1	Tropical Climate Variability: Interactions across the Pacific, Indian, and Atlantic Oceans
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3	Jules B. Kajtar ^{*1,2} , Agus Santoso ^{1,2} , Matthew H. England ^{1,2} , and Wenju Cai ³
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5	¹ Australian Research Council's Centre of Excellence for Climate System Science, Australia
6	² Climate Change Research Centre, University of New South Wales, NSW, Australia
7	³ CSIRO Marine and Atmospheric Research, Aspendale, Victoria, Australia
8	
9	* Corresponding author: Jules B. Kajtar; <u>j.kajtar@unsw.edu.au</u> ; Telephone: +61 (0)2 9385 9766; Facsimile:
10	+61 (0)2 9385 8969
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30 Abstract

Complex interactions manifest between modes of tropical climate variability across the Pacific, Indian, and Atlantic Oceans. For example, the El Niño-Southern Oscillation (ENSO) extends its influence on modes of variability in the tropical Indian and Atlantic Oceans, which in turn feed back onto ENSO. Interactions between pairs of modes can alter their strength, periodicity, seasonality, and ultimately their predictability, yet little is known about the role that a third mode plays. Here we examine the interactions and relative influences between pairs of climate modes using ensembles of 100-year partially coupled experiments in an otherwise fully coupled general circulation model. In these experiments, the air-sea interaction over each tropical ocean basin, as well as pairs of ocean basins, is suppressed in turn. We find that Indian Ocean variability has a net damping effect on ENSO and Atlantic Ocean variability, and conversely they each promote Indian Ocean variability. The connection between the Pacific and the Atlantic is most clearly revealed in the absence of Indian Ocean variability, and our model runs suggest a weak damping influence by Atlantic variability on ENSO, and an enhancing influence by ENSO on Atlantic variability.

59 1. Introduction

60 Modes of tropical climate variability, such as the El Niño-Southern Oscillation (ENSO), the Indian Ocean Basinwide Mode (IOBM), the Indian Ocean Dipole (IOD), and the Atlantic Equatorial Mode (sometimes 61 62 referred to as the Atlantic Zonal Mode, or Atlantic Niño) interact most readily via the atmosphere. Sea 63 surface temperature anomalies (SSTAs) in the tropics drive changes in the Walker Circulation, which in turn influence SSTAs in remote regions, hence forming a teleconnection (e.g., Lau and Nath 1996; Klein et al. 64 1999; Alexander et al. 2002). Through this atmospheric teleconnection mechanism, modes in one ocean 65 basin can be damped, enhanced, or even entirely generated by a mode in another ocean basin. Oceanic 66 67 pathways also provide the means for inter-basin interactions, although at lag times typically beyond several months. Whilst the literature is rich on the interactions between ENSO and each of the tropical modes in the 68 Indian and Atlantic Oceans, few studies have examined the possible interactions between the Indian and 69 Atlantic modes, apart from the Atlantic influence on the Indian monsoon (e.g. Kucharski et al. 2008, Losada 70 71 et al. 2010). Furthermore, if all three tropical basins are strongly coupled, then interactions between any two basins can be influenced or modulated by the third ocean basin. This aspect, which has previously been 72 neglected in the literature, will be considered here. A better understanding of these highly complex inter-73 basin interactions is relevant for improving climate prediction, especially for ENSO (e.g. Izumo et al. 2010, 74 75 Keenlyside et al. 2013).

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ENSO is a manifestation of complex internal dynamics within the Pacific Ocean, but it is now widely 77 recognised that modes of variability in the Indian and Atlantic oceans also influence the air-sea feedback 78 79 processes that govern ENSO characteristics (e.g., Dommenget et al. 2006; Izumo et al. 2010; Ding et al. 2012; Santoso et al. 2012; McGregor et al. 2014; Polo et al. 2014; Terray et al. 2015; Kucharski et al. 2016). 80 81 It is not possible to determine the interaction dynamics from observations and standard climate models alone due to the coupled nature of these modes of variability. Instead, the problem needs to be studied in model 82 83 experiments whereby individual modes are nullified. This can be achieved by eliminating air-sea interactions 84 over a region of variability, so that the atmosphere does not respond to the SSTAs associated with these remote modes. Such earlier 'partial coupling' studies concluded that Indian Ocean variability tends to 85 enhance ENSO (Barsugli and Battisti 1998; Yu at al. 2002; Wu and Kirtman 2004). However, many of these 86 studies were based on a single experiment, with short run-time (approximately 50 years), or using simplified 87

GCMs. The robustness of this conclusion was put into doubt with higher resolution models and longer experiments (Dommenget et al. 2006; Santoso et al. 2012; Terray et al. 2015). While the IOD may play a role in initiating ENSO events (e.g., Luo et al. 2010; Izumo et al. 2010), Santoso et al. (2012) showed that Indian Ocean variability as a whole exerts a net damping influence on ENSO via the IOBM.

93 The climatic connections between the Pacific and Indian Oceans are further complicated by the presence of the Indonesian Throughflow (ITF). The ITF typically transports a large volume of water (Potemra 1999; 94 Gordon 2005; Wijffels et al. 2008) and heat (Vranes et al. 2002; England and Huang 2005) from the Pacific 95 96 to the Indian Ocean, but it exhibits interannual variability which is linked to ENSO and the Indian Ocean modes (Meyers 1996; England and Huang 2005; van Sebille et al. 2014; Sprintall and Revelard 2014). Its 97 significance in the global context is exhibited by model experiments with a blocked ITF, where the mean 98 climate and modes of variability are greatly altered (Song et al. 2010; Santoso et al. 2011). It has been 99 100 recently suggested that Indian Ocean variability can influence ENSO via Kelvin wave propagation through 101 the ITF (Yuan et al. 2013) at longer time lags. However, a recent study by Izumo et al. (2015) argued that the 102 atmospheric bridge mechanism is more dominant for Indo-Pacific interactions. The robustness of the 103 atmospheric bridge is attested by the fact that the Indo-Pacific feedback interactions persist even in the 104 absence of the ITF (Kajtar et al. 2015).

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106 On interannual time scales, ENSO and variability in the tropical Atlantic basin interact via the atmospheric bridge. The Atlantic Equatorial Mode (AEM), which is the dominant mode of variability in the tropical 107 108 Atlantic Ocean (Zebiak 1983), displays ENSO-like characteristics (Keenlyside and Latif 2007; Jansen et al. 2009), with SSTAs across the central to eastern equatorial Atlantic Ocean. The relationship between ENSO 109 and the AEM is complex, and predicting the state of the Atlantic Ocean based on the precedence of an ENSO 110 event is not reliable (Saravanan and Chang 2000; Chang et al. 2006; Rodrigues et al. 2011; Lübbecke and 111 112 McPhaden 2012; Taschetto et al. 2015). In contrast, knowledge of the Atlantic Ocean state can improve ENSO prediction (Frauen and Dommenget 2012; Keenlyside et al. 2013). Frauen and Dommenget (2012) 113 used a GCM, albeit with a simplified ocean model, to demonstrate that the Atlantic Ocean has no net 114 115 discernible influence on ENSO characteristics, but does influence the state of the Pacific Ocean that is 116 relevant for ENSO prediction. Other studies have shown that an Atlantic Niño (the warm phase of the AEM)

tends to favour the development of a La Niña in the Pacific (Ding et al. 2012; Polo et al. 2014). Furthermore,
it appears that this relationship has strengthened in recent decades (Rodríguez-Fonseca et al. 2009) and is
likely associated with multi-decadal variability (Latif 2001; Martín-Rey et al. 2014; McGregor et al. 2014).

Despite the extensive literature on the Indian and Atlantic Ocean influence on ENSO, few studies (e.g. 121 Dommenget et al. 2006; Frauen and Dommenget 2012; Terray et al. 2015) have examined the role of each 122 within the same modelling framework. These studies agree that Indian Ocean variability damps ENSO, but 123 124 Terray et al. (2015) point to a weak damping influence by the Atlantic on ENSO, whereas the other studies 125 found none. Terray et al. (2015) also show that decoupling either basin tends to shift ENSO to longer periods. It is also important to note, however, that many coupled models suffer from strong SST biases in the 126 127 equatorial Atlantic (Ritcher et al. 2014), and hence any comparison between coupled and partially coupled 128 experiments may be compromised by internal model biases.

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In this study we will examine, for the first time, the interactions of tropical models of variability between 130 131 each of the Pacific, Atlantic, and Indian Oceans, all within the same coupled GCM. In particular we expand 132 on very small pool of literature on the Indo-Atlantic coupling (e.g. Kucharski et al. 2008, McGregor et al. 133 2014), which may play a role in modulating the interactions between ENSO and other modes of variability. We ran sets of five-member, 100-year, partially coupled experiment ensembles in an otherwise fully coupled 134 135 GCM. In addition to sequentially nullifying the air-sea interactions over each tropical ocean basin individually, we ran further experiments with pairs of ocean basins decoupled. The rationale behind 136 137 decoupling pairs of ocean basins is to eliminate the influence that a third ocean may play on interactions between modes in the first two, thus helping to infer the role of the third ocean basin. In essence, this study 138 aims to build a global picture of interactions between climate modes across the tropics with potential 139 implications for their predictability. We focus on the dominant modes in the tropics, namely ENSO, the IOD, 140 the IOBM, and the AEM, since they readily interact via induced changes to the Walker Circulation. We 141 focus our analysis on changes to the strength and period of the modes by examining the monthly standard 142 deviations and power spectral densities of the relevant SST indices (Section 3). We then demonstrate the 143 144 zonal wind stress influences by which the modes interact (Section 4).

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146 **2. The climate simulations**

147 **2.1 Model description and experimental setup**

The simulations were performed with version 1.2 of the Commonwealth Scientific and Industrial Research 148 Organisation (CSIRO) Mk3L general circulation model (GCM; Phipps 2010; Phipps et al. 2013). The 149 150 atmospheric GCM (AGCM) has a resolution of ~ 5.6° longitude × ~ 3.2° latitude, with 18 levels in the hybrid vertical coordinate. The oceanic GCM (OGCM) has resolution ~ 2.8° longitude \times ~ 1.6° latitude, and 21 151 152 levels in the vertical z-coordinate. The OGCM was first spun up for 7000 years, and then the ocean surface 153 state was used to spin up the AGCM for 100 years. To maintain a more realistic climatology and minimise 154 drifts, the AGCM and OGCM were then coupled with constant, but seasonally varying, flux adjustment terms applied to the surface heat flux, surface salinity tendency, and surface momentum fluxes. The terms 155 156 are derived at the end of the spin-up phase, and not restored towards observations during the coupled run.

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158 Following the initial spin-up, the coupled model was integrated for 1550 years, with CO₂ fixed at the preindustrial level of 280 ppm, since here we are focussing on the dynamics without anthropogenic forcing. 159 The last 300 years of this run is referred to as the control simulation (CTRL). This 300-year run was split 160 161 into five 100-year ensemble members, with 50-year intervals for the starting year of each set, i.e. the first set 162 starts at year 1, the second at year 51, and so on, until year 201 for the fifth set. The 100-year partial coupling experiments were then initialized with the climate state at each of these epochs. In these runs, the air-sea 163 164 interaction over a single or pair of ocean basins was suppressed by fixing SST to the climatological seasonal mean field from the first 200 years of the model control run (as per the methodology of Baguero-Bernal et al. 165 166 2002; Behera et al. 2005; Dommenget et al. 2006; Santoso et al. 2012; Kajtar et al. 2015). In this way, the atmosphere responds only to the seasonally varying climatological SST over the decoupled region, and hence 167 168 any modes of variability in that region are nullified.

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We note that one may also choose to nudge SST toward observed climatology. However, as shown by Terray et al. (2015), the decoupling effect on tropical climate variability in a given ocean basin will also contain changes resulting from an altered mean state within that basin. In the case of Terray et al. (2015), nudging the Indian or Atlantic toward their respective observed SST climatology results in a more realistic Pacific climatology. At the same time, the ENSO response becomes stronger, but still exhibits similar tendencies as in the case of nudging toward model climatology. To isolate the effects of only the remote forcing, we chose to perform the decoupling by nudging toward model SST climatology. The flux adjustments in our model assist with maintaining a more realistic climatology, and are applied globally and consistently throughout all experiments, thus ensuring that any of the diagnosed changes are not due to flux adjustments, but to the decoupling of remote SSTAs.

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181 The decoupled regions in our experiments are bounded by 30°S and 30°N, and by the coast to the east and 182 west in each ocean basin. As discussed by Santoso et al. (2012), choosing a particular boundary between the 183 Pacific and Indian Oceans may affect the conclusions reached. Nevertheless, in this study we follow their approach in which the western side of the Maritime Continent is considered part of the eastern Indian Ocean, 184 and the eastern side as part of the western Pacific Ocean. The decoupled Pacific, Indian, and Atlantic Ocean 185 186 experiments are denoted DCPL_{PO}, DCPL_{IO}, and DCPL_{AO} respectively. The experiments where pairs of ocean 187 basins were decoupled follow a similar nomenclature, i.e. DCPL_{PO+AO}, DCPL_{PO+IO}, and DCPL_{AO+IO}. Note that throughout the text, "decoupling" will refer to the suppression of SST variability over a particular ocean 188 basin. We ran 100-year experiments so that the significance of the interactions between low-frequency 189 190 modes could be statistically assessed. The mean climate drift over this period in the Mk3L model is 191 negligible.

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193 **2.2 Model validation**

Mk3L performs relatively well in capturing the mean climatology and tropical modes of variability, albeit 194 195 with some notable biases. As with many GCMs, Mk3L suffers from the "cold tongue" bias, with overly strong trade winds and lower than observed rainfall, associated with anomalously cold SST extending 196 westward from the eastern equatorial region of each ocean basin (Fig. 1). The dry bias appears to be 197 exacerbated over the Maritime Continent by a shallower than observed thermocline depth in the eastern 198 199 Indian Ocean which causes overly cold SST in that region (Santoso et al. 2012). Somewhat expected given the overly strong easterly winds in the Pacific, the mean Indonesian Throughflow rate in the model 200 (approximately 21 Sv with a standard deviation of 1.3 Sv; see also Santoso et al. 2011) is substantially larger 201 202 than the observed estimate of 15 Sv (Sprintall et al. 2014). This is also partly attributed to the coarse model resolution, likely through the joint effect of baroclinicity and relief (JEBAR; England et al. 1992; Santoso et 203

204 al. 2011).

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206 In order to evaluate the Mk3L model performance in simulating the relevant modes of tropical climate variability, we compare against observations (HadISST) and a set of the historical experiments, over the 207 208 period 1900-1999, from the Coupled Model Intercomparison Project, phase 5 models (CMIP5; bcc-csm1-1, CanESM2, CCSM4, CNRM-CM5, FGOALS-g2, FGOALS-s2, GFDL-CM3, GFDL-ESM2G, GFDL-209 210 ESM2M, HadCM3, HadGEM2-CC, HadGEM2-ES, IPSL-CM5A-LR, IPSL-CM5A-MR, IPSL-CM5B-LR, 211 MIROC4h, MIROC5, MIROC-ESM-CHEM, MPI-ESM-LR, MPI-ESM-MR, MPI-ESM-P, MRI-CGCM3, 212 NorESM1-M, and NorESM1-ME). The modes and their associated characteristic SST indices are given in Table 1. The Mk3L ENSO SST variability has weaker magnitude and peaks 2-3 months earlier than the 213 observed, although it falls well within the overall CMIP5 model range (Fig. 2a). Its period in Mk3L is also 214 slightly longer, with the strongest signal in the 5-7 year band (Fig. 2b). The Atlantic Equatorial Mode (AEM) 215 216 is also weaker than observed, but has a similar seasonal cycle to the observed, with peak variability around May to August (Fig. 2c). This seasonal cycle is not captured by the CMIP5 multi-model mean, and may be 217 attributed to either strong model biases in the tropical Atlantic (Ritcher et al. 2014) or large model diversity, 218 or both. The observed Atl-3 index shows variability at a range of time-scales with an increasing tendency 219 220 toward interdecadal timescales; a tendency also seen in the CMIP5 ensemble and Mk3L. Spectral peaks common to those of ENSO are seen in the observed and Mk3L, although in Mk3L the signal is clearly 221 occurring in the 5-7 year band (Fig. 2d). Higher-frequency variability is muted in Mk3L, and is captured by 222 only a few of the CMIP5 models. 223

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Over the Indian Ocean, the IOD in Mk3L is stronger than observed, as with many GCMs (Fig. 2e), which is 225 associated with the shallower than observed thermocline in the eastern Indian Ocean (Cai and Cowan 2013). 226 Also associated with a shallower thermocline are overly persistent cool SSTAs in a limited region of the 227 228 eastern Indian Ocean during a positive IOD, thus slightly affecting the representation of the warm phase of 229 the IOBM, and vice versa for the negative IOD (see Fig. 2 of Santoso et al. 2012). Nevertheless, the variability and seasonality of the model IOBM agrees well with observations (Fig. 2g). Like observations, 230 231 the IOD and IOBM in Mk3L exhibit variability that coincides with ENSO time-scales, but again in Mk3L this is in the 5-7 year band (Figs. 2f,h). This correspondence is not clear in the CMIP5 ensemble due to the 232

large model diversity. The common periodicity signal across the indices, apparent in the observations and
Mk3L, indicates a coupling across the modes. These interactions will be disentangled and studied using our
partial coupling experiments.

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237 We now focus further on evaluating the model's performance in simulating the tropical Atlantic climate, 238 since the performance in the Indo-Pacific region has been previously documented in detail (Santoso et al. 239 2011; 2012). Firstly, note that the SST bias across the tropical Atlantic is slightly different to most other 240 coupled models. The CMIP3 and CMIP5 models tend to show a temperature gradient across the basin that is 241 warmer in the east, opposite to what is observed (see Fig. 1 of Richter et al. 2014). Yet the Mk3L model correctly captures the sign of this gradient (apparent in Fig. 1a, b). The spatial pattern of the model's first 242 243 empirical orthogonal function (EOF) of June to August SST (Fig. 3a), i.e., the AEM peak season, shows 244 some agreement with the observed (Fig. 3b). The southeastern branch is absent in the model, but it largely 245 captures the zonal mode structure along the equator. Many CMIP5 models either do not exhibit the AEM as the first EOF, or at all (Richter et al. 2014). The first EOF explains a similar percentage of the total variance 246 in the model (32%; 24% for EOF-2) and in the observed (47%; 25% for EOF-2). The correlation between the 247 Atl-3 index and the SST field, both averaged over June to August, reveals a good agreement overall in the 248 249 variability pattern between the model and observed (Fig. 3c,d), except for a negative correlation at $\sim 10^{\circ}$ N that is not seen in the reanalysis fields. Despite the shortcomings of the Mk3L model, it broadly captures the 250 251 tropical climate modes, at least within the expected range of state-of-the-art CMIP5 model performance. Finally, we note that its coarse resolution makes it suitable for running large ensembles of century-scale 252 253 experiments, such as those undertaken here.

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255 3. Changes in amplitude and periodicity of the modes

We begin by examining the changes to the amplitude, and then later the periodicity, of the dominant tropical modes in the partial coupling experiments. The modes are each characterised by area-averaged SST indices that display distinct seasonality (outlined in Table 1), and so we examine changes to the monthly standard deviation of the SST indices.

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261 Suppressing tropical Indian Ocean SST variability in the DCPL_{IO} experiments increases the monthly

standard deviation of the Niño-3.4 index relative to CTRL (Fig. 4a), indicating stronger ENSO in DCPL_{IO}.
The increase is statistically significant throughout most of the calendar year (October to July), as indicated
by the separation between the confidence intervals associated with the CTRL and DCPL_{IO} experiments. This
change indicates that the presence of SST variability over the tropical Indian Ocean acts to damp ENSO
variability, in agreement with earlier studies (Dommenget et al. 2006; Terray et al. 2015), and with Santoso
et al. (2012) who used a larger ensemble in the same MK3L model. The mechanisms by which the Indian
Ocean climate modes influence ENSO evolution are described in detail in Section 4.

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270 The influence of the Atlantic Ocean on ENSO is less clear. The confidence interval of Niño-3.4 in the DCPL_{AO} experiments overlaps that of CTRL in each calendar month (Fig. 4a). Closer inspection of the 271 272 individual ensemble members reveals that the variability is weaker relative to CTRL in two members, but 273 stronger in the remaining three (figure not shown). This is in contrast to $DCPL_{IO}$, which exhibits stronger 274 ENSO variability across all ensemble members. The inconsistent changes underscore the importance of ensemble experiments. Terray et al. (2015) found a weak enhancement of ENSO variability, although they 275 perform twin partial coupling experiments, nudged towards either observed or model climatological SST (an 276 issue not assessed here). Dommenget et al. (2006) found no clear overall change in ENSO variance when the 277 278 Atlantic is decoupled in their single 500-year experiment, but they did not examine the change in shorter sub-periods. Nevertheless, it is important to keep in mind that this inconsistent response does not mean that 279 the Atlantic Ocean exerts no influence on ENSO. Frauen and Dommenget (2012), for instance, concluded 280 that the Atlantic Ocean plays a role in the predictability of ENSO, despite having no clear impact on its 281 282 dynamics. The inconsistency can also arise due to multi-decadal variability in Atlantic-Pacific connection (Latif 2001; Rodríguez-Fonseca et al. 2009; Martín-Rey et al. 2014; McGregor et al. 2014). 283

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Care should be taken when inferring the actual impact of the Atlantic Ocean on ENSO based on decoupling the Atlantic alone, without considering changes that could occur in the Indian Ocean resulting from a decoupled Atlantic. The same can be said about decoupling the Indian Ocean to diagnose its isolated effect on ENSO when Atlantic variability is still present in DCPL_{IO}. Further insights into this interplay can be garnered when both the Atlantic and Indian Oceans are decoupled (DCPL_{AO+IO}). As shown in Fig. 4a, decoupling both the Indian and Atlantic Oceans enhances the variability of ENSO more strongly than

291 decoupling the Atlantic or Indian Ocean alone, which is consistent with Dommenget et al. (2006) and Frauen 292 and Dommenget (2012). Interestingly, the enhanced variability seen in $DCPL_{AO+IO}$ relative to $DCPL_{IO}$ is strongest during May to August, coinciding with the peak season of AEM variability. This could imply that 293 Atlantic variability exerts a more consistent effect on ENSO growth than what is inferred from $DCPL_{AO}$ in 294 295 which Indian Ocean variability is present. It is also possible that the impact of the Indian Ocean on ENSO is in actuality greater without the presence of Atlantic variability. In other words, in DPCL₁₀, the Atlantic 296 297 variability may be altered upon decoupling the Indian Ocean in such a way that it limits the enhancement of ENSO amplitude due to the removal of Indian Ocean variability. Concurrent removal of Atlantic variability 298 then leads to more amplified ENSO in DCPL_{AO+IO} compared to DCPL_{IO}. This would imply that the Atlantic 299 variability has a net damping effect on ENSO in our model – different to what is inferred from the DCPL_{AO} 300 301 runs alone. A damping influence on ENSO by the Atlantic, but smaller than the damping influence by 302 Indian Ocean variability, is consistent with Terray et al. (2015) based on their individual basin decoupling 303 experiments. In any case, our DCPL_{AO+IO} result compared with DCPL_{IO} and DCPL_{AO} suggests that there are potential interactions occurring between the Atlantic and Indian Ocean variability. This will become clearer 304 305 below.

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307 As shown in Fig. 4b, the change in AEM variability is significantly different from CTRL due to removal of the Indian Ocean alone (DCPL_{IO}), implying a potential damping role of Indian Ocean variability on AEM in 308 CTRL. However, it is also possible that the AEM enhancement is due to the enhanced ENSO arising from a 309 decoupled Indian Ocean (Fig. 4a). We argue that it is both, but the Indian Ocean damping is more dominant, 310 311 for the following reasons. Decoupling the Pacific Ocean (DCPLPO) does not result in net significant changes to the AEM, except weakened Atl-3 index outside the boreal summer peak of the AEM (Fig. 4b). If the 312 enhancing effect of ENSO were dominant, then we would expect the weakening of the AEM to be 313 significant. The fact that it is not can be explained by a reduced damping effect of the Indian Ocean (in 314 DCPL_{PO}), since Indian Ocean variability is weakened when the Pacific is decoupled (Fig. 4c,d). Removing 315 Indian Ocean variability while the Pacific remains decoupled (DCPL_{PO+IO}; i.e., no ENSO enhancing effect 316 and no Indian Ocean damping) exhibits a slight increase in AEM amplitude, but the change is not 317 318 statistically significant from either CTRL or DCPL_{IO} duing the peak season. This confirms that the significant AEM enhancement seen in DCPL_{IO} is due to primarily an absence of Indian Ocean variability and 319

- to a lesser extent the presence of a stronger ENSO. Thus, the Indian Ocean variability in our model has a net
 damping effect on AEM, while the ENSO tends to enhance AEM.
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323 Decoupling the Pacific Ocean reduces the amplitude of the Indian Ocean Dipole, consistent with previous 324 studies (Fischer et al. 2005; Behera et al. 2006). The reduction of DMI standard deviation in DCPL_{PO} relative to CTRL is statistically significant for all months, but weakest during boreal autumn when IOD 325 326 peaks (Fig. 4c). Decoupling both the Atlantic and the Pacific (DCPL_{PO+AO}) shows similar effect as DCPL_{PO}, indicating that the Atlantic has little influence on IOD amplitude. This is confirmed by DCPL_{AO}, which 327 328 shows a weak change to the DMI standard deviation. Thus, in our model the Pacific Ocean partly drives IOD variability, but there is still substantial Pacific-independent component. This supports the notion that 329 while ENSO can generate IOD, the IOD itself is intrinsically a mode internal to the Indian Ocean. 330

332 The enhancing influence of Pacific Ocean variability on the IOBM is more pronounced than its influence on the IOD, since the variability of the BWI is reduced during most months, but most importantly throughout its 333 peak season when the Pacific is decoupled (DCPL_{PO} and DCPL_{PO+AO}). There is also a noticeable change in 334 the seasonality of the BWI, which can be seen by comparing DCPL_{PO} and DCPL_{PO+AO} relative to CTRL (Fig. 335 336 4d). While this at a first glance supports the suggestion that the IOBM is largely a response to ENSO (Klein et al. 1999; Lau and Nath 2003; Du et al. 2009), the fact that decoupling the Pacific does not entirely remove 337 the IOBM also suggests that the IOBM can occur without ENSO. There is also an indication of an Atlantic 338 influence. Decoupling the Atlantic while the Pacific remains decoupled further reduces the IOBM 339 340 amplitude, although the associated changes occur outside the peak season of the IOBM. This is supported by DCPL_{AO} result showing reduced variability in those months, thus revealing that Atlantic variability enhances 341 342 the IOBM in CTRL.

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The weak change in seasonality of all modes under each partial coupling experiment (Fig. 4), suggests they are, to a varying extent, internally generated modes in their respective basins. This demonstrates a certain degree of independence between the modes. However, the weak change in seasonality may also be a result of replacing the SSTAs with the model SST climatology. It is important to note that the response may be different if using observed SST climatology for the partial coupling. For example, Terray et al. (2015) find

that the seasonality of ENSO is more pronounced when nudged toward the observed SST climatology.

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351 The changes described above involve shifts in the periodicity of each mode. This is clearly shown in the power spectral densities (PSDs; Fig. 5). The most striking feature is the dominant variability in the 5-7 year 352 353 band in each index in the control experiments, which is also the dominant ENSO frequency in Mk3L. This 354 common periodicity across the three basins suggests the timescale at which the inter-basin coupling occurs. 355 The partial coupling experiments reveal the collapse of variability over this 5-7 year frequency band, 356 resulting in variability tending to be skewed towards longer periodicity. The shift in ENSO variability 357 toward longer periods in the partial coupling experiments relative to CTRL (Fig. 5a) is in agreement with previous studies (e.g., Dommenget et al. 2006; Santoso et al. 2012; Terray et al. 2015). Decoupling the 358 359 Indian Ocean alone enhances variability in the 7-9 year band, thus accounting for the stronger ENSO 360 amplitude as seen in the monthly standard deviation of the Niño-3.4 index (Fig. 5a). The Niño-3.4 PSD 361 reddens in DCPL_{AO} without exhibiting any characteristic frequency, while DCPL_{AO+IO} exhibits a peak in the 9-15 year band. These results suggest that the ENSO evolution in Mk3L could be more sluggish without 362 Indian and Atlantic variability. Both the Indian and Atlantic Ocean variability appear to play a role in setting 363 that 5-7 year ENSO periodicity. Shift toward longer periodicity is also seen in the AEM in the absence of 364 365 Indian and Pacific variability (Fig. 5b), while it is not as apparent for the IOD and IOBM, which involve primarily a collapse in variability, particularly at the 5-7 year timescales (Fig. 5c,d). 366 367 368 Our findings can be summarised as follows: 369 There are interactions between modes of variability across the three tropical oceans. • 370 Indian and Atlantic variability has a net damping effect on ENSO magnitude and increases the • 371 rapidity of ENSO evolution. Indian Ocean variability has a net damping effect on the AEM, while ENSO tends to enhance the 372 • AEM. 373 IOD variability is enhanced by ENSO, but there is little influence by Atlantic variability. 374 • IOBM variability is enhanced by ENSO and to a weaker extent by Atlantic variability. 375 ٠ Decoupling any ocean basin collapses variability in the 5-7 year primary ENSO frequency band and 376 • 377 tends to shift modes toward longer periodicity.

379 4. Interbasin feedback interactions via atmospheric bridge

As has been shown in many previous studies (e.g., Lau and Nath 1996; Klein et al. 1999; Alexander et al. 380 2002; Dayan et al. 2015; Kajtar et al. 2015), interactions between modes of variability across different basins 381 readily occur via the atmospheric bridge. SSTAs across the equatorial seas associated with these modes of 382 383 variability drive anomalies in the Walker Circulation. These atmospheric disturbances generate wind stress 384 anomalies over remote seas, which impinge on the oceanic dynamics by, for example, forcing Kelvin waves. In this section we examine the composite evolution of the SST anomalies (Fig. 6) and the equatorial zonal 385 386 wind stress anomalies (Fig. 7, 8) in each experiment to show how the changes in each of the modes of variability under partial coupling can be explained via alterations to the Walker Circulation that connects the 387 388 climate in the three tropical basins. The composites in Fig. 6-8 are produced by first constructing annual 389 time-series of each SST index, averaged over the corresponding peak seasons in each ensemble and 390 experiment (Table 1), and then selecting the years that exceed one standard deviation of that time-series. For 391 example, El Niño or La Niña events are defined as when the September-December (SOND) average of Niño-392 3.4 is greater or less than one standard deviation. In cases where the threshold is exceeded in two or more 393 consecutive years, only the year with strongest anomaly is included in the composite, to avoid the inclusion 394 of double events. Note that this compositing approach is valid since there are no significant changes to the 395 seasonality of the modes under the different partial coupling experiments (Fig. 4).

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Analysis of the correlations between pairs of modes in the control simulation is also necessary to provide 397 398 information on the typical interactions, helping to interpret the decoupling experiment results. For example, if the warm phase of the IOBM acts to damp El Niño, but the cool phase is equally likely to co-occur with El 399 400 Niño (thus rendering a weak correlation between IOBM and ENSO), then the net influence of the decoupled 401 Indian Ocean on ENSO would be expected to be minimal. Hence in the following analysis we show the 402 correlation coefficients between a range of the characteristic SST indices, alongside the observed and CMIP5 403 values for comparison (Fig. 9). As noted below, some aspects of the CMIP5 results in Fig. 9 clearly illustrate the need for partial coupling experiments in diagnosing inter-basin interactions, which is not 404 405 otherwise possible from the multi-model statistics.

407 4.1 El Niño-Southern Oscillation

408 The Niño-3.4 composite evolution for El Niño and La Niña (Fig. 6a,b; solid lines) reinforces the changes seen in the monthly standard deviation (Fig. 4a). Focussing firstly on the peak season of variability (i.e. 409 Sep(0)-Dec(0), yellow shaded region in Fig. 6a,b), the SSTAs are enhanced in both phases in DCPL₁₀ and 410 411 $DCPL_{AO+IO}$ relative to CTRL (Fig. 6a,b; red and orange solid lines compared to black), but not in $DCPL_{AO}$ (green solid line compared to black). After the peak ENSO season, the enhancement of DCPLIO relative to 412 413 CTRL (Fig. 6a,b) reflects the strongest shift in the monthly standard deviation, occurring during the boreal 414 winter and spring (Fig. 4a). Other changes are seen outside of the peak season in all partial coupling 415 experiments, most notably in DCPL $_{AO+IO}$, but these are consistent with lengthening of the periodicity (Fig. 416 5a).

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418 The enhanced SSTAs are consistent with the absence of the IOBM damping effect (Santoso et al. 2012). 419 The notably stronger Niño-3.4 anomalies in the months following the ENSO peak season (Fig. 6a,b; red solid line compared to black) coincide with IOBM peak occurrence (in a non-decoupled Indian Ocean). It is 420 421 known that the IOBM induces zonal wind stress (τ^x) anomalies over the western Pacific that oppose the westerly or easterly anomalies associated with eastern Pacific El Niño warming or La Niña cooling (e.g., 422 423 Kug and Kang 2006; Santoso et al. 2012). The τ^x signature of IOBM is apparent in the CTRL composites 424 between 100°E-160°E during Jan(1) to May(1) (Fig. 7a,e). In response to these western Pacific wind anomalies, oceanic Kelvin waves also act to promote ENSO phase turnabout (Wang et al. 1999a). 425 Consistently, in the absence of these negative feedback processes in DCPL₁₀ and DCPL_{A0+10}, ENSO 426 variability is enhanced and prolonged in those runs (see Section 3). Note that the western Pacific wind 427 428 anomalies are not solely a remote response to IOBM but are also part of ENSO evolution that is internal to 429 the equatorial Pacific (Watanabe and Jin 2002; Wang et al. 1999b). Thus, even when the IOBM is 430 completely absent in DCPLIO and DCPLAO+IO, these western Pacific wind anomalies associated with ENSO still prevail. These western Pacific τ^x anomalies are significantly weaker in the decoupled experiments than 431 in CTRL, and this is particularly so in DCPL_{IO} and DCPL_{AO+IO} in which the IOBM is completely absent 432 (Figs. 7b-d,f-h). The weakened τ^x occurs despite ENSO anomalies being enhanced in those decoupled 433 experiments (Fig. 4a), thus underscoring the impact of the missing IOBM in DCPL_{IO} and DCPL_{AO+IO}. 434

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The effect of the IOBM on western Pacific τ^x cannot be clearly inferred in BWI composites from DCPL_{PO} or DCPL_{PO+AO} (Fig. 8j,l,n,p) since its associated SSTAs, although present, become substantially weaker in the absence of ENSO (Fig. 6g,h). Nevertheless, the expected easterly τ^x anomalies are still visible near the Indo-Pacific boundary for positive IOBM (Fig. 8l), and westerly τ^x for negative IOBM (Fig. 8p), confirming the presence of a weak IOBM without ENSO.

441

442 The IOBM damping on ENSO is an intrinsic feature of the Indo-Pacific feedback interactions. This stems 443 from the co-occurrence of El Niño with warm IOBM and La Niña with cool IOBM, which is underscored by 444 the strong positive correlation between Niño-3.4 and BWI as seen in the observations, CMIP5 models and Mk3L (Fig. 9g,h). In this way the IOBM wind anomalies tend to consistently damp ENSO. The CMIP5 445 model spread reveals that there is a tendency for stronger positive correlation between ENSO and IOBM 446 447 with stronger ENSO and IOBM amplitude. Such tendency makes it challenging to diagnose the effect of 448 IOBM on ENSO using statistical analysis alone, hence reinforcing the need for partially coupled model experiments. 449

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451 The more severe weakening of western Pacific τ^x anomalies in DCPL_{AO+IO} than in DCPL_{IO} and the lesser 452 weakening in DCPL_{AO} relative to CTRL indicates the potential role of Atlantic variability in enhancing the western Pacific τ^x anomalies that have a damping effect on ENSO in CTRL. This result is consistent with the 453 454 amplification of Niño-3.4 variability in DCPL_{AO+10} outside the ENSO peak season (Fig. 4a; Fig. 6a,b). Albeit weak, the τ^x response to Atlantic Niño is visible in the composite plot of Atl-3 with the Indian Ocean 455 456 decoupled (Fig. 7j,l), with westerly anomalies in the equatorial Atlantic and easterly anomalies in the western Indian Ocean which appear to correspond with further anomalies in the western Pacific towards the 457 end of the calendar year (at Dec(0) between 120°E and 160°E). These anomalies are of the opposite sign for 458 Atlantic Niña (Fig. 7n,p). The western Pacific anomalies vary in strength, timing, and position across the 459 ensemble members (figure not shown), and hence appear weak in the ensemble mean. These τ^x signals are 460 not clear with air-sea interactions occurring in the Indian Ocean (Fig. 7k, o), presumably due to interference 461 with Indian Ocean internal variability. With the Atlantic decoupled, such wind anomalies are absent, thus 462 tending to enhance ENSO growth, as seen in DCPLAOHO. 463

465 Another way in which Atlantic variability might influence ENSO is via an alteration to Indian Ocean 466 variability. As seen earlier, the Atlantic appears to enhance the IOBM (Fig. 4d), and the IOBM has correspondingly been shown to damp ENSO. Therefore, if the IOBM becomes weaker in the absence of 467 468 Atlantic variability (DCPL_{AO}), the ENSO variability is expected to increase, although it should not be 469 stronger than when the IOBM is completely removed in DCPLAO+IO. The composites of Atl-3 indeed show 470 that in correspondence with Atlantic Niño, there is a warm IOBM response (Fig. 6c), and conversely for 471 Atlantic Niña (Fig. 6d). This Atlantic influence on ENSO via the IOBM may explain the enhanced Niño-3.4 anomalies in DCPL_{AO} after Jul(1) (Fig. 6a,b). Such Indian Ocean warming or cooling response can be 472 473 achieved through Atlantic forced wind stress anomalies in the western tropical Indian Ocean (Fig. 7j,l,n,p) over which ocean advection and entrainment are the dominant factors that generate interannual surface 474 temperature anomalies (Santoso et al. 2010). These easterly/westerly anomalies force a 475 downwelling/upwelling signal in the western Indian Ocean that then propagates eastward as a Kelvin wave 476 477 (figure not shown), thus promoting the occurrence of a basin-wide warming/cooling pattern. Decoupling the Atlantic alone does not remove the IOBM signal entirely as the IOBM warming/cooling is part of ENSO 478 evolution (Fig. 6a,b, green line with dots), and so any enhancement of ENSO amplitude in DCPLAO is 479 480 expectedly weaker than when the Indian Ocean is also decoupled (i.e., $DPCL_{AO+IO}$; Fig. 4a, 6a,b). ENSO 481 enhancement in DCPL_{IO} compared to DCPL_{AO+IO} is also limited because decoupling the Indian Ocean tends to enhance the AEM (Fig. 4b), which in turn has a damping effect on ENSO via its tendency to enhance 482 western Pacific τ^x anomalies. 483

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485 The AEM effects on ENSO described above would be maximal when an El Niño condition co-occurs with a warm AEM, and likewise for co-occurring cool events. However, such a combination has a weak tendency 486 of occurring in our model, and similarly across the CMIP5 models (Fig. 9a-d), unlike the robust ENSO-487 IOBM relationship (Fig. 9g,h). When we examine the correlation between the two indices with Atl-3 leading 488 by 12 months (figure not shown), we find that the two ensemble members showing a statistically significant 489 negative correlation exhibit a damped ENSO when the Atlantic is decoupled. For the remaining three 490 members, where the correlation is weak or positive, the ENSO is enhanced when the Atlantic is decoupled. 491 492 These inconsistent connections seem to be in line with recent studies that claim there is varying Atlantic-Pacific connection in observations due to decadal variability (Rodríguez-Fonseca et al. 2009; Martín-Rey et 493

494 al. 2014; Sasaki et al. 2014; McGregor et al. 2014).

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496 Another factor that can contribute to the lack of consistency is asymmetry between the warm and cool phases of the modes. The SSTA and wind stress composites for Atl-3 (Fig. 6c,d; 7i,m) show that there is an 497 498 asymmetry in CTRL, namely that Atlantic Niño tends to precede La Niña, but Atlantic Niña tends to follow a La Niña. Indian Ocean variability seems to be the source of this asymmetry, since both phases of the AEM 499 500 tends to follow ENSO (warm/cool AEM follows El Niño/La Niña) when the Indian Ocean is decoupled. The 501 correlation of Atl-3 averaged over June to August leading December to February Niño-3.4 is close to zero 502 (Fig. 9c,d). This however appears to disagree with observations, which shows a statistically significant negative correlation. The CMIP5 models slightly favour a negative correlation, but many models also display 503 a positive correlation (Kucharski et al. 2015), and the correlation does not appear to be related to the strength 504 505 in variability of either index across the models.

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The IOD generates only weak τ^x anomalies over the Pacific Ocean in DCPL_{PO+AO} (Fig. 8d,h). Santoso et al. (2012) noted that in this way the IOD is conducive for ENSO growth given the dominant IOBM damping effect. The significant positive correlation between Niño-3.4 and DMI at near zero lag is consistent with observations (Annamalai et al. 2005; Santoso et al. 2012), and the CMIP5 models (Fig. 9e,f). The intermodel relationships also show the tendency for stronger ENSO and IOD amplitude to reinforce this positive correlation. Again, analysis of partial coupling experiments suggests that care should be taken when inferring IOD impact on ENSO and vice versa.

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515 4.2 Atlantic Equatorial Mode

The model SST composites reveal that weak Atlantic Niño conditions coincide with the El Niño, and are followed by warm IOBM (Fig. 6a). Similarly, a weak Atlantic Niña coincides with La Niña, followed by cool IOBM (Fig. 6b). The Atlantic Niño is accompanied by westerly τ^x anomalies over the central Atlantic Ocean during April to July (Fig. 7i), and easterly τ^x anomalies with the Atlantic Niña (Fig. 7m). Wind stress anomalies opposing the Atlantic Niño are revealed in composites of El Niño, strongest at one-year lag (Fig. 7a), and also in composites of positive IOBM in CTRL (Fig. 8i). The same holds for Atlantic Niña, and composites of La Niña (Fig. 7e) and negative IOBM (Fig. 8m).

524 The influence of the IOBM is evidenced by the more significant weakening of the τ^x signal over the Atlantic Ocean in DCPL_{IO} (Fig. 7b,f) compared to any of the other experiments (Fig. 7c,d,g,h). Although weak, these 525 526 τ^{x} anomalies are still present even when the Indian Ocean is decoupled (DCPL_{IO}), but are more prominent 527 when the Atlantic is also decoupled (DCPL_{AO+IO}) – which is expected given the now absent anomalies 528 associated with the AEM. Note that these τ^x anomalies occur during the growth phase, as well as the decay 529 phase of ENSO in the model. As argued by Latif and Grötzner (2000), easterly wind anomalies over 530 equatorial Atlantic during El Niño force downwelling that lead to the formation of an Atlantic Niño six 531 months later. Similarly, westerly wind anomalies during La Niña promote an Atlantic Niña. Thus, here we are looking at both ENSO and IOBM processes that enhance and damp the AEM. Given the results in 532 533 Section 3 (Fig. 4b), the IOBM damping is the more dominant factor. Interestingly, this Indian Ocean 534 damping effect appears to be supported by the CMIP5 inter-model correlations: there is a negative 535 correlation between the BWI leading Atl-3 correlation and the Atl-3 standard deviation (Fig. 9o). This shows the tendency for models that simulate more occurrences of warm IOBM with Atlantic Niño to have weaker 536 AEM amplitude. This tendency is also supported by the relationship between the ENSO leading AEM 537 correlation with the Atl-3 standard deviation (Fig. 9b), given that El Niño induces warm IOBM. 538 539 Furthermore, since the IOBM is to a large extent a response to ENSO, it can also be inferred that ENSO damps the AEM indirectly via the IOBM. The effect of ENSO in enhancing the AEM may also be conveyed 540 via the IOD. A positive IOD is associated with easterly τ^x anomalies in the equatorial Atlantic during boreal 541 winter (Fig. 8d), and the opposite for negative IOD (Fig. 8h). However, this is outside the peak season of the 542

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545 4.3 Indian Ocean Modes

AEM, so the effect is expected to be weaker.

The positive IOD is associated with easterly τ^x anomalies across the Indian Ocean basin (Fig. 8a), and the negative IOD with westerly τ^x anomalies (Fig. 8e). Some of these τ^x anomalies can be induced by ENSO, as is clear in the decoupled Indian Ocean runs (Fig. 7b,d). In DCPL_{AO+IO} (Fig. 7d,h), the τ^x anomalies of the opposite sign near the Indo-Pacific boundary are stronger than in DCPL_{IO} (Fig. 7b,f). Nevertheless, ENSO drives τ^x anomalies over the western side of the Indian Ocean that promote the IOD (Annamalai et al. 2003; Fischer et al. 2005). The ENSO-IOD relationship is highlighted by a positive correlation in observations and

across models (Fig. 9e,f). The CMIP5 inter-model correlations also suggest a tendency for stronger ENSO-IOD correlation with stronger ENSO or stronger IOD. The AEM also drives τ^x anomalies that are favourable for the IOD (Fig. 7l,p), but the inter-model correlation is weak for Atl-3 leading DMI (Fig. 9i,j). Given the Atl-3 and Niño-3.4 correlation is also not strong, the AEM appears to be secondary in the influence on the IOD compared to ENSO.

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As mentioned above, the western Indian Ocean τ^x anomalies associated with the IOD, which are enhanced by ENSO and the AEM, can facilitate the formation of IOBM. Such dynamical association is highlighted by the strong positive correlation between the BWI and DMI (Fig. 9q,r), implying that with enhanced IOD events, stronger IOBM will ensue.

562

563 **5. Summary**

564 This study investigated the interactions between the dominant modes of climate variability across the tropics, namely ENSO, the AEM, the IOD and the IOBM. Using a series of coupled and partially coupled GCM 565 experiments we inferred the impact between modes on their strength, period, and seasonality. In agreement 566 with earlier studies, we found that Indian Ocean variability acts to damp ENSO via the IOBM (Santoso et al. 567 568 2012). Conversely the Pacific enhances both the IOD and IOBM (e.g., Behera et al. 2006). We have highlighted other findings that have not previously been explored in great depth, for instance, the connection 569 between tropical Indian Ocean and Atlantic variability. We found that Atlantic Ocean variability has little 570 influence on the IOD, but enhances the IOBM amplitude. Conversely, Indian Ocean variability has a net 571 572 damping effect on the AEM. As suggested by the wide range of conflicting literature, the connection between the Pacific and Atlantic Niños and Niñas is complex. Our study has shown that the coupling to the 573 tropical Indian Ocean is a factor that needs to be considered in inferring the Pacific-Atlantic interactive 574 feedbacks. After accounting for the effect of the Indian Ocean, our model experiments reveal that the AEM 575 576 has a net damping effect on ENSO magnitude in our model, whilst ENSO tends to enhance the AEM. 577

578 The dominant ENSO period in the Mk3L model is in the 5-7 year band. We found that decoupling either or 579 both of the Indian and Atlantic Ocean basins shifts the ENSO to longer periods, implying that variability in 580 each plays a role in the faster switching between ENSO phases. The 5-7 year signal is also dominant in the

power spectral densities of the other SST indices. The signal vanishes or is reddened in the absence of ENSO variability, thus demonstrating the coupling between ENSO and each of the tropical modes. Nevertheless, our results also show that each mode persists when variability in other basins, in turn and in combination, is removed. Furthermore, apart from slight changes to the IOBM, the overall seasonality of these modes is unchanged. This suggests that ENSO, the IOD, the AEM, and to a lesser extent the IOBM, are largely internally generated modes, despite the fact that the coupling between them influences their overall behaviour.

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589 Although the Atlantic Ocean appears to have a weak damping effect on ENSO amplitude, individual ensemble runs showed varying results. Earlier studies present conflicting reports on this matter. Dommenget 590 591 et al. (2006) showed that Atlantic variability damps ENSO, Frauen and Dommenget (2012) showed no net 592 influence, and Sasaki at al. (2014) showed that ENSO amplitude is reduced when the equatorial Atlantic is 593 decoupled. We see both ENSO damping and enhancement in different 100-year runs when the Atlantic alone is decoupled. It is likely that the interaction may be related to multi-decadal variations in the Atlantic-Pacific 594 connection (Rodríguez-Fonseca et al. 2009; McGregor et al. 2014; Sasaki et al. 2014; Kucharski et al. 2016). 595 596 The inconsistency in the Atlantic-Pacific Niño/Niña relationship is also exhibited in the large CMIP5 multi-597 model spread clustering around zero (Fig. 9c,d). This spread highlights the need for multiple ensemble experiments and, in light of model biases, the necessity to repeat experiments with several different models. 598 599

Since coupled models tend to suffer from the pervasive Indo-Pacific cold tongue bias and strong 600 601 climatological biases in the tropical Atlantic (Ritcher et al. 2014), alternative decoupling techniques may be necessary to further explore these connections. For example, one could nudge SST over a decoupled region 602 toward observed SST, as conducted by Terray et al. (2015). In this way, biases would be eliminated, 603 however the associated analyses would be confounded by additional mean-state changes, which introduce 604 further complexity. Adding to the complexity is that climatological biases also translate to biases in the 605 modes of variability. One clear example is the shallow thermocline bias in the eastern Indian Ocean that 606 tends to make simulated IOD events notably stronger than observed. However, an error-compensating effect 607 608 may also occur. Namely, as air-sea coupling tends to be more active in the strong convective region of the Indo-Pacific warm pool, the cold bias may underestimate the remote effect of the IOD, but this should be to a 609

610 certain extent compensated by the overly large IOD amplitude (Santoso et al. 2012). Thus, while the result 611 may not be greatly affected by this particular bias due to such an error-compensating tendency, a multimodel approach with less biased models seems to be the way forward. The Mk3L model used here exhibits 612 reasonable skill in simulating the tropical modes of climate variability, especially relative to the CMIP5 613 614 models that are of higher resolution (Fig. 2). The ensemble spread of index correlations also lies within the CMIP5 multi-model spread in each case (Fig. 9). Therefore, combined with its coarser resolution, the model 615 616 we employed is particularly useful for studying these types of problems over centennial and millennial 617 timescales. However, it should be further noted that in our case flux adjustments have been employed in 618 order to maintain a more realistic climatology.

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620 Our study highlights the coupling across tropical modes of variability, linked by the atmospheric Walker 621 Circulation. This carries an important implication in that understanding, predicting, and projecting each 622 mode of variability would require a careful consideration of other remote modes of variability. Such coupling also implies that diagnosing interactive feedback, relying on statistical inferences alone, is 623 challenging. We illustrated this challenge by utilising an analysis of CMIP5 models (Fig. 9). For instance, 624 there is a tendency for stronger ENSO and IOBM amplitude to be associated with higher ENSO-IOBM 625 626 coherence across the models (Fig. 9g,h). At best, this relationship would suggest the IOBM is a mere slave to ENSO. However, with the aid of decoupling experiments in this study and others, the IOBM has been 627 shown to have a damping effect on ENSO. This study provides a basis for understanding the interactions 628 between the dominant modes of variability in the tropics. We note that weaker modes in the tropical domain, 629 630 for example, the Madden-Julian Oscillation or the Atlantic Meridional Mode, may also influence the interactions that impact on strengths, periods, or predictability of these modes. Furthermore, modes outside 631 of the tropical domain, such as the Southern Annular Mode or the North Atlantic Oscillation, may also have 632 an influence. How these other modes impact on the tropical interactions is a complex topic of investigation 633 634 that should be explored in future studies.

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1	Tropical Climate Variability: Interactions across the Pacific, Indian, and Atlantic Oceans
2	Jules B. Kajtar ^{*1,2} , Agus Santoso ^{1,2} , Matthew H. England ^{1,2} , and Wenju Cai ³
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4	¹ ARC Centre of Excellence for Climate System Science, University of New South Wales, NSW, Australia
5	² Climate Change Research Centre, University of New South Wales, NSW, Australia
6	³ CSIRO Marine and Atmospheric Research, Aspendale, Victoria, Australia
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8	* Corresponding author: Jules B. Kajtar, j.kajtar@unsw.edu.au
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10	Tables and Figures
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29 TABLE 1. Summary of the tropical modes of variability considered in this study.

Mode	Ocean	Characteristic SST Index	SST Averaging Area	Peak Season
				September to
El Niño-Southern	D : C		50G 500 1 1 7 00 1 0 000 1	
Oscillation (ENSO)	Pacific	Niño-3.4	5°S-5°N, 170°-120°W	December
Oscillation (E1350)				(SOND)
				(50112)
Atlantic Equatorial		Atl-3	5°S-5°N, 20°W-0°	June to August
	Atlantic			
Mode (AEM)				(JJA)
			west (10°S-10°N, 50°-	August to
Indian Ocean Dipole		Dipole Mode Index		C
	Indian	/	70°E) minus east (10°S-	November
(IOD)		(DMI)	00,000,1100E	
			0,90-110E)	(ASON)
Indian Ocean				
				January to April
Basinwide Mode	Indian	Basinwide Index (BWI)	20°S-20°N, 40°-100°E	
(IOBM)				(JFMA)
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Fig. 1 Comparison of the Mk3L model SST, winds, and rainfall with observed data. (a) Model SST, (b) HadISST (averaged over 1900-2010), and (c) difference between model and observed SST. (d) Model zonal wind stress (τ^x), (e) NCEP-NCAR reanalysis τ^x (averaged over 1948-2010), and (f) difference between model and observed τ^x . (g) Model rainfall, (h) Climate Prediction Center Merged Analysis of Precipitation (CMAP) reanalysis rainfall (averaged over 1979-2010), and (i) difference between model and observed rainfall. The model climatology is averaged over the entire 300 years of the original CTRL run.

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Fig. 2 Monthly standard deviation and power spectral densities (PSDs) of the SST indices characterizing the dominant modes of tropical climate variability. (a,b) Niño-3.4, which characterizes the El Niño-Southern Oscillation. (c,d) Atl-3, for the Atlantic Equatorial Mode. (e,f) DMI, for the Indian Ocean Dipole. (g,h) BWI, for the Indian Ocean Basinwide Mode. The faint red curves denote the five ensemble members of the CTRL experiment in Mk3L, with the bold red curve denoting the ensemble mean. The faint blue curves denote individual historical CMIP5 model runs (Section 2.2), with the bold blue curve denoting the CMIP5 mean. The green curve denotes the observed from the HadISST set. For easier comparison, the time-series are normalised to unity variance before computing the PSDs. Note that the individual CMIP5 models are not distinguished, since they are presented purely to indicate the multi-model spread.



Fig. 3 First empirical orthogonal function (EOF) of June to August Atlantic SST presented as a regression map, showing the pattern for the Atlantic equatorial mode for (a) the model and (b) HadISST. The variances explained by the first EOFs are given in the panel titles. (c,d) Correlation coefficients of Atl-3 index with SST field, using June to August mean in each case. The EOF analysis is for the entire 300 years of the original model CTRL run, and over 1900 to 2010 for HadISST.





Fig. 4 Monthly standard deviation of SST indices representing the dominant modes of tropical climate variability: (a) Niño-3.4 for the Pacific Ocean, (b) Atl-3 for the Atlantic Ocean, (c) Dipole Mode Index, and (d) Basinwide Index for the Indian Ocean. The results for the CTRL and the various partial coupling experiments are shown. A low-pass filter with a 4-month cut-off was applied to smooth each time-series before computing the monthly standard deviations. The color-shaded areas indicate the 95% confidence intervals, which were estimated based on 1000 bootstrapped means from the five 100-year ensemble members. The borders of the shaded regions are outlined for clarity. The months shaded in yellow indicate the peak season of variability of each mode, upon which further analysis in the text is focussed.



105 Fig. 5 Power spectral densities of the same SST indices shown in Fig. 4.



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Fig. 6 Composite evolution of tropical climate mode events for each SST index over a 36-month period. Events are composited as described in Section 4. The different lines styles denote the different indices throughout the figure: solid lines for Niño-3.4, lines with squares for Atl-3, lines with crosses for DMI, and lines with circles for BWI. The colours denote the different experiments. (a) Composite evolution of the Niño-3.4 index for El Niño events (thick, solid lines), in CTRL (black) and each of the partially coupled experiments (DCPL_{IO}: red, DCPL_{AO}: green, DCPL_{AO+IO}: orange). Alongside is the co-evolution of each of

126	the other indices (thin lines, and again, line styles and colors indicate the different indices and experiments).
127	(b) Composite evolution of the Niño-3.4 index for La Niña events. (c,d) Composite evolution of Atl-3 for
128	warm and cool AEM events, respectively. (e,f) Composite evolution of DMI for positive and negative IOD
129	events, respectively. (g,h) Composite evolution of BWI for warm and cool IOBM events, respectively. Apart
130	from the evolutions of main index in each panel (thick lines), which are plotted for the entire 36-month span,
131	only those periods for which the other indices are significantly different from zero at the 95% confidence
132	level under a <i>t</i> -test are shown (thin lines). Jul(0) denotes July in the year of the event, and -1 or 1 denotes the
133	year prior or ahead. Year 0 is relative to the peak of each mode. The months shaded in yellow indicate the
134	peak seasons of variability for the main index in each panel (Table 1).
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Fig. 7 Composites of equatorial zonal surface wind stress anomalies over a 36-month period associated with ENSO and AEM events. (a) El Niño composites in CTRL. (b-d) El Niño composites in each of the relevant decoupling experiments. (e-h) La Niña composites. (i-l) Atlantic Niño composites. (m-p) Atlantic Niña composites. Events are composited as described in Section 4. Only the wind stresses that are significantly different from zero at the 95% confidence level under a t-test are shaded. The zonal wind stress is averaged over 5°S-5°N. The vertical dashed lines indicate the approximate boundaries between the ocean basins. The box in each panel indicates the spatial and temporal extent of relevant composited index. On the DCPL panels, the pink contours indicate where |DCPL| - |CTRL| is positive and significant above the 90% level. The green contours indicate where |DCPL| - |CTRL| is negative and significant above the 90% level.



Fig. 8 As in Fig. 7 but for IOD and IOBM events. (a) Positive IOD composites in CTRL. (b-d) Positive IOD
composites in each of the relevant decoupling experiments. (e-h) Negative IOD composites. (i-l) Warm

¹⁷² IOBM composites. (m-p) Cool IOBM composites.



Fig. 9 Correlation coefficients between pairs of indices plotted against the standard deviation of each in the 185 pair. Annual averages are taken for each index over the following months: DJF for Niño-3.4, JJA for Atl-3, 186 SON for DMI, and JFM for BWI. The months here were chosen based on the observed peak season of 187 188 variability, rather than the peak seasons of the modes in the model. The blue dots denote individual CMIP5 historical runs (over the period 1900-1999) with blue lines of best-fit, the red dots denote the 100-year Mk3L 189 CTRL ensemble members, and the green crosses denote the HadISST observations (1900-1999). The CMIP5 190 inter-model correlation coefficient is given in each panel, where a black value indicates that the correlation is 191 192 significant at the 95% level, and grey is not significant.