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SPACE TELESCOPE AND OPTICAL REVERBERATION MAPPING PROJECT.
IV. ANOMALOUS BEHAVIOR OF THE BROAD ULTRAVIOLET EMISSION LINES IN NGC 5548

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ABSTRACT

During an intensive *Hubble Space Telescope* (*HST*) Cosmic Origins Spectrograph (COS) UV monitoring campaign of the Seyfert 1 galaxy NGC 5548 performed from 2014 February to July, the normally highly correlated far-UV continuum and broad emission-line variations decorrelated for ~ 60 –70 days, starting ~ 75 days after the first *HST*/COS observation. Following this anomalous state, the flux and variability of the broad emission lines returned to a more normal state. This transient behavior, characterised by significant deficits in flux and equivalent width of the strong broad UV emission lines, is the first of its kind to be unambiguously identified in an active galactic nucleus reverberation mapping campaign. The largest corresponding emission-line flux deficits occurred for the high-ionization collisionally excited lines, C IV and Si IV(+O IV)], and also He II(+O III)], while the anomaly in Ly α was substantially smaller. This pattern of behavior indicates a depletion in the flux of photons with $E_{\text{ph}} > 54$ eV, relative to those near 13.6 eV. We suggest two plausible mechanisms for the observed behavior: (i) temporary obscuration of the ionizing continuum incident upon BLR clouds by a moving veil of material lying between the inner accretion disk and inner BLR, perhaps resulting from an episodic ejection of material from the disk, or (ii) a temporary change in the intrinsic ionizing continuum spectral energy distribution resulting in a deficit of ionizing photons with energies > 54 eV, possibly due to a transient restructuring of the Comptonizing atmosphere above the disk. Current evidence appears to favor the latter explanation.

Subject headings: galaxies: active — galaxies: individual (NGC 5548) — galaxies: nuclei — galaxies: Seyfert

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1. INTRODUCTION

One of the most secure correlations in studies of active galactic nuclei (AGN) is that found between the ultraviolet (UV)–optical continuum variations and the UV–optical broad emission-line (BEL) variations. This simple relation is causal in nature. The UV–optical broad emission lines arise in high-density, fast-moving gas that is photoionized by an intense source of ionizing contin-

uum radiation originating from a disk of material feeding the central supermassive black hole. Intrinsic variations in the incident ionizing continuum flux translate into correlated changes in the BEL strengths and their ratios. That an AGN’s observed UV/optical continuum variations closely track the largely unobservable driving continuum is demonstrated by the observed positive correlations between the flux variations in the emission lines and the observed UV/optical continuum. Positive correlations are also observed between the UV/optical and X-ray variations (Clavel et al. 1992; Marshall et al. 1997; Peterson et al. 2000; Uttley et al. 2003; Edelson et al. 2015, hereafter Paper II).

This ubiquitous property has proven extremely useful because the correlated continuum–emission-line variations may be used to probe the spatial distribution and kinematics of the BEL gas, probing size scales (roughly a few microarcseconds) inaccessible via more conventional means. This technique, commonly referred to as reverberation mapping (RM; see Blandford & McKee 1982), in its simplest form can reveal the characteristic “size” of the broad emission-line region (BLR) for lines of very different ionization state and arising in gas with a wide range in density, thereby mapping out the density and ionization structure of the BLR. However, its real power manifests when the measured delays between the continuum and emission-line variations are coupled with velocity information. Time-resolved spectroscopy can reveal not only the bulk motion of the line-emitting gas, but has been usefully exploited to measure the mass of the central black hole in ~ 60 AGN (e.g., Denney et al. 2010; Bentz et al. 2010b; Grier et al. 2012; Kaspi et al. 2000, 2007; Peterson et al. 2002; Peterson 2014, and references therein; see also Bentz & Katz 2015 for a comprehensive list of RM sources).

RM is now a mature field. Notable improvements have been made in target selection, campaign design, and observing strategies. Coupled with new techniques developed to isolate the variable broad emission lines from contaminating components and major advances in methods for recovering the emission-line response function, these have facilitated improved dynamical mass estimates for several AGN (Barth et al. 2011; Bentz et al. 2014; Brewer et al. 2011; Denney et al. 2006, 2009a,

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2010; Pancoast et al. 2012, 2014a,b; Zu et al. 2011, 2013a,b), as well as black hole mass determinations for possibly super-Eddington sources (Grier et al. 2012; Du et al. 2014, 2015; Hu et al. 2015; Wang et al. 2014), objects with purported slim disks and exhibiting a different mode of accretion. Significantly, RM campaigns are now producing velocity-delay maps of sufficient fidelity that the gas spatial configuration and dynamics can be discerned (Denney et al. 2009b; Bentz et al. 2010a,b; Skielboe et al. 2015).

In 2014, the AGN Space Telescope and Optical Reverberation Mapping (STORM) collaboration undertook the most intensive RM experiment to date. Awarded 179 orbits in Cycle 22 with *HST*/COS, the aim was to provide the first high-fidelity velocity-resolved delay maps for the strong, broad UV emission lines in an AGN. The chosen target, NGC 5548, has a well-documented history of correlated large-amplitude continuum and BEL variations at both UV and optical wavelengths (Clavel et al. 1991; Korista et al. 1995; Peterson et al. 2002, and references therein). The *HST*/COS observations are presented by De Rosa et al. (2015, hereafter Paper I) and are of an unprecedented quality. These data were supported by photometric observations with *Swift* (Paper II), as well as by ground-based photometric (Fausnaugh et al. 2016, hereafter Paper III) and spectroscopic (Pei et al., in prep.; hereafter Paper V) observations.

This remarkable dataset has revealed a number of interesting and unexpected results. Chief among these is a temporary breakdown in one of the fundamental tenets of RM. Approximately 75 days after the start of the *HST*/COS campaign, the continuum and emission-line variations appeared to decorrelate. This transient phenomenon, characterized by a significant and anomalous depression in the BEL flux and a notable reduction in emission-line variability lasting ~ 65 –70 days, is the subject of this investigation.

2. REVERBERATION MAPPING

The primary goal of RM is to use correlated continuum and BEL variations to solve the transfer equation,

$$L(V, t) = \int_0^\infty \Psi(V, \tau) C(t - \tau) d\tau, \quad (1)$$

where $\Psi(V, \tau)$, the transfer function (or convolution kernel), maps the *driving* continuum light curve at earlier times $C(t - \tau)$ onto the observed emission-line light curve $L(V, t)$ at the current epoch.

To eliminate nuisance background components (e.g., the host-galaxy contribution to the UV/optical continuum bands and contaminating narrow emission-line contributions), this equation has traditionally been solved in its linearized form

$$\Delta L(V, t) = \int_0^\infty \Psi'(V, \tau) \Delta C(t - \tau) d\tau, \quad (2)$$

where $\Delta L(V, t)$ and $\Delta C(t - \tau)$ respectively represent the variable component of $L(V, t)$ and $C(t - \tau)$ relative to some average background value. Here, the quantity of interest is $\Psi'(V, \tau)$, the emission-line “response function.” In general, $\Psi'(V, \tau) \neq \Psi(V, \tau)$.

Underpinning this simple relation, and a prerequisite for a successful RM campaign, is that the continuum

and emission-line variations are causally related. More specifically the basic assumptions of RM are as follows.

1. At the distance of the BLR, the driving continuum source appears point-like.
2. The ionizing continuum and emission-line photons propagate freely at the speed of light — i.e., the BLR filling factor is small.
3. There exists a simple causal (though not necessarily linear) relation between the observed continuum and emission-line variations.
4. The observed continuum variations, modulo some scale factor and temporal smoothing function, are a suitable proxy for the unobserved driving ionizing continuum variations.
5. The dominant timescale is the light-crossing time.

Violation of any of these assumptions will impact our ability to recover $\Psi'(V, \tau)$. Fortunately, in *the vast majority of previous RM experiments*, these assumptions are broadly consistent with the data.

3. FOLLOWING THE ENERGY

3.1. Anomalous Broad Emission-Line Variations

The strongest evidence to date for the violation of one of the above assumptions (specifically, assumption 4) appeared during the middle of the AGN STORM campaign, upon comparison of the UV continuum and BEL light curves (Paper I)⁶⁸. While the UV continuum and BEL variations were well-correlated at the start and end of the *HST* observing campaign, they appeared to decorrelate starting approximately midway through the campaign. It was for this reason that in Paper I a separate analysis was performed on the data for the two halves of the campaign. We illustrate this in Figure 1c, where we compare a modified version of the C IV BEL light curve (colored points) with the continuum light curve at 1157 Å. We use emission-line light curves based on the *direct integrations* over the bulk of the broad emission line reported in Paper I, and modified as described below.

For C IV and all other emission lines, we have *shifted the entire emission-line light curve* according to the measured delay between the emission line and the 1157 Å continuum band as determined for *the first 75 days* of the *HST* campaign. We assume that this delay provides an adequate description of the delay between the continuum and emission-line variations for the latter half of the campaign. For the measured delays we use the peak and centroid of the cross-correlation function (hereafter CCF, see Table 1). The CCF centroids are calculated over the range in delay for which the CCF coefficient exceeds 80% of the peak value. Uncertainties are estimated using the model-independent FR/RSS Monte Carlo method (Peterson et al. 1998), computing 1000 realizations of each light curve, assuming random sampling and full replacement. Next, we remove an estimate for contaminating

⁶⁸ The only anomalies reported to date have been attributed solely to changes in the continuum reprocessing efficiency of BLR clouds (e.g., Maoz 1992; Sparke 1993; Grier et al. 2008). In these cases, while the line responsivity changed, the emission-line and continuum variations remain significantly correlated.

narrow emission-line flux (Table 1, column 4; based on measurements of low flux-state spectra of NGC 5548 first reported by Crenshaw et al. 1993; see also Goad & Koratkar 1998), and then rescale the emission-line light curve to the mean continuum flux over the first 75 days of the *HST* campaign. Finally, we scale the emission-line variability amplitude by dividing through by the measured responsivity for each line (§3.3 and Table 2, column 2). In so doing, the continuum and emission-line variations should then be of similar amplitude (they will be identical if the response function were simply a delta function shifted in time).

Figures 1a–1c indicate that starting from day 75 (cyan points), the C IV flux relative to our proxy measure of the ionizing continuum flux drops significantly and then remains anomalously low for ~ 40 days (red points), before recovering (magenta points) and again showing normal, correlated behavior (green points). Additionally, during the anomalous period (red points), the emission-line variability amplitude, relative to the observed continuum variations, is also significantly suppressed. Similar behavior, though of varying strength, is seen *for all of the strong, broad UV emission lines* (§3.4), and is indicative of a dramatic decline in the continuum reprocessing efficiency of the line-emitting gas, as inferred by the observed continuum at 1157 Å. We have examined the 1989 *IUE*/SWP,LWP (Clavel et al. 1991) and 1993 *HST*/FOS (Korista et al. 1995) monitoring campaigns on NGC 5548 to look for similar effects and found none. To our knowledge, the AGN STORM campaign is the first for which a decorrelation between the observed UV continuum and UV BEL variations of such long duration has been reported.

3.2. BEL Responsivity and Reprocessing Efficiency

Significant changes in the continuum reprocessing efficiency of an emission line are best revealed through studying the emission-line responsivity and/or equivalent width (EW). The time-averaged emission-line responsivity η_{eff} is the power-law index relating the observed continuum, preferably the band closest in wavelength to the ionizing continuum, and BEL fluxes, F_{cont} and F_{line} :

$$F_{\text{line}} \propto F_{\text{cont}}^{\eta_{\text{eff}}}. \quad (3)$$

The exponent η_{eff} is normally measured after first removing contaminating contributions from nonvariable background components (e.g., the narrow emission lines and starlight from the host galaxy), and after correcting for the mean delay between the continuum and emission-line variations (Pogge & Peterson 1992; Gilbert & Peterson 2003; Goad, Korista & Knigge 2004), in order to minimize the effects of geometrical dilution. Alternatively, the emission-line EW and the continuum may be related by

$$\text{EW}_{\text{line}} \propto F_{\text{cont}}^{\beta}, \quad (4)$$

where $\beta = \eta_{\text{eff}} - 1$. The relationship between line responsivity and emission-line EW is thus made clear. For a strictly linear (1:1) response, $\eta_{\text{eff}} = 1$, $\beta = 0$, and the emission-line EW remains constant relative to the continuum flux variations. If $\eta_{\text{eff}} = 0$, the emission line does not respond to variations in the driving continuum. In general, $\eta_{\text{eff}} < 1$ for most emission lines (i.e., a nonlinear

response), and hence $\beta < 0$. Thus, the majority of emission lines will show an intrinsic Baldwin effect (Kinney et al. 1990). The measured value of η_{eff} will also depend upon the reference continuum band (since the amplitudes of the continuum variations are larger at shorter wavelengths — e.g., Wamsteker et al. 1990; Clavel et al. 1991; Papers II & III) and the degree of geometrical dilution (Gilbert & Peterson 2003; Goad & Korista 2014, 2015).

3.3. Measuring η_{eff} and Identifying the Anomaly

For the continuum light curve against which the emission-line delays and their responsivities will be measured, we use the *HST* continuum-band measurements centered at 1157 Å (Paper III), as this is the wavelength band accessible with *HST*/COS that lies nearest to the driving ionizing continuum and has negligible host-galaxy contamination⁶⁹. For illustrative purposes, we focus on the strong, broad C IV emission line, but we treat all other emission lines in the same fashion. To measure η_{eff} we follow the procedure of Goad, Korista, & Knigge (2004). Using the delay-corrected BEL light curve we reconstruct the continuum flux associated with the BEL flux at the current epoch by interpolation, adopting a weighted average of the two bracketing continuum points. The appropriate weights are derived from the first-order structure function of the continuum light curve (Kawaguchi et al. 1998; Paltani 1999; Goad, Korista, & Knigge 2004). Continuum data associated with emission-line data shifted to epochs before the start of the campaign are determined by a linear extrapolation. Emission-line points beyond the end of the campaign can be determined in a similar fashion. We do not use these extrapolated points (either emission line or continuum) when measuring the time-averaged BEL responsivity η_{eff} , preferring instead to exclude these data from the analysis. Uncertainties in the reconstructed continuum and emission-line points were determined from a structure-function analysis of their respective light curves (for details, see Goad, Korista, & Knigge 2004).

Having corrected the data for the mean delay $\langle \tau \rangle$, we then fit the relation

$$\log F_{\text{line}} = A + \eta_{\text{eff}} [\log F_{\text{cont}}(1157\text{Å})] \quad (5)$$

to points lying outside of the anomaly (blue and green points only), assuming uncertainties in both the ordinate and abscissa. For C IV, the best-fit slope gives $\eta_{\text{eff}} = 0.25 \pm 0.01$ (Figure 1a, solid red line). In Table 2, column 2, we report the best-fit slopes for the time-averaged responsivity η_{eff} and their 1σ uncertainties for all of the broad emission lines reported in Paper I. Quoted values are the centroid and 1σ uncertainties determined from bootstrap resampling of the (corrected) light curves (10,000 realizations) with full replacement.

Figure 1a indicates that during the anomaly (epochs 75–140) there is a dramatic and uncorrelated reduction in the C IV emission-line flux and response amplitude to continuum variations. Note in particular the island

⁶⁹ Relative to the driving ionizing continuum variations, the observed UV and optical continuum variations are of lower amplitude and smeared in time (e.g., Papers II & III), and may at some wavelengths be significantly contaminated by more slowly varying background components.

of red points well below the main $\log F_{\text{line}} - \log F_{\text{cont}}$ relation with $\eta_{\text{eff}} \approx 0$, indicating either (i) an absence of emission-line response to continuum variations, or (ii) that the ionizing continuum responsible for driving the observed BEL variations varies less than the 1157 Å continuum. Also, note that the journey away from the main relation (cyan points) into the anomaly forms a continuous path (as indicated by the cyan arrows). Curiously, when exiting the anomaly (magenta points), the data follow a similar path but in the reverse direction (magenta arrows), before rejoining the normal $F_{\text{line}} - F_{\text{cont}}$ behavior⁷⁰. There are indications that the $\log F_{\text{line}} - \log F_{\text{cont}}$ relation is somewhat steeper at the end of the campaign (green points) than during the first 75 days (blue points), though some of this effect may result from the use of a single lag when correcting for the continuum–emission-line delays. Larger emission-line responsivities (η_{eff}) are usually associated with the emergence from lower ionizing continuum flux states (Korista & Goad 2004; Goad & Korista 2014, 2015).

Figure 1b illustrates the relationship between the 1157 Å continuum flux and C IV equivalent width, referenced to the 1157 Å continuum. The best-fit slope in this relation (fitting to data on either side of the anomaly; i.e., blue and green points only) is $\beta = -0.75 \pm 0.01$, as expected since $\beta = \eta_{\text{eff}} - 1$. Once again, the color-coded behavior described above is evident. There is an island of red points lying below the main relation and for which $\beta \approx -1$, while the cyan and magenta points taken together appear as extensions of the same relation with a slope intermediate between that of the main relation and that of the anomaly. The same general behavior, though differing in detail and with differing noise levels, is seen for Ly α , Si IV(+O IV)], and He II(+O III)].

3.4. Measuring the Amplitude of the Anomaly

In previous AGN RM campaigns, the strong positive correlation observed between the UV/optical continuum and BEL variations provides substantive supporting evidence that the UV and optical continuum variations are reasonable proxies, though smaller in amplitude and smeared in time, for the variable driving ionizing continuum. The strong positive correlation observed between the continuum and emission-line variations pre- and post-anomaly for this campaign suggests that over these time periods, the UV/optical continuum variations are also suitable proxies for the variable driving ionizing continuum. If this interpretation is correct, these data may then be used to infer the “expected” form of the emission-line light curve during the anomaly — that is, as if the anomaly had not occurred.

Using the β values reported in Table 2, column 3, we reconstruct the emission-line flux corresponding to the observed continuum flux at 1157 Å and shifted in time according to the delay reported in Table 1, column 2⁷¹. In Figure 1d we display a comparison between the observed (colored points) and reconstructed (black trian-

gles) emission-line light curve for C IV. In Figures 2a–2d we show the observed (colored points) and reconstructed (black triangles) emission-line light curves for all of the strong UV emission lines reported in Paper I, here normalized to their mean values over the first 75 days of the campaign. There are indications (e.g., Figures 2a–2d) that the start times, stop times, and duration of the anomaly differs among the emission lines. If real, timing differences among the various lines may provide additional clues about the BLR structure. We defer investigation of this effect to a future paper.

Using the observed and reconstructed emission-line light curves, we can measure the fraction of line emission lost f_{lost} during the period of the anomaly, $f_{\text{lost}} = [(f_{\text{rec}} - f_{\text{obs}})/f_{\text{rec}}] \times 100\%$ (e.g., Figure 1e), where f_{obs} and f_{rec} represent the observed and reconstructed emission-line fluxes, respectively. We find time-averaged fractional losses of $\sim 23\%$ (Si IV(+O IV)], $\sim 21\%$ (He II(+O III)], $\sim 18\%$ (C IV), and $\sim 9\%$ (Ly α) during the time period spanning the anomaly. The largest flux deficit is found for Si IV and He II, followed closely by C IV, while Ly α shows the smallest (by a factor 2) flux deficit (Figures 2a–2d). In Table 3, we report the measured time-averaged emission-line EWs over the period spanning the anomaly (Figure 1a – red points) and compare it with the “expected” EW, here defined as the EW measured from the main relation (solid red line in Figure 1b) that corresponds to the average continuum level measured over the duration of the anomalous period [$F_{\text{cont}}(1157\text{Å}) = (4.84 \pm 0.37) \times 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Å}^{-1}$]. The measured deficit in EW is largest for Si IV(+O IV)] ($\sim 24\%$), followed by He II(+O III)] ($\sim 21\%$), C IV ($\sim 19\%$), and then Ly α ($\sim 11\%$), in good agreement with the estimates given above.

The analyses described here and illustrated in Figures 1 and 2 can be further extended to include the strong optical recombination lines (Paper V, in prep.), and it may also be applied as a function of projected line-of-sight velocity for all of the strong, broad UV and optical emission lines. In particular, it may be possible to see the effects of the transition into, through, and out of the anomaly in the velocity-resolved data, thereby constraining the BLR geometry and dynamics. This will require the highest-fidelity light curves, free from contaminating narrow absorption and emission lines. These are typically obtained from spectral decomposition of the emission lines and is work in progress. Finally, we note that any variability in the equivalent widths of the weak absorption lines associated with the outflow is unrelated to the anomaly, as clearly indicated by the fact that He II(+O III)] is strongly affected by the anomaly but has no absorption lines.

4. DISCUSSION

The observed behavior of the broad emission lines during the middle of the *HST*/COS campaign is consistent with a temporary yet significant softening of the ionizing continuum spectral energy distribution (SED) *incident* upon the BLR lasting ~ 40 days (Figure 1, red points). The cyan and magenta points may represent transition phases (each lasting ~ 1 –2 weeks) straddling the anomalous period. In the following, we suggest two plausible scenarios to account for the anomalous behavior of the broad emission lines.

⁷⁰ An animated GIF of this behavior is available in the online material.

⁷¹ A more thorough treatment would instead solve for the transfer function using those epochs bracketing the anomaly, and then apply forward modeling to infer the expected emission-line fluxes during the anomalous period. This will be explored in the future but is beyond the scope of the current work.

In Figure 2e, we show the light curve for a powerful diagnostic BEL flux ratio involving the two strongest emission lines, C IV/Ly α (e.g., Shields et al. 1995; Netzer 1990). This temperature-sensitive ratio is a measure of the heating-cooling balance within the BEL clouds, centered on the average ionizing photon energy of ~ 100 eV, estimated for the central source in NGC 5548 (Medhipour et al. 2015). This flux ratio is thus particularly sensitive to the SED of the incident ionizing continuum⁷². During the first ~ 80 days, the C IV/Ly α flux ratio varies (weakly) inversely with the continuum variability, owing to the somewhat larger responsivity of Ly α (Table 2, column 2). Approximately coincident with the onset of the anomaly, this emission-line flux ratio then exhibits a gradual decline, and continues to decline (on average) for the duration of the anomalous period, before eventually rising again as the source behavior returns to normal.

We also show in Figure 2f the light curve for the BEL flux ratio, He II(+O III)/Ly α . This flux ratio is sensitive to the ratio of photons responsible for ionizing the He⁺ ion ($E_{\text{ph}} > 54.4$ eV) to those responsible for ionizing hydrogen ($E_{\text{ph}} > 13.6$ eV). During the first ~ 87 days, this flux ratio correlates directly with the observed continuum variations, owing to the substantially larger responsivity in He II when compared to Ly α . It then quickly drops and remains approximately constant throughout the anomalous period. This behavior is indicative of a reduction in the incident flux of photons capable of ionizing He⁺ relative to the flux of photons responsible for ionizing hydrogen over the time period of anomalous emission-line behavior. We note that the time-dependent behavior of this pair of BEL flux ratios is also reflected in the comparison of panels a, b, c of Figure 2.

Contemporaneous observations in the soft (0.3–0.8 keV) and hard (0.8–10 keV) X-ray bands with *Swift*/XRT (Figure 1 in Paper II) indicate a gradual decline in the soft and hard X-ray photon fluxes starting around day 75 (the onset of the anomaly), beyond which the measured flux in the observed UV continuum bands is trending upward. In Figure 2g, we show the hardness ratio, $\text{HR} = (\text{H}-\text{S})/(\text{H}+\text{S})$, constructed from the soft (S) and hard (H) X-ray bands as defined in Paper II (see also Medhipour et al. 2016, in press). Figure 2g indicates that on entering the anomalous period, the X-ray HR initially increases, thereafter remaining approximately constant with an average value similar to that measured at the start and end of the campaign. That is, there appears to be nothing unusual about the HR during the anomalous period. In general, the HR shows a marginal trend downward over the course of the campaign interrupted by two significant drops (indicative of a softer spectrum), starting ~ 50 days prior to the onset of the anomaly and lasting a combined total of ~ 60 days before returning to the more general trend. It is unclear whether these large excursions in HR prior to the anomaly are in any way related to the anomaly. Significantly, we see no color-dependent changes in the UV–optical continuum band during the time of the anomaly.

On exiting the anomaly, the soft and hard X-ray pho-

ton fluxes appear to trend upward, coincident with the increasing trend in the C IV/Ly α flux ratio. While some of the observed X-ray variability is likely caused by changes in the absorption of X-ray photons in outflowing material lying along our line of sight (Kaastra et al. 2015), these data are also consistent with a gradual softening of the X-ray–UV continuum, notably in our line of sight.

All of these observed temporal behaviors are consistent with a temporary yet significant softening of the incident ionizing continuum both along our line of sight and toward the BLR. The larger emission-line flux and EW deficits exhibited by the collisionally excited lines (C IV and Si IV(+O IV)) are consistent with a drop in electron temperature that accompanied the softening of the ionizing SED. Together with the larger deficit in the broad He II emission line compared to Ly α , the evidence points to most of the softening occurring for $E_{\text{ph}} > 54$ eV.

We suggest two mechanisms to explain the temporary softening of the incident ionizing continuum.

- (i) A temporary obscuration of the ionizing continuum source incident upon BLR clouds by a moving veil of gas lying between the inner accretion disk and the inner BLR. This could represent an episodic ejection of material from the accretion disk. This veiling gas should be moderately highly ionized and of modest column density to produce the phenomena described above. These details are left for future investigation.
- (ii) A temporary change in the intrinsic ionizing SED leading to a reduction in the number of ionizing photons, particularly above ~ 54 eV. NGC 5548 apparently has a hard ionizing continuum, in which Comptonization plays an important role (Magdziarz et al. 1998; Medhipour et al. 2015). A temporary restructuring of the Comptonizing atmosphere situated above the inner accretion disk could significantly reduce the production of EUV continuum photons.

Since all of the broad emission lines (UV and optical) display a deficit in flux during the anomaly, then if temporary obscuration is the mechanism responsible, this imposes constraints on the nature of the obscuring veil of material and/or the BLR geometry. For example, for the obscuring veil of material to affect all lines simultaneously, it must cover a large solid angle, thus constraining its location to be close to the continuum source. Alternatively, if the BLR is in a flattened configuration, the obscuring veil of material need not be raised to very large scale heights above the disk in order to affect all of the emission lines simultaneously, though such configurations for the BLR may be precluded on energetic grounds. Indeed, it may even be possible to distinguish between these differing scenarios by investigating the variation in the start and end times of the anomalous behavior among different lines.

Evidence for a change in the disk geometry can be found in high-resolution *Chandra* and *XMM-Newton* X-ray spectra of NGC 5548 taken during 2013–2014 (Di Gesu et al. 2015; Medhipour et al. 2016). Based on an observed reduction in the covering fraction of the X-ray obscurer at higher source luminosities, Di Gesu et

⁷² In the absence of changes in the ionizing SED, variations in the delay-corrected emission-line flux ratios are mainly due to intrinsic differences between their individual line responsivities, and any remaining residual reverberation effects (e.g., geometrical dilution).

al. (2015) and Medhipour et al. (2016) argue that the continuum-emitting region producing the soft excess becomes larger as the source brightens. If this brightening coincides with a puffing up of a disk whose outer regions are optically thick at photon energies $E_{\text{ph}} > 54$ eV, then this gas may provide the shielding necessary to explain the observations.

However, while obscuration can plausibly explain the significant drop in flux and EW of the broad emission lines during the anomaly, it does not readily explain their reduced amplitude of variability. While an obscuring screen of material will reduce the flux of ionizing continuum photons incident upon BLR clouds, the fractional variability amplitude of the continuum will remain unchanged. By contrast, physical changes in the Comptonizing atmosphere altering the SED at especially EUV photon energies might also lead naturally to suppressed variability at these energies. The second scenario can thereby explain not only the significant emission-line flux deficits, but also the absence of significant emission-line variations ($\eta_{\text{eff}} \approx 0$) during the anomalous period. Note that our emission-line reconstructions (Figures 2a–2d) indicate that if the ionizing continuum had varied as the 1157 Å continuum did, we would have easily detected the corresponding emission-line variations. In addition, the first scenario might be expected to be accompanied by a significant increase in the X-ray HR, with greater photoabsorption occurring for $E_{\text{ph}} < 1$ keV, while observations indicate that for the majority of the anomalous period the HR remains approximately constant (Figure 2g). On the other hand, a reorganization of the Comptonizing atmosphere could manifest itself differently. Speculating further, the decline in the variability amplitude of the 1157 Å continuum during days ~ 75 –120 resembles a damped oscillation, and could be a manifestation of a redistribution of energy within the Comptonizing atmosphere. A more detailed analysis of the *HST*/COS data and contemporaneous X-ray spectroscopy from *Chandra* and *Swift* (Mathur et al. 2016, in prep.) are currently underway. When combined, such data could potentially distinguish these (and possibly other) scenarios, thereby leading to a deeper understanding of the inner workings of AGN accretion disks.

From the outset of AGN STORM, we had assumed that correlated continuum and emission-line variations were the primary (and perhaps only) means of probing the spatial distribution and kinematics of the BLR and the nature of the central engine. Instead, by good fortune, the presence of the anomaly, far from compromising the current RM campaign, has provided a different and rather unique window into the behavior of the central engine and the processes that drive the continuum and BEL variations.

5. SUMMARY

Analysis of the time-dependent behavior of the BEL reprocessing efficiency has proven to be a powerful tool for investigating the relationship between the ionizing continuum source and line-emitting gas. Here we have used it to reveal the origin of the anomalous behavior exhibited by the strong, broad UV emission lines during the middle of the 2014 *HST*/COS monitoring campaign of the nearby Seyfert 1 galaxy NGC 5548: a significant yet temporary softening of the ionizing continuum inci-

dent upon BLR clouds. By following the energy, we are able to recover the form for the expected emission-line light curve during the anomaly, as though the anomaly had not occurred. In so doing, we identified a potential route for utilizing all of the available data from this campaign in future attempts at recovering the emission-line response function for the strongest broad UV emission lines.

Finally, we note that effects of this nature have been looked for, but not found, in previous UV spectroscopic monitoring campaigns of NGC 5548. At the level at which these effects are present in the 2014 *HST* campaign, they would have easily been detected if present, in both the 1989 and 1993 UV monitoring campaigns on NGC 5548. Nor do we find such behavior in the far more extensive 13 yr optical monitoring data available for NGC 5548 (Grier & Peterson 2003; Goad, Korista, & Knigge 2004). Thus, anomalous behavior of this kind is likely rare in this source. Since analyses similar to that presented here have yet to be performed on the archival UV and optical datasets available for other AGN, we are unable to comment on the occurrence rate of this phenomenon among the RM AGN.

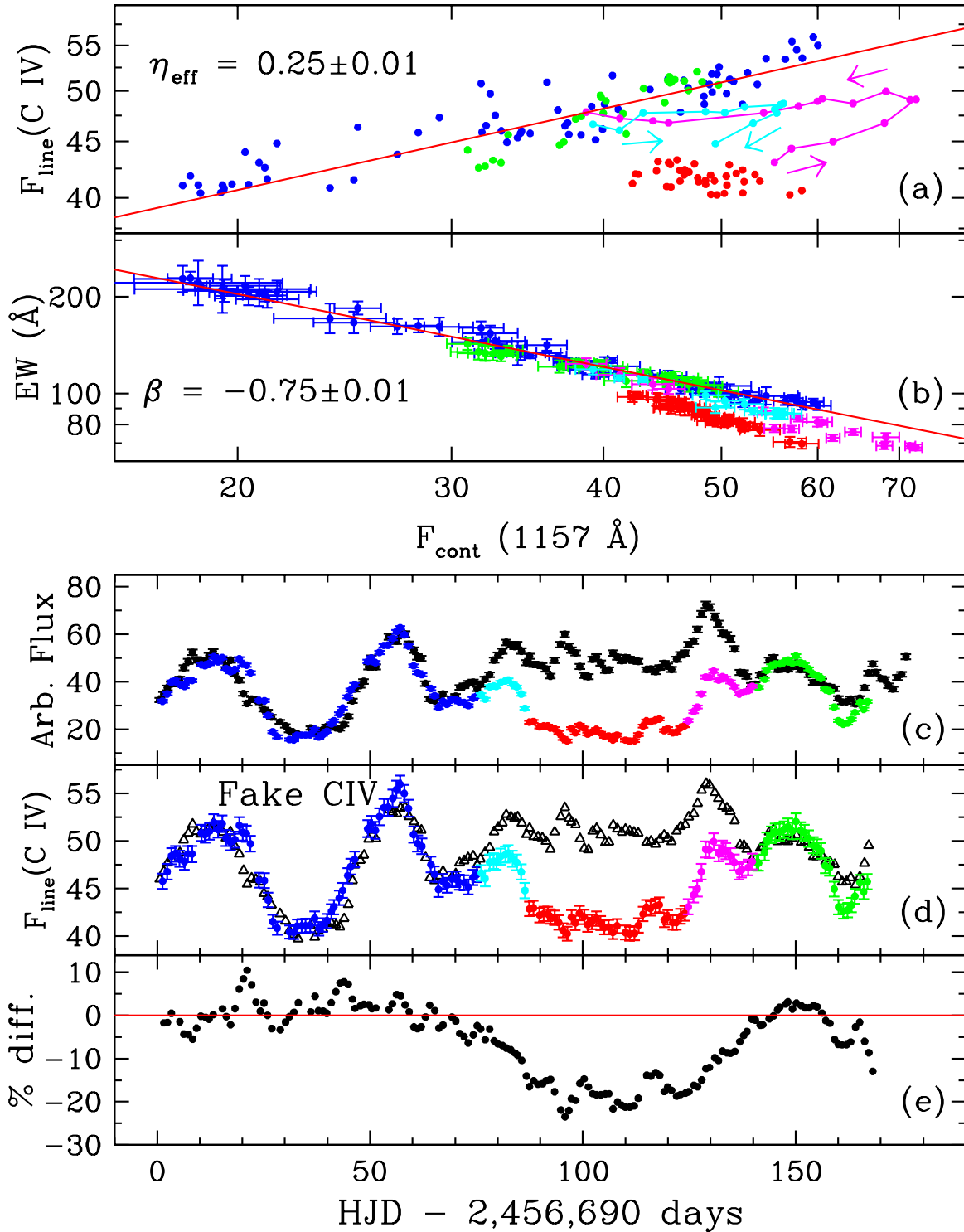


FIG. 1.— Upper two panels – (a) $\log F(\text{C IV})$ versus $\log F_{\text{cont}}(1157 \text{ \AA})$, color coded as follows: HJD–2,400,000: 56,691.2813–56,765.1133 (blue), 56,766.1094–56,776.3984 (cyan), 56,777.4336–56,813.9688 (red), 56,814.7891–56,829.8281 (magenta), 56,830.8242–56,856.8438 (green). Error bars have been omitted for clarity. In particular, note the island of red points lying away from the main relation, as well as the cyan and magenta points which trace a contiguous path to and from the main relation. The red line indicates the best-fit slope to the blue and green points only. (b) $\log \text{EW}(\text{C IV})$ versus $\log F_{\text{cont}}(1157 \text{ \AA})$. As above, the red line indicates the best-fit slope to the blue and green points only. Bottom three panels – (c) A comparison between the scaled and shifted version of the C IV emission-line light curve (colored points) and the observed 1157 Å continuum light curve (black points), highlighting the anomalous flux and response in this emission line midway through the campaign. (d) A comparison between the observed C IV emission-line light curve (colored points) and the reconstructed emission-line light curve (black triangles; see text for details). (e) The estimated percentage deficit in the C IV flux.

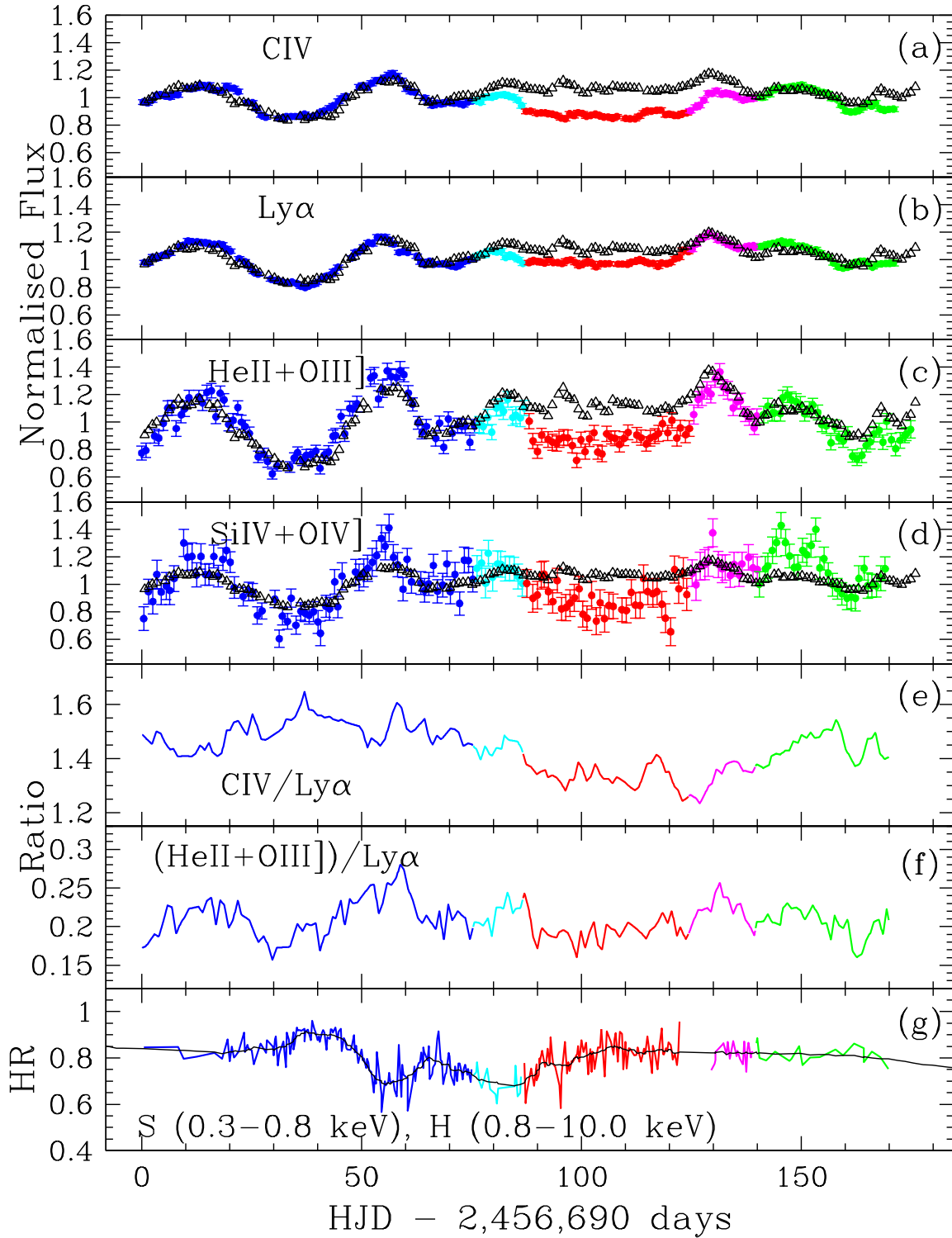


FIG. 2.— Panels (a)–(d), a comparison between the observed BEL light curves (color coded as described in the text) and their respective *reconstructed* light curves (black triangles). Each has been shifted backward in time according to the measured delay between the 1157 Å continuum and emission-line light curve for that particular line, determined over the first 75 days. Panels (e)–(f), time-variable emission-line ratios C IV/Ly α and (He II+O III)]/Ly α . (g) The time-variable *Swift*/XRT X-ray hardness ratio (H–S/H+S) color coded as above). The solid black line indicates the hardness ratio after smoothing with a 15-point boxcar filter.

TABLE 1
 BEL DELAYS AND ADOPTED NARROW EMISSION-LINE FLUXES

Line ID	CCF(cent) (days)	CCF(lag) (days)	F(narrow) (10^{-13} erg s $^{-1}$ cm $^{-2}$)
Ly α	6.69 ± 0.41	6.41 ± 0.51	8.9
Si IV + O IV]	5.80 ± 0.58	5.86 ± 0.89	1.2
C IV	4.97 ± 0.39	4.95 ± 0.50	7.0
He II+ OIII]	2.42 ± 0.44	2.14 ± 0.58	1.2

TABLE 2
TIME-AVERAGED BEL RESPONSIVITIES AND THE SLOPE OF THE
INTRINSIC BALDWIN EFFECT

Line ID	η_{eff}	β
Ly α	0.30 ± 0.01	-0.73 ± 0.02
Si IV + O IV]	0.45 ± 0.01	-0.58 ± 0.03
C IV	0.25 ± 0.01	-0.75 ± 0.01
He II + OIII]	0.58 ± 0.04	-0.48 ± 0.04

TABLE 3
TIME-AVERAGED BEL EWs FOR THE ANOMALY AND THE MAIN
RELATION

Line ID	EW (Å) (anomaly)	EW (Å) (main relation)
Ly α	64.5 ± 5.1	72.5 ± 5.9
Si IV + O IV]	6.2 ± 0.9	8.2 ± 1.2
C IV	86.2 ± 6.9	106.3 ± 8.3
He II + OIII]	12.5 ± 1.3	15.9 ± 3.3

The EWs are here measured with respect to the average UV continuum flux at 1157 Å for those epochs spanning the time of the anomaly (Figure 1a – red points); $F_{\text{cont}}(1157 \text{ \AA}) = (4.84 \pm 0.37) \times 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ \AA}^{-1}$.

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