

# The Isotope Geochemistry of Ni

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## 1. INTRODUCTION

Nickel is an iron-peak element with 5 stable isotopes (see Table 1) which is both cosmochemically abundant and rich in the information carried in its isotopic signature. Significantly,  $^{60}\text{Ni}$  is the radiogenic daughter of  $^{60}\text{Fe}$ , a short-lived nuclide ( $t_{1/2} = 2.62\text{Ma}$ ; Rugel et al., 2009) of a major element.  $^{60}\text{Fe}$  has the potential to be both an important heat source and chronometer in the early solar system.  $^{60}\text{Ni}$  abundances serve to document the prior importance of  $^{60}\text{Fe}$  and this is a topic of on-going debate (section 3). The four other stable Ni nuclides span a sizeable relative mass range of  $\sim 10\%$ , including the notably neutron-rich nuclide  $^{64}\text{Ni}$ . The relative abundances of these isotopes vary with diverse stellar formation environments and provide a valuable record of the nucleosynthetic heritage of Ni in the solar system (section 2). Ni occurs widely as both elemental and divalent cationic species, substituting for Fe and Mg in common silicate structures and forming Fe/Ni metal alloys. The Ni isotope chemistry of all the major planetary reservoirs and fractionations between them can thus be characterised (section 4). Ni is also a bio-essential element and its fractionation during low-temperature biogeochemical cycling is a topic that has attracted recent attention (section 4).

### 1.1 Notation

Much of the work into Ni has been cosmochemical, focussing on the nucleosynthetic origins of different meteoritic components. Such studies have primarily investigated mass-independent isotopic variations, both radiogenic and non-radiogenic, which require choosing a reference isotope pair for normalisation. Throughout this work we use  $^{58}\text{Ni}$ - $^{61}\text{Ni}$  as the normalising pair, in keeping with current practice in the field. An alternative  $^{58}\text{Ni}$ - $^{62}\text{Ni}$  normalisation scheme has previously been used for bulk analyses (Shimamura and Lugmair 1983; Shukolyukov and Lugmair, 1993a,b; Cook et al., 2007; Cook et al., 2008; Quitté et al., 2006; Quitté et al., 2011; Chen et al., 2009) and one early study used  $^{58}\text{Ni}$ - $^{60}\text{Ni}$  (Morand and Allègre, 1983). Although the large isotopic variability accessible by *in situ* analyses often makes external normalisation a viable option for mass-independent measurements by secondary ionisation mass-spectrometry (SIMS), some have employed internal normalisation in determinations of  $^{60}\text{Ni}/^{61}\text{Ni}$  (Tachibana and Huss, 2003; Tachibana et al., 2006; Mishra and Chaussidon, 2014; Mishra et al., 2016). Given unresolvable Fe and Zn interferences on masses 58 and 64, this requires normalising to  $^{62}\text{Ni}/^{61}\text{Ni}$ . In this review, all data have been renormalised to  $^{58}\text{Ni}$ - $^{61}\text{Ni}$ , where possible. Some studies have only reported normalised data and so such conversion cannot be made. Fortunately, the subtle differences resulting from different normalisations do not affect the inferences being made in these cases and we simply indicate the normalisation scheme used. To be clear about these potentially important details, we use a notation proposed by Steele et al. (2011), which includes this information. For example:

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$$\epsilon^{60}\text{Ni}_{58/61} = \left( \frac{{}^{60}\text{Ni}/{}^{58}\text{Ni}^{\text{sample}}}{\text{norm } 58/61} / \frac{{}^{60}\text{Ni}/{}^{58}\text{Ni}^{\text{standard}}}{\text{norm } 58/61} - 1 \right) \times 10000 \quad (\text{Eq 1})$$

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or the parts per ten thousand variation of  ${}^{60}\text{Ni}/{}^{58}\text{Ni}$  (internally normalised to a reference  ${}^{58}\text{Ni}/{}^{61}\text{Ni}$ ) relative to a standard measured in the same way. The established isotopic standard for Ni is the National Institute of Standards and Technology Standard Reference Material (NIST SRM) 986, Gramlich et al. (1989). Reference Ni isotope ratios for this standard are reported in Table 1. This NIST SRM has been widely used, providing a valuable common datum in all but the earliest work. If it is necessary to clarify which reference standard has been used, the notation above can be augmented, e.g.  $\epsilon^{60}\text{Ni}_{58/61}$  (NIST 986). For elements such as Ni, however, where the same standard is conventionally used, we feel this additional information can be omitted without too much confusion, provided it is imparted elsewhere (as we do here). We use the epsilon notation solely for mass-independent isotopic data (internally normalised). This approach is typical although not universal and the presence of the subscript in our notation (Eq 1) makes the use of internal normalisation evident.

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We report mass-dependent variations in the delta notation:

$$\delta^{60/58}\text{Ni} = \left( \frac{{}^{60}\text{Ni}/{}^{58}\text{Ni}^{\text{sample}}}{\text{norm } 58/61} / \frac{{}^{60}\text{Ni}/{}^{58}\text{Ni}^{\text{standard}}}{\text{norm } 58/61} - 1 \right) \times 1000 \quad (\text{Eq 2})$$

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As for the mass-independent work, NIST SRM 986 is extensively used as the Ni isotope reference standard and is implicit in Eq (2). In Eq (2), we follow another proposal made in Steele et al. (2011) to report the isotope ratio used, i.e.  $\delta^{60/58}\text{Ni}$  instead of  $\delta^{60}\text{Ni}$ . This removes any ambiguity over which nuclide is used as the denominator. For an element such as Ni, with more than two stable isotopes, such qualification is valuable. We suggest this systematic notion could be useful more generally.

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Moynier et al. (2007) reported their mass-dependent Ni isotope data as  $\bar{\delta}\text{Ni}$ , an error weighted, average fractionation per unit mass difference, using the three measured ratios  $({}^{60}\text{Ni}/{}^{58}\text{Ni})/2$ ,  $({}^{61}\text{Ni}/{}^{58}\text{Ni})/3$  and  $({}^{62}\text{Ni}/{}^{58}\text{Ni})/4$ . This is an interesting idea (see also Albalat et al., 2012), which reduces the error for the sample-standard bracketing technique by using all measured data. The alternative method for determining mass-dependent isotopic fractionation is by double-spiking (see review by Rudge et al., 2009). Double-spiking requires measurements of four isotopes, thus yielding three independent isotope ratio determinations. Given a troublesome Zn interference on mass 64, all double-spiked studies to date have used a  ${}^{61}\text{Ni}$ - ${}^{62}\text{Ni}$  double spike and employed  ${}^{58}\text{Ni}$ ,  ${}^{60}\text{Ni}$ ,  ${}^{61}\text{Ni}$  and  ${}^{62}\text{Ni}$  in the data reduction. With ostensible similarity to the sample-standard bracketing approach of Moynier et al. (2007), the combined measurements of these four isotopes yield a single value of natural isotopic fractionation, which is normally converted into a more tangible delta value for a specific but arbitrary isotope ratio (e.g.  $\delta^{60/58}\text{Ni}$ ). The key difference in double-spiking is that the additional isotope ratios are used to constrain explicitly instrumental mass-bias. Namely this procedure improves accuracy, whereas in sample-standard bracketing instrumental mass bias is assumed to be identical for sample and standard and the additional isotope measurements are used to improve precision.

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Both methods described above use measurements of  ${}^{60}\text{Ni}$  for determining mass-dependent Ni isotope variability. It is germane to consider whether or not it makes good sense to use a radiogenic isotope for such a purpose. For terrestrial samples, there should be no variability in the relative abundance of radiogenic  ${}^{60}\text{Ni}$ , given likely terrestrial isotopic homogenisation after parental  ${}^{60}\text{Fe}$  became extinct. For extra-terrestrial samples, this is potentially a consideration, but bulk variations in  $\epsilon^{60}\text{Ni}_{58/61}$  are typically small ( $\sim 0.1$ ), dominantly nucleosynthetic rather than radiogenic and associated with mass-independent variability of other isotopes (see section 2). For the most accurate mass-dependent measurements, a second mass-independent isotopic determination is therefore required (e.g. Steele et al., 2012), but such mass-independent

variability does not have a significant impact on more typical mass-dependent determinations at the delta unit level (section 4.1).

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In section 3 we address the presence of  $^{60}\text{Fe}$  in the early solar system from  $^{60}\text{Ni}$  measurements of meteoritic samples. Such determinations yield initial  $^{60}\text{Fe}/^{56}\text{Fe}$  for the objects analysed, denoted  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$ . Given samples that yield precise values of  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  may have different ages, it is useful to calculate  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  at a common reference time, typically the start of the solar system as marked by calcium aluminium rich inclusion formation. Such a solar system initial value is abbreviated to  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$ . Throughout this review, the uncertainties quoted for various average measurements are two standard errors, unless otherwise stated.

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## 2. NUCELOSYNTHETIC Ni ISOTOPIC VARIATIONS

Nickel is significant element in stellar nucleosynthesis. Nickel-62 has the highest binding energy per nucleon of any nuclide; no nuclear reaction involving heavier nuclides can produce more energy than it consumes. It is often, incorrectly, said that  $^{56}\text{Fe}$  has the highest binding energy per nucleon, likely due to the anomalously high abundance of  $^{56}\text{Fe}$ . In fact,  $^{56}\text{Fe}$  is dominantly produced in stars as the decay product  $^{56}\text{Ni}$ , which is the result of the last energetically favourable reaction during Si burning.

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It is thought the Ni isotopes are dominantly produced during nuclear statistical equilibrium (NSE or the e-process) in supernovae (Burbidge, Burbidge Fowler and Hoyle 1957). There are two main astrophysical environments in which the majority of Ni is thought to be produced, these are the type Ia (SN Ia) and type II (SN II) supernovae. SN Ia are thought to be the violent explosions of carbon-oxygen white dwarves which accrete material from a binary host to reach the Chandrasekher limit ( $<1.39 M_\odot$  non-rotating). SN Ia are highly neutron enriched environments and have been hypothesized to be the source of some important neutron-rich nuclides, including  $^{48}\text{Ca}$ ,  $^{60}\text{Fe}$ ,  $^{62}\text{Ni}$  and  $^{64}\text{Ni}$ . Due to the size of the progenitor ( $<1.39 M_\odot$ ) the stars are old,  $\sim 1$  Ga, as they have burned their fuel slowly. This means they out-live the stellar nurseries (lifetime  $\sim 10$  Ma) in which they formed. Therefore, they would make an unlikely, and so very interesting, source for isotope anomalies in the Solar System. SNII are the terminal explosions of much larger stars ( $>12 M_\odot$ ) which consequently have much shorter main sequence lifetimes. Their shorter lifetimes make them much more likely to return material to star forming regions, meaning they are a more probable source for nucleosynthetic anomalies in the Solar System.

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Trying to identify nucleosynthetic contributions from specific stellar events, such as those described above, within the average composition of the solar system is very difficult. However, the presence of 'isotopic anomalies' within meteoritic material provides opportunities to try to fingerprint individual sources. Isotopic anomalies are mass-independent isotopic compositions that differ from anthropocentric terrestrial values, and can indicate a different blend nucleosynthetic products to those on Earth. Such heterogeneities reveal either imperfect mixing of contrasting stellar inputs to the solar system (e.g. Reynolds, 1960), or else the unmixing of components in a generally well mixed, but locally heterogeneous nebula (e.g. Trinquier et al., 2009). These observations are of great interest to understanding the processes occurring in the early solar nebula and the stellar contributors to our solar system.

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Calcium aluminium rich inclusions (CAIs), found in some primitive meteorites, document isotopic anomalies in the elements hosted in the refractory minerals from which they are formed. This was initially evident in the mass-independent oxygen isotopic compositions (Clayton et al., 1973) of the abundant CAIs in the meteorite Allende, although this signature is now largely attributed to gas-phase processes in the nebula (e.g. Yurimoto and Kuramoto, 2004). However, isotopic anomalies were also discovered for a range of refractory metals (e.g. McCulloch and Wasserburg, 1978 a, b) which are still believed to document, at the macroscopic scale ( $>1\text{mm}$ ),

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nucleosynthetic mixtures that contrast with the bulk solar system. Following reports of mass-independent variations of other iron-peak nuclides (e.g. Lee et al., 1978; Heydegger et al., 1979), the first Ni isotopic analyses of Allende CAIs (Morand and Allègre, 1983; Shimamura and Lugmair, 1983) failed to resolve signatures that differed from terrestrial values, in all but a single, highly anomalous 'FUN' inclusion (Shimamura and Lugmair, 1983). Subsequent improvements in precision allowed Birck and Lugmair (1988) to resolve excesses of  $\sim 1 \epsilon^{62}\text{Ni}_{61/58}$  and  $\sim 3 \epsilon^{64}\text{Ni}_{61/58}$  within Allende CAIs (Fig. 1), which they noted was in keeping with a neutron-rich, equilibrium process nucleosynthesis. These findings were pleasingly compatible with anomalies in neutron-rich isotopes of Ca (Jungck et al., 1984) and Ti (Heydegger et al., 1979; Niederer et al., 1980; Niemeyer and Lugmair, 1981) from previous studies and excesses of  $^{54}\text{Cr}$  in their own work (Birck and Lugmair, 1988). This landmark contribution identified the key mass-independent variations in Ni isotopes that would subsequently become apparent in bulk meteorite analyses.

The CAIs analysed by Birck and Lugmair (1988) also displayed  $^{60}\text{Ni}$  enrichments,  $\epsilon^{60}\text{Ni}_{61/58} \sim 1$  (Fig. 1a), potentially related to the decay of  $^{60}\text{Fe}$  co-produced with the neutron-rich Ni isotopes (see section 3). If these  $\epsilon^{60}\text{Ni}_{61/58}$  values are taken solely as the consequence of *in situ*  $^{60}\text{Fe}$  decay, they imply initial  $^{60}\text{Fe}/^{56}\text{Fe} \sim 1 \times 10^{-6}$ , but Birck and Lugmair (1988) cautioned against such an inference, given associated nucleosynthetic variations of comparable magnitude. Further analyses of CAIs by Quitté et al. (2007) similarly showed positive  $\epsilon^{60}\text{Ni}_{61/58}$  and  $\epsilon^{62}\text{Ni}_{61/58}$  (see Fig 1a); the method used in this study suffered from too large  $^{64}\text{Zn}$  interferences to make precise  $\epsilon^{64}\text{Ni}_{61/58}$  measurements. The authors attributed their observations to synthesis of  $^{60}\text{Fe}$  and  $^{62}\text{Ni}$  (and  $^{96}\text{Zr}$ ) in a neutron burst event, followed by the decay of  $^{60}\text{Fe}$ . Birck and Lugmair (1988) had argued against this style of model, given the absence of predicted, associated  $^{46}\text{Ca}$  anomalies in CAIs. However, Quitté et al. (2007) tentatively inferred  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}} > 1 \times 10^{-6}$  from their nucleosynthetic model, a two point internal CAI isochron and the interpretation of  $\epsilon^{60}\text{Ni}_{61/58}$  excesses in two CAIs without  $\epsilon^{62}\text{Ni}_{61/58}$  anomalies.

Better precision was required to investigate mass-independent Ni isotopic variability between bulk meteorite samples. This came with improvements in mass-spectrometry, including multi-collection systems and their coupling with plasma sources (MC-ICPMS). The latter allows intense beams to be runs with relative ease, permitting counting statistical limitations to be overcome given sufficient sample availability. Nonetheless, the technique does require careful monitoring of a wide range of potential sample and plasma related interferences that may be significant at high precision (see common interferences listed in Steele et al., 2011 and Gall et al., 2012). Some of the earlier MC-ICPMS studies focussed on iron meteorites or the metallic phases of iron-bearing chondrites (Cook et al., 2006; Quitté et al., 2006; Moynier et al., 2007; Dauphas et al., 2008). This approach usefully exploits the natural concentration of Ni in a form readily purified to provide sufficient material for high precision analysis. In comparison, work on silicate samples requires more involved separation of Ni from a wider range of elements, typically accomplished using the highly Ni-specific complexing agent dimethylglyoxime either in solvent extraction (e.g. Morand and Allegre, 1983; Shimamura and Lugmair, 1983), mobile phase (e.g. Wahlgreen et al. 1970; Victor, 1986; Steele et al., 2011; Gall et al., 2012, Chernozhkin et al, 2015) or stationary phase of ion-chromatography (e.g. Quitté and Oberli, 2006; Cameron and Vance, 2011). Alternatively, cation chromatography using a mixed HCl-acetone eluent (Strelow et al 1971) has also been successfully used (e.g. Tang and Dauphas, 2012).

Initial MC-ICPMS studies on meteoritic metal phases (Cook et al., 2006; Quitté et al., 2006; Moynier et al., 2007; Dauphas et al., 2008) dominantly argued against Ni isotope anomalies in bulk meteorites (Figs. 1-3), as did Chen et al. (2009) for their multi-collector thermal ionisation mass-spectrometry (MC-TIMS) analyses (Fig. 3). As exceptions to these overall observations of bulk Ni isotope homogeneity, Quitté et al. (2006) reported correlated, negative values of  $\epsilon^{60}\text{Ni}_{61/58}$  and  $\epsilon^{62}\text{Ni}_{61/58}$  in many of the sulfide inclusions they analysed from iron meteorites (Fig. 2a). Subsequently, Cook et al. (2008) reported more modest anomalies in troilites from iron

meteorites (Fig. 2). Both studies argued for the preservation of a pre-solar component in these sulfide inclusions, although how this occurred mechanistically was problematic. The MC-TIMS work of Chen et al. (2009) provided a different measurement perspective. This study argued against resolvable differences in  $\epsilon^{60}\text{Ni}_{61/58}$  and  $\epsilon^{62}\text{Ni}_{61/58}$  in either bulk or sulfide samples, at a level of  $\pm 0.2\epsilon$  and  $\pm 0.5\epsilon$  respectively (Fig. 2a). Since then, no one has further pleaded for the case of anomalous sulfides and the original analyses seem likely to have been measurement artefacts. However, there has been on-going debate about the presence of Ni isotopic anomalies in bulk meteorites.

In striking contrast to the bulk meteorite analyses described above, Bizzarro et al. (2007) reported a dataset with near constant negative  $\epsilon^{60}\text{Ni}_{61/58}$  (and  $\epsilon^{62}\text{Ni}_{61/58}$ ) in differentiated meteorites but  $\epsilon^{60}\text{Ni}_{61/58} \sim 0$  and positive  $\epsilon^{62}\text{Ni}_{61/58}$  in chondrites. These data were used to invoke a late super-nova injection of  $^{60}\text{Fe}$  into the solar system. Subsequent studies were unable to reproduce these results (Dauphas et al., 2008; Regelous et al., 2008; Chen et al., 2009; Steele et al., 2011) and noted that the systematics of the Bizzarro et al. (2007) dataset were consistent with an interference on  $^{61}\text{Ni}$ . In reporting the results of further analyses, Bizzarro et al. (2010) commented that their new data were inconsistent with Bizzarro et al. (2007) but agreed with the observations of Regelous et al. (2008). The data from Bizzarro et al. (2007) will thus not be further considered.

Regelous et al. (2008) presented bulk analyses of  $\epsilon^{60}\text{Ni}_{61/58}$  and  $\epsilon^{62}\text{Ni}_{61/58}$  on a suite of chondrites and iron meteorites with precisions of around  $\pm 0.02 \epsilon$  and  $\pm 0.04 \epsilon$  respectively. By making higher precision measurements, in part by pooling multiple repeats of the same sample and by examining a wider range of meteorites than earlier studies, Regelous et al. (2008) were able to resolve differences in bulk meteorite compositions (Fig. 3a). They illustrated that variability between different chondrite groups is largely echoed by that in iron meteorites (Fig. 3a). Notably the IVB irons have Ni isotopic compositions similar to carbonaceous chondrites (positive  $\epsilon^{62}\text{Ni}_{61/58}$ ), whilst the other magmatic irons resemble ordinary chondrites (with negative  $\epsilon^{62}\text{Ni}_{61/58}$ ). As for a number of other isotopic systems, enstatite chondrites were largely within error of terrestrial values.

These observations were further refined at higher precision (Fig. 4a) and with the inclusion of  $\epsilon^{64}\text{Ni}_{61/58}$  data (Fig. 4b) by Steele et al. (2011, 2012) and Tang and Dauphas (2012, 2014). These data revealed a continuous, well defined array in  $\epsilon^{62}\text{Ni}_{61/58}$  vs  $\epsilon^{64}\text{Ni}_{61/58}$  from ordinary chondrites and most magmatic irons, through terrestrial values in EH chondrites to carbonaceous chondrites and IVB irons (Fig. 4b). This ordering of meteorite groups is the same as observed in the mass-independent isotopic compositions of other first row, transition elements (Trinquier et al., 2007, 2009). Although the total isotopic variability in Ni is smaller than for Ti or Cr, its notable strength is that both chondrites and iron meteorites can be analysed to high precision, allowing genetic associations to be made from bulk compositions of iron meteorites rather than from occasional oxygen-bearing inclusions they contain (e.g. Clayton et al., 1983). Strikingly, the bulk Ni isotope array in Figure 4b points towards the CAI values presented by Birck and Lugmair (2008), see Figure 1b. For the same arguments as made by Trinquier et al. (2009), however, the bulk meteorite array is not formed by simple mixing between a single bulk composition and CAI, which would create a strongly curved array (see Figure 5). Carbonaceous chondrites are enriched, relative to ordinary chondrites, in the same isotopic component that is manifest more strongly in CAIs. Yet, it is not addition of CAIs themselves that causes the trend in Figure 4b but presumably variable abundances of specific pre-solar grains in chondrite matrices (see Dauphas et al., 2010 and Qin et al., 2011 for the Cr isotope case).

The well-defined array in non-radiogenic isotopes (Fig. 4b) provides key constraints on the nucleosynthetic origins of this important nebular component. At face value it represents coupled enrichments in the neutron rich isotopes of  $\epsilon^{62}\text{Ni}_{61/58}$  and  $\epsilon^{64}\text{Ni}_{61/58}$ . However, the 3:1 slope of the array (Fig. 4b) can also be reproduced by variable meteoritic values of  $^{58}\text{Ni}/^{61}\text{Ni}$ , the

normalising isotope ratio. Indeed, from high-precision mass-dependent isotopic measurements (see section 4), Steele et al. (2012) showed that variability in  $^{58}\text{Ni}$  best explains all observations. This implies the source of this anomalous Ni is from the Si-S zone of a SNII. This contrasts with material from the O-Ne zone required to account for a similarly constrained component in the Ti isotopic system. Steele et al. (2012) suggested ways in which the different, contributing zones for these different elements might be reconciled by grains from different zones being sorted (homogenised or unmixed) by solar system processes. However, these issues remain unresolved.

The correlation of bulk analyses of  $\epsilon^{60}\text{Ni}_{61/58}$  with the other isotope ratios ( $\epsilon^{62}\text{Ni}_{61/58}$  or  $\epsilon^{64}\text{Ni}_{61/58}$ ) is much less systematic (Regelous et al., 2008; Steele et al., 2012; Tang and Dauphas, 2014), see Figure 4a. In general, carbonaceous chondrites have lower  $\epsilon^{60}\text{Ni}_{61/58}$  than ordinary chondrites and enstatite chondrites, but CI chondrites are a notable exception. It is tempting to attribute these different relationships to the radiogenic nature of  $\epsilon^{60}\text{Ni}_{61/58}$ , but given the abundance of  $^{60}\text{Fe}$  inferred from bulk meteorite studies (see section 3 below) this seems unlikely. Instead, this presumably reflects part of a more complex nucleosynthetic signature, with several components, potentially one which incorporates high fossil  $^{60}\text{Fe}$ , involved in generating the Ni isotopic compositions of ordinary and carbonaceous chondrites.

The carriers of the exotic Ni isotopic components that shape variable bulk meteorite compositions remain to be identified. The highly anomalous isotopic compositions found in separated SiC grains from the CM2 meteorite, Murchison (Marhas et al., 2008) do not readily account for the mass-independent Ni isotope variations seen in bulk samples. These ion-probe analyses show  $^{61}\text{Ni}/^{58}\text{Ni}$  and  $^{62}\text{Ni}/^{58}\text{Ni}$  ratios  $>1000\epsilon$  higher than terrestrial values (for these extreme ratios there is no internal normalisation) in X-grains believed to be derived from SNII sources. These Ni isotopic signatures suggest derivation from outer He/N and He/C zones of an SNII event, different to those inferred by Steele et al. (2012) to be necessary to account for bulk isotopic variability. As with other analyses of pre-solar SiC, these measurements bear striking testimony to the diversity of stellar sources that contribute to the bulk composition of the solar system, but do not identify a carrier for the signature that causes variability between different bulk, meteoritic objects (Steele and Boehnke, 2016).

### 3. EXTINCT $^{60}\text{Fe}$ AND RADIOGENIC $^{60}\text{Ni}$

As alluded to in the preceding section, much of the initial interest in Ni isotope cosmochemistry was focussed on trying to identify the presence of live  $^{60}\text{Fe}$  in the early solar system. A salient property of  $^{60}\text{Fe}$  is that it is only made during stellar nucleosynthesis and cannot be produced by particle irradiation in our own solar system, unlike  $^{26}\text{Al}$  (Lee et al. 1998). Thus an initial solar system value of  $^{60}\text{Fe}/^{56}\text{Fe}$  ( $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}}$ ) significantly higher than the calculated steady state background of the interstellar medium (ISM), recently estimated at  $\sim 3 \times 10^{-7}$  by Tang and Dauphas (2012), would provide strong evidence for injection of short-lived nuclides by a proximal stellar explosion. Such an explanation, linked to a trigger for solar system formation, has long been invoked to account for the elevated initial solar system  $^{26}\text{Al}/^{27}\text{Al}$  (e.g. Cameron and Turan, 1977), but alternative means of generating  $^{26}\text{Al}$  (e.g. Clayton and Jin, 1995) have made a more definitive test using  $^{60}\text{Fe}$  highly desirable (see Wasserburg et al., 1998). As explored below, values of  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}}$  determined from Ni isotopic measurements of meteorites, remain contentious although our own perspective is that values above those of the ISM are questionable.

In a pioneering contribution, Shukolyukov and Lugmair (1993a) reported evidence for live  $^{60}\text{Fe}$  from analyses of  $^{60}\text{Ni}/^{58}\text{Ni}$  on samples of the eucrite Chervony Kut. Eucrites represent a good target for detecting *in situ* decay of  $^{60}\text{Fe}$ , given their extremely high Fe/Ni ( $>10000$ ). Such high Fe/Ni ratios are a consequence of core formation at low pressures (i.e. on a small planetary body) and subsequent magmatic fractionation. Given the formation and differentiation of the

325 eucrite parent body within the first few million years of solar system history (e.g. Papanastassiou  
and Wasserburg, 1969) the resultant Fe/Ni were likely to generate resolvable  $^{60}\text{Ni}$  anomalies  
given significant initial  $^{60}\text{Fe}$ . Different components of Chervony Kut showed elevated  $\epsilon^{60}\text{Ni}_{62/58}$   
(up to 50) and three bulk samples with a range of Fe/Ni defined a straight line with a slope that  
330 implied  $^{60}\text{Fe}/^{56}\text{Fe}^\circ \sim 4 \times 10^{-9}$  (Fig. 6a). The authors used an age of  $10 \pm 2$  Ma post CAI, from the  
similarity of the meteorites'  $^{87}\text{Sr}/^{86}\text{Sr}$  to dated angrites (Lugmair and Galer, 1992), to calculate  
 $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}} = 1.6 \times 10^{-6}$ . Since then, determinations of the timing of differentiation of the eucrite  
parent body give older ages,  $\sim 3$  Ma post CAI (Lugmair and Shukolyukov, 1998; Bizzarro et al.  
2005; Trinquier et al., 2008; Schiller et al., 2010) and the accepted half-life of  $^{60}\text{Fe}$  has increased  
335 from 1.49 Ma (Kutchera et al., 1984) to 2.62 Ma (Rugel et al., 2009). As a result of these changes  
a more contemporary interpretation of these results would yield  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}} \sim 9 \times 10^{-9}$ .

Shukolyukov and Lugmair (1993a) originally noted that their value of  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  was in  
keeping with  $\epsilon^{60}\text{Ni}_{62/58}$  measurements made on CAIs by Birck and Lugmair (1988), assuming that  
such  $^{60}\text{Ni}$  excesses were radiogenic. From the discussion of nucleosynthetic variability in section  
340 2, it should be clear that this assumption is by no means valid, nor is the value thus derived  
consistent with the revised  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$ . The magnitude of possibly radiogenic  $\epsilon^{60}\text{Ni}_{61/58}$  in CAIs  
is no bigger than  $\epsilon^{62}\text{Ni}_{61/58}$ , which must be nucleosynthetic (Fig. 1a). Although Kruijier et al.  
(2014) demonstrated that it is possible to disentangle nucleosynthetic from radiogenic  
contributions to  $^{182}\text{W}/^{184}\text{W}$  in CAIs in order to define  $^{182}\text{Hf}/^{180}\text{Hf}^\circ_{\text{SSI}}$ , neither initial parent nor  
345 parent-daughter fractionation in CAIs are sufficiently large to make this approach viable for  
determining  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$ . Indeed, most CAIs have sub-chondritic  $^{56}\text{Fe}/^{58}\text{Ni}$  (e.g. Quitté et al.,  
2007), as a result of the slightly more refractory cosmochemical behaviour of Ni relative to Fe.  
Thus at best, minor  $\epsilon^{60}\text{Ni}_{61/58}$  deficits in bulk CAI would be expected if  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  were large  
enough for the bulk nebula to evolve to more radiogenic  $\epsilon^{60}\text{Ni}_{61/58}$ , as is the case for the  $^{53}\text{Mn}$ - $^{53}\text{Cr}$   
350 system (Birck and Allègre, 1985). Moreover, internal CAI isochrons offer little scope for  
constraining  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  given both modest Fe/Ni fractionation between common phases in  
CAIs and the frequent disturbance of Fe/Ni ratios in CAIs specifically (Quitté et al., 2007) and  
chondritic meteorites in general (e.g. Telus et al., 2016). In contrast to the  $^{26}\text{Al}$ - $^{26}\text{Mg}$  and  $^{182}\text{Hf}$ -  
 $^{182}\text{W}$  systems, but similar to the  $^{53}\text{Mn}$ - $^{53}\text{Cr}$  pair, CAIs therefore offer scant opportunity to  
355 determine the initial parent abundance of the  $^{60}\text{Fe}$ - $^{60}\text{Ni}$  system. Instead, basaltic achondrites, with  
their highly fractionated Fe-Ni, have been the focus of much additional work to constrain  
 $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$ .

In a follow-on study, Shukolyukov and Lugmair (1993b) analysed the eucrite Juvinas and  
360 noted an order of magnitude lower  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  ( $\sim 4 \times 10^{-10}$ ), which the authors attributed to  $\sim 4$  Ma  
evolution of the eucrite parent body mantle source between generation of the basaltic melts  
represented by Chervony Kut and Juvinas. Using the newer  $^{60}\text{Fe}$  half-life this period would be  $\sim 9$   
Ma, a value which resonates with the age difference of  $\sim 11$  Ma inferred from Hf-W analyses  
365 between a group of eucrites including Juvinas and an older group of eucrites (Touboul et al.,  
2015); sadly Chervony Kut itself was not analysed for its Hf-W systematics. As with Chervony  
Kut (Shukolyukov and Lugmair, 1993a), separated mineral phases from Juvinas did not show  
meaningful isochronous relations and Shukolyukov and Lugmair (1993b) argue that the  
mobility/diffusivity of Ni made it susceptible to resetting on this shorter length scale, especially  
370 given the complex history of eucrites, with well documented thermal metamorphism (e.g. Takeda  
and Graham, 1991).

Subsequent studies of eucrites have been made with increasingly precise analyses using  
MC-ICPMS, but in essence show comparable features. Quitté et al. (2011) reported data for  
Bouvante and further analyses of Juvinas. Whole rock sub-samples of Bouvante define two  
375 arrays and if these are taken to be isochrons they imply  $^{60}\text{Fe}/^{56}\text{Fe}^\circ = 5 \times 10^{-9}$  (but with an  
implausibly negative intercept of  $\epsilon^{60}\text{Ni}_{62/58} = -24 \pm 3$ ) and  $^{60}\text{Fe}/^{56}\text{Fe}^\circ = 5 \times 10^{-10}$  (with a reasonable,  
near zero  $\epsilon^{60}\text{Ni}_{62/58}$  intercept). Instead, the authors argue that the data represent two mixing lines  
between clasts of different compositions. Whilst the second array gives a similar value of

380  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  to Juvinas, as obtained by Shukolyukov and Lugmair (1993b), additional  
measurements of Juvinas by Quitté et al. (2011) are inconsistent with the array of Shukolyukov  
and Lugmair (1993b). Quitté et al. (2011) obtain a higher  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  ( $2 \times 10^{-9}$ ), using their two  
unwashed bulk samples and two unwashed, lower Fe/Ni samples from the Shukolyukov and  
Lugmair (1993b) array (ignoring the washed samples which they argue may have suffered Fe-Ni  
385 fractionation as a result of this preparation). As with the work of Shukolyukov and Lugmair  
(1993b), the Juvinas mineral analyses of Quitté et al. (2011) are scattered and presumably  
perturbed by secondary processes. There is clear difficulty in distinguishing primary from  
secondary signatures in these brecciated, thermally metamorphosed meteorites, but none of the  
eucrite arrays define  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  greater than  $4 \times 10^{-9}$ .

390 Tang and Dauphas (2012) comprehensively reassessed this issue using a collection of bulk  
eucrites and diogenites. This dataset gave a large range in Fe/Ni, which coupled with their high  
precision measurements (typical  $\epsilon^{60}\text{Ni}_{61/58}$  better than  $\pm 0.3$ ), provided a more definitive  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$   
=  $(3.5 \pm 0.3) \times 10^{-9}$  for the eucrite parent body. This result is notably in keeping with the higher  
values obtained from ‘internal isochrons’ on sub-samples of meteorites discussed above (e.g. Fig.  
395 6a). Using an age of  $2.4 \pm 1.1$  Ma post CAI for silicate differentiation of 4 Vesta (Trinquier et al.,  
2008; Connelly et al., 2012), which presumably sets the variable Fe/Ni seen in the eucrites, these  
data yield a  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$   $(6.6 \pm 2.5) \times 10^{-9}$ . This work of Tang and Dauphas (2012) also yields a  
bound of  $4 \pm 2$  Ma on the age of core formation on 4 Vesta, from a two-stage evolution model of  
the mantle. Namely, the  $\epsilon^{60}\text{Ni}_{61/58}$  of the mantle, determined by the intercept of the eucrite-  
400 diogenite array with an estimated bulk mantle  $^{56}\text{Fe}/^{58}\text{Ni} \sim 2700$  informs on the time since core  
formation. Tang and Dauphas (2014) subsequently used a similar approach to constrain the  
timing of core formation and hence growth of Mars. They argued that the planet reached 44% of  
its size no earlier than 1.2 Ma post CAI or otherwise the  $\epsilon^{60}\text{Ni}_{61/58}$  of the SNC meteorites they  
measured would be more radiogenic.

405 Quenched angrites provide a more petrologically robust sample for determining  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$ ,  
even if their Fe/Ni are not quite as extreme as the eucrites (cf. Figs. 6a & b). Moreover, the well-  
defined ages for these samples determined using extant isotope chronometry (Amelin 2008a,b;  
Connolly et al., 2008, Brennecka and Wadhwa, 2012), potentially provide more accurate decay  
410 correction in calculating  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$ . Three independent studies (Quitté et al., 2010; Spivak-  
Birndorf et al., 2011; Tang and Dauphas 2012) obtained consistent values for internal isochrons  
of d’Orbigny (e.g. Fig. 6b) which give a weighted average  $^{60}\text{Fe}/^{56}\text{Fe}^\circ = (3.3 \pm 0.5) \times 10^{-9}$ . Two  
internal isochrons from a second quenched angrite, Sahara 99555, are also in mutual agreement  
but yield a lower weighted mean  $^{60}\text{Fe}/^{56}\text{Fe}^\circ = (1.9 \pm 0.4) \times 10^{-9}$  (Quitté et al., 2010; Tang and  
415 Dauphas, 2015). Tang and Dauphas (2015) convincingly argue this difference relative to  
d’Orbigny likely reflects terrestrial weathering experienced by Sahara 99555.

In all, these angrite  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  are in keeping with those of eucrites, suggesting a similar  
initial  $^{60}\text{Fe}$  and timing of planetary differentiation on these two bodies. So a reassuringly  
420 consistent value has emerged from these various TIMS/MC-ICPMS studies which span several  
decades of work. The data from d’Orbigny provides the best constrained value and using an age  
of 3.9 Ma post CAI (Amelin 2008a; Brennecka and Wadhwa, 2012; Connelly et al., 2012) we  
calculate  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}} = (9.8 \pm 4.5) \times 10^{-9}$ . We note this value is lower than the equivalent cited by  
Tang and Dauphas (2012) as a consequence of our using the Pb-Pb ages rather than Mn-Cr  
425 chronometry.

Further bulk analyses of a range of meteoritic materials are supportive of the low  
 $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  determined from angrite and eucrite analyses, albeit from less well constrained  
scenarios. Shukolyukov and Lugmair (1993b) and Quitté et al. (2010) reported no systematic  
430 differences in  $\epsilon^{60}\text{Ni}_{62/58}$  for various bulk samples (ureilites) and separated phases (e.g. troilite)  
with high but variable Fe/Ni ratios. Moynier et al. (2011) placed a maximum upper bound on  
 $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  of  $3 \times 10^{-9}$  from the absence of  $^{60}\text{Ni}$  isotope anomalies in measurements of troilite



435 from the iron meteorite Muonionalusta. These troilites have Pb-Pb model ages as old as the  
quenched angrites (Blichert-Toft et al., 2010), which coupled with their high Fe/Ni (up to 1500)  
should result in radiogenic  $\epsilon^{60}\text{Ni}_{61/58}$  given sufficiently high  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$ . Although appealing  
440 targets for analysis, re-equilibration of the troilites with the surrounding Ni-rich metal during  
parent body during cooling would tend to erase any  $\epsilon^{60}\text{Ni}_{61/58}$  anomalies. The authors briefly  
argue against such an interpretation on the basis of the preservation of ancient Pb-Pb ages, but the  
potential for diffusional exchange of Ni with the Ni-rich host metal (see Chernozhukin et al.,  
2016) seems much greater than for Pb. Whilst the conclusions of Moynier et al (2011) are thus  
compatible with other studies, whether or not the measurements represent an independent  
constraint on  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  remains open to debate.

445 Analysis of chondrules from the CB<sub>a</sub> meteorite Gujba and ungrouped 3.05 ordinary  
chondrite NWA 5717 by Tang and Dauphas (2012) form near horizontal arrays that yield  
 $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  from  $1\text{-}3\times 10^{-9}$ . However, elemental mapping by Telus et al. (2016) showed that  
chondrules in all chondrites they studied had experienced some open system behaviour of Fe and  
Ni. Only the most pristine, LL3.0 meteorite, Semarkona, retained undisturbed chondrules, about  
450  $\sim 40\%$  of those studied. Prompted by these findings, Tang and Dauphas (2015) made  
measurements of single chondrules from Semarkona, to yield a valuable but still relatively poorly  
defined  $^{60}\text{Fe}/^{56}\text{Fe}^\circ = (5\pm 3)\times 10^{-9}$ . The lack of significant differences between  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  for these  
meteorites of different metamorphic grade suggests that they are not unduly compromised by this  
open system behaviour. Given an average chondrule age of 2Ma post CAI (see recent  
455 compilation of data in Budde et al., 2016), the  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}} \sim 9\times 10^{-9}$  derived from these individual  
chondrule measurements is notably compatible with the studies from achondrite meteorites. In a  
grand compilation of various determinations of  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  from bulk measurements, Tang and  
Dauphas (2015) derived a weighted average  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}} = (1.0\pm 0.3)\times 10^{-8}$ .

460 In contrast to the work described above on Ni separated from bulk samples, much higher  
 $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  have been inferred from *in situ* work by SIMS. Initially, the absence of detectable  
differences in the  $^{60}\text{Ni}/^{61}\text{Ni}$  of olivines from type II chondrules from Semarkona, relative to  
terrestrial olivines, was used to place an upper limit of  $3.4\times 10^{-7}$  on their  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  (Kita et al.,  
2000). However, later studies documented correlated  $\epsilon^{60}\text{Ni}$  and Fe/Ni in matrix sulfides and  
465 oxides from primitive ordinary chondrites suggesting  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  from  $1\times 10^{-7}$  to  $1\times 10^{-6}$  (Tachibana  
and Huss, 2003; Mostefaoui et al., 2004, 2005; Guan et al., 2007). As the time of formation of  
these phases is uncertain, the significance of these data arrays for inferring  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  was open  
to question. In an elegant study, Tachibana et al. (2006) subsequently analysed different phases  
from the chondrules of Semarkona. Although the ranges in correlated Fe/Ni and  $\epsilon^{60}\text{Ni}_{62/61}$  were  
470 lower than in the sulfide work, interpretation of the arrays as constraining  $^{60}\text{Fe}/^{56}\text{Fe}^\circ_{\text{SSI}}$  ( $5\text{-}10$ ) $\times 10^{-7}$   
seemed less equivocal. Similar results were reported by Mishra et al. (2010) for single  
chondrule analyses from a wider range of unequilibrated ordinary chondrites. Yet, all such  
analyses are controlled by the errors on the very small  $^{61}\text{Ni}$  and  $^{62}\text{Ni}$  beams used for determining  
 $\epsilon^{60}\text{Ni}_{62/61}$  (the larger  $^{58}\text{Ni}$  beam cannot be used as it is interfered by  $^{58}\text{Fe}$ ). These measurements  
475 thus critically require accurate background determination and interference free spectra.  
Moreover, the use of the minor Ni isotope in the denominator of such low intensity  
measurements can lead to a statistical bias in calculated ratios (Ogliore et al., 2011, see also  
Coath et al., 2013). This artefact resulted spurious correlations between Fe/Ni and  $\epsilon^{60}\text{Ni}_{62/61}$  in all  
earlier work (Telus et al., 2012). Yet subsequent work has continued to report high inferred  
480  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  (see Fig. 6c) from *in situ* analyses of chondrules in studies for which such statistical  
bias is argued to be insignificant (Mishra and Goswami, 2014; Mishra and Chaussidon, 2014;  
Mishra et al., 2016).

485 Hence, inferred  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  from TIMS/MC-ICPMS studies and SIMS analyses of individual  
chondrules are markedly different. This contrast in conclusions from bulk and *in situ* approaches  
also extends to other MC-ICPMS work. Regelous et al. (2008) and Steele et al. (2012) obtained  
the loose constraint that the  $^{60}\text{Fe}/^{56}\text{Fe}^\circ$  of carbonaceous chondrites was  $<1\times 10^{-7}$ , given their

indistinguishable compositions relative to IVB iron meteorites. Tang and Dauphas (2012) noted that the constant  $^{58}\text{Fe}/^{54}\text{Fe}$  in all their analyses was incompatible with nucleosynthetic models that could account for  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}} \sim 1 \times 10^{-6}$ . Explicit comparison can be made between SIMS and MC-ICPMS data for individual chondrule analyses from Semarkona, reported in the studies of Mishra and Chaussidon (2014) and Tang and Dauphas (2015) respectively. The former inferred  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ} (3 \pm 2) \times 10^{-7}$  whereas the latter, as discussed above, argued for  $(1 \pm 3) \times 10^{-9}$ . Although a detailed comparison requires knowledge of individual chondrule ages (e.g. Mishra et al., 2010) which can be variable (see Connelly et al. 2012), this cannot account for the two orders of magnitude difference in the results of the two studies. It is also worth noting that a SIMS study of basaltic achondrites reported  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ} = (6 \pm 9) \times 10^{-9}$  for quenched angrites (Siguirra et al 2006), in good agreement with the MC-ICPMS studies (Quitté et al, 2010; Spivak-Birndorf, 2011; Tang and Dauphas, 2012; Tang and Dauphas, 2015). Thus the divergent results between SIMS and MC-ICPMS are seem restricted to analyses of chondrules.

The difference in inferred  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}}$  between the bulk and *in situ* studies continues to be debated, see Mishra et al. (2016) and Tang and Dauphas (2015). Although we cannot offer an unbiased opinion, we argue strongly for the validity of the interpretations based on MC-ICPMS and TIMS analyses. Not only does this work show consistency in values obtained on a range of materials, from chondrules to bulk achondrites, but the approach removes sample matrix before analysis. The high Fe/Ni SIMS analyses typically collect only  $\sim 10000$  counts of the minor Ni isotopes and apply a background correction determined from a single point on the mass spectrum. Perhaps most critically, the *in situ* approach has not documented accurate  $^{60}\text{Ni}$  measurement for materials with the highest Fe/Ni, that define the isochrons (see Fig. 6c).

In this light, we note there may be a possible interference problem for SIMS analyses from the presence of  $^{59}\text{CoH}^+$ . Even for high resolution SIMS measurements,  $^{59}\text{CoH}^+$  overlaps considerably with  $^{60}\text{Ni}^+$  (Mishra and Chaussidon, 2014) and depending on measurement mass, it may not be resolved at all. The high Fe/Ni portions of chondrules are the result of fractional removal of olivine and/or low-Ca pyroxene and troilite from the cooling melt droplet. Ni is considerably more compatible than Fe in these fractionating phases and so the residual melt (and phases that subsequently crystallise from it) acquire high Fe/Ni. Ni is also more compatible than Co under the same conditions and so as crystallisation proceeds, Co/Ni will also increase. Thus high Fe/Ni in chondrules is inevitably associated with high Co/Ni. We have quantified this process in Fig 7, using the partition coefficients from experiments designed to mimic crystallisation of chondritic liquids at low pressure and which provide simultaneously determined Fe, Ni and Co data (Gaetani and Grove, 1995). In detail, the evolution of Fe/Ni and Co/Ni depend on the amount of fractionating sulfide. We examine two scenarios using relatively low (0.2%) and high (1%) amounts of sulfide for Semarkona Type II chondrules, taken from the work of Jones (1990). Although this parameter changes the amount of crystallisation required to reach a given Fe/Ni, the amount of sulfide little affects the overall trend. We show that the extreme  $^{56}\text{Fe}/^{58}\text{Ni} \sim 7000$  of the key SIMS analyses will be associated with  $^{59}\text{Co}/^{60}\text{Ni} \sim 200$  (Fig. 7). This is the model value for a ferro-magnesian phase (specifically olivine) grown from the residual melt, as this was the main target in the study of Mishra and Chaussidon (2014). Assuming that the fraction of hydride production determined for iron ( $^{56}\text{FeH}/^{56}\text{Fe} \sim 5 \times 10^{-5}$ , Mishra and Chaussidon, 2014) can be used for cobalt, this implies a  $^{59}\text{CoH}$  contribution of 1% on the mass 60 peak (i.e.  $\Delta^{60}\text{Ni} \sim 10$ ), sufficient to explain the anomalous  $^{60}\text{Ni}$  reported. This suggested explanation of the divergence of bulk and *in situ* inferences can be tested by measuring Co intensities during SIMS analysis.

In summary we infer  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}} \sim 1 \times 10^{-8}$ , although this is not a consensus view. This conclusion implies that the solar system did not form with  $^{60}\text{Fe}/^{56}\text{Fe}$  discernibly different from ISM. Thus if the elevated  $^{26}\text{Al}$  of the solar system is to be explained by recent injection of stellar material, this must come from a source which does not also generate significant  $^{60}\text{Fe}$ . Tang and Dauphas (2012) suggest that winds from the outer portions of a giant, Wolf-Rayet star could

constitute an appropriate explanation. These results argue against a supernova trigger of solar system collapse (see Wasserburg et al., 1998).

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#### 4. MASS-DEPENDENT Ni ISOTOPIC VARIABILITY

##### 4.1 Magmatic Systems

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Mass-dependent Ni isotopic variations in magmatic systems have received relatively little interest. This may be due to the lack of redox variability of Ni in common silicate phases on Earth, which has been a prime driver for exploration of several other isotopic systems, based on the notion that redox related changes in bonding environment will lead to large mass-dependent fractionations. However, the notable change in Ni oxidation state from 2+ to 0 and partitioning during core formation provides a significant point of interest.

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In common with many other isotopic systems recently investigated, commercially available metals can show marked mass-dependent isotopic variability. For example, Tanizuma and Hirata (2006) reported a range of 0.6‰ in the  $\delta^{60/58}\text{Ni}$  of a selection of high purity commercial Ni reagents. The authors noted the potential of isotopic fractionation during purification of Ni by the Mond process, which involves separation of Ni in the vapour phase as the volatile  $\text{Ni}(\text{CO})_4$  species. However, two Ni ores analysed by Tanizuma and Hirata (2006) showed an even greater variability ( $\delta^{60/58}\text{Ni} = 0.5$  to  $-0.4$ ), spanning the compositions of the laboratory Ni reagents and making it hard to distinguish the influences of anthropogenic from natural fractionation. An important consequence of the fractionated values of  $\delta^{60/58}\text{Ni}$  seen in purified Ni is that NBS SRM 986 is sadly not representative of bulk terrestrial reservoirs, either in mass-dependent (Cameron et al., 2009; Steele et al., 2011; Gall et al., 2012; Gueguen et al., 2013) or mass-independent Ni isotopic compositions (Steele et al., 2011). The latter likely stems from imperfect mass-bias correction using a single exponential form for combined instrumental, natural and industrial fractionations. Separation of Ni as a carbonyl species in the Mond process but its measurement as an elemental ion in mass spectrometry offers a tangible mechanism to account for this feature (Steele et al., 2011). Although cosmetically unappealing, the slight deviation of the Ni reference standard from bulk terrestrial values does not diminish its value as a common datum.

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As a planetary reference Moynier et al. (2007) reported  $\bar{\delta}\text{Ni}$  for a wide range of iron meteorites and metal phases in chondrites. Planetary cores dominate the Ni budgets of planets and planetesimals, so iron meteorites should provide a representative mass-dependent Ni isotopic composition of their parent body. Likewise a metal phase will control the Ni bulk composition of host chondrites. These sample-standard bracketing data of Moynier et al. (2007) are replotted as  $\delta^{60/58}\text{Ni}$  in Fig. 8a and vary from 0-0.8‰ with an average  $\sim 0.3$  ‰ (individual errors  $\sim \pm 0.15$  ‰). Cameron et al. (2009) presented  $\delta^{60/58}\text{Ni}$  for a representative selection of bulk chondrite and iron meteorites by double spiking, which showed less scatter than the Moynier et al (2007) data but with a similar average of  $0.27 \pm 0.06$  ‰. A set of iron meteorites measured by sample-standard bracketing (Chernonozhkin et al., 2016) and two higher precision double-spiked measurements of chondrites Steele et al. (2012) fall within this range. These studies are combined to give a planetary reference of  $\delta^{60/58}\text{Ni} = 0.24 \pm 0.04$  (Table 2). Whilst there may be systematic differences in  $\delta^{60/58}\text{Ni}$  between different meteorite types (cf. Luck et al., 2003), given the limited number of samples currently available and their small variability relative to precision, we have done no more than to compile this average, which implicitly assumes a uniform bulk Ni mass-dependent isotopic composition for different planetary objects in the solar system.

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To date, most literature  $\delta^{60/58}\text{Ni}$  analyses of peridotitic material or mantle derived melts, appropriate for estimating the composition of the silicate Earth, are reference materials. In their assessment of a baseline value for silicates, Cameron et al. (2009) found that samples from peridotite to granite, in addition to a range of continental sediments varied little,  $\delta^{60/58}\text{Ni} = 0.15 \pm 0.24$ . Combining two higher precision measurements of fresh peridotites with the data of

595 Cameron et al. (2009), Steele et al. (2011) calculated a weighted average of  $0.18 \pm 0.04$  ‰ for the  
silicate Earth. Additional analyses of mafic and ultra-mafic samples by Gall et al. (2012),  
Gueguen et al. (2013) and Chernozhkin et al. (2015) support such isotopically lighter values of  
600 the terrestrial mantle relative to the meteorite reference (Fig. 8). Gueguen et al. (2013) compiled  
analyses of silicate samples to calculate  $\delta^{60/58}\text{Ni} = 0.05 \pm 0.05$  for bulk silicate Earth. The latter  
includes measurements of deep sea clays, as representative of the continental crust. Given the  
likely pelagic component of some Ni in these sediments and possible associated isotopic  
fractionations (see below), coupled with the trivial contribution of crustal Ni to bulk silicate  
Earth, the inclusion of these sediments is unwarranted but does not unduly bias the estimate.  
605 Revising this value to include only samples of the mantle and mantle derived melts, a better  
defined  $\delta^{60/58}\text{Ni} = 0.11 \pm 0.04$  is obtained for the bulk silicate Earth (Table 2).

These data thus imply that the bulk silicate Earth is slightly isotopically light relative to a  
bulk planetary reference, by  $\sim 0.15$  ‰. Most plausibly this is the result of core formation, as Ni is  
a siderophile element and strongly partitioned into the metal phase. The sense of isotopic  
610 fractionation between the silicate Earth and chondrites is consistent with the isotope partitioning  
experiments of Lazar et al. (2012), who determined a fractionation factor between metal and talc,  
 $\Delta^{62/58}\text{Ni}_{\text{metal-silicate}} = (0.25 \pm 0.02) \times 10^6 / T^2$ . Using this expression to calculate isotopic fractionation  
between metal and silicate during core formation predicts a rather more modest than observed  
615 difference between bulk and silicate Earth of  $\sim 0.02$  ‰ for a single core-mantle equilibrium at  
2500 K. Yet, as the authors note, larger net fractionations are possible for a Rayleigh process  
with a well mixed mantle magma ocean. Clearly, there is also scope for measuring isotopic  
fractionation under conditions that more closely approximate core formation, although these are  
difficult experiments. Empirical constraints on high temperature fractionation factors could  
620 potentially be gleaned from the analyses of some metal and silicate separates from mesosiderites  
and pallasites presented by Chernozhkin et al. (2016). However, the sense of fractionation  
between metal and silicate in the two meteorite types is inconsistent and an important role for  
kinetic (diffusive) fractionation during cooling seems likely (Chernozhkin et al., 2016).

In contrast to the relative constancy of  $\delta^{60/58}\text{Ni}$  in many silicate samples noted above, Steele  
625 et al. (2011), Gueguen et al. (2013) and Hofmann et al. (2014) reported a wide range of  
isotopically light values ( $\delta^{60/58}\text{Ni} = -0.1$  to  $-1$ ) in Ni-bearing sulfides from komatiitic magmas, see  
Fig. 8b. Such ores are believed to be formed by sulfide saturation following crustal assimilation  
by these ultra-mafic magmas (Huppert et al., 1984). The Ni in the sulfides is dominantly from  
the mantle derived melts and so the data imply a significant Ni isotopic fractionation between  
630 sulfide and silicate magma, although kinetic effects may also play a role. These low  $\delta^{60/58}\text{Ni}$  echo  
the value of  $-0.34$ ‰ measured previously for a millerite ore from the Thompson Ni Belt  
(Manitoba, Canada) by Taminura and Hirata (2006). Despite experiencing some secondary  
reworking, these Ni deposits were similarly formed by the interaction of ultra-mafic melts with  
crustal sulfides (e.g. Bleeker and Macek, 1996) and so this process generally seems to result in  
635 low  $\delta^{60/58}\text{Ni}$ . On the other hand, an isotopically heavy pentlandite from Sudbury ( $\delta^{60/58}\text{Ni} \sim 0.5$ )  
was reported by Taminura and Hirata (2006). These sulfides were formed from a meteorite  
impact induced melt sheet, with the Ni largely sourced from the molten continental crust (Walker  
et al., 1991). This different source of Ni for the Sudbury ore relative to the komatiite hosted  
samples is a plausible explanation of their contrasting  $\delta^{60/58}\text{Ni}$ . Another style of Ni ore deposit,  
640 produced from lateritic weathering of ultramafic lithologies in Brazil, exhibits less anomalous  
compositions,  $\delta^{60/58}\text{Ni} = 0.02$ - $0.2$  (Ratié et al., 2016), implying little net fractionation during their  
formation.

#### 4.2 Ni isotope fractionation during weathering and the hydrological cycle.

645 As discussed above, there is currently only a single experiment to underpin understanding of  
the controls on Ni isotopic fractionation at magmatic temperatures (Lazar et al., 2012). In  
contrast, isotopic fractionations between Ni species relevant for the hydrological cycle have  
received more extensive, theoretical investigation (Fujii et al. 2011; Fujii et al. 2014). Despite

650 the absence of redox changes at the Earth's surface, differences in the bonding environment  
between common, aqueous Ni species result in significant isotopic fractionations (Fig. 9). These  
first principles calculations for a homogenous aqueous system provide a useful guide to potential  
magnitudes of Ni isotopic fractionation but do not currently provide direct information on the  
effects of sorption of Ni onto Fe/Mn oxide surfaces, for example, which likely play an important  
655 role in the surface cycle of Ni (e.g. Peacock and Sherman, 2007). However, experimental work  
by Wasylenki et al. (2015) documents that sorption of Ni onto the ferrihydrite favours the lighter  
isotopes, leaving an isotopically heavier residual Ni in solution (Fig. 9).

The processes of weathering, transport to and removal from the ocean can result in major  
isotopic fractionations, such that the isotopic composition of seawater is a notable end-member  
660 for a number of elements (e.g. Li, B, Mg, Mo). In a comprehensive study of seawater, the first to  
define its Ni isotopic composition, Cameron and Vance (2014) reported  $\delta^{60/58}\text{Ni} = 1.44 \pm 0.15$   
which is near constant across all ocean basins and at various depths. Seawater is thus markedly  
isotopically heavy relative to typical silicate rocks (Fig. 8b,c). Cameron and Vance (2014) further  
665 documented that the dissolved Ni loads of rivers are dominantly isotopically heavy relative to  
unaltered silicates (Fig. 8b,c), showing that weathering processes result in significant Ni isotopic  
fractionation. Similarly, the  $\delta^{60/58}\text{Ni} \sim -0.5$  of organic-rich marine shales (Fig. 8d) is argued to  
reflect the end product of the weathering process (Porter et al., 2014), as is also likely the case for  
the more modestly elevated  $\delta^{60/58}\text{Ni}$  of coal and banded ironstone formation (BIF) standards  
(Gueguen et al., 2013). As a corollary, several studies have shown that soils (Fig. 8d), the  
670 residues of weathering, are commensurately isotopically light (Estrade et al., 2015, Ratié et al.  
2014, Gall et al., 2013).

A major sink of Ni in the oceans is its incorporation into Fe-Mn crusts (e.g. Peacock and  
Sherman, 2007). A detailed study of these materials by Gall et al. (2013) showed  $\delta^{60/58}\text{Ni}$  that  
675 ranged from 0.9-2.5 ‰, with modern samples yielding an average of  $1.6 \pm 0.8$  ‰ (Fig. 8c). This  
work found little systematic variation in the  $\delta^{60/58}\text{Ni}$  of samples between ocean basins or indeed  
over an  $\sim 80$  Ma record in the Pacific. In large part, it appears that the  $\delta^{60/58}\text{Ni}$  of Fe-Mn crusts  
reflects the isotopically heavy composition of seawater and so this sink does not represent a  
major source of isotopic fractionation as is the case in the Mo isotope cycle (e.g. Barling et al.,  
680 2001). Earlier in geological history, sorption of Ni to iron oxide phases, preserved in BIF, may  
have been a significant sink for oceanic Ni. The near constant  $\Delta^{60/58}\text{Ni}_{\text{dissolved-sorbed}} = 0.35 \pm 0.10$   
determined by for Ni sorption onto ferrihydrite over a range of conditions (Wasylenki et al.,  
2015) gives encouragement for recovering information on past ocean compositions from BIF. It  
is worth noting that this fractionation, however, is in the wrong direction to account for the  
685 isotopic difference between seawater and modern Fe-Mn crusts (Fig. 8c) by such a sorption  
process. On the other hand, the sense of fractionation could help account for the isotopically  
heavy composition of Ni in rivers as the product of weathering.

Although the first order observations on the hydrological cycling of Ni isotopes have thus  
690 recently been outlined, the isotopic balance of Ni is still not obviously complete (Cameron and  
Vance, 2014) and this further echoes a long identified elemental imbalance between oceanic  
inputs and outputs for Ni (e.g. Krishnaswami, 1976). Whilst rivers are isotopically heavy inputs  
to the oceans, the oceans are  $\sim 0.6$  ‰ heavier than the mean riverine input. There always remains  
an uncertainty over the representativeness of the river sample, but the magnitude of the Ni  
695 isotopic mis-match is troublesome. Since Fe-Mn crusts are isotopically heavier, on average, they  
do not provide a ready explanation to the problem. A tempting scenario would be to invoke the  
role of kinetic isotopic fractionation of Ni resulting from its essential role in several biological  
pathways in the marine environment (see below). Yet the lack of difference in  $\delta^{60/58}\text{Ni}$  between  
surface and deep water limits the magnitude of this light isotope sink (Cameron and Vance,  
700 2014). Potential solutions to resolve this imbalance are discussed by Cameron and Vance (2014)  
who tentatively propose the release of Ni carried on Fe-Mn coatings from the suspended load  
during estuarine processes. The authors argue against the role of serpentinisation, given an

705 absence of vertical  $\delta^{60/58}\text{Ni}$  variation in the water column. Yet, a single, oceanic serpentine sample reported by Gueguen et al. (2013) is notably isotopically light ( $\delta^{60/58}\text{Ni} = -0.13$ ), at high  $[\text{Ni}] \sim 1500\text{ppm}$  and so tantalisingly hints at a plausible reservoir to complement the isotopically heavy oceans.

710 Most recently, however, measurements of isotopically heavy Ni ( $\delta^{60/58}\text{Ni} \sim 2$ ) in euxinic deep waters of the Black Sea (Vance et al., in review) appear to be the complement to partial removal to the sediment pile of a sorped, Ni sulfide species, which Fujii et al. (2011) predicted to be isotopically light (Fig. 9). If such a process occurred more generally in the sediment interface on upwelling margins of the Pacific, which become euxinic a few centimetres below the surface, this might provide an answer to oceanic Ni isotopic balance (Vance et al., in review).

### 715 4.3 Biological Mass-dependent Fractionations

720 As with some other first row transition elements, Ni plays an important role in several biogeochemical cycles. It is an essential trace element for key enzymes in methane production and decomposition of urea. Kinetic fractionation in metabolic pathways will favour the least strongly bound, lightest isotopes and given this basic notion it might be anticipated that Ni incorporated into biological systems would have low  $\delta^{60/58}\text{Ni}$ . On the other hand, the first principles calculations (Fujii et al., 2011; Fujii et al., 2014) show substantial equilibrium fractionation of Ni isotopes between aqueous species complexed with different naturally occurring organic ligands (Fig. 9). Fractionation factors at 298K, calculated between the hexa-aqua ion, the dominant Ni species over a range of conditions in natural waters and complexes with oxalate, phosphate and citrate ions are shown in Figure 9. These examples are indicative of the range of fractionations reported by Fujii and co-workers. Moreover they serve as model representations of possible bonding environments for Ni in soils and plants (Fujii et al., 2011; Fujii et al., 2014). If biological pathways efficiently discriminate between such species, macroscopic Ni isotopic variability can result. Yet, some caution needs to be exercised in using simple equilibria to predict the effects of a complex biochemical reaction sequence. For example, the inorganic oxidation of iron results in a 2.75‰  $^{56}\text{Fe}/^{54}\text{Fe}$  isotopic fractionation at room temperature (Johnson et al., 2002), whereas if this processes is biologically mediated through iron oxidising bacteria, a much smaller 1.5‰ fractionation is observed (Croal et al., 2004).

735 Given the critical role of Ni in biological methanogenesis, it might be anticipated that Ni isotopic fractionation would be associated with species that had evolved to incorporate environmental Ni to these ends. In a reduced environment, such as the early Earth, methanogens likely represent a major component of the surface Ni cycle and have the potential to impart an isotopic record of their presence. In the first study to look at the effects of biological processes on Ni isotope compositions, Cameron et al. (2009) examined Ni isotopic fractionations during laboratory growth of 3 species of methanogens (2 mesophiles, *Methanosarcina barkeri* and *Methanosarcina acetivorans* and one hyperthermophile, *Methanococcus jannaschii*) and one non-methanogen, a heterotrophic archaea (*P. calidifontis*) as reference. Cameron et al. (2009) showed that Ni incorporated into all three methanogens was isotopically light compared to the media in which they grew,  $\Delta^{60/58}\text{Ni}_{\text{cells-media}} \sim -1$  but  $\Delta^{60/58}\text{Ni}_{\text{cells-media}} \sim 0$  for the heterotrophic archaea, *P. Calidifontis* (Fig. 9).

750 These substantial fractionations show promise as a means for identifying the active role of methanogens in the rock record of the early Earth. As noted earlier, there is a near constant and smaller fractionation when Ni is sorbed from aqueous solution to ferrihydrite (Wasylenki et al., 2015) and further experiments suggest that such Ni remains immobile during subsequent diagenesis (Robbins et al., 2015). Hence the isotopic composition of Ni in BIF, originally sorbed from Archaean seawater is anticipated to be isotopically heavy for a system in which methanogens dominated the marine Ni cycle. Analyses of BIF of different ages therefore should provide a possible means to investigate the proposed demise of the methanogens prior to the rise in

atmospheric oxygen ~2.4 Ga (Konhauser et al., 2009). Yet other fractionations could possibly produce a similar signature (Fig. 7). As discussed above, the partial removal of an isotopically light Ni sulfide species has been invoked to account for modern ocean Ni mass balance. This process can impart a similar sense and magnitude of Ni isotopic fractionation to residual ocean water as Ni sequestration by methanogens (Fig. 9). Such *partial* removal of sulfide species, however, seems unlikely to have occurred under the fully reduced oceanic conditions of the Archean and should have become more important at the onset of oxidation- the converse to that expected from the activity of methanogens.

Ni is an essential micro-nutrient in higher plants (e.g. Brown et al. 1987) given its functional presence in the enzyme urease (Dixon et al., 1975). However, Ni abundances as low as 90ng/g can be sufficient to maintain this function (Dixon et al., 1975) and so plants do not represent a large Ni reservoir, with Ni concentrations typically less than 10µg/g. Exceptions are the so-called Ni hyperaccumulators, that have evolved to be tolerant of Ni rich soils that develop on ultramafic lithologies. Many plants deal with the toxicity of metals by sequestering them in their roots, but hyperaccumulators transport large quantities of metals to their leaves. Such Ni hyperaccumulators make interesting subjects for Ni isotopic investigation, as high Ni concentrations make them amenable to analysis and their specialisation to cope with high Ni concentrations might lead to characteristic isotopic fractionations.

Laboratory experiments (Deng et al., 2014) and a field study (Estrade et al. 2015) reported isotopic analyses for several Ni hyperaccumulators, together with a non-accumulator species for comparison. In all but one case, the plants have net isotopically light compositions relative to dissolved Ni in growth media or inferred for soil water ( $\Delta^{60/58}\text{Ni}_{\text{bulk plant-solution}} = -0.06$  to  $-0.8$ ), see Figure 9. In the case of field study, fractionation factors were derived relative to the bio-available Ni obtained by extraction from the soil with diethylenetriaminepentaacetic acid (Estrade et al., 2015) and this process adds some uncertainty to the values derived. However, the isotopically heavy composition of this extractable Ni would be consistent with predictions if  $\text{Ni}(\text{C}_2\text{O}_4)_3^{2-}$  is used as a proxy for Ni bound to humic acid, as suggested by Fujii et al. (2011), and removed from the bio-available pool. The largest Ni isotopic fractionations are observed in some laboratory experiments with Ni-hyperaccumulators but in general there is not a clear distinction between hyperaccumulators and non-hyperaccumulators (Fig. 9). The overall light isotopic composition of the plants is likely the result of kinetic fractionation during Ni up-take by 'low affinity' transport systems across the cell membranes of plant roots. The dominant role of such ion channels, as opposed to 'high affinity' carrier-mediated transport, had previously been invoked by Weiss et al. (2005) to account for the light bulk Zn isotopic compositions of plants they studied.

Isotopically lighter Zn in shoots relative to roots was first reported by Weiss et al., (2005) and similar features are observed in the Ni isotopic studies of plants (Deng et al., 2014; Estrade et al., 2015). In a comparison of shoot-root isotopic differences of Zn and Ni in the same species, Deng et al. (2014) noted a more modest fractionations for Ni. They attributed this to higher Ni mobility in plants relative to Zn. It should be re-emphasised that most of the species analysed for their Ni isotopic compositions are hyperaccumulators, in which Ni is anomalously enriched shoots relative to roots. The two non-accumulators studied so far show contrasting behaviour, with both positive, *Thlaspi arvensa* (Deng et al., 2014) and negative, *Euphorbia spinosa* (Estrade et al., 2015) values of  $\Delta^{60/58}\text{Ni}_{\text{shoots-roots}}$ , see Figure 9. Some variability in  $\Delta^{60/58}\text{Ni}$  between roots, stems and leaves may be related to the growth stage of the plants (Estrade et al., 2015). A full explanation of the diverse results requires more detailed work on the transport of Ni in plants, which is currently poorly understood. However, these data do provide a means to assess the generality of a hypothesis proposed to explain plant isotopic variability Zn system. Fujii and Albarède (2012) invoked isotopic fractionations between Zn phosphate species concentrated in the roots and citrates and malates in the leaves. In the case of Ni, however, Fujii et al. (2014) calculated that no discernible fractionations are predicted for this mechanism (Fig. 9) and so this cannot explain the root-shoot variability in  $\delta^{60/58}\text{Ni}$ .

815 Not only is Ni incorporated by living plants, but also during diagenesis of plant matter Ni  
frequently becomes bound by porphyrins, derived from chlorophyll, that originally complexed  
Mg. As an important trace element in oils, the abundance of Ni has, especially in relation to V,  
long been used as a tracer of source and evolution of oil reservoirs (e.g. Hodgson, 1954).  
820 Adding the additional constraint of the Ni isotopic composition of crude oils is an intriguing new  
means to understand oil formation. A preliminary study by Ventura et al. (2015) report  $\delta^{60/58}\text{Ni}$   
(together with V and Mo isotopic compositions) for a global selection of crude oils. These show a  
limited range of  $\delta^{60/58}\text{Ni}$  from 0.4 to 0.7, which overlaps the composition of organic-rich shales  
(Fig. 8c).

## 5. OUTLOOK

825 There are exciting prospects for Ni isotopic measurements across the fields outlined above.

830 Although the most anomalous macroscopic objects in the solar system, CAIs have not been  
systematically studied for their Ni isotopic compositions and rather few analyses currently exist.  
A recent study has reported extreme  $\epsilon^{64}\text{Ni}_{61/58}$  (up to 55) in two Allende CAIs (Chen and  
Papanastassiou, 2014). Further investigation of Ni isotopic compositions of CAIs and their  
relation to anomalies of other neutron rich nuclides would be valuable (e.g. recent discussion by  
Chen et al., 2015). More generally, finding the carrier that is responsible for bulk meteoritic Ni  
isotopic variation is a key quest.

835 It is clearly critical to resolve definitively the value of  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}}$ , given its implications for  
the seeding of the solar system with exotic isotopic signatures. We strongly favour low values of  
 $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}_{\text{SSI}}$ , consistent with background ISM, but a community consensus is required. We  
suggest that exploring the potential role of interfering  $\text{CoH}^+$  as highlighted in section 3 would be  
a useful step forward.

840 The measurable abundances of Ni in both silicate and iron meteorites allows valuable  
comparison of different meteoritic objects. To date this has shown that in contrast to the other  
magmatic irons so far analysed, the IVB iron meteorites have affinities with carbonaceous rather  
than ordinary chondrites. It would be useful to make a more comprehensive investigation of  
845 these genetic relationship in other, so far unanalysed meteorite types.

850 As with Mg, for which Ni typically substitutes, there appears to be little mass-dependent Ni  
isotopic fractionation during silicate partial melting and differentiation. This observation would  
benefit from more detailed study and unpublished work by Gall (2011) suggests that there might  
be some systematic silicate melt-solid fractionations at magmatic temperatures. Publication of  
these and additional empirical data in conjunction with experimental determinations of  
fractionation factors would be helpful. In the absence of major fractionations during silicate  
differentiation, the small, but apparently significant fractionation of Ni isotopes during core  
855 formation becomes of interest, but a more robust estimate of bulk silicate Earth is required. The  
current value relies heavily on a few reference standards. Additional observations of  $\delta^{60/58}\text{Ni}$  on  
the silicate mantles of smaller planetary bodies, where the effect might be larger (due to greater  
Ni partitioning into the core at low pressure), should prove informative. Interpretation of such  
results would greatly benefit from further experimental determinations of silicate-metal  
fractionation factors.

860 Better understanding the seemingly significant fractionation of Ni isotopes during  
weathering processes and a clear resolution of the apparent isotope imbalance in the ocean are  
obvious targets for future research. Comparisons with the better studied, non-redox sensitive,  
divalent cations, Mg and Ca, would be welcome. The contrast in oceanic sinks for these different



865 elements (Fe-Mn crusts, versus seafloor alteration versus carbonate production) makes  
differences in their behaviour through time of considerable interest.

870 As with other novel, isotopic systems, their nascent application to biological systems carries  
unknown potential. The large Ni isotopic fractionation of methanogens is a specifically  
appealing prospect, but whether or not such Ni isotopic fractionations remains locked in the rock  
record and sufficiently distinct over background variability, including abiotic reactions in the  
weathering environment, remains to be seen.

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## ACKNOWLEDGEMENTS

In the writing of this review TE was supported by the ERC (AdG 321209 ISONEB) and RCJS by ETH. This funding is gratefully acknowledged. The authors wish to thank the Editor for his patience and two anonymous reviewers for numerous useful suggestions, kindly pointing to some egregiously missed references and encouragement to expand the final section to speculate in a field in which we have limited expertise. A timely, final sanity check by Derek Vance and access to his unpublished manuscript was appreciated. The authors would like to doff their caps to the many who have helped in the development of their ideas about the Ni system, not least Chris Coath, Marcel Regelous, Hart Chen, Matthias Willbold, Derek Vance, Vyll Cameron and Corey Archer,

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1340 **TABLE AND FIGURE CAPTIONS**

Table 1 Ni isotopic abundances (Gramlich et al. 1989), nuclide masses (Wang et al., 2012) and atomic weight (Wieser et al., 2013).

1345 Table 2 Averages of mass-dependent and mass-independent Ni isotope compositions for meteorites and terrestrial reservoirs. Averages are weighted means and uncertainties of  $2 SE_w$  of multiple analyses unless indicated with \* which denotes arithmetic mean uncertainty of  $2 SD_w$ . The weighted mean is the preferred estimator of central tendency when data have differing uncertainties, as in the case of combining different studies. The weighted standard error has been  
 1350 estimated using the expression,  $SE_w = \sqrt{\frac{1}{\sum_i^n \sigma_i^{-2}}}$ . The weighted standard deviation is given by,  

$$SD_w = \sqrt{\frac{\sum w_i(x_i - \bar{x})^2}{(N-1)\sum w_i n^{-1}}}$$
. Both the  $SE_w$  and  $SD_w$  have been adjusted by multiplying by the expression  $1 + \frac{1}{4(n-1)}$  which compensates for the bias in standard deviation (and consequently standard error) of small sample sizes (Gurland and Tripathi 1971). Data sources indicated as follows: [A] Birck and Lugmair (1988) [b] Cameron et al. (2009) [c] Cameron and Vance (2014)  
 1355 [D] Chen et al. (2009) [e] Chernozhkin et al. (2015) [f] Chernozhkin et al. (2016) [G] Cook et al. (2006) [H] Cook et al. (2006) [I] Cook et al. (2008) [J] Dauphas et al. (2008) [k] Estrade et al. (2015) [l] Gall et al. (2013) [m] Gueguen et al. (2013) [n] Kita et al. (1998) [o] Porter et al. (2014) [P] Quitté et al. (2011) [Q] Quitté et al. (2006) [R] Quitté et al. (2007) [s] Ratié et al. (2015) [T] Regelous et al. (2008) [U] Shukolyukov and Lugmair (1993a) [V] Shukolyukov and  
 1360 Lugmair (1993b) [W] Steele et al. (2011) [X] Steele et al. (2012) [Y] Tang and Dauphas (2012) [Z] Tang and Dauphas (2014) [AA] Tang and Dauphas (2015) [AB] Tanimizu and Hirata (2006)

Figure 1. Mass independent nickel isotope data for meteorites and CAIs from earlier studies, (a)  $\epsilon^{60}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$  and (b)  $\epsilon^{64}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$ . Data from Birck and Lugmair 1988; Cook et al. 2006, 2008; Quitté et al. 2006, 2007. The dashed line indicates the vector of change in composition (from the origin) as a result of error or interference (moving to the lower left) in  $^{61}\text{Ni}$ . Some scatter appears to be along this trajectory. For ready comparison with the higher precision data, the blue boxes indicate the dimensions of parameter space shown in Figure 4. An extrapolation of a least squares York regression and the  $2\sigma$  uncertainty error envelope (York 1968; Mahon 1995) through the bulk meteorite data of Steele et al. 2012 and Tang and Dauphas (2012, 2014) from the origin through the CAI data (see Fig. 4) is shown in Figure 1b for reference.  
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Figure 2. Mass independent nickel isotope data of sulfide inclusions from iron meteorites (Quitté et al. 2006; Cook et al. 2008; Chen and Papanastassiou., 2009), (a)  $\epsilon^{60}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$  and (b)  $\epsilon^{64}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$ . Again the dashed line indicates trajectories caused by error or interference (moving to the lower left) on  $^{61}\text{Ni}$ . It appears that  $^{61}\text{Ni}$  error may be a dominant (but not sole) cause of variability in the highly anomalous sulfides. The blue boxes indicate the dimensions of parameter space shown in Figure 4.  
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Figure 3. Mass independent nickel isotope data a)  $\epsilon^{60}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$  and b)  $\epsilon^{64}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$  of bulk meteorite analyses from ‘second generation’ studies (Dauphas et al., 2008; Regelous et al., 2008; Chen and Papanastassiou, 2009). The samples of Regelous et al. 2008 show small but resolved anomalies in both chondritic and iron meteorites. These analyses are consistent with the less precise data of Dauphas et al. (2008), who conversely argued against bulk Ni isotopic variability (as did Chen and Papanastassiou, 2009). Lines representing  $^{61}\text{Ni}$  error and best fit array of Steele et al. (2012) and Tang and Dauphas (2012, 2014) as discussed in caption  
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to Figure 1. Again (dashed) blue boxes indicate the dimensions of parameter space shown in Figure 4.

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Figure 4. Mass independent Ni isotope measurements a)  $\epsilon^{60}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$  and b)  $\epsilon^{64}\text{Ni}_{58/61}$  vs.  $\epsilon^{62}\text{Ni}_{58/61}$  on bulk meteorites from the highest precision studies (Steele et al., 2011, 2012; Tang and Dauphas, 2012, 2014). Data from these studies show a consistent set of anomalies between carbonaceous, ordinary and enstatite chondrites and different groups of iron meteorites. There is a strong positive correlation between  $\epsilon^{62}\text{Ni}_{58/61}$  and  $\epsilon^{64}\text{Ni}_{58/61}$  with a slope of  $3.003 \pm 0.116$  (fitted as described in the caption to Figure 1, and illustrated with black line and grey error envelope), first identified by Steele et al. (2012) and confirmed by Tang and Daphaus (2012). This is array quite distinct from the dashed  $^{61}\text{Ni}$  error line.

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Figure 5 Plot of  $\epsilon^{64}\text{Ni}_{58/61}$  vs  $\epsilon^{50}\text{Ti}_{47/49}$  for bulk meteorites (Ti isotope data from Trinquier et al. 2009, Ni data from Steele et al 2012 and Birck and Lugmair 1988). A linear array is defined by analyses of ordinary, enstatite and CI chondrites and this extends intersect with CAI compositions. However, CAI-rich meteorite groups (CV CO) plot above the array, likely reflecting analyses of unrepresentative sub-samples of these meteorites with an excess of CAI ~3-4% by mass (illustrated with a mixing curve from a CI composition). The high Ti/Ni CAIs mean mixing trajectories with bulk chondrite compositions are highly convex up. Thus simple mixing between CAIs and ordinary chondrites (black curve) cannot explain the straight bulk meteorite array.

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Figure 6. Representative isochron diagrams for the  $^{60}\text{Fe}$ - $^{60}\text{Ni}$  system measured by different analytical techniques. (a) TIMS Isochron of different bulk samples of the eucrite Chervony Kut (Shukolyukov and Lugmair 1993a). (b) MC-ICPMS Isochron of mineral separates from the quenched angrite D'Orbigny (Tang and Dauphas 2012). (c) SIMS Isochron of a single Efremovka chondrule (Ch 1), Mishra and Chaussidon (2014). Also shown are recalculated least squares linear York regressions (York 1968, Mahon 1995) for each dataset, yielding  $^{60}\text{Fe}/^{56}\text{Fe}^{\circ}$ . Note the highly contrasting values obtained by bulk (a) and (b) versus *in situ* analysis (c) for objects of comparable age.

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Figure 7. A plot of the evolution of Co/Ni (expressed as  $^{59}\text{Co}/^{60}\text{Ni}$ ) and Fe/Ni (as  $^{56}\text{Fe}/^{58}\text{Ni}$ ) with fractional (Rayleigh) crystallisation of a model chondrule. The values shown are for an olivine crystal that is grown from the evolving melt. We have used representative partition coefficients from Gaetani and Grove (1995),  $D(\text{Ni})_{\text{olivine-melt}} = 5.9$ ,  $D(\text{Co})_{\text{olivine-melt}} = 1.96$ ,  $D(\text{Ni})_{\text{sulfide-melt}} = 0.916$  and  $D(\text{Co})_{\text{sulfide-melt}} = 510$ ,  $D(\text{Fe})_{\text{sulfide-melt}} = 17.7$ ,  $D(\text{Fe})_{\text{sulfide-melt}} = 1.53$ . We assume the chondrule starts with chondritic elemental abundances  $[\text{Ni}] = 10500\text{ppm}$ ,  $[\text{Co}] = 500\text{ppm}$ ,  $[\text{Fe}] = 181000\text{ppm}$  and crystallises only olivine and sulfide. Since sulfide has a strong influence on partitioning, we show cases with masses of (modally) crystallising sulfide at the low (0.2%) and high (1%) ends of the ranges of troilite abundances in chondrules reported by Jones (1990). Tick marks represent 0.12, 0.14, 0.16 and 0.18 solid fractionation for the high sulfide case and 0.35, 0.4, 0.45 and 0.5 for the low sulfide case. Also shown are two points at the high Fe/Ni end of isochrons from two chondrule analysed by SIMS (Mishra and Chaussidon 2014). The values of these two points on the y-axis correspond to the  $^{59}\text{Co}/^{60}\text{Ni}$  required to explain their associated  $^{60}\text{Ni}/^{58}\text{Ni}$  anomaly by  $^{59}\text{CoH}^+$  interference alone. This value is calculated assuming  $^{59}\text{CoH}/^{59}\text{Co}$  production is the same as  $^{56}\text{FeH}/^{56}\text{Fe}$  reported in Mishra and Chaussidon (2014) and the  $^{59}\text{Co}/^{60}\text{Ni}$  calculated at requisite  $^{56}\text{Fe}/^{58}\text{Ni}$ . The range reported reflects the different  $^{59}\text{Co}/^{60}\text{Ni}$  calculated using the two values for  $^{56}\text{FeH}/^{56}\text{Fe}$  evident in the appendix of Mishra and Chaussidon (2014) and includes the error reported on the  $^{60}\text{Ni}/^{58}\text{Ni}$ . One of the  $^{60}\text{Ni}$  anomalies (Ch12), and much of the other (Ch1), can be entirely explained by  $^{59}\text{CoH}^+$  according to these calculations, although more work is required to constrain better all input values.

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Figure 8. Mass-dependent Ni isotope data,  $\delta^{60/58}\text{Ni}$  for different geochemical reservoirs. Data for (a) bulk meteorites (Cameron et al., 2009; Chernozhkin et al., 2015, 2016; Gueguen et al.,

1445 2013; Moynier et al., 2007; Steele et al., 2011, 2012), (b) terrestrial igneous rocks and Ni ores, formed by sulfide segregation in igneous systems (Cameron et al., 2009; Chernonozhkin et al., 2015; Gueguen et al., 2013; Rati'e et al., 2015; Steele et al., 2011, 2012; Tanimizu and Hirata, 2006). (c) the ocean and major input and output fluxes (Cameron and Vance, 2014; Gall et al., 2013; Gueguen et al., 2013)and (d) biosphere processed materials (Cameron et al., 2009; Estrade et al., 2015; Gueguen et al., 2013; Porter et al., 2014; Rati'e et al., 2015)

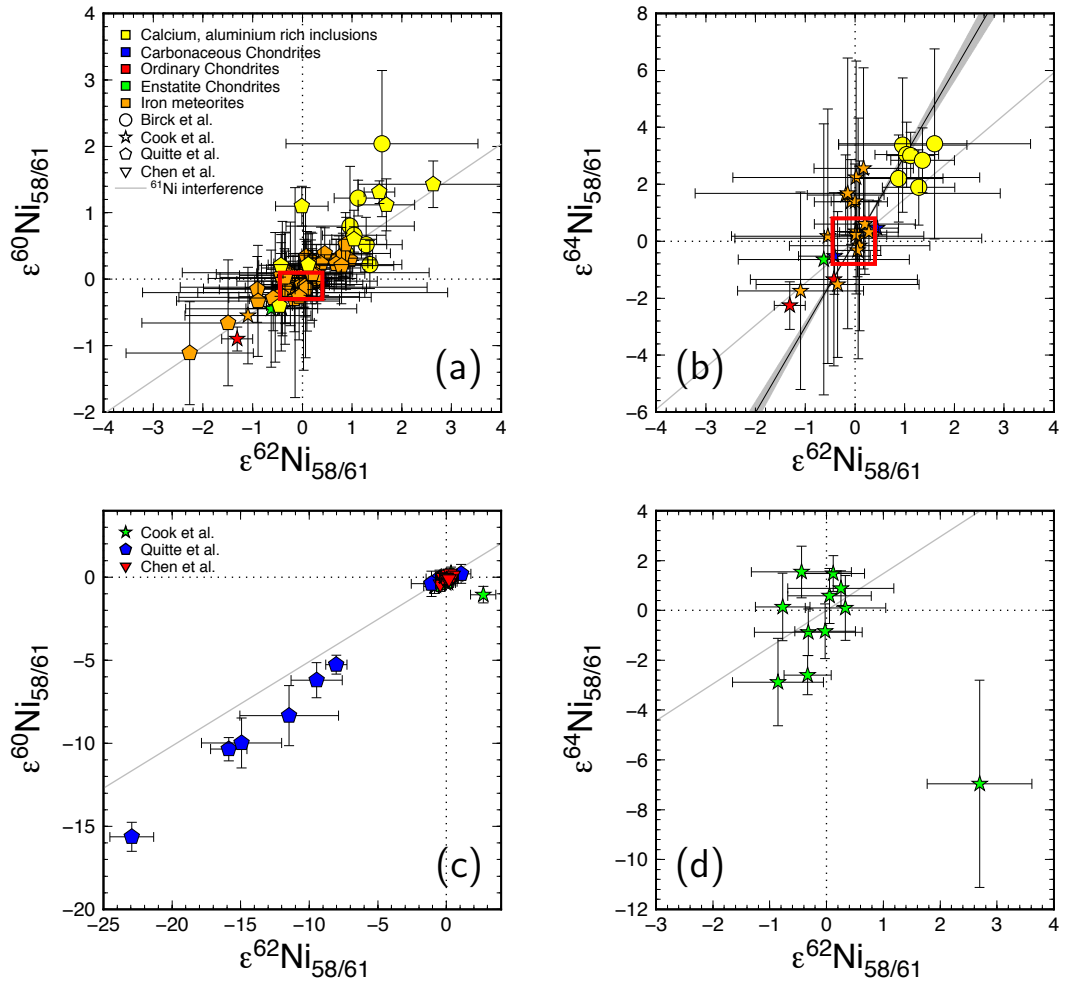


Fig 1

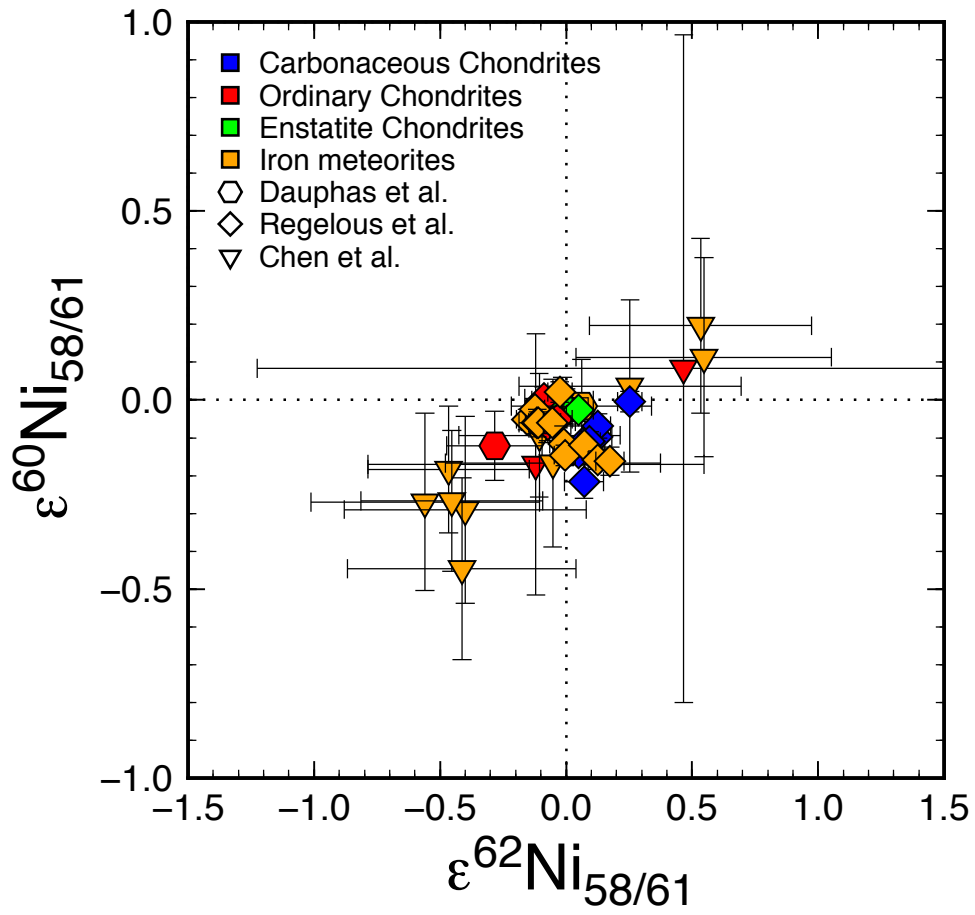


Fig 2

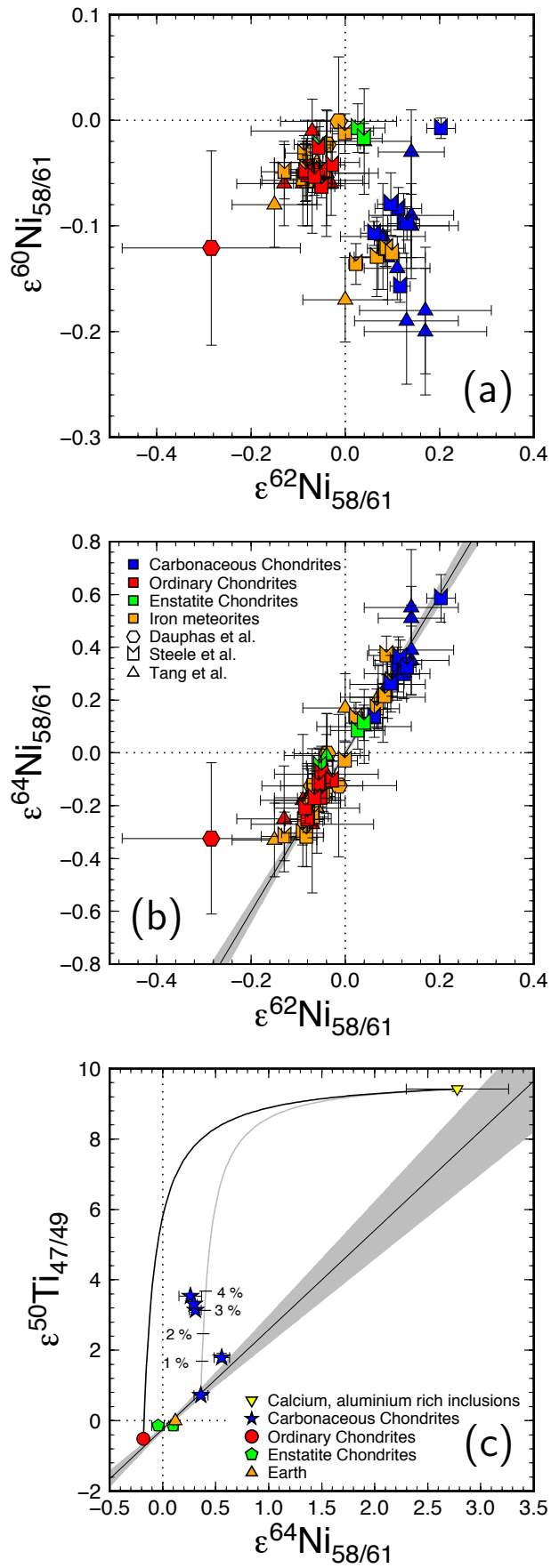


Fig 3

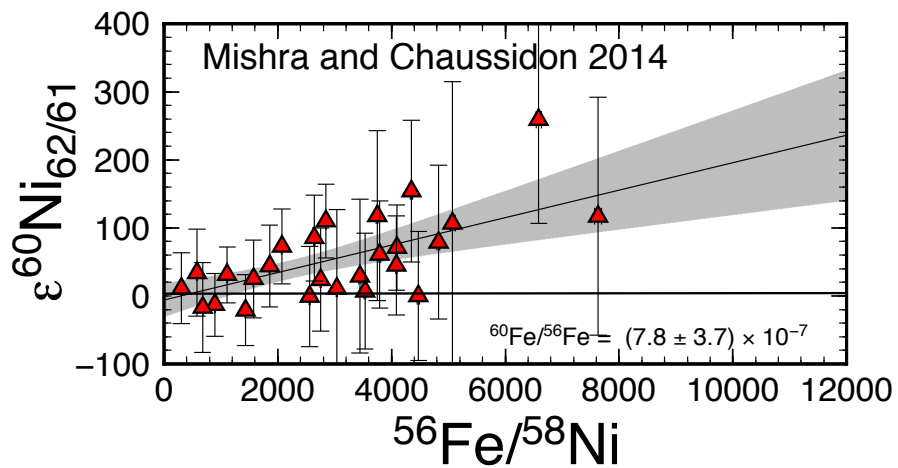
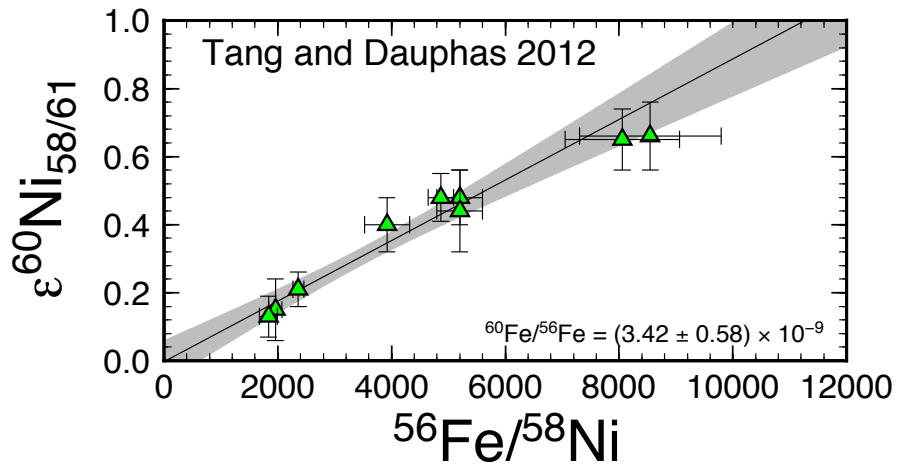
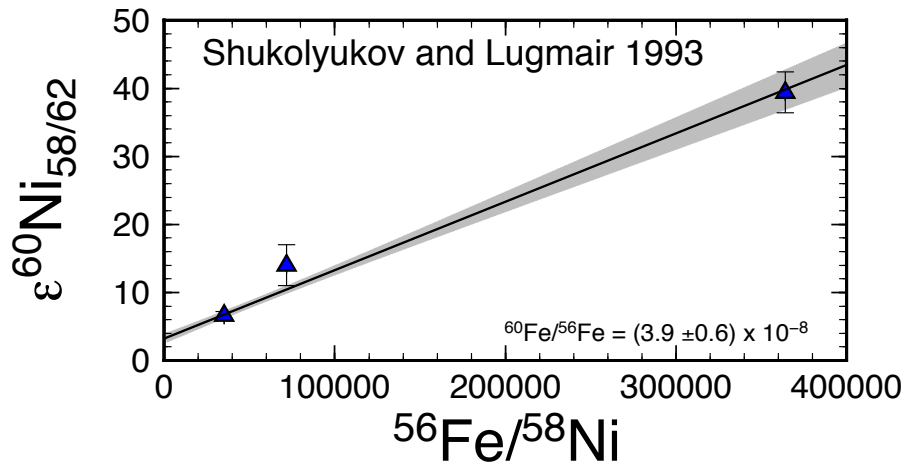


Fig 4



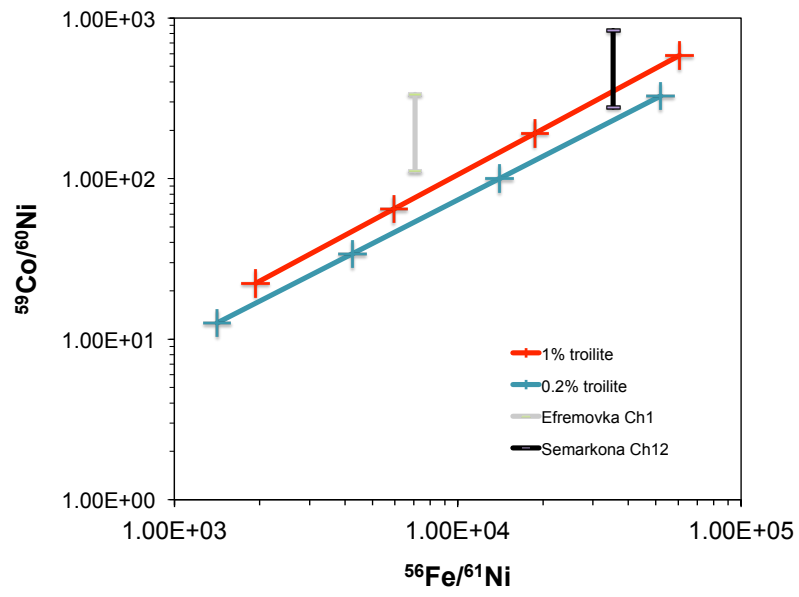
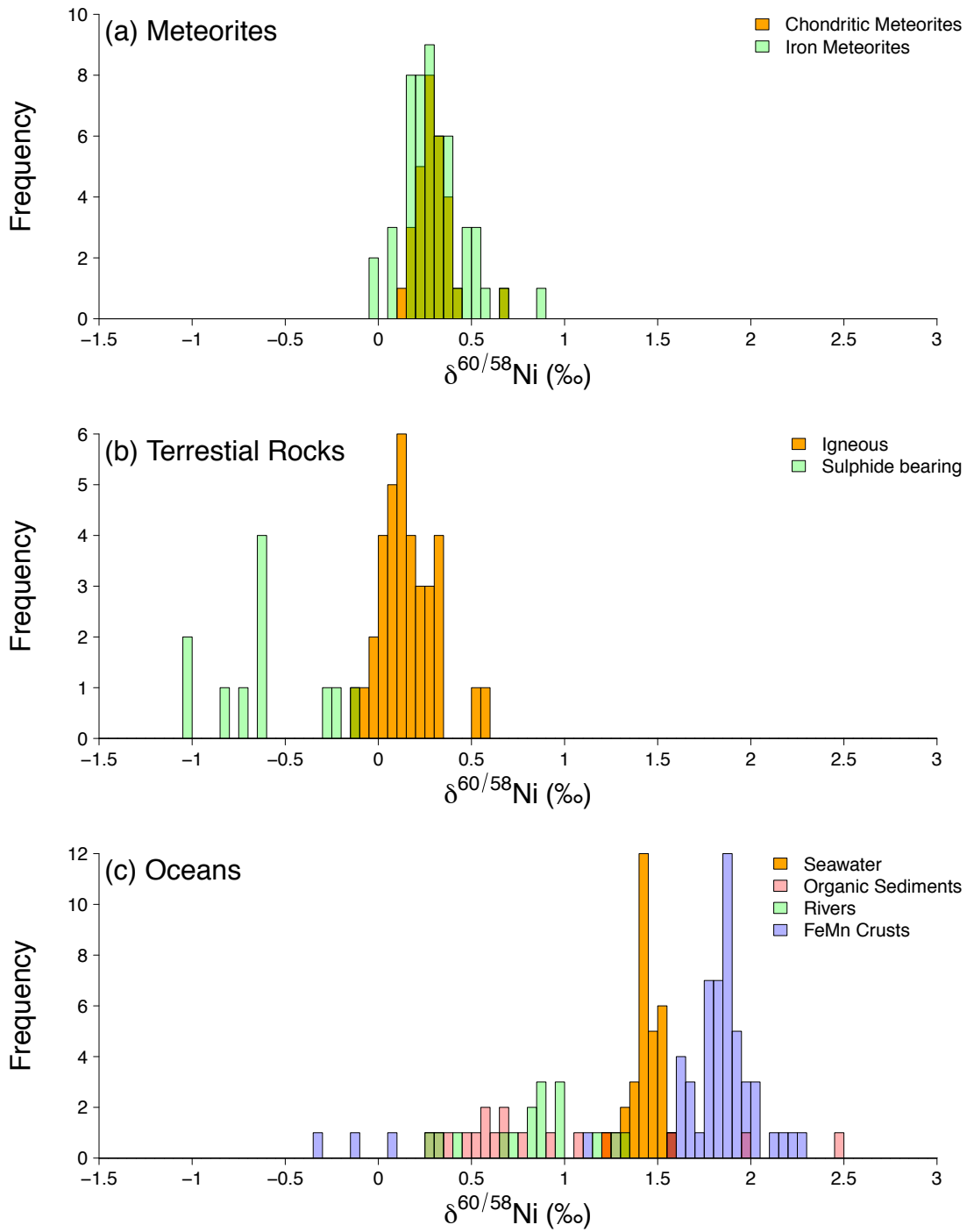


Fig 5



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Fig 6