orbiting companions could bias the Gaia astrometric measurements and therefore the resulting parallaxes if they are assumed single stars (see e.g. [Anderson et al. 2016b](#)). Here we report the determination of the geometrical distance of the Galactic binary Cepheid V1334 Cyg with a very high accuracy of 1%, providing the most accurate anchor point of the P–L relation for short period Cepheids. Our observations also enable us to confirm and measure the bias due to the companion on P–L calibrations.

Binary Cepheids offer the opportunity to measure accurate distances and masses through the combination of astrometric (visual orbit) and spectroscopic (radial velocities) measurements. But this is a challenging task because of 1) the small angular size of the orbits that hampers astrometric measurements using single-dish telescopes, and 2) the high contrast between the Cepheid and its companion, that complicates the detection of the companion’s spectral lines and the measurement of its radial velocity. Despite these difficulties, we report in the present work the detection of both the Cepheid and secondary component of V1334 Cyg using interferometric, as well as visible and ultraviolet spectroscopic observations. This system was the first target chosen for our interferometric observing campaign, first because the expected H -band contrast between the components is reachable by the current interferometric combiners, and also because the spectroscopic orbit is

particularly well known (Evans 1992, 2000). We present in Sect. 2 all the data we collected for the V1334 Cyg system. In Sect. 3 we introduce the models we used to fit our observations and derive our orbital solutions. We then discuss and conclude our results in Sect. 4.

2. OBSERVATIONS AND DATA REDUCTION

2.1. Interferometry

The relative astrometric positions of the companion of V1334 Cyg were measured repeatedly using interferometric observations performed with the Michigan InfraRed Combiner (Monnier et al. 2004), installed at the Center for High Angular Resolution Astronomy array (ten Brummelaar et al. 2005) located on Mount Wilson, California. The array consists of six 1 m aperture telescopes with an Y-shaped configuration (two telescopes on each branch), oriented to the east, west and south, and so offering a good coverage of the (u, v) plane. The baselines range from 34 m to 331 m, providing an angular resolution down to 0.5 mas in H . The MIRC instrument combines the light coming from all six telescopes in the H band, with three spectral resolutions ($R = 42, 150$ and 400). The recombination of six telescopes gives simultaneously 15 fringe visibilities and 20 closure phase measurements, that are our primary observables. Our observations were carried out from 2012 July to 2016 October using either four, five or six telescopes. We used the low spectral resolution mode, where a prism splits the light on the detector into 8 narrow spectral channels. We followed a standard observing procedure, i.e. we monitored the interferometric transfer function by observing a calibrator before and after our Cepheids. The calibrators were selected using the SearchCal software¹ (Bonneau et al. 2006) provided by the Jean-Marie Mariotti Center (JMMC). They are listed in Table 1.

Table 1. H -band uniform disk sizes of calibrators.

Star	θ_{UD} (mas)	V (mag)	H (mag)
HD192985	0.450 ± 0.032	5.87	4.87
HD199956	0.603 ± 0.043	6.65	4.46
HD202850	0.566 ± 0.049	4.22	3.86
HD204153	0.440 ± 0.031	5.59	4.86
HD207978	0.571 ± 0.040	5.52	4.44
HD212487	0.434 ± 0.031	6.18	4.97
HD214200	0.704 ± 0.050	6.11	4.16
HD218470	0.477 ± 0.034	6.68	4.67

¹ Available at <http://www.jmmc.fr/searchcal>.

The data were reduced with the standard MIRC pipeline (Monnier et al. 2007). The main procedure is to compute squared visibilities and triple products for each baseline and spectral channel, and to correct for photon and readout noises.

For each epoch, astrometric position were determined by fitting squared visibilities (V^2) and closure phases (CP) with our dedicated tool CANDID (Gallenne et al. 2015), with errors estimated using the bootstrapping technique (with replacement and 1000 bootstrap sample). At each epoch, the fitted parameters are the relative astrometric position $(\Delta\alpha, \Delta\delta)$, the uniform disk (UD) angular diameter of the primary θ_{UD} , and the flux ratio between the components f . The equation for the complex visibility for a binary system with an unresolved companion is:

$$\begin{aligned}
 V(u, v) &= \frac{V_{\star}(u, v) + fV_c(u, v)}{1 + f} \\
 V_{\star}(u, v) &= \frac{2J_1(\pi\theta_{\text{UD}}B/\lambda)}{\pi\theta_{\text{UD}}B/\lambda} \\
 V_c(u, v) &= \exp\left[-\frac{2i\pi}{\lambda}(u\Delta\alpha + v\Delta\delta)\right] \\
 \phi_{ijk} &= \arg[V(u_i, v_i)V(u_j, v_j)V^*(u_k, v_k)],
 \end{aligned}$$

where $V_{\star}(u, v)$ and $V_c(u, v)$ are the complex visibilities of the Cepheid and companion, respectively, J_1 the first-order Bessel function, (u, v) the spatial frequencies, B the baseline between two telescopes (i, j) , λ the observing wavelength, and ϕ_{ijk} the closure phase for each closed baseline triangle (i, j, k) . More details about the companion search formalism with CANDID are available in Gallenne et al. (2015, 2013b). Briefly, the tool provides a 2D-grid of fit using a least-squares algorithm. For each starting point in the grid, a multi-parameter fit is performed, i.e. adjusting $\Delta\alpha$, $\Delta\delta$, θ_{UD} , and f . Each position of the grid leads to a local minimum. If the starting grid is fine enough (estimated a posteriori), multiple starting points lead to the same local minima, guaranteeing that all the local minima are explored. The deepest location identifies the best local minimum. Fig. 1 shows the map of the local minima of the reduced χ^2 for our first observing epoch, in which the companion is detected at more than 50σ . The measured astrometric positions are listed in Table 2.

These observations also provides us the diameter of the Cepheid at various epochs of its pulsation phase. However, the angular diameter variation is at most 0.02 mas for such a short period Cepheid. Visibility measurements with an accuracy better than 2% would be required to detect such a variation, and this is well below the average MIRC accuracy for a 4.7 mag star. We therefore determined the average UD diameter simply by cal-

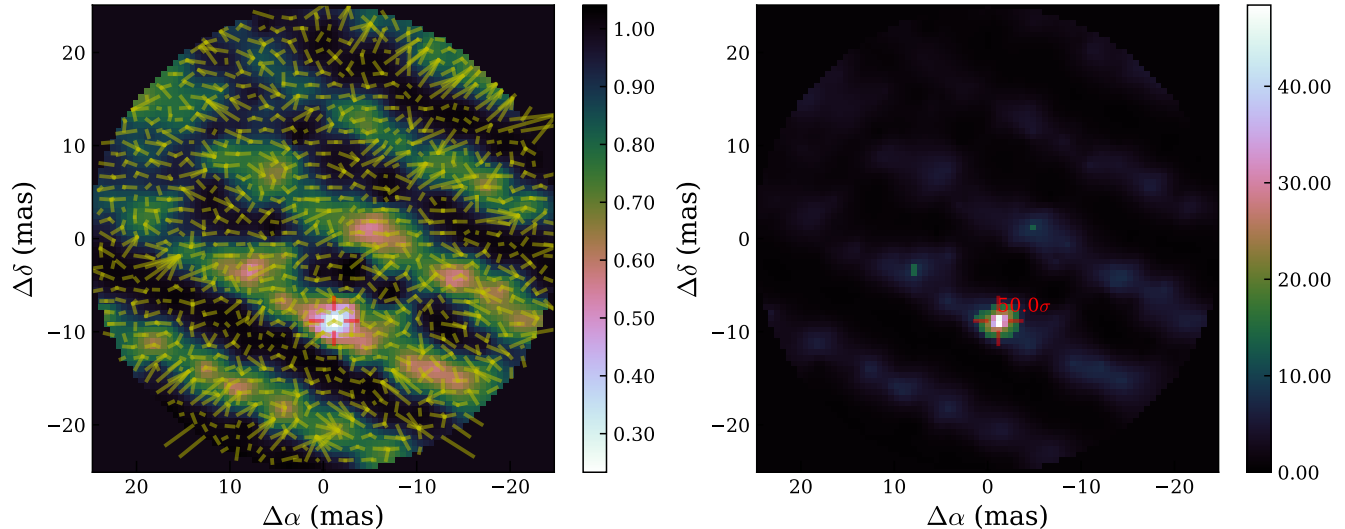


Figure 1. χ_r^2 map of the local minima (left) and detection level map (right) of V1334 Cyg for the July 2012 observations. The yellow lines represent the convergence from the starting points to the final fitted position, as explained in Gallenne et al. (2015). The maps were reinterpolated in a regular grid for clarity. The axis limits were chosen according to the location of the companion.

culating the mean and standard deviation of our measurements. The conversion from UD to limb-darkened (LD) angular diameter is done by using a linear-law parametrization with the LD coefficient $u_\lambda = 0.2423$, as explained in Gallenne et al. (2013b). We measured $\theta_{LD} = 0.524 \pm 0.039$ mas. Combined with our measured distance, we estimate the linear radius of the Cepheid to be $R_1 = 40.6 \pm 3.0 R_\odot$.

Table 2. Measured astrometric positions and flux ratio (H -band) of the companion of V1334 Cyg.

MJD	$\Delta\alpha$ (mas)	$\Delta\delta$ (mas)	f (%)
56135.461	-1.140 ± 0.142	-8.832 ± 0.043	3.39 ± 0.60
56201.211	-0.099 ± 0.029	-8.365 ± 0.023	3.32 ± 0.37
56505.474	4.450 ± 0.116	-3.695 ± 0.050	2.99 ± 0.52
56560.226	4.935 ± 0.008	-2.615 ± 0.010	3.19 ± 0.16
56855.307	5.818 ± 0.023	3.786 ± 0.024	2.89 ± 0.26
56860.413	5.815 ± 0.033	3.844 ± 0.028	3.31 ± 0.46
56862.406	5.744 ± 0.028	3.937 ± 0.025	3.31 ± 0.15
56931.307	5.094 ± 0.027	5.062 ± 0.018	3.45 ± 0.35
57228.411	-1.280 ± 0.064	4.306 ± 0.037	3.08 ± 0.25
57235.412	-1.398 ± 0.022	4.130 ± 0.014	3.12 ± 0.15
57587.410	-6.008 ± 0.035	-4.987 ± 0.042	4.26 ± 0.58

2.2. Visible spectroscopy

High quality optical spectra were collected using the fiber-fed high-resolution SOPHIE and HERMES spectrographs (Bouchy et al. 2009; Raskin et al. 2011). A

total of 14 measurements have been obtained from 2013 to 2016 with SOPHIE, mounted on the 1.93 m telescope of the Observatoire de Haute Provence (France). The instrument covers visible wavelengths with a spectral resolution $R \sim 75\,000$ (high-resolution mode). We also collected 56 spectra from 2010 to 2017 with HERMES ($R \sim 85\,000$), mounted on the Flemish 1.2 m telescope of the Roque de los Muchachos Observatory (La Palma, Canary Islands). Exposure times of a few minutes allowed a signal-to-noise ratio per pixel at 550nm > 50 . Data were reduced using the dedicated HERMES and SOPHIE pipelines.

Radial velocities (RVs) were estimated using the cross-correlation method. We created our own weighted binary mask by selecting unblended lines from high-resolution synthetic spectra ($R \sim 120\,000$) covering the wavelength range 4500-6800 Å. The cross-correlation function is then fitted by a Gaussian whose minimum value gives an estimate of the radial velocity. Uncertainties include photon noise and internal drift. In addition, the standard deviation of the residual of the pulsation fit was also quadratically added. Measurements are listed in Table 3. The zero point difference between the instruments has been estimated using the systemic velocity as follows. We first performed an orbital fit using only the SOPHIE data, then with only the HERMES ones. We measured a difference of 30 m s^{-1} , which we added to the SOPHIE measurements, taking Hermes data as the reference.

2.3. Ultraviolet spectroscopy

Table 3. Radial velocities of V1334 Cyg.

BJD-2400000.5 (day)	V_1 (km s ⁻¹)	Instr.	BJD-2400000.5 (day)	V_1 (km s ⁻¹)	Instr.
56496.018	0.683 ± 0.013	SOPHIE	56493.168	-2.082 ± 0.017	HERMES
56497.065	-7.105 ± 0.014	SOPHIE	56493.229	-2.651 ± 0.017	HERMES
56498.038	-5.153 ± 0.013	SOPHIE	56493.917	-7.496 ± 0.017	HERMES
56500.021	-4.386 ± 0.014	SOPHIE	56494.000	-7.555 ± 0.017	HERMES
56586.910	-9.291 ± 0.024	SOPHIE	56494.201	-7.175 ± 0.017	HERMES
56587.880	-8.900 ± 0.013	SOPHIE	56494.235	-7.101 ± 0.017	HERMES
56588.955	-2.648 ± 0.012	SOPHIE	56494.908	-3.740 ± 0.016	HERMES
56829.089	-10.946 ± 0.011	SOPHIE	56495.235	-1.540 ± 0.016	HERMES
56830.100	-17.562 ± 0.011	SOPHIE	56495.958	0.736 ± 0.017	HERMES
56831.091	-17.999 ± 0.021	SOPHIE	56496.055	0.588 ± 0.017	HERMES
57200.059	-15.583 ± 0.018	SOPHIE	56496.237	-0.136 ± 0.017	HERMES
57201.084	-14.880 ± 0.012	SOPHIE	56855.146	-14.091 ± 0.008	HERMES
57644.919	10.503 ± 0.011	SOPHIE	56856.156	-13.070 ± 0.008	HERMES
57649.968	5.864 ± 0.018	SOPHIE	56857.153	-20.297 ± 0.008	HERMES
55522.925	7.657 ± 0.017	HERMES	56858.150	-16.430 ± 0.008	HERMES
56096.209	10.746 ± 0.017	HERMES	56859.150	-11.935 ± 0.008	HERMES
56097.196	2.917 ± 0.017	HERMES	56860.154	-18.842 ± 0.008	HERMES
56098.189	5.372 ± 0.016	HERMES	56861.155	-18.484 ± 0.008	HERMES
56099.203	10.961 ± 0.016	HERMES	56862.149	-12.708 ± 0.007	HERMES
56488.993	0.782 ± 0.016	HERMES	56864.136	-20.193 ± 0.007	HERMES
56489.095	1.116 ± 0.016	HERMES	56864.219	-19.919 ± 0.007	HERMES
56489.193	1.258 ± 0.016	HERMES	56865.146	-14.547 ± 0.006	HERMES
56489.978	-3.563 ± 0.017	HERMES	56865.223	-14.083 ± 0.007	HERMES
56490.062	-4.360 ± 0.017	HERMES	56866.142	-13.293 ± 0.007	HERMES
56490.226	-5.724 ± 0.017	HERMES	56866.220	-13.844 ± 0.008	HERMES
56490.951	-7.091 ± 0.017	HERMES	57178.090	-14.217 ± 0.012	HERMES
56491.001	-6.878 ± 0.016	HERMES	57179.098	-9.793 ± 0.011	HERMES
56491.039	-6.721 ± 0.016	HERMES	57180.091	-16.895 ± 0.009	HERMES
56491.198	-5.855 ± 0.016	HERMES	57181.155	-15.587 ± 0.007	HERMES
56491.934	-1.301 ± 0.016	HERMES	57182.187	-9.629 ± 0.007	HERMES
56492.004	-0.835 ± 0.016	HERMES	58026.968	2.570 ± 0.026	HERMES
56492.041	-0.592 ± 0.016	HERMES	58035.037	9.640 ± 0.014	HERMES
56492.952	-0.433 ± 0.017	HERMES	58070.845	3.797 ± 0.022	HERMES
56493.000	-0.770 ± 0.017	HERMES	58071.917	9.969 ± 0.015	HERMES
56493.037	-1.013 ± 0.017	HERMES	58073.878	2.356 ± 0.017	HERMES

Measuring the orbital velocity of the companion is also necessary to resolve the degeneracy between the distance and the masses. Due to its early spectral type, the companion is best detected at ultraviolet (UV) wavelengths, which is only possible with the HST because of the UV absorption by the Earth's atmosphere. Four observations were obtained with the Space Telescope Imaging Spectrograph (STIS) on board HST. We used the high resolution echelle mode E140H in the wavelength region 1163 to 1357Å. They were taken over three annual HST cycles. Two were made very close in time to determine

whether the companion is itself a short period binary. The exposure time for each was 1800 seconds. Data reductions was made in the same way as in [Evans et al. \(2018a\)](#). Briefly, for each epoch, the 49 echelle orders were combined into a single one dimension spectrum using an appropriate blaze function. Then, interstellar lines, which are recognizably narrow compared with the stellar lines, were removed by interpolation.

The absolute stability of the STIS E140H wavelength scale over time can be checked by comparing the wavelengths of interstellar features in different observations

of the same star. Using these observations of V1334 Cyg and similar observations of the hot companion to S Mus, Proffitt et al. (2017) found the dispersion in the zero-point of E140H velocity measurements to be about 0.24 km s^{-1} .

Radial velocities of the companion were determined by cross-correlation. No synthetic template spectrum was used, instead we estimated a velocity difference with respect to the first observation by cross-correlating each of the spectra against it. This was typically done in 11 segments of approximately 10 \AA , which were inspected to optimize the features so that a feature was not divided by a boundary. Mild smoothing was included (50 point smoothing) since the stellar features are broad. These velocity differences are directly used in our orbital model fitting. The uncertainties were derived from the differences between segments. Measurements are listed in Table 4.

Table 4. Differential radial velocities of the companion of V1334 Cyg.

MJD	ΔV_2 (km s^{-1})
56916.757	0.00 ± 0.51
56918.548	0.47 ± 0.43
57251.901	-8.70 ± 0.39
57645.272	-28.36 ± 0.71

3. MODEL FITTING AND ORBITAL SOLUTIONS

We performed a combined fit of the interferometric orbit and RVs of the Cepheid and its companion. The fit includes three models with shared parameters for our three datasets. The first models the Cepheid radial velocities only (SOPHIE and HERMES data), i.e. the orbit around the system barycentre and the pulsation. The second models the companion's relative radial velocities (STIS data), i.e. the orbit of the secondary around the system barycentre. The last model fits the relative astrometric orbit of the companion determined from interferometry (MIRC data).

Our radial velocity model of the primary star, the Cepheid, is defined as

$$V_1(t) = V_{1,\text{orb}} + V_{\text{puls}},$$

with the orbital radial velocity $V_{1,\text{orb}}$ and the pulsation velocity V_{puls} , which are expressed with

$$V_{1,\text{orb}} = K_1 [\cos(\omega + \nu) + e \cos \omega] + v_\gamma,$$

$$V_{\text{puls}} = \sum_{i=1}^n [A_i \cos(2\pi i \phi_{\text{puls}}) + B_i \sin(2\pi i \phi_{\text{puls}})],$$

with K_1 the semi-amplitude of the Cepheids orbit due to the companion, ω the argument of periastron of the companions orbit, ν the true anomaly of the companion, e the eccentricity of the orbit, v_γ the systemic velocity, the pulsation phase $\phi_{\text{puls}} = (t - T_0)/P_{\text{puls}}$ (modulo 1), and (A_i, B_i) the amplitude of the Fourier series. The true anomaly is defined implicitly with the following three Keplerian parameters P_{orb}, T_p and e and Kepler's equation:

$$\tan \frac{\nu(t)}{2} = \sqrt{\frac{1+e}{1-e}} \tan \frac{E(t)}{2}$$

$$E(t) - e \sin E(t) = \frac{2\pi(t - T_p)}{P_{\text{orb}}},$$

where $E(t)$ is the eccentric anomaly, T_p the time of periastron passage, t the time of radial velocity observations, and P_{orb} the orbital period.

Following the linear parametrization developed in Wright & Howard (2009) for $V_{1,\text{orb}}$, and including now the pulsation, our model can be simplified to

$$V_1(t) = C_1 \cos \nu + C_2 \sin \nu + v_0$$

$$+ \sum_{i=1}^2 [A_i \cos(2\pi i \phi_{\text{puls}}) + B_i \sin(2\pi i \phi_{\text{puls}})] \quad (1)$$

where we restricted the Fourier series to $n = 2$ as the pulsation curve is sinusoidal-like for such short-period Cepheid (see also Fig. 4). The parameters C_1, C_2 and v_0 are related to the Keplerian parameters through the relations (Wright & Howard 2009):

$$C_1 = K_1 \cos \omega,$$

$$C_2 = -K_1 \sin \omega,$$

$$v_0 = v_\gamma + K_1 e \cos \omega,$$

which can be converted back with (ω chosen so that $\sin \omega$ has the sign of the numerator):

$$K_1 = \sqrt{C_1^2 + C_2^2},$$

$$\tan w = \frac{-C_2}{C_1},$$

$$v_\gamma = v_0 - K_1 e \cos \omega.$$

The fitted parameters are therefore defined as $(C_1, C_2, v_0, A_1, B_1, A_2, B_2, P_{\text{orb}}, T_p, e, P_{\text{puls}})$. T_0 is kept fixed in the fitting process to avoid degeneracy with T_p . The value is chosen such that $V_{\text{puls}}(0) = 0$.

For the second model, as we cross-correlated each STIS spectra with respect to the first observation, our relative radial velocities of the companion can be parametrized with:

$$\Delta V_{2,\text{orb}}(t) = V_{2,\text{orb}}(t) - V_{2,\text{orb}}(t_0), \quad (2)$$

$$V_{2,\text{orb}}(t) = \frac{K_1}{q} [\cos(\omega + \nu) + e \cos \omega] + v_\gamma \quad (3)$$

with the mass ratio $q = M_2/M_1$, and t_0 the time of the first STIS measurement. This adds the new parameter, q as fitted parameter.

Finally, the astrometric positions of the companion as measured from interferometry are modelled with the following equations (Heintz 1978):

$$\Delta\alpha = BX + GY, \quad (4)$$

$$\Delta\delta = AX + FY, \quad (5)$$

where (X, Y) are the elliptical rectangular coordinates defined as:

$$\begin{aligned} X(t) &= \cos E(t) - e \\ Y(t) &= \sqrt{1 - e^2} \sin E(t) \end{aligned}$$

The Thiele-Innes constants are parametrized in term of the orbital elements with:

$$\begin{aligned} A &= a_{\text{as}}(\cos \Omega \cos \omega - \sin \Omega \sin \omega \cos i), \\ B &= a_{\text{as}}(\sin \Omega \cos \omega + \cos \Omega \sin \omega \cos i), \\ F &= a_{\text{as}}(-\cos \Omega \sin \omega - \sin \Omega \cos \omega \cos i), \\ G &= a_{\text{as}}(-\sin \Omega \sin \omega + \cos \Omega \cos \omega \cos i), \end{aligned}$$

with a_{as} the angular semi-major axis in arcsecond (as), Ω the position angle of the ascending node, and i the orbital inclination, which are determined through the inversion relations:

$$\begin{aligned} \tan(\Omega + \omega) &= \frac{B - F}{A + G}, \\ \tan(\Omega - \omega) &= \frac{B + F}{A - G}, \\ k_1 &= \frac{1}{2}(A^2 + B^2 + F^2 + G^2), \\ k_2 &= AG - BF, \\ a_{\text{as}}^2 &= \sqrt{k_1^2 - k_2^2} + k_1, \\ \cos i &= \frac{k_2}{a_{\text{as}}^2}, \end{aligned}$$

with the quadrant of $\Omega \pm \omega$ determined from the signs of the numerators. These Thiele-Innes constants are included to the parameter list, giving a total of 16 variables to fit.

The distance and masses are then given by:

$$\begin{aligned} K_2 &= \frac{K_1}{q}, \\ a_{\text{au}} &= \frac{9.191966 \times 10^{-5} (K_1 + K_2) P_{\text{orb}} \sqrt{1 - e^2}}{\sin^3 i}, \\ d &= \frac{a_{\text{au}}}{a_{\text{as}}}, \\ M_T &= M_1 + M_2 = \frac{a^3 d^3}{P_{\text{orb}}^2}, \\ M_1 &= \frac{M_T}{1 + q}, \\ M_2 &= q M_1, \end{aligned}$$

with K_2 the semi-amplitude of the companion, a_{au} the linear semi-major axis in astronomical unit (au), d the distance to the system, and M_1 and M_2 the Cepheid and companion mass, respectively.

We then applied a classical Monte Carlo Markov Chain (MCMC) technique to fit the 16 model parameters characterizing the standard orbital elements and the pulsation of the Cepheid. Our best fit values were obtained from the median values of the MCMC distribution (100 000 samples with uniform priors), while uncertainties were derived from the maximum value between the 16 % and 84 % percentiles. Our best fit is shown in Fig. 2. We determined a distance $d = 720.35 \pm 7.84$ pc ($\pm 1.1\%$), which is the most accurate model-independent distance of a Cepheid. This corresponds to a parallax of $\pi = 1.388 \pm 0.015$ mas. The two component masses were also determined with a high accuracy: $M_1 = 4.288 \pm 0.133 M_\odot$ ($\pm 3.1\%$) and $M_2 = 4.040 \pm 0.048 M_\odot$ ($\pm 1.2\%$), with the primary star being the Cepheid. The other derived orbital and pulsation parameters are listed in Table 5.

Note that our orbit solutions are rather good agreement with the previous estimates of the spectroscopic elements determined by Evans (2000, also listed in Table 5). It is worth mentioning that we did not use additional RVs measurements from the literature for several reasons: 1) they are usually not very precise, 2) we wanted to use a dataset as uniform as possible (i.e. RVs estimated the same way), 3) the effect on the RVs of a possible third component is reduced, and 4) we also avoid possible bias from the changing pulsation period of the Cepheid by limiting the time range.

4. DISCUSSION AND CONCLUSIONS

This is the first time that a binary Cepheid is resolved both spatially and spectroscopically. This provides the most accurate model-independent distance of a Cepheid, and the most precise masses for a Galactic Cepheid.

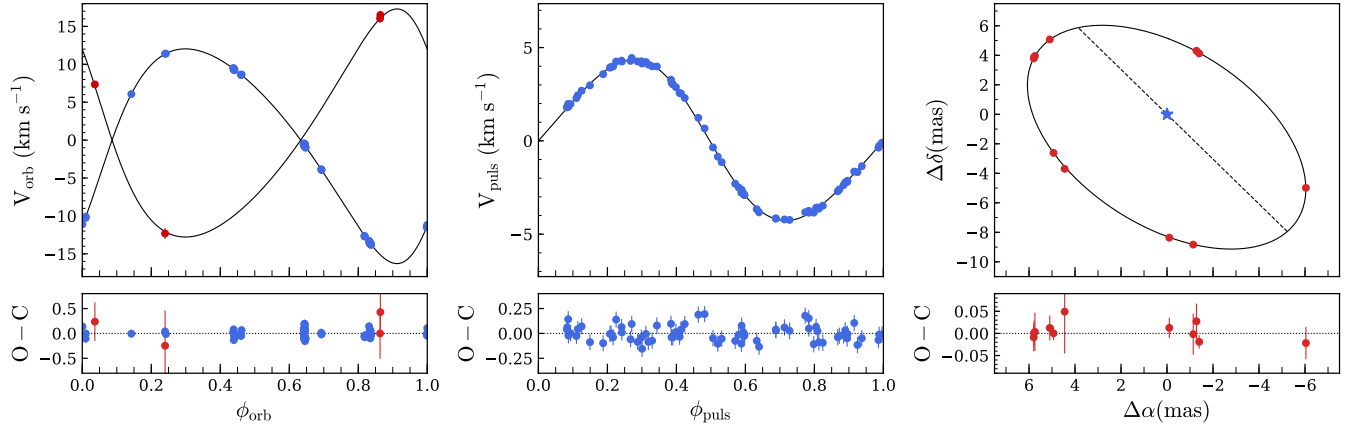


Figure 2. Result of our combined fit. Left: fitted (solid lines) and extracted primary (blue dots) and secondary (red dots) orbital velocity. Middle: fitted (solid line) and extracted (blue dots) pulsation velocity. Right: relative astrometric orbit of V1334 Cyg Ab.

Table 5. Final estimated parameters of the V1334 Cyg system. Index 1 designates the Cepheid and index 2 its main sequence companion. Note that T_0 is kept fixed.

Pulsation		
P_{puls} (days)	3.33242 ± 0.00002	
T_0 (JD)	2 445 000.550	
A_1 (km s ⁻¹)	-0.20 ± 0.35	
B_1 (km s ⁻¹)	4.22 ± 0.04	
A_2 (km s ⁻¹)	-0.01 ± 0.08	
B_2 (km s ⁻¹)	-0.42 ± 0.01	
Orbit		
	Evans+ (2000)	This work
P_{orb} (days)	1937.5 ± 2.1	1932.8 ± 1.8
T_p (JD)	2443607 ± 14	2453316.75 ± 4.1
e	0.197 ± 0.009	0.233 ± 0.001
ω (°)	226.3 ± 2.9	229.8 ± 0.3
K_1 (km s ⁻¹)	14.1 ± 0.1	14.168 ± 0.014
K_2 (km s ⁻¹)	–	15.036 ± 0.304
v_γ (km s ⁻¹)	-1.8 ± 0.1	-2.65 ± 0.01
Ω (°)	–	213.17 ± 0.35
i (°)	–	124.94 ± 0.09
a (mas)	–	8.54 ± 0.04
a (au)	–	6.16 ± 0.07
d (pc)	–	720.35 ± 7.84
q	–	0.942 ± 0.020
M_1 (M_\odot)	–	4.288 ± 0.133
M_2 (M_\odot)	–	4.040 ± 0.048

Here we do not discuss or compare the evolutionary status of this system using our mass measurements as this can be found in [Evans et al. \(2018b\)](#).

With such a high distance accuracy, it is possible to test the agreement of V1334 Cyg with the existing cal-

ibrations of the P–L relations. We selected a sample of published relations in the V - and K filters, as they are the most frequently used photometric bands ([Sandage et al. 2004](#); [Fouqué et al. 2007](#); [Benedict et al. 2007](#); [Storm et al. 2011](#); [Groenewegen 2013](#); [Gaia Collaboration et al. 2017](#)).

This Cepheid is classified as pulsating in its first overtone, with a slightly shorter period than in the fundamental mode. We therefore converted the period from the fundamental to the first overtone mode following the method used in [Pilecki et al. \(2018\)](#) with a pulsation model, and described in [Taormina et al. \(2018\)](#), resulting in an equivalent fundamental period of $P_0 = 4.7897 \pm 0.0322$ days. The comparison with the empirical relation from [Alcock et al. \(1995\)](#), $P_1/P_0 = 0.720 - 0.027 \log P_0$, is very good with a relative difference $< 1\%$ (or within 1.3σ).

For each tested P–L relation, we determined the predicted absolute magnitudes for P_0 , that are represented in [Fig. 3](#) as black dots. We adopted as uncertainties the scatter of each relation given by the authors, which better represents the observed intrinsic dispersion (the width of the instability strip is the dominating uncertainty in the P–L relations, see e.g [Anderson et al. 2016a](#), for a recent study). We derived the absolute magnitude with the usual relation $M_\lambda = m_\lambda - A_\lambda - 5 \log d + 5$, with m_λ the apparent magnitude at wavelength λ , A_λ the interstellar absorption parameter, and d the distance to the star. We corrected for the interstellar extinction, using $A_V = 3.23E(B - V)$ and $A_K = 0.119A_V$ ([Fouqué et al. 2007](#)), with $E(B - V) = 0.025 \pm 0.009$ ([Kovtyukh et al. 2008](#)). For the V band, we estimated the weighted mean apparent magnitude using a periodic cubic spline fit of four different published light curves ([Henden 1980](#); [Berdnikov 2008](#); [Szabados 1977](#);

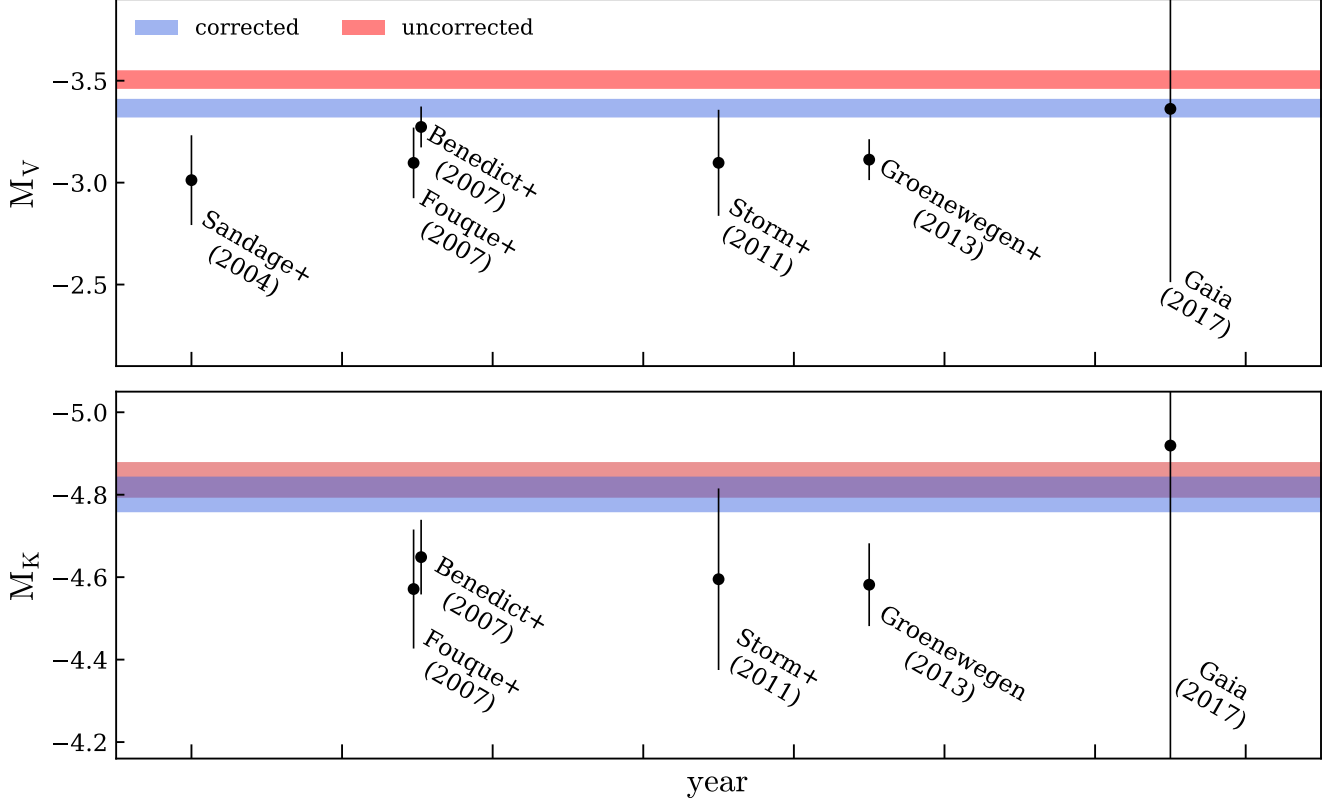


Figure 3. Comparison between the absolute magnitudes of V1334 Cyg predicted from literature P-L relations (black dots) and the present distance measurement (red and blue areas), in two photometric bands. The red and blue areas represent the measured absolute magnitude without and with the subtraction of the companion’s flux contamination, respectively. In the K band, these two regions overlap as the contribution of the companion is smaller.

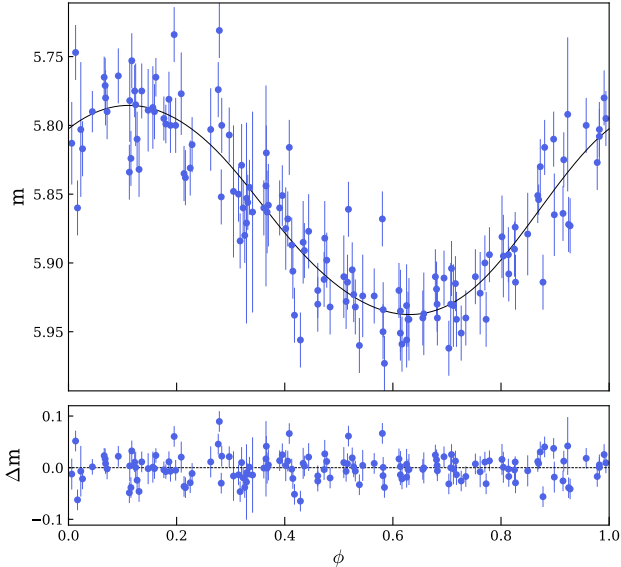


Figure 4. Interpolated V -band light curve of V1334 Cyg.

Kiss 1998). The curve is shown in Fig. 4. We determined the weighted mean dereddened magnitude

$m_{0,V} = 5.781 \pm 0.026$ mag, where the total uncertainty is estimated from the standard deviation of the residual values. There are no light curves available in the near-infrared, so to estimate $m_{0,K}$ we used the value measured by 2MASS (Cutri et al. 2003). These observations were obtained at the pulsation phase 0.36, very close to the phase of the mean radius (0.37), so that the magnitude variation is negligible (within the uncertainty). We estimated $m_{0,K} = 4.451 \pm 0.036$ mag for the mean apparent magnitude in the K band. We then combined with our measured distance to estimate the absolute magnitudes M_λ , for which the total error bar includes the uncertainties on d , A_λ , and m_λ . Fig. 3 shows the difference between M_λ of V1334 Cyg (red area) and the predicted values from literature P-L relations (black dots). However, the contribution of the companion must be subtracted from the combined magnitudes. To correct from the flux contamination of the companion, we used the magnitude difference between the components estimated in V by Evans (1995), $\Delta V = 2.18$ mag, and from our measured H -band average flux ratio from interferometry for the K band, i.e. we assumed $\Delta K = \Delta H = 3.70 \pm 0.11$ mag.

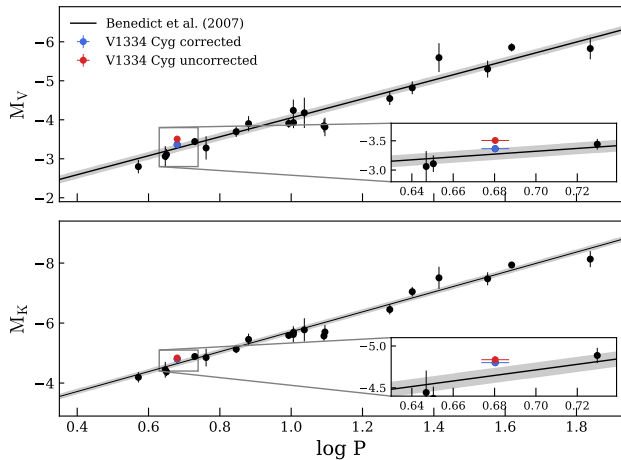


Figure 5. Period-Luminosity relations constructed from direct parallax measurements. This includes our measured distance and those from Casertano et al. (2016); Riess et al. (2018); Benedict et al. (2007, 2002). We also included the distance estimated from light echoes by Kervella et al. (2014). The blue and red dot represent the location of V1334 Cyg for companion-flux corrected and uncorrected magnitude, respectively. Uncertainty on the “fundamentalized” period is also shown. The relation derived by Benedict et al. (2007) is also plotted, with the rms of the residuals in shaded area.

With our distance we can estimate the mean corrected absolute magnitudes of the Cepheid to be $M_V(\text{cep}) = -3.37 \pm 0.05$ mag, $M_H(\text{cep}) = -4.60 \pm 0.05$ mag and $M_K(\text{cep}) = -4.80 \pm 0.04$ mag. Note that this is a rough estimate in K as we assumed the flux ratio to be the same as in H . From the magnitude differences, absolute magnitudes for the companion can also be derived, we found $M_V(\text{comp}) = -1.19 \pm 0.11$ mag, $M_H(\text{comp}) = -0.90 \pm 0.12$ mag and $M_K(\text{comp}) = -1.10 \pm 0.12$ mag. The comparison of the $(V-H)_0$ colour with the intrinsic colour table of Pecaut & Mamajek (2013)² corresponds to a $\sim B7V$ companion of $\sim 3.9 M_\odot$ (the $(V-K)_0$ colour is more uncertain as we assumed the same flux ratio in K and H). This is in good agreement with the spectral type estimate of Evans (1995, mean spectrum of a B7V star) and consistent with our measured mass.

The blue area in Fig. 3 shows the absolute magnitudes of the Cepheid corrected from the flux contamination. We see in the optical band that the relations based on trigonometric parallax measurements (Benedict et al. 2007; Gaia Collaboration et al. 2017) are in agreement with the measured corrected absolute magnitude of the Cepheid. In the infrared, the agreement is better, but still points to an offset, and again direct

distance estimates seem to be more consistent. This can be explained by the fact that photometric measurements are more sensitive to a companion’s flux contamination than are parallax observations, and therefore more strongly biased. The relation from Storm et al. (2011) is in marginal agreement, but also has a large uncertainty. In general, we see a systematic offset in P-L relations derived through photometric measurements (Baade-Wesselink method). In V , the average offset for photometric-based absolute magnitudes only is $\sim 1.5\sigma$, and $\sim 1.2\sigma$ for the K band. It is worth mentioning that increasing the value of $E(B-V)$ would have the effect of increasing the inconsistency of our estimated absolute magnitude and the ones from the P-L relations plotted in Fig. 3. We will also note that as V1334 Cyg is a first overtone pulsator, therefore the amplitude in temperature can be neglected (~ 240 K, Luck et al. 2008).

V1334 Cyg is a clear demonstration that Cepheid companions in binary systems are a significant source of a systematically positive photometric bias to the Cepheid fluxes, as well as contributors to the scattering of the P-L relations. In Fig. 5 we show that correcting for the companion’s flux brings the star closer to the position predicted by the P-L relations, with a stronger effect in V for this specific system. The same effect has been observed in the V and I band by Pilecki et al. (2018) in six LMC Cepheids in eclipsing binary systems.

Finally, the comparison of our orbital parallax with the second Gaia data release ($\pi = 1.151 \pm 0.066$ mas Gaia Collaboration et al. 2018) shows a 3.6σ deviation with our value. Correcting for the Gaia systematic offset of ~ -0.029 mas (estimated for quasars, Lindgren et al. 2018) still leaves a disagreement at 3.2σ . This is however not surprising as binarity is not yet taken into account in the Gaia data reduction process. V1334 Cyg stands out as a demonstration that binarity can significantly impact Cepheid parallaxes. The selection criteria of Gaia Collaboration et al. (2017) of rejecting the known binary Cepheids seems a good approach, as seen in our Fig. 3 with a good agreement in V band between the expected value and our measured corrected absolute magnitude. The final Gaia data release with hundreds of Galactic Cepheid parallaxes promises to constrain the zero point of P-L relations to an exquisite precision. But reddening corrections, biases from orbiting companions (astrometric and photometric), as well as the flux contribution of circumstellar envelopes will limit the accuracy of the P-L relation calibration if not properly taken into account.

This research is based on observations made with SOPHIE spectrograph on the 1.93-m telescope at Ob-

² see also http://www.pas.rochester.edu/~emamajek/EEM_dwarf_UBVIJHK_colors_Teff.txt

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Software: Candid (Gallenne et al. 2015), Emcee (Foreman-Mackey et al. 2013), SearchCal (Bonneau et al. 2006), MIRC pipeline (Monnier et al. 2007)

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