# PDF hosted at the Radboud Repository of the Radboud University Nijmegen 

The following full text is a publisher's version.

For additional information about this publication click this link.
http://hdl.handle.net/2066/92359

Please be advised that this information was generated on 2017-12-06 and may be subject to change.

# Spherical Fourier Transforms on Locally Compact Quantum Gelfand Pairs ${ }^{\star}$ 

Martijn CASPERS<br>Radboud Universiteit Nijmegen, IMAPP, Heyendaalseweg 135, 6525 AJ Nijmegen, The Netherlands<br>E-mail: caspers@math.ru.nl<br>URL: http://www.math.ru.nl/~caspers/

Received April 14, 2011, in final form August 30, 2011; Published online September 06, 2011
http://dx.doi.org/10.3842/SIGMA.2011.087


#### Abstract

We study Gelfand pairs for locally compact quantum groups. We give an operator algebraic interpretation and show that the quantum Plancherel transformation restricts to a spherical Plancherel transformation. As an example, we turn the quantum group analogue of the normaliser of $S U(1,1)$ in $S L(2, \mathbb{C})$ together with its diagonal subgroup into a pair for which every irreducible corepresentation admits at most two vectors that are invariant with respect to the quantum subgroup. Using a $\mathbb{Z}_{2}$-grading, we obtain product formulae for little $q$-Jacobi functions.


Key words: locally compact quantum groups; Plancherel theorem; Fourier transform; spherical functions
2010 Mathematics Subject Classification: 16T99; 43A90

## 1 Introduction

In the classical setting of locally compact groups, a Gelfand pair consists of a locally compact group $G$, together with a compact subgroup $K$ such that the convolution algebra of bi- $K$ invariant $L^{1}$-functions on $G$ is commutative. See [5] or [8] for a comprehensive introduction. Gelfand pairs give rise to spherical functions and a spherical Fourier transform which decomposes bi- $K$-invariant functions on $G$ as an integral of spherical functions, see [5, Theorem 6.4.5] or [8, Théorème IV.2].

For many examples, this decomposition is made precise [5]. The examples include the group of motions of the plane together with its diagonal subgroup and the pair $\left(S O_{0}(1, n), S O(n)\right)$, where $S O_{0}(1, n)$ is the connected component of the identity of $S O(1, n)$. In particular the spherical functions are determined and one can derive product formulae for these type of functions.

Since the introduction of quantum groups, Gelfand pairs were studied in a quantum context, see for example [9, 26, 39, 40] and also the references given there. These papers consider pairs of quantum groups that are both compact. For such pairs it suffices to stay with a purely (Hopf-)algebraic approach. Under the assumption that every irreducible unitary corepresentation admits only one matrix element that is invariant under both the left and right action of the subgroup, these quantum groups are called (quantum) Gelfand pairs. Classically, this is equivalent to the commutativity assumption on the convolution algebra of bi- $K$-invariant elements. If the matrix coefficients form a commutative algebra one speaks of a strict (quantum) Gelfand pair. In the group setting every Gelfand pair is automatically strict and as such strictness is a purely non-commutative phenomenon.

[^0]For quantum groups, many deformations of classical Gelfand pairs do indeed form a quantum Gelfand pair that moreover is strict. As a compact example, $\left(S U_{q}(n), U_{q}(n-1)\right)$ forms a strict Gelfand pair [40]. In a separate paper [38] Vainerman introduces the quantum group of motions of the plane, together with the circle as a subgroup as an example of a Gelfand pair of which the larger quantum group is non-compact. As a result a product formula for the Hahn-Exton $q$-Bessel functions, also known as $1 \varphi_{1} q$-Bessel functions, is obtained [38, Corollary, p. 324], see also [16, Corollary 6.4]. However, a comprehensive general framework of quantum Gelfand pairs in the non-compact operator algebraic setting was unavailable at that time.

At the turn of the millennium, locally compact (l.c.) quantum groups have been put in an operator algebraic setting by Kustermans and Vaes in their papers [21, 22], see also [19, 31, 42]. The definitions give a $\mathrm{C}^{*}$-algebraic and a von Neumann algebraic interpretation of locally compact quantum groups. Many aspects of abstract harmonic analysis have found a suitable interpretation in this von Neumann algebraic framework. In particular, Desmedt proved in his thesis [4] that there is an analogue of the Plancherel theorem, which gives a decomposition of the left regular corepresentation of a l.c. quantum group.

From this perspective, it is a natural question if the study of Gelfand pairs can be continued in the l.c. operator algebraic setting. In this paper we give this interpretation. Motivated by Desmedt's proof of the quantum Plancherel theorem, we define the necessary structures to obtain a classical Plancherel-Godement theorem [5, Theorem 6.4.5] or [8, Théorème IV.2]. For this the operator algebraic interpretation of Gelfand pairs is essential.

We keep the setting a bit more general than one would expect. For a classical Gelfand pair of groups, one can prove that the larger group is unimodular from the commutativity assumption on bi- $K$-invariant elements. Here we will study pairs of quantum groups for which the smaller quantum group is compact and we assume that the larger group is unimodular. We will not impose the classically stronger commutativity assumption. The reason for this is that we would like to study $S U_{q}(1,1)_{\text {ext }}$ together with its diagonal subgroup. However, the natural analogue of the commutativity assumption would exclude this example.

We mention that it is known that the notion of a quantum subgroup is in a sense too restrictive. Using Koornwinder's twisted primitive elements, it is possible to define double coset spaces associated with $S U_{q}(2)$ and get so called $(\sigma, \tau)$-spherical elements, see [15] for this particular example. See also [14] for a similar study of $S U_{q}(1,1)$ on an algebraic level. The subgroup setting then corresponds to the limiting case $\sigma, \tau \rightarrow \infty$. In the present paper we do not incorporate such a general setting.

Motivated by the Hopf-algebraic framework, we introduce the non-compact analogues of bi-$K$-invariant functions and its dual $[9,40]$ and equip these with weights. We will do this in a von Neumann algebraic manner and for the dual structure also in a $\mathrm{C}^{*}$-algebraic manner. We prove that the $\mathrm{C}^{*}$-algebraic weight lifts to the von Neumann algebraic weight. Moreover, we establish a spherical analogue of a theorem by Kustermans [18] which establishes a correspondence between representations of the (universal) $\mathrm{C}^{*}$-algebraic dual quantum group and corepresentations of the quantum group itself. Eventually, this structure culminates in a quantum PlancherelGodement theorem, as an application of [4, Theorem 3.4.5]. This illustrates the advantage of an operator algebraic interpretation above the Hopf algebraic approach. In particular, we get a spherical $L^{2}$-Fourier transform, or spherical Plancherel transformation, and we show in principle that this is a restriction of the non-spherical Plancherel transformation.

As an example, we treat the first example of a quantum Gelfand pair involving a $q$-deformation of $S U(n, 1)$. Namely, we treat the quantum analogue of the normalizer of $S U(1,1)$ in $S L(2, \mathbb{C})$, which we denote by $S U_{q}(1,1)_{\text {ext }}$, see [13] and [10]. We identify the circle as its diagonal subgroup and study the spherical properties of this pair. We see that the classical commutativity assumption on the convolution algebra is too restrictive to capture $S U_{q}(1,1)_{\text {ext }}$ with its diagonal subgroup. Nevertheless, the pair exhibits properties reminiscent of classical

Gelfand pairs. In particular, we see how one can derive product formulae using gradings on this quantum group and its dual.

We mention that the $q$-deformation of $S U(1,1)$ was first established on the operator algebraic level in [13]. The construction heavily relies on $q$-analysis. More recently, de Commer [3] was able to obtain $S U_{q}(1,1)_{\text {ext }}$ using Galois co-objects. So far, the higher dimensional $q$-deformations of $S U(n, 1)$ remain undefined on the von Neumann algebraic level.

## Structure of the paper

In Section 3 we study the homogeneous space of left and right invariant elements. We also give their dual spaces. Main goal here is the introduction of the von Neumann algebras $N$ and $\hat{N}$ as well as the C ${ }^{*}$-algebras $\hat{N}_{c}$ and $\hat{N}_{u}$. These are homogeneous counterparts of the von Neumann algebra of a quantum group and its dual as well as the underlying reduced and universal dual C*-algebra.

In Section 4, we study the natural weights on these homogeneous spaces. Since the weight on the larger quantum group is generally not a state, the analysis is much more intricate compared to the compact Hopf-algebraic approach. We prove that the $\mathrm{C}^{*}$-algebraic weights defined here lift to the von Neumann algebraic (dual) weight. This result is a major ingredient for the quantum Plancherel-Godement theorem, since it allows us to apply Desmedt's auxiliary theorem [4, Theorem 3.4.5]. Next, we introduce the necessary terminology of corepresentations that admit a vector that is invariant under the action of a subgroup. This is worked out in Section 5.

In Section 6 we elaborate on a spherical version of Kustermans' result [18]: there is a 1-1 correspondence between representations of the universal dual of a quantum group and corepresentations of the quantum group itself. This will form the essential bridge between [4, Theorem 3.4.5] and the quantum Plancherel-Godement Theorem 7.1. Eventually, Section 7 combines the results of Sections 3-6 to prove a quantum version of the Plancherel-Godement theorem.

In Section 8 we work out the example of $S U_{q}(1,1)_{\text {ext }}$ together with its diagonal subgroup. We determine all the objects defined in Sections 3-7. As an application of the theory we find product formulae for little $q$-Jacobi-functions that appear as matrix coefficients of irreducible corepresentations.

## 2 Preliminaries and notation

We briefly recall the definition and essential results from the theory of locally compact quantum groups. The results can be found in [21, 22] and [18]. For an introduction we refer to [31] and [19]. For the theory of weights on von Neumann algebras we refer to [30].

We use the following notational conventions. If $\pi$ and $\rho$ are linear maps, we write $\pi \rho$ for the composition $\pi \circ \rho$. $\iota$ denotes the identity homomorphism. The symbol $\otimes$ will be used for either the tensor product of two elements, of linear maps, the von Neumann algebraic tensor product or the tensor product of representations. We use the leg-numbering notation for operators. For example, if $W \in B(\mathcal{H}) \otimes B(\mathcal{H})$, we write $W_{23}$ for $1 \otimes W$ and $W_{13}=(\Sigma \otimes 1) W_{23}(\Sigma \otimes 1)$, where $\Sigma: \mathcal{H} \otimes \mathcal{H} \rightarrow \mathcal{H} \otimes \mathcal{H}$ is the flip. For a linear map $A$, we denote $\operatorname{Dom}(A)$ for its domain.

Let $B$ be a Banach $*$-algebra. With a representation of $B$, we mean a $*$-homomorphism from $B$ to the bounded operators on a Hilbert space, which is referred to as the representation space. If $B$ is a $\mathrm{C}^{*}$-algebra, we write $\operatorname{Rep}(B)$ for the equivalence classes of representations of $B$ and $\operatorname{IR}(B)$ for the equivalence classes of irreducible representations of $B$. With equivalence, we mean unitary equivalence. With slight abuse of notation we sometimes write $\pi \in \operatorname{Rep}(B)$ (or $\pi \in \operatorname{IR}(B))$ to mean that $\pi$ is an (irreducible) representation of $B$ instead of looking at its class.

If $M$ is a von Neumann algebra and $\omega \in M_{*}$, we denote $\bar{\omega}$ for the normal functional defined by $\bar{\omega}(x)=\overline{\omega\left(x^{*}\right)}$. For $\omega \in M_{*}$ and $x, y \in M$, we denote $x \omega, \omega y, x \omega y$ for the functionals defined by $(x \omega)(z)=\omega(z x),(\omega y)(z)=\omega(y z),(x \omega y)(z)=\omega(y z x)$, where $z \in M$.

Let $\phi$ be a weight on $M$. Let $\mathfrak{n}_{\phi}=\left\{x \in M \mid \phi\left(x^{*} x\right)<\infty\right\}, \mathfrak{m}_{\phi}=\mathfrak{n}_{\phi}^{*} \mathfrak{n}_{\phi}$. Let $\sigma^{\phi}$ be the modular automorphism group of $\phi$. We denote $\mathcal{T}_{\phi}$ for the Tomita algebra defined by

$$
\mathcal{T}_{\phi}=\left\{x \in M \mid x \text { is analytic w.r.t. } \sigma^{\phi} \text { and } \forall z \in \mathbb{C}: \sigma_{z}^{\phi}(x) \in \mathfrak{n}_{\phi} \cap \mathfrak{n}_{\phi}^{*}\right\} .
$$

For $x, y \in \mathcal{T}_{\phi}$, we write $x \phi y$ for the normal functional determined by $(x \phi y)(z)=\phi(y z x), z \in M$.

## Quantum groups

We use the Kustermans-Vaes definition of a locally compact quantum group [21, 22], see also [19, 31, 42].

Definition 2.1. A locally compact quantum group $(M, \Delta)$ consists of the following data:

1. A von Neumann algebra $M$;
2. A unital, normal $*$-homomorphism $\Delta: M \rightarrow M \otimes M$ satisfying the coassociativity relation $(\Delta \otimes \iota) \Delta=(\iota \otimes \Delta) \Delta ;$
3. Two normal, semi-finite, faithful weights $\varphi, \psi$ on $M$ so that

$$
\begin{array}{lll}
\varphi((\omega \otimes \iota) \Delta(x))=\varphi(x) \omega(1), & \forall \omega \in M_{*}^{+}, \forall x \in \mathfrak{m}_{\varphi}^{+} & \text {(left invariance); } \\
\psi((\iota \otimes \omega) \Delta(x))=\psi(x) \omega(1), & \forall \omega \in M_{*}^{+}, \forall x \in \mathfrak{m}_{\psi}^{+} & \text {(right invariance). }
\end{array}
$$

$\varphi$ is the left Haar weight and $\psi$ the right Haar weight.
Note that we suppress the Haar weights in the notation. A locally compact quantum group ( $M, \Delta$ ) is called compact if $\varphi$ and $\psi$ are states. $(M, \Delta)$ is called unimodular if $\varphi=\psi$. Compact quantum groups are unimodular.

The triple $(\mathcal{H}, \pi, \Lambda)$ denotes the GNS-construction with respect to the left Haar weight $\varphi$. We may assume that $M$ acts on the GNS-space $\mathcal{H}$. We use $J$ and $\nabla$ for the modular conjugation and modular operator of $\varphi$ and $\sigma$ for the modular automorphism group of $\varphi$. Recall that there is a constant $\nu \in \mathbb{R}^{+}$called the scaling constant such that $\psi \sigma_{t}=\nu^{-t} \psi$. By applying [34], we see that there is a positive, self-adjoint operator $\delta$, called the modular element, such that $\psi=\varphi_{\delta}$, i.e. formally $\psi(\cdot)=\varphi\left(\delta^{\frac{1}{2}} \cdot \delta^{\frac{1}{2}}\right)$. For compact quantum groups, the scaling constant and the modular element are trivial.

## Multiplicative unitary

There exists a unique unitary operator $W \in B(\mathcal{H} \otimes \mathcal{H})$ defined by

$$
W^{*}(\Lambda(a) \otimes \Lambda(b))=(\Lambda \otimes \Lambda)(\Delta(b)(a \otimes 1)), \quad a, b \in \mathfrak{n}_{\varphi}
$$

$W$ is known as the multiplicative unitary. It satisfies the pentagon equation $W_{12} W_{13} W_{23}=$ $W_{23} W_{12}$ in $B(\mathcal{H} \otimes \mathcal{H} \otimes \mathcal{H})$. Here we use the usual leg numbering notation. Moreover, $W$ implements the comultiplication, i.e. $\Delta(x)=W^{*}(1 \otimes x) W$ and $W \in M \otimes B(\mathcal{H})$.

## The unbounded antipode

To $(M, \Delta)$ one can associate an unbounded map called the antipode $S: \operatorname{Dom}(S) \subseteq M \rightarrow M$. It can be defined as the $\sigma$-strong-* closure of the map $(\iota \otimes \omega)(W) \mapsto(\iota \otimes \omega)\left(W^{*}\right)$, where $\omega \in B(\mathcal{H})_{*}$. One can prove that there exists a unique $*$-anti-automorphism $R: M \rightarrow M$ and a unique strongly continuous one-parameter group of $*$-automorphisms $\tau: \mathbb{R} \rightarrow \operatorname{Aut}(M)$ such that

$$
S=R \tau_{-i / 2}, \quad R^{2}=\iota, \quad \tau_{t} R=R \tau_{t}, \quad \varphi \tau_{t}=\nu^{-t} \varphi, \quad \forall t \in \mathbb{R}
$$

$R$ is called the unitary antipode and $\tau$ is called the scaling group. Moreover,

$$
\begin{equation*}
\Delta R=\Sigma_{M, M}(R \otimes R) \Delta, \quad \psi=\varphi R, \tag{2.1}
\end{equation*}
$$

where $\Sigma_{M, M}: M \otimes M \rightarrow M \otimes M$ is the flip. Using the relative invariance property of the left Haar weight with respect to the scaling group, we define $P$ to be the positive operator on $\mathcal{H}$ such that $P^{i t} \Lambda(x)=\nu^{\frac{t}{2}} \Lambda\left(\tau_{t}(x)\right), t \in \mathbb{R}, x \in \mathfrak{n}_{\varphi}$. We use the notation

$$
M_{*}^{\sharp}=\left\{\omega \in M_{*} \mid \exists \theta \in M_{*}:(\theta \otimes \iota)(W)=(\omega \otimes \iota)(W)^{*}\right\} .
$$

For $\omega \in M_{*}^{\sharp}, \omega^{*}$ is defined by $\left(\omega^{*} \otimes \iota\right)(W)=(\omega \otimes \iota)(W)^{*}$. In that case $\omega^{*}(x)=\bar{\omega}(S(x))$, $x \in \operatorname{Dom}(S)$. For $\omega \in M_{*}^{\sharp}$, we set $\|\omega\|_{*}=\max \left\{\|\omega\|,\left\|\omega^{*}\right\|\right\}$. $M_{*}^{\sharp}$ becomes a Banach-*-algebra with this norm.

## The dual quantum group

In [21, 22], it is proved that there exists a dual locally compact quantum group $(\hat{M}, \hat{\Delta})$, so that $(\hat{\hat{M}}, \hat{\hat{\Delta}})=(M, \Delta)$. The dual left and right Haar weight are denoted by $\hat{\varphi}$ and $\hat{\psi}$. Similarly, all other dual objects will be denoted by a hat, i.e. $\hat{\nabla}, \hat{J}, \hat{\delta}, \hat{\sigma}_{t}, \hat{W}, \ldots P$ is self-dual, i.e. $\hat{P}=P$. By construction,

$$
\hat{M}=\overline{\left\{(\omega \otimes \iota)(W) \mid \omega \in M_{*}\right\}^{\sigma \text {-strong-* }} .}
$$

Furthermore, $\hat{W}=\Sigma W^{*} \Sigma$, where $\Sigma$ denotes the flip on $\mathcal{H} \otimes \mathcal{H}$. This implies that $W \in M \otimes \hat{M}$ and

$$
M=\overline{\left\{(\iota \otimes \omega)(W) \mid \omega \in \hat{M}_{*}\right\}}{ }^{\sigma \text {-strong-* }} .
$$

For $\omega \in M_{*}$, we use the standard notation $\lambda(\omega)=(\omega \otimes \iota)(W) \in \hat{M}$. Then $\lambda: M_{*}^{\sharp} \rightarrow \hat{M}$ is a representation.

We denote $\mathcal{I}$ for the set of $\omega \in M_{*}$, such that $\Lambda(x) \mapsto \omega\left(x^{*}\right), x \in \mathfrak{n}_{\varphi}$, extends to a bounded functional on $\mathcal{H}$. By the Riesz theorem, for every $\omega \in \mathcal{I}$, there is a unique vector denoted by $\xi(\omega) \in \mathcal{H}$ such that $\omega\left(x^{*}\right)=\langle\xi(\omega), \Lambda(x)\rangle, x \in \mathfrak{n}_{\varphi}$. The dual left Haar weight $\hat{\varphi}$ is defined to be the unique normal, semi-finite, faithful weight on $\hat{M}$, with GNS-construction ( $\mathcal{H}, \iota, \hat{\Lambda})$ such that $\lambda(\mathcal{I})$ is a $\sigma$-strong- $* /$ norm core for $\hat{\Lambda}$ and $\hat{\Lambda}(\lambda(\omega))=\xi(\omega), \omega \in \mathcal{I}$. By definition, this means that $\{(\lambda(\omega), \xi(\omega)) \mid \omega \in \mathcal{I}\}$ is dense in the graph of $\hat{\Lambda}$ with respect to the product of the $\sigma$-strong-* topology on $M$ and the norm topology on $\mathcal{H}$.

## Corepresentations

A (unitary) corepresentation is a unitary operator $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ that satisfies the relation $(\Delta \otimes \iota)(U)=U_{13} U_{23}$. In this paper all corepresentations are assumed to be unitary. It follows from the pentagon equation that the multiplicative unitary $W$ is a corepresentation on the GNSspace $\mathcal{H}$. Two corepresentations $U_{1}, U_{2}$ are equivalent if there is a unitary $T: \mathcal{H}_{U_{1}} \rightarrow \mathcal{H}_{U_{2}}$ such that $T(\omega \otimes \iota)\left(U_{1}\right) v=(\omega \otimes \iota)\left(U_{2}\right) T v$ for every $\omega \in M_{*}, v \in \mathcal{H}_{U_{1}}$. We denote $\operatorname{IC}(M)$ for the set of equivalence classes of all unitary corepresentations. For any corepresentation $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ and $\omega \in B\left(\mathcal{H}_{U}\right)_{*},(\iota \otimes \omega)(U) \in \operatorname{Dom}(S)$ and

$$
S(\iota \otimes \omega)(U)=(\iota \otimes \omega)\left(U^{*}\right) .
$$

## Reduced quantum groups

To every locally compact quantum group $(M, \Delta)$ one can associate a reduced $\mathrm{C}^{*}$-algebraic quantum group. We define $M_{c}$ to be the norm closure of $\left\{(\iota \otimes \omega)(W) \mid \omega \in \hat{M}_{*}\right\}$, which is a C ${ }^{*}$-algebra. We restrict $\Delta, \varphi$ and $\psi$ to $M_{c}$ and denote the respective restrictions by $\Delta_{c}, \varphi_{c}, \psi_{c}$. In fact, $\Delta_{c}$ should be considered as a map into the multiplier algebra of the minimal tensor product of $M_{c}$ with itself. The GNS-construction representation $(\mathcal{H}, \Lambda, \pi)$ then restricts to a GNSrepresentation of the $\mathrm{C}^{*}$-algebraic weight $\varphi_{c}$, which is denoted by $\left(\mathcal{H}, \Lambda_{c}, \pi_{c}\right)$. It is proven in [21] that $\left(M_{c}, \Delta_{c}\right)$ forms a reduced $\mathrm{C}^{*}$-algebraic quantum group. Similarly, one defines the reduced dual $\mathrm{C}^{*}$-algebraic quantum group $\left(\hat{M}_{c}, \hat{\Delta}_{c}\right)$. The associated objects are denoted with a hat.

## Universal quantum groups

Universal quantum groups were introduced by Kustermans [18]. For $\omega \in M_{*}^{\sharp}$, we define

$$
\|\omega\|_{u}=\sup \left\{\|\pi(\omega)\| \mid \pi \text { a representation of } M_{*}^{\sharp}\right\} .
$$

Recall that with a representation, we mean a $*$-homomorphism to the bounded operators on a Hilbert space. Note that this defines a norm since the representation $\lambda$ is injective. Let $\hat{M}_{u}$ be the completion of $M_{*}^{\sharp}$ with respect to $\|\cdot\|_{u}$. We let $\lambda_{u}: M_{*}^{\sharp} \rightarrow \hat{M}_{u}$ denote the canonical embedding. $\hat{M}_{u}$ carries the following universal property: if $\pi$ is a representation of $M_{*}^{\sharp}$, then there is a unique representation $\rho$ of $\hat{M}_{u}$ such that $\pi=\rho \lambda_{u}$. In particular, from the representation $\lambda$ we get a surjective map $\hat{\vartheta}: \hat{M}_{u} \rightarrow \hat{M}_{c}$. We define a universal weight on $\hat{M}_{u}$ by setting $\hat{\varphi}_{u}=\hat{\varphi}_{c} \hat{\vartheta}$, and $\hat{\psi}_{u}=\hat{\psi}_{c} \hat{\vartheta}$. The GNS-representation of $\hat{\varphi}_{u}$ is given by $\left(\mathcal{H}, \hat{\Lambda}_{u}=\hat{\Lambda}_{c} \hat{\vartheta}, \hat{\pi}_{u}=\hat{\pi}_{c} \hat{\vartheta}\right)$. If $U \in$ $M \otimes B\left(\mathcal{H}_{U}\right)$ is a corepresentation of $M$, then the map $M_{*}^{\sharp} \rightarrow B\left(\mathcal{H}_{U}\right): \omega \mapsto(\omega \otimes \iota)(U)$ determines a representation of $\hat{M}_{u}$, which we denote by $\pi_{U}$. In fact, it is shown in [18] that this establishes a 1-1 correspondence between corepresentations of $M$ and non-degenerate representations of $\hat{M}_{u}$.

For completeness, we mention that $\hat{M}_{u}$ can be equipped with a comultiplication $\hat{\Delta}_{u}$, which is a map from $\hat{M}_{u}$ to the multiplier algebra of the minimal tensor product of $\hat{M}_{u}$ with itself, such that $\left(\hat{M}_{u}, \hat{\Delta}_{u}\right)$ is a universal C*-algebraic quantum group in the sense of [18]. Similarly, one defines $M_{u}, \Delta_{u}, \vartheta, \varphi_{u}, \psi_{u}, \ldots$.

## 3 Spherical Fourier transforms

Let $(G, K)$ be a pair consisting of a locally compact group $G$ and a compact subgroup $K$. If $L^{1}(K \backslash G / K)$, the bi- $K$-invariant $L^{1}$-functions on $G$ equipped with the convolution product is a commutative algebra, then $(G, K)$ is a called a Gelfand pair. Examples of Gelfand pairs can be found in [5, Chapter 7] and [8].

A notion of Gelfand pairs for compact quantum groups was introduced by Koornwinder [17]. We briefly recall the definition. Consider two unital Hopf-algebras $H, H_{1}$. Denote $\Delta$ for the comultiplication of $H$, denote $\varphi_{1}$ for the Haar functional on $H_{1}$. Suppose that there exists a surjective morphism $\pi: H \rightarrow H_{1}$, so that $H_{1}$ is identified as a quantum subgroup of $H$. Now consider the left and right coactions $\Delta^{l}=(\pi \otimes \iota) \Delta, \Delta^{r}=(\iota \otimes \pi) \Delta$. Define

$$
H_{1} \backslash H=\left\{h \in H \mid \Delta^{l}(h)=1 \otimes h\right\}, \quad H / H_{1}=\left\{h \in H \mid \Delta^{r}(h)=1 \otimes h\right\}
$$

and set $H_{1} \backslash H / H_{1}=\left(H_{1} \backslash H\right) \cap\left(H / H_{1}\right)$. Classically, $H_{1} \backslash H / H_{1}$ corresponds to the algebra of bi- $K$-invariant elements. Set

$$
\tilde{\Delta}=\left(\iota \otimes \varphi_{1} \pi \otimes \iota\right)(\iota \otimes \Delta) \Delta .
$$

Now, the following definition characterizes a quantum Gelfand pair. In fact, there are more equivalent definitions. We state the one which is closest to the theory we develop in the present section.

Definition 3.1. Let $\left(H, H_{1}\right)$ be as above. $\left(H, H_{1}\right)$ is called a Gelfand pair if $\tilde{\Delta}$ is cocommutative, i.e. $\tilde{\Delta}=\Sigma_{H, H} \tilde{\Delta}$. The pair is called a strict Gelfand pair if moreover $H_{1} \backslash H / H_{1}$ is commutative. Here $\Sigma_{H, H}$ denotes the flip.

A pair of compact groups $(G, K)$ is a Gelfand pair if and only if the Hopf-algebra of matrix coefficients of unitary finite dimensional representations form a quantum Gelfand pair (which is automatically strict). Many deformations of classical Gelfand pairs form strict Gelfand pairs in the Hopf-algebraic setting. Examples can be found in for instance [9, 15, 24, 39] and [40].

The aim of this section is to give a general framework of Gelfand pairs in the locally compact quantum group setting as introduced by Kustermans and Vaes [21, 22]. This puts the earlier studies as mentioned in the introduction in a non-compact, von Neumann algebraic setting.

One of the main motivations for the operator algebraic approach is that we can define spherical Fourier transforms. In particular, we show that the structure presented here allows us to prove a decomposition theorem analogous to the classical Plancherel-Godement theorem $[8$, Théorème IV.2] or [5, Section 6]. The proof is an application of Desmedt's auxiliary result [4, Theorem 3.4.5].

As explained in the introduction, we do not assume a natural quantum analogue of the classical commutativity assumption on the convolution algebra of bi-invariant functions. Instead, we assume unimodularity of the larger quantum group, which is classically a weaker assumption. This allows us to cover the example of $S U_{q}(1,1)_{\text {ext }}$, see Section 8 .

Notation 3.2. Throughout Sections $3-7$, we fix a locally compact quantum group $(M, \Delta)$ together with a closed quantum subgroup $\left(M_{1}, \Delta_{1}\right)$ which we assume to be the compact. Recall [37, Definition 2.9] that this means that we have a surjective $*$-homomorphism $\pi: M_{u} \rightarrow\left(M_{1}\right)_{u}$ on the level of universal C*-algebras and the induced dual $*$-homomorphism $\hat{\pi}:\left(\hat{M}_{1}\right)_{u} \rightarrow \hat{M}_{u}$ lifts to a map on the level of von Neumann algebras $\hat{\pi}: \hat{M}_{1} \rightarrow \hat{M}$, which with slight abuse of notation is denoted by $\hat{\pi}$ again. When we encounter $\hat{\pi}$ in this paper, we always mean the von Neumann algebraic map.

Note $\pi$ and $\hat{\pi}$ are in principal also used for the GNS-representations of $M$ and $\hat{M}$. However, we omit the maps most of the time, since $M$ and $\hat{M}$ are identified with their GNS-representations. In that case we explicitly need the GNS-representations, we will mention this.

We mention that from a certain point, see Notation 3.17, we will assume that $(M, \Delta)$ is a unimodular quantum group.

We use $\Sigma_{M_{1}, M}: M_{1} \otimes M \rightarrow M \otimes M_{1}$ to denote the flip. The objects associated with ( $M_{1}, \Delta_{1}$ ) will be equipped with a subscript 1 , i.e. $S_{1}, R_{1}, \tau_{1}, \nu_{1}, \varphi_{1}, \ldots$.

## Homogeneous spaces

Due to [36, Proposition 3.1], there are canonical right coactions of $\left(M_{1}, \Delta_{1}\right)$ on $M$, denoted by $\beta, \gamma: M \rightarrow M \otimes M_{1}$, which are normal $*$-homomorphisms uniquely determined by

$$
\begin{equation*}
(\beta \otimes \iota)(W)=W_{13}\left((\iota \otimes \hat{\pi})\left(W_{1}\right)\right)_{23}, \quad(\gamma \otimes \iota)(W)=\left(\left(R_{1} \otimes \hat{\pi}\right)\left(W_{1}\right)\right)_{23} W_{13} . \tag{3.1}
\end{equation*}
$$

We have the relation $\gamma=(R \otimes \iota) \beta R$. The map $\beta$ corresponds classically to right translation, whereas $\gamma$ corresponds to left translation.
Remark 3.3. In [36, Section 3], the roles of $\left(M_{1}, \Delta_{1}\right)$ and $\left(\hat{M}_{1}, \hat{\Delta}_{1}\right)$ are interchanged. Using our conventions for the roles of $\left(M_{1}, \Delta_{1}\right)$ and $\left(\hat{M}_{1}, \hat{\Delta}_{1}\right)$, recall the left action $\mu$ of $\left(M_{1}, \Delta_{1}\right)$ on $M$ and the left action $\theta$ of $\left(M_{1}, \Delta^{\mathrm{op}}\right)$ on $M$ from [36, Proposition 3.1]. By this proposition, $\beta$ equals the right coaction $\Sigma_{M_{1}, M} \theta . \gamma$ equals the right coaction $\Sigma_{M_{1}, M}\left(R_{1} \otimes \iota\right) \mu$.
Lemma 3.4. As maps $M \rightarrow M \otimes M \otimes M_{1}$, we have an equality

$$
\begin{equation*}
\left(\iota \otimes \Sigma_{M_{1}, M}\right)(\beta \otimes \iota) \Delta=\left(\iota \otimes \iota \otimes R_{1}\right)(\iota \otimes \gamma) \Delta . \tag{3.2}
\end{equation*}
$$

Proof. For the left hand side, using the pentagon equation and [36, Proposition 3.1],

$$
(\beta \otimes \iota \otimes \iota)(\Delta \otimes \iota) W=(\beta \otimes \iota \otimes \iota) W_{13} W_{23}=W_{14}\left((\iota \otimes \hat{\pi})\left(W_{1}\right)\right)_{24} W_{34} .
$$

For the right hand side, using again [36, Proposition 3.1],

$$
\begin{aligned}
(\iota \otimes \iota & \left.\otimes R_{1} \otimes \iota\right)(\iota \otimes \gamma \otimes \iota)(\Delta \otimes \iota) W=\left(\iota \otimes \iota \otimes R_{1} \otimes \iota\right)(\iota \otimes \gamma \otimes \iota) W_{13} W_{23} \\
\quad & =\left(\iota \otimes \iota \otimes R_{1} \otimes \iota\right) W_{14}\left(\left(R_{1} \otimes \iota\right)(\iota \otimes \hat{\pi})\left(W_{1}\right)\right)_{34} W_{24}=W_{14}\left((\iota \otimes \hat{\pi})\left(W_{1}\right)\right)_{34} W_{24} .
\end{aligned}
$$

The lemma follows by the fact that the elements $\left\{(\iota \otimes \omega)(W) \mid \omega \in \hat{M}_{*}\right\}$ are $\sigma$-strong-* dense in $M$.

Definition 3.5. We denote $M^{\beta}$ for the fixed point algebra $\{x \in M \mid \beta(x)=x \otimes 1\}$. Similarly, $M^{\gamma}$ denotes the fixed point algebra of $\gamma$. By definition of $\gamma$ we find $M^{\gamma}=R\left(M^{\beta}\right)$. We define

$$
N=M^{\beta} \cap M^{\gamma} .
$$

Note that $M^{\beta}, M^{\gamma}$ and $N$ are von Neumann algebras. Furthermore, $\Delta\left(M^{\beta}\right) \subseteq M \otimes M^{\beta}$. Also, by $M^{\gamma}=R\left(M^{\beta}\right)$, (2.1) and (3.2) it follows that $\Delta\left(M^{\gamma}\right) \subseteq M^{\gamma} \otimes M$.

We recall from [35] that we have normal, faithful operator valued weights,

$$
T_{\beta}: M^{+} \rightarrow\left(M^{\beta}\right)^{+}: x \mapsto\left(\iota \otimes \varphi_{1}\right) \beta(x) ; \quad T_{\gamma}: M^{+} \rightarrow\left(M^{\gamma}\right)^{+}: x \mapsto\left(\iota \otimes \varphi_{1}\right) \gamma(x) .
$$

Since $\left(M_{1}, \Delta_{1}\right)$ is compact, $T_{\beta}$ and $T_{\gamma}$ are finite. We extend the domains of $T_{\beta}$ and $T_{\gamma}$ to $M$ in the usual way. We denote the extensions again by $T_{\beta}$ and $T_{\gamma}$. The composition of $T_{\beta}$ and $T_{\gamma}$ forms a well-defined map on $M$. Note that $T_{\beta}\left(x^{*}\right)=T_{\beta}(x)^{*}$ and $T_{\gamma}\left(x^{*}\right)=T_{\gamma}(x)^{*}$, where $x \in M$.

Remark 3.6. The spaces $M^{\gamma}, M^{\beta}$ where already introduced in [33] as homogeneous spaces. They also fall within the definition of a homogeneous space as introduced by Kasprzak [12, Remark 3.3]. Moreover, we stretch that $T_{\beta}$ and $T_{\gamma}$ are conditional expectation values, which properties have been studied in the related papers [28] and [32].

Lemma 3.7. $T_{\gamma}: M \rightarrow M^{\gamma}$ and $T_{\beta}: M \rightarrow M^{\beta}$ satisfy the following properties:

1. $T_{\beta} T_{\gamma}=T_{\gamma} T_{\beta}$;
2. $\Delta T_{\beta}=\left(\iota \otimes T_{\beta}\right) \Delta$ and $\Delta T_{\gamma}=\left(T_{\gamma} \otimes \iota\right) \Delta$;
3. $\left(\iota \otimes T_{\gamma}\right) \Delta=\left(T_{\beta} \otimes \iota\right) \Delta$;
4. $T_{\gamma} S \subseteq S T_{\beta}$ and $T_{\beta} S \subseteq S T_{\gamma}$.

Proof. (1) This follows from the fact $\left(\iota \otimes \Sigma_{M_{1}, M_{1}}\right)(\gamma \otimes \iota) \beta=(\beta \otimes \iota) \gamma$, which can be established as in the proof of Lemma 3.4.
(2) We prove that $\Delta T_{\gamma}=\left(T_{\gamma} \otimes \iota\right) \Delta$, the other equation can be proved similarly using $\beta=(R \otimes \iota) \gamma R$ and (2.1). We find:

$$
\begin{aligned}
\left(\Delta T_{\gamma} \otimes \iota\right)(R \otimes \iota)(W) & =\left(\iota \otimes \iota \otimes \varphi_{1} \otimes \iota\right)(\Delta \otimes \iota \otimes \iota)(R \otimes \iota \otimes \iota)(\beta \otimes \iota)(W) \\
& =\left(\Sigma_{M, M} \otimes \iota\right)(R \otimes R \otimes \iota)(\Delta \otimes \iota)\left(W\left(1 \otimes \hat{\pi}\left(\left(\varphi_{1} \otimes \iota\right)\left(W_{1}\right)\right)\right)\right) \\
& =(R \otimes R \otimes \iota) W_{23} W_{13}\left(1 \otimes 1 \otimes \hat{\pi}\left(\left(\varphi_{1} \otimes \iota\right)\left(W_{1}\right)\right)\right) \\
& =\left(\iota \otimes \varphi_{1} \otimes \iota \otimes \iota\right)(R \otimes \iota \otimes \iota \otimes \iota)(\beta \otimes R \otimes \iota) W_{23} W_{13} \\
& =\left(T_{\gamma} \otimes \iota \otimes \iota\right)\left(\Sigma_{M, M} \otimes \iota\right)(R \otimes R \otimes \iota) W_{13} W_{23} \\
& =\left(T_{\gamma} \otimes \iota \otimes \iota\right)(\Delta \otimes \iota)(R \otimes \iota)(W) .
\end{aligned}
$$

Now, the equation follows by taking slices of the second leg of $W$.
(3) Since $\left(M_{1}, \Delta_{1}\right)$ is compact, it is unimodular and hence $\varphi_{1}=\varphi_{1} R_{1}$. Using Lemma 3.4 we find:

$$
\begin{aligned}
\left(\iota \otimes T_{\gamma}\right) \Delta & =\left(\iota \otimes \iota \otimes \varphi_{1}\right)(\iota \otimes \gamma) \Delta=\left(\iota \otimes \iota \otimes \varphi_{1} R_{1}\right)(\iota \otimes \gamma) \Delta \\
& =\left(\iota \otimes \iota \otimes \varphi_{1}\right) \Sigma_{23}(\beta \otimes \iota) \Delta=\left(T_{\beta} \otimes \iota\right) \Delta .
\end{aligned}
$$

(4) It follows from [22, Proposition 2.1, Corollary 2.2] that $\left(\tau_{t} \otimes \iota\right)(W)=\left(\iota \otimes \hat{\tau}_{-t}\right)(W), t \in \mathbb{R}$. Therefore,

$$
\left(\beta \tau_{t} \otimes \iota\right)(W)=\left(\beta \otimes \hat{\tau}_{-t}\right)(W)=\left(\iota \otimes \iota \otimes \hat{\tau}_{-t}\right)\left(W_{13}\left((\iota \otimes \hat{\pi})\left(W_{1}\right)\right)_{23}\right) .
$$

By [21, Proposition 5.45], we know that $\hat{\pi}\left(\hat{\tau}_{1}\right)_{t}=\hat{\tau}_{t} \hat{\pi}$. Here $\left(\hat{\tau}_{1}\right)_{t}$ denotes the scaling group of ( $\hat{M}_{1}, \hat{\Delta}_{1}$ ). Continuing the equation, we find

$$
\begin{aligned}
\left(\beta \tau_{t} \otimes \iota\right)(W) & =\left(\iota \otimes \hat{\tau}_{-t}\right)(W)_{13}\left(\left(\iota \otimes \hat{\pi}\left(\hat{\tau}_{1}\right)_{-t}\right)\left(W_{1}\right)\right)_{23} \\
& \left.=\left(\tau_{t} \otimes \iota\right)(W)_{13}\left(\left(\tau_{1}\right)_{t} \otimes \hat{\pi}\right)\left(W_{1}\right)\right)_{23}=\left(\tau_{t} \otimes\left(\tau_{1}\right)_{t} \otimes \iota\right)(\beta \otimes \iota)(W) .
\end{aligned}
$$

We have $\varphi_{1}\left(\tau_{1}\right)_{t}=\varphi_{1}$, since $\left(M_{1}, \Delta_{1}\right)$ is compact. Hence, $T_{\beta} \tau_{t}=\tau_{t} T_{\beta}$. With a similar computation, involving the relation $(R \otimes \hat{R})(W)=W^{*}$, see [22, Proposition 2.1, Corollary 2.2], we find $T_{\beta} R=R T_{\gamma}$. The proposition follows, since $S=R \tau_{-i / 2}$.

Definition 3.8. For $x \in N$, we define

$$
\begin{equation*}
\Delta^{\natural}(x)=\left(\iota \otimes T_{\gamma}\right) \Delta(x)=\left(T_{\beta} \otimes \iota\right) \Delta(x), \tag{3.3}
\end{equation*}
$$

see Lemma 3.7 (3). This is the von Neumann algebraic version of [40, equation (4)].
Recall that $\Delta(N) \subseteq M^{\gamma} \otimes M^{\beta}$, so that $\Delta^{\natural}(N) \subseteq N \otimes N$. Moreover, using (1)-(3) of the previous lemma, $\Delta^{\natural}$ is coassociative, i.e.

$$
\begin{align*}
\left(\iota \otimes \Delta^{\natural}\right) \Delta^{\natural} & =\left(\iota \otimes \iota \otimes T_{\gamma}\right)(\iota \otimes \Delta)\left(\iota \otimes T_{\gamma}\right) \Delta=\left(\iota \otimes T_{\beta} T_{\gamma} \otimes \iota\right)(\iota \otimes \Delta) \Delta \\
& =\left(\iota \otimes T_{\gamma} T_{\beta} \otimes \iota\right)(\Delta \otimes \iota) \Delta=\left(\Delta^{\natural} \otimes \iota\right) \Delta^{\natural} . \tag{3.4}
\end{align*}
$$

Note that $\Delta^{\natural}$ is unital, but generally not multiplicative.
Definition 3.9. We define a convolution product $*^{\natural}$ on $N_{*}$,

$$
\omega_{1} *^{\natural} \omega_{2}=\left(\omega_{1} \otimes \omega_{2}\right) \Delta^{\natural}, \quad \omega_{1}, \omega_{2} \in N_{*} .
$$

This convolution product is associative, since by (3.4) $\Delta^{\natural}$ is coassociative.
Definition 3.10. If $\omega \in N_{*}$, then we define

$$
\tilde{\omega}=\omega T_{\gamma} T_{\beta}=\omega T_{\beta} T_{\gamma} \in M_{*} .
$$

For $\omega \in M_{*}$, we put $\tilde{\omega}=\left(\left.\omega\right|_{N}\right)^{\sim}$.
Remark 3.11. Note that $\Delta^{\natural}$ is the von Neumann algebraic version of $\tilde{\Delta}$ [40], which was used to define hypergroup structures. See also the remarks at the beginning of this section. Here, we will not focus on hypergroups for two reasons. First of all, we stay mostly at the measurable von Neumann algebraic setting, which does not allow one to directly define the generalized shift operators [40]. Moreover, we will not assume that $N$ is Abelian, i.e. what is called a strict Gelfand pair in [40].

Proposition 3.12. The map $N_{*} \ni \omega \mapsto \tilde{\omega}$ defines a bijective, norm preserving correspondence between $N_{*}$ and the functionals $\theta \in M_{*}$ that satisfy the invariance properties:

1. $(\theta \otimes \iota) \beta(x)=\theta(x) 1_{M_{1}}$ for all $x \in M$;
2. $(\theta \otimes \iota) \gamma(x)=\theta(x) 1_{M_{1}}$ for all $x \in M$.

Moreover, for $\omega_{1}, \omega_{2} \in N_{*}$ we have $\left(\omega_{1} *^{\natural} \omega_{2}\right)^{\sim}=\tilde{\omega}_{1} * \tilde{\omega}_{2}$.
Proof. Using right invariance of $\varphi_{1}$, for $x \in M$,

$$
\begin{aligned}
\left(T_{\beta} \otimes \iota\right) \beta(x) & =\left(\iota \otimes \varphi_{1} \otimes \iota\right)(\beta \otimes \iota) \beta(x)=\left(\iota \otimes \varphi_{1} \otimes \iota\right)\left(\iota \otimes \Delta_{1}\right) \beta(x) \\
& =\left(\iota \otimes \varphi_{1}\right) \beta(x) \otimes 1_{M_{1}}=T_{\beta}(x) \otimes 1_{M_{1}}
\end{aligned}
$$

Similarly, the right invariance of $\varphi_{1}$ implies $\left(T_{\gamma} \otimes \iota\right) \gamma(x)=T_{\gamma}(x) \otimes 1_{M_{1}}, x \in M$. Using (1) of Lemma 3.7 one easily verifies that for $\omega \in N_{*}, \tilde{\omega}$ satisfies the invariance properties (1) and (2).

Let $\omega \in N_{*}$. For $x \in N$ we have $\omega(x)=\tilde{\omega}(x)$, so that $N_{*} \ni \omega \mapsto \tilde{\omega}$ is injective. If $\theta \in M_{*}$ satisfies (1), then for $x \in M$,

$$
\theta T_{\beta}(x)=\theta\left(\iota \otimes \varphi_{1}\right) \beta(x)=\varphi_{1}(\theta \otimes \iota) \beta(x)=\theta(x)
$$

Similarly, if $\theta \in M_{*}$ satisfies (2), then $\theta T_{\gamma}(x)=\theta(x)$. We find that $\theta=\left(\left.\theta\right|_{N}\right)^{\sim}$ if $\theta$ satisfies (1) and (2). So $N_{*} \ni \omega \mapsto \tilde{\omega}$ ranges over the normal functionals on $M$ that satisfy the invariance properties (1) and (2).

Using the left invariance of $\varphi_{1}$, it is a straightforward check that $T_{\gamma} T_{\gamma}=T_{\gamma}$. Then, using (1)-(3) of Lemma 3.7, we find:

$$
\begin{aligned}
\left(\omega_{1} *^{\natural} \omega_{2}\right)^{\sim} & =\left(\omega_{1} \otimes \omega_{2}\right)\left(\iota \otimes T_{\gamma}\right) \Delta T_{\gamma} T_{\beta}=\left(\omega_{1} \otimes \omega_{2}\right)\left(T_{\gamma} \otimes T_{\gamma} T_{\beta}\right) \Delta \\
& =\left(\omega_{1} \otimes \omega_{2}\right)\left(T_{\gamma} \otimes T_{\gamma}^{2} T_{\beta}\right) \Delta=\left(\omega_{1} \otimes \omega_{2}\right)\left(T_{\gamma} T_{\beta} \otimes T_{\gamma} T_{\beta}\right) \Delta=\tilde{\omega_{1}} * \tilde{\omega_{2}} .
\end{aligned}
$$

Proposition 3.13. For $\omega \in M_{*}^{\sharp}$, we have $\tilde{\omega} \in M_{*}^{\sharp}$ and $(\tilde{\omega})^{*}=\left(\omega^{*}\right)^{\sim}$.
Proof. Using (4) of Lemma 3.7, for $x \in \operatorname{Dom}(S)$, we find

$$
\overline{\tilde{\omega}}(S(x))=\overline{\omega\left(T_{\beta} T_{\gamma}\left(S(x)^{*}\right)\right)}=\overline{\omega\left(\left(T_{\beta} T_{\gamma}(S(x))\right)^{*}\right)}=\overline{\omega\left(S\left(T_{\beta} T_{\gamma}(x)\right)^{*}\right)}=\left(\omega^{*}\right)^{\sim}(x) .
$$

So $\left(\omega^{*}\right)^{\sim} \in M_{*}$ has the property $\left(\left(\omega^{*}\right)^{\sim} \otimes \iota\right)(W)=(\tilde{\omega} \otimes \iota)(W)^{*}$. This proves that $\tilde{\omega} \in M_{*}^{\sharp}$ and $\tilde{\omega}^{*}=\left(\omega^{*}\right)^{\sim}$.

The following proposition is proved in [36, Proposition 3.1].
Proposition 3.14. If $x \in \mathfrak{n}_{\varphi}$, then $T_{\gamma}(x) \in \mathfrak{n}_{\varphi}$ and $\Lambda\left(T_{\gamma}(x)\right)=\hat{\pi}\left(\left(\varphi_{1} \otimes \iota\right)\left(W_{1}\right)\right) \Lambda(x)$.
Proposition 3.14 defines an orthogonal projection $\hat{\pi}\left(\left(\varphi_{1} \otimes \iota\right)\left(W_{1}\right)\right)$ for which we simply write

$$
P_{\gamma}=\hat{\pi}\left(\left(\varphi_{1} \otimes \iota\right)\left(W_{1}\right)\right)
$$

Classically, it corresponds to projecting onto the space of functions that are left invariant with respect to the compact subgroup. Note that $P_{\gamma} \in \hat{M}$ and

$$
N=T_{\beta} T_{\gamma}(M)=\left\{\left(\iota \otimes \omega_{P_{\gamma} v, P_{\gamma} w}\right)(W) \mid v, w \in \mathcal{H}\right\}
$$

We need a similar result as Proposition 3.14 for $T_{\beta}$. For this we need unimodularity of $(M, \Delta)$. Classically, if $G$ is a group with compact subgroup $K$ such that $(G, K)$ forms a Gelfand pair, one can prove that $G$ is unimodular, see [8, Proposition I.1]. The natural definition of a quantum Gelfand pair would be to require that $\Delta^{\natural}$ is cocommutative. However, we like to stretch the
definition of a Gelfand pair a bit to handle the example of $S U_{q}(1,1)_{\text {ext }}$. The following essential result, see Proposition 3.16, is the motivation of assuming the (classically) weaker condition that $(M, \Delta)$ is unimodular, see Notation 3.17.

First, we need the following lemma. Note that for $a, b \in \mathcal{T}_{\varphi}$, we have $a \varphi b \in M_{*}$ and hence $\mathcal{T}_{\varphi} \varphi \mathcal{T}_{\varphi}$ is a subset of $M_{*}$. Recall that for $\omega \in \mathcal{I}, \xi(\omega) \in \mathcal{H}$ is defined using the Riesz theorem as the unique vector such that $\langle\xi(\omega), \Lambda(x)\rangle=\omega\left(x^{*}\right), x \in \mathfrak{n}_{\varphi}$. By [21, Lemma 8.5], $\mathcal{T}_{\varphi} \varphi \mathcal{T}_{\varphi}$ is included in $\mathcal{I}$.

Lemma 3.15. Let $\omega \in \mathcal{I}$. There exists a net $\left(\omega_{j}\right)_{j}$ in $\mathcal{T}_{\varphi} \varphi \mathcal{T}_{\varphi}$ such that $\omega_{j} \rightarrow \omega$ in norm and $\xi\left(\omega_{j}\right) \rightarrow \xi(\omega)$ in norm.

Proof. For $\omega \in \mathcal{I}$ we define the norm $\|\omega\|_{\mathcal{I}}=\max \{\|\omega\|,\|\xi(\omega)\|\}$. We have to prove that $\mathcal{T}_{\varphi} \varphi \mathcal{T}_{\varphi}$ is dense in $\mathcal{I}$ with respect to this norm. This is exactly what is obtained in the proof of [1, Proposition 3.4]. Indeed, let $L$ and $k$ be as in [1]. As indicated in the introduction of [1], $\mathcal{T}_{\varphi}^{2} \subseteq L$ and for $a, b \in \mathcal{T}_{\varphi}, k(a b)=\sigma_{i}(b) \varphi a$, see also [1, Corollary 2.15]. So $k\left(\mathcal{T}_{\varphi}^{2}\right)=\mathcal{T}_{\varphi} \varphi \mathcal{T}_{\varphi}$. The proposition yields that $k\left(\mathcal{T}_{\varphi}^{2}\right)$ is dense in $\mathcal{I}$.

Proposition 3.16. Suppose that $(M, \Delta)$ is unimodular. For $x \in \mathfrak{n}_{\varphi}$, we have $T_{\beta}(x) \in \mathfrak{n}_{\varphi}$. The map $\Lambda(x) \mapsto \Lambda\left(T_{\beta}(x)\right)$ is bounded and it extends continuously to the projection $\hat{J} P_{\gamma} \hat{J}$.

Proof. By Proposition 3.14, we see that $\Lambda\left(T_{\gamma}\left(x^{*}\right)\right)=\Lambda\left(T_{\gamma}(x)^{*}\right), x \in \mathfrak{n}_{\varphi} \cap \mathfrak{n}_{\varphi}^{*}$. Denote $T$ for the closure of the map $\Lambda(x) \mapsto \Lambda\left(x^{*}\right), x \in \mathfrak{n}_{\varphi} \cap \mathfrak{n}_{\varphi}^{*}$. We see that $P_{\gamma} \mathcal{H}$ is an invariant subspace for $T$. Since $T=J \nabla^{1 / 2}$, we find that $\nabla^{i t}, t \in \mathbb{R}$ commutes with $P_{\gamma}$.

Recall [21, Lemma 8.8, Proposition 8.9] that $\hat{\nabla}^{i t}=P^{i t} J \delta^{i t} J$ and by Pontrjagin duality $\nabla^{i t}=\hat{P}^{i t} \hat{J} \hat{\delta}^{i t} \hat{J}$. Using $\delta=1$ and the self-duality $\hat{P}=P$, we see that $\hat{\nabla}^{i t}=\nabla^{i t} \hat{J} \hat{\delta}^{-i t} \hat{J}$. Since $\hat{\delta}$ is affiliated with $\hat{M}$ and using the previous paragraph, this shows that $P_{\gamma} \hat{\nabla}^{i t}=\hat{\nabla}^{i t} P_{\gamma}$. Hence $\hat{\sigma}_{t}\left(P_{\gamma}\right)=P_{\gamma}$.

Now, let $a, b \in \mathcal{T}_{\hat{\varphi}}$ and put $\omega=a \hat{\varphi} b \in \hat{M}_{*}$. Then,

$$
T_{\beta}\left((\iota \otimes \omega)\left(W^{*}\right)\right)=(\iota \otimes \omega)\left(\left(1 \otimes P_{\gamma}\right) W^{*}\right) .
$$

Since $P_{\gamma} \in \hat{M}$ is invariant under $\hat{\sigma}_{t}$, it follows from [30, Chapter VIII, Lemma 2.4 (ii) and Lemma 2.5] that $\omega P_{\gamma} \in \hat{\mathcal{I}}$ and

$$
\hat{\xi}\left(\omega P_{\gamma}\right)=\hat{\Lambda}\left(a \hat{\sigma}_{-i}(b) \hat{\sigma}_{-i}\left(P_{\gamma}\right)\right)=\hat{J} P_{\gamma} \hat{J} \hat{\Lambda}\left(a \hat{\sigma}_{-i}(b)\right) .
$$

Now, let $\omega \in \hat{\mathcal{I}}$. We prove the proposition for $x=\hat{\lambda}(\omega) \in \mathfrak{n}_{\varphi}$. Let $\left(\omega_{j}\right)_{j \in J}$ be a net in $\mathcal{T}_{\hat{\varphi}} \hat{\varphi} \mathcal{T}_{\hat{\varphi}}$ that converges to $\omega$ and such that $\hat{\xi}\left(\omega_{j}\right)$ converges to $\hat{\xi}(\omega)$, c.f. Lemma 3.15. Then, $T_{\beta}\left(\hat{\lambda}\left(\omega_{j}\right)\right) \in$ $\mathfrak{n}_{\varphi}$ and $T_{\beta}\left(\hat{\lambda}\left(\omega_{j}\right)\right) \rightarrow T_{\beta}(x)$ in the $\sigma$-weak topology. Furthermore, $\Lambda\left(T_{\beta}\left(\hat{\lambda}\left(\omega_{j}\right)\right)\right)=\hat{J} P_{\gamma} \hat{J} \hat{\xi}\left(\omega_{j}\right)$ is norm convergent. Since $\Lambda$ is $\sigma$-weak/weak closed, and $\operatorname{Dom}(\Lambda)=\mathfrak{n}_{\varphi}$, this proves that $T_{\beta}(x) \in \mathfrak{n}_{\varphi}$ and $\Lambda\left(T_{\beta}(x)\right)=\hat{J} P_{\gamma} \hat{J} \Lambda(x)$.

Since the elements in $\hat{\lambda}(\hat{\mathcal{I}})$ form a $\sigma$-strong- $* /$ norm core for $\Lambda$, this proves the proposition.
We will write $P_{\beta}$ for the projection $\hat{J} P_{\gamma} \hat{J}$. In particular $P_{\beta} \in \hat{M}^{\prime}$. Under the assumption that $(M, \Delta)$ is unimodular, we see that $P_{\beta}$ projects onto the elements that are right invariant with respect to the closed quantum subgroup $\left(M_{1}, \Delta_{1}\right)$. Since we will need this interpretation of $P_{\beta}$, i.e. Proposition 3.16, we assume unimodularity from now on.

Notation 3.17. From this point we assume that $(M, \Delta)$ is unimodular, i.e. $\varphi=\psi$.
We are ready to define the dual structures associated with $N$. We define left and right invariant analogues of the dual von Neumann algebraic quantum group and the universal dual

C*-algebraic quantum group. These duals can be constructed by means of the multiplicative unitary $W$ associated with $(M, \Delta)$. We define

$$
N_{*}^{\sharp}=\left\{\omega \in N_{*} \mid \exists \theta \in N_{*}:(\tilde{\omega} \otimes \iota)(W)^{*}=(\tilde{\theta} \otimes \iota)(W)\right\} .
$$

For $\omega \in N_{*}^{\sharp}$, we set $\|\omega\|_{*}=\max \left\{\|\omega\|,\left\|\omega^{*}\right\|\right\}$. Then, $N_{*}^{\sharp}$ becomes a Banach-*-algebra with respect to this norm. Proposition 3.13 shows that $\omega \in N_{*}^{\sharp}$ if and only if $\tilde{\omega} \in M_{*}^{\sharp}$. Note that $N_{*}^{\sharp}$ is dense in $N_{*}$. Indeed, the restriction map $M_{*} \rightarrow N_{*}:\left.\omega \mapsto \omega\right|_{N}$ is continuous and the image of the subset $M_{*}^{\sharp} \subseteq M_{*}$ is contained in $N_{*}^{\sharp}$ by Proposition 3.13. The inclusion $M_{*}^{\sharp} \subseteq M_{*}$ is dense, see [22, Lemma 2.5]. Hence $N_{*}^{\sharp}$ is dense in $N_{*}$. Using this together with Propositions 3.12 and 3.13, we see that

$$
\hat{N}={\overline{\left\{(\tilde{\omega} \otimes \iota)(W) \mid \omega \in N_{*}\right\}}}^{\sigma \text {-strong-* }}
$$

is a $*$-subalgebra of $\hat{M}$. Since we can conveniently write $(\tilde{\omega} \otimes \iota)(W)=P_{\gamma}(\omega \otimes \iota)(W) P_{\gamma}$ by (3.1), we see that $\hat{N}$ is a von Neumann algebra if considered as acting on $P_{\gamma} \mathcal{H}$, so that $P_{\gamma}$ is its unit. In particular $\hat{N}=P_{\gamma} \hat{M} P_{\gamma}$.

We define $\hat{N}_{c}$ to be the norm closure of the set $\left\{(\tilde{\omega} \otimes \iota)(W) \mid \omega \in N_{*}\right\}$. Then $\hat{N}_{c}$ is a C ${ }^{*}-$ subalgebra of the reduced dual C*-algebra $\hat{M}_{c}$.

For $\omega \in N_{*}^{\sharp}$, we define

$$
\|\omega\|_{u}^{\natural}=\sup \left\{\|\pi(\omega)\| \mid \pi \text { a representation of } N_{*}^{\sharp}\right\} .
$$

Note that this defines a norm since the representation $\omega \mapsto(\tilde{\omega} \otimes \iota)(W)$ is injective as follows using the bijective correspondence established in Proposition 3.12. Let $\hat{N}_{u}$ be the completion of $N_{*}^{\sharp}$ with respect to $\|\cdot\|_{u}^{\natural}$.

Recall the $\operatorname{map} \lambda: M_{*}^{\sharp} \rightarrow \hat{M}: \omega \mapsto(\omega \otimes \iota)(W)$. We set

$$
\lambda^{\natural}: N_{*}^{\sharp} \rightarrow \hat{N}: \omega \mapsto(\tilde{\omega} \otimes \iota)(W) .
$$

Note that the image of $\lambda^{\natural}$ is contained in $\hat{N}_{c}$ and we will use this implicitly. $\lambda_{u}: M_{*}^{\sharp} \rightarrow \hat{M}_{u}$ is the canonical inclusion and similarly

$$
\lambda_{u}^{\natural}: \quad N_{*}^{\sharp} \rightarrow \hat{N}_{u}
$$

denotes the canonical inclusion.
Recall that $\hat{N}_{u}$ is a C ${ }^{*}$-algebra with the following universal property: if $\pi$ is a representation of $N_{*}^{\sharp}$ on a Hilbert space, then there is a unique representation $\rho$ of $\hat{N}_{u}$ such that $\pi=\rho \lambda_{u}^{\natural}$. By this universal property, the map $N_{*}^{\sharp} \rightarrow \hat{M}_{u}: \omega \mapsto \lambda_{u}(\tilde{\omega})$ extends to a representation

$$
\iota_{u}: \hat{N}_{u} \rightarrow \hat{M}_{u}
$$

Similarly, the map $\lambda^{\natural}: N_{*}^{\sharp} \rightarrow \hat{N}_{c}$ gives rise to a surjective map

$$
\hat{\vartheta}^{\natural}: \quad \hat{N}_{u} \rightarrow \hat{N}_{c}
$$

In particular, $\hat{\vartheta} \iota_{u}=\hat{\vartheta} \natural$, where $\hat{\vartheta}: \hat{M}_{u} \rightarrow \hat{M}_{c}$ was the canonical projection induced by the representation $\lambda: M_{*}^{\sharp} \rightarrow \hat{M}_{c}$.

Remark 3.18. Note that we do not claim that $\iota_{u}: \hat{N}_{u} \rightarrow \hat{M}_{u}$ is injective. In fact, this is generally not true, see the comments made in Remark 6.4 and the paragraph before this remark.

## 4 Weights on homogeneous spaces

We introduce weights on the von Neumann algebras and $\mathrm{C}^{*}$-algebras as were introduced in Section 3. We study their GNS-representations and prove Proposition 4.7, which is essential for implementing [4, Theorem 3.4.5].

Recall that the $\mathrm{C}^{*}$-algebraic weights $\hat{\varphi}_{u}, \hat{\varphi}_{c}$ were defined in Section 2. The weights on the von Neumann algebras $M, \hat{M}$ and the C*-algebras $\hat{M}_{c}, \hat{M}_{u}$ restrict to weights on $N, \hat{N}$ and $\hat{N}_{c}, \hat{N}_{u}$ by setting

$$
\varphi^{\natural}=\left.\varphi\right|_{N}, \quad \hat{\varphi}^{\natural}=\left.\hat{\varphi}\right|_{\hat{N}} \quad \text { and } \quad \hat{\varphi}_{c}^{\natural}=\left.\hat{\varphi}_{c}\right|_{\hat{N}_{c}}, \quad \hat{\varphi}_{u}^{\natural}=\hat{\varphi}_{u} \iota_{u}=\hat{\varphi}_{c} \hat{\vartheta}_{\iota_{u}}=\hat{\varphi}_{c}^{\natural} \hat{\vartheta}^{\natural},
$$

respectively. We prove that $\varphi^{\natural}$ and $\hat{\varphi}^{\natural}$ are normal, semi-finite, faithful weights and $\hat{\varphi}_{c}^{\natural}$ and $\hat{\varphi}_{u}^{\natural}$ are lower semi-continuous, densely-defined, non-zero weights. Here the assumption made in Notation 3.17 becomes essential.

Proposition 4.1. $\varphi^{\natural}$ is a normal, semi-finite, faithful weight on N. Its GNS-representation is given by $\left(P_{\gamma} P_{\beta} \mathcal{H},\left.\Lambda\right|_{N \cap \mathfrak{n}_{\varphi}},\left.\pi\right|_{N}\right)$.

Proof. Trivially, $\varphi^{\natural}$ is normal and faithful. Since $\mathfrak{n}_{\varphi}$ is $\sigma$-weakly dense in $M, T_{\beta}$ and $T_{\gamma}$ are $\sigma$-weakly continuous and $N=T_{\gamma} T_{\beta}(M)$, Propositions 3.14 and 3.16 prove that $\varphi^{\natural}$ is semifinite. It is straightforward to check that $\left(P_{\gamma} P_{\beta} \mathcal{H},\left.\Lambda\right|_{N \cap n_{\varphi}},\left.\pi\right|_{N}\right)$ satisfies all the properties of a GNS-representation [30, Section VII.1].

Proposition 4.2. For $\omega \in \mathcal{I}$, we find $\tilde{\omega} \in \mathcal{I}$ and $\xi(\tilde{\omega})=P_{\beta} P_{\gamma} \xi(\omega)$. The set $\mathcal{I}_{N}=\left\{\omega \in N_{*} \mid\right.$ $\tilde{\omega} \in \mathcal{I}\}$ is dense in $N_{*}$. The set $\left\{\xi(\tilde{\omega}) \mid \omega \in \mathcal{I}_{N}\right\}$ is a dense subset of $P_{\beta} P_{\gamma} \mathcal{H}$.

Proof. For $x \in \mathfrak{n}_{\varphi}$, using Propositions 3.14 and 3.16,

$$
\Lambda(x) \mapsto \omega\left(T_{\beta} T_{\gamma}\left(x^{*}\right)\right)=\omega\left(\left(T_{\beta} T_{\gamma}(x)\right)^{*}\right)=\left\langle\xi(\omega), \Lambda\left(T_{\beta} T_{\gamma}(x)\right)\right\rangle=\left\langle P_{\beta} P_{\gamma} \xi(\omega), \Lambda(x)\right\rangle .
$$

The first claim now follows by the definitions of $\mathcal{I}$ and $\xi(\cdot)$. Moreover, we find $\mathcal{I}_{N}=\left\{\left.\omega\right|_{N} \mid \omega \in\right.$ $\mathcal{I}\}$, so that the second claim follows by [21, Lemma 8.5]. The last claim also follows from [21, Lemma 8.5].

Proposition 4.3. $\hat{\varphi}^{\natural}$ is a normal, semi-finite, faithful weight on $\hat{N}$. Its GNS-representation is given by $\left(P_{\gamma} P_{\beta} \mathcal{H},\left.\hat{\Lambda}\right|_{\hat{N} \cap \mathfrak{n}_{\hat{\varphi}}},\left.\hat{\pi}\right|_{\hat{N}}\right)$.

Proof. By Proposition 4.2, $\left\{(\tilde{\omega} \otimes \iota)(W) \mid \omega \in \mathcal{I}_{N}\right\} \subseteq \hat{N}$ is a $\sigma$-strong-* dense subset of $\hat{N}$ contained in $\mathfrak{n}_{\hat{\varphi}}$. This proves that $\hat{\varphi}^{\natural}$ is semi-finite. Trivially, $\hat{\varphi}^{\natural}$ is normal and faithful.

To prove the claim about the GNS-representation, we only need to show that the image of $\left.\hat{\Lambda}\right|_{\hat{N} \cap n_{\varphi}}$ is contained in $P_{\beta} P_{\gamma} \mathcal{H}$ and that the inclusion is dense. For $\omega \in \mathcal{I}_{N}$, we have $\lambda(\tilde{\omega}) \in$ $\hat{N} \cap \mathfrak{n}_{\hat{\varphi}}$ and $\hat{\Lambda}(\lambda(\tilde{\omega}))=\xi(\tilde{\omega}) \in P_{\beta} P_{\gamma} \mathcal{H}$ by Proposition 4.2. Now, let $x \in \hat{N} \cap \mathfrak{n}_{\hat{\varphi}}$. Since the elements $\lambda(\mathcal{I})$ form a $\sigma$-strong-*/norm core for $\hat{\Lambda}$, we can find a net $\left(\omega_{j}\right)_{j \in J}$ in $\mathcal{I}$ such that $\lambda\left(\omega_{j}\right) \rightarrow x$ in the $\sigma$-strong-* topology and $\xi\left(\omega_{j}\right) \rightarrow \hat{\Lambda}(x)$ in norm. Consider the net $\left(\tilde{\omega}_{j}\right)_{j \in J}$. We find $\lambda\left(\tilde{\omega}_{j}\right)=P_{\gamma} \lambda\left(\omega_{j}\right) P_{\gamma} \rightarrow P_{\gamma} x P_{\gamma}=x$. And $\xi\left(\tilde{\omega}_{j}\right)=P_{\beta} P_{\gamma} \xi\left(\omega_{j}\right)$ is norm convergent to $P_{\beta} P_{\gamma} \hat{\Lambda}(x)$. Since $\hat{\Lambda}$ is $\sigma$-strong-*/norm closed, it follows that $\hat{\Lambda}(x)=P_{\beta} P_{\gamma} \hat{\Lambda}(x) \in P_{\beta} P_{\gamma} \mathcal{H}$. Moreover, it follows from Proposition 4.2 that the range of $\left.\hat{\Lambda}\right|_{\hat{N} \cap \mathfrak{n}_{\varphi}}$ is dense in $P_{\beta} P_{\gamma} \mathcal{H}$.

We refer to [21, Section 1.1] for the definition of a GNS-representation for a $\mathrm{C}^{*}$-algebraic weight.

Proposition 4.4. $\hat{\varphi}_{c}^{\dagger}$ is a proper (i.e. densely defined, lower semi-continuous, non-zero) weight on $\hat{N}_{c}$. Its GNS-representation is given by $\left(P_{\gamma} P_{\beta} \mathcal{H},\left.\hat{\Lambda}\right|_{\hat{N}_{c} \cap \mathfrak{n}_{\hat{\varphi}_{c}}},\left.\hat{\pi}\right|_{\hat{N}_{c}}\right)$.

Proof. By Proposition 4.2, $\left\{(\tilde{\omega} \otimes \iota)(W) \mid \omega \in \mathcal{I}_{N}\right\} \subseteq \hat{N}$ is a norm dense subset of $\hat{N}_{c}$ contained in $\mathfrak{n}_{\hat{\varphi}}$. The lower semi-continuity and non-triviality follow since $\hat{\varphi}_{c}^{\natural}$ is a restriction of the faithful weight $\hat{\varphi}_{c}$. The claim on the GNS-representation follows exactly as in the proof of Proposition 4.3.

The following lemma can be found as [18, Lemma 4.2]
Lemma 4.5. $\mathcal{I} \cap M_{*}^{\sharp}$ is dense in $M_{*}^{\sharp}$ with respect to the norm $\|\cdot\|_{*}$.
Proposition 4.6. $\hat{\varphi}_{u}^{h}$ is a proper (i.e. densely defined, lower semi-continuous, non-zero) weight on $\hat{N}_{u}$. Its GNS-representation is given by $\left(P_{\gamma} P_{\beta} \mathcal{H}, \hat{\Lambda}_{u} \iota_{u}, \hat{\pi}_{u} \iota_{u}\right)$.

Proof. Lemma 4.5 shows that $\mathcal{I} \cap M_{*}^{\sharp}$ is dense in $M_{*}^{\sharp}$. Since $\mathcal{I}_{N}$ consists of the restrictions to $N$ of functionals in $\mathcal{I}$, see Proposition 4.2 , and $N_{*}^{\sharp}$ consists of the restrictions to $N$ of functionals in $M_{*}^{\sharp}$, see Proposition 3.13 , it follows that $\mathcal{I}_{N} \cap N_{*}^{\sharp}$ is dense in $N_{*}^{\sharp}$. Hence, $\lambda_{u}^{\sharp}\left(\mathcal{I}_{N} \cap N_{*}^{\sharp}\right)$ is dense in $\hat{N}_{u}$. Moreover, for $\omega \in \mathcal{I}_{N} \cap N_{*}^{\sharp}$,

$$
\begin{aligned}
\hat{\varphi}_{u}^{\natural}\left(\lambda_{u}^{\natural}(\omega)^{*} \lambda_{u}^{\natural}(\omega)\right) & =\hat{\varphi}_{c} \hat{\vartheta}_{\iota}\left(\lambda_{u}^{\natural}(\omega)^{*} \lambda_{u}^{\natural}(\omega)\right) \\
& =\hat{\varphi}_{c} \hat{\vartheta}\left(\lambda_{u}(\tilde{\omega})^{*} \lambda_{u}(\tilde{\omega})\right)=\hat{\varphi}\left((\tilde{\omega} \otimes \iota)(W)^{*}(\tilde{\omega} \otimes \iota)(W)\right)<\infty .
\end{aligned}
$$

So $\lambda_{u}^{\natural}\left(\mathcal{I}_{N} \cap N_{*}^{\sharp}\right)$ is contained in $\mathfrak{n}_{\hat{\varphi}_{u}^{\natural}}$. Thus, $\hat{\varphi}_{u}^{\natural}$ is densely defined. That $\hat{\varphi}_{u}^{\natural}$ is lower semicontinuous follows from [21, Definition 1.5]. Take any $\omega \in N_{*}^{\sharp}$ such that $\tilde{\omega} \neq 0$, which exists since all functionals in $N_{*}$ are given by restrictions of functionals in $M_{*}$, see Proposition 3.12. Then, $\hat{\varphi}_{u}^{\natural}\left(\lambda_{u}^{\natural}\left(\omega^{*} *^{\natural} \omega\right)\right)=\hat{\varphi}_{u}\left(\lambda_{u}\left(\tilde{\omega}^{*} * \tilde{\omega}\right)\right)=\hat{\varphi}_{c}\left(\lambda\left(\tilde{\omega}^{*} * \tilde{\omega}\right)\right) \neq 0$, where $\hat{\varphi}_{c}$ is the faithful left invariant weight on the reduced $\mathrm{C}^{*}$-algebraic dual $\left(\hat{M}_{c}, \hat{\Delta}_{c}\right)$.

Finally, we have to prove that $\hat{\Lambda}_{u} \iota_{u} \operatorname{maps} \hat{N}_{u}$ densely into $P_{\beta} P_{\gamma} \mathcal{H}$. The proof is similar to the one of Proposition 4.3, but since the difference is subtle we state it here. Observe that $\lambda\left(\mathcal{I} \cap M_{*}^{\sharp}\right)$ is a $\sigma$-strong-*/norm core for $\hat{\Lambda}$ as follows from [22, Proposition 2.6]. So for every $x \in \mathfrak{n}_{\hat{\varphi}_{u}^{\natural}}$ we have $\left(\hat{\vartheta} \iota_{u}\right)(x) \in \mathfrak{n}_{\varphi}$ and hence there exists a net $\left(\omega_{j}\right)_{j \in J}$ in $\mathcal{I} \cap M_{*}^{\sharp}$ such that $\lambda\left(\omega_{j}\right) \rightarrow\left(\hat{\vartheta} \iota_{u}\right)(x)$ in the $\sigma$-strong-* topology and $\xi\left(\omega_{j}\right) \rightarrow \Lambda\left(\left(\hat{\vartheta} \iota_{u}\right)(x)\right)=\hat{\Lambda}_{u}\left(\iota_{u}(x)\right)$. Consider $\left(\tilde{\omega}_{j}\right)_{j \in J}$. Then, $\lambda\left(\tilde{\omega}_{j}\right) \rightarrow\left(\hat{\vartheta} \iota_{u}\right)(x)$ and $\xi\left(\tilde{\omega}_{j}\right)=P_{\beta} P_{\gamma} \xi\left(\omega_{j}\right) \rightarrow \hat{\Lambda}_{u}\left(\iota_{u}(x)\right)$. Hence $\hat{\Lambda}_{u}\left(\iota_{u}(x)\right) \in P_{\beta} P_{\gamma} \mathcal{H}$. The range of $\hat{\Lambda}_{u} \iota_{u}$ is dense in $P_{\beta} P_{\gamma} \mathcal{H}$ by Proposition 4.2.

Next, we like to prove that $\hat{\varphi}^{\natural}$ is essentially the $W^{*}$-lift of $\hat{\varphi}_{u}^{\natural}$, see [21, Definition 1.31]. A priori this question is ill-defined, since these weights are defined on different von Neumann algebras. Indeed, $\hat{\varphi}^{\natural}$ is a weight on $\hat{N}$, which by definition acts on $P_{\gamma} \mathcal{H}$, whereas the $\mathrm{W}^{*}$-lift of $\hat{\varphi}_{u}^{\natural}$ is a weight on $\left(\hat{\pi}_{u} \iota_{u}\left(\hat{N}_{u}\right)\right)^{\prime \prime}$. Since $P_{\beta} \in \hat{M}^{\prime}$, we see that $P_{\gamma} P_{\beta} \mathcal{H}$ is an invariant subspace of $\hat{N}$. By Proposition 4.6, the von Neumann algebra $\left(\hat{\pi}_{u} \iota_{u}\left(\hat{N}_{u}\right)\right)^{\prime \prime}$ equals the restriction of $\hat{N}$ to $P_{\beta} P_{\gamma} \mathcal{H}$, i.e. $\left(\hat{\pi}_{u} \iota_{u}\left(\hat{N}_{u}\right)\right)^{\prime \prime}=\hat{N} P_{\beta}$ acting on $P_{\beta} P_{\gamma} \mathcal{H}$.

The point is that that $\hat{N}$ and $\hat{N} P_{\beta}$ are in fact isomorphic. This follows in fact from Proposition 4.3, but we give a different argument here. We claim, more precisely, that the map $\hat{N} \rightarrow \hat{N} P_{\beta}: x \mapsto x P_{\beta}$ is an isomorphism. Indeed, for any $x$ in the center of $\hat{M}, \hat{J} x \hat{J}=x$. Since $P_{\beta}=\hat{J} P_{\gamma} \hat{J}$, every projection in the center of $\hat{M}$ majorizes $P_{\beta}$ if and only if it majorizes $P_{\gamma}$. Therefore, the central supports of $P_{\beta}$ and $P_{\gamma}$ are equal. It follows from [11, Theorem 10.3.3] that $\hat{N}$ is isomorphic to $P_{\beta} \hat{N} P_{\beta}=\hat{N} P_{\beta}$, where the isomorphism is given by the map $\hat{N} \rightarrow \hat{N} P_{\beta}: x \mapsto x P_{\beta}$.

We emphasize that $\hat{N}$ will always be considered as a von Neumann algebra acting on $P_{\gamma} \mathcal{H}$, whereas if we encounter $\left(\hat{\pi}_{u} \iota_{u}\left(\hat{N}_{u}\right)\right)^{\prime \prime}$ we assume that it acts on $P_{\beta} P_{\gamma} \mathcal{H}$. The described isomorphism $\hat{N} \simeq\left(\hat{\pi}_{u} \iota_{u}\left(\hat{N}_{u}\right)\right)^{\prime \prime}$, makes the following proposition well-defined. A similar argument holds on the reduced level.

Proposition 4.7. The $W^{*}$-lifts of $\hat{\varphi}_{c}^{\natural}$ and $\hat{\varphi}_{u}^{\natural}$ to $\hat{N}$ both equal $\hat{\varphi}^{\natural}$.
Proof. We prove the proposition for $\hat{\varphi}_{u}^{\natural}$, the proof for $\hat{\varphi}_{c}^{\natural}$ is similar. Recall from Proposition 4.3 that $\left(P_{\gamma} P_{\beta} \mathcal{H},\left.\hat{\Lambda}\right|_{\hat{N} \cap \mathfrak{n}_{\hat{\varphi}}},\left.\hat{\pi}\right|_{\hat{N}}\right)$ gives the GNS-representation of $\hat{\varphi}^{\natural}$. Denote the $\mathrm{W}^{*}$-lift of $\hat{\varphi}_{u}^{\natural}$ by $\phi$. We denote its GNS-representation by $\left(\mathcal{H}_{\phi}, \pi_{\phi}, \Lambda_{\phi}\right)$. Recall from Proposition 4.6 that the GNSrepresentation of $\hat{\varphi}_{u}^{\natural}$ was given by $\left(P_{\gamma} P_{\beta} \mathcal{H}, \hat{\Lambda}_{u} \iota_{u}, \hat{\pi}_{u} \iota_{u}\right)$.

From [21, Proposition 1.32], the elements $\hat{\pi}_{u} \iota_{u}(x), x \in \mathfrak{n}_{\hat{\varphi}_{u}^{\natural}}$, form a $\sigma$-strong- $* /$ norm core for $\Lambda_{\phi}$ and we may take $\mathcal{H}_{\phi}=P_{\gamma} P_{\beta} \mathcal{H}$. If we can prove that the elements $\hat{\pi}_{u} \iota_{u}(x), x \in \mathfrak{n}_{\hat{\varphi}_{u}^{\natural}}$, also form a $\sigma$-strong- $/$ /norm core for $\left.\hat{\Lambda}\right|_{\hat{N} \cap \mathfrak{n}_{\varphi}}$, then $\phi$ and $\hat{\varphi}^{\natural}$ have identical GNS-representations and hence they are equal. So let $x \in \mathfrak{n}_{\hat{\varphi}} \cap \hat{N}$. By [22, Proposition 2.6], $\lambda\left(\mathcal{I} \cap M_{*}^{\sharp}\right)$ forms a $\sigma$ -strongly-*/norm core for $\hat{\Lambda}$. So let $\left(\omega_{j}\right)_{j \in J}$ be a net in $\mathcal{I} \cap M_{*}^{\sharp}$ such that $\lambda\left(\omega_{j}\right) \rightarrow x$ in the $\sigma$-strong-* topology and $\xi\left(\omega_{j}\right) \rightarrow \hat{\Lambda}(x)$ in norm. Then, $\xi\left(\tilde{\omega}_{j}\right)=P_{\gamma} P_{\beta} \xi\left(\omega_{j}\right)$ converges in norm. Furthermore,

$$
\lambda\left(\tilde{\omega}_{j}\right)=\left(\tilde{\omega}_{j} \otimes \iota\right)(W)=P_{\gamma}\left(\omega_{j} \otimes \iota\right)(W) P_{\gamma}=P_{\gamma} \lambda\left(\omega_{j}\right) P_{\gamma} \rightarrow P_{\gamma} x P_{\gamma}=x,
$$

where the convergence is in the $\sigma$-strong-* topology. Note that $\lambda\left(\tilde{\omega}_{j}\right)=\left(\hat{\pi}_{u} \iota_{u} \lambda_{u}^{\natural}\right)\left(\left.\omega_{j}\right|_{N}\right)$. All in all, we conclude that $\left(y_{j}\right)_{j \in J}:=\left(\left(\lambda_{u}^{\natural}\right)\left(\left.\omega_{j}\right|_{N}\right)\right)_{j \in J}$ is a net in $\mathfrak{n}_{\hat{\varphi}_{u}^{\natural}}$ such that $\hat{\pi}_{u} \iota_{u}\left(y_{j}\right)$ is $\sigma$-strong-* convergent to $x$ and $\left.\hat{\Lambda}\right|_{\hat{N} \cap \mathfrak{n}_{\varphi}}\left(\hat{\pi}_{u} \iota_{u}\left(y_{j}\right)\right)$ is convergent.

Remark 4.8. In particular, Proposition 4.3 implies that $\hat{\varphi}_{c}^{\natural}$ and $\hat{\varphi}_{u}^{\natural}$ are approximately KMSweights, see [21, Proposition 1.35].

Remark 4.9. Note that on one hand, we have a map $N_{*} \rightarrow \hat{N}: \omega \mapsto \lambda(\tilde{\omega})=(\tilde{\omega} \otimes \iota)(W)$. On the other hand, we can define a map $\hat{N}_{*} \rightarrow N: \omega \mapsto \hat{\lambda}\left(P_{\gamma} \omega P_{\gamma}\right)=(\iota \otimes \omega)\left(1 \otimes P_{\gamma}\right)\left(W^{*}\right)\left(1 \otimes P_{\gamma}\right)$. These can be considered as the spherical $L^{1}$-Fourier transforms. Since both $\varphi^{\natural}$ and $\hat{\varphi}^{\natural}$ are normal, semi-finite, faithful weights, one can proceed as in [1] to obtain a spherical $L^{p}$-Fourier transform which then is a restriction of the $L^{p}$-Fourier transform as defined in [1, Theorem 5.6].

## 5 Spherical corepresentations

We introduce the necessary terminology for corepresentations that admit vectors that are invariant under the action of a quantum subgroup. These corepresentations can be considered as spherical corepresentations.

Definition 5.1. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be a corepresentation. Then $v \in \mathcal{H}_{U}$ is called a $M_{1}$ invariant vector if

$$
(\omega \otimes \iota)\left(\left(T_{\beta} \otimes \iota\right)(U)\right) v=(\omega \otimes \iota)(U) v, \quad \forall \omega \in M_{*}
$$

We denote

$$
\mathcal{H}_{U}^{M_{1}}=\left\{v \in \mathcal{H}_{U} \mid v \text { is } M_{1} \text {-invariant }\right\} .
$$

Note that $\mathcal{H}_{U}^{M_{1}}$ is a closed subspace of $\mathcal{H}_{U}$. We denote $\operatorname{IC}\left(M, M_{1}\right)$ for the equivalence classes of irreducible corepresentations of $M$ that admit non-trivial $M_{1}$-invariant vectors. We refer to such corepresentations as spherical corepresentations. If $\left\{(\omega \otimes \iota)(U) \mathcal{H}_{U}^{M_{1}} \mid \omega \in M_{*}\right\}$ is dense in $\mathcal{H}_{U}$, then we call $U$ homogeneously cyclic. It should be clear that every irreducible corepresentation of $M$ that admits a non-trivial $M_{1}$-invariant vector is homogeneously cyclic.

Proposition 5.2. $v \in P_{\gamma} \mathcal{H}$ if and only if $v$ is $M_{1}$-invariant for $W$.

Proof. Observe that $\left(T_{\beta} \otimes \iota\right)(W)=W\left(\iota \otimes \hat{\pi}\left(\left(\varphi_{1} \otimes \iota\right)\left(W_{1}\right)\right)\right)=W\left(\iota \otimes P_{\gamma}\right)$. This yields the only if part. The other implication follows by taking a net $\left(\omega_{j}\right)_{j \in J}$ in $M_{*}$ such that $\lambda\left(\omega_{j}\right) \rightarrow 1$ in the $\sigma$-strong-* topology. Then,

$$
v \leftarrow\left(\omega_{j} \otimes \iota\right)(W) v=\left(\omega_{j} \otimes \iota\right)\left(T_{\beta} \otimes \iota\right)(W) v=\left(\omega_{j} \otimes \iota\right)(W) P_{\gamma} v \rightarrow P_{\gamma} v
$$

where the convergence is in norm.
Lemma 5.3. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be a corepresentation. Let $v \in \mathcal{H}_{U}$ and $\omega \in N_{*}$. The vector $(\tilde{\omega} \otimes \iota)(U) v$ is $M_{1}$-invariant.
Proof. This follows from the following series of equalities for which we use Lemma 3.7 (3) and $T_{\gamma}^{2}=T_{\gamma}$. For $\theta \in M_{*}$,

$$
\begin{aligned}
& (\theta \otimes \iota)\left(\left(T_{\beta} \otimes \iota\right)(U)\right)(\tilde{\omega} \otimes \iota)(U) v=\left(\theta T_{\beta} \otimes \omega T_{\beta} T_{\gamma} \otimes \iota\right) U_{13} U_{23} v \\
& \quad=\left(\theta T_{\beta} \otimes \omega T_{\beta} T_{\gamma} \otimes \iota\right)(\Delta \otimes \iota)(U) v=\left(\theta \otimes \omega T_{\beta} T_{\gamma} \otimes \iota\right)(\Delta \otimes \iota)(U) v \\
& \quad=(\theta \otimes \iota)(U)(\tilde{\omega} \otimes \iota)(U) v
\end{aligned}
$$

We mention the following two propositions in order to compare our framework with the setting of classical Gelfand pairs of groups. The proof of the first one is completely analogous to the proof of [8, Proposition II.6] or [5, Lemma 6.2.3]. For Proposition 5.5, one proves that the representation $N_{*}^{\sharp} \rightarrow B\left(\mathcal{H}_{U}^{M_{1}}\right):\left.\omega \rightarrow(\tilde{\omega} \otimes \iota)(U)\right|_{\mathcal{H}_{U}^{M_{1}}}$ is irreducible. This can be done along the lines of [8, Théorème III.1] or [5, Proposition 6.31].

Proposition 5.4. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be a corepresentation. Suppose that $U$ admits a $M_{1-}$ invariant, cyclic vector $v$. If $\operatorname{dim}\left(\mathcal{H}_{U}^{M_{1}}\right)=1$, then $U$ is irreducible.

Proposition 5.5. Assume that $\hat{N}$ is Abelian. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be an irreducible corepresentation. Then, $\operatorname{dim}\left(\mathcal{H}_{U}^{M_{1}}\right) \leq 1$.

In particular, suppose that $\hat{N}$ is Abelian, and that $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ is an irreducible corepresentation with $M_{1}$-invariant unit vector $v$. Then $x=\left(\iota \otimes \omega_{v, v}\right)(U)$ satisfies

$$
\begin{equation*}
\Delta^{\natural}(x)=x \otimes x . \tag{5.1}
\end{equation*}
$$

Hence, $x$ is a character of the convolution algebra $N_{*}^{\sharp}$. It can be considered as a quantum spherical function or quantum spherical element. The equality (5.1) allows one to derive product formulae as is done in for example $[38,39,40]$. Here we keep to a more general setting and do not assume that $\hat{N}$ is Abelian in order to include the example of $S U_{q}(1,1)_{\text {ext }}$ in Section 8.
Remark 5.6. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be a corepresentation. It follows in particular that $\mathcal{H}_{U}^{M_{1}}$ is a closed invariant subspace for the representation $N_{*}^{\sharp} \rightarrow B\left(\mathcal{H}_{U}\right): \omega \mapsto(\tilde{\omega} \otimes \iota)(U)$. By the universal property of $\hat{N}_{u}$, we see that this gives rise to a representation of $\hat{N}_{u}$ on $\mathcal{H}_{U}^{M_{1}}$. Of course, this representation can be trivial. If $U$ is homogeneously cyclic, then the corresponding representation of $\hat{N}_{u}$ is non-degenerate. Indeed, suppose that there exist $w \in \mathcal{H}_{U}^{M_{1}}$, such that for all $v \in \mathcal{H}_{U}^{M_{1}}$ and $\omega \in M_{*}^{\sharp},\langle(\tilde{\omega} \otimes \iota)(U) v, w\rangle=0$. Then, using (4) of Lemma 3.7,

$$
\begin{aligned}
0 & =\left\langle\left(\omega T_{\gamma} T_{\beta} \otimes \iota\right)(U) v, w\right\rangle=\left\langle\left(\omega T_{\gamma} \otimes \iota\right)(U) v, w\right\rangle=\left\langle v,\left(\omega T_{\gamma} \otimes \iota\right)(U)^{*} w\right\rangle \\
& =\left\langle v,\left(\overline{\omega T_{\gamma}} \otimes \iota\right)\left(U^{*}\right) w\right\rangle=\left\langle v,\left(\bar{\omega} T_{\gamma} \otimes \iota\right)(S \otimes \iota)(U) w\right\rangle \\
& =\left\langle v,(\bar{\omega} \otimes \iota)(S \otimes \iota)\left(T_{\beta} \otimes \iota\right)(U) w\right\rangle=\left\langle v,\left(\omega^{*} \otimes \iota\right)\left(T_{\beta} \otimes \iota\right)(U) w\right\rangle \\
& =\left\langle v,\left(\omega^{*} \otimes \iota\right)(U) w\right\rangle=\langle(\omega \otimes \iota)(U) v, w\rangle
\end{aligned}
$$

We see that for all $v \in \mathcal{H}_{U}^{M_{1}}$ and $\omega \in M_{*}^{\sharp},\langle(\omega \otimes \iota)(U) v, w\rangle=0$, which proves that $w=0$, since $U$ is homogeneously cyclic. Hence, every non-zero homogeneously cyclic corepresentation $U$ of $M$ gives rise to a non-degenerate representation of $\hat{N}_{u}$.

Definition 5.7. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be a homogeneously cyclic corepresentation of $M$ on a Hilbert space $\mathcal{H}_{U}$. Then, we get a representation $\pi_{U}^{\natural}$ of $\hat{N}_{u}$ determined by

$$
\pi_{U}^{\natural}:\left.\lambda_{u}^{\natural}(\omega) \mapsto(\tilde{\omega} \otimes \iota)(U)\right|_{\mathcal{H}_{U}^{M_{1}}}, \quad \omega \in N_{*}^{\sharp} .
$$

We emphasize that the representation Hilbert space of $\pi_{U}^{\natural}$ is $\mathcal{H}_{U}^{M_{1}}$.
Remark 5.8. Let $U \in M \otimes B\left(\mathcal{H}_{U}\right)$ be a homogeneously cyclic corepresentation of $M . U$ is irreducible if and only if $\pi_{U}^{\natural}$ is irreducible. Indeed, if $U$ is reducible, then clearly $\pi_{U}^{\natural}$ is reducible. The only if part follows from a computation similar to the one in Remark 5.6.

Recall that we denote $\pi_{U}$ for the representation of $\hat{M}_{u}$ given by $\lambda_{u}(\omega) \mapsto(\omega \otimes \iota)(U)$. Then, $\pi_{U}^{\natural}$ equals the restriction of $\pi_{U} \iota_{u}$ to $\mathcal{H}_{U}^{M_{1}}$.

We will need the following result for Theorem 7.1. We refer to [6] for the theory of direct integration and the definition of a fundamental sequence.

Proposition 5.9. Let $X$ be a measure space, with standard measure $\mu$. Suppose that for every $x \in X$, we have a homogeneously cyclic corepresentation $U_{x}$ of $M$ on a Hilbert space $\mathcal{H}_{x}$ such that $\left(\pi_{U_{x}}^{\natural}\right)_{x \in X}$ is a $\mu$-measurable field of representations of $\hat{N}_{u}$. Suppose that $\hat{M}_{u}$ is separable. Then, $\left(\mathcal{H}_{x}\right)_{x \in X}$ is a $\mu$-measurable field of Hilbert spaces such that $\left(\mathcal{H}_{x}^{M_{1}}\right)_{x \in X}$ is a $\mu$-measurable field of subspaces and $\left(U_{x}\right)_{x \in X}$ is a $\mu$-measurable field of corepresentations.

Proof. Let $\omega_{i} \in N_{*}^{\sharp}, i \in \mathbb{N}$ be a such that $\lambda_{u}\left(\omega_{i}\right), i \in \mathbb{N}$ is dense in $\hat{M}_{u}$. Since $\left(\pi_{U_{x}}^{\natural}\right)_{x \in X}$ is $\mu$-measurable, we have a fundamental sequence $\left(e_{x}^{j}\right)_{x \in X}, j \in \mathbb{N}$ for the $\mu$-measurable field of Hilbert spaces $\left(\mathcal{H}_{x}^{M_{1}}\right)_{x \in X}$. For $i, j \in \mathbb{N}, x \in X$, define $f_{x}^{i, j} \in \mathcal{H}_{x}$ by

$$
f_{x}^{i, j}=\left(\omega_{i} \otimes \iota\right)\left(U_{x}\right) e_{x}^{j} .
$$

We claim that $\left(f_{x}^{i, j}\right)_{x \in X}$ is a fundamental sequence for $\left(\mathcal{H}_{x}\right)_{x \in X}$. Indeed, since $U_{x}$ is homogeneously cyclic, the span of $\left(\omega_{i} \otimes \iota\right)\left(U_{x}\right) e_{x}^{j}, i, j \in \mathbb{N}$ is dense in $\mathcal{H}_{x}$. Moreover,

$$
\begin{align*}
\left\langle f_{x}^{i, j}, f_{x}^{i^{\prime}, j^{\prime}}\right\rangle & =\left\langle\left(\omega_{i} \otimes \iota\right)\left(U_{x}\right) e_{x}^{j},\left(\omega_{i^{\prime}} \otimes \iota\right)\left(U_{x}\right) e_{x}^{j^{\prime}}\right\rangle=\left\langle\left(\left(\omega_{i^{\prime}}^{*} * \omega_{i}\right) \otimes \iota\right)\left(U_{x}\right) e_{x}^{j}, e_{x}^{j^{\prime}}\right\rangle \\
& =\left\langle\left(\left(\omega_{i^{\prime}}^{*} * \omega_{i}\right)^{\sim} \otimes \iota\right)\left(U_{x}\right) e_{x}^{j}, e_{x}^{j^{\prime}}\right\rangle=\left\langle\pi_{U_{x}}^{\natural}\left(\lambda_{u}^{\natural}\left(\left.\left(\omega_{i^{\prime}}^{*} * \omega_{i}\right)\right|_{N}\right)\right) e_{x}^{j}, e_{x}^{j^{\prime}}\right\rangle, \tag{5.2}
\end{align*}
$$

where the third equality follows by a computation similar to the one in Remark 5.6. Since $\left(\pi_{U_{x}}^{\natural}\right)_{x \in X}$ is a $\mu$-measurable field of representations, we see that (5.2) is a $\mu$-measurable function of $x$. Hence $\left(f_{x}^{i, j}\right)_{x \in X}$ is a fundamental sequence. Moreover, by a similar computation as (5.2), for any $\omega \in M_{*}^{\sharp}$,

$$
\begin{aligned}
\left\langle(\omega \otimes \iota)\left(U_{x}\right) f_{x}^{i, j}, f_{x}^{i^{\prime}, j^{\prime}}\right\rangle & =\left\langle\left(\left(\omega_{i^{\prime}}^{*} * \omega * \omega_{i}\right) \otimes \iota\right)\left(U_{x}\right) e_{x}^{j}, e_{x}^{j^{\prime}}\right\rangle \\
& =\left\langle\pi_{U_{x}}^{\natural}\left(\lambda_{u}^{\natural}\left(\left.\left(\omega_{i^{\prime}}^{*} * \omega * \omega_{i}\right)\right|_{N}\right)\right) e_{x}^{j}, e_{x}^{j^{\prime}}\right\rangle,
\end{aligned}
$$

is a $\mu$-measurable function of $x$. For $\omega=\omega_{v, w}$ with $v, w \in \operatorname{Dom}\left(\hat{\nabla}^{\frac{1}{2}}\right) \cap \operatorname{Dom}\left(\hat{\nabla}^{-\frac{1}{2}}\right)$, we have $\omega \in M_{*}^{\sharp}$ by [4, Proposition 1.10]. Using [6, Proposition II.1.10 and II.2.1] it is straightforward to prove that $\left(U_{x}\right)_{x \in X}$ is a $\mu$-measurable field of operators.

Proposition 5.10. Let $U_{1}$ and $U_{2}$ be homogeneously cyclic corepresentations of $M$ on Hilbert spaces $\mathcal{H}_{1}$ and $\mathcal{H}_{2}$ respectively. Suppose that the representations of $N_{*}^{\sharp}$ given by $\pi_{i}: \omega \mapsto$ $\left.(\tilde{\omega} \otimes \iota)\left(U_{i}\right)\right|_{\mathcal{H}_{U_{i}}^{M_{1}}}, \omega \in N_{*}^{\sharp}, i \in\{1,2\}$ are equivalent. Then $U_{1}$ and $U_{2}$ are equivalent.

Proof. Let $T: \mathcal{H}_{U_{1}}^{M_{1}} \rightarrow \mathcal{H}_{U_{2}}^{M_{1}}$ be the unitary intertwiner between $\pi_{1}$ and $\pi_{2}$. Let $Q_{0}$ be the mapping

$$
\begin{aligned}
& \left\{(\omega \otimes \iota)\left(U_{1}\right) v \mid \omega \in M_{*}^{\sharp}, v \in \mathcal{H}_{U_{1}}^{M_{1}}\right\} \rightarrow\left\{(\omega \otimes \iota)\left(U_{2}\right) w \mid \omega \in M_{*}^{\sharp}, w \in \mathcal{H}_{U_{2}}^{M_{1}}\right\}, \\
& (\omega \otimes \iota)\left(U_{1}\right) v \mapsto(\omega \otimes \iota)\left(U_{2}\right) T v .
\end{aligned}
$$

This map is well-defined and isometric. Indeed, for $\omega \in M_{*}^{\sharp}$ and $v \in \mathcal{H}_{U_{1}}^{M_{1}}$,

$$
\begin{aligned}
& \left\|(\omega \otimes \iota)\left(U_{2}\right) T v\right\|^{2}=\left\langle\left(\omega^{*} * \omega \otimes \iota\right)\left(U_{2}\right) T v, T v\right\rangle=\left\langle\left(\left(\omega^{*} * \omega\right)^{\sim} \otimes \iota\right)\left(U_{2}\right) T v, T v\right\rangle \\
& =\left\langle\left(\left(\omega^{*} * \omega\right)^{\sim} \otimes \iota\right)\left(U_{1}\right) v, v\right\rangle=\left\langle\left(\omega^{*} * \omega \otimes \iota\right)\left(U_{1}\right) v, v\right\rangle=\left\|(\omega \otimes \iota)\left(U_{1}\right) v\right\|^{2},
\end{aligned}
$$

where the second equality follows from a similar calculation as in Remark 5.6. Since $U_{1}$ and $U_{2}$ are homogeneously cyclic, $Q_{0}$ is densely defined and has dense range. Let $Q: \mathcal{H}_{U_{1}} \rightarrow \mathcal{H}_{U_{2}}$ be the unitary extension of $Q_{0}$. Let $\omega, \omega_{1}, \omega_{2} \in M_{*}^{\sharp}$ and $v, w \in \mathcal{H}_{U_{1}}^{M_{1}}$,

$$
\begin{aligned}
&\langle(\omega\left.\otimes \iota)\left(U_{1}\right)\left(\omega_{1} \otimes \iota\right)\left(U_{1}\right) v,\left(\omega_{2} \otimes \iota\right)\left(U_{1}\right) w\right\rangle=\left\langle\left(\omega_{2}^{*} * \omega * \omega_{1} \otimes \iota\right)\left(U_{1}\right) v, w\right\rangle \\
& \quad=\left\langle\left(\left(\omega_{2}^{*} * \omega * \omega_{1}\right)^{\sim} \otimes \iota\right)\left(U_{1}\right) v, w\right\rangle=\left\langle\left(\left(\omega_{2}^{*} * \omega * \omega_{1}\right)^{\sim} \otimes \iota\right)\left(U_{2}\right) T v, T w\right\rangle \\
& \quad=\left\langle\left(\omega_{2}^{*} * \omega * \omega_{1} \otimes \iota\right)\left(U_{2}\right) T v, T w\right\rangle=\left\langle(\omega \otimes \iota)\left(U_{2}\right)\left(\omega_{1} \otimes \iota\right)\left(U_{2}\right) T v,\left(\omega_{2} \otimes \iota\right)\left(U_{2}\right) T w\right\rangle \\
& \quad=\left\langle(\omega \otimes \iota)\left(U_{2}\right) Q\left(\omega_{1} \otimes \iota\right)\left(U_{1}\right) v, Q\left(\omega_{2} \otimes \iota\right)\left(U_{1}\right) w\right\rangle
\end{aligned}
$$

where the second equality follows again from a similar calculation as in Remark 5.6. Since $U_{1}$ is homogeneously cyclic, this proves that $Q$ intertwines $U_{1}$ with $U_{2}$.

Note that the converse of the previous proposition is clear: if $U_{1}$ and $U_{2}$ are equivalent corepresentations, then the corresponding representations as considered in Remark 5.6 are equivalent.

Remark 5.11. Proposition 5.10 and its converse also hold on the universal level. So let $U_{1}$ and $U_{2}$ be homogeneously cyclic corepresentations of $M . \pi_{U_{1}}^{\natural}$ and $\pi_{U_{2}}^{\natural}$ are equivalent if and only if $U_{1}$ and $U_{2}$ are equivalent.

## 6 Representations of $\hat{M}_{u}, \hat{M}_{c}, \hat{N}_{u}$ and $\hat{N}_{c}$

In this section, we compare the representations of the $\mathrm{C}^{*}$-algebras defined in Sections 2 and 3. Main objective is to prove that the representations of $\hat{N}_{c}$ 'lift' to representations of $\hat{M}_{c}$.

Let us give a more elaborate discussion. There are three special types of representations within $\operatorname{Rep}\left(\hat{N}_{u}\right)$.

1. As explained in Remark 5.6, the corepresentations of $M$ give rise to representations of $\hat{N}_{u}$. Recall [18] that the corepresentations of $M$ are in 1-to-1 correspondence with non-degenerate representations of $\hat{M}_{u}$. Hence, the representations of $\hat{M}_{u}$ give rise to representations of $\hat{N}_{u}$. This correspondence can be described more directly: if $\pi$ is a representation of $\hat{M}_{u}$ on a Hilbert space $\mathcal{H}_{\pi}$, then $\pi \iota_{u}$ is the corresponding representation of $\hat{N}_{u}$ on the closure of $\left(\left(\pi \iota_{u}\right)\left(\hat{N}_{u}\right)\right) \mathcal{H}_{\pi}$. Note that by Remark 5.11 , this assignment descends to a well-defined, injective map on the equivalence classes of representations,

$$
\operatorname{Rep}\left(\hat{M}_{u}\right) \hookrightarrow \operatorname{Rep}\left(\hat{N}_{u}\right): \pi \mapsto \pi \iota_{u}
$$

2. If $\pi$ is a representation of $\hat{N}_{c}$, then $\pi \hat{\vartheta}^{\natural}$ is a representation of $\hat{N}_{u}$. These representations correspond to the representations that are weakly contained in the GNS-representation of $\hat{\varphi}_{u}^{\natural}$. Indeed this follows, since by Proposition 4.6, this GNS-representation is given by

$$
\hat{\pi}_{u} \iota_{u}=\hat{\pi}_{c} \hat{\vartheta} \iota_{u}=\hat{\pi}_{c} \hat{\vartheta}^{\natural}=\left.\hat{\pi}\right|_{\hat{N}_{c}} \hat{\vartheta}^{\natural} .
$$

Hence, every representation $\pi \hat{\vartheta}^{\natural}$, with $\pi \in \operatorname{Rep}\left(\hat{N}_{c}\right)$ is weakly contained in the GNSrepresentation of $\hat{N}_{u}$. The other way around, any representation of $\hat{N}_{u}$ that is weakly contained in the GNS-representation of $\hat{\varphi}_{u}^{\natural}$ factors through the canonical projection $\hat{\vartheta}^{\natural}$.
3. If $\pi$ is a representation of $\hat{M}_{c}$, then $\pi \hat{\vartheta}^{\natural}=\pi \hat{\vartheta}_{\iota_{u}}$ is a representation of $\hat{N}_{u}$. Here, we used that $\hat{N}_{c} \subseteq \hat{M}_{c}$.

The main results of this section will be the following. We prove that every representation of $\hat{N}_{c}$ comes from a representation of $\hat{M}_{c}$, i.e. the representations of $\hat{N}_{u}$ obtained in (2) and (3) are the same ones.

Theorem 6.1. For every non-degenerate representation $\rho \in \operatorname{Rep}\left(\hat{N}_{c}\right)$, there exists a nondegenerate representation $\pi \in \operatorname{Rep}\left(\hat{M}_{c}\right)$ on a Hilbert space $\mathcal{H}_{\pi}$ such that $\rho$ is equivalent to the restriction of $\left.\pi\right|_{\hat{N}_{c}}$ to the closure of $\pi\left(\hat{N}_{c}\right) \mathcal{H}_{\pi}$.

Proof. Let $M_{*}^{\sharp, \beta}$ denote space of functionals $\theta \in M_{*}^{\sharp}$ such that $\theta T_{\beta}=\theta$. By a similar argument as in the proof of Proposition 3.13, we see that if $\theta \in M_{*}^{\sharp}$, then $\theta T_{\beta} \in M_{*}^{\sharp, \beta}$ and $\left(\theta T_{\beta}\right)^{*}=\theta^{*} T_{\gamma}$. In particular, $\left(N_{*}^{\sharp}\right)^{\sim} \subseteq M_{*}^{\sharp, \beta}$. Note that for $\theta_{1}, \theta_{2} \in M_{*}^{\sharp, \beta}$, we have that $\left(\theta_{1}^{*} * \theta_{2}\right)^{\sim}=\theta_{1}^{*} * \theta_{2}$.

We complete $M_{*}^{\sharp, \beta}$ into a (right) Hilbert $\hat{N}_{c}$-module. For $\theta, \theta_{1}, \theta_{2} \in M_{*}^{\sharp, \beta}$ and $\omega \in N_{*}^{\sharp}$, we put

$$
\begin{align*}
& \theta \cdot \omega=\theta * \tilde{\omega} \in M_{*}^{\sharp, \beta},  \tag{6.1}\\
& \left\langle\theta_{1}, \theta_{2}\right\rangle_{N_{*}^{\sharp}}=\left.\left(\theta_{1}^{*} * \theta_{2}\right)\right|_{N} \in N_{*}^{\sharp} . \tag{6.2}
\end{align*}
$$

The fact that (6.1) is in $M_{*}^{\sharp, \beta}$ follows from Lemma 3.7. That (6.2) is in $N_{*}^{\sharp}$ follows from Proposition 3.13. This gives a right $N_{*}^{\sharp}$-module structure on $M_{*}^{\sharp, \beta}$. We will apply [27, Lemma 2.16] to get a Hilbert $\hat{N}_{c}$-module $X$. Here we consider $N_{*}^{\sharp}$ is a subalgebra of $\hat{N}_{c}$ by means of the map $\lambda^{\natural}$. To continue, note that conditions (a) and (c) of [27, Definition 2.1] are indeed satisfied. (b) follows, since for $\theta_{1}, \theta_{2} \in M_{*}^{\sharp, \beta}$ and $\omega \in N_{*}^{\sharp}$,

$$
\begin{aligned}
& \left\langle\theta_{1}, \theta_{2} \cdot \omega\right\rangle_{N_{*}^{\sharp}}=\left.\left(\theta_{1}^{*} \otimes \theta_{2} \otimes \omega\right)\left(\iota \otimes \iota \otimes T_{\beta} T_{\gamma}\right)(\Delta \otimes \iota) \Delta\right|_{N}, \\
& \left\langle\theta_{1}, \theta_{2}\right\rangle_{N_{*}^{\sharp}} *^{\natural} \omega=\left.\left(\theta_{1}^{*} \otimes \theta_{2}\right) \Delta\right|_{N} *^{\natural} \omega=\left.\left(\theta_{1}^{*} \otimes \theta_{2} \otimes \omega\right)\left(\iota \otimes \iota \otimes T_{\gamma}\right)(\Delta \otimes \iota) \Delta\right|_{N} .
\end{aligned}
$$

Since $\left.\left(\iota \otimes T_{\beta}\right) \Delta\right|_{N}=\left.\Delta T_{\beta}\right|_{N}=\left.\Delta\right|_{N}$, these expressions are equal. (d) follows, since for $\theta \in M_{*}^{\sharp, \beta}$,

$$
\lambda^{\natural}\left(\left.\left(\theta^{*} * \theta\right)\right|_{N}\right)=\lambda\left(\theta^{*} * \theta\right)=\lambda(\theta)^{*} \lambda(\theta) \geq 0,
$$

in $\hat{N}_{c}$. This defines the right Hilbert $\hat{N}_{c}$-module $X$ and we denote its norm by $\|\cdot\|_{X}$.
We are able to define an action of $M_{*}^{\sharp}$ on $X$ by means of adjointable operators which extends the convolution product on $M_{*}^{\sharp}$. Indeed, for $\omega \in M_{*}^{\sharp}$, and $\theta_{1}, \theta_{2} \in M_{*}^{\sharp, \beta}$,

$$
\left\langle\omega * \theta_{1}, \theta_{2}\right\rangle_{N_{*}^{\sharp}}=\left.\theta_{1}^{*} * \omega^{*} * \theta_{2}\right|_{N}=\left\langle\theta_{1}, \omega^{*} * \theta_{2}\right\rangle_{N_{*}^{\sharp}} .
$$

Furthermore, for $\omega \in M_{*}^{\sharp}$ and $\theta \in M_{*}^{\sharp, \beta}$,

$$
\begin{equation*}
\|\omega * \theta\|_{X}^{2}=\left\|\lambda\left(\theta^{*} * \omega^{*} * \omega * \theta\right)\right\| \leq\|\lambda(\omega)\|^{2}\left\|\lambda\left(\theta^{*} * \theta\right)\right\|=\|\lambda(\omega)\|^{2}\|\theta\|_{X}^{2} . \tag{6.3}
\end{equation*}
$$

So the action of $\omega \in M_{*}^{\sharp}$ is bounded. This allows us to extend the action of $\omega$ to an adjointable operator on $X$. We denote this operator by $L_{\omega}$. Moreover, from (6.3) we get a representation of $\hat{M}_{c}$ on the Hilbert module $X$ by means of adjointable operators. This representation is uniquely determined by $\lambda^{\natural}(\omega) \mapsto L_{\omega}, \omega \in N_{*}^{\sharp}$.

Now, let $\rho$ be a representation of $\hat{N}_{c}$ on a Hilbert space $\mathcal{H}_{\rho}$. By [27, Proposition 2.66] we get an induced representation of $\hat{M}_{c}$ on $X \otimes_{\hat{N}_{c}} \mathcal{H}_{\rho}$. Let us denote the latter by Ind $\rho$. Let $\left(a_{j}\right)_{j \in J}$ be
an approximate unit for $\hat{M}_{c}$. Then, for $\theta \in M_{*}^{\sharp, \beta}, v \in \mathcal{H}_{\rho}$, we see that $(\operatorname{Ind} \rho)\left(a_{j}\right)(\theta \otimes v) \rightarrow(\theta \otimes v)$ in $X$. So Ind $\rho$ is non-degenerate.

Note that the completion of $\left(N_{*}^{\sharp}\right) \sim \otimes_{\hat{N}_{c}} \mathcal{H}_{\rho}$ is a closed subspace of the Hilbert space $X \otimes_{\hat{N}_{c}} \mathcal{H}_{\rho}$ that is isomorphic to $\mathcal{H}_{\rho}$ via the unitary extension $T$ of

$$
\begin{equation*}
\left(N_{*}^{\sharp}\right) \sim \otimes_{\hat{N}_{c}} \mathcal{H}_{\rho} \rightarrow \mathcal{H}_{\rho}: \tilde{\omega} \otimes v \mapsto \rho\left(\lambda^{\natural}(\omega)\right) v . \tag{6.4}
\end{equation*}
$$

The map (6.4) extends unitarily since $\rho$ is non-degenerate. For $\omega, \theta \in N_{*}^{\sharp}, v \in \mathcal{H}_{\rho}$,

$$
(\operatorname{Ind} \rho)\left(\lambda^{\natural}(\omega)\right)(\tilde{\theta} \otimes v)=\tilde{\omega} * \tilde{\theta} \otimes v=\left(\omega *^{\natural} \theta\right)^{\sim} \otimes v
$$

so that $\left(N_{*}^{\sharp}\right) \sim \otimes \mathcal{H}_{\rho}$ is an invariant subspace for Ind $\rho$. We denote its closure by $Y$. Moreover, for $\omega, \theta \in N_{*}^{\sharp}$ and $v \in \mathcal{H}_{\rho}$,

$$
\begin{aligned}
T(\operatorname{Ind} \rho)\left(\lambda^{\natural}(\omega)\right)(\tilde{\theta} \otimes v) & =T((\tilde{\omega} * \tilde{\theta}) \otimes v)=T\left(\left(\omega *^{\natural} \theta\right)^{\sim} \otimes v\right) \\
& =\rho\left(\lambda^{\natural}\left(\omega *^{\natural} \theta\right)\right) v=\rho\left(\lambda^{\natural}(\omega)\right) \rho\left(\lambda^{\natural}(\theta)\right) v=\rho\left(\lambda^{\natural}(\omega)\right) T(\tilde{\theta} \otimes v),
\end{aligned}
$$

so that $T$ intertwines $\left.(\operatorname{Ind} \rho)\right|_{\hat{N}_{c}}$ restricted to the Hilbert space $Y$ with $\rho$. Finally, we claim that $Y$ equals the closure of $(\operatorname{Ind} \rho)\left(\hat{N}_{c}\right) X$. Indeed, for any $\omega \in N_{*}^{\sharp}, \theta \in M_{*}^{\sharp, \beta}$ and $v \in \mathcal{H}_{\rho}$, we see that $(\operatorname{Ind} \rho)\left(\lambda^{\natural}(\omega)\right)(\theta \otimes v)=(\tilde{\omega} * \theta) \otimes v=\left(\left.\omega *^{\natural} \theta\right|_{N}\right)^{\sim} \otimes v \in Y$. Since $\hat{N}_{c} \hat{N}_{c} \subseteq \hat{N}_{c}$ is dense, it is straightforward to prove that $Y$ equals the closure of $(\operatorname{Ind} \rho)\left(\hat{N}_{c}\right) X$ in $X$. This concludes the proof by choosing $\pi=$ Ind $\rho$.

Corollary 6.2. For every representation $\rho$ of $\hat{N}_{u}$ that factors through $\hat{\vartheta}^{\natural}$, there is a homogeneously cyclic corepresentation $U$ of $M$ such that $\rho$ is equivalent to $\pi_{U}^{\natural}$.

Remark 6.3. An essential ingredient for the proof of the quantum version of the PlancherelGodement theorem is to see to which corepresentation the GNS-map of $\hat{N}_{u}$ corresponds. Recall that the GNS-representation of $\hat{\varphi}_{u}^{\natural}$ was given by the triple $\left(P_{\gamma} P_{\beta} \mathcal{H}, \hat{\Lambda}_{u} \iota_{u}, \hat{\pi}_{u} \iota_{u}\right)$, see Proposition 4.6. Since $\hat{\pi}_{u} \iota_{u}=\left.\hat{\pi}\right|_{\hat{N}_{c}} \hat{\vartheta}^{\natural}$, we can apply Corollary 6.2. We define the closed subspace

$$
\begin{equation*}
\mathcal{E}=\overline{\left\{(\tilde{\omega} \otimes \iota)(W) P_{\beta} P_{\gamma} \mathcal{H} \mid \omega \in N_{*}^{\sharp}\right\}}=\overline{\left\{\left(\hat{\pi}_{u} \iota_{u} \lambda_{u}^{\natural}(\omega) P_{\beta} P_{\gamma} \mathcal{H} \mid \omega \in N_{*}^{\sharp}\right\}\right.} \subseteq \mathcal{H} \tag{6.5}
\end{equation*}
$$

where the closure is with respect to the norm in $\mathcal{H}$. It is clear that the representation of $\hat{N}_{u}$ that corresponds to the restriction of $W$ to $\mathcal{E}$ equals $\hat{\pi}_{u} \iota_{u}$.

We use the notation $\operatorname{IR}\left(\hat{M}_{u}, M_{1}\right)$ to denote the irreducible representations $\pi$ of $\hat{M}_{u}$ such that the representation $\pi \iota_{u}$ is non-trivial. Under the 1-1 correspondence between $\operatorname{IR}\left(\hat{M}_{u}\right)$ and $\operatorname{IC}(M)$, see [18], we see from Remark 5.11 and the remarks following Definition 5.7 that $\operatorname{IR}\left(\hat{M}_{u}, M_{1}\right)$ corresponds to $\operatorname{IC}\left(M, M_{1}\right)$. Let $\operatorname{IR}\left(\hat{M}_{c}, M_{1}\right)$ denote the irreducible representations of $\hat{M}_{c}$ such that the restriction to $\hat{N}_{c}$ is non-trivial. We find the following diagram of inclusions.


Note that the map $\operatorname{IR}\left(\hat{M}_{u}, M_{1}\right) \hookrightarrow \operatorname{IR}\left(\hat{N}_{u}\right)$ indeed maps into the representations of $\hat{N}_{u}$ that are irreducible, c.f. Remark 5.8. Hence, also the vertical inclusion on the right hand side of (6.6) preserves irreducibility.

The example in Section 8 shows that the inclusion $\operatorname{IR}\left(\hat{M}_{c}, M_{1}\right) \hookrightarrow \operatorname{IR}\left(\hat{M}_{u}, M_{1}\right)$ is not surjective. We briefly comment on the fact that also the inclusion $\operatorname{IR}\left(\hat{M}_{u}, M_{1}\right) \hookrightarrow \operatorname{IR}\left(\hat{N}_{u}\right)$ is generally
not surjective. This is a consequence of the fact that there are Lie groups $G$ with compact subgroup $K$ for which the there are non-unitary representations whose restriction to the bi- $K$ invariant functions forms a representation, i.e. a homomorphism that preserves the $*$-operation, whereas the representation of all $L^{1}$-functions on $G$ does not preserve the *. This happens for example for $S L(2, \mathbb{R})$, see [44, Example 1.1.2 on p. 37 and p. 40]. This shows that the induction argument contained in the proof of Theorem 6.1 does not work in general on the universal level.

Remark 6.4. Assume the map $\iota_{u}: \hat{N}_{u} \rightarrow \hat{M}_{u}$ to be injective. Then by general C*-algebra theory, it is isometric. With this additional assumption Theorem 6.1 holds on the universal level. So for every $\rho \in \operatorname{Rep}\left(\hat{N}_{u}\right)$, there is a representation $\pi \in \operatorname{Rep}\left(\hat{M}_{u}\right)$ on a Hilbert space $\mathcal{H}_{\pi}$, such that $\rho$ is equivalent to the restriction of $\pi \iota_{u}$ to $\left(\pi \iota_{u}\left(\hat{N}_{u}\right)\right) \mathcal{H}_{\pi}$.

The proof is completely analogous to the one of Theorem 6.1, where one takes the universal norm instead of the reduced norm on $N_{*}^{\sharp}$. The injectivity of $\iota_{u}$ plays an essential role at two places. First of all, the injectivity of $t_{u}$ can be used to prove positivity of the inner product (6.2), since in this case an element in $\hat{N}_{u}$ is positive if and only if it is positive in $\hat{M}_{u}$. Secondly, the universal analogue of (6.3) can be recovered from the injectivity of $\iota_{u}$, since for $\omega \in M_{*}^{\sharp}$ and $\theta \in M_{*}^{\sharp, \beta}$,

$$
\begin{aligned}
\left\|\lambda_{u}^{\natural}\left(\theta^{*} * \omega^{*} * \omega * \theta\right)\right\|_{u}^{\natural} & =\left\|\lambda_{u}\left(\theta^{*} * \omega^{*} * \omega * \theta\right)\right\|_{u} \\
& \leq\left\|\lambda_{u}\left(\omega^{*} * \omega\right)\right\|_{u}\left\|\lambda_{u}\left(\theta^{*} * \theta\right)\right\|_{u}=\left\|\lambda_{u}(\omega)\right\|_{u}\left\|_{u}^{\natural}\left(\theta^{*} * \theta\right)\right\|_{u}^{\natural} .
\end{aligned}
$$

The rest of the prove of Theorem 6.1 can be copied mutatis mutandis.

## 7 A quantum group analogue of the Plancherel-Godement theorem

Here we prove a decomposition theorem that may be considered as a locally compact quantum group version of the Plancherel-Godement theorem as can be found in [8, Théorème IV.2]. The proof is different from the one given in [8] and follows the line of the Plancherel theorem as proved by Desmedt [4]. We show that the C*-algebra $\hat{N}_{u}$ together with the weight $\varphi_{u}^{\natural}$ that we introduced and studied so far fit into the framework of [4, Theorem 3.4.5]. Then we use Theorem 6.1 to translate the results in terms of corepresentations of $M$ that admit a $M_{1}$-invariant vector.

Recall that we defined the space $\mathcal{E}$ in (6.5). Let $\mathcal{L}$ be any Hilbert space and let $\mathcal{L}_{0} \subseteq \mathcal{L}$ be a closed subspace. We denote the conjugate Hilbert space of $\mathcal{L}$ by $\overline{\mathcal{L}}$. Note that the space $\mathcal{L} \otimes \overline{\mathcal{L}_{0}}$ can canonically be identified with the Hilbert-Schmidt operators in $B\left(\mathcal{L}_{0}, \mathcal{L}\right)$. We denote the latter space by $B_{2}\left(\mathcal{L}_{0}, \mathcal{L}\right)$. For results on direct integration, we refer to [6, 23] and [25]. In particular, we will use [23, Theorem 1.10] implicitly several times. If $A$ and $B$ are unbounded operators such that $A B$ is closable, we denote $A \cdot B$ for the closure of $A B$.

Theorem 7.1 (Plancherel-Godement). Let $(M, \Delta)$ be a unimodular locally compact quantum group and let $\left(M_{1}, \Delta_{1}\right)$ be a compact (closed) quantum subgroup. Let $\hat{N}$ and $\hat{N}_{u}$ be the von Neumann algebra and $C^{*}$-algebra as defined earlier in this section. Suppose that $\hat{N}$ is a type $I$ von Neumann algebra and that $\hat{N}_{u}$ and $\hat{M}_{u}$ are separable.

Then, there exists a standard measure $\mu^{M_{1}}$ on $\mathrm{IC}\left(M, M_{1}\right)$, a $\mu^{M_{1}}$-measurable field of Hilbert spaces $\left(\mathcal{H}_{U}\right)_{U \in \operatorname{IC}\left(M, M_{1}\right)}$ of which $\left(\mathcal{H}_{U}^{M_{1}}\right)_{U \in \operatorname{IC}\left(M, M_{1}\right)}$ forms a measurable field of subspaces, a measurable field of self-adjoint, strictly positive operators $\left(D_{U}^{M_{1}}\right)_{U \in \operatorname{IC}\left(M, M_{1}\right)}$ acting on $\mathcal{H}_{U}^{M_{1}}$ and an isomorphism $\mathcal{Q}^{M_{1}}$ of $\mathcal{E}$ onto $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U} \otimes \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$ with properties:

1. For $\omega \in \mathcal{I}_{N}$ and $\mu^{M_{1}}$-almost all $U \in \operatorname{IC}\left(M, M_{1}\right)$, the operator $(\tilde{\omega} \otimes \iota)(U)\left(D_{U}^{M_{1}}\right)^{-1}$ is bounded and $(\tilde{\omega} \otimes \iota)(U) \cdot\left(D_{U}^{M_{1}}\right)^{-1}$ is in $B_{2}\left(\mathcal{H}_{U}^{M_{1}}\right)$.
2. For $\omega_{1}, \omega_{2} \in \mathcal{I}_{N}$, we have

$$
\begin{aligned}
& \left\langle\xi\left(\tilde{\omega}_{1}\right), \xi\left(\tilde{\omega}_{2}\right)\right\rangle \\
& \quad=\int_{\mathrm{IC}\left(M, M_{1}\right)} \operatorname{Tr}\left(\left(\left(\tilde{\omega}_{2} \otimes \iota\right)(U) \cdot\left(D_{U}^{M_{1}}\right)^{-1}\right)^{*}\left(\left(\tilde{\omega}_{1} \otimes \iota\right)(U) \cdot\left(D_{U}^{M_{1}}\right)^{-1}\right)\right) d \mu^{M_{1}}(U)
\end{aligned}
$$

and we let $\mathcal{Q}_{0}^{M_{1}}: P_{\beta} P_{\gamma} \mathcal{H} \rightarrow \int_{\operatorname{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U}^{M_{1}} \otimes \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$ be the isometric extension of

$$
\begin{aligned}
& \hat{\Lambda}\left(\lambda\left(\tilde{\mathcal{I}}_{N}\right)\right) \rightarrow \int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} B_{2}\left(\mathcal{H}_{U}^{M_{1}}\right) d \mu(U): \\
& \xi(\tilde{\omega}) \mapsto \int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus}(\tilde{\omega} \otimes \iota)(U) \cdot\left(D_{U}^{M_{1}}\right)^{-1} d \mu^{M_{1}}(U)
\end{aligned}
$$

3. $\mathcal{Q}_{0}^{M_{1}}$ intertwines $\hat{\pi}_{u} \iota_{u}$ and $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \pi_{U}^{\natural} \otimes 1 \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$.
4. $\mathcal{Q}^{M_{1}}$ intertwines the restriction of $W$ to $\mathcal{E}$ with $\int_{\operatorname{IC}\left(M, M_{1}\right)}^{\oplus} U \otimes 1 \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$. Moreover, the restriction of $\mathcal{Q}^{M_{1}}$ to $P_{\beta} P_{\gamma} \mathcal{H} \rightarrow \int_{\operatorname{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U}^{M_{1}} \otimes \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$ equals $\mathcal{Q}_{0}^{M_{1}}$.
5. Assume moreover that $\hat{M}$ is a type I von Neumann algebra and that $\hat{M}_{u}$ is separable. Let $\mu, D_{U}, \mathcal{Q}$ be defined as in [4, Theorem 3.4.1]. Then $\mathcal{H}_{U}^{M_{1}}$ is an invariant subspace for $\mu$-almost all $D_{U}$ and $\mathrm{IC}\left(M, M_{1}\right)$ is a $\mu$-measurable subset of $\mathrm{IC}(M)$. If one takes:

- $\mu^{M_{1}}$ equal to the restriction of $\mu$ to $\operatorname{IC}\left(M, M_{1}\right)$;
- $D_{U}^{M_{1}}$ equal to the restriction of $D_{U}$ to $\mathcal{H}_{U}^{M_{1}}$;
- $\mathcal{Q}^{M_{1}}$ the restriction of $\mathcal{Q}$ to $\mathcal{E}$.

Then, $\mu^{M_{1}}, D_{U}^{M_{1}}, \mathcal{Q}^{M_{1}}$ satisfy the properties (1)-(4).
Proof. By Propositions 4.6 and $4.7, \hat{\varphi}_{u}^{h}$ is a proper approximate KMS-weight. Therefore, we can apply [4, Theorem 3.4.5], so that we obtain a measure $\mu^{M_{1}}$ on $\operatorname{IR}\left(\hat{N}_{u}\right)$, a measurable field of Hilbert spaces $\left(\mathcal{K}_{\sigma}^{M_{1}}\right)_{\sigma \in \operatorname{IR}\left(\hat{N}_{u}\right)}$, a measurable field of representations $\left(\pi_{\sigma}\right)_{\sigma \in \operatorname{IR}\left(\hat{N}_{u}\right)}$, a measurable field of self-adjoint, strictly positive operators $\left(D_{\sigma}^{M_{1}}\right)_{\sigma \in \operatorname{IR}\left(\hat{N}_{u}\right)}$ and an isomorphism $\mathcal{Q}_{0}^{M_{1}}$ of $P_{\gamma} P_{\beta} \mathcal{H}$ onto $\int_{\operatorname{IR}\left(\hat{N}_{u}\right)}^{\oplus} \mathcal{K}_{\sigma}^{M_{1}} \otimes \overline{\mathcal{K}_{\sigma}^{M_{1}}} d \mu^{M_{1}}(\sigma)$ satisfying the properties of this theorem.

Let $\rho \in \operatorname{IR}\left(\hat{N}_{u}\right)$ be in the support of $\mu^{M_{1}}$. We claim that $\pi_{\rho}$ is weakly contained in $\hat{\pi}_{u} \iota_{u}$. Suppose that this is not the case, so that there exists $x \in \hat{N}_{u}$ such that $\pi_{\rho}(x) \neq 0$ but $\hat{\pi}_{u} \iota_{u}(x)=0$. Let $X=\left\{\sigma \in \operatorname{IR}\left(\hat{N}_{u}\right) \mid \pi_{\sigma}(x) \neq 0\right\}$. Then $X$ is an open neighbourhood of $\rho$. Moreover, it follows form [4, Theorem 3.4.5] that $\mathcal{Q}_{0}^{M_{1}}$ intertwines $\hat{\pi}_{u} \iota_{u}$ with $\int_{\operatorname{IR}\left(\hat{N}_{u}\right)}^{\oplus} \pi_{\sigma} \otimes \overline{1_{\mathcal{K}_{\sigma}}} d \mu^{M_{1}}(\sigma)$ from which it follows that $\mu^{M_{1}}(X)=0$. This contradicts the fact that $\rho$ is in the support of $\mu^{M_{1}}$, so $\pi_{\rho}$ is weakly contained in $\hat{\pi}_{u} \iota_{u}$.

Since $\hat{\pi}_{u} \iota_{u}=\left.\hat{\pi}\right|_{\hat{N}_{c}} \hat{\vartheta}^{h}$, we see that $\rho$ is in $\operatorname{IR}\left(\hat{N}_{c}\right)$, where $\operatorname{IR}\left(\hat{N}_{c}\right)$ is considered as a subset of $\operatorname{IR}\left(\hat{N}_{u}\right)$ by the inclusion (6.6). Now we use Theorem 6.1 to identify $\operatorname{IR}\left(\hat{N}_{c}\right)$ with $\operatorname{IR}\left(\hat{M}_{c}, M_{1}\right)$, which we consider as a subspace of $\operatorname{IC}\left(M, M_{1}\right)$ by (6.6). We consider $\mu^{M_{1}}$ as a measure on $\operatorname{IC}\left(M, M_{1}\right)$ by defining the complement of $\operatorname{IR}\left(\hat{N}_{c}\right)$ in $\operatorname{IC}\left(M, M_{1}\right)$ to be negligible. Let $U_{\sigma} \in \operatorname{IC}\left(M, M_{1}\right)$ denote the corepresentation corresponding to $\sigma \in \operatorname{IR}\left(\hat{N}_{c}\right)$. So, $\pi_{\sigma}=\pi_{U_{\sigma}}^{\natural}$. We write $D_{U_{\sigma}}^{M_{1}}$ for $D_{\sigma}^{M_{1}}$ and set $D_{U}=0$ for $U \in \mathrm{IC}\left(M, M_{1}\right)$ not in the support of $\mu^{M_{1}}$. We denote $\mathcal{H}_{U}$ for the corepresentation Hilbert space of $U \in \mathrm{IC}\left(M, M_{1}\right)$, and we get $\mathcal{H}_{U_{\sigma}}^{M_{1}}=\mathcal{K}_{\sigma}^{M_{1}}$. Therefore, since the support of $\mu$ is contained in $\operatorname{IR}\left(\hat{N}_{c}\right)$, we see that $\mathcal{Q}_{0}^{M_{1}}$ is a map from $P_{\beta} P_{\gamma} \mathcal{H} \rightarrow \int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U}^{M_{1}} \otimes \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$.
(1). It follows from Corollary 6.2, Remark 5.6 and the properties of $D_{\sigma}^{M_{1}}$ described in (1) of [4, Theorem 3.4.5], that for $\omega \in \mathcal{I}_{N}$, the operator $(\tilde{\omega} \otimes \iota)(U)\left(D_{U}^{M_{1}}\right)^{-1}$ is bounded and its closure is Hilbert-Schmidt for $\mu^{M_{1}}$-almost every $U \in \operatorname{IC}\left(M, M_{1}\right)$.
(2) and (3). We make two observations. First note that by Proposition 4.6, for $\omega \in \mathcal{I}_{N}$, $\hat{\Lambda}_{u} \iota_{u}\left(\lambda_{u}^{\natural}(\omega)\right)=\xi(\tilde{\omega})$. Secondly, we have proved that for every $\rho$ in the support of $\mu^{M_{1}}$, there is a $U \in \operatorname{IC}\left(M, M_{1}\right)$, such that $\pi_{\rho}=\pi_{U}^{\natural}$. Then (2) of [4, Theorem 3.4.5] yields (2) of the present theorem. The second observation also yields (3). Note that by Proposition 5.9, we see that $\left(\mathcal{H}_{U}\right)_{U \in \operatorname{IC}\left(M, M_{1}\right)}$ is a measurable field of Hilbert spaces of which $\left(\mathcal{H}_{U}^{M_{1}}\right)_{U \in \operatorname{IC}\left(M, M_{1}\right)}$ forms a measurable field of subspaces. Here we used that $\hat{M}_{u}$ is separable.

To prove (4), we make the following observations. First of all, by Remark 6.3, we see that $\hat{\pi}_{u} \iota_{u}=\pi_{W_{\mathcal{E}}}^{\natural}$, where $W_{\mathcal{E}}$ denotes the restriction of the multiplicative unitary $W$ to the Hilbert space $\mathcal{E}$. Secondly,

$$
\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \pi_{U}^{\natural} \otimes 1 \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)=\pi_{\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\natural} U \otimes 1 \overline{\mathcal{H}_{U}^{M_{1}}} d \mu(U),},
$$

where we use Proposition 5.9 to infer that the direct integral on the right hand side exists. Hence, by Remark 5.11 we see that $W_{\mathcal{E}}$ and $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} U \otimes 1 \overline{\mathcal{H}_{U}^{M_{1}}} d \mu^{M_{1}}(U)$ are equivalent. Moreover, we define $\mathcal{Q}^{M_{1}}$ to be the intertwiner as constructed in the proof of Proposition 5.10. Then, $\mathcal{Q}^{M_{1}}$ satisfies (4).

We now prove (5). The proof of Proposition 3.16 shows that $\hat{\sigma}_{t}\left(P_{\gamma}\right)=P_{\gamma}$. Write $D=$ $\int_{\mathrm{IC}(M)}^{\oplus} D_{U} d \mu(U)$. Since $D^{-2}$ is the Radon-Nikodym derivative of $\hat{\varphi}$ with respect to a trace on $\hat{M}$, see the proof of [4, Theorem 3.4.5], we have $D \eta \hat{M}$ and $\hat{\sigma}_{t}(x)=D^{-2 i t} x D^{2 i t}$. Thus, $P_{\gamma}$ commutes with $D^{i t}$ for all $t \in \mathbb{R}$. Since $P_{\gamma} \in \hat{M} \simeq \int_{\operatorname{IC}(M)}^{\oplus} B\left(\mathcal{H}_{U}\right) d \mu(U)$, we have a direct integral decomposition $P_{\gamma}=\int_{\mathrm{IC}(M)}^{\oplus}\left(P_{\gamma}\right)_{U} d \mu(U)$ and $\left(P_{\gamma}\right)_{U}$ commutes with $D_{U}^{i t}$ for $\mu$-almost all $U \in \operatorname{IC}(M)$.
$P_{\gamma}$ is the projection of $\mathcal{H}$ onto $\mathcal{H}^{M_{1}}$, see Proposition 5.2. A vector $v=\int_{\mathrm{IC}(M)}^{\oplus} v_{U} d \mu(U)$ is $M_{1}$-invariant for $W$ if and only if $v_{U}$ is $M_{1}$-invariant for $\mu$-almost all $U \in \operatorname{IC}(M)$, as follows directly from Definition 5.1. Hence $\left(P_{\gamma}\right)_{U}$ is the projection of $\mathcal{H}_{U}$ onto $\mathcal{H}_{U}^{M_{1}}$ for $\mu$-almost all $U \in \operatorname{IC}(M)$. We record three conclusions:
(P1) $\mathcal{H}_{U}^{M_{1}}$ is an invariant subspace of $D_{U}$ for $\mu$-almost all $U \in \operatorname{IC}(M)$;
(P2) $\operatorname{IC}\left(M, M_{1}\right)$ is a $\mu$-measurable subset of $\operatorname{IC}(M)$ by [6, Proposition II.1.1 (i)];
(P3) The image of $P_{\gamma} \mathcal{H}$ under $\mathcal{Q}$ equals $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U}^{M_{1}} \otimes \overline{\mathcal{H}_{U}} d \mu(U)$.
For the choice of $\mu^{M_{1}}, D_{U}^{M_{1}}$ and $\mathcal{Q}^{M_{1}}$ made in (5), properties (1) and (2) follow from the properties of $\mu, D_{U}$ and $\mathcal{Q}$ as described in [4, Theorem 3.4.1] using Proposition 4.2. By [4, Theorem 3.4.1], $\mathcal{Q}$ intertwines $W$ with $\int_{\operatorname{IC}(M)}^{\oplus} U \otimes 1_{\overline{\mathcal{H}_{U}}} d \mu(U)$. Using (P3) together with [4, Theorem 3.4.5 (3)] and the fact that $P_{\beta}=\hat{J} P_{\gamma} \hat{J}$, we see that $\mathcal{Q}$ restricts to a unitary map from $P_{\beta} P_{\gamma} \mathcal{H}$ to $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U}^{M_{1}} \otimes \overline{\mathcal{H}_{U}^{M_{1}}} d \mu(U)$. Hence $\mathcal{Q}$ restricts to a unitary map from $\mathcal{E}$ to $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} \mathcal{H}_{U} \otimes \overline{\mathcal{H}_{U}^{M_{1}}} d \mu(U)$, which then intertwines the restriction of $W$ to $\mathcal{E}$ with $\int_{\mathrm{IC}\left(M, M_{1}\right)}^{\oplus} U \otimes 1 \overline{\mathcal{H}_{U}^{M_{1}}} d \mu(U)$. This proves (4), from which (3) follows by the construction in Remarks 5.6 and 6.3.

Remark 7.2. If $\Delta^{\natural}$ is cocommutative, then $\hat{N}$ is Abelian. By Proposition 5.5 we see that for all $U \in \operatorname{IC}(M), \operatorname{dim}\left(\mathcal{H}_{U}^{M_{1}}\right) \leq 1$. Hence the operators $D_{U}^{M_{1}}$ are scalars. We may assume that $D_{U}^{M_{1}}=1$ by replacing the measure $\mu^{M_{1}}$ if necessary. In particular, we see that for a classical Gelfand pair $(G, K)$, the map $\mathcal{Q}_{0}^{M_{1}}$ is the spherical Fourier transform.

Remark 7.3. The support of $\mu^{M_{1}}$ is given by $\operatorname{IR}\left(\hat{N}_{c}\right)$. Here, $\operatorname{IR}\left(\hat{N}_{c}\right)$ is a subspace of $\operatorname{IC}\left(M, M_{1}\right)$ as in (6.6). The prove can be done in exactly the same manner as [4, Theorem 3.4.8], see also [7, Proposition 8.6.8]. Note that in the course of the proof of Theorem 7.1 we have already proved that the support of $\mu^{M_{1}}$ is contained in $\operatorname{IR}\left(\hat{N}_{c}\right)$.

Remark 7.4. As pointed out in remark [2, Remark 3.3], the Theorem 7.1 also holds if one assumes that $\hat{N}_{c}$ and $\hat{M}_{c}$ are separable, instead of $\hat{N}_{u}$. Moreover, note that if $\hat{M}_{c}$ is separable, then so is $\hat{N}_{c} \subseteq \hat{M}_{c}$. If $\hat{M}$ is a type I von Neumann algebra, then so is $\hat{N}=P_{\gamma} \hat{M} P_{\gamma}$. In particular, if $\hat{M}$ is type I and $\hat{M}_{c}$ is separable, then then the result of Theorem 7.1 holds for any closed quantum subgroup of $(M, \Delta)$.

## 8 Example: $S U_{q}(1,1)_{\text {ext }}$

Let $(M, \Delta)$ be the quantum group analogue of the normaliser of $S U(1,1)$ in $S L(2, \mathbb{C})$ as introduced in the operator algebraic framework in [13] and further studied in [10]. In this section we identify the circle as a closed quantum subgroup of $(M, \Delta)$. We show that for this pair the map $\Delta^{\natural}$ defined in (3.3) is not cocommutative. Moreover, the von Neumann subalgebra $N$ of $M$ consisting of bi-invariant elements is not commutative. However, we show how the von Neumann algebras $N$ and $\hat{N}$ as defined in the previous section can be equipped with a $\mathbb{Z}_{2}$-grading. The grading allows us to derive similar results as for (quantum) Gelfand pairs. In particular, we make the Fourier transform explicit and show that it preserves the $\mathbb{Z}_{2}$-grading. Moreover, we derive product formulae for little $q$-Jacobi functions appearing as matrix coefficients of corepresentations which admit invariant vectors.

Remark 8.1. In [2, Propositions B. 2 and B.3] it is proved that $S U_{q}(1,1)_{\text {ext }}$ satisfies the hypotheses of the Plancherel theorem, [4, Theorem 3.4.1]. By Remark 7.4 we can also apply Theorem 7.1.

Notation 8.2. In this section we adopt all the notational conventions made in [10]. In particular, in this section we write $\mathcal{K}$ instead of $\mathcal{H}$ to denote the GNS-space of $(M, \Delta)$. For $z \in \mathbb{C}$, we denote $\mu(z)=\left(z+z^{-1}\right) / 2$. For a set $X$ and $x \in X$, we will write $\delta_{x}$ for the function on $X$ that equals 1 in $x$ and 0 elsewhere. It should always be clear from the context what the domain of this function is. Recall in particular that $I_{q}=-q^{\mathbb{N}} \cup q^{\mathbb{Z}}$, where $\mathbb{N}$ denotes the natural numbers excluding 0 .

For the reader's convenience, we summarize the necessary results on the corepresentation theory of $(M, \Delta)$ from [10]. Let $W$ be the multiplicative unitary of $(M, \Delta)$ and recall [10] that we have a direct integral decomposition

$$
\begin{equation*}
W=\bigoplus_{p \in q^{\mathbb{Z}}}\left(\int_{[-1,1]}^{\oplus} W_{p, x} d x \oplus \bigoplus_{x \in \sigma_{d}\left(\Omega_{p}\right)} W_{p, x}\right) \tag{8.1}
\end{equation*}
$$

Here $\sigma_{d}\left(\Omega_{p}\right)$ is the discrete spectrum of the Casimir operator [10, Definition 4.5, Theorem 4.6] restricted to the subspace given in [10, Theorem 5.7]. $W_{p, x}$ is a corepresention that is a direct sum of at most 4 irreducible corepresentations [10, Propositions 5.3 and 5.4 ]. We will simply write $W=\int^{\oplus} W_{p, x} d(p, x)$ for the integral decomposition (8.1).

The corepresentations in the continuous part of the decomposition are called principal series corepresentations, the corepresentations that appear as a direct summand are called the discrete series corepresentations. In addition the complementary series corepresentations $W_{p, x}, x \in$ $(\mu(-q),-1) \cup(1, \mu(q))$, are defined by analytic continuation of matrix coefficients, see [10, Section 10.3]. We mention that it remains unproved that these make up all the corepresentations.

Using the notation of [10, Sections 10.2 and 10.3], an orthonormal basis for the corepresentation Hilbert space $\mathcal{L}_{p, x}$ of the principal and complementary series $W_{p, x}$ is given by the vectors

$$
\begin{equation*}
e_{m}^{\varepsilon, \eta}(p, x), \quad \varepsilon, \eta \in\{-,+\}, \quad m \in \mathbb{Z} \tag{8.2}
\end{equation*}
$$

For the discrete series corepresentation $W_{p, x}$ a subset of the vectors (8.2) gives a basis for the corepresentation space $\mathcal{L}_{p, x}$, see [10, Proposition 5.2]. For our purposes, it is convenient to use the following notational convention.

Notation 8.3. Let $p \in q^{\mathbb{Z}}, x \in \sigma_{d}\left(\Omega_{p}\right)$, so that $W_{p, x}$ is a discrete series corepresentation. We denote $e_{m}^{\varepsilon, \eta}(p, x), \varepsilon, \eta \in\{-,+\}, m \in \mathbb{Z}$ for the zero vector in case $e_{m}^{\varepsilon, \eta}(p, x)$ is not in one of the sets defined in cases 1-3 of [10, Proposition 5.2]. In particular, the non-zero vectors of (8.2) form an orthonormal basis of $\mathcal{L}_{p, x}$.

For any $p \in q^{\mathbb{Z}}, x \in(\mu(-q), \mu(q)) \cup \sigma_{d}\left(\Omega_{p}\right)$, we set $f_{m}^{\varepsilon, \eta}(p, x)=e_{m-\frac{1}{2} \chi(p)}^{\varepsilon, \eta}(p, x), \varepsilon, \eta \in\{-,+\}$, where $m \in \mathbb{Z}$ if $p \in q^{2 \mathbb{Z}}$ and $m \in \frac{1}{2}+\mathbb{Z}$ if $p \in q^{1+2 \mathbb{Z}}$. We remind the reader that the direct integrals over of the vectors $f_{m}^{\varepsilon, \eta}(p, x)$ are vectors in $\mathcal{K}$. Recall that for $x \in[-1,1]$ the actions of the (unbounded) generators of the dual quantum group, see [10, equation (92)], are given by

$$
\begin{align*}
& K f_{m}^{\varepsilon, \eta}(p, x)=q^{m} f_{m}^{\varepsilon, \eta}(p, x), \\
& \left(q^{-1}-q\right) E f_{m}^{\varepsilon, \eta}(p, x)=q^{-\frac{1}{2}-m}\left|1+\varepsilon \eta q^{2 m+1} e^{i \theta}\right| f_{m+1}^{\varepsilon, \eta}(p, x), \\
& U_{0}^{+-} f_{m}^{\varepsilon, \eta}(p, x)=\eta(-1)^{v(p)} f_{m}^{\varepsilon,-\eta}(p, x), \\
& U_{0}^{-+} f_{m}^{\varepsilon, \eta}(p, x)=\varepsilon \eta^{\chi(p)}(-1)^{m-\frac{1}{2} \chi(p)} f_{m}^{-\varepsilon, \eta}(p, x), \tag{8.3}
\end{align*}
$$

where $\theta$ is such that $x=\mu\left(e^{i \theta}\right)$. Similar expressions can be obtained for the discrete and complementary series corepresentations from the expressions in [10, Lemma 10.1 and Section 10.3].

Remark 8.4. The corepresentations appearing in (8.1) are not mutually inequivalent. We give a complete list of equivalences in Proposition 8.11.
Remark 8.5. For every $p \in q^{\mathbb{Z}}, x \in \mu\left(-q^{2 \mathbb{Z}+1} p \cup q^{2 \mathbb{Z}+1} p\right)$, one can define a corepresentation $W_{p, x}$ by defining the action of the generators of $\hat{M}$ by means of [10, Lemma 10.1]. Since the actions of the generators of $W_{p r, x}$ are equal for any $r \in q^{4 \mathbb{Z}}$, these corepresentations are all equivalent. Using [10, Proposition 5.2] one can check that every such corepresentation is equivalent to at least one corepresentation in the decomposition (8.1). Since every discrete series corepresentation is infinite dimensional, it occurs infinitely many times in the decomposition (8.1) by the Plancherel theorem [4, Theorem 3.4.1], or see Proposition 8.11 below for a direct proof. Therefore, we see that

$$
W \simeq \bigoplus_{p \in q^{\mathbb{Z}}}\left(\int_{[-1,1]}^{\oplus} W_{p, x} d x \oplus \bigoplus_{x \in \mu\left(-q^{2 Z+1} p \cup q^{2 Z}+1_{1} p\right)} W_{p, x}\right) .
$$

We will also need the following expressions for the matrix coefficients. By [10, Lemma 10.9 and Section 10.3], for $p \in q^{\mathbb{Z}}, x \in(\mu(-q), \mu(q))$,

$$
\begin{align*}
& \left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right) f_{m_{0}, p_{0}, t_{0}}  \tag{8.4}\\
& =C\left(\varepsilon \eta x ; m^{\prime}-\frac{1}{2} \chi(p), \varepsilon^{\prime}, \eta^{\prime} ; \varepsilon \varepsilon^{\prime}\left|p_{0}\right| q^{-m-m^{\prime}}, p_{0}, m-m^{\prime}\right) \delta_{\operatorname{sgn}\left(p_{0}\right), \eta \eta^{\prime}} f_{m_{0}-m+m^{\prime}, \varepsilon \varepsilon^{\prime}\left|p_{0}\right| q^{-m-m^{\prime}, t_{0}}} .
\end{align*}
$$

We refer to [10] for the precise definitions of the function $C(\cdot)$ in terms of basic hypergeometric series. (8.4) also holds for the discrete series corepresentations.

## The diagonal subgroup

Let $\mathbb{T}$ denote the circle group and let $\left(L^{\infty}(\mathbb{T}), \Delta_{\mathbb{T}}\right)$ be the usual locally compact quantum group associated with $\mathbb{T}$ with dual quantum group $\left(L^{\infty}(\mathbb{Z}), \Delta_{\mathbb{Z}}\right)$. We identify $\left(L^{\infty}(\mathbb{T}), \Delta_{\mathbb{T}}\right)$ as a closed quantum subgroup of $(M, \Delta)$. Recall from [10, Definition 4.3] that the spectrum of the operator $K$ equals $0 \cup q^{\frac{1}{2} \mathbb{Z}}$.

Definition 8.6. We define a normal, injective $*$-homomorphism $\hat{\pi}: L^{\infty}(\mathbb{Z}) \rightarrow \hat{M}$ by setting $\hat{\pi}\left(\delta_{k}\right)=\delta_{q^{\frac{k}{2}}}(K)$.

Note that $\hat{\pi}$ preserves the comultiplication, since

$$
\begin{aligned}
(\hat{\pi} \otimes \hat{\pi}) \Delta_{\mathbb{Z}}\left(\delta_{k}\right) & =(\hat{\pi} \otimes \hat{\pi})\left(\sum_{l \in \mathbb{Z}} \delta_{l} \otimes \delta_{k-l}\right) \\
& =\sum_{l \in \mathbb{Z}} \delta_{q^{\frac{l}{2}}}(K) \otimes \delta_{q^{\frac{k-l}{2}}}(K)=\delta_{q^{\frac{k}{2}}}(K \otimes K)=\delta_{q^{\frac{k}{2}}}(\hat{\Delta}(K)) .
\end{aligned}
$$

Therefore, $\hat{\pi}$ identifies $\left(\mathbb{T}, \Delta_{\mathbb{T}}\right)$ as a closed quantum subgroup of $(M, \Delta)$. Furthermore, $\hat{\pi}$ induces a morphism $\pi$ between the universal quantum groups $\left(M_{u}, \Delta_{u}\right)$ and $\left(C(\mathbb{T}), \Delta_{\mathbb{T}}\right)$, where here with slight abuse of notation $\Delta_{\mathbb{T}}$ is restricted to a map $C(\mathbb{T}) \rightarrow C(\mathbb{T} \times \mathbb{T})$.

## Spherical corepresentations

We compute the actions $\gamma$ and $\beta$ of left and right translation and determine which of the corepresentations found in [10] admit a $L^{\infty}(\mathbb{T})$-invariant vector.

Proposition 8.7. For all $p \in q^{\mathbb{Z}}, x \in(\mu(-q), \mu(q)) \cup \sigma_{d}\left(\Omega_{p}\right)$,

$$
\begin{align*}
& \beta\left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\epsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right)=\left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right) \otimes \zeta^{2 m}  \tag{8.5}\\
& \gamma\left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right)=\left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right) \otimes \zeta^{2 m^{\prime}} \tag{8.6}
\end{align*}
$$

Here, $\zeta$ is the identity function on the complex unit circle $\mathbb{T}$.
Proof. Note that $\hat{\pi}$ has a direct integral decomposition $\hat{\pi}=\int{ }^{\oplus} \hat{\pi}_{p, x} d(p, x)$. Here $\hat{\pi}_{p, x}$ : $L^{\infty}(\mathbb{Z}) \rightarrow B\left(\mathcal{L}_{p, x}\right)$ is determined by $\hat{\pi}_{p, x}\left(\delta_{k}\right)=\delta_{q^{\frac{k}{2}}}(K)$, where the action of $K$ on the representation space is given by (8.3) and similarly for the discrete series corepresentations.

By the definition of $\beta$, see (3.1), and the decomposition (8.1), we see that

$$
\int^{\oplus}(\beta \otimes \iota)\left(W_{p, x}\right) d(p, x)=\int^{\oplus}\left(W_{p, x}\right)_{13}\left(\iota \otimes \hat{\pi}_{p, x}\right)\left(W_{\mathbb{T}}\right)_{23} d(p, x) .
$$

By (8.3) and the definition of $\hat{\pi}$, for almost all pairs $(p, x)$ in the decomposition (8.1),

$$
\begin{equation*}
(\beta \otimes \iota)\left(W_{p, x}\right)=\left(W_{p, x}\right)_{13}\left(\iota \otimes \hat{\pi}_{p, x}\right)\left(W_{\mathbb{T}}\right)_{23} . \tag{8.7}
\end{equation*}
$$

This proves (8.5) for the discrete series corepresentations. It follows from (8.4) and the fact that the function $C$ given there is analytic on a neighbourhood of $(\mu(-q), \mu(q))$, see [10, Section 10.3], that for every $p \in q^{\mathbb{Z}}$,

$$
\begin{equation*}
(\mu(-q), \mu(q)) \rightarrow M \otimes L^{\infty}(\mathbb{T}): x \mapsto \beta\left(\left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right)\right), \tag{8.8}
\end{equation*}
$$

extends to an analytic function on a neighbourhood of $(\mu(-q), \mu(q))$. Here, we use the fact that $\beta$ is normal and the fact that a function is $\sigma$-weak analytic if and only if it is analytic with respect to the norm [20, Result 1.2]. Similarly, it follows that for $p \in q^{\mathbb{Z}}$,

$$
\begin{equation*}
(\mu(-q), \mu(q)) \rightarrow M \otimes L^{\infty}(\mathbb{T}): x \mapsto\left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right) \otimes \zeta^{2 m} \tag{8.9}
\end{equation*}
$$

extends to an analytic function on a neighbourhood of $(\mu(-q), \mu(q))$. Note that for all $p \in q^{\mathbb{Z}}$, $x \in(\mu(-q), \mu(q))$,

$$
\begin{aligned}
& \left(\iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(W_{p, x}\right) \otimes \zeta^{2 m} \\
& \quad=\left(\iota \otimes \iota \otimes \omega_{f_{m}^{\varepsilon, \eta}(p, x), f_{m^{\prime}}^{\varepsilon^{\prime}, \eta^{\prime}}(p, x)}\right)\left(\left(W_{p, x}\right)_{13}\left(\iota \otimes \hat{\pi}_{p, x}\right)\left(W_{\mathbb{T}}\right)_{23}\right)
\end{aligned}
$$

Now, (8.7) yields that (8.8) and (8.9) agree on a dense subset of $[-1,1]$. Since (8.8) and (8.9) have analytic extensions, they are equal for every $x \in(\mu(-q), \mu(q))$. The proof of (8.6) is similar.

Remark 8.8. Note that in the preceding proof, we cannot apply [29, Theorem IV.8.25] directly to the map $\hat{\pi}$ on the von Neumann algebraic level since $L^{\infty}(\mathbb{T})$ is not separable. Therefore, we have defined the representations $\hat{\pi}_{p, x}$ explicitly.

Recall that we defined $L^{\infty}(\mathbb{T})$-invariant vectors in Definition 5.1.
Corollary 8.9. Let $p \in q^{\mathbb{Z}}, x \in(\mu(-q), \mu(q)) \cup \sigma_{d}\left(\Omega_{p}\right)$.

1. If $p \in q^{2 \mathbb{Z}}$, then the space of $L^{\infty}(\mathbb{T})$-invariant vectors of $W_{p, x}$ is spanned by $f_{0}^{\varepsilon, \eta}(p, x)$, with $\varepsilon, \eta \in\{+,-\}$.
2. If $p \in q^{2 \mathbb{Z}}, x \in \sigma_{d}\left(\Omega_{p}\right)$. The space of $L^{\infty}(\mathbb{T})$-invariant vectors of $W_{p, x}$ is one dimensional.
3. If $p \in q^{1+2 \mathbb{Z}}$, then $W_{p, x}$ has no $L^{\infty}(\mathbb{T})$-invariant vectors.
4. $P_{\gamma}=\delta_{1}(K)=\hat{\pi}\left(\delta_{0}\right)$ and hence the von Neumann algebras $N$ and $\hat{N}$ that are constructed in Section 3 are given by

$$
N=\overline{\left\{\left(\iota \otimes \omega_{\delta_{1}(K) v, \delta_{1}(K) w}\right)(W) \mid v, w \in \mathcal{K}\right\}^{\sigma-\text { strong-* }}, \quad \hat{N}=\delta_{1}(K) \hat{M} \delta_{1}(K) . . . . ~ . ~}
$$

Proof. From the considerations in Sections 3 and 5, (1), (3) and (4) follow. For (2) consider $\lambda \in-q^{2 \mathbb{Z}+1} \cup q^{2 \mathbb{Z}+1}$ be such that $x=\mu(\lambda)$ and $|\lambda| \geq 1$. Put $j^{\prime}, l^{\prime} \in \mathbb{Z}$ by setting $|\lambda|=q^{1-2 j^{\prime}}=$ $q^{1+2 l^{\prime}}$. In particular, $l^{\prime}<j^{\prime}$. For case (i) of [10, Proposition 5.2], $f_{0}^{-,+}(p, x)$ is zero if and only if $0>l^{\prime}$ if and only if $|\lambda| \geq 1$, which is true by assumption. Similarly, $f_{0}^{+,-}(p, x)=0$. Hence the only $L^{\infty}(\mathbb{T})$-invariant vector is $f_{0}^{+,+}(p, x)$. Cases (ii) and (iii) of [10, Proposition 5.2] follow in exactly the same manner.

Remark 8.10. Note that the space of $L^{\infty}(\mathbb{T})$-invariant vectors for the irreducible components of $W_{p, x}$ is not necessarily of dimension $\leq 1$. For example $W_{p, x}, p \in q^{2 \mathbb{Z}}, x \in[-1,1] \backslash\{0\}$ splits as a sum of 2 irreducible corepresentations of which the $L^{\infty}(\mathbb{T})$-invariant vectors form a 2 dimensional vector space, see [10, Section 10.2]. This implies that $\Delta^{\natural}$ is not cocommutative and we can not use (5.1) directly to obtain product formulae. However, we will define gradings on the spaces $N$ and $\hat{N}$ which still allows us to find such formulae.

In Corollary 8.9 we determined the corepresentations appearing in the decomposition that admit a $L^{\infty}(\mathbb{T})$-invariant vector. These corepresentations are not mutually inequivalent. Here, we give a list of the equivalences. We only consider the spherical corepresentations and consider the principal, discrete as well as the complementary series. Only the principal and discrete series
are important to determine the spaces $N, \hat{N}$ and the spherical Fourier transform, see Theorem 7.1. Nevertheless, the complementary series is still important, since the product formulae we derive later still hold for the spherical matrix elements of the complementary series.

Motivated by [10, Lemma 10.11], we introduce the following basis vectors. For $p \in q^{2 \mathbb{Z}}$, $x \in(\mu(-q), \mu(q))$, set

$$
\begin{aligned}
g_{m}^{1,+}(p, x) & =\frac{1}{2} \sqrt{2}\left(f_{m}^{+,+}(p, x)+i^{\chi(p)} f_{m}^{-,-}(p, x)\right), \\
g_{m}^{1,-}(p, x) & =\frac{1}{2} \sqrt{2}\left(f_{m}^{+,-}(p, x)-i^{\chi(p)} f_{m}^{-,+}(p, x)\right), \\
g_{m}^{2,+}(p, x) & =\frac{1}{2} \sqrt{2}\left(f_{m}^{+,+}(p, x)-i^{\chi(p)} f_{m}^{-,-}(p, x)\right), \\
g_{m}^{2,-}(p, x) & =\frac{1}{2} \sqrt{2}\left(f_{m}^{+,-}(p, x)+i^{\chi(p)} f_{m}^{-,+}(p, x)\right) .
\end{aligned}
$$

For every $j \in\{1,2\}, p \in q^{2 \mathbb{Z}}, x \in(\mu(-q), \mu(q))$, the vectors $g_{m}^{j, \sigma}(p, x), \sigma \in\{+,-\}, m \in \mathbb{Z}$ form an orthonormal basis for $\mathcal{L}_{p, x}^{j}$, the corepresentation space of one of the summands of $W_{p, x}$, see [10, equation (95)]. For $p \in q^{2 \mathbb{Z}}, x \in \sigma_{p}\left(\Omega_{d}\right)$, set

$$
\begin{align*}
g_{m}^{1,+}(p, x) & =f_{m}^{+,+}(p, x)+i^{\chi(p)} f_{m}^{-,-}(p, x), \\
g_{m}^{1,-}(p, x) & =f_{m}^{+,-}(p, x)-i^{\chi(p)} f_{m}^{-,+}(p, x), \\
g_{m}^{2,+}(p, x) & =f_{m}^{+,+}(p, x)-i^{\chi(p)} f_{m}^{-,-}(p, x), \\
g_{m}^{2,-}(p, x) & =f_{m}^{+,-}(p, x)+i^{\chi(p)} f_{m}^{-,+}(p, x) . \tag{8.10}
\end{align*}
$$

Recall that we made the convention that $f_{m}^{\varepsilon, \eta}(p, x)=0$ in case $f_{m}^{\varepsilon, \eta}(p, x)$ is not in the basis given in [10, Proposition 5.2]. Hence, for $x \in \sigma_{d}\left(\Omega_{p}\right)$, we see from [10, Proposition 5.2] that the vectors defined in (8.10) are dependent and any of them is equal to a vector of the form $f_{m}^{\varepsilon, \eta}(p, x)$ for some $\varepsilon, \eta \in\{+,-\}$ modulo a phase factor.

Now, we determine which of the discrete, principal and complementary series corepresentations are equivalent. By considering the action of the Casimir operator as in the proof of $[2$, Proposition B.1], it is clear that any two corepresentations that fall within a different series are inequivalent. We restrict ourselves to the spherical corepresentations.

## Proposition 8.11.

1. Let $x, x^{\prime} \in(\mu(-q), \mu(q)) \backslash\{0\}, p, p^{\prime} \in q^{2 \mathbb{Z}}, j, j^{\prime} \in\{1,2\}$. $W_{p, x}^{j} \simeq W_{p^{\prime}, x^{\prime}}^{j^{\prime}}$ if and only if either $x= \pm x^{\prime}, j=j^{\prime}, p / p^{\prime} \in q^{4 \mathbb{Z}}$ or $x= \pm x^{\prime}, j \neq j^{\prime}, p / p^{\prime} \in q^{2+4 \mathbb{Z}}$.
2. Let $p, p^{\prime} \in q^{2 \mathbb{Z}}, j, j^{\prime}, k, k^{\prime} \in\{1,2\}$. $W_{p, 0}^{j, k} \simeq W_{p^{\prime}, 0}^{j^{\prime}, k^{\prime}}$ if and only if either $j=j^{\prime}, k=k^{\prime}$, $p / p^{\prime} \in q^{4 \mathbb{Z}}$ or $j \neq j^{\prime}, k=k^{\prime}, p / p^{\prime} \in q^{2+4 \mathbb{Z}}$.
3. Let $x, x^{\prime} \in \mu\left(-q^{\mathbb{Z}} \cup q^{\mathbb{Z}}\right), p, p^{\prime} \in q^{2 \mathbb{Z}} . W_{p, x} \simeq W_{p^{\prime}, x^{\prime}}$ if and only if $|x|=\left|x^{\prime}\right|$.

Proof. The proposition follows from a careful comparison of the action of the generators, see (8.3) for the principal series, [10, Proposition 5.2 and Lemma 10.1] for the discrete series and [10, Section 10.3] for the complementary series. We prove (1). By considering the Casimir operator [10, Definition 4.5] one sees that if an irreducible component of $W_{p, x}$ is equivalent to an irreducible component of $W_{p^{\prime}, x^{\prime}}$, then $|x|=\left|x^{\prime}\right|$.

In case $x=x^{\prime}$, an intertwiner must send $g_{m}^{j, \pm}(p, x)$ to a non-zero scalar multiple of $g_{m}^{j^{\prime}, \pm}\left(p^{\prime}, x\right)$ as follows by considering the actions of $K$ and $E$. Writing out the actions of $U_{0}^{+-}$and $U_{0}^{-+}$one sees that there exists such an intertwiner only in the following two cases:
(i) $p / p^{\prime} \in q^{4 \mathbb{Z}}, j=j^{\prime}$, for which it sends $g_{m}^{j, \pm}(p, x)$ to $g_{m}^{j, \pm}\left(p^{\prime}, x\right)$;
(ii) $p / p^{\prime} \in q^{2+4 \mathbb{Z}}, j \neq j^{\prime}$, in which case it sends $g_{m}^{j, \pm}(p, x)$ to $\pm g_{m}^{j^{\prime}, \pm}\left(p^{\prime}, x\right)$.

Similarly, in case $x=-x^{\prime}$, an intertwiner must send $g_{m}^{j, \pm}(p, x)$ to $g_{m}^{j^{\prime}, \mp}\left(p^{\prime},-x\right)$ as follows from the actions of $K$ and $E$. The actions of $U_{0}^{+-}$and $U_{0}^{-+}$show that this is only possible if
(i) $p / p^{\prime} \in q^{4 \mathbb{Z}}, j=j^{\prime}$ for which it sends $g_{m}^{j, \pm}(p, x)$ to $\pm g_{m}^{j, \mp}\left(p^{\prime},-x\right)$;
(ii) $p / p^{\prime} \in q^{2+4 \mathbb{Z}}, j \neq j^{\prime}$, in which case it sends $g_{m}^{j, \pm}(p, x)$ to $g_{m}^{j^{\prime}, \pm}\left(p^{\prime},-x\right)$.

This proves (1), the other cases follows similarly.

Summarizing Corollary 8.9 and Proposition 8.11, we find that $\operatorname{IC}\left(M, M_{1}\right)$, the space of equivalence classes of irreducible spherical corepresentations is partly given by

$$
\begin{array}{rlrl}
\operatorname{IC}\left(M, M_{1}\right) \supseteq & (0,1,1) \cup(0,1,2) \cup(0,2,1) \cup(0,2,2) \cup((0,1] \times\{1,2\}) & & \text { (principal) } \\
\cup \mu\left(q^{2 N+1}\right) & & \text { (discrete) }  \tag{8.11}\\
\cup(1, \mu(q)) \times\{1,2\} & & \text { (complementary) }
\end{array}
$$

Here, we identify the points $(0,1,1),(0,1,2),(0,2,1),(0,2,2)$ with the respective irreducible corepresentations $W_{1,0}^{1,1}, W_{1,0}^{1,2}, W_{1,0}^{2,1}, W_{1,0}^{2,2}$, see [10, Proposition 10.14 (ii)]. We let a point $(x, j) \in(0,1] \times\{1,2\}$ correspond to $W_{1, x}^{j}$, see [10, Proposition 10.12]. The points $x \in \mu\left(q^{2 \mathbb{N}+1}\right)$ corresponds to $W_{1, x}$, see [10, Proposition 5.2]. We emphasize that it is not known if the corepresentations described in [10] are all the corepresentations, therefore we do not know if this description completely describes $\operatorname{IC}\left(M, M_{1}\right)$. For the von Neumann algebras $N$ and $\hat{N}$ as well as the spherical Fourier transform, only the principal and discrete (spherical) series matter. These are fully identified within (8.11). For completeness, we illustrate the right hand side of (8.11) by means of Fig. 1.


III

Figure 1. Known part of $\operatorname{IC}\left(M, M_{1}\right)$ for $S U_{q}(1,1)_{\text {ext }}$ with the diagonal subgroup, c.f. (8.11). Part I: Principal spherical series. Part II: Complementary spherical series. Part III: Discrete spherical series.

## The von Neumann algebra $N$

The next step is to make $N$ and $\hat{N}$ more explicit and meanwhile define a grading on these spaces. Therefore, we will first find an alternative formula for the mapping $T_{\beta} T_{\gamma}$ which is convenient for computations. This is also going to play a role when we derive product formulae. Recall that the spectrum of $K$ is given by $q^{\frac{1}{2} \mathbb{Z}}$. For $k, l \in \mathbb{Z}$, set

$$
M_{k, l}=\overline{\left\{\left.\left(\iota \otimes \omega_{\delta^{\frac{1}{2}(k-l)}}(K) v, \delta_{q^{\frac{1}{2}(k+l)}}(K) w\right)(W) \right\rvert\, v, w \in \mathcal{K}\right\}^{\sigma \text {-strong-* }} .}
$$

By (8.3) and (8.4), the spaces $M_{k, l}$ have mutually trivial intersection. Moreover, $M$ equals the $\sigma$-strong-* closure of the direct sum of vector spaces $\oplus_{k, l \in \mathbb{Z}} M_{k, l}$. Note that $N=M_{0,0}$.

Definition 8.12. For $m \in \mathbb{Z}, p, t \in-q^{\mathbb{Z}} \cup q^{\mathbb{Z}}$, let $P_{m, p, t}$ be the orthogonal rank one projection of $\mathcal{K}$ onto the space spanned by the vector $f_{m, p, t}$. Here we define $f_{m, p, t}$ to be the zero vector if either $p \notin I_{q}$ or $t \notin I_{q}$. Define mappings $T^{+}$and $T^{-}$defined by

$$
T^{ \pm}: M \rightarrow M: x \mapsto \sum_{m \in \mathbb{Z}, p, t \in I_{q}} P_{m, \pm p, t} x P_{m, p, t}
$$

where the sum converges in the strong topology.
Proposition 8.13. $T^{+}+T^{-}=T_{\beta} T_{\gamma}$.
Proof. For a normal functional $\omega=\sum_{i \in I} \omega_{\xi_{i}, \eta_{i}} \in M_{*}$, with $\sum_{i \in I}\left\|\xi_{i}\right\|^{2}, \sum_{i \in I}\left\|\eta_{i}\right\|^{2}<\infty$, the linear map

$$
M \rightarrow \mathbb{C}: x \mapsto \omega T^{ \pm}(x)=\sum_{i \in I} \sum_{m \in \mathbb{Z}, p, t \in I_{q}}\left\langle x P_{m, p, t} \xi_{i}, P_{m, \pm p, t} \eta_{i}\right\rangle
$$

is normal. Hence, the maps $T^{ \pm}$are normal. Moreover, for $x \in M_{k, l}$, using (8.4),

$$
T^{+}(x)+T^{-}(x)= \begin{cases}x, & \text { if } k=l=0  \tag{8.12}\\ 0, & \text { otherwise }\end{cases}
$$

For $x \in M_{k, l}$ we see by using (8.5) and (8.6) that $T_{\beta} T_{\gamma}(x)$ also equals the right hand side of (8.12). Since also $T_{\beta} T_{\gamma}$ is normal, this proves that $T^{+}+T^{-}=T_{\beta} T_{\gamma}$.

Let $u_{0} \in B\left(L^{2}\left(I_{q}\right)\right)$ be the partial isometry determined by $u_{0}: \delta_{p} \mapsto \delta_{-p}$, and recall that $u$ was defined as $u=1 \otimes u_{0} \in B\left(L^{2}(\mathbb{Z}) \otimes L^{2}\left(I_{q}\right)\right)$. By [13, Lemma 2.4] $M \simeq L^{\infty}(\mathbb{T}) \otimes B\left(L^{2}\left(I_{q}\right)\right)$. Proposition 8.13 yields,

$$
N=T_{\beta} T_{\gamma}(M)=\left(T^{+}+T^{-}\right)(M)
$$

which is isomorphic to the von Neumann subalgebra of $M$ generated by $1 \otimes L^{\infty}\left(I_{q}\right)$ and $u$. Therefore, introduce the following identifications.

Definition 8.14. We identify $N$ with the von Neumann algebra acting on $L^{2}\left(I_{q}\right)$ being generated by $L^{\infty}\left(I_{q}\right)$ and $u_{0}$. We split $N$ as a direct sum of vector spaces

$$
N=N_{+} \oplus N_{-}, \quad \text { where } \quad N_{+}=L^{\infty}\left(I_{q}\right), \quad \text { and } \quad N_{-}=L^{\infty}\left(I_{q}\right) u_{0}=L^{\infty}\left(I_{q} \cap(-1,1)\right) u_{0}
$$

This turns $N$ into a $\mathbb{Z}_{2}$-graded algebra.
We find that $\varphi^{\natural}$, i.e. the restriction of the Haar weight $\varphi[13$, Definition 4.1$]$ to $N$, equals the measure given by:

$$
\varphi^{\natural}(f)=\sum_{p_{0} \in I_{q}} f\left(p_{0}\right) p_{0}^{2}, \quad f \in N^{+}=L^{\infty}\left(I_{q}\right)^{+}
$$

For $f \in \mathfrak{m}_{\varphi^{\natural}}^{+}$we find that $u_{0} f \in \mathfrak{m}_{\varphi^{\natural}}$ and it follows by [13, Definition 4.1] that $\varphi^{\natural}\left(u_{0} f\right)=0$, so that $\Lambda\left(\mathfrak{n}_{\varphi} \cap N_{+}\right)$and $\Lambda\left(\mathfrak{n}_{\varphi} \cap N_{-}\right)$are orthogonal spaces. The discussion so far allows us to make the following identifications.

Definition 8.15. Identify the closure of $\Lambda\left(\mathfrak{n}_{\varphi} \cap N_{+}\right)$with $L^{2}\left(I_{q}\right)$. We will write $L^{2}\left(N_{+}\right)$for $L^{2}\left(I_{q}\right)$ to indicate explicitly that $L^{2}\left(I_{q}\right)$ is considered as part of the GNS-space of $\varphi^{\natural}$. Similarly, we identify the closure of $\Lambda\left(\mathfrak{n}_{\varphi} \cap N_{-}\right)$with $L^{2}\left(I_{q} \cap(-1,1)\right)$, by identifying $f u_{0} \in N_{-}$, where $f \in L^{\infty}\left(I_{q} \cap(-1,1)\right) \cap L^{2}\left(I_{q} \cap(-1,1)\right)$ with $f \in L^{2}\left(I_{q} \cap(-1,1)\right)$. We write short hand $L^{2}\left(N_{-}\right)$for $L^{2}\left(I_{q} \cap(-1,1)\right)$. We write $L^{2}(N)=L^{2}\left(N_{+}\right) \oplus L^{2}\left(N_{-}\right)$. We emphasize that here the spaces $L^{2}\left(I_{q}\right)$ (respectively $L^{2}\left(I_{q} \cap(-1,1)\right)$ ) should be understood with respect to the integral given by the weighted $\operatorname{sum} \sum_{p_{0} \in I_{q}}(\cdot) p_{0}^{2}\left(\right.$ respectively $\left.\sum_{p_{0} \in I_{q} \cap(-1,1)}(\cdot) p_{0}^{2}\right)$.

## The von Neumann algebra $\hat{N}$

We now turn our attention to the von Neumann algebra $\hat{N}$ as defined in Section 3. Considered as a subalgebra of $\hat{M}$, it inherits the $\mathbb{Z}_{2}$-grading defined in [10, Definition 4.7].

Definition 8.16. We set

$$
\hat{N}_{+}=\hat{N} \cap \hat{M}_{+}, \quad \hat{N}_{-}=\hat{N} \cap \hat{M}_{-},
$$

see [10, Definition 4.7] for $\hat{M}_{+}$and $\hat{M}_{-}$. This turns $\hat{N}$ into a $\mathbb{Z}_{2}$-graded algebra.
Proposition 8.17. We have an isomorphism of von Neumann algebras:

$$
\begin{equation*}
\hat{N} \simeq \int_{[0,1]}^{\oplus} M_{2}(\mathbb{C}) \oplus M_{2}(\mathbb{C}) d x \oplus \bigoplus_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)} \mathbb{C} \tag{8.13}
\end{equation*}
$$

The isomorphism is determined by the map

$$
\begin{equation*}
(\tilde{\omega} \otimes \iota)(W) \mapsto \int_{[0,1]}^{\oplus}(\tilde{\omega} \otimes \iota)\left(W_{1, x}^{1}\right) \oplus(\tilde{\omega} \otimes \iota)\left(W_{1, x}^{2}\right) d x \oplus \bigoplus_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)}(\tilde{\omega} \otimes \iota)\left(W_{1, x}\right) . \tag{8.14}
\end{equation*}
$$

Under this isomorphism $\hat{N}_{+}$corresponds to the direct integrals over matrices whose entries vanish off the diagonal. $\hat{N}_{-}$corresponds to the direct integrals over matrices with values vanishing on the diagonal.
Proof. It follows from Corollary 8.9 that the map defined in (8.14) indeed maps into the proper matrix algebras. Let $P \in \hat{M} \cap \hat{M}^{\prime}$ be the projection as in the proof of [2, Proposition B.2]. By considering the Casimir operator, one sees that $P$ is the spectral projection on the interval $[-1,1]$ of the Casimir operator. In [2, Proposition B.2], it is shown that

$$
P \hat{M} P \simeq \int_{x \in[0,1]}^{\oplus} \hat{M}_{x} d x
$$

where $\hat{M}_{x}$ is generated by $\left\{(\omega \otimes \iota)\left(W_{x}\right) \mid \omega \in M_{*}\right\}$ and $W_{x}=\left(\oplus_{p \in q^{z}} W_{p, x}\right) \oplus\left(\oplus_{p \in q^{z}} W_{p,-x}\right)$. By similar techniques as in [2, Proposition B.2], one can show that

$$
\begin{equation*}
\hat{M} \simeq \int_{x \in[0,1]}^{\oplus} \hat{M}_{x} d x \oplus \bigoplus_{x \in \sigma_{d}(\Omega) \cap(1, \infty)} \hat{M}_{x} \tag{8.15}
\end{equation*}
$$

where $\hat{M}_{x}, x \in \sigma_{d}(\Omega) \cap(1, \infty)$ is generated by $\left\{(\omega \otimes \iota)\left(W_{x}\right) \mid \omega \in M_{*}\right\}$ and

$$
W_{x}=\left(\bigoplus_{p \in q^{Z}, \text { s.t. } x \in \sigma_{d}\left(\Omega_{p}\right)} W_{p, x}\right) \oplus\left(\bigoplus_{p \in q^{\mathbb{Z}}, \text { s.t. }-x \in \sigma_{d}\left(\Omega_{p}\right)} W_{p,-x)}\right) .
$$

Since $P_{\gamma} \in \hat{M}$, it has a direct integral decomposition as in (8.15), i.e.

$$
P_{\gamma} \hat{M} P_{\gamma} \simeq \int_{x \in[0,1]}^{\oplus}\left(P_{\gamma}\right)_{x} \hat{M}_{x}\left(P_{\gamma}\right)_{x} d x \oplus \bigoplus_{x \in \sigma_{d}(\Omega) \cap(1, \infty)}\left(P_{\gamma}\right)_{x} \hat{M}_{x}\left(P_{\gamma}\right)_{x}
$$

Similarly, $K$ decomposes with respect to the direct integral decomposition (8.15) as a direct integral over a field $\left(K_{x}\right)_{x \in[0,1] \cup\left(\sigma_{d}(\Omega) \cap(1, \infty)\right)}$ and by Proposition 8.7, we see that $\left(T_{\beta} T_{\gamma} \otimes \iota\right)\left(W_{x}\right)=$ $\left(1 \otimes \delta_{1}\left(K_{x}\right)\right) W_{x}\left(1 \otimes \delta_{1}\left(K_{x}\right)\right)=\left(1 \otimes\left(P_{\gamma}\right)_{x}\right) W_{x}\left(1 \otimes\left(P_{\gamma}\right)_{x}\right)$. Hence, $\left(P_{\gamma}\right)_{x} \hat{M}_{x}\left(P_{\gamma}\right)_{x}$ is generated by $\left\{(\tilde{\omega} \otimes \iota)\left(W_{x}\right) \mid \omega \in M_{*}\right\}$. Due to Corollary 8.9 and Proposition 8.11, the latter von Neumann algebra is isomorphic to $M_{2}(\mathbb{C}) \oplus M_{2}(\mathbb{C})$ in case $x \in(0,1]$ and to $\mathbb{C}$ in case $x \in \sigma_{d}(\Omega) \cap(1, \infty)$. This proves (8.13). That (8.14) gives the isomorphism follows directly from this proof. The claim on the gradings follows from [10, Proposition 4.8].

Remark 8.18. The von Neumann algebra $\hat{M}_{0}$ in the previous proof is isomorphic to 4 copies of $\mathbb{C}$ as follows from [10, Proposition 10.13]. Since this does not matter for the integral decomposition (8.13), and to avoid some redundant extra notation we have not treated the point $x=0$ separately.

We claim that $\hat{\varphi}^{\natural}$, i.e. the restriction of the dual left Haar weight to $\hat{N}$, is a trace. Indeed, it follows from [2, Section 5] and [4, Proposition 3.5.5] that for any $p \in q^{\mathbb{Z}}, x \in[-1,1] \cup \sigma_{d}\left(\Omega_{p}\right)$, there is a constant $c(p, x)$ such that

$$
D_{p, x} f_{m}^{\epsilon, \eta}(p, x)=p^{1 / 2} q^{m} c(p, x) f_{m}^{\varepsilon, \eta}(p, x)
$$

Here $D=\int^{\oplus} D_{p, x} d(p, x)$, with short hand notation $D_{p, x}=D_{W_{p, x}}$, is the Duflo-Moore operator arising from the Plancherel theorem, see [4, Theorem 3.4.1] and also Theorem 7.1. Since the eigenvalues of $D_{p, x}$ are independent of the signs $\varepsilon, \eta$, we see that $\int^{\oplus} D_{p, x}^{L^{\infty}(\mathbb{T})} d(p, x)=$ $P_{\gamma} \int^{\oplus} D_{p, x} d(p, x) P_{\gamma}$ is central in $\hat{N}$. Since $\int^{\oplus}\left(D_{p, x}^{L^{\infty}(\mathbb{T})}\right)^{-2} d(p, x)$ is the Radon-Nikodym derivative of $\hat{\varphi}^{\natural}$ with respect to a trace, see the proof of [4, Theorem 3.4.5], we see that $\hat{\varphi}^{\natural}$ is a trace. Therefore, under the identification (8.13), we see that

$$
\begin{equation*}
\hat{\varphi}^{\natural}=\int_{[0,1]}^{\oplus}\left(\operatorname{Tr}_{M_{2}(\mathbb{C})} \oplus \operatorname{Tr}_{M_{2}(\mathbb{C})}\right) d(x) d x \oplus \bigoplus_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)} d(x) \operatorname{Tr}_{\mathbb{C}} \tag{8.16}
\end{equation*}
$$

where $d$ is a scalar valued function, which we leave undetermined.
Remark 8.19. It follows from [4, Theorem 3.4.1 (6)] that the function $c(1, x)$ depends on the Plancherel measure. In case the Plancherel measure is choosen as in (8.1), one has $d(x)=$ $c(1, x)^{-2}$. For the discrete series these constants are computed in [4, Section 3.5]. For the principal series, the exact values cannot be found in the literature. We will not derive them here, since this is beyond the scope of our example.

We identify the GNS-space of $\varphi^{\natural}$ using the isomorphism (8.13). That is, the GNS-space is given by

$$
L^{2}(\hat{N})=\int_{[0,1]}^{\oplus} M_{2}(\mathbb{C}) \oplus M_{2}(\mathbb{C}) d x \oplus \bigoplus_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)} \mathbb{C}
$$

where the direct integral and direct sums are taken as Hilbert spaces (as opposed to the direct integrals and sums of von Neumann algebras given in (8.13)), where the inner product comes from the traces $d(x)\left(\operatorname{Tr}_{M_{2}(\mathbb{C})} \oplus \operatorname{Tr}_{M_{2}(\mathbb{C})}\right)$ for the integral part and $d(x) \operatorname{Tr}_{\mathbb{C}}$ for the direct summands.

It follows from (8.13) and (8.16) that $\hat{\Lambda}\left(\mathfrak{n}_{\hat{\varphi}} \cap \hat{N}_{+}\right)$and $\hat{\Lambda}\left(\mathfrak{n}_{\hat{\varphi}} \cap \hat{N}_{-}\right)$are orthogonal spaces, we denote their closures interpreted within $L^{2}(\hat{N})$ by $L^{2}\left(\hat{N}_{+}\right)$and $L^{2}\left(\hat{N}_{-}\right)$. They consist of direct integrals of diagonal matrices and off-diagonal matrices, respectively.

## The spherical Fourier transform

Here, we determine the spherical Fourier transform, i.e. we determine the map $\mathcal{Q}_{0}^{L^{\infty}(\mathbb{T})}$ for Theorem 7.1. In particular, we are interested in the kernels of the integral transformations that appear in this transform, since we will need them later. Using the identifications of the GNSspaces for $\varphi^{\natural}$ and $\hat{\varphi}^{\natural}$ with $L^{2}(N)$ and $L^{2}(\hat{N})$, we may consider $\mathcal{Q}_{0}^{L^{\infty}(\mathbb{T})}$ as a map from $L^{2}(N)$ to $L^{2}(\hat{N})$ and use the short hand notation $\mathcal{F}_{2}: L^{2}(N) \rightarrow L^{2}(\hat{N})$ for this map.

If $f \in L^{1}\left(I_{q}\right) \cap L^{\infty}\left(I_{q}\right) \subseteq N_{+}$, then there is a functional $f \cdot \varphi^{\natural} \in N_{*}$ given by $\left(f \cdot \varphi^{\natural}\right)(x)=$ $\varphi^{\natural}(x f), x \in N$. Then, $\xi\left(f \cdot \varphi^{\natural}\right)=\Lambda(f)$ by definition of $\xi$, see Section 2 . Thus, for such $f$, we find by Theorem 7.1 that,

$$
\begin{equation*}
\mathcal{F}_{2}: \Lambda(f) \mapsto\left(\left(f \cdot \varphi^{\natural}\right)^{\sim} \otimes \iota\right)\left(\int_{[0,1]}^{\oplus} W_{1, x} d x \oplus \bigoplus_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)} W_{1, x}\right) . \tag{8.17}
\end{equation*}
$$

Here, the right hand side indeed is an element of $L^{2}(\hat{N})$ by Corollary 8.9. Note that we choose the inner product on $L^{2}(\hat{N})$ to be the direct integral over the traces $d(x)\left(\operatorname{Tr}_{M_{2}(\mathbb{C})} \oplus \operatorname{Tr}_{M_{2}(\mathbb{C})}\right)$ for the integral part and $d(x) \operatorname{Tr}_{\mathbb{C}}$ for the direct summands. Hence, the Duflo-Moore operators which are direct integrals of scalar multiples of the identity are contained in the inner product by means of the function $d$. Hence, they do not appear in (8.17). Similarly, for $f \in L^{1}\left(I_{q}\right) \cap L^{\infty}\left(I_{q}\right)$, so that $f u_{0} \in N_{-}$,

$$
\begin{equation*}
\mathcal{F}_{2}: \Lambda\left(f u_{0}\right) \mapsto\left(\left(f u_{0} \cdot \varphi^{\natural}\right)^{\sim} \otimes \iota\right)\left(\int_{[0,1]}^{\oplus} W_{1, x} d x \oplus \bigoplus_{x \in \mu\left(q^{2 N+1}\right)} W_{1, x}\right) . \tag{8.18}
\end{equation*}
$$

We see from (8.17) that the spherical Fourier transformation is in fact a combination of integral transformations with the spherical matrix elements of the corepresentations $W_{1, x}$ in its kernel. The next step is to make these kernels explicit.

For any $m_{0} \in \mathbb{Z}, t_{0} \in I_{q}, x \in[-1,1] \cup \mu\left(-q^{2 \mathbb{N}+1} \cup q^{2 \mathbb{N}+1}\right), j \in\{1,2\}, \sigma, \tau \in\{+,-\}$, put

$$
K_{j}^{\sigma, \tau}\left(p_{0} ; x\right)=\left\langle\left(\iota \otimes \omega_{g_{0}^{j, \sigma}(1, x), g_{0}^{j, \tau}(1, x)}\right)\left(W_{1, x}\right) f_{m_{0}, p_{0}, t_{0}}, f_{m_{0}, \sigma \tau p_{0}, t_{0}}\right\rangle .
$$

We emphasize that $K_{j}^{\sigma, \tau}\left(p_{0} ; x\right)$ for $x$ not in $\sigma_{d}\left(\Omega_{1}\right)$ is defined by Remark 8.5. This expression is independent of $m_{0}$ and $t_{0}$. Due to Proposition 8.11, we have the following symmetry

$$
\begin{equation*}
K_{j}^{\sigma, \tau}\left(p_{0} ; x\right)=\sigma \tau K_{j}^{-\sigma,-\tau}\left(p_{0} ;-x\right) . \tag{8.19}
\end{equation*}
$$

From (8.4), (8.17) and (8.18), we see that we get a graded Fourier transform $\mathcal{F}_{2,+} \oplus \mathcal{F}_{2,-}$ : $L^{2}\left(N_{+}\right) \oplus L^{2}\left(N_{-}\right) \rightarrow L^{2}\left(\hat{N}_{+}\right) \oplus L^{2}\left(\hat{N}_{-}\right)$, which is defined by sending $f \oplus g u_{0} \in\left(L^{2}\left(I_{q}\right) \cap\right.$ $\left.L^{1}\left(I_{q}\right)\right) \oplus\left(L^{2}\left(I_{q} \cap(-1,1)\right) \cap L^{1}\left(I_{q} \cap(-1,1)\right)\right) u_{0} \subseteq L^{2}\left(N_{+}\right) \oplus L^{2}\left(N_{-}\right)$to the matrix valued function on $\operatorname{IC}\left(M, L^{\infty}(\mathbb{T})\right)$ determined by sending $(x, j) \in[0,1] \times\{0,1\}$ to

$$
\sum_{p_{0} \in I_{q}}\left(\begin{array}{cc}
K_{j}^{+,+}\left(p_{0} ; x\right) f\left(p_{0}\right) & K_{j}^{-,+}\left(p_{0} ; x\right) g\left(p_{0}\right)  \tag{8.20}\\
K_{j}^{+,-}\left(p_{0} ; x\right) g\left(p_{0}\right) & K_{j}^{-,-}\left(p_{0} ; x\right) f\left(p_{0}\right)
\end{array}\right) p_{0}^{2},
$$

and $x \in \mu\left(q^{2 \mathbb{Z}+1}\right)$ to

$$
\sum_{p_{0} \in I_{q}} K_{0}^{+,+}\left(p_{0} ; x\right) f\left(p_{0}\right) p_{0}^{2}=\sum_{p_{0} \in I_{q}} K_{1}^{+,+}\left(p_{0} ; x\right) f\left(p_{0}\right) p_{0}^{2} .
$$

Note that for $x=0$, the matrix appearing in (8.20) is actually a direct sum of two matrix blocks after a basis transformation, since $W_{1,0}$ splits as a direct sum of four irreducible corepresentations, see [10, Proposition 10.13]. By Theorem 7.1, this map is unitary.

For completeness, we give the analogous result for the inverse Fourier transform $\mathcal{F}_{2}^{-1}$ : $L^{2}\left(\hat{N}_{+}\right) \oplus L^{2}\left(\hat{N}_{-}\right) \rightarrow L^{2}\left(N_{+}\right) \oplus L^{2}\left(N_{-}\right)$. Since $\mathcal{Q}_{0}^{L^{\infty}(\mathbb{T})}$ is a restriction of the Plancherel transformation $\mathcal{Q}$, see Theorem 7.1, we see that $\mathcal{F}_{2}^{-1}$ can be considered as the restriction of $\mathcal{Q}^{-1}$. The transform $\mathcal{Q}^{-1}$ is described in [1] on the operator algebraic level, see in particular [1, Example 3.13]. See also [41] for the algebraic counterpart. More explicitly, the spherical inverse Fourier transform is determined by

$$
\mathcal{F}_{2}^{-1}: f \mapsto(\iota \otimes(f \cdot \hat{\varphi}))\left(W^{*}\right),
$$

where $f \in \hat{N} \cap \mathfrak{n}_{\hat{\varphi}}$ is such that there is a normal functional on $\hat{M}$, denoted by $(f \cdot \hat{\varphi})$, which is determined by $(f \cdot \hat{\varphi})(x)=\hat{\varphi}(x f), x \in \mathfrak{n}_{\hat{\varphi}}^{*}$. By the decomposition of $W$ (8.1), we find the following theorem.

Theorem 8.20. For $\sigma, \tau \in\{+,-\}$, let $f_{1}^{\sigma, \tau}, f_{2}^{\sigma, \tau} \in L^{1}([0,1]) \cap L^{2}([0,1])$, where $L^{1}([0,1])$ and $L^{2}([0,1])$ should be understood with respect to the integral $\int_{[0,1]} d(x) d x$. Let $f_{d} \in L^{1}\left(\mu\left(q^{2 \mathbb{N}+1}\right)\right) \cap$ $L^{2}\left(\mu\left(q^{2 \mathbb{N}+1}\right)\right)$, where $L^{1}\left(\mu\left(q^{2 \mathbb{N}+1}\right)\right.$ ) and $L^{2}\left(\mu\left(q^{2 \mathbb{N}+1}\right)\right)$ should be understood with respect to the integral $\sum_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)} d(x)$. Define the function $f \in L^{2}(\hat{N})$ by

$$
f(x)= \begin{cases}\left(\begin{array}{ll}
f_{1}^{+,+}(x) & f_{1}^{-,+}(x) \\
f_{1}^{+,-}(x) & f_{1}^{-,-}(x)
\end{array}\right) \oplus\left(\begin{array}{ll}
f_{2}^{+,+}(x) & f_{2}^{-,+}(x) \\
f_{2}^{+,-}(x) & f_{2}^{-,-}(x)
\end{array}\right), & x \in[0,1] \\
f_{d}(x), & \end{cases}
$$

Then,

$$
\begin{align*}
\int_{[0,1]} & \left(f_{1}^{+,+}(x) \overline{K_{1}^{+,+}}\left(p_{0} ; x\right)\right. \\
& \left.+f_{1}^{-,-}(x) \overline{K_{1}^{-,-}\left(p_{0} ; x\right)}\right) d(x) d x \\
& +\sum_{x \in \mu\left(q^{2 \mathbb{N}+1}\right)}\left(f_{2}^{+,+}(x) \overline{K_{2}^{+,+}\left(p_{0} ; x\right)}+f_{2}^{-,-}(x) \overline{K_{2}^{-,-}\left(p_{0} ; x\right)}\right) d(x) d x \\
& \bigoplus\left(\int_{[0,1]}\left(f_{1}^{-,+,+}(x) \overline{K_{1}^{+,-}\left(p_{0} ; x\right)} d(x)\right)\right. \\
& \left.+\int_{[0,1]}\left(f_{2}^{-,+}(x) \overline{K_{2}^{+,-}\left(p_{0} ; x\right)}+f_{2}^{+,-}(x) \overline{K_{2}^{-,+}\left(p_{0} ; x\right)}\right) d(x) d x\right) u_{0} \tag{8.21}
\end{align*}
$$

exists for every $p_{0} \in I_{q}$. Moreover, (8.21) considered as a direct sum of functions in $p_{0}$ forms an element of $L^{2}\left(N_{+}\right) \oplus L^{2}\left(N_{-}\right)$. This mapping extends to a unitary map $L^{2}\left(\hat{N}_{+}\right) \oplus L^{2}\left(\hat{N}_{-}\right) \rightarrow$ $L^{2}\left(N_{+}\right) \oplus L^{2}\left(N_{-}\right)$, which is inverse to $\mathcal{F}_{2}$.

We explicitly state the formulae for the kernels $K_{j}^{\sigma, \tau}\left(p_{0} ; x\right)$ for $x \in[-1,1] \cup \mu\left(-q^{2 \mathbb{N}+1} \cup q^{2 \mathbb{N}+1}\right)$, which can be expressed in terms of little $q$-Jacobi functions. Using the notation of [10, Section 9 ] for $S(\cdot)$ and $A(\cdot)$, we find

$$
\begin{aligned}
& K_{1}^{+,+}\left(p_{0} ; x\right)=S\left(-\lambda, p_{0}, p_{0}, 0\right) \times \begin{cases}1, & p_{0}<0 \\
\frac{A(\lambda, 1,0,+,+)}{A(\lambda, 1,0,-,-)}, & p_{0}>0\end{cases} \\
& K_{2}^{+,+}\left(p_{0} ; x\right)=S\left(-\lambda, p_{0}, p_{0}, 0\right) \times \begin{cases}1, & p_{0}<0 \\
-\frac{A(\lambda, 1,0,+,+)}{A(\lambda, 1,0,-,-)}, & p_{0}>0\end{cases} \\
& K_{1}^{+,-}\left(p_{0} ; x\right)=-S\left(\lambda,-p_{0}, p_{0}, 0\right) \times \begin{cases}\frac{A(\lambda, 1,0,+,+)}{A(-\lambda, 1,0,+,-)}, & p_{0}<0 \\
\frac{A(\lambda, 1,0,+,+)}{A(-\lambda, 1,0,-,+)}, & p_{0}>0\end{cases} \\
& K_{2}^{+,--}\left(p_{0} ; x\right)=-S\left(\lambda,-p_{0}, p_{0}, 0\right) \times \begin{cases}\frac{A(\lambda, 1,0,+,+)}{A(-\lambda, 1,0,+,-)}, & p_{0}<0 \\
\frac{A(\lambda, 1,0,+,+)}{A(-\lambda, 1,0,-,+)}, & p_{0}>0\end{cases}
\end{aligned}
$$

The fractions of the functions $A(\cdot)$ are phase factors. Here, $\lambda \in \mathbb{T}$ such that $\mu(\lambda)=x$. And, for $\lambda \in \mathbb{C} \backslash\{0\}$, simplifying [10, Lemma 9.1 (ii)],

$$
\begin{aligned}
& S\left( \pm \lambda, \mp p_{0}, p_{0}, 0\right)=\left|p_{0}\right|^{2} \nu\left(p_{0}\right)^{2} c_{q}^{2} \sqrt{\left( \pm \kappa\left(p_{0}\right),-\kappa\left(p_{0}\right) ; q^{2}\right)_{\infty}}\left(\mp q^{2} ; q^{2}\right)_{\infty} \\
& \quad \times \frac{\left(q^{2},-q^{2} / \kappa\left(p_{0}\right), \lambda q^{2}, 1 / \lambda,-q / \lambda ; q^{2}\right)_{\infty}}{\left(\operatorname{sgn}\left(p_{0}\right) 1 / \lambda, \operatorname{sgn}\left(p_{0}\right) \lambda q^{2}, \pm q / \lambda ; q^{2}\right)_{\infty}} 2 \varphi_{1}\left(\begin{array}{c}
-q / \lambda,-\lambda q \\
-q^{2}
\end{array} ; q^{2},-q^{2} / \kappa\left(p_{0}\right)\right) .
\end{aligned}
$$

The other kernels occurring in (8.20) can be expressed explicitly by means of the symmetry relation (8.19). The matrix coefficients $K_{j}^{\sigma, \tau}\left(p_{0}, x\right)$ are special types of little $q$-Jacobi functions, see also [10, Appendix B.5] and references given there.

## Product formulae for little $\boldsymbol{q}$-Jacobi functions

Let $N_{*,+}$ be the space of normal functionals in $N_{*}$ which are zero on $N_{-}$. Let $N_{*,-}$ be the space of normal functionals in $N_{*}$ which are zero on $N_{+}$. Note that $\varphi^{\natural}$ is a trace on $N$. Therefore, every functional in $N_{*,+}$ is given by $f \cdot \varphi^{\natural}$, where $f$ is a function on $I_{q}$ so that $\sum_{p_{0} \in I_{q}} f\left(p_{0}\right) p_{0}^{2}<\infty$. Every functional in $N_{*,-}$ is given by $f u_{0} \cdot \varphi^{\natural}$, where $f$ is a function on $I_{q} \cap(-1,1)$ so that $\sum_{p_{0} \in I_{q} \cap(-1,1)} f\left(p_{0}\right) p_{0}^{2}<\infty$. So $N_{*}=N_{*,+} \oplus N_{*,-}$ as vector spaces.

For $p \in I_{q}$, recall that $\delta_{p}$ denotes the function on $I_{q}$ whose value is 1 in $p$ and 0 elsewhere. We write $\delta_{p,+}$ for $\delta_{p} \cdot \varphi^{\natural} \in N_{*,+}$. For $p \in I_{q}$, we write $\delta_{p,-}$ for the functional $\delta_{p} u_{0} \cdot \varphi^{\natural} \in N_{*,-}$. Note that only for $p \in I_{q} \cap(-1,1)$, this functional is non-zero, but it is convenient to keep this notation.

Remark 8.21. Note that if one identifies $N$ with the subalgebra $\left(1 \otimes L^{\infty}\left(I_{q}\right) \otimes 1\right) \cup(1 \otimes$ $L^{\infty}\left(I_{q}\right) u_{0} \otimes 1$ ) of $M$ acting on the GNS-space $\mathcal{K}$, then $\delta_{p,+}=\omega_{f_{m_{0}, p, t_{0}}, f_{m_{0}, p, t_{0}}}$ and $\delta_{p,-}=$ $\omega_{f_{m_{0},-p, t_{0}}, f_{m_{0}, p, t_{0}}}$, where $m_{0} \in \mathbb{Z}, t_{0} \in I_{q}$. This functional is independent of $m_{0}$ and $t_{0}$ if considered as a functional on $N$.

Using the fact that the Fourier transform preserves the $\mathbb{Z}_{2}$-gradings on $N$ and $\hat{N}$, we see that for $p_{1}, p_{2} \in I_{q}$ :

$$
\begin{array}{ll}
\left(\delta_{p_{1},+} \otimes \delta_{p_{2},+}\right) \Delta^{\natural} \in N_{*,+}, & \left(\delta_{p_{1},-} \otimes \delta_{p_{2},-}\right) \Delta^{\natural} \in N_{*,+}, \\
\left(\delta_{p_{1},+} \otimes \delta_{p_{2},-}\right) \Delta^{\natural} \in N_{*,-}, & \left(\delta_{p_{1},-} \otimes \delta_{p_{2},+}\right) \Delta^{\natural} \in N_{*,-} .
\end{array}
$$

Hence, we see that for $f, g \in L^{\infty}\left(I_{q}\right)$, there exist constants $A_{p_{0}}\left(p_{1}, p_{2}\right), B_{p_{0}}\left(p_{1}, p_{2}\right), p_{0}, p_{1}, p_{2} \in I_{q}$ and $C_{p_{0}}\left(p_{1}, p_{2}\right), D_{p_{0}}\left(p_{1}, p_{2}\right), p_{0} \in I_{q} \cap(-1,1), p_{1}, p_{2} \in I_{q}$ such that for any $p_{1}, p_{2} \in I_{q}$ the four sums

$$
\begin{array}{ll}
\sum_{p_{0} \in I_{q}}\left|A_{p_{0}}\left(p_{1}, p_{2}\right)\right| p_{0}^{2}, & \sum_{p_{0} \in I_{q}}\left|B_{p_{0}}\left(p_{1}, p_{2}\right)\right| p_{0}^{2}, \\
\sum_{p_{0} \in I_{q} \cap(-1,1)}\left|C_{p_{0}}\left(p_{1}, p_{2}\right)\right| p_{0}^{2}, & \sum_{p_{0} \in I_{q} \cap(-1,1)}\left|D_{p_{0}}\left(p_{1}, p_{2}\right)\right| p_{0}^{2},
\end{array}
$$

are finite and

$$
\begin{align*}
& \left(\delta_{p_{1},+} \otimes \delta_{p_{2},+}\right) \Delta^{\natural}(f)=\sum_{p_{0} \in I_{q}} A_{p_{0}}\left(p_{1}, p_{2}\right) f\left(p_{0}\right) p_{0}^{2},  \tag{8.22}\\
& \left(\delta_{p_{1},-} \otimes \delta_{p_{2},-}\right) \Delta^{\natural}(f)=\sum_{p_{0} \in I_{q}} B_{p_{0}}\left(p_{1}, p_{2}\right) f\left(p_{0}\right) p_{0}^{2},  \tag{8.23}\\
& \left(\delta_{p_{1},+} \otimes \delta_{p_{2},-}\right) \Delta^{\natural}\left(f u_{0}\right)=\sum_{p_{0} \in I_{q} \cap(-1,1)} C_{p_{0}}\left(p_{1}, p_{2}\right) f\left(p_{0}\right) p_{0}^{2}, \tag{8.24}
\end{align*}
$$

$$
\begin{equation*}
\left(\delta_{p_{1},-} \otimes \delta_{p_{2},+}\right) \Delta^{\natural}\left(f u_{0}\right)=\sum_{p_{0} \in I_{q} \cap(-1,1)} D_{p_{0}}\left(p_{1}, p_{2}\right) f\left(p_{0}\right) p_{0}^{2} . \tag{8.25}
\end{equation*}
$$

Putting $f=K_{j}^{\sigma, \tau}\left(p_{0} ; x\right)$ for any $x \in[-1,0) \cup(0,1] \cup \mu\left(-q^{2 \mathbb{N}+1} \cup q^{2 \mathbb{N}+1}\right), j \in\{1,2\}$, this yields a product formula for ${ }_{2} \varphi_{1}$-series. For $p_{1}, p_{2} \in I_{q}$ and $p_{3}, p_{4} \in I_{q} \cap(-1,1)$,

$$
\begin{align*}
& K_{j}^{+,+}\left(p_{1} ; x\right) K_{j}^{+,+}\left(p_{2} ; x\right)=\sum_{p_{0} \in I_{q}} A_{p_{0}}\left(p_{1}, p_{2}\right) K_{j}^{+,+}\left(p_{0} ; x\right) p_{0}^{2},  \tag{8.26}\\
& K_{j}^{+,-}\left(p_{3} ; x\right) K_{j}^{-,+}\left(p_{4} ; x\right)=\sum_{p_{0} \in I_{q}} B_{p_{0}}\left(p_{3}, p_{4}\right) K_{j}^{-,-}\left(p_{0} ; x\right) p_{0}^{2},  \tag{8.27}\\
& K_{j}^{+,+}\left(p_{1} ; x\right) K_{j}^{-,+}\left(p_{3} ; x\right)=\sum_{p_{0} \in I_{q} \cap(-1,1)} C_{p_{0}}\left(p_{1}, p_{3}\right) K_{j}^{-,+}\left(p_{0} ; x\right) p_{0}^{2},  \tag{8.28}\\
& K_{j}^{-,+}\left(p_{3} ; x\right) K_{j}^{-,-}\left(p_{1} ; x\right)=\sum_{p_{0} \in I_{q} \cap(-1,1)} D_{p_{0}}\left(p_{3}, p_{1}\right) K_{j}^{-,+}\left(p_{0} ; x\right) p_{0}^{2} . \tag{8.29}
\end{align*}
$$

Remark 8.22. Note that the gradings on $N$ and $\hat{N}$ make that the left hand sides of (8.26)-(8.29) consists of a single product of two ${ }_{2} \varphi_{1}$-functions.
Remark 8.23. If $x \in \mu\left(-q^{2 \mathbb{N}+1} \cup q^{2 \mathbb{N}+1}\right)$, so that $W_{1, x}$ is a discrete series corepresentation, then equations (8.27)-(8.29) are trivial, i.e. they equate 0 to 0 . If in addition $x<0$, then (8.26) is also trivial.

Remark 8.24. From (8.22)-(8.25), we can also get formulae for the products

$$
\begin{array}{ll}
K_{j}^{-,-}\left(p_{1} ; x\right) K_{j}^{-,-}\left(p_{2} ; x\right), & K_{j}^{+,-}\left(p_{3} ; x\right) K_{j}^{-,+}\left(p_{4} ; x\right), \\
K_{j}^{+,+}\left(p_{1} ; x\right) K_{j}^{-,+}\left(p_{3} ; x\right), & K_{j}^{-,+}\left(p_{3} ; x\right) K_{j}^{-,-}\left(p_{1} ; x\right),
\end{array}
$$

where $j \in\{1,2\}, p_{1}, p_{2} \in I_{q}$ and $p_{3}, p_{4} \in I_{q} \cap(-1,1)$. However, using the symmetry (8.19), these formulae are already contained in (8.26)-(8.29).

In the remainder of this section we determine the coefficient functions $A, B, C$, and $D$. We explicitly show how to find $A$, the other coefficients can be found by the same method. Let $f=\delta_{p_{0}}$, so that the right hand side of (8.22) is equal to $A_{p_{0}}\left(p_{1}, p_{2}\right) p_{0}^{2}$. To determine the left hand side, note that $\delta_{p_{0}}=\delta_{\operatorname{sgn}\left(p_{0}\right)}(e) \delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)$. Recall that $\Delta(e)=e \otimes e \in N \otimes N$, so that $\Delta\left(\delta_{\operatorname{sgn}\left(p_{0}\right)}(e)\right)=\delta_{\operatorname{sgn}\left(p_{0}\right)}(e \otimes e) \in N \otimes N$. Hence,

$$
\Delta^{\natural}\left(\delta_{\operatorname{sgn}\left(p_{0}\right)}(e)\right)=\Delta\left(\delta_{\operatorname{sgn}\left(p_{0}\right)}(e)\right)=\delta_{1}(e) \otimes \delta_{\operatorname{sgn}\left(p_{0}\right)}(e)+\delta_{-1}(e) \otimes \delta_{-\operatorname{sgn}\left(p_{0}\right)}(e) \in N \otimes N .
$$

Note that by the relation $\Delta(x)=W^{*}(1 \otimes x) W, x \in M$, we find that $\Delta\left(\delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)\right)$ equals the projection onto the closure of

$$
\operatorname{span}\left\{W^{*} f_{m^{\prime}, p^{\prime}, t^{\prime}} \otimes f_{m, p_{0}, t} \mid m, m^{\prime} \in \mathbb{Z}, p^{\prime}, t, t^{\prime} \in I_{q}\right\} .
$$

This projection is given by the formula

$$
\mathcal{K} \otimes \mathcal{K} \rightarrow \mathcal{K} \otimes \mathcal{K}: v \mapsto \sum_{m, m^{\prime} \in \mathbb{Z}, p^{\prime}, t, t^{\prime} \in I_{q}}\left\langle v, W^{*} f_{m^{\prime}, p^{\prime}, t^{\prime}} \otimes f_{m, p_{0}, t}\right\rangle W^{*} f_{m^{\prime}, p^{\prime}, t^{\prime}} \otimes f_{m, p_{0}, t},
$$

where the sum is norm convergent. Note that for $x \in N, \Delta(x) \in M^{\gamma} \otimes M^{\beta}$, so that $\Delta^{\natural}(x)=$ $\left(\iota \otimes T_{\beta} T_{\gamma}\right) \Delta(x)=\left(\iota \otimes\left(T^{+}+T^{-}\right)\right) \Delta(x)$. Since $\Delta^{\natural}\left(\delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)\right) \in N \otimes N$, we find that for any $m_{1}, m_{2} \in \mathbb{Z}$ and $t_{1}, t_{2} \in I_{q}$,

$$
\left(\delta_{p_{1}} \otimes \delta_{p_{2}}\right) \Delta^{\natural}\left(\delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)\right)
$$

$$
\begin{align*}
& =\left\langle\left(\iota \otimes\left(T^{+}+T^{-}\right)\right)\left(\Delta\left(\delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)\right)\right) f_{m_{1}, p_{1}, t_{1}} \otimes f_{m_{2}, p_{2}, t_{2}}, f_{m_{1}, p_{1}, t_{1}} \otimes f_{m_{2}, p_{2}, t_{2}}\right\rangle \\
& =\left\langle\Delta\left(\delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)\right) f_{m_{1}, p_{1}, t_{1}} \otimes f_{m_{2}, p_{2}, t_{2}}, f_{m_{1}, p_{1}, t_{1}} \otimes f_{m_{2}, p_{2}, t_{2}}\right\rangle \\
& =\sum_{m^{\prime}, m \in \mathbb{Z}, p^{\prime}, t, t^{\prime} \in I_{q}}\left|\left\langle f_{m_{1}, p_{1}, t_{1}} \otimes f_{m_{2}, p_{2}, t_{2}}, W^{*} f_{m^{\prime}, p^{\prime}, t^{\prime}} \otimes f_{m, p_{0}, t}\right\rangle\right|^{2}, \tag{8.30}
\end{align*}
$$

where the equations follow from Remark 8.21 and Proposition 8.13, the definition of $T^{+}$and $T^{-}$ and the discussion above. Recall the functions $a_{p}(x, y), p, x, y \in I_{q}$ from [10, Definition 6.2]. From [13, Propositions 4.5 and 4.10], we see that (8.30) equals

$$
\begin{aligned}
\sum_{t \in I_{q}} & \left(\frac{t}{t_{2}}\right)^{2} a_{t}\left(\operatorname{sgn}\left(p_{0} p_{2} t_{2}\right) p_{1} q^{m_{2}} t, t_{2}\right)^{2} a_{p_{0}}\left(p_{1}, p_{2}\right)^{2} \\
& =\sum_{t \in I_{q}} a_{t_{2}}\left(\operatorname{sgn}\left(t_{2}\right)\left|p_{1}\right| q^{m_{2}} t, t\right)^{2} a_{p_{0}}\left(p_{1}, p_{2}\right)^{2}=a_{p_{0}}\left(p_{1}, p_{2}\right)^{2} .
\end{aligned}
$$

Here, the first equality follows from [10, equation (24)] and the fact that by definition $a_{p_{0}}\left(p_{1}, p_{2}\right)^{2}$ $=0$ if $\operatorname{sgn}\left(p_{0} p_{2}\right) \neq \operatorname{sgn}\left(p_{1}\right)$. The last equality follows from [10, Proposition 6.3]. Hence, using [30, Section IX.4, equation (4)] in the second equality,

$$
\begin{aligned}
A_{p_{0}}\left(p_{1}, p_{2}\right) p_{0}^{2} & =\left(\delta_{p_{1}} \otimes \delta_{p_{2}}\right) \Delta^{\natural}\left(\delta_{\operatorname{sgn}\left(p_{0}\right)}(e) \delta_{p_{0}^{-2}}\left(\gamma^{*} \gamma\right)\right) \\
& =\left(\delta_{p_{1}} \otimes \delta_{p_{2}}\right)\left(\iota \otimes T_{\gamma}\right)\left(\delta_{\operatorname{sgn}\left(p_{0}\right)}(\Delta(e)) \delta_{p_{0}^{-2}} \Delta\left(\gamma^{*} \gamma\right)\right) \\
& =\left(\delta_{p_{1}} \otimes \delta_{p_{2}}\right) \delta_{\operatorname{sgn}\left(p_{0}\right)}(\Delta(e))\left(\iota \otimes T_{\gamma}\right)\left(\delta_{p_{0}^{-2}} \Delta\left(\gamma^{*} \gamma\right)\right) \\
& =\delta_{\operatorname{sgn}\left(p_{0}\right), \operatorname{sgn}\left(p_{1} p_{2}\right)} a_{p_{0}}\left(p_{1}, p_{2}\right)^{2}=a_{p_{0}}\left(p_{1}, p_{2}\right)^{2} .
\end{aligned}
$$

Recall [13, Definition 3.1] that by definition $a_{p_{0}}\left(p_{1}, p_{2}\right)=0$ if $\delta_{\operatorname{sgn}\left(p_{0}\right), \operatorname{sgn}\left(p_{1} p_{2}\right)}=0$, so indeed the last equality follows. By a similar computation we get,

$$
\begin{aligned}
& B_{p_{0}}\left(p_{1}, p_{2}\right) p_{0}^{2}=a_{p_{0}}\left(p_{1}, p_{2}\right) a_{p_{0}}\left(-p_{1},-p_{2}\right), \\
& C_{p_{0}}\left(p_{1}, p_{2}\right) p_{0}^{2}=a_{p_{0}}\left(p_{1},-p_{2}\right) a_{-p_{0}}\left(p_{1}, p_{2}\right), \\
& D_{p_{0}}\left(p_{1}, p_{2}\right) p_{0}^{2}=a_{p_{0}}\left(-p_{1}, p_{2}\right) a_{-p_{0}}\left(p_{1}, p_{2}\right) .
\end{aligned}
$$

Remark 8.25. In particular, the sums in (8.26)-(8.29) run only through either the positive or the negative numbers, since $a_{z}(x, y)=0$ if $\operatorname{sgn}(x y z)=-1$.

Remark 8.26. It is also possible to obtain (8.22)-(8.25) from [10, Proposition 4.10], using the pairing between $M$ and $\left\{\lambda(\omega) \mid \omega \in M_{*}\right\} \subseteq \hat{M}$, defined by $\langle x, \lambda(\omega)\rangle=\omega(x)$.

Remark 8.27. In case of the group $S U(1,1)$ we know [43, Chapter 6] that there exists an addition formula corresponding to the product formula, i.e. the product formula corresponds to the constant term in the addition formula. It would be of interest to obtain addition formulae corresponding to (8.26)-(8.29).

## Acknowledgement

The author likes to thank Erik Koelink for the useful discussions and Noud Aldenhoven for providing Fig. 1. Also, the author benefits from a detailed referee report.

## References

[1] Caspers M., The $L^{p}$-Fourier transform on locally compact quantum groups, J. Operator Theory, to appear, arXiv:1008.2603.
[2] Caspers M., Koelink E., Modular properties of matrix coefficients of corepresentations of a locally compact quantum group, J. Lie Theory 21 (2011), 905-928, arXiv:1003.2278.
[3] De Commer K., On a correspondence between quantum $S U(2)$, quantum $E(2)$ and extended quantum $S U(1,1)$, Comm. Math. Phys. 304 (2011), 187-228, arXiv:1004.4307.
[4] Desmedt P., Aspects of the theory of locally compact quantum groups: Amenability-Plancherel measure, Ph.D. Thesis, Katholieke Universiteit Leuven, 2003.
[5] van Dijk G., Introduction to harmonic analysis and generalized Gelfand pairs, de Gruyter Studies in Mathematics, Vol. 36, Walter de Gruyter \& Co., Berlin, 2009.
[6] Dixmier J., Les algèbres d'opérateurs dans l'espace hilbertien, Gauthiers-Villars, Paris, 1957.
[7] Dixmier J., Les $C^{*}$-algèbres et leurs représentations, Gauthiers-Villars, Paris, 1969.
[8] Faraut J., Analyse harmonique sur les paires de Guelfand et les espaces hyperboliques, in Analyse Harmonique (Universite de Nancy I, 1980), Editors J.L. Clerc et al., Les Cours du C.I.M.P.A., Centre International de Mathematiques Pures et Appliquees, Nice, France, 1983, 315-446.
[9] Floris P.G.A., Gel'fand pair criteria for compact matrix quantum groups, Indag. Math. (N.S.) 6 (1995), 83-98.
[10] Groenevelt W., Koelink E., Kustermans J., The dual quantum group for the quantum group analogue of the normalizer of $S U(1,1)$ in $S L(2, \mathbb{C})$, Int. Math. Res. Not. 2010 (2010), no. 7, 1167-1314, arXiv:0905.2830.
[11] Kadison R.V., Ringrose J.R., Fundamentals of the theory of operator algebras, Vol. II, Advanced theory, Graduate Studies in Mathematics, Vol. 16, American Mathematical Society, Providence, RI, 1997.
[12] Kasprzak P., On a certain approach to quantum homogeneous spaces, arXiv:1007.2438.
[13] Koelink E., Kustermans J., A locally compact quantum group analogue of the normalizer of $\mathrm{SU}(1,1)$ in SL(2, C), Comm. Math. Phys. 233 (2003), 231-296, math.QA/0105117.
[14] Koelink E., Stokman J.V., Fourier transforms on the quantum $S U(1,1)$ group, With an appendix by Mizan Rahman, Publ. Res. Inst. Math. Sci. 37 (2001), 621-715, math.QA/9911163.
[15] Koelink H.T., Askey-Wilson polynomials and the quantum $S U(2)$ group: survey and applications, Acta Appl. Math. 44 (1996), 295-352.
[16] Koelink H.T., The quantum group of plane motions and the Hahn-Exton $q$-Bessel function, Duke Math. J. 76 (1994), 483-508.
[17] Koornwinder T.H., Positive convolution structures associated with quantum groups, in Probability Measures on Groups, X (Oberwolfach, 1990), Plenum, New York, 1991, 249-268.
[18] Kustermans J., Locally compact quantum groups in the universal setting, Internat. J. Math. 12 (2001), 289-338, math.OA/9902015.
[19] Kustermans J., Locally compact quantum groups, in Quantum Independent Increment Processes. I, Lecture Notes in Math., Vol. 1865, Springer, Berlin, 2005, 99-180.
[20] Kustermans J., One-parameter representations of C*-algebras, funct-an/9707009.
[21] Kustermans J., Vaes S., Locally compact quantum groups, Ann. Sci. École Norm. Sup. (4) 33 (2000), 837-934.
[22] Kustermans J., Vaes S., Locally compact quantum groups in the von Neumann algebraic setting, Math. Scand. 92 (2003), 68-92, math.OA/0005219.
[23] Lance C., Direct integrals of left Hilbert algebras, Math. Ann. 216 (1975), 11-28.
[24] Noumi M., Yamada H., Mimachi K., Finite-dimensional representations of the quantum group $G L_{q}(n, \mathbb{C})$ and the zonal spherical functions on $U_{q}(n-1) \backslash U_{q}(n)$, Japan. J. Math. (N.S.) 19 (1993), 31-80.
[25] Nussbaum A.E., Reduction theory for unbounded closed operators in Hilbert space, Duke Math. J. 31 (1964), 33-44.
[26] Podkolzin G.B., Vainerman L.I., Quantum Stiefel manifold and double cosets of quantum unitary group, Pacific J. Math. 188 (1999), 179-199.
[27] Raeburn I., Williams D.P., Morita equivalence and continuous-trace $C^{*}$-algebras, Mathematical Surveys and Monographs, Vol. 60, American Mathematical Society, Providence, RI, 1998.
[28] Salmi P., Skalski A., Idempotent states on locally compact quantum groups, Quart. J. Math., to appear, arXiv:1102.2051.
[29] Takesaki M., Theory of operator algebras. I, Springer-Verlag, New York - Heidelberg, 1979.
[30] Takesaki M., Theory of operator algebras. II, Springer-Verlag, Berlin, 2003.
[31] Timmermann T., An invitation to quantum groups and duality. From Hopf algebras to multiplicative unitaries and beyond, EMS Textbooks in Mathematics, European Mathematical Society (EMS), Zürich, 2008.
[32] Tomatsu R., A characterization of right coideals of quotient type and its application to classification of Poisson boundaries, Comm. Math. Phys. 275 (2007), 271-296, math.OA/0611327.
[33] Vaes S., A new approach to induction and imprimitivity results, J. Funct. Anal. 229 (2005), 317-374, math.OA/0407525.
[34] Vaes S., A Radon-Nikodym theorem for von Neumann algebras, J. Operator Theory 46 (2001), 477-489, math.OA/9811122.
[35] Vaes S., The unitary implementation of a locally compact quantum group action, J. Funct. Anal. 180 (2001), 426-480, math.OA/0005262.
[36] Vaes S., Vainerman L.I., Extensions of locally compact quantum groups and the bicrossed product construction, Adv. Math. 175 (2003), 1-101, math.OA/0101133.
[37] Vaes S., Vainerman L.I., On low-dimensional locally compact quantum groups, in Locally Compact Quantum Groups and Groupoids (Strasbourg, 2002), IRMA Lect. Math. Theor. Phys., Vol. 2, de Gruyter, Berlin, 2003, 127-187, math.QA/0207271.
[38] Vainerman L.I., Gel'fand pair associated with the quantum group of motions of the plane and $q$-Bessel functions, Rep. Math. Phys. 35 (1995), 303-326.
[39] Vainerman L.I., Gel'fand pairs of quantum groups, hypergroups and $q$-special functions, in Applications of Hypergroups and Related Measure Algebras (Seattle, WA, 1993), Contemp. Math., Vol. 183, Amer. Math. Soc., Providence, RI, 1995, 373-394.
[40] Vainerman L.I., Hypergroup structures associated with Gel'fand pairs of compact quantum groups, in Recent Advances in Operator Algebras (Orléans, 1992), Asterisque (1995), no. 232, 231-242.
[41] Van Daele A., The Fourier transform in quantum group theory, in New Techniques in Hopf Algebras and Graded Ring Theory, K. Vlaam. Acad. Belgie Wet. Kunsten (KVAB), Brussels, 2007, 187-196, math.RA/0609502.
[42] Van Daele A., Locally compact quantum groups. A von Neumann algebra approach, math.OA/0602212.
[43] Vilenkin N.Ja., Klimyk A.U., Representations of Lie groups and special functions, Vol. 1, Simplest Lie groups, special functions and integral transforms, Kluwer Academic Publishers Group, 1991.
[44] Vogan D.A., Representations of real reductive Lie groups, Progress in Mathematics, Vol. 15, Birkhäuser, Boston, Mass., 1981.


[^0]:    *This paper is a contribution to the Special Issue "Relationship of Orthogonal Polynomials and Special Functions with Quantum Groups and Integrable Systems". The full collection is available at http://www.emis.de/journals/SIGMA/OPSF.html

