

Resonance fluorescence from an atomic-quantum-memory compatible single photon source based on GaAs droplet quantum dots

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(Dated: June 19, 2018)

Single photon sources, which are compatible with quantum memories are an important component of quantum networks. In this article, we show optical investigations on isolated GaAs/Al_{0.25}Ga_{0.75}As quantum dots grown via droplet epitaxy, which emit single photons on resonance with the Rb-87-D₂ line (780 nm). Under continuous wave resonant excitation conditions, we observe bright, clean and narrowband resonance fluorescence emission from such a droplet quantum dot. Furthermore, the second-order correlation measurement clearly demonstrates the single photon emission from this resonantly driven transition. Spectrally resolved resonance fluorescence of a similar quantum dot yields a linewidth as narrow as 660 MHz (2.7 μ eV), which corresponds to a coherence time of 0.482 ns. The observed linewidth is the smallest reported so far for strain free GaAs quantum dots grown via the droplet method. We believe that this single photon source can be a prime candidate for applications in optical quantum networks.

PACS numbers: Valid PACS appear here

Keywords: Single photon source, quantum memory, droplet quantum dots, Resonance Fluorescence, droplet epitaxy

Long distance quantum communication based on quantum repeater schemes¹ can be realized by a combination of high-performance single photon sources² and suitable quantum memories.³⁻⁶ For this application single photon sources should possess not only high photon fluxes and very narrow linewidth, but also match spectrally with other components of the system⁷ such as long-lived quantum memories⁸ and which may be utilized in building blocks of quantum repeater nodes.^{9,10} Semiconductor quantum dots (QDs) have been shown to be good candidates for efficient and highly indistinguishable single photon sources.¹¹⁻¹³ Moreover, first prototypes of a hybrid semiconductor QD-atomic interface¹⁴ have been already realized to demonstrate the principle feasibility of this approach. Coupling of atomic clouds of Rubidium (Rb) with QDs allowed for demonstration of compact, tunable and spectrally selective delay lines for single photons.¹⁵ Furthermore, the emission wavelength of solid-state GaAs/AlGaAs QDs were tuned in the spectral range of Rb-87-D₂ lines (780 nm) by introducing strain^{16,17} and even operated under electrical pumping.¹⁸

A crucial resource for applications of single photons in quantum networks¹⁹ and linear optical computing,^{20,21} are high degrees of indistinguishability of the emitted photons. The photons should be identical in color, polarization and their coherence should be Fourier limited.²² However, the quality of single photons emitted from the semiconductor QDs critically depends on excitation conditions. Highest degrees of indistinguishability of single photons have been observed so far under resonant excitation conditions¹¹⁻¹³ as opposed to non-resonant excitation. Furthermore, continuous wave (CW) resonance fluorescence provided a paradigm approach to demonstrate high degrees of indistinguishability and measure single QD coherence time by probing linewidth of a QD transition, reaching close to lifetime limited values.^{23,24}

In this work, we focus on a resonance fluorescence measurements of single GaAs/Al_{0.25}Ga_{0.75}As QDs grown by droplet epitaxy. We show that the energy of emitted single photons from selected QDs match the Rb-87-D₂ lines (780 nm).²⁵ Furthermore, the resonance fluorescence measurements yield a spectral linewidth of only 660 MHz corresponding to a coherence time (τ_{coh}) of 0.482 ns (using the relation, $linewidth = \frac{1}{\pi\tau_{coh}}$). These measurements strongly outline the feasibility to implement droplet epitaxy grown QDs as single photon sources in

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quantum information schemes.

Our sample consists of a low-density layer of GaAs/ $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ QDs grown via droplet epitaxy. The layer of the QDs is embedded in a Schottky-diode structure. The sample was grown by solid-source molecular beam epitaxy (MBE) on a GaAs(100) semi-insulating a substrate.^{26,27} Ga and Al were supplied by standard effusion cells while As_4 molecules were provided by a valved cracker source. After the growth of a GaAs buffer layer, a 50 nm Si-doped GaAs layers and a 150 nm thick, Si-doped $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ layer were deposited on the substrate at the temperature of 550 °C. The doping concentration was kept constant to be $3 \times 10^{18} \text{ cm}^{-3}$. Further, a 20 nm undoped $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ tunnel barrier was grown below the QDs at substrate temperature of 350 °C and the As valve was closed to reduce the background As pressure. Moreover, Ga was deposited at a flow rate of 0.071 monolayer/s for 2 minutes and crystallized with As at 3×10^{-5} Torr beam equivalent pressure. Followed by a rise of substrate temperature to 400 °C, the GaAs QDs were *in-situ* annealed for 10 min to cure crystal defects. Next, the growth was continued with the AlGaAs QD capping layer. The overgrowth of QDs with 10 nm $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ was done by migration enhanced epitaxy,²⁷ followed by raising the substrate temperature to 550 °C and 10 periods of a 2 nm AlAs/2 nm $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ superlattice. The sample growth was completed with 60 nm $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ and a 15 nm GaAs cap layer. The areal density of the QDs was found to be $5 \times 10^9 \text{ cm}^{-2}$.

For μ photoluminescence (PL) and resonance fluorescence (RF) measurements the sample was kept in a closed-cycle cryostat at 5K. Non-resonant μ PL measurements were done using CW 532 nm laser, while resonantly excited by a 780.076 nm laser. The emitted light from the sample was collected into a spectrometer (1200 lines, grating resolution 0.01 nm) by a 50X objective (NA=0.42).

Figure 1b depicts a photoluminescence (PL) spectrum of a droplet GaAs/ $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ QDs, which was excited under non-resonant excitation (532 nm). Apart from a distinct QD line at 780 nm, we observe several other QD-attributed emission features, which most likely stem from other QDs excited within laser spot size. The diameter of the laser spot is $\sim 0.9 \mu\text{m}$ for 532 nm and we used 50X magnification objective. Considering the areal density of QDs which is $\sim 5 \times 10^9 \text{ cm}^{-2}$ we estimate that about 35 quantum dots might be excited. In order to single out the emission line of interest, we resonantly drive the QD with the fundamental transition at 780 nm with a narrow-band CW laser. As shown in the spectrum in Fig. 1(c), this method clearly yields a purely monochromatic emission by selectively driving the ground state of the QD. The excitation scheme and energy levels for a resonant excitation is shown in Fig. 1(a).

More importantly, the resonance fluorescence spectrum shows that the emitter of interest can be coherently driven at a frequency which overlaps with the Rb-87- D_2 lines (780 nm). Further, it is also seen that the emitter

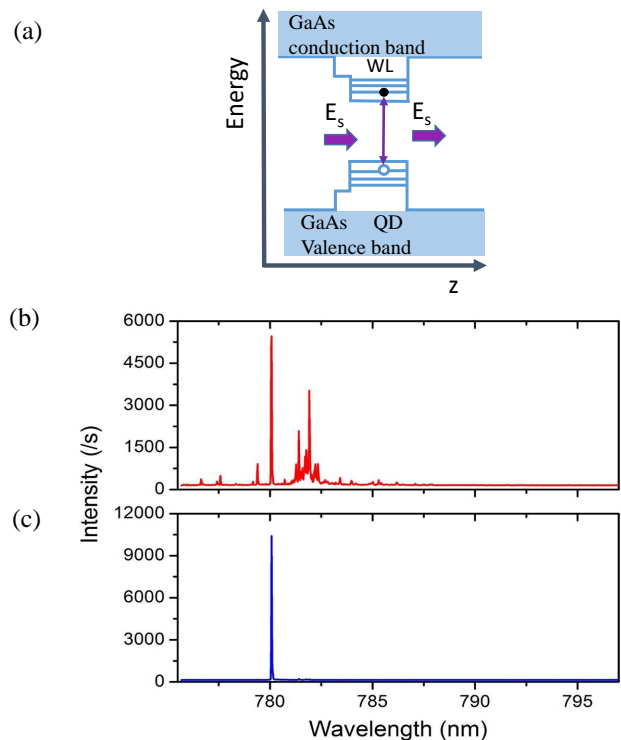


Figure 1. a) A schematic drawing for the resonant excitation (E_s) scheme of the QD. b) Photoluminescence spectrum of a droplet GaAs/ $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ QD under non-resonant excitation at 532 nm with excitation power of $2 \mu\text{W}$. c) Resonant photoluminescence spectrum, yielding a single transition line at 780 nm.

under resonance fluorescence is twice brighter than the one excited under non-resonant scheme under similar excitation power ($2 \mu\text{W}$).

Our method to acquire a pure, nearly background-free resonant emission spectrum from a solid-state quantum emitter is shown in Fig. 2(a). The driving laser is polarized by a combination of a linear polarizer and a half-wave plate, and is focused on the sample surface by a high NA objective. The resonance fluorescence signal from the QD is then collected by the same objective. A second perpendicular linear polarizer filters out the laser signal. Spatial filtering is achieved by focusing the collected RF signal into a single mode fiber, before the signal is dispersed in a monochromator and recorded on a CCD camera. In order to accurately deduce the linewidth of our droplet QDs, we performed resonant laser scans under CW excitation [Fig. 2(b)]. RF counts were recorded as a function of the laser frequency. There, we chose a pump power of $0.1 \mu\text{W}$, well below the saturation of the QD transition and an integration time of 20 s. Two resonance fluorescence peaks evolve in the spectrum, which correspond to the fine structure splitting (FSS) of 1.25 GHz of the QD emitter. The FSS is relatively small, due to a greater degree of structural symmetry^{28,29} in droplet quantum dots than a Stranski-Krastanov QD.³⁰⁻³² For example, Huo *et al.*³⁰ reported FSS values ranging from

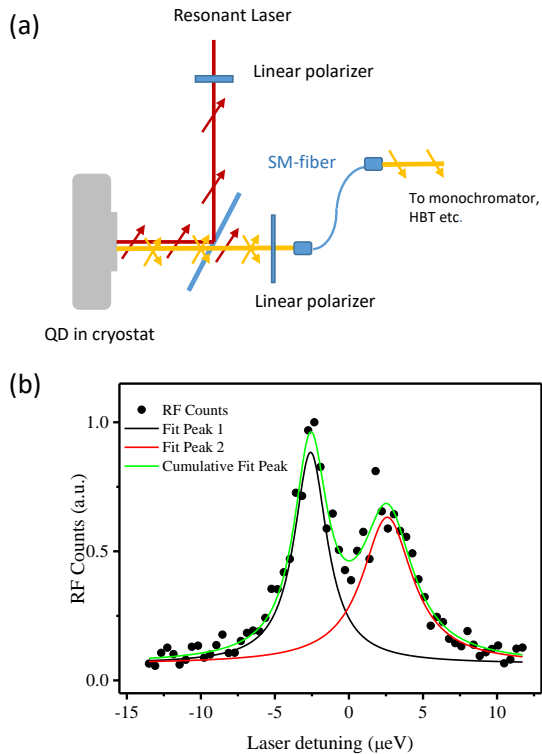


Figure 2. a) A scheme of cross-polarization resonance fluorescence setup. A laser beam (CW 780.076 nm, resonant to the ground state transition of the QD) is passed through a linear polarizer and a microscope objective to excite the QD. The laser beam reflected from the sample surface is then filtered by a second linear polarizer oriented perpendicular to the first one. The resonance fluorescence signal emitted from the QD is then collected by the same objective and coupled to a single mode (SM) fiber working as a spatial filter. Finally the optical signal is introduced to the monochromator equipped with a charged coupled camera detector and Hanbury Brown-Twiss (HBT) interferometer. b) The resonance fluorescence spectrum recorded by a laser scan. Solid lines present Lorentzian fits showing two peaks originating from the same QD with a fine structure splitting of about 1.25 GHz. The peak 1 has a linewidth of (0.66 ± 0.05) GHz and peak 2 has a linewidth of (1.00 ± 0.09) GHz. The signal to noise ratio is $\sim 11:1$.

2.4 GHz to 0.85 GHz.

In order to confirm the quantum nature of our narrow-band emission feature, we performed power dependent measurements shown in Fig. 3(a). The excitation power dependent PL intensity (Fig. 3a) is fitted with the equation for a two-level system,³³

$$I = I_{sat} / (1 + P_n / P_{exc}), \quad (1)$$

where, I_{sat} is the saturation intensity, P_{exc} is the excitation power and P_n is the normalization power. A second-order correlation function for a two-level system is given by³⁴

$$g^2(\tau) = 1 - a \times e^{-\frac{|\tau - \tau_0|}{\tau_{c1}}}, \quad (2)$$

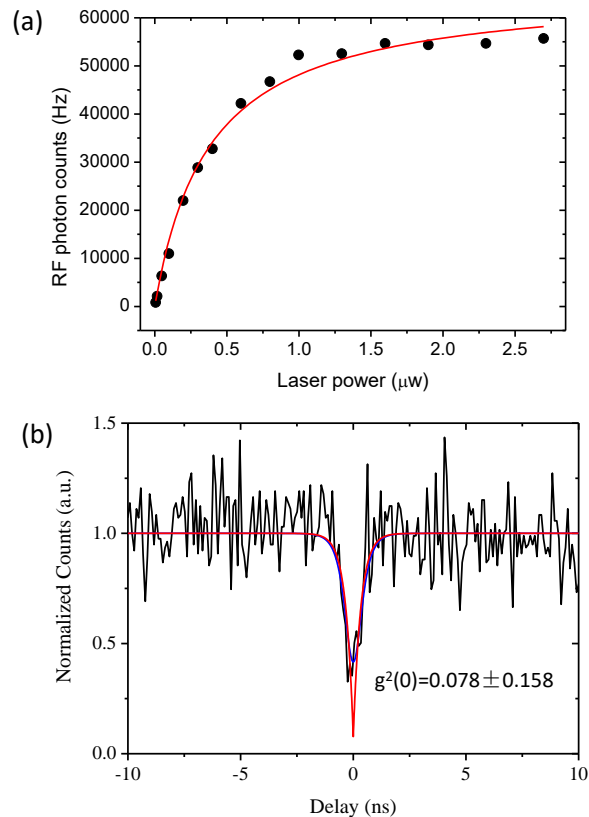


Figure 3. (a) Resonance fluorescence photon (RF) counts as a function of increasing power of the excitation CW laser. The red line is a fit (using Eq. 1), showing a saturation of emission intensity with increase in excitation power. The best-fit normalizations power, $P_n = 0.376 \mu W$. b) Second order correlation function $g^2(\tau)$ as a function of time delay between emitted photons. The blue line is the fit to the intensity autocorrelation measurement data without deconvolution. The red line shows deconvolution fit. The best-fit second-order correlation function $g^2(0)$ at zero time delay was found to be 0.078 ± 0.158 .

here, τ_{c1} accounts for the exciton relaxation and decay rates of a two-level system and a is a fitting parameter. τ , τ_0 are the time delay and the time delay offset, respectively.

The finite-time resolution of our single-photon detector (350 ps) is accounted for by introducing a Gaussian response function G_d given by³⁴

$$G_d = \frac{1}{\sigma\sqrt{\pi}} e^{-\frac{1}{2}(\frac{\tau - \tau_0}{\sigma})^2}. \quad (3)$$

Here, $2\sqrt{2\ln(2\sigma)}$ represent 350 ps timing resolution of our avalanche photo diode. Hence, the convoluted fitting function, $g_{real}^2(\tau)$ is given by

$$g_{real}^2(\tau) = G_d \otimes g^2(\tau). \quad (4)$$

Here, \otimes represent the convolution operator. In Fig. 3 (b) $g_{real}^2(\tau)$ was used to fit the data to obtain the value of

the second-order correlation function at zero delay, $g^2(0)$ to be 0.078 ± 0.158 as shown in Fig. 3(b). The non-zero $g^2(0)$ could be due to presence of a non-filtered laser background.

The Schottky diode geometry used in our study enables us to charge tune the emission and spectrally fine tune the emission line of quantized energy levels of the QDs²⁷. We observed resonance fluorescence in both contacted and non-electrically-contacted geometry, and about 50 % of the investigated QDs showed resonance fluorescence. We note that in contrast to investigations on high quality InGaAs quantum dots in high quality micropillar cavities,^{11,36} the application of additional weak 532 nm laser light on the investigated GaAs/AlGaAs turned out to improve the QD properties similar to other studies, for instance, Gazzano *et al.*³⁵ The measured linewidths of the doublets in the RF spectrum (660 MHz and 1 GHz) are larger than the Fourier-transform limit of 488 MHz for the exciton relaxation time of 330 ps. This difference can be attributed to an inhomogeneous broadening due to charge noise via the d.c. Stark shift.³⁷

In summary, we report a solid state semiconductor single photon device based on isolated charge-tunable droplet $\text{Al}_{0.25}\text{Ga}_{0.75}\text{As}$ QDs. Our resonance fluorescence measurement yield important information about linewidth and fine structure splitting. Most importantly the QD is shown to emit single photons at 780 nm in resonance with Rb-87-D₂ lines (780 nm). Continuous wave laser power dependent fluorescence measurement demonstrate two-level behavior of probed emitter. We suggest, that the QD device may find application in quantum communication devices.

This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 721394. We furthermore gratefully acknowledge support by the State of Bavaria.

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