

# **Engineering Lipid-stabilized Microbubbles for Magnetic Resonance Imaging guided Focused Ultrasound Surgery**

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# **ABSTRACT**

## **Engineering Lipid-stabilized Microbubbles for Magnetic Resonance Imaging guided Focused Ultrasound Surgery**

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Lipid-stabilized microbubbles are gas-filled microspheres encapsulated with a phospholipid monolayer shell. Because of the high echogenicity provided by its highly compressible gas core, these microbubbles have been adapted as ultrasound contrast agents for a variety of applications such as contrast-enhanced ultrasonography (CEUS), targeted drug delivery and metabolic gas transport. Recently, these lipid-stabilized microbubbles have demonstrated increased potential as theranostic (therapy + diagnostics) agents for non-invasive surgery with focused ultrasound (FUS). For instance, their implementation has reduced the acoustic intensity threshold needed to open the blood-brain-barrier (BBB) with FUS, which potentially allows for the localized delivery of drugs to treat neurodegenerative diseases such as Alzheimer's, Parkinson's and Huntington's diseases. However, the effectiveness of microbubbles for this application is dependent on successful microbubble engineering. One necessary improvement is the development and utilization of monodisperse microbubbles of varying size classes. Another design improvement is the development of a microbubble construct whose fragmentation state during or after FUS surgery can be tracked by magnetic resonance imaging (MRI).

Thus, in this thesis, we describe a method to generate and select lipid-coated gas-filled microbubbles of specific size fractions based on their migration in a centrifugal field. We also detail the design and characterization of size-selected lipid-coated microbubbles with shells

containing the magnetic resonance (MR) contrast media Gadolinium (Gd(III)), for utility in both MR and ultrasound imaging. Initial characterization of the lipid headgroup labeled Gd(III)-microbubbles by MRI revealed that the Gd(III) relaxivity increased after microbubble fragmentation into non-gas-containing lipid vesicles. This behavior was explained to stem from an increase in interaction between water protons and the Gd(III)-bound lipid fragments due to an increase in lipid headgroup area after microbubble fragmentation. To explore this hypothesis, an alternative construct consisting of Gd(III) preferentially bound to the protective poly(ethylene glycol) (PEG) brush of the lipid shell architecture was also designed and compared to the lipid headgroup-labeled Gd(III) microbubbles. Nuclear magnetic resonance (NMR) analysis revealed that, in contrast to the headgroup labeled Gd(III) microbubbles, the relaxivity of the PEG-labeled Gd(III)-microbubbles decreased after microbubble fragmentation. NMR analysis also revealed an independent concentration-dependent enhancement of the transverse MR signal by virtue of the microbubble gas core. The results of this study illustrated the roles that Gd(III) placement on the lipid shell and the presence of the gas core may play on the MR signal when monitoring Gd(III)-microbubble cavitation during non-invasive surgery with FUS.

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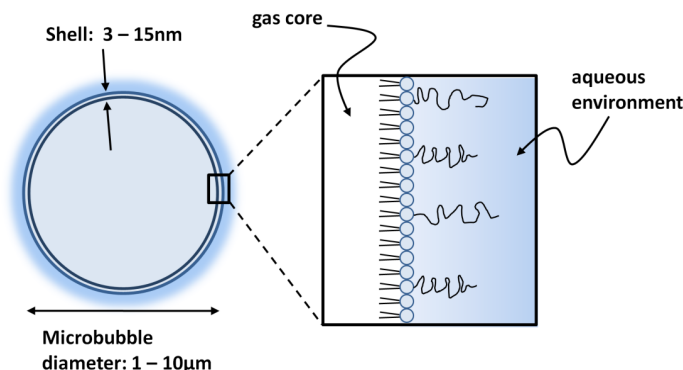
## Chapter 1: Introduction

Gas-filled microbubbles are 1 to 100  $\mu\text{m}$  spheres that create a colloidal suspension consisting of gas globules dispersed in a liquid medium. These microbubbles occur naturally in fresh or seawater, in the fluid within our bodies, or are synthetically manufactured. Although adaptable in a wide range of applications, their use in the medical field is currently gaining ground. For example, microbubbles have been studied as tools for contrast-enhanced ultrasonography (CEUS), targeted drug/gene delivery and metabolic gas delivery. Another developing medical application with microbubbles involves their application as theranostic (therapy + diagnostic) agents for non-invasive surgery using focused ultrasound (FUS). Ultrasound in combination with microbubble cavitation can focus the mechanical energy to microscale events distributed throughout the insonified vasculature, resulting in violent or subtle changes to the local environment. A combination of this effect with the capability of microbubbles to simultaneously offer contrast enhancement for diagnostic imaging and drug loading, as well as delivering a payload for targeted therapy, makes for a promising medical tool (Sirsi and Borden 2009). However, there is still a need to improve the design of microbubbles for their successful application in FUS surgery. These improvements may include the choice of surfactant shell type, reducing microbubble size polydispersity, gas choice and the ability to monitor therapy with techniques like MRI. Although a variety of microbubble shell materials may be selected, lipid-shelled microbubbles are preferred due to their ultrasound compliance (Bloch, Wan et al. 2004; Ferrara, Pollard et al. 2007). Thus, in this chapter, we discuss the structure and composition of lipid-stabilized microbubbles, their fabrication methods, the importance of size monodispersity and the need for a MRI-detectable microbubble design to aid in tracking of the microbubble or its lipid shell debris during non-invasive surgery with FUS.



### 1.1. Lipid-shelled Microbubble Structure and Composition

Lipid-shelled microbubbles are typically composed of two key components: the amphiphilic phospholipid moiety and a poly(ethylene glycol) (PEG) labeled lipid emulsifier or lipopolymer (Figure 1.1). The inclusion of the PEGylated lipid into the shell improves microbubble stability.



**Figure 1.1. Cartoon illustrating the microbubble lipid-shelled monolayer**

In a comparable engineering design to long-circulation liposomes, the brush layer of PEG chains is an important part of the lipid microbubble shell structure because it forms a steric barrier against coalescence and adsorption of macromolecules to the microbubble surface (Klibanov, Maruyama et al. 1990; Klibanov 1999). The protective role of PEG is assumed to be due to the steric hindrance effect of the polymer brush, which forms a semi-impenetrable barrier over part of the microbubble surface that partly inhibits certain molecules from diffusing into the brush layer (Needham, McIntosh et al. 1998).

The main phospholipid component of the shell consists of long hydrophobic fatty acid tails and a hydrophilic polar headgroup. During microbubble formation, these phospholipids spontaneously absorb from lipid vesicles in solution such as micelle and liposomes, to form a highly oriented and densely packed monolayer at the gas-liquid interface, such that their hydrophobic tail faces the gas and their hydrophilic headgroups faces the aqueous medium.

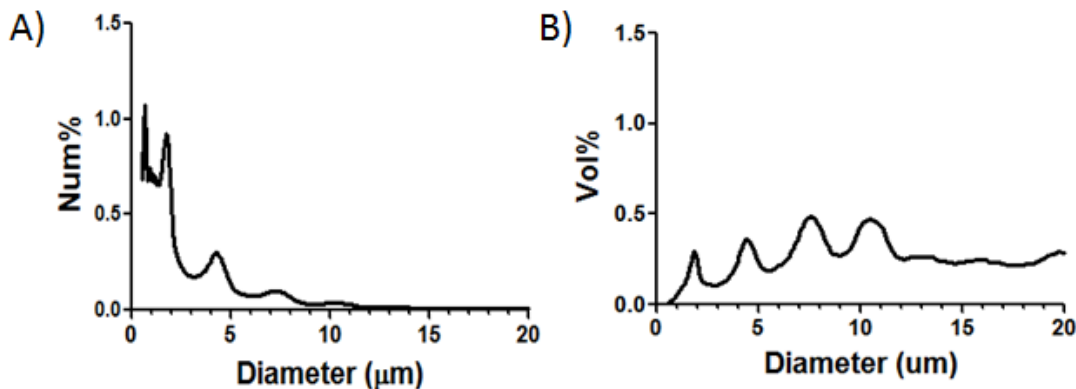
Phospholipids below their main phase transition temperature ( $T_m$ ), defined as the temperature at which the membrane transforms from a crystalline state to a liquid crystalline, can achieve lateral compression within the monolayer plane that results in very low surface tension (Ferrara, Pollard et al. 2007). The length of the lipid chain is integral to microbubble stability and affects its mechanical properties. For instance, increasing the length of the lipid chain has been found to reduce surface tension and increase shell surface viscosity, buckling stability and gas permeation resistance (Borden and Longo 2002; Kim, Costello et al. 2003; Borden and Longo 2004; Duncan and Needham 2004; Pu, Borden et al. 2006; Kwan and Borden 2010).

Research has also revealed the complex interaction of the lipid shell constituents stabilizing the microbubble. Kim et al. (Kim, Costello et al. 2003) revealed the lipid shell to consist of planar microdomain phases separated by regions of defects. Borden et al. (Borden, Martinez et al. 2006) later elucidated the nature of the phase separation on the shell. The microdomains were composed mainly of lecithin and were segregated by grain boundary regions consisting of a miscible phase rich in other monolayer constituents such as lipopolymers. Both phases are integral in the stability of lipid-stabilized microbubbles (Pu, Borden et al. 2006; Talu, Lozano et al. 2006; Ferrara, Pollard et al. 2007).

## **1.2. Fabrication Techniques**

The most common technique for generating high yields of lipid-coated microbubbles is by mechanical agitation (Feinstein, Shah et al. 1984; Feinstein 2004). This process can be subdivided into two techniques, shaking or acoustic emulsification (sonication) at the interface of a gas and liquid. Both techniques can produce microbubble suspensions rapidly (~10-40 seconds) with relatively high concentration ( $10^9$ — $10^{10}$  mL<sup>-1</sup>) and containing a polydisperse size

distribution, ranging between 1-100 microns in diameter. For biomedical applications requiring small quantities of microbubbles, the shaking method (Unger *et al.* 1992), allows for on-site generation of microbubbles by vigorously shaking a serum vial (2 – 3 mL) containing lipid solution and gas headspace. This technique has been adapted in the development of commercially available lipid microbubbles such as Definity<sup>TM</sup> (Lantheus Medical Imaging, North Billerica, MA, USA). The advantage of this technique is that the pre-microbubble lipid solution can be stored, and the microbubble suspension can be later prepared using equipment such as a dental amalgamator. Alternatively, the sonication technique is generally used for laboratory applications requiring large volumes of polydisperse microbubbles (figure 1.2). In this methodology, microbubbles are generated by mechanical perturbation produced by an ultrasonic probe sonifier (at about 10 kHz), whose tip is placed close to the interface of the lipid suspension and the overhead gas space. The technique has been improved in recent years. For instance, Swanson *et al.* (Swanson, Mohan *et al.* 2010) adapted the sonication technique to semi-continuously create lipid-coated oxygen microbubble suspensions at concentrations up to  $10^{10}$  mL<sup>-1</sup> within a few minutes and at volumes up to 1 L.



**Figure 1.2. Representative size distribution of freshly made microbubbles (prepared by sonication) measured by light obscuration.**

Other more refined techniques have been developed to produce microbubbles. These include the adaptation of T-junctions (Xu, Nie et al. 2005), flow focusing (Talu, Lozano et al. 2006; Talu, Hettiarachchi et al. 2008), ink-jet printing (Bohmer, Schroeders et al. 2006) and coaxial electrohydrodynamic atomization (Stride and Edirisinghe 2008). These techniques mainly make use of microfluidic processing methods to produce microbubbles with narrow size distributions. However, these methodologies lack the microbubble production capacity when compared to the sonication method.

### **1.3. Microbubble Size-dependent behavior**

Microbubble size plays a big role on its behavior under ultrasound. For example, Apfel and Holland (Apfel and Holland 1991) developed a Blake threshold, for the onset of inertial cavitation of a bubble during pulsed insonation at higher frequencies, by solving mathematical models of bubble dynamics. Their calculation demonstrated that the Blake mechanism causes unstable growth of smaller bubbles, while liquid inertia restricts the growth of larger bubbles. Chomas et al. (Chomas, Dayton et al. 2001) also linked the ultrasound parameters needed to

cause bubble fragmentation with bubble size by driving optically observed microbubbles of different sizes under a range of ultrasound pressures. In their analysis, a fragmentation threshold was developed to delineate microbubbles that remained intact during ultrasound-induced cavitation from those that were destroyed.

There is also a size-dependent behavior of microbubbles for diagnostic imaging applications. In 1933, Minnaert (Minnaert 1933) developed a relationship between acoustic resonant frequency,  $f$  to bubble radius,  $R$  in an infinite medium neglecting surface tension and viscosity.

$$f = \frac{1}{2\pi R} \left( \frac{3\gamma P_A}{\rho} \right)^{\frac{1}{2}} \quad (1.2)$$

where  $\gamma$  represents the polytropic coefficient,  $\rho$  represents the density of water, and  $P_A$  represents ambient pressure. One example of an implication of this relation is that increasing the microbubble diameter from 1 to 5 microns will change the resonance frequency of an unencapsulated microbubble from 4.7 to 0.72 MHz. (Wu and Nyborg 2008). In a related study, Goertz et al. (Goertz, de Jong et al. 2007) showed that, as predicted by a Rayleigh-Plesset model, the acoustic response of small (< 2 microns) lipid-coated microbubble led to greater acoustic attenuation (extinction) at higher frequencies because they were closer to resonance.

Microbubble size also affects the biodistribution and pharmacodynamics after intravenous injection, the bioeffects during ultrasound insonification, the gas release profile, and other related behaviors. For instance, Bouakaz et al. (Bouakaz, de Jong et al. 1998) demonstrated that microbubbles larger than the lung mean capillary (~5 microns) diameter were filtered out. Other studies using positron emission tomography have shown that a large portion of microbubbles accumulate in the spleen via a potential size dependent filtration mechanism

(Tartis, Kruse et al. 2008). Clearly, for both diagnostic imaging and therapeutic applications microbubbles size control is necessary for their successful development as theranostic agents.

#### **1.4. Attempts to Reduce Size Polydispersity**

Most commercially available microbubble contrast agents are highly polydisperse in size. For example, the commercially available lipid-based microbubble, Definity, contains microbubbles ranging from 1 to 10 microns in diameter, with most of the microbubbles below 2 microns (Goertz, de Jong et al. 2007). The high polydispersity is a consequence of the emulsification methods (as described in section 1.2) used to generate microbubbles in high quantity, such as sonication and shaking. In most cases, the size distribution is broad over a range of sub-micrometer to tens of micrometer in diameter. Consequently, a number of methods are being developed to improve the control over microbubble size and other characteristics. These techniques include flow focusing (Talu, Lozano et al. 2006; Talu, Hettiarachchi et al. 2007; Talu, Hettiarachchi et al. 2008), T-junctions (Xu, Nie et al. 2005) and electrohydrodynamic atomization (Farook, Stride et al. 2007; Farook, Zhang et al. 2007). While these techniques provide very low polydispersity, they are rather slow at generating microbubbles. Using flow focusing, for example, requires several hours to produce microbubbles at sufficient numbers for even a single small-animal trial ( $\sim 0.1 \text{ mL} \times 10^9 \text{ mL}^{-1}$ ). Additionally, dust particles can plug microchannels, thus requiring fabrication and calibration of a new device. Although engineering breakthroughs may eventually allow efficient and robust generation of monodisperse microbubbles via microfluidic strategies, these techniques currently remain untenable for biomedical studies.

Another strategy takes advantage of microbubble buoyancy to isolate monodisperse microbubble populations. In principle, larger microbubbles are more buoyant and rise faster, thus allowing separation possible based on different migration rates in a gravitational field. Kvale et al. (Kvale, Jakobsen et al. 1996) described a model for the size fractionation of air-filled microbubbles by simple flotation. Microbubbles were injected at the bottom of a stagnant water column and allowed to rise under normal gravity. The model predicted the size distribution of microbubbles at certain distances from the bottom of the column as a function of time. The form of the model was a second-order PDE (equation 1.3) that used Stokes velocity (equation 1.4) to account for the convective motion of the bulk dispersed phase (liquid moved down the column as microbubbles moved up) as well as the Stokes-Einstein diffusion constant (equation 1.5) to account for Brownian (thermal) diffusive motion of the particles.

$$\frac{\partial n_i}{\partial t} = -u_i \frac{\partial n_i}{\partial x} + D_{iB} \frac{\partial^2 n_i}{\partial x^2} \quad (1.3)$$

$$u_i^r = \frac{2(\rho - \rho_{il})}{9\mu^*} R_i^2 g \quad (1.4)$$

$$D_{iB} = \frac{kT}{6\pi\mu R} \quad (1.5)$$

where  $u_i^r$  represents the relative velocity between the microbubble and the fluid,  $n$  represents the number of microbubbles,  $t$  represents time,  $x$  represents vertical position,  $R_i$  represents microbubble radius,  $g$  represents gravity,  $D_{iB}$  represents diffusion constant,  $\mu^*$  represents effective viscosity,  $\rho - \rho_{il}$  represents the density difference between the bubble and the fluid,  $R$  represents gas constant,  $T$  represents ambient fluid temperature and  $k$ , represents Boltzmann's constant.

Wheatley et al. (Wheatley, Forsberg et al. 2006), described a methodology to separate submicron population of bubbles from larger ones by using predetermined centrifugal speeds. This centrifugation technique reduced the separation time as compared to the gravity separation protocol. By adapting these protocols, it is conceivable to develop a technique to select different microbubble size classes by varying centrifugal speed.

## **1.5. Applications of Microbubbles in Focused Ultrasound Surgery**

### **1.5.1. Focused Ultrasound Surgery with Microbubbles**

Focused ultrasound (FUS) surgery has become a promising medical technique since it provides a truly non-invasive surgical method without the unwanted side effects associated with other treatment option like chemotherapy and radiotherapy (Kennedy, ter Haar et al. 2003). When ultrasonic waves are focused using transducers with spherical geometries, or beam steering techniques, the ultrasonic waves align to form a focal region a certain distance from the plane of the transducer leading to localization of the acoustic energy to a target region without significantly perturbing the surrounding media. The ability of high intensity FUS to induce lesion formation for non-invasive surgery is based on the observation that above a certain thermal threshold and sonication time, irreversible cell death occurs through “coagulative necrosis” (Kennedy, ter Haar et al. 2003). Ultrasound can be used to generate non-thermal bioeffects like the acoustic cavitation generated by bubbles (Kennedy, ter Haar et al. 2003; Jolesz and McDannold 2008). At high acoustic intensities, these gas bubbles can form in the focal area during the rarefaction phase of ultrasound. These inception bubbles oscillate in size or undergo inertial collapse, causing mechanical stresses and generating temperatures exceeding 2000 K in



the microenvironment (Kennedy, ter Haar et al. 2003; Hindley, Gedroyc et al. 2004), which further aids the therapy.

While heating with inception microbubble cavitation may provide some advantages over heating alone, it has its disadvantages. For instance, the forces generated from inertial cavitation from an inception bubble can cause significant tissue damage. Moreover, the uncertainty in formation and location of inception cavitation may result in tissue damage outside the desired region of interest. One method to mitigate the potential problems with cavitation-enhanced FUS is to introduce preformed lipid-stabilized microbubbles into the vasculature (Jolesz and McDannold 2008). Since these microbubbles, are already commercially available as ultrasound contrast agents, they can substantially reduce the threshold for inertial cavitation and further reduce the acoustic power requirements needed to ablate tissue. These preformed microbubbles have been demonstrated in animals to aid lesion formation through non-thermal mechanisms. For example, Fujishiro et al. (Fujishiro, Mitsumori et al. 1998) demonstrated an increase ultrasound-induced heating efficiency of a 2 centimeter beef tissue phantom with the administration of air-filled protein shelled microbubbles. In a related study, Miller et al. (Miller and Gies 1998) demonstrated the importance of microbubble persistence in their efficacy for ultrasound-induced cavitation. Generally, perflouorocarbon-containing microbubbles were reported to be more effective in inducing hemolysis in the canine blood than air-containing microbubbles. McDannold et al. (McDannold, Vykhodtseva et al. 2006) showed that, upon administration of protein-shelled microbubbles (Optison, GE Healthcare, Milwaukee, WI), lesion formation in the rabbit brain could be attained a temperature below the threshold for thermal damage. Finally, in a study comparing different microbubbles shell architectures, Takegami et al. (Takegami, Kaneko et al. 2005) demonstrated that lipid-shelled microbubbles (MRX-133,

ImaRX, Tuscon, Arizona) were more efficient in inducing lesions in the rabbit liver than protein-shelled counterparts. Clearly, varying parameters in microbubble design such as the choice of shell, gas-type and size are important considerations in the development of microbubbles for non-invasive surgery with FUS.

### **1.5.2. Magnetic Resonance Imaging (MRI)-guided Focused Ultrasound Surgery with Microbubbles**

Another important consideration in the clinical application of FUS surgery is the ability to monitor treatment accurately. In modern practice this is achieved by either using real-time ultrasound (Madersbacher, Kratzik et al. 1994; Wu, Chen et al. 2001), or MRI (Cline, Schenck et al. 1992; Hynynen, Darkazanli et al. 1993; Cline, Hynynen et al. 1995; Hynynen, Freund et al. 1996). The advantage in using MRI to guide FUS therapy is the capability of producing better image quality and the ability to monitor changes in temperature. This treatment is called MRI-guided focused ultrasound surgery (MRIgFUS). In this method, the target region is initially identified with “sublesioning” ultrasound exposures (Jolesz and McDannold 2008). Subsequently, the local rise in temperature is used to validate the position of the ultrasound focus before higher intensity pulses are used for the actual treatment. The Exablate 2000 is an example MRIgFUS device, developed by Insightec Ltd, that is currently FDA approved to treat uterine fibroids, a benign tumor found in the uterus. This device allows the operator to monitor and treat the tumor in real time via a thermal feedback control loop. McDannold *et al.* (McDannold, Vykhodtseva et al. 2006) demonstrated that the initiation and level of lesion formation in rabbits, created using the combination of FUS and an ultrasound contrast agent (Optison), correlated with temperature contours observed in MR thermometry. Moreover, the study revealed that the

administration of the microbubbles reduced the acoustic threshold needed to produce the lesions by 91 % as compared to without microbubbles. While MRIgFUS is currently under development to treat a variety of other diseases found in the liver, uterus, kidney and bone, one of the most ideal applications is to treat diseases found in the brain (Jolesz and McDannold 2008).

### **1.5.3. Blood-Brain-Barrier Opening with Focused Ultrasound and Microbubbles**

A method that allows non-invasive and efficient delivery of drugs to anatomically desired regions of the brain currently does not exist. The key challenge in this development is severe limitation imposed by action of the blood-brain-barrier (BBB). In general, the BBB is a permeability barrier that prevents most large compounds ( $> 400$  Da) from crossing into the interstitial space of the brain (Pardridge 2005). Thus, it serves as a major bottleneck for drug delivery to the brain to treat neurodegenerative diseases such as Parkinson's, Huntington's and Alzheimer's, as most effective drugs are larger than the BBB size cut-off. While other invasive or non-localized strategies are under development to overcome the BBB, FUS with pre-administered microbubbles represents one of the only non-invasive and localized methodologies available (Choi, Pernot et al. 2007). The intravenous addition of microbubble before sonication has improved safety of FUS-induced BBB opening since it is not associated with significant neuronal damage. BBB opening with FUS and microbubbles relies mainly on non-thermal mechanical perturbation mechanisms caused by the interaction of the microbubbles with ultrasound as opposed to the predominantly thermal effects used to induce tissue ablation in MRIgFUS. Furthermore, the BBB disruption is reversible since it recovers within a few hours. The presence of the microbubbles in the blood supply also offers the advantage of reduction in

the acoustic driving pressure needed for non-thermal cell disruption. Hynynen *et al.* (Hynynen, McDannold *et al.* 2001) demonstrated that FUS sonication in the presence of pre-injected Optison microbubbles temporarily opens the BBB in rabbits at acoustic intensity levels lower than that required for thermal ablation. Choi *et al.* (Choi, Pernot *et al.* 2007) further demonstrated the non-invasive, localized and transient opening of the BBB in mice at lower acoustic driving pressures, with the administration of the lipid-stabilized microbubble Definity.

#### **1.5.4. The Need for MRI-detectable Microbubbles in MRI-guided Focused Ultrasound Surgery and Blood-Brain Barrier Opening with Focused Ultrasound**

Until now, the microbubble sizes used for FUS-induced BBB opening were generally polydisperse in nature (Figure 1.2), similar to those formed by sonication or shaking, thus resulting in the unknown dependence of bubble size on FUS-induced BBB opening. Additionally, owing to the different cavitation modes of microbubbles during FUS sonications, the state of the microbubble construct during or after therapy is not known. Providing solutions to these problems begins with proper microbubble engineering. For instance, the dependence of microbubble size on BBB opening may be elucidated by developing and testing the effects of different-sized monodisperse microbubbles. The monodisperse microbubbles size classes could be prepared using the centrifugation technique proposed in section 1.4. Furthermore, the question of how to track the microbubble fragmentation state during or after FUS-induced BBB opening or MRIGFUS may be answered by taking advantage of the techniques used to determine the opening and closing of the BBB. In most BBB opening studies, this involves monitoring MR signal intensity changes after administration of MRI contrast media, typically Gadolinium Gd(III) based (Hynynen, McDannold *et al.* 2001; Choi, Pernot *et al.* 2007). With an open BBB,

the injection of the MRI contrast media after microbubble-mediated BBB opening results in enhancement in the MRI signal intensity in the sonicated region. This is because these Gd(III) molecules are able to diffuse into the extracellular space, previously unreachable when the BBB is closed. One may gain insight on the final state of the lipid-stabilized microbubble by developing a contrast agent that combines both imaging modalities, ultrasound with microbubbles and MRI with Gd(III). Thus, the development of size-selected microbubbles with Gd(III) bound lipid shells could be beneficial to advance both MRIgFUS and BBB opening with FUS. The development of this dual-modal contrast agent may also give information on shell-associated biodistribution and pharmacokinetics for drug delivery to the brain, since the Gd(III)-bound lipid shell can serve as a potential drug surrogate. Finally, by monitoring the effect of microbubble destruction on the MRI signal, one may conceivably use Gd(III) microbubble fragmentation as a guide for FUS surgery as opposed to MR thermometry used in MRIgFUS.

## 1.6. Gadolinium as an MRI contrast agent

Before development of a dual-modal construct for ultrasound and MRI, it is necessary to understand the functionality of the contrast media conventionally used in MRI. Historically, MRI contrast agents play a significant role in the development of MRI for medical applications. They are commonly referred to as  $T_1$ - or  $T_2$ -agents based on whether the relative reduction in relaxation times caused by the agent is larger for the longitudinal ( $T_1$ ) or transverse ( $T_2$ ) relaxation (Toth, Helm et al. 2002). MR signal intensity increases with increasing  $1/T_1$  (brighter  $T_1$ -weighted images) and reduces with increasing  $1/T_2$  (darker  $T_2$ -weighted images). While a variety of other contrast agents exist, such as iron-oxide based particles, the most popularly used, especially for  $T_1$ -weighted MRI, are derived from stable chelates of Gadolinium (Gd(III)).

Gd(III) is the ideal contrast agent because of its slow electron spin relaxation and because its seven unpaired electrons makes it the most paramagnetic of all the stable metal ions (Caravan, Ellison et al. 1999; Kubicek and Toth 2009).

### 1.6.1. Relaxivity of Gd(III) based contrast agents

Solomon, Bloembergen and others (Bloembergen, Purcell et al. 1948; Solomon and Bloembergen 1956; Bloembergen 1957; Bloembergen and Morgan 1961; Connick and Fiat 1966) provided detailed description of the relaxation of solvent nuclei around a paramagnetic center. As equation 1.6 below shows, the observed solvent relaxation rate,  $1/T_{i, obs}$ , is the summation of the diamagnetic term  $1/T_{i, d}$ , corresponding to the relaxation rate of the solvent nuclei without the paramagnetic solute, and a paramagnetic term  $1/T_{i, p}$  which is the relaxation rate augmentation caused by the paramagnetic substance:

$$\frac{1}{T_{i, obs}} = \frac{1}{T_{i, d}} + \frac{1}{T_{i, p}} \quad i = 1, 2 \quad (1.6)$$

This paramagnetic contribution is linearly proportional to the concentration of paramagnetic, [Gd(III)] :

$$\frac{1}{T_{i, obs}} = \frac{1}{T_{i, d}} + r_i [Gd(III)] \quad i = 1, 2 \quad (1.7)$$

where the parameter  $r_i$  is the longitudinal or transverse proton relaxivity ( $\text{mM}^{-1}\text{s}^{-1}$ ), defined as the efficiency of a paramagnetic substance to enhance the relaxation rate of water protons. As

described by Toth *et al.* (Toth, Helm *et al.* 2002), the two key contributors to overall paramagnetic relaxation rate enhancement and/or relaxivity are the inner sphere, due to the interaction between the Gd(III) electron spins and the water protons in the first coordination sphere, and the outer sphere contribution, arising from random translational diffusion of bulk solvent molecules. The total Gd(III) relaxivity is therefore expressed by the equations below.

$$r_i^p = r_i^{IS} + r_i^{OS} \quad , i = 1, 2 \quad (1.8)$$

$$\left( \frac{1}{T_{i,p}} \right) = \left( \frac{1}{T_{i,p}} \right)^{IS} + \left( \frac{1}{T_{i,p}} \right)^{OS} \quad (1.9)$$

where  $r_i^{IS} = 1/T_{i,p}^{IS}$  and  $r_i^{OS} = 1/T_{i,p}^{OS}$  and IS and OS stand for inner and outer sphere respectively. Efforts to improve the relaxivity of contrast agents generally involve increasing the inner sphere term since the outer sphere contribution can hardly be modified (Toth, Helm *et al.* 2002). The most important factors governing relaxivity that can be modified are the rotational correlation time, proton exchange and the hydration number,  $q$  (Toth, Helm *et al.* 2002; Kubicek and Toth 2009). However, most efforts to improve relaxivity involve increasing the rotational correlation time by binding Gd(III) to macromolecular compounds like proteins, polymers and lipids structures (Toth, Helm *et al.* 2002; Strijkers, Mulder *et al.* 2005; Hermann, Kotek *et al.* 2008; Hak, Sanders *et al.* 2009).

## 1.7. Microbubbles as MRI Contrast Agents

### 1.7.1. Microbubbles as Contrast Agents for $T_2^*$ -weighted MRI

Gas-filled microbubbles have been previously adapted mainly as  $T_2^*$  MR contrast agents. Studies have demonstrated that microbubbles (without attached MRI contrast media) can enhance  $T_2^*$  shortening (Alexander, McCreery et al. 1996; Wong, Huang et al. 2004; Cheung, Chow et al. 2009).  $T_2^*$  is the additional enhancement of the transverse magnetization signal above of baseline provided by  $T_2$  and stems from inhomogeneities in the magnetic field. Previous results have demonstrated this  $T_2^*$  shortening effect to be microbubble concentration-dependent. Alexander *et al.* (Alexander, McCreery et al. 1996) reported  $T_2^*$  shortening from 9 different gas-types and the potential to use the gas-volume dependent susceptibility effect as a pressure sensor for evaluating cardiovascular function. Cheung *et al.* (Cheung, Chow et al. 2009) reported  $T_2^*$  shortening at 7 T in the rat brain as a function of microbubble volume fraction for both sulfur hexafluoride and air-filled microbubbles. Wong *et al.* (Wong, Huang et al. 2004) used the  $T_2^*$  shortening induced by Optison (GE Healthcare) microbubbles for intravascular imaging of the rat liver with MRI. Subsequent studies by Chow *et al.* (Chow, Chan et al.) and Yang *et al.* (Yang, Li et al. 2009) demonstrated additional  $T_2^*$  shortening by loading iron oxide into the shell of polymeric microbubbles (Chow, Chan et al. ; Yang, Li et al. 2009). Finally, Liu *et al.* (Liu, Lammers et al. 2011) reported an additional  $T_2^*$  shortening after polymeric microbubble fragmentation with ultrasound. The additional shortening after microbubble fragmentation was attributed to greater interaction of water protons with iron oxide in the shell fragments.



### 1.7.2. Microbubbles as Contrast Agents for $T_1$ -weighted MRI

Currently, the adaptation of gas-filled microbubbles for  $T_1$ -weighted MRI has been limited. Ao *et al.* (Ao, Wang *et al.* 2010) loaded Gd(III) onto the shell of 1.5- $\mu\text{m}$  diameter polymeric microbubbles and demonstrated  $T_1$  shortening as a function Gd(III)-loaded microbubble concentration. Additionally, Liu *et al.* (Liu, Lammers *et al.* 2011) demonstrated that fragmentation of iron oxide loaded polymer-shelled microbubbles with ultrasound, not only shortened  $T_2^*$ , but also  $T_1$ . In another study, Liao *et al.* (Liao, Liu *et al.* 2012) developed albumin microbubbles containing a shell surface-conjugated with Gd(III)-DTPA to serve as dual mode contrast agents for ultrasound and MR imaging. However, the problem with previously designed MRI contrast media loaded microbubbles for FUS surgery is that shell material has mainly been polymer or protein based, which is not ideal for FUS surgery due to their stiffness and limited ultrasound compliance (Bloch, Wan *et al.* 2004; Ferrara, Pollard *et al.* 2007; Sirsi and Borden 2009). Adaptation of Gd(III)-loaded lipid-shelled microbubbles may potentially solve this problem.

## 1.8. Research Motivations and Specific Aims

Microbubbles are ubiquitous in nature and are important for many industrial and clinical applications. Over the past several decades, microbubbles have been developed for biomedical applications such as CEUS, targeted drug delivery and metabolic gas delivery. Due to their high echogenicity, acoustic response and biocompatibility, lipid-stabilized microbubbles are most commonly desired for these applications. More recently, these lipid-stabilized microbubbles have been adapted for non-invasive surgery with FUS. For instance, microbubbles have been demonstrated to induce the opening of the BBB *in vivo* in a non-invasive and transient manner,

and at a lower acoustic driving pressure than with convention high intensity FUS and inception microbubbles (Hynynen, McDannold et al. 2001; Choi, Pernot et al. 2007). However, the microbubbles used in previous studies tended to be highly polydisperse in size, which poses a problem since their behavior under ultrasound is very size dependent. Thus, in order to understand the dependence of microbubble size on FUS-induced BBB opening, a technique that efficiently selects different microbubbles classes is necessary. Additionally, since the fragmentation state of a microbubble during or after FUS varies depending on microbubble size and ultrasound driving parameters, a means to detect the final fragmentation state of a microbubble is desired. To address these issues this research project was proposed to accomplish to following aims:

**Specific Aims:**

1. *Develop a method to size select different microbubble size classes rapidly and efficiently with centrifugation.*
2. *Develop a means to conjugate an MRI contrast agent, Gd(III,) to the shell of the lipid microbubble architecture without sacrificing microbubble size monodispersity.*
3. *Characterize the MR behavior of the Gd(III)-bound size-selected lipid-shelled microbubbles before and after fragmentation with ultrasound.*

The following hypotheses were made regarding these specific aims:

1. Microbubbles are large enough such that Brownian forces are neglected during centrifugal size separation.

2. Hydrodynamic interaction of multiple microbubbles during centrifugal separation can be accounted for by a second order correlation of the effective fluid viscosity.
3. Lipid-shelled microbubbles are robust enough to withstand the centrifugation speeds (up to 300 times gravity) without significant degradation.
4. Individual Gd(III)-ligand molecules cannot bind to multiple binding sites at the same time.
5. Post-labeling of the Gd(III)-ligand chelate to the lipid groups on the microbubble is more efficient and more cost effective than pre-labeling.
6. The heating from Gd(III) ligand chelation reaction does not significantly degrade microbubble architecture provided it is below the phase transition temperature of the main lipid headgroup.
7. Fragmentation/destruction of Gd(III)-bound lipid-shelled microbubbles with a combination of ultrasound and heating above the phase transition temperature of the main lipid headgroup, will produce a suspension of lipid vesicles.

## Chapter 2: Microbubble Size Isolation by Differential Centrifugation

### 2.1. Introduction

Microbubbles are being employed for several biomedical applications, including contrast enhanced ultrasound (Feinstein 2004; Lindner 2004), drug and gene delivery (Ferrara, Pollard et al. 2007; Hernot and Klibanov 2008) and metabolic gas delivery (Burkard and Vanliew 1994; Kheir, Zurakowski et al. 2007). Microbubbles react strongly to ultrasonic pressure waves by virtue of their compressible gas cores, which resonate at the MHz-frequencies used by current clinical scanners. Oscillation of the gas core allows re-radiation (backscatter) of ultrasound energy to the transducer at harmonic frequencies and nonlinear modes, thus providing exquisite sensitivity in detection with current contrast-enhanced pulse sequences and signal processing algorithms. Additionally, microbubbles may cavitate stably or inertially to facilitate drug release (Borden, Kruse et al. 2005; Lum, Borden et al. 2006) and extravascular delivery (Choi, Pernot et al. 2007; Stieger, Caskey et al. 2007) within the transducer focus.

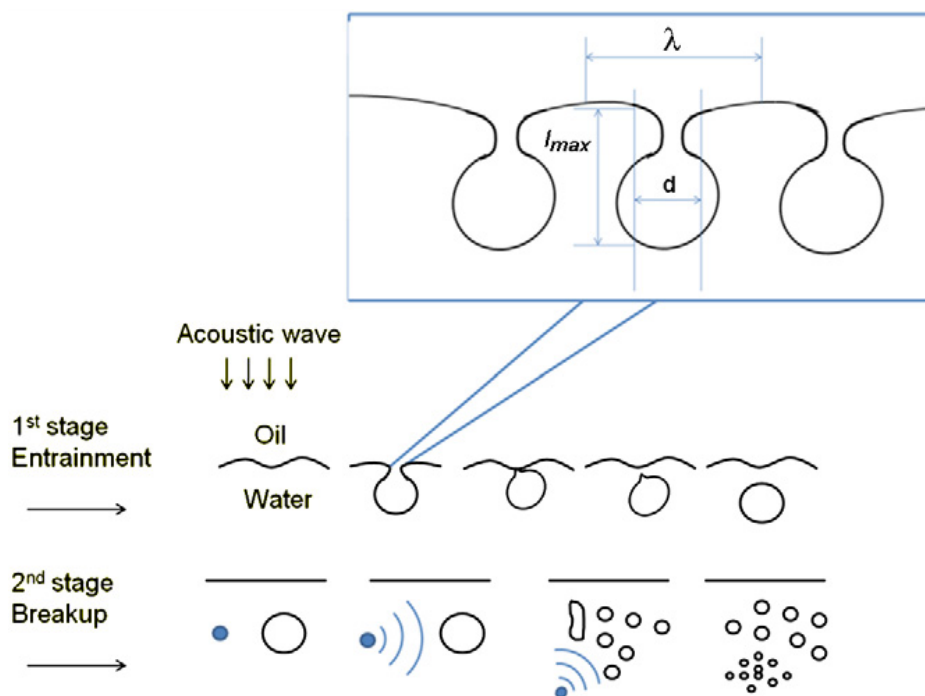
Current commercially available microbubble formulations are polydisperse in size. In most cases, the size distribution is broad over a range of submicron to tens of microns in diameter. This is problematic because microbubble behavior depends very strongly on size. For example, increasing the microbubble diameter from 1 to 5  $\mu\text{m}$  will change the resonance frequency of an unencapsulated microbubble from 4.7 to 0.72 MHz (Wu and Nyborg 2008). Microbubble size also affects the biodistribution and pharmacodynamics after intravenous injection, the bioeffects during ultrasound insonification, the gas release profile, and other related behaviors. Clearly, microbubbles of a specific size with low polydispersity are desired for advanced biomedical applications (Talu, Hettiarachchi et al. 2007).

Efforts to engineer monodisperse microbubble suspensions have mainly focused on microfluidic technologies. These techniques include flow focusing (Talu, Lozano et al. 2006; Talu, Hettiarachchi et al. 2007; Talu, Hettiarachchi et al. 2008), T-junctions (Xu, Li et al. 2006) and electrohydrodynamic atomization (Farook, Stride et al. 2007; Farook, Zhang et al. 2007). While these techniques provide very low polydispersity, they are rather slow at generating microbubbles (Pancholi, Farook et al. 2008). Using flow focusing, for example, requires several hours to produce microbubbles at sufficient numbers for even a single small-animal trial ( $\sim 0.1 \text{ mL} \times 10^9 \text{ mL}^{-1}$ ). Additionally, dust particles can plug micro-channels, thus requiring fabrication and calibration of a new device. While engineering breakthroughs may eventually allow efficient and robust generation of monodisperse microbubbles via microfluidic strategies, these techniques currently remain untenable for biomedical studies.

Mechanical agitation has been the main method to create encapsulated microbubbles for biomedical applications, since their inception by Feinstein et al. (Feinstein, Shah et al. 1984). Mechanical agitation is a common emulsification procedure in which a hydrophobic phase (i.e., gas) is dispersed within an aqueous surfactant solution by disruption of the interface. Acoustic emulsification (sonication), for example, generates large quantities of microbubbles ( $100 \text{ mL} \times 10^{10} \text{ mL}^{-1}$ ) rapidly and reproducibly within just a few seconds. Shaking a serum vial with a device similar to a dental amalgamator produces a sufficient dose of microbubbles ( $2 \text{ mL} \times 10^{10} \text{ mL}^{-1}$ ) for a single patient study, at the bedside in under a minute (Unger, Fritz et al. 1999). While mechanical agitation is highly efficient at generating microbubbles, the size distributions tend to be highly polydisperse and thus are not optimal for biomedical applications.

The origins of polydispersity in acoustically generated emulsions were elucidated three decades ago by Li and Fogler (Li and Fogler 1978; Li and Fogler 1978). Emulsification was

reported to occur in two stages. Instability at the water surface results in entrainment of drops (or bubbles) into the aqueous medium, and subsequent cavitation in the medium results in droplet breakup to a critical size (Figure 2.1).



**Figure 2.1. Cartoon showing origins of polydispersity during acoustic emulsification. Initial bubble entrainment occurs as a capillary instability. Inset shows relevant length scales. Subsequent cavitation in the suspension (shown as filled circle with propagating waves) induces breakup of the larger bubbles to a critical diameter, where surface forces and inertial forces balance. Figure adapted from Li and Fogler (Li and Fogler 1978; Li and Fogler 1978).**

The first stage, entrainment, occurs as the unstable growth and eventual eruption of interfacial capillary waves produced by sonication. The second stage of acoustic emulsification involves the continual cavitation-induced breakdown of larger particles as a function of sonication time until a stable size is reached. The breakdown mechanism depends on the type of deformation

and flow pattern around the droplet. The stable size results when surface tension forces balance the inertial forces on the droplet. Initial and final droplet size is difficult to predict *a priori*. The analysis provided by Li and Fogler for liquid droplets points to the origin of polydispersity as a consequence of multiple mechanisms acting simultaneously on the multi-body system. Given that emulsion polydispersity is inherent in mechanical agitation processes, it is desirable to find a means of separating subpopulations of the particles based on size. This will allow improved microbubble formulations for advanced biomedical applications.

Previous reports have described the use of flotation to isolate subpopulations from polydisperse microbubble suspensions. In principle, larger microbubbles are more buoyant and rise faster, thus allowing separation based on different migration rates in a gravitational field. Kvale et al. described a model for the size fractionation of air-filled microbubbles by simple flotation (Kvale, Jakobsen et al. 1996). Microbubbles were injected at the bottom of a stagnant water column and allowed to rise under normal gravity. The model predicted the size distribution of microbubbles at certain distances from the bottom of the column as a function of time. The form of the model was a second-order PDE that accounted for the convective motion of the bulk dispersed phase (liquid moved down the column as microbubbles moved up) as well as the Brownian (thermal) diffusive motion of the particles. The crowding effect of the microspheres was accounted for by using a modified version of Einstein's derivation for the effective viscosity in a dilute suspension (Batchelor and Green 1972).

Wheatley *et al.* reported the isolation of submicron bubbles using differential centrifugation (Wheatley, Forsberg et al. 2006). Isolation was accomplished by flotation at normal gravity, or centrifugation at a relative centrifugal force (RCF) of 16 or 45 for pre-determined time intervals. This method allowed isolation of the submicron bubble fractions.

The use of centrifugation reduced the flotation time, but led to destabilization of the surfactant-stabilized microbubbles during subsequent insonification. Microbubbles centrifuged at 45 RCF for 1 minute were not stable, whereas those spun at 16 RCF for the same time were relatively stable. Destabilization was attributed to the extra hydrostatic pressure exerted on the microbubbles, which increased towards the bottom of the column and in proportion to centrifugation speed. Flotation at normal gravity was more time consuming, but less detrimental to microbubble stability.

In contrast to surfactant-coated microbubbles, lipid-coated microbubbles have been shown to be stable after centrifugation up to several hundred RCF (Takalkar, Klibanov et al. 2004; Zhao, Borden et al. 2004). The lipid shell is highly viscous (Kim, Costello et al. 2003) and relatively impermeable to gases (Borden and Longo 2004). We therefore sought to further develop the differential centrifugation method of Wheatley et al. (Wheatley, Forsberg et al. 2006), but as a rapid and facile means to isolate sub-populations of lipid-coated microbubbles. Below, we report on the experimental characterization of the initial polydisperse suspension, the development of a method to isolate size fractions of interest for biomedical applications, and characterization of the long-term stability of the isolated fractions.

## **2.2. Methods and Materials**

### **2.2.1. Materials**

All solutions were prepared using filtered, 18M $\Omega$  deionized water (Direct-Q, Millipore, Billerica, MA). All glassware was cleaned with 70 vol% ethyl alcohol solution (Sigma-Aldrich; St. Louis, MO) and rinsed with deionized water. The gas used to form microbubbles was perfluorobutane



(PFB) at 99 wt% purity obtained from FluoroMed (Round Rock, TX). 1,2-distearoyl-*sn*-glycero-3-phosphocholine (DSPC) was purchased from Avanti Polar Lipids (Alabaster, AL) and dissolved in chloroform (Sigma-Aldrich) for storage. Polyoxyethylene-40 stearate (PEG40S) was obtained from Sigma-Aldrich and dissolved in deionized water. The fluorophore probe 3,3'-dioctadecyloxycarbocyanine perchlorate (DiO) solution (Invitrogen; Eugene, OR) was used to label the microbubbles for part of the experiments.

### **2.2.2. Microbubble Generation**

Microbubbles were coated with DSPC and PEG40S at molar ratio of 9:1. The indicated amount of DSPC was transferred to a separate vial, and the chloroform was evaporated with a steady nitrogen stream during vortexing for about ten minutes followed by several hours under house vacuum. 0.01 M phosphate buffered saline (PBS) solution (Sigma-Aldrich) was filtered using 0.2- $\mu$ m pore size polycarbonate filters (VWR, West Chester, PA). The dried lipid film was then hydrated with filtered PBS and mixed with PEG40S (25 mg/mL in filtered PBS) to a final lipid/surfactant concentration of 1.0 mg/mL. The lipid mixture was first sonicated with a 20-kHz probe (model 250A, Branson Ultrasonics; Danbury, CT) at low power (power setting dialed to 3/10; 3 Watts) in order to heat the pre-microbubble suspension above the main phase transition temperature of the phospholipid ( $\sim 55$  °C for DSPC) and further disperse the lipid aggregates into small, unilamellar liposomes (Kim and Franses 2005). PFB gas was introduced by flowing it over the surface of the lipid suspension. Subsequently, higher power sonication (power setting dialed to 10/10; 33 Watts) was applied to the suspension for about 10 seconds at the gas-liquid interface to generate microbubbles. For flow cytometry and fluorescence microscopy

experiments, DiO solution (1 mM) was added prior to high-power sonication at an amount of 1  $\mu$ L DiO solution per mL of lipid mixture.

### **2.2.3. Microbubble Washing & Lipid Recycling**

The microbubble suspension was collected into 30-mL syringes (Tyco Healthcare, Mansfield, MA), which were used as the flotation columns. Washing and size fractionation by centrifugation was performed with a bucket-rotor centrifuge (model 5804, Eppendorf, Westbury, NY), which had a radius of approximately 16 cm from the center to the syringe tip and operated between 10 and 4500 RPM. Centrifugation (10 minutes, 300 RCF) was performed to collect all microbubbles from the suspension into a cake resting against the syringe plunger. The remaining suspension (infranatant), which contained residual lipids and vesicles that did not form part of the microbubble shells, was recycled to produce the next batch of microbubbles. All resulting cakes were combined and re-suspended in PBS to improve total yield.

### **2.2.4. Size and Concentration Measurements**

Microbubble size distribution was determined by laser light obscuration and scattering (Accusizer 780A, NICOMP Particle Sizing Systems, Santa Barbara, CA). 2- $\mu$ L samples of each microbubble suspension were diluted into a 30-mL flask under mild mixing during measurement. Size distribution was also determined using the electrozone sensing method (Coulter Multisizer III, Beckman Coulter, Opa Locka, Fl). A 4- $\mu$ L sample of microbubble suspension was diluted into a 60-mL flask and stirred continuously to prevent flotation-induced error. A 30- $\mu$ m aperture

(size range of 0.6-18  $\mu\text{m}$ ) was used for the measurements. All samples were measured at least three times by either instrument and analyzed for both number- and volume-weighted size distributions.

### **2.2.5. Optical Microscopy**

Direct visual confirmation of microbubble size was performed 48 hours after the samples were prepared using an Olympus 1X71 inverted microscope (Olympus; Center Valley, PA). The microbubble samples were taken directly from the serum vials and imaged at room temperature. Images were captured in both bright-field and epi-fluorescence modes using a high-resolution digital camera (Orca HR, Hamamatsu, Japan) and processed with Simple PCI software (C-Imaging, Cranberry Township, PA). A 40X objective was used to capture the images of size-isolated microbubbles of 4-5  $\mu\text{m}$  diameter, while a 100X oil-immersion objective was used for polydispersed microbubbles and size-isolated microbubbles of 1-2  $\mu\text{m}$  diameter. Subsequent image analysis was done using ImageJ 1.4g (<http://rsb.nih.gov/ij/>).

### **2.2.6. Flow Cytometry**

A FACScan Cell Analyzer (Becton-Dickinson, Franklin Lakes, NJ) was used to characterize microbubble fluorescence intensity (FL) and light scattering profiles (FSC and SSC). Voltage and gain settings for FSC, SSC and FL were adjusted to delineate the microbubble populations from instrument and sample noise. 10  $\mu\text{L}$  samples were diluted with 3 mL deionized water prior

to each measurement. Subsequent data analysis was done using CellQuest Pro (Becton-Dickinson, Franklin Lakes, NJ).

### 2.3. Size Isolation

Differential centrifugation was used to isolate size-selected microbubbles based on their migration in a centrifugal field (Fig. 2.2). The initial microbubble size distribution and concentration was measured and imported into a spreadsheet (Excel, Microsoft, Bellevue, WA) in order to determine the number density for each size channel and the total gas volume fraction. The spreadsheet was used to calculate the relative centrifugal force (RCF) needed for a microbubble size class to rise through the column of length L for a fixed centrifugation time. Following Kvale (Kvale, Jakobsen et al. 1996), *Stokes'* equation for the rise velocity of a buoyant particle relative to the bulk fluid under creeping flow conditions was used as follows:

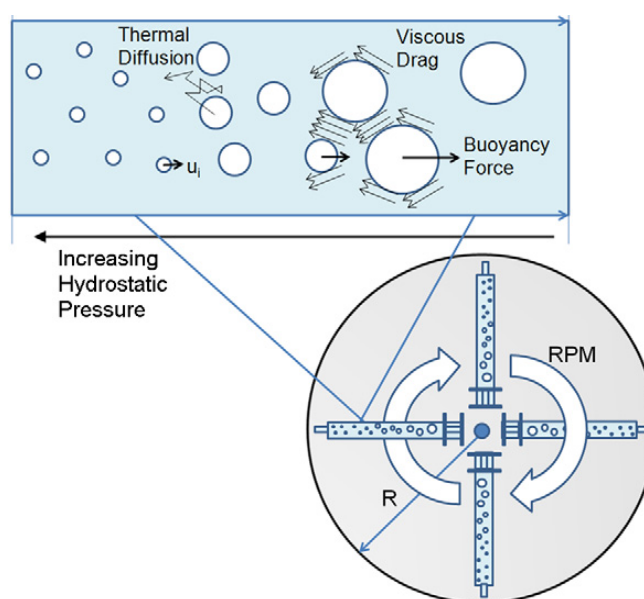
$$u_i = \frac{2(\rho_2 - \rho_{1i})}{9\eta_2} r_i^2 g.RCF, \quad (2.1)$$

where subscript  $i$  refers to the particle size class,  $r_i$  is the particle radius and  $g$  is the gravitational force. The effective viscosity,  $\eta_2^*$ , of the microbubble suspension was calculated using Batchelor and Greene's (Batchelor and Green 1972) correlation for the modified fluid viscosity:

$$\frac{\eta_2^*}{\eta_2} = 1 + 2.5\Phi + 7.6\Phi^2, \quad (2.2)$$

$$\Phi = \sum_{i=1}^{N_d} \Phi_i, \quad (2.3)$$

where  $\Phi$  is total the microbubble volume fraction for  $N_d$  size classes. Equations 11-13 were used to calculate the strength of the centrifugal field (in RCF) for a given initial size distribution, time period and syringe column length. Volume fraction was assumed to be constant over the entire column, and acceleration/deceleration effects were neglected.



**Figure 2.2.** Schematic of differential centrifugation for the size isolation of microbubbles.

Following production, microbubbles were collected into 30-mL syringes ( $L = 8.2$  cm) and washed, as above. Production-washing was repeated 3-5 times, each time saving the microbubble cake and recycling the lipid infranatant. The cakes were combined and re-dispersed into 30 mL of filtered PBS. We noted that in order to ensure a high yield, the concentration of microbubbles after this step should be at least  $\sim 1$  volume %. The following protocol was then used to isolate microbubbles of 1-2  $\mu\text{m}$  and 4-5  $\mu\text{m}$  diameter. At least three separate experimental runs were performed for each isolation, and size distributions were measured at

least three times each. A two-tailed parametric unpaired student t-test was used to determine to significance between the polydispersity indices (PI) of 1-2  $\mu\text{m}$  and 4-5  $\mu\text{m}$  samples versus that of freshly made samples.

### **2.3.1. Isolation of 4-5 $\mu\text{m}$ Diameter Microbubbles**

Before beginning the isolation process, care was taken to remove large, visible bubbles that may have formed during production or subsequent handling. Microbubbles of greater than 10- $\mu\text{m}$  diameter were removed by performing one centrifugation cycle at 30 RCF for 1 min. The infranatant consisting of less than 10- $\mu\text{m}$  diameter microbubbles was saved and re-dispersed in 30 mL PBS, while the cake was discarded. Next, microbubbles of greater than 6- $\mu\text{m}$  diameter were removed by performing one centrifugation cycle at 70 RCF for 1 min. The infranatant consisting of less than 6- $\mu\text{m}$  diameter microbubbles was saved and re-dispersed to 30 mL PBS; the cake was discarded. Finally, microbubbles of less than 4- $\mu\text{m}$  diameter were removed by centrifuging at 160 RCF for 1 min. This was repeated about 5-10 times, while each time the infranatant was discarded and the cake was re-dispersed in filtered PBS. Alternatively, 12-mL syringes ( $L = 6.3 \text{ cm}$ ) were employed and centrifuged at 120 RCF for 1 min to improve yield. These cycles were repeated until the infranatant was no longer turbid, indicating complete removal of microbubbles less than 4  $\mu\text{m}$ . The final cake was concentrated to a 1-mL volume of 20 vol% glycerol solution in PBS and stored in a 2-mL scintillation vial with PFB headspace.

### **2.3.2. Isolation of 1-2 $\mu\text{m}$ Diameter Microbubbles**

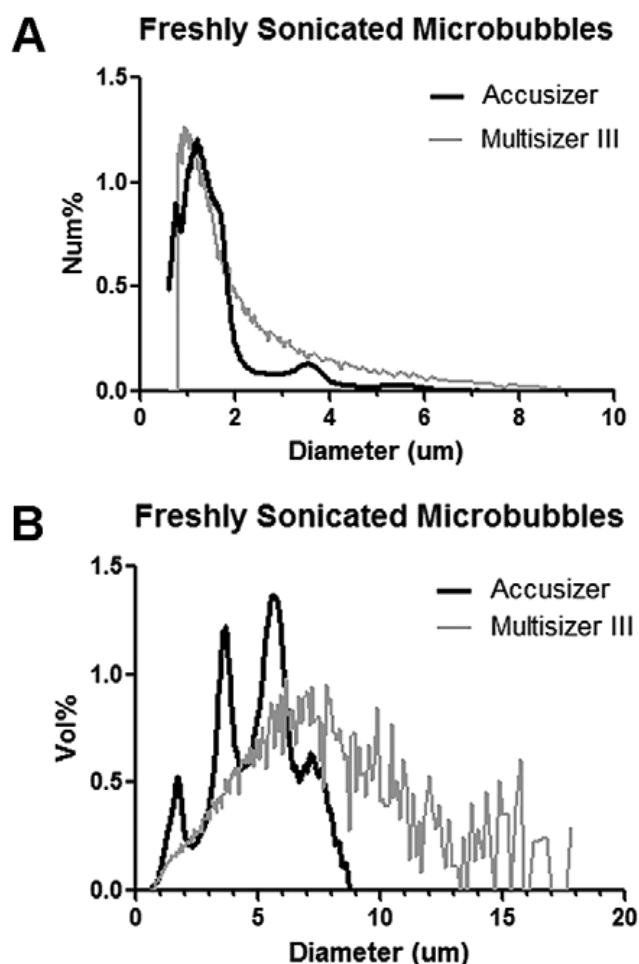
The infranatant collected from the 4-5 micron isolation was centrifuged at 270 RCF for 1 min for one cycle in order to remove microbubbles of approximately 3- $\mu\text{m}$  diameter and above by collecting them into the cake. The infranatant consisted mostly of microbubbles 1-2  $\mu\text{m}$  diameter. The target microbubbles were collected into a concentrated cake by centrifuging at 300 RCF for 10 min. The final cake was re-dispersed to a 1-mL volume of 20 vol% glycerol solution in PBS and stored in a 2-mL serum vial with PFB headspace.

## **2.4. Results**

### **2.4.1. Polydispersity of Freshly Sonicated Microbubbles**

Preparation of microbubbles by sonication of a 50 mL lipid mixture resulted in a polydisperse suspension of approximately  $10^9$  to  $10^{10}$  particles  $\text{mL}^{-1}$ . Particle sizing with the Accusizer and Multisizer showed a distribution ranging from the lower limit of resolution,  $\sim 0.5 \mu\text{m}$ , to greater than  $15 \mu\text{m}$  diameter (Fig. 2.3). A significant portion of the freshly generated suspension contained submicron microbubbles, as previously reported (Borden, Martinez et al. 2006). Submicron microbubbles also have been observed by static light scattering (Wheatley, Forsberg et al. 2006) and freeze-fracture transmission electron microscopy (Brancewicz, Rasmussen et al. 2006). For larger microbubbles, the number-weighted distribution tailed off near 6-8  $\mu\text{m}$  diameter (Fig. 2.3A). The volume-weighted distribution, however, showed a significant population out to greater than  $10 \mu\text{m}$  diameter (Fig. 2.3B). Microbubbles with larger diameters tended to skew the volume-weighted size distribution. Median volume-weighted diameters

therefore were chosen in order to judge the samples during size isolation, since this gives a more rigorous indication of the central tendency than arithmetic mean in a skewed distribution.



**Figure 2.3. Size distributions for freshly sonicated microbubbles.**

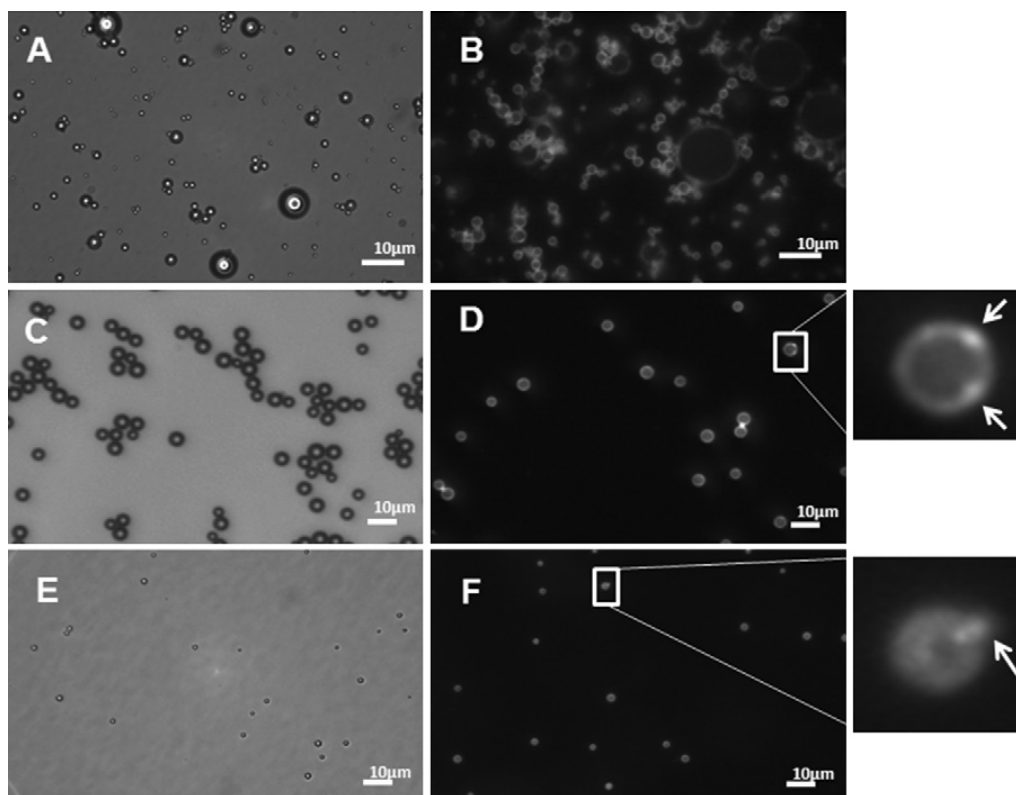
Interestingly, the Accusizer consistently measured distinct peaks centered on approximately 1-2, 4-5, 7-8 and 9-11  $\mu\text{m}$  diameter for each batch of lipid-coated microbubbles. These peaks were evident on the volume-weighted distribution, but they also could be discerned from the number-weighted distribution. Similar results were reported previously (Borden, Zhang et al. 2008). In the laboratory, we have observed these peaks for a variety of gas and lipid combinations (data not shown). We also measured size distribution using a Multisizer III. While the Accusizer



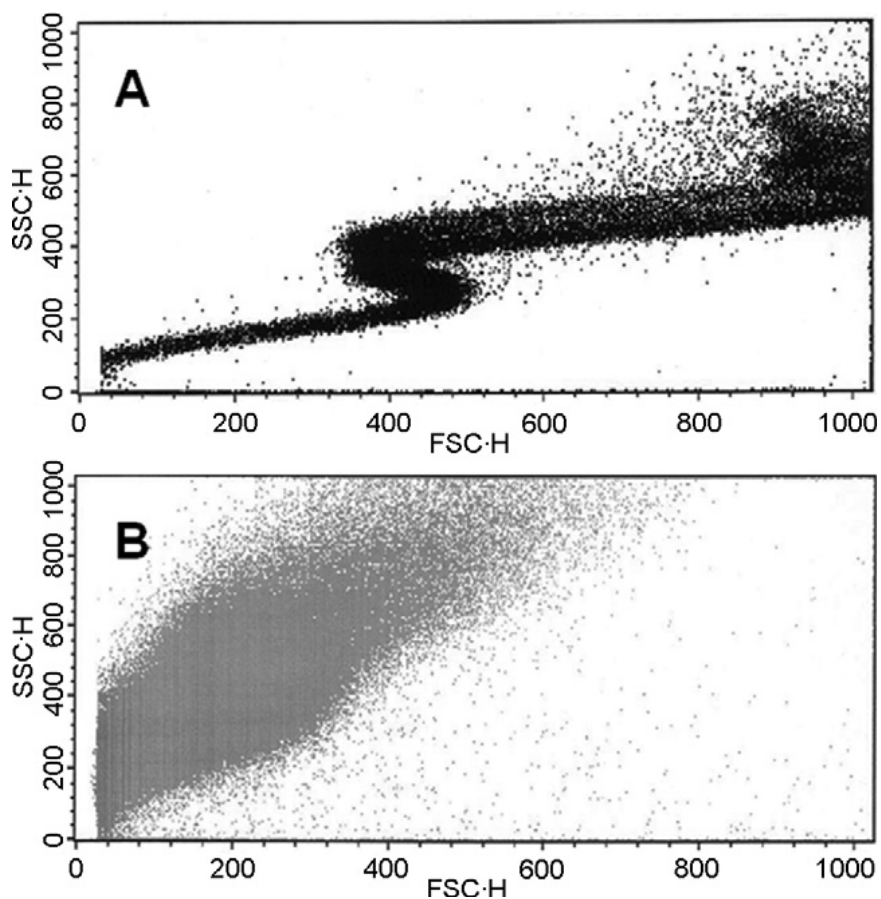
measures size based on light obscuration and scattering, the Multisizer utilizes electrical impedance sensing of the volume of electrolyte displaced by the microbubble as it passes through an orifice. Interestingly, the multimodal distribution was not observed on the Multisizer, which gave a broad distribution with a single peak located at  $\sim 1 \mu\text{m}$  for the number-weighted distribution and  $\sim 8 \mu\text{m}$  for the volume-weighted distribution. From this data, it was unclear whether the multimodal distribution was real, and could not be resolved by the Multisizer, or if it was an artifact of the Accusizer. We therefore sought to better characterize the microbubble distribution.

Microscopy allowed direct visual inspection of individual microbubbles from the suspension. Bright-field and epi-fluorescence microscopy images are shown in Figure 2.4. In fluorescence mode, microbubbles appeared as bright rings with dark centers, clearly showing uptake of DiO into the shell. In bright-field mode, microbubbles appeared as dark spheres with bright centers. Diffraction rings were particularly prevalent for the smaller microbubbles. This confirmed the predominance of gas-filled microbubbles in the suspension. Analysis of the bright-field images using ImageJ indicated that the distribution of the freshly generated microbubbles was multimodal, with a mean diameter of  $4.0 \pm 3.0 \mu\text{m}$  for the image shown in Figure 2.4A.

Flow cytometry was used to further characterize the polydisperse microbubbles (Fig. 2.5). Forward- (FSC) and side- (SSC) light scattering measurements were taken. Interestingly, a serpentine shape was observed on the dot plot of FSC versus SSC for the polydisperse microbubble suspensions, as shown in Figure 2.5A. The serpentine shape appeared to correlate with the distinct peaks found by the Accusizer, lending more support to the validity of a multimodal distribution in the freshly generated microbubble suspension.



**Figure 2.4. Microscopy images of initial polydisperse and final size-isolated microbubbles. Shown are bright-field images (left) and fluorescence images (right) of the membrane probe DiO. Arrows point to microstructural features of high probe density. The initial, polydisperse microbubble suspensions (A, B) are shown for comparison to the isolated 4–5 μm diameter microbubbles (C, D) and 1–2 μm diameter microbubbles (E, F).**



**Figure 2.5. Dot plots of side scatter (SSC) versus forward scatter (FSC) for polydisperse and size-isolated microbubbles as determined by flow cytometry. (A) shows the serpentine trend for the initial polydisperse microbubble suspension (cytometer settings: detector P1, voltage E01, amp 1.98; detector P2, voltage 287, amp 1.49). (B) shows no serpentine trend for the isolated 4–5  $\mu\text{m}$  microbubbles (cytometer settings: detector P1, voltage E01, amp 2.30; detector P2, voltage 173, amp 2.88).**

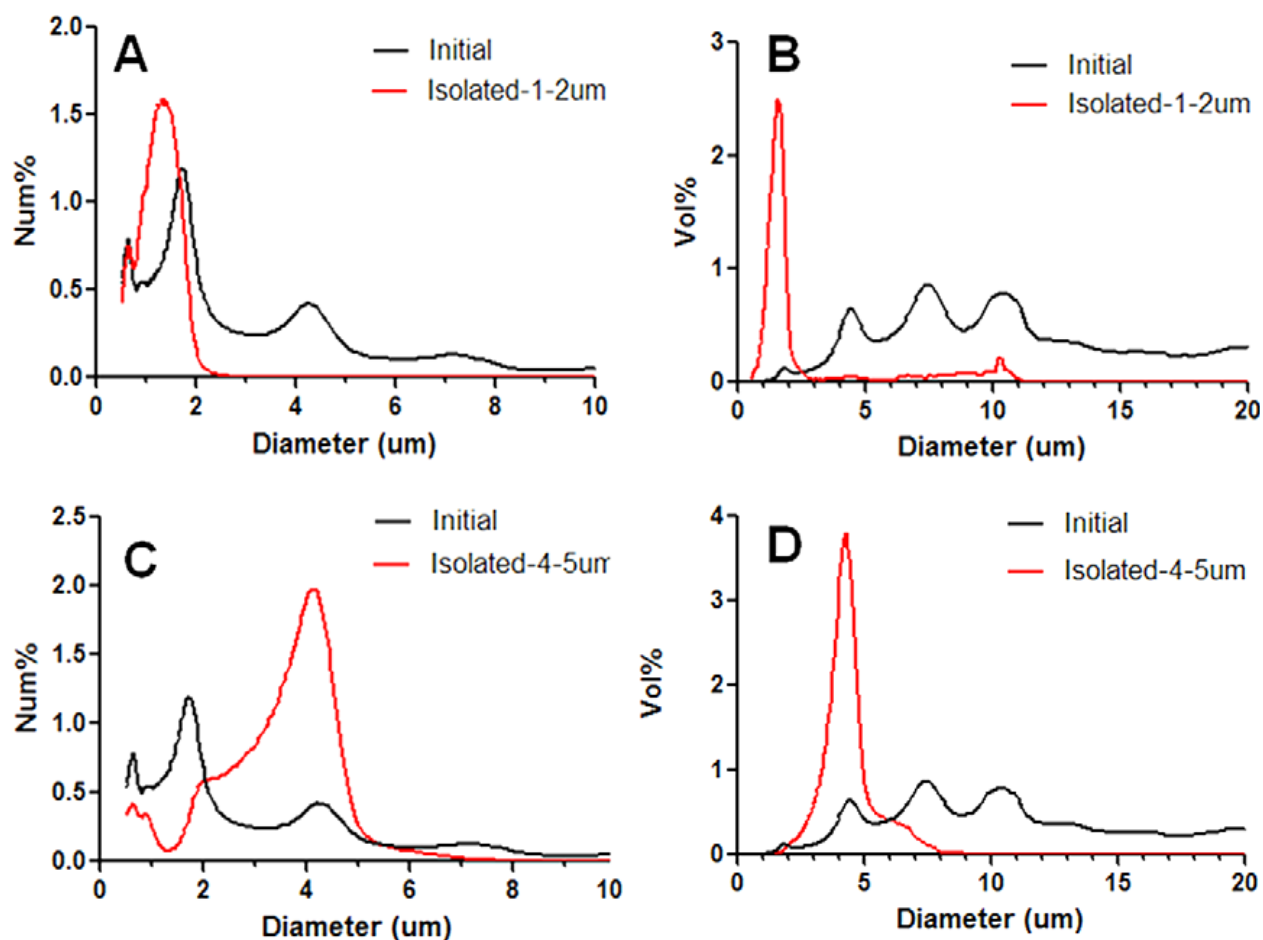
The origins of polydispersity in the freshly generated suspension of lipid-coated microbubbles observed here may be explained by the multiple interacting mechanisms occurring during entrainment and cavitation-induced disintegration, as described above. The fact that the microbubbles themselves may be oscillating in the acoustic field and may act as cavitation nuclei

adds further complexity to analysis. Additionally, the dynamics of lipid adsorption and spreading and monolayer shell formation are expected to play a role in determining the apparent surface tension and, for the lipids used here, may be expected to add additional surface viscosity and elasticity terms. Marangoni effects are expected to play a central role in both entrainment and cavitation-induced breakdown (Edwards, Brenner et al. 1991). While polydispersity may be unavoidable, the ability of mechanical agitation to rapidly generate large numbers of microbubbles brings this technique to the forefront of current microbubble creation methods. Given the excellent stability of lipid-coated microbubbles and the apparent presence of distinct peaks in the multimodal distribution, size isolation by differential centrifugation appeared to be a feasible approach. In what follows, we describe experiments set at isolating narrow distributions and characterizing their size distribution and long-term stability.

#### **2.4.2. Size Isolation of Microbubbles**

Experiments were performed to isolate relevant subpopulations of the multimodal distribution and reduce polydispersity. Submicron microbubbles were found to be relatively unstable and therefore were not isolated. Instead, microbubbles in the 1-2  $\mu\text{m}$  and 4-5  $\mu\text{m}$  diameter ranges were isolated. These ranges are interesting for biomedical applications. While both sizes are comparable to that of an erythrocyte, they may possibly yield different biodistributions, resonance frequencies, and acoustically induced bioeffects. In general, the 1-2  $\mu\text{m}$  microbubbles were approximately 100-fold more abundant than the 4-5  $\mu\text{m}$  microbubbles in the initial dispersion.

Microbubbles in the larger diameter range (4-5  $\mu\text{m}$ ) were isolated first, while the smaller microbubbles were saved for the subsequent isolation of the 1-2  $\mu\text{m}$  fraction. After repeated centrifugation and re-concentration according to the simple model, microbubbles with diameters of 4-5  $\mu\text{m}$  were successfully isolated from the initial polydisperse suspension (Fig. 2.6). Multiple centrifugation steps were needed to expel smaller microbubbles ( $< 4 \mu\text{m}$ ), which were more abundant in the initial suspension. The final 4-5  $\mu\text{m}$  microbubble suspension typically had a total volume of 1 mL with concentration in the order of  $10^8$  to  $10^9 \text{ mL}^{-1}$ , as determined by the Accusizer. Table 2.1 summarizes both averaged number-weighted and volume-weighted mean and median values for each size fraction.



**Figure 2.6. Size distributions for initial polydisperse and final size-isolated microbubbles. Shown are isolated subpopulations at the 1–2 μm (A, B) and 4–5 μm (C, D) diameter size ranges. Comparison of number-weighted (A, C) and volume-weighted (B, D) size distributions allows inspection of polydispersity. Results are summarized in Table 2.1.**

**Table 2.1: 1-2 micron and 4-5 micron microbubble size distribution data.**

Sample	Total Concentration (#/mL)	Number-Weighted Diameter ( $\mu\text{m}$ )		Volume-Weighted Diameter ( $\mu\text{m}$ )		PI
		(Mean $\pm$ SD)	(Median $\pm$ SD)	(Mean $\pm$ SD)	(Median $\pm$ SD)	
Initial	1.35E+10	2.1 $\pm$ 0.5	1.4 $\pm$ 0.6	19.4 $\pm$ 17.1	19.0 $\pm$ 17.3	10.5 $\pm$ 10.6
1-2 $\mu\text{m}$ Isolated	2.8E+09	1.3 $\pm$ 0.2	1.3 $\pm$ 0.2	1.8 $\pm$ 0.1	1.5 $\pm$ 0.2	1.5 $\pm$ 0.2
4-5 $\mu\text{m}$ Isolated	1.01E+08	2.8 $\pm$ 0.2	3.0 $\pm$ 0.2	4.2 $\pm$ 0.1	4.1 $\pm$ 0.1	1.5 $\pm$ 0.1

Microbubbles of 1-2  $\mu\text{m}$  diameter were isolated in fewer steps than for the 4-5  $\mu\text{m}$  microbubbles. For instance, separation of microbubbles less than 2  $\mu\text{m}$  diameter in the infranatant was typically completed by a single centrifugation step. However, the final step of concentrating microbubbles greater than 1- $\mu\text{m}$  diameter required substantially higher centrifugal force than for the 4-5  $\mu\text{m}$  microbubbles, which is consistent with their lower buoyancy. The final 1-2  $\mu\text{m}$  microbubble suspension typically had a total volume of 1 mL with concentration on the order of  $10^9$  to  $10^{10}$   $\text{mL}^{-1}$ , as determined by the Accusizer.

In order to assess size uniformity, we defined the polydispersity index (PI) as the volume-weighted mean diameter divided by the number-weighted mean diameter. Table 2.1 gives the average PI value for the freshly generated microbubbles and the size-isolated microbubbles. The initial suspension was highly polydisperse, with PI values as high as 18 but no lower than 6. The PI for the 4-5  $\mu\text{m}$  fraction was only  $1.5 \pm 0.1$  ( $p < 0.05$ ), while that of the 1-2  $\mu\text{m}$  fraction was only  $1.5 \pm 0.2$  ( $p < 0.05$ ).

Bright-field and epi-fluorescence microscopy images provided direct visual confirmation for the narrow size distribution of size-isolated microbubbles (Fig. 2.4). Analysis of the bright-field

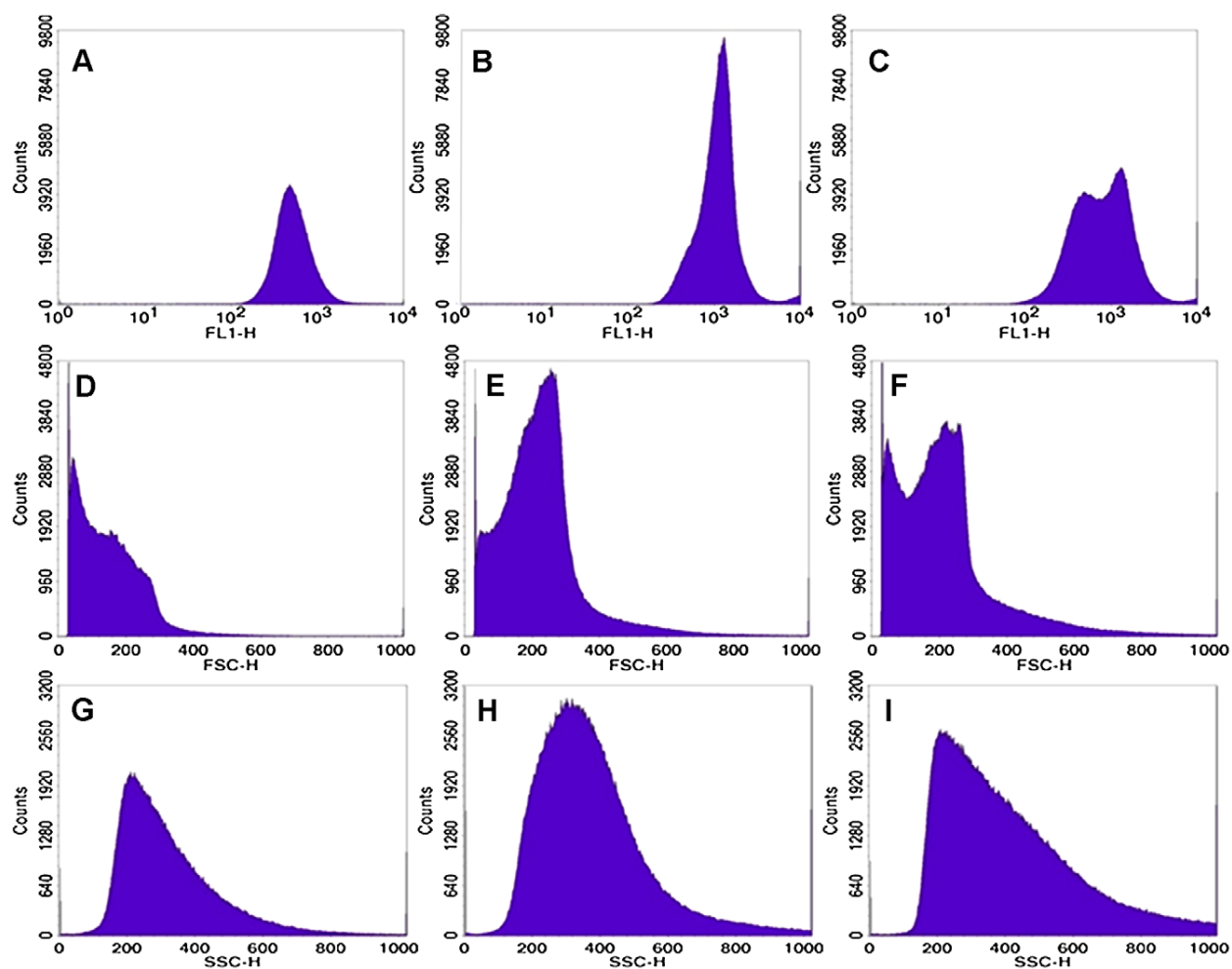
images using ImageJ gave mean diameters of  $4.6 \pm 0.3 \mu\text{m}$  and  $1.8 \pm 0.3 \mu\text{m}$  for microbubbles seen in Figs. 2.4C and 2.4E, respectively. These results are in agreement with the size distributions determined by the Accusizer and Multisizer. Fluorescence images also showed microstructural features within the lipid shell. For example, brighter spots indicating non-uniform DiO distribution were often observed (see arrows in Figures 2.4D and 2.4F). Previous results have shown heterogeneous membrane probe distribution that indicated lateral phase separation (Borden, Pu et al. 2004; Borden, Martinez et al. 2006) and collapse in the lipid shell (Borden, Kruse et al. 2005; Pu, Borden et al. 2006).

Flow cytometry was performed to characterize the size-isolated fractions (Fig. 2.7). Fluorescence intensity (FL), FSC and SSC measurements were all taken under the same cytometer settings. The serpentine shape was not observed for the size-isolated suspensions, as it was for the polydisperse case. Instead, the data points were found to be clustered in one region of the dot plot. The lack of the serpentine shape in the size-isolated samples indicated that they were indeed subpopulations of the initial multimodal sample. Table 2.2 lists the median values of three cytometry tests for each microbubble sample. Comparison of the FSC and SSC results for individual, size-isolated fractions and their mixture supported the existence of two distinct microbubble subpopulations. Monomodal distributions were observed for the individual size-isolated suspensions, with a lower median value corresponding to the 1-2  $\mu\text{m}$  microbubbles. When the size-isolated microbubbles were subsequently mixed together, a bimodal distribution appeared with two distinct peaks that agreed with the respective median values for the individual suspensions.

As expected, a single peak was observed on the FL histogram for the size isolated microbubbles (Fig. 2.7). The median FL value for 1-2  $\mu\text{m}$  microbubbles was lower than that for the 4-5  $\mu\text{m}$



microbubbles. Upon mixing, a bimodal distribution was observed with peak median FL values corresponding to those of the individual suspensions. Assuming that each microbubble is a perfect sphere and the fraction of fluorophores in the lipid shell is the same for all microbubbles, regardless of size, the number of fluorophores per microbubble should be directly proportional to the surface area, or the square of the diameter. This was confirmed when comparing the averaged FL values versus microbubble squared diameter. The fluorescence intensity value for the mixture of 1-2  $\mu\text{m}$  and 4-5  $\mu\text{m}$  microbubble samples ( $775 \pm 18$ ) agreed with the average between the FL values measured for each individual monodisperse suspension ( $510 \pm 16$  and  $1110 \pm 35$ ).



**Figure 2.7. Fluorescence intensity and light scattering profiles for microbubble suspensions after size isolation as determined by flow cytometry. Three different tests (fluorescence intensity FL, forward- (FSC) and side- (SSC) light scattering versus particle count) were performed for the same sample as represented by each column of plots. Column 1 (A, D, G) and Column 2 (B, E, H) samples had median volume-weighted diameters of 1.8 and 4.6  $\mu\text{m}$ , respectively. Column 3 (C, F, I) was a mixture of these two suspensions. Results are summarized in Table 2.2.**

**Table 2.2: Summary of flow cytometry measurements.**

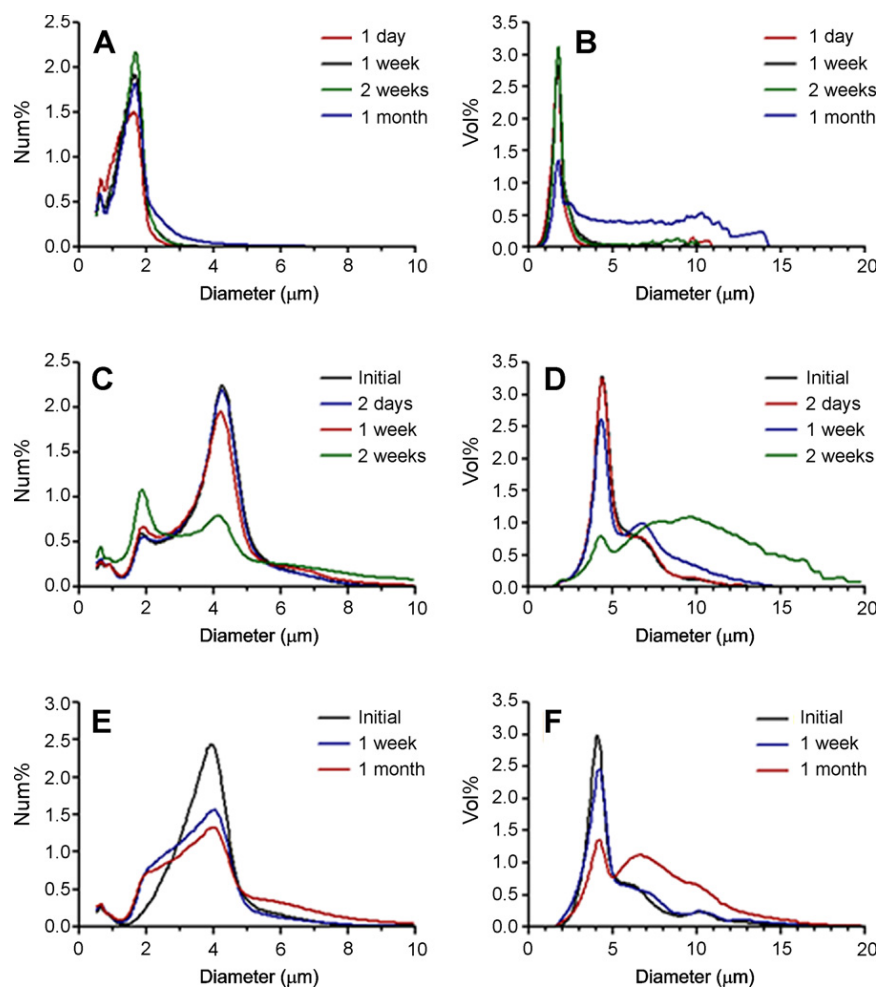
Microbubble Diameter	Fluorescence Intensity (Median $\pm$ SD)	Forward-Light Scattering (Median $\pm$ SD)	Side-Light Scattering (Median $\pm$ SD)
1 - 2 $\mu\text{m}$	509.54 $\pm$ 16.02	137.00 $\pm$ 5.29	289.33 $\pm$ 4.93
4 - 5 $\mu\text{m}$	1114.33 $\pm$ 35.03	214.33 $\pm$ 7.37	417.33 $\pm$ 23.50
Mixture	775.17 $\pm$ 17.51	193.67 $\pm$ 0.58	375.67 $\pm$ 1.53

Results from particle sizing, microscopy and flow cytometry showed the ability to isolate distinct fractions of the microbubbles at the desired size ranges. Next, we determined how stable these size-isolated suspensions were when stored in the refrigerator. What follows is an analysis of the shelf-life for size-isolated microbubbles.

### 2.4.3. Stability of Size-Isolated Microbubbles

For biomedical applications, it is desired that the microbubbles be stable for at least 48 hours at their respective size distributions. A test of microbubble stability was performed using samples concentrated to  $10^{10} \text{ mL}^{-1}$  for 1-2  $\mu\text{m}$  microbubbles, and  $10^8 \text{ mL}^{-1}$  or  $10^9 \text{ mL}^{-1}$  for 4-5  $\mu\text{m}$  microbubbles, in a 1-mL volume of 20 vol% glycerol in PBS and stored in a sealed 2-mL serum vial with PFB headspace (Fig. 2.8). Table 2.3 shows the concentration and PI for the 1-2 and 4-5  $\mu\text{m}$  microbubbles at various time points following size isolation. Both size fractions were stable over two days. Microscopy after two days storage also indicated the persistence of intact microbubbles at their isolated size range over this timeframe. However, results indicated that the microbubbles underwent ripening during longer-term storage. For 1-2  $\mu\text{m}$  microbubbles, the

concentration decreased from by an order of magnitude, and PI nearly doubled over a period of 28 days. For 4-5  $\mu\text{m}$  microbubbles at less than 1 vol% encapsulated gas, the concentration decreased by more than half, and PI nearly doubled over a period of 14 days. Higher microbubble concentrations provided much greater stability, as seen for the comparison of the 4-5  $\mu\text{m}$  microbubbles in Fig. 2.8. In general, we found that encapsulated gas fractions greater than 1 vol% were necessary for good stability, particularly when the vial is intermittently opened to the atmosphere as typically occurs for an *in vivo* study (data not shown). Interestingly, when measuring the number-weighted distribution with the Accusizer, the monomodal peak for the 4-5  $\mu\text{m}$  microbubbles changed to a bimodal peak during storage. Figure 2.8C shows that 4-5  $\mu\text{m}$  microbubbles degraded into 1-2  $\mu\text{m}$  microbubbles as well as larger microbubbles over the test period. The formation of the larger microbubbles is consistent with Ostwald ripening, in which small microbubbles shrink at the expense of larger ones (Taylor 1998), as a major factor affecting microbubble stability.



**Figure 2.8. Stability of size-isolated microbubbles.** Shown are distributions at various time points for the 1–2 μm (A, B) and 4–5 μm (C, D, E, F) diameter microbubbles. Number-weighted (A, C, E) and volume-weighted (B, D, F) distributions are shown for inspection of polydispersity. The suspensions initially were dispersed in 1-mL volume of PBS with 20 vol% glycerol, with a concentration of  $\sim 10^{10}$  mL $^{-1}$  for the 1–2 μm diameter microbubbles and either <1 vol% (C, D) or >1 vol% (E, F) for the 4–5 μm diameter microbubbles. Each curve is the average of three experiments with three measurements each. Results are summarized in Table 2.3.

**Table 2.3: Summary of stability of size-isolated microbubble suspensions.**

Time	Total Concentration (#/mL)	Volume Fraction (mL/mL)	Number Weighted Diameter ( $\mu\text{m}$ )		Volume Weighted Diameter ( $\mu\text{m}$ )		PI
			(Mean $\pm$ SD)	(Median $\pm$ SD)	(Mean $\pm$ SD)	(Median $\pm$ SD)	
1-2 micron bubble stability							
Initial	$2.3\text{E}+10 \pm 1.4\text{E}+09$	$3.7 \pm 0.4\%$	$1.3 \pm 0.1$	$1.2 \pm 0.1$	$1.9 \pm 0.1$	$1.8 \pm 0.1$	$1.5 \pm 0.1$
1 day	$8.3\text{E}+09 \pm 2.9\text{E}+09$	$1.1 \pm 0.4\%$	$1.2 \pm 0.1$	$1.2 \pm 0.1$	$1.7 \pm 0.1$	$1.6 \pm 0.1$	$1.4 \pm 0.1$
7 days	$7.6\text{E}+09 \pm 1.3\text{E}+09$	$1.0 \pm 0.2\%$	$1.3 \pm 0.1$	$1.4 \pm 0.1$	$1.8 \pm 0.1$	$1.7 \pm 0.1$	$1.4 \pm 0.1$
14 days	$4.3\text{E}+09 \pm 4.2\text{E}+08$	$0.79 \pm 0.09\%$	$1.4 \pm 0.1$	$1.4 \pm 0.1$	$2.0 \pm 0.1$	$1.7 \pm 0.1$	$1.4 \pm 0.1$
28 days	$1.1\text{E}+09 \pm 1.3\text{E}+08$	$0.39 \pm 0.06\%$	$1.5 \pm 0.1$	$1.4 \pm 0.1$	$4.1 \pm 0.1$	$2.6 \pm 0.1$	$2.8 \pm 0.1$
4-5 micron bubble stability for suspensions containing less than 1% gas volume							
Initial	$1.2\text{E}+08 \pm 2.8\text{E}+07$	$0.34 \pm 0.08\%$	$3.2 \pm 0.1$	$3.4 \pm 0.1$	$4.8 \pm 0.1$	$4.5 \pm 0.1$	$1.5 \pm 0.1$
2 days	$1.0\text{E}+08 \pm 2.5\text{E}+06$	$0.29 \pm 0.01\%$	$3.2 \pm 0.1$	$3.4 \pm 0.1$	$4.8 \pm 0.1$	$4.4 \pm 0.1$	$1.5 \pm 0.1$
7 days	$9.6\text{E}+07 \pm 1.1\text{E}+06$	$0.28 \pm 0.03\%$	$3.1 \pm 0.1$	$3.2 \pm 0.1$	$5.3 \pm 0.1$	$4.6 \pm 0.1$	$1.7 \pm 0.1$
14 days	$3.9\text{E}+07 \pm 8.6\text{E}+06$	$0.15 \pm 0.02\%$	$2.8 \pm 0.1$	$2.1 \pm 0.1$	$8.1 \pm 0.1$	$7.6 \pm 0.1$	$3.0 \pm 0.1$
4-5 micron bubble stability on suspensions containing greater than 1% gas volume							
Initial	$6.30\text{E}+08 \pm 2.9\text{E}+07$	$1.9 \pm 0.1\%$	$3.3 \pm 0.1$	$3.5 \pm 0.1$	$4.9 \pm 0.1$	$4.3 \pm 0.1$	$1.5 \pm 0.1$
7 days	$5.01\text{E}+08 \pm 4.6\text{E}+07$	$1.1 \pm 0.1\%$	$2.9 \pm 0.1$	$2.9 \pm 0.1$	$4.8 \pm 0.3$	$4.1 \pm 0.1$	$1.6 \pm 0.1$
28 days	$2.98\text{E}+08 \pm 2.5\text{E}+07$	$1.0 \pm 0.2\%$	$3.1 \pm 0.1$	$3.0 \pm 0.1$	$6.2 \pm 0.1$	$6.0 \pm 0.1$	$2.0 \pm 0.1$

The greater stability of 1-2  $\mu\text{m}$  microbubble suspensions versus the 4-5  $\mu\text{m}$  microbubble suspensions observed here could be explained by the higher concentration of microbubbles, which itself is a direct result of the more abundant population of the smaller microbubbles. We found in general that more concentrated suspensions tended to be more stable, regardless of microbubble size. This presumably was due to a thicker cake coating the top surface that buffered against film rupture and gas release at the surface (“popping”) that could diminish the population. However, this does not explain the greater number of 1-2  $\mu\text{m}$  microbubbles in the initial formulation, nor the observation that the 4-5  $\mu\text{m}$  microbubbles broke down over time into 1-2  $\mu\text{m}$  microbubbles. These two observations clearly point to greater stability for the 1-2  $\mu\text{m}$  microbubbles in these formulations. The presence of a stable microbubble size is not predicted by any current theory of microbubble stability. Models that account for microbubble dissolution, such as Epstein and Plesset’s original formulation (Epstein and Plesset 1950), clearly indicate that smaller microbubbles should be less stable, due to higher curvature and surface-to-volume ratio. Yet we observed microbubbles that were more stable at the smaller diameter! This data correlates well with the stable diameter previously observed during acoustically driven dissolution (Borden, Kruse et al. 2005; Wrenn 2008), and also the stable diameter reached for lipid-coated microbubbles made via microfluidics (Talu, Hettiarachchi et al. 2008). However, submicron bubbles were found to be less stable than 1-2  $\mu\text{m}$  microbubbles, which indicated the existence of an optimal microbubble size.

What could explain the presence of an optimally stable diameter for these lipid-coated microbubbles? We speculate that the greater stability is due to microstructural features of the lipid shell and lipid aggregates present in the continuous phase. Dressaire *et al.* recently described the stability of nanopatterned surfactant microbubbles arising from the balance

between domain bending and pressure-volume work (i.e., Laplace overpressure) (Dressaire, Bee et al. 2008). In their analysis, however, size was fixed by the blending conditions. It may be that size was fixed here by cavitation during acoustic emulsification, according to Li and Fogler's analysis (Li and Fogler 1978; Li and Fogler 1978). However, this does not explain the same stable diameter observed by Talu *et al.* (Talu, Hettiarachchi et al. 2008) for microbubbles made by microfluidics, in which no cavitation shock waves were produced, nor does it explain the stability against acoustically driven dissolution at this size (Borden, Kruse et al. 2005). Future work on the relationships between lipid nanostructural features and microbubble stability is clearly warranted to better explain this phenomenon.

## **2.5. Conclusions**

Lipid-coated microbubbles formed by acoustic emulsification were found to be polydisperse and appeared to be multimodal, with distinct peaks centered near 1-2 and 4-5  $\mu\text{m}$  diameter. Differential centrifugation was used successfully to isolate narrowly dispersed fractions at these size ranges. These size ranges are stable over a period of at least two days, although the 4-5  $\mu\text{m}$  microbubbles were found to disintegrate into 1-2  $\mu\text{m}$  microbubbles after two weeks. This latter observation indicates a stable diameter in the 1-2  $\mu\text{m}$  range for these microbubbles, which is supported by observations of microbubbles that underwent acoustically driven dissolution and those formed by microfluidics. Overall, differential centrifugation appears to be a useful approach to generate narrow distributions at relevant sizes for biomedical applications and may lend itself to the study of the surface properties and colloidal behavior of different microbubble size classes.



## **Chapter 3. Theranostic Gd(III)-Lipid Microbubbles for MRI-Guided Focused Ultrasound Surgery**

### **3.1. Introduction**

MRI-guided focused ultrasound therapy is a rapidly developing medical technique that utilizes high intensity focused ultrasound (FUS) to ablate tissue and MRI to monitor the applied thermal dosage (Cline, Schenck et al. 1992; Hynynen, Darkazanli et al. 1993; Jolesz and McDannold 2008). MRI-guided FUS therapy is approved by the US Food and Drug Administration for the treatment of uterine fibroids, and it is currently being developed to treat liver, bone, prostate and brain-related diseases (Hudson and Stewart 2008; Jolesz and McDannold 2008; Jolesz 2009). At high acoustic intensities, gas-filled microbubbles (MBs) may form and undergo inertial cavitation, producing jets and shockwaves that enhance the heating of tissue. However, the formation of these inception microbubbles is unpredictable, and their cavitation can result in tissue damage outside of the desired target region. Pre-formed microbubbles, which are currently FDA-approved as intravascular contrast agents for echocardiography, can be used as cavitation nuclei to lower the acoustic intensity threshold required for tissue ablation with FUS, thereby lowering the thermal buildup in surrounding tissue (Tran, Seo et al. 2003; Kaneko, Maruyama et al. 2005; Tung, Liu et al. 2006; Yu, Hu et al. 2008). Intravenously administered microbubbles also may be used to enhance vascular permeability for targeted drug and gene delivery (Unger, Porter et al. 2004; Ferrara, Pollard et al. 2007). For example, microbubbles have been used to lower the acoustic intensity threshold needed for FUS-induced blood-brain barrier (BBB) opening (Hynynen, McDannold et al. 2001; Choi, Pernot et al. 2007). For these applications, it would be advantageous to use an MRI-

detectable microbubble formulation, which could be used to measure microbubble concentration, image cavitation events and determine the biodistribution of microbubble shell debris (a potential surrogate for an attached drug) following FUS. Such microbubbles may also be useful as general dual modality US/MRI contrast agents.

Previously, microbubbles were demonstrated to enhance the  $T_2^*$ -weighted MRI contrast *in vivo* by virtue of the change in magnetic susceptibility at the gas-liquid interface (Wong, Huang et al. 2004; Cheung, Chow et al. 2009). This negative enhancement was reportedly a linear function of gas concentration and was further increased in subsequent studies by loading  $T_2$ -weighted MRI contrast agents (iron oxide particles) onto the shell of polymeric microbubbles (Chow, Chan et al. ; Yang, Li et al. 2008; Yang, Li et al. 2009). Recently, Lui *et al.* (Liu, Lammers et al. 2011) reported that ultrasonic fragmentation of magnetite-loaded polymeric microbubbles resulted in greater proton relaxation than for the intact microbubbles. This effect was attributed to the greater interaction of peripheral water to released iron oxide particles. These superparamagnetic microbubbles offer significant potential as theranostic agents for MRI-guided FUS.

An alternative means to produce dual modality US/MRI contrast agents is to load microbubbles with paramagnetic contrast agents, such as gadolinium ions. Gd(III) enhances the positive contrast of blood by shortening both the longitudinal and transverse proton relaxation times,  $T_1$  and  $T_2$  (Caravan, Ellison et al. 1999; Aime, Botta et al. 2005; Hermann, Kotek et al. 2008). Previously, Gd(III)-DTPA was loaded into the shell of 1.5- $\mu\text{m}$  diameter polymeric microbubbles (Ao, Wang et al. 2010). The enhancement of the  $T_1$ -weighted MRI signal was reportedly a linear function of Gd(III)-DTPA-loaded microbubble concentration. However, polymeric microbubbles tend to be much stiffer than lipid-coated microbubbles, providing less

echogenicity for ultrasound imaging and requiring greater acoustic intensity to induce microbubble fragmentation and sonoporation for therapy (Hoff, Sontum et al. 2000; Bloch, Wan et al. 2004).

Lipid-coated microbubbles with Gd(III)-bound shells have not been reported in literature previously. However, Gd(III)-bound liposomes have been designed and characterized as  $T_1$ -weighted MRI contrast agents for applications in cellular and small animal imaging (Ghaghada, Hawley et al. 2008; Terreno, Castelli et al. 2008; Ghaghada, Ravoori et al. 2009; Hak, Sanders et al. 2009; Kamaly, Kalber et al. 2010). Liposomes and lipid-coated microbubbles are similar with respect to their lipid composition and formulation (Ferrara, Borden et al. 2009). The main difference between the two is that microbubbles consist of a condensed monolayer with a gas core and are typically a few microns in size, while liposomes consist of a lipid bilayer with an aqueous core and are usually several hundred nanometers in size. Gd(III) can be loaded into the liposomal aqueous core and/or conjugated to the lipid polar headgroups in the bilayer (Ghaghada, Hawley et al. 2008; Ghaghada, Ravoori et al. 2009; Hak, Sanders et al. 2009). While both strategies were reported to increase the  $T_1$ -weighted MRI relaxation rate, surface conjugation resulted in greater relaxation enhancement than encapsulation, owing to greater access of bulk water protons to the Gd(III) ions (Ghaghada, Ravoori et al. 2009). Owing to the presence of the gas core and thin monolayer shell, lipid-coated microbubbles can only be functionalized with Gd(III) ions using a surface conjugation methodology. Since the lipid coating of microbubble shells self-assembles into liposomal bilayers in the absence of the gas core, a comparison of the MRI relaxation rates of Gd(III)-bound microbubbles before and after fragmentation would be necessary.

Below, we report the fabrication and characterization of lipid-coated microbubbles that were surface-conjugated with the paramagnetic MRI contrast agent, Gd(III). Gadolinium was selected over iron oxide since it works primarily to enhance positive contrast through increasing longitudinal proton relaxation (Caravan, Ellison et al. 1999). The chelation ligand DOTA was chosen over DTPA since it forms a more thermodynamically stable complex with Gd(III) (De Leon-Rodriguez and Kovacs 2008). A stronger chelator is preferred since free Gd(III) ions *in vivo* have been associated with nephrogenic systemic fibrosis (Sherry, Caravan et al. 2009). Additionally, the 4-5  $\mu\text{m}$  microbubble size range was selected because of its increased acoustic signal and circulation persistence (Sirsi, Feshitan et al. 2010; Streeter, Gessner et al. 2010) and more effective BBB opening capability compared to polydisperse and smaller size-ranged microbubbles (Choi, Feshitan et al. ; Tung, Marquet et al. 2011). The resulting microbubbles were tested for ultrasound contrast *in vivo* and MRI contrast *in vitro*. Finally, we report both longitudinal ( $R_1 = 1/T_1$ ) and transverse ( $R_2 = 1/T_2$ ) relaxation rates of the Gd(III)-bound microbubbles before and after they were fragmented into lipid bilayers by sonication.

## **3.2. Materials and Methods**

### **3.2.1. Materials**

All solutions were prepared using filtered, 18 M $\Omega$ -cm deionized water (Direct-Q, Millipore, Billerica, MA). Glassware was cleaned with 70 vol% ethyl alcohol solution (Sigma-Aldrich; St. Louis, MO) and rinsed with deionized water. The gas used to form microbubbles was perfluorobutane (PFB) at 99 wt% purity obtained from FluoroMed (Round Rock, TX). 1,2-distearoyl-sn-glycero-3-phosphoethanolamine (DSPE) was purchased from Avanti Polar Lipids

(Alabaster, AL). 1,2-distearoyol-sn-glycero-3-phosphoethanolamine-N-[methoxy(polyethylene glycol)2000] (DSPE-PEG(2000)) was obtained from NOF America Corporation (White Plains, NY). 5/6-carboxyfluorescein succinimidyl ester (FITC-NHS) was purchased from Pierce (Rockford, IL). 1,4,7,10-Tetraazacyclododecane-1,4,7,10-tetraacetic acid mono(*N*-hydroxysuccinimide ester) (DOTA-NHS) was purchased from Macrocyclics (Dallas, TX) and dissolved in *N,N*-dimethylformamide (DMF; Sigma-Aldrich) prior to use. Gadolinium (III) chloride ( $\text{GdCl}_3$ ) was purchased from Sigma Aldrich and dissolved in 0.2 M, pH 5.6 acetate buffer (VWR, Radnor, PA).

### 3.2.2. Microbubble Generation and Size Isolation

Microbubbles were formulated using a lipid suspension of 90 mol% DSPE and 10 mol% DSPE-PEG(2000) at 2 mg/mL in 100 mL PBS (pH 7.2; 0.15 M NaCl, 0.2 M phosphate). The solution was degassed by applying house vacuum with constant stirring. The solution was then preheated to 80 °C, which is 6 °C above the main phase transition temperature ( $T_m$ ) of DSPE (Cevc and Marsh 1985). The lipid mixture was sonicated with a 20-kHz probe (model 250A, Branson Ultrasonics; Danbury, CT) at low power (3 watts) in order to further disperse the lipid aggregates into small, unilamellar liposomes. PFB gas was introduced by flowing it over the surface of the lipid suspension. Higher power sonication (33 watts) was applied to the suspension for about 10 seconds at the gas-liquid interface to generate microbubbles.

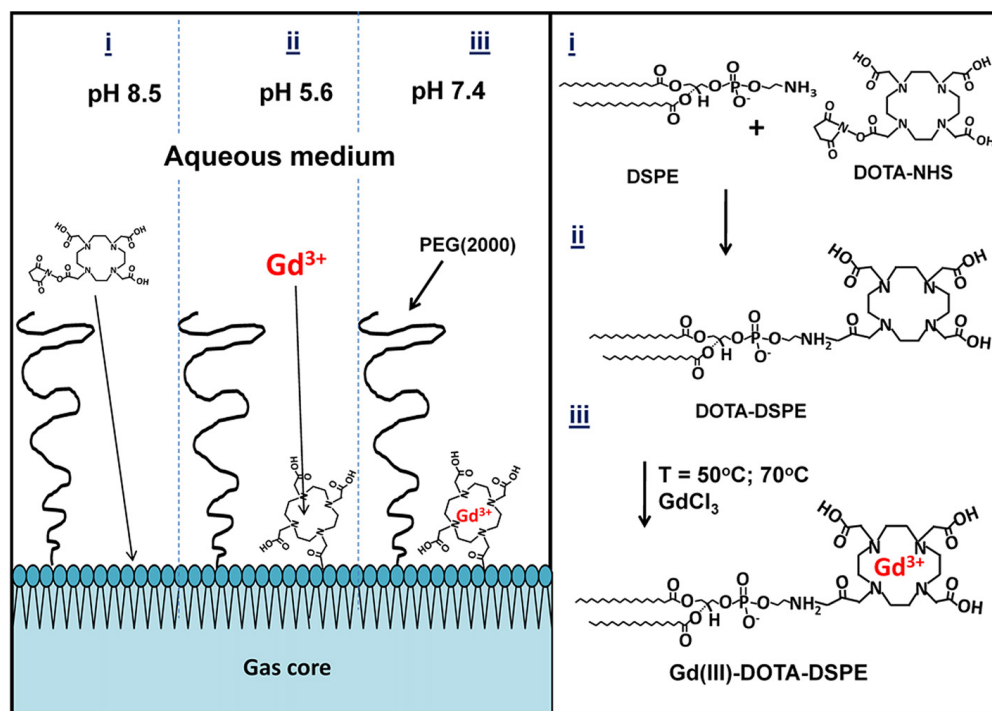
The microbubble suspension was collected into 30 mL syringes (Tyco Healthcare, Mansfield, MA), which were used as the flotation columns. Washing and size-selection by centrifugation was performed with a bucket-rotor centrifuge (model 5804, Eppendorf, Westbury, NY).

Centrifugation at 300 RCF (relative centrifugal force) for 5 min was performed to collect all microbubbles from the suspension into a cake resting against the syringe plunger. The remaining suspension (infranatant), which contained residual lipids and vesicles, was recycled to produce the next batch of microbubbles. All resulting cakes were combined and re-suspended in PBS to improve total yield.

Microbubble size distribution was determined by laser light obscuration and scattering (Accusizer 280A, NICOMP Particle Sizing Systems, Santa Barbara, CA). During measurements, 2  $\mu$ L samples of each microbubble suspension were diluted into a flask containing 30 mL of distilled water under mild mixing. All samples were measured in triplicate and analyzed for both number- and volume-weighted size distributions. The 4-5  $\mu$ m size class was isolated as described in chapter 2 and reconstituted in pH 8.5 PBS.

### **3.2.3. Synthesis of Gd(III)-bound Microbubbles**

Microbubbles with Gd(III)-bound shells were fabricated using a post-labeling technique (Chen and Borden 2010; Chen and Borden 2011). The macrocyclic ligand DOTA-NHS was conjugated to the amine group of the DSPE in the shell of size-selected microbubbles, followed by chelation of Gd(III). The NHS ester contains an electrophilic active group that couples rapidly with the primary amine on DSPE to create a stable amide bond. Figure 3.1 shows a schematic of the overall conjugation process.



**Figure 3.1.** Synthesis of the Gd(III)-DOTA-DSPE microbubble shells using the post-labeling technique: (i) 100 M excess DOTA-NHS, pH 8.5; (ii) 20 M excess GdCl<sub>3</sub>, pH 5.6, T = 50 or 70 °C; (iii) storage at pH 7.4.

### 3.2.4. Surface Functionalization with FITC-NHS or DOTA-NHS

Each 4-5 μm microbubble sample was diluted to  $2 \times 10^9$  MB/mL using pH 8.5 PBS. Following Chen and Borden, the total amount of available functional lipid groups (DSPE) on the microbubble surface was calculated assuming that the microbubbles were spherical with an average molecular area of 0.4 nm<sup>2</sup>. To test the post-labeling headgroup conjugation method, FITC-NHS was added to a 100:1 molar ratio of NHS to amine, and the mixture was continually stirred at room temperature for 2 hours using a benchtop rotary mixer. To chelate Gd(III), DOTA-NHS was added to a 100:1 molar ratio of NHS to amine, and the suspension was mixed

as above. Unreacted FITC-NHS or DOTA-NHS was removed by several cycles of flotation using 0.2 M, pH 5.6 acetate buffer. The microbubble cake was then analyzed for size with the Accusizer.

### **3.2.5. Complexation of Gd(III) to DOTA on Microbubble Shells**

Based on the initial concentration and size distribution calculated from the Accusizer, each sample of DOTA-bound microbubbles was diluted to at least  $2 \times 10^9$  MB/mL using pH 5.6 acetate buffer. Assuming 100% binding of DOTA to available functional DSPE lipid groups, the amount of  $\text{GdCl}_3$  needed for a 20:1 molar ratio of Gd(III) to DOTA was determined and mixed with the microbubble suspension. The sample mixture was sealed in a 3 mL serum vial then immersed under continuous stirring in a water bath whose temperature was controlled at 50 °C or 70 °C for 2 hours. After reaction, the sample mixture was cooled to room temperature by running the vial under cold tap water for 10 min. Excess Gd(III) ions were removed by several cycles of washing/centrifuging (1 minute, 100 RCF) using pH 5.6 acetate buffer followed by several cycles of washing/centrifuging using pH 7.4 PBS as solvent medium. The final microbubble cake was reconstituted to a volume of 1 mL and a concentration of at least  $1 \times 10^9$  MB/mL using pH 7.4 PBS. The size distribution and concentration of microbubbles after chelation reaction were determined by the Accusizer. The concentration of Gd(III)-bound to the microbubble shell was determined by inductively coupled plasma optical emission spectroscopy (ICP-OES, ACTIVA, HORIBA, Edison, NJ). Destruction/fragmentation of microbubbles in suspension was accomplished by simultaneous bath sonication and heating to 80 °C for 5 min.



### 3.2.6. Ultrasound Characterization of Gd(III)-bound Microbubbles

All animal experiments were conducted according to the National Institutes of Health guidelines and approved by the University of Colorado Institutional Animal Care and Use Committee. Ultrasound imaging was performed using a high-frequency ultrasound scanner (Vevo 2100, Visualsonics, Toronto, Ontario, Canada) with a MS-250 transducer. Images were acquired using the contrast mode setting at 18 MHz transmit frequency and 4 % power. The transducer was positioned at the mouse midsection along the long axis of the kidney. B-mode ultrasound images were acquired using a field of view of 13 x 16 mm<sup>2</sup>. Mice were anesthetized with 3 % isoflurane and tail veins were catheterized for injections, as previously described (Sirsi, Feshitan et al. 2010). A 100- $\mu$ L bolus ( $1 \times 10^8$  MB/mL) followed by a 15- $\mu$ L saline flush was injected while imaging at the maximum frame rate for respiratory gating (~14 frames/second). B-mode images captured before and after microbubble injection were used to detect signal enhancement using a background reference subtraction method.

### 3.2.7. MRI Characterization of Gd(III)-bound Microbubbles

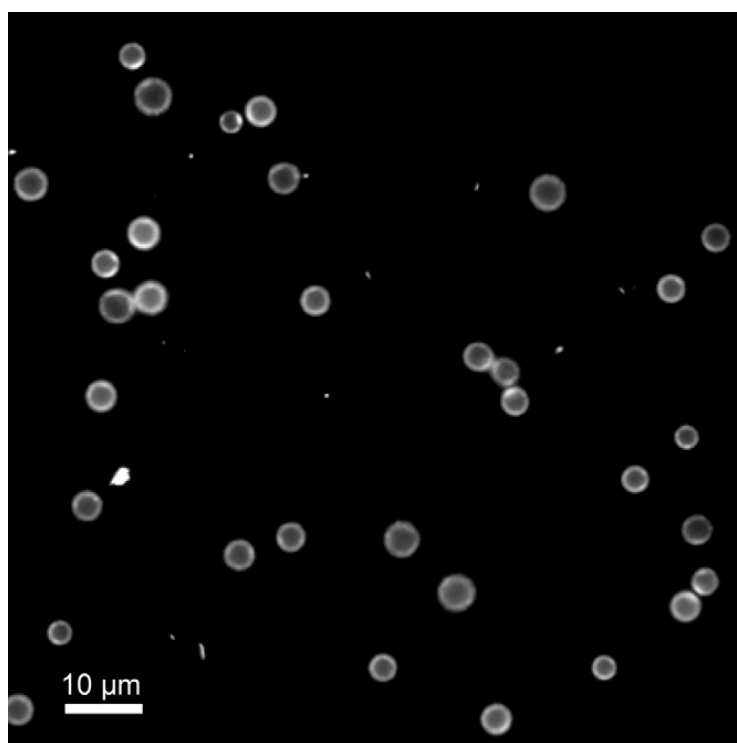
The effect of Gd(III)-bound microbubbles on the  $T_1$  and  $T_2$  relaxation times was determined using MRI relaxometry. Intact and fragmented Gd(III)-bound microbubbles were mixed with saline in four different volume ratios (0, 25, 50 and 100%) creating 200  $\mu$ L solutions, which were placed in MR-compatible tubes with an inner diameter of 5 mm. Intact and fragmented 4-5  $\mu$ m DOTA-bound microbubbles without Gd(III) binding were used as controls. A 9.4 T vertical MRI system (Bruker Biospin, Billerica, MA) was used to acquire turbo spin echo (RARE-VTR) images with variable repetition times (from 300 to 12,500 ms) and multi-slice multi-echo

(MSME) images with variable echo times (from 20 to 320 ms) for  $T_1$  and  $T_2$  mapping, respectively. This spin-echo sequence reportedly lacks sensitivity to inhomogeneities in the magnetic field compared to gradient-echo sequences used in susceptibility-weighted imaging of microbubbles (Berns, Ross et al. 1991; Cheung, Chow et al. 2009). Eight 1.5 mm-thick, axial slices with a field of view (FOV) of  $15 \times 15 \text{ mm}^2$  (matrix size:  $96 \times 96$ ) covered the entire solution in each tube. Each slice depicted a slab of all four solutions at a specific height.  $T_1$  and  $T_2$  relaxation maps of each slice were derived using the Image Processing Toolbox of MATLAB R2008b (MathWorks Inc., Natick, MA). The first and last slice were not taken into account in the relaxation measurements, since the MR signal coming from these slices was contaminated by the void below and over the solution. The pixel-by-pixel estimations were used to generate  $T_1$  and  $T_2$  maps. Four pre-defined, identical, circular regions of interest (ROI) of 2.35 mm in diameter were selected on each slice, in order to measure the relaxation rate of each solution throughout the tube. Each ROI covered a large surface area within the limits of the tube. Six measurements were made for each tube (from slice 1 to 6) and the mean value yielded the  $T_1$  or  $T_2$  relaxation times for each solution. Molar relaxivities ( $r_1$  and  $r_2$ ) were calculated from slopes of regression lines of the plot of  $R_1$  and  $R_2$  versus Gd(III) concentration. A two-tailed unpaired student t-test was used to determine the significance between  $r_1$  and  $r_2$  of fragmented versus intact samples assuming a Gaussian distribution.

### 3.3. Results

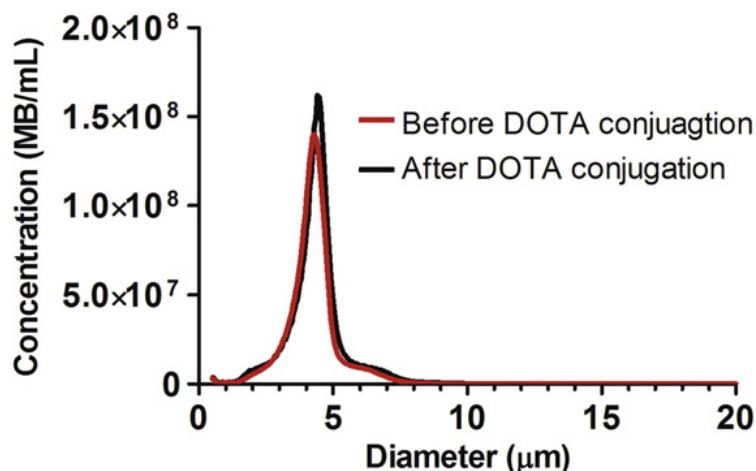
#### 3.3.1. Preparation of Gd(III)-bound Microbubbles

The size isolation protocol yielded 4-5  $\mu\text{m}$  diameter microbubbles at a concentration a least  $2 \times 10^9$  MB/mL. Figure 3.2 shows visual confirmation of FITC-NHS coupling to the DSPE shell using epi-fluorescence microscopy. The fluorescence intensity level was not uniform across individual microbubble, which suggested a heterogeneous coverage of FITC on lipid monolayer. However, this result confirmed that small molecules ( $< 1$  kDa) could diffuse through the PEG brush layer to react with the polar lipid headgroups. This is an extension of previous post-labeling work, which showed reactions occurring on PEG-tethered active groups (Klibanov 2005; Chen and Borden 2010; Chen and Borden 2011), but not with the underlying lipid.



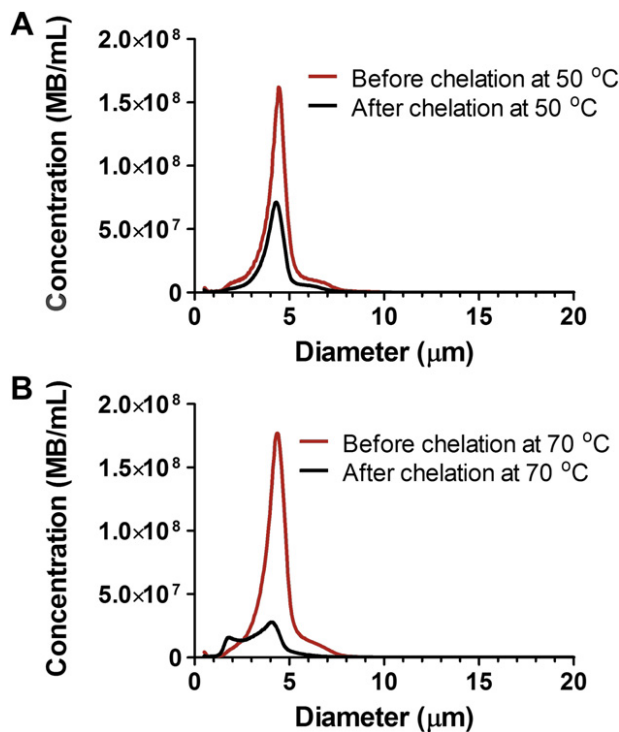
**Figure 3.2. Fluorescence microscopy image of 4-5  $\mu\text{m}$  DSPE-coated microbubbles modified with FITC-NHS using the post-labeling technique. Scale bar represents 10  $\mu\text{m}$ .**

Figure 3.3 shows the change in size distribution of microbubbles before and after conjugation of DOTA (at room temperature) to the primary amines on the DSPE shell via NHS coupling. After DOTA conjugation, microbubble concentration and number-weighted median diameter deviated by less than 1%. Thus, the DOTA reaction did not appear to increase the lipid headgroup area sufficiently to affect lipid packing and thereby change microbubble size or stability. We did not detect Gd(III) binding to DOTA-microbubbles after incubation at room temperature for several hours (data not shown). The Gd(III)-DOTA complex has been reported to take several days to complete at room temperature (De Leon-Rodriguez and Kovacs 2008). This is because the rate-determining step involves the slow, base-assisted rearrangement and deprotonation of an intermediate before formation of the final complex (Sherry, Caravan et al. 2009). Previous researchers have completed the Gd(III)-DOTA chelation reaction in 5 min by heating reactants to 90 °C, or in 20 min by heating to 80 °C (De Leon-Rodriguez and Kovacs 2008). However, these temperatures are above the main phase transition temperature ( $T_m$ ) of the lipid component DSPE (74 °C) and may have resulted in significant microbubble destabilization. We therefore tested Gd(III) chelation onto microbubbles incubated at 50 °C and 70 °C for 2 hr.



**Figure 3.3. Number-weighted size distributions of DSPE microbubbles before and after conjugation with DOTA-NHS.**

Figure 3.4 shows the change in size distribution of microbubbles before and after chelation of Gd(III) under these conditions. After chelation at 50 °C, microbubble concentration decreased by ~50% while the number-weighted median diameter deviated by less than 1%. After chelation at 70 °C, however, microbubble concentration decreased by ~65% while the number-weighted median diameter also decreased by ~30%. From ICP-OES analysis, the Gd(III) chelation on the microbubble shell occurring at 70 °C and 50 °C was  $7.0 \times 10^5 \pm 1.6 \times 10^5$  (mean  $\pm$  standard deviation) and  $7.5 \times 10^5 \pm 3.0 \times 10^5$  ions/ $\mu\text{m}^2$ , respectively. Therefore, all subsequent chelation reactions were carried out at 50 °C since the size distribution of microbubbles was maintained at this temperature without affecting the degree of Gd(III) binding. Under these conditions, the average Gd(III) loading was  $3.6 \times 10^7 \pm 1.0 \times 10^7$  ions/microbubble. ICP-OES analysis also determined that negligible amounts of Gd(III) bound to lipid-coated microbubbles without DOTA (data not shown).

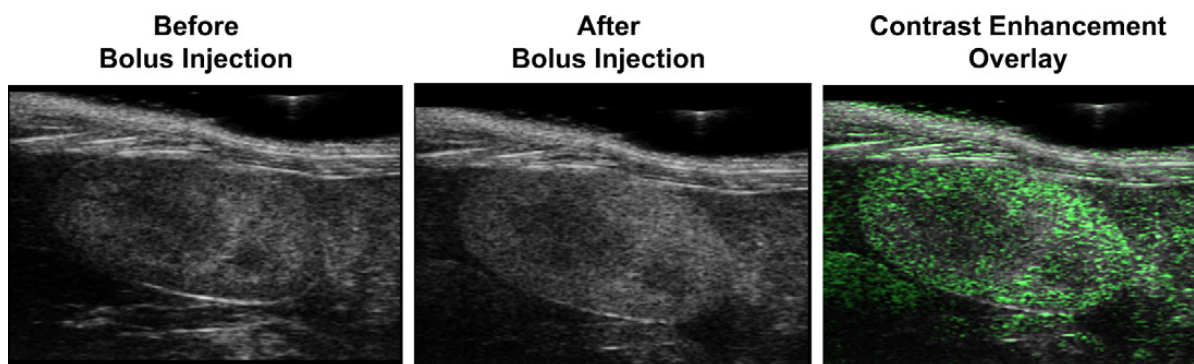


**Figure 3.4. Number-weighted size distributions of microbubbles before and after Gd(III) chelation at A) 50 °C and B) 70 °C.**

Thus, the post-labeling methodology provided a robust means of generating size-selected, Gd-DOTA-lipid microbubbles. Previous work showed that small molecules are capable of diffusing through a PEG overbrush on the microbubble surface to bind to functional groups tethered by shorter PEG chains (Chen and Borden 2010; Chen and Borden 2011). Here, we showed that the small molecule DOTA-NHS is capable of diffusing through the PEG brush to bind to a functional amine on the lipid headgroup. The average molecular area was  $\sim 1\text{-}2 \text{ nm}^2$  per Gd-DOTA complex. This value was higher than that of the minimum molecular area for a lipid ( $\sim 0.4 \text{ nm}^2$ ), indicating that roughly 20-40% of the DSPE was conjugated to Gd-DOTA. This fraction is similar to previous reports for Gd-DOTA-DSPE liposomes (Hak, Sanders et al. 2009) and Gd-DTPA bis(stearylamide) liposomes (Ghaghada, Ravoori et al. 2009).

### 3.3.2. Ultrasound Characterization of Gd(III)-microbubbles

Lipid-coated microbubbles labeled with Gd(III) were tested for echogenicity in the mouse kidney using a preclinical ultrasound scanner. Figure 3.5 shows the B-mode images before and after microbubble injection. A bolus injection of  $1 \times 10^7$  Gd(III)-bound microbubbles significantly increased the fundamental mode backscatter, as was evident by an increase in video intensity and speckling throughout the kidney region. Higher microbubble doses (e.g.,  $5 \times 10^8$ ) led to strong contrast enhancement in the upper portion of the kidney and shadowing in the lower portion (data not shown). These results show that the Gd(III)-bound microbubbles are highly echogenic and suitable for contrast-enhanced US imaging.



**Figure 3.5. Ultrasound images of the mouse kidney before and after bolus injection of  $1 \times 10^7$  Gd(III)-bound microbubbles: (Left) pre-injection, (Middle) post-injection, (Right) contrast enhancement overlay determined using signal subtraction.**

### 3.3.3. MRI Characterization of Gd(III)-microbubbles

Figures 3.6a and 3.6b show the  $T_1$ -weighted and  $T_2$ -weighted MRI maps of fragmented and intact Gd(III)-bound and control (DOTA without Gd(III)) microbubbles. Fragmented microbubbles were produced by the removal of the gas core of intact microbubbles through bath sonication and heating. The color-coding (from red to blue) indicates a greater relaxation effect and therefore an MRI signal intensity increase. Figures 3.7a-d show plots of the longitudinal and transversal relaxation rates ( $R_1$  and  $R_2$ ) of intact and fragmented Gd(III)-bound microbubbles as a function of microbubble concentration, normalized to total Gd(III) concentration using ICP-OES results, for 4 independent trials. Results also are shown for intact and fragmented control microbubbles as a function of surface area.

### 3.3.4. Relaxation Rates of Control Microbubbles and their Fragments

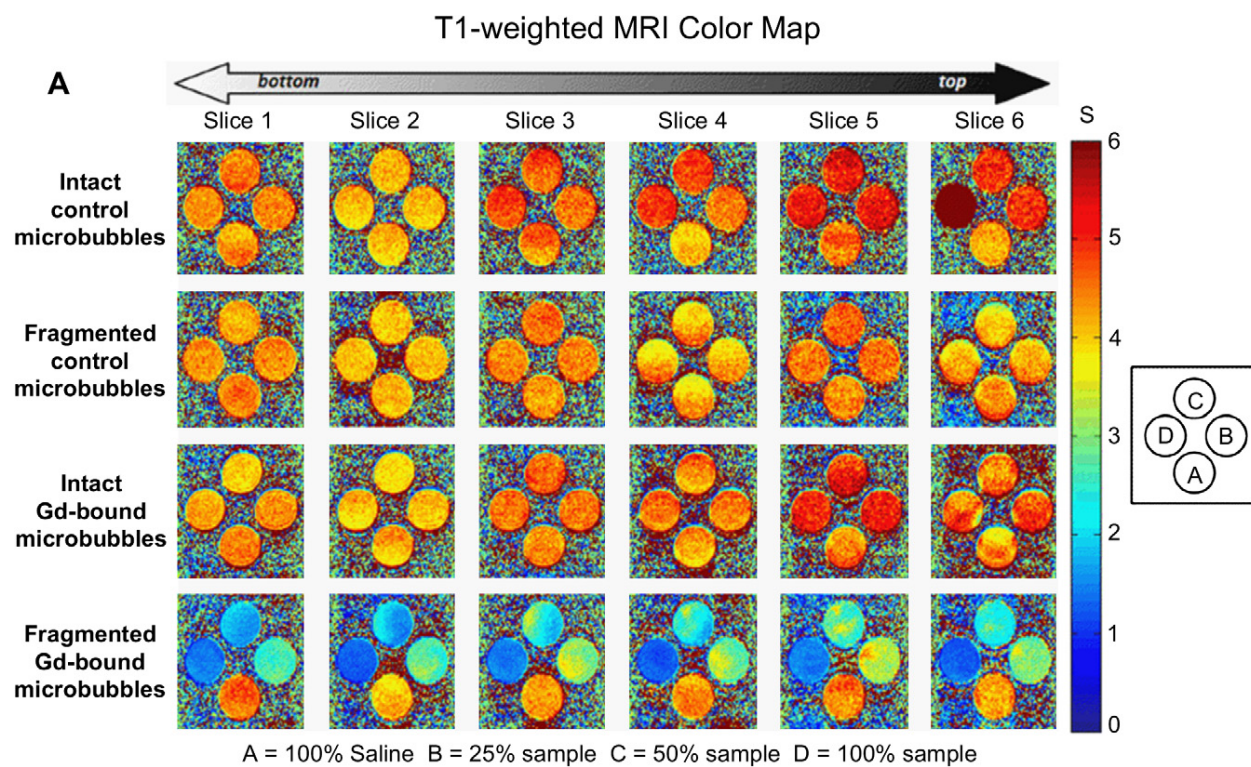
As observed from both the  $T_1$ - and  $T_2$ -weighted color-coded MRI maps (Fig. 3.6), the control microbubbles (DOTA without Gd(III)) produced an MRI signal similar to baseline (saline), which did not deviate significantly with an increase in microbubble concentration. This is further evident in the plot of the relaxation rate versus increasing microbubble surface area for both  $R_1$  and  $R_2$  (Fig. 3.7a,b). The lack of a significant relaxation effect was found for both intact and destroyed control microbubbles.

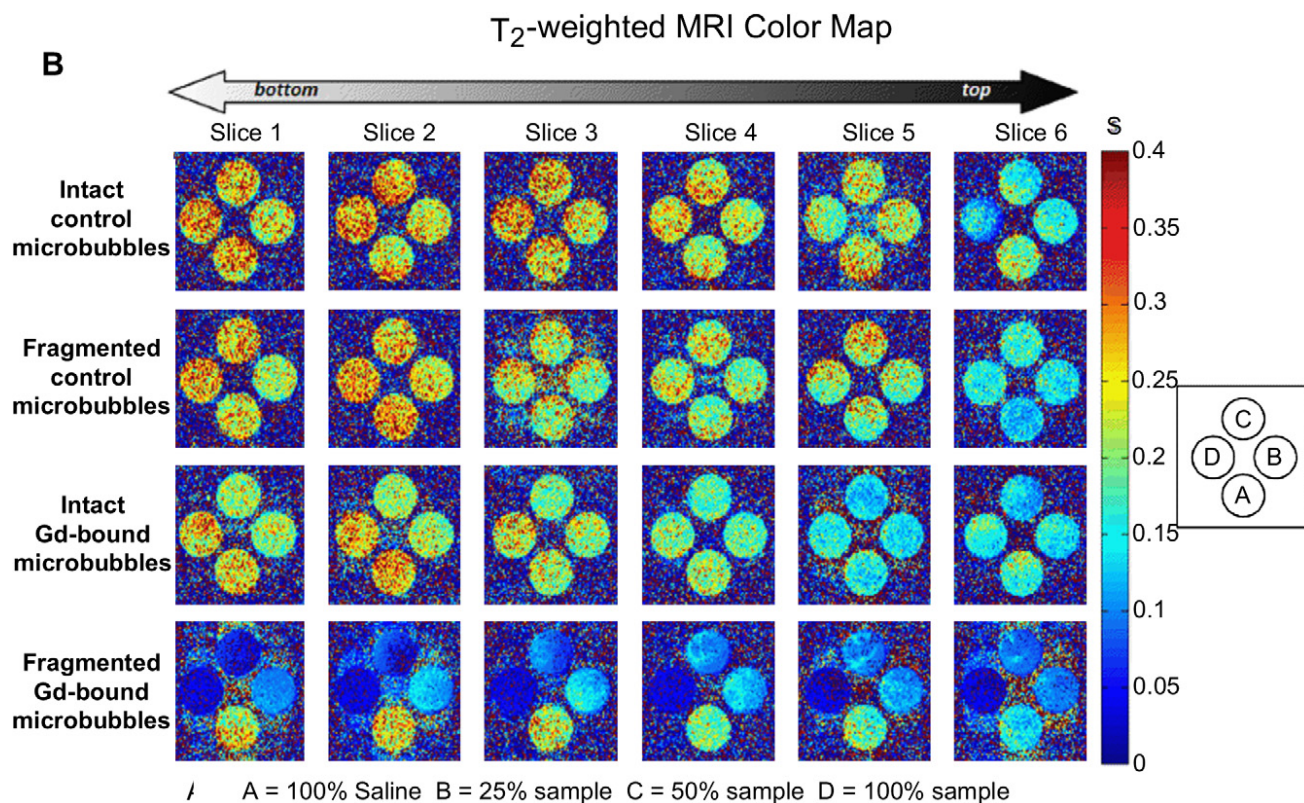


### 3.3.5. Relaxation Rates of Gd(III)-bound Microbubbles and their Fragments

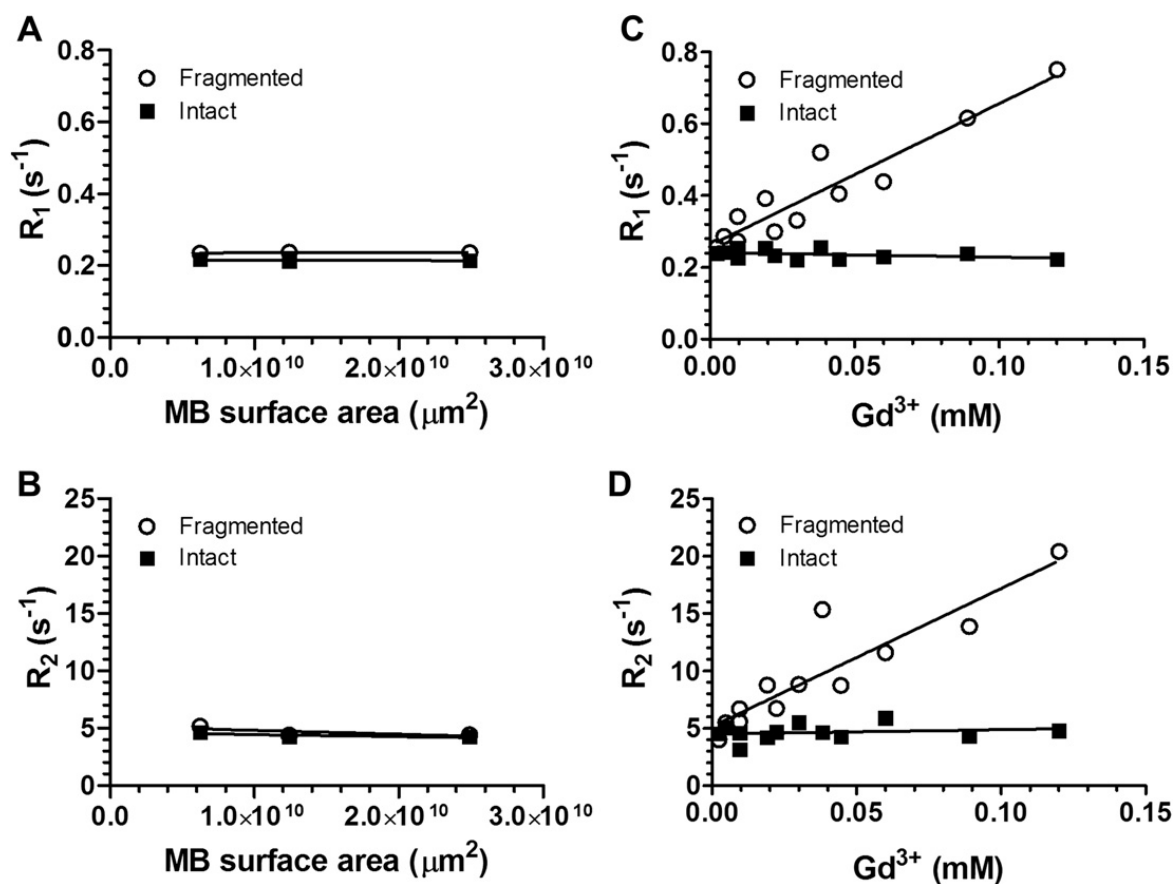
Figure 3.6 also shows that the intact Gd(III)-bound samples produced similar MR signal intensities as saline and control microbubbles, and the signal intensity was not dependent on an increase in sample concentration. Similarly, Figures 3.7c and 3.7d show that the relaxation rate did not increase with increasing intact Gd(III)-bound microbubble concentration (the fitted slope was slightly negative); the MRI signal was similar to that of control samples. This was surprising, as we expected the MRI signal to increase with increasing Gd(III) as has been observed in liposomal suspensions (Ghaghada, Hawley et al. 2008; Terreno, Castelli et al. 2008; Ghaghada, Ravoori et al. 2009; Hak, Sanders et al. 2009; Kamaly, Kalber et al. 2010).

Interestingly, the fragmented Gd(III)-bound microbubbles resulted in a noticeable increase in color-coded MRI signal intensity compared to saline, control and intact Gd(III)-bound microbubbles (Figs. 3.6 and 3.7). Additionally, the effect was concentration-dependent, with an increase in fragmented Gd(III)-bound sample concentration leading to an increase in MRI signal intensity. These results suggest that the MR signal came primarily from the Gd(III) groups and not the other components of the lipid microbubble shell, and the relaxation rate appeared to be most strongly related to the state of the microbubble (i.e., intact vs. fragmented).





**Figure 3.6.** Color maps of MRI relaxation time for intact and fragmented microbubble samples. Longitudinal relaxation time (A) and transverse relaxation time (B) increases from blue to red, as shown. Samples are arranged shown: row 1) intact DOTA-bound control microbubbles; 2) fragmented DOTA-bound control microbubbles; row 3) intact Gd(III)-bound microbubbles; row 4) fragmented Gd(III)-bound microbubbles. Students t-tests showed that, for a given microbubble sample,  $T_1$  and  $T_2$  were not statistically different between vertical slices.



**Figure 3.7. Relaxation rates of intact and fragmented microbubble samples. A)  $R_1$  versus MB surface area; B)  $R_2$  versus MB surface area; C)  $R_1$  versus Gd(III) concentration; D)  $R_2$  versus Gd(III) concentration.**

### 3.3.6. Molar Relaxivities of Gd(III)-bound Microbubbles

Molar relaxivities ( $\text{mM}^{-1}\text{s}^{-1}$ ) were calculated from the slopes of the linear trendlines in Figures 3.7c and 3.7d and are shown in Table 3.1. Fragmentation of the intact Gd(III)-bound microbubble samples led to a 40-fold increase in longitudinal molar relaxivities  $r_1$  ( $p < 0.05$ ) and a 32-fold increase in transverse molar relaxivities  $r_2$  ( $p < 0.05$ ). Thus, both  $r_1$  and  $r_2$  for the fragmented Gd(III)-bound microbubbles were greater than the corresponding values for the

intact Gd(III)-bound microbubbles. A potential mechanism for this surprising phenomenon is discussed below.

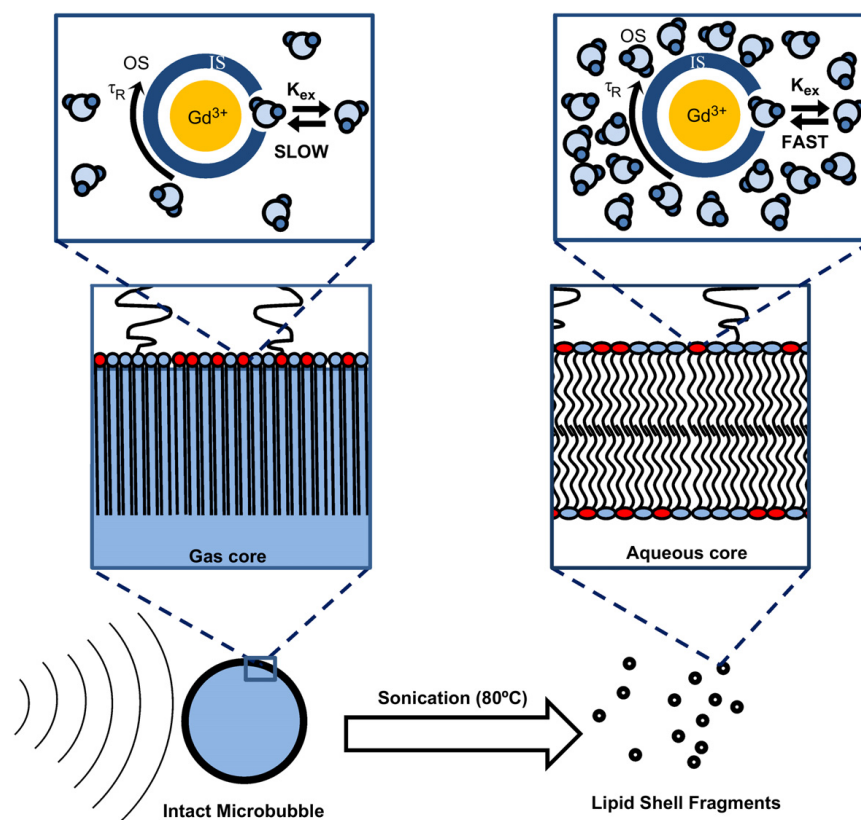
**Table 3.1: Relaxivity of Intact and Fragmented Gd(III)-bound Microbubbles**

	$r_1$ (mM <sup>-1</sup> s <sup>-1</sup> )	$r_2$ (mM <sup>-1</sup> s <sup>-1</sup> )	$r_1$ /MB (s <sup>-1</sup> /MB)	$r_2$ /MB (s <sup>-1</sup> /MB)
Intact MBs	<b>-0.1 ± 0.10</b>	<b>3.8 ± 5.8</b>	<b>-3.6 x 10<sup>6</sup></b>	<b>1.4 x 10<sup>8</sup></b>
Destroyed MBs	<b>4.0 ± 0.4</b>	<b>120.2 ± 17.7</b>	<b>1.4 x 10<sup>8</sup></b>	<b>4.3 x 10<sup>9</sup></b>

### 3.4. Discussion

The relaxivity values in table 3.1 were determined from the quantifying the relative changes in signal intensity from MRI images. This relaxivity values do not necessarily represent actual theoretical values based on relaxation theory. This is because the MRI protocol used is designed primarily for making qualitative comparisons between images whose quality may change depending on machine sensitivity. One way to determine accurate values of relaxivity would be to employ standardized techniques like Nuclear magnetic resonance (NMR) spectroscopy. Surprisingly, we found that the relaxivity of Gd(III)-lipid monolayer-coated microbubbles increased significantly after destruction by sonication to form bilayer fragments. One explanation is that the presence of the microbubble gas core weakened the MRI signal intensity owing to susceptibility effects. However, the small difference in relaxation rates between intact and fragmented control microbubbles, and the lack of concentration dependence for these

samples (Fig. 3.7), does not support this explanation. We propose an alternative explanation for this phenomenon based on the difference in bulk water access to the Gd(III)-DOTA-DSPE complex for microbubbles versus liposomes (Fig. 3.8). Gd(III) is a paramagnetic ion that must interact with and exchange nearby water protons via its inner core (first hydration layer) in order to have a measurable effect on relaxivity (Toth, Helm et al. 2002). Intact microbubbles, which comprise a highly condensed monolayer shell held under compression by Laplace pressure-driven dissolution (Kim, Costello et al. 2003; Duncan and Needham 2004), may have restricted access of aqueous protons to the Gd(III) ion. Fragmentation of the microbubble converted the lipid to a more relaxed liposomal bilayer configuration, which may have allowed for greater access of water molecules to the Gd(III) complex, thus allowing a greater relaxation enhancement. The average area per lipid molecule for a fully compressed monolayer may be as low as  $0.32 \text{ nm}^2$  (Saad, Policova et al. 2009), which is 25 % less than that for a typical gel-phase bilayer of  $0.48 \text{ nm}^2$  (Lewis and Engelman 1983; Israelachvili 1992; Nagle and Tristram-Nagle 2000; Petrache, Dodd et al. 2000). We therefore propose that the tighter lipid packing in the monolayer configuration silences the relaxation effect by inhibiting water proton exchange between the Gd(III) complex and the bulk, whereas looser packing in the bilayer configuration provides sufficient exchange to significantly affect relaxation. This mechanism is supported by recent results for magnetite-bearing polymeric microbubbles, in which a rise in longitudinal and transversal relaxivities was found following microbubble destruction and release of the iron oxide particles (Liu, Lammers et al. 2011).



**Figure 3.8.** Cartoon showing proposed mechanism for MRI relaxivity increase with the cavitation-induced conversion of lipid from the compressed monolayer form on the intact microbubble to the relaxed bilayer form of the fragments. The lipid molecular area and hydrocarbon membrane thickness are estimated to be  $0.32 \text{ nm}^2$  and  $2.2 \text{ nm}$  for the condensed monolayer (Saad, Policova et al. 2009; Israelachvili 2011) and  $0.48 \text{ nm}^2$  and  $2.4 \text{ nm}$  for the relaxed, gel-state bilayer (Lewis and Engelman 1983; Nagle and Tristram-Nagle 2000; Petrache, Dodd et al. 2000; Israelachvili 2011). The parameters shown in the schematic are the outer hydration shell of the Gd(III) ion, OS; the inner hydration shell, IS; the molecular tumbling time,  $\tau_R$ ; and the proton exchange rate from the OS to the IS,  $k_{ex}$ . It is proposed that the ability of the Gd(III) ion to magnetize bulk water protons, i.e., the value of  $k_{ex}$ , increases as the lipid packing relaxes during the monolayer-to-bilayer transition.

Regardless of the underlying mechanism, this behavior of Gd(III)-bound microbubbles may have useful implications for MRI-guided FUS therapy. Using the Gd(III)-bound microbubbles fabricated here, one may envision that microbubble cavitation within the ultrasound focus can be spatially and temporally controlled *in situ* via monitoring of the MRI signal increase as the Gd(III)-DOTA-DSPE is converted from monolayer to bilayer. Cavitation detection during focused ultrasound surgery may serve as a method to guide and monitor therapeutic effects and prevent unwanted bioeffects (Farny, Holt et al. 2009; O'Reilly and Hynynen 2010; Hsieh, Smith et al. 2011). For example, Huang et al. (Huang, Xu et al. 2010) recently proposed to use phase-change agents, such as perfluorocarbon-liquid emulsion droplets that vaporize upon heating, to detect the margins of ablation during high-intensity focused ultrasound. Here, we propose an alternative strategy, in which Gd(III)-lipid microbubbles may serve as both a source and MRI beacon for acoustic cavitation. Following the proposed mechanism given above, the MRI signal would increase from baseline tissue contrast to positive contrast, in a dose-dependent manner, as microbubbles are fragmented. The change in signal intensity would provide a measure of the microbubble cavitation dose within the region of interest. Thus, the Gd(III)-microbubbles developed here may serve as a theranostic agent to monitor treatment and minimize the side effects associated with FUS.



### 3.5. Conclusions

Chelation of the paramagnetic lanthanide Gd(III) to the DOTA ligand on the surface of lipid-shelled microbubbles was achieved at a reaction temperature of 50 °C without degrading the 4-5  $\mu\text{m}$  microbubble size distribution. The microbubbles were echogenic and provided contrast during high-frequency ultrasound imaging *in vivo*. Surprisingly, MRI relaxometry showed that intact Gd(III)-bound microbubbles did not significantly enhance longitudinal or transverse proton relaxation. However, the bilayer fragments of Gd(III)-bound microbubbles formed by cavitation resulted in a significant increase  $r_1$  and  $r_2$ . A mechanism based on bulk water access to the Gd(III) complex was proposed to explain the increase in MRI signal intensity observed upon conversion of the condensed monolayer form to the relaxed bilayer form. Gd(III)-bound microbubbles could find use as cavitation probes for MRI-guided FUS therapy applications.

## **Chapter 4. Magnetic Resonance Properties of Gd(III)-Bound Lipid-Coated Microbubbles and their Cavitation Fragments**

### **4.1. Introduction**

Magnetic resonance imaging guided focused ultrasound surgery (MRIGFUS) is a rapidly developing medical technique that uses high intensity focused ultrasound to ablate tissue and magnetic resonance (MR) thermometry to monitor treatment (Cline, Schenck et al. 1992; Jolesz and Hynynen 2002; Jolesz and McDannold 2008). Formation and transient cavitation of gas-filled microbubbles nucleated from dissolved gases in tissue and blood plays an integral role in the efficacy of this therapy (Jolesz and McDannold 2008). However, the nucleation of cavitation sites is unpredictable and can lead to deleterious effects outside the targeted region. Thus, it is preferable to use preformed, stabilized microbubbles (1-10  $\mu\text{m}$  diameter) that can interact with ultrasound waves in a more predictable manner (Jolesz and McDannold 2008). Lipid-coated microbubbles are currently approved by the U.S. Food and Drug administration (FDA) for echocardiography and are currently being developed for expanded imaging capabilities (Lindner 2004; Qin, Caskey et al. 2009) and therapeutic applications in drug, gene and gas delivery (Ferrara, Pollard et al. 2007; Coussios and Roy 2008; Swanson, Mohan et al. 2010; Sirsi, Hernandez et al. 2012). One potential therapeutic application of microbubbles is the non-invasive, localized and transient opening of the blood-brain-barrier (BBB) for targeted drug delivery to the brain. Previously, lipid-coated microbubbles were shown to reduce the acoustic threshold needed for opening of the BBB *in vivo* (Choi, Feshitan et al. ; Hynynen, McDannold et al. 2001; Choi, Pernot et al. 2007; Xie, Boska et al. 2008; Marquet, Tung et al. 2011). In addition to disrupting vasculature, microbubbles also have been designed to deliver a therapeutic

payload (Ferrara, Pollard et al. 2007; Lentacker, De Smedt et al. 2009). However, methods are unavailable to use MRI for tracking microbubbles and their interactions with ultrasound. It therefore is desirable to develop MR-detectable microbubbles, so that MRI can be used to monitor and control not only thermal ablation, but also pharmaceutical delivery.

In chapter 3, the MRI-contrast agent Gd(III)-DOTA was conjugated to the lipid shell of size-selected gas-filled microbubbles using a post-labeling technique. Gd(III)-DOTA was conjugated to the primary amine on the headgroup of phosphatidylethanolamine (PE). The effect of Gd(III)-bound microbubble cavitation on the MR signal (at 9.4 T) was determined by comparing  $r_1$  and  $r_2^*$  of 4-5 micron-sized gas core-containing Gd(III)-bound microbubbles to those of microbubbles that were fragmented by inertial cavitation into lipid fragments. Surprisingly, both  $r_1$  and  $r_2^*$  increased after the fragmentation of intact Gd(III)-bound microbubbles. The explanation for this phenomenon was based on the difference in bulk water access to the lipid headgroup-labeled Gd(III)-complex for microbubbles versus the lipid fragments. Paramagnetic Gd(III) ions reportedly enhance MR signal by interacting with nearby water protons (Toth, Helm et al. 2002). Intact microbubbles, which comprise a highly condensed monolayer shell, may restrict access of aqueous protons to the Gd(III) bound lipid headgroup. Microbubble fragmentation converts the Gd(III) bound PE monolayer to a more relaxed liposomal bilayer structure, which may allow for greater access of water molecules to the Gd(III), increasing  $R_1$  and enabling higher signal intensities in  $T_1$ -weighted images.

One method of testing this hypothesis is to characterize the MR relaxation properties of intact and fragmented lipid microbubbles comprising a shell with Gd(III) preferentially conjugated to a long, flexible linker, such as the distal group on the polyethylene glycol (PEG)

brush on the lipid monolayer encapsulation. Microbubble (or liposome) design requires the PEG brush to provide stability against coalescence (or fusion) and to protect against an immunogenic response *in vivo* (Allen 1994; Ferrara, Pollard et al. 2007). Although, liposomes with Gd(III)-labeled lipid headgroups have been designed for cellular and vascular imaging applications (Strijkers, Mulder et al. 2005; Ghaghada, Ravoori et al. 2009; Hak, Sanders et al. 2009), neither liposomes nor microbubbles with Gd(III) conjugated to the PEG brush have been documented. Comparison between the MR signal intensity induced by PEG-labeled Gd(III)-microbubbles and their liposomal fragments may shed some light into the exchange rate hypothesis that was originally proposed to explain the behavior of lipid headgroup-labeled Gd(III)-microbubbles under fragmentation.

The microbubble gas core may also affect longitudinal and transverse relaxation. Previously, gas-filled microbubbles without Gd(III) reportedly enhanced  $R_2^*$  in a dose-dependent manner (Alexander, McCreery et al. 1996; Wong, Huang et al. 2004; Cheung, Chow et al. 2009).  $R_2^*$  is the additional enhancement of the transverse magnetization signal above of baseline  $R_2$  that stems from inhomogeneities in the magnetic field. The  $R_2^*$  enhancement produced by the microbubbles was attributed to the differences in magnetic susceptibility at the gas-liquid interface, which creates local inhomogeneities in the magnetic field that affect  $R_2^*$  decay. The equation given below relates the approximate magnetic field perturbation caused by a single, isolated spherical gas-filled microbubble at a position described by cylindrical coordinates  $(r, \theta)$  (Chu, Xu et al. 1990; Xu and Haacke 2006):

$$\Delta B(r, \theta) / B_o = (4/3)\pi\Delta\chi(R/r)^3(3\cos^2\theta - 1) \quad (4.1)$$

where  $\Delta B$  is the magnetic field perturbation,  $B_o$  is the static magnetic field vector,  $\Delta\chi$  represents the magnetic susceptibility difference between gas and liquid and  $R$  is the sphere radius. Thus, the degree of magnetic field perturbation caused by the presence of a single gas bubble during MR analysis is theoretically dependent on gas volume and the susceptibility differences at the gas-liquid interface. Alexander et al. (Alexander, McCreery et al. 1996) reported  $R_2^*$  enhancement from 9 different gas-types and the potential to use the gas-volume dependent susceptibility effect as a pressure sensor for evaluating cardiovascular function. Cheung et al. (Cheung, Chow et al. 2009) reported an increase of  $R_2^*$  at 7 T in the rat brain as a function of microbubble volume fraction for both sulfur hexafluoride and air microbubbles. Wong et al. (Wong, Huang et al. 2004) used the  $R_2^*$  enhancement induced by Optison (GE Healthcare) microbubbles for intravascular imaging of the rat liver with MRI. Subsequent studies by Chow *et al.* (Chow, Chan et al.) and Yang et al. (Yang, Li et al. 2009) demonstrated additional enhancement of  $R_2^*$  by loading iron oxide into the shell of polymeric microbubbles (Chow, Chan et al. ; Yang, Li et al. 2009). Finally, Liu et al. (Liu, Lammers et al. 2011) reported an additional enhancement of both  $R_1$  and  $R_2^*$  after microbubble fragmentation with ultrasound. The additional MR enhancement after fragmentation was attributed to greater interaction of water protons with iron oxide in the shell fragments. For lipid microbubbles loaded with paramagnetic Gd(III), it is desirable to determine the effects the paramagnetic ions and gas core will have on the MR signal, and ultimately what the implications are for applications in MRIgFUS.

In this report, the effects of microbubble gas core and Gd(III) conjugation on  $R_1$  and  $R_2^*$  were determined quantitatively by NMR. The microbubble shell used in this study is lipid-based because of the advantages in biocompatibility and ultrasound compliance when compared to polymeric agents (Bloch, Wan et al. 2004). A post-labeling protocol was used to generate lipid headgroup-labeled or PEG-labeled Gd(III)-microbubbles, and microbubbles without Gd(III) conjugated to the shell served as controls. The changes in  $r_1$  and  $r_2^*$  of both lipid headgroup-labeled and PEG-labeled Gd(III)-microbubbles after fragmentation were used to determine the potential mechanisms responsible for an increase in MR signal of fragmented lipid headgroup-labeled Gd(III) microbubbles.

## 4.2. Materials and Methods

### 4.2.1. Materials

All solutions were prepared using filtered, 18 M $\Omega$  cm deionized water (Direct-Q, Millipore, Billerica, MA). All glassware were cleaned with 70 vol% ethyl alcohol solution (Sigma-Aldrich; St. Louis, MO) and rinsed with deionized water. The gas used to form microbubbles was perfluorobutane (PFB) at 99 wt% purity obtained from FluoroMed (Round Rock, TX). 1,2-distearoyl-*sn*-glycero-3-phosphoethanolamine (DSPE) and 1,2-distearoyl-*sn*-glycero-3-phosphoethanolamine-N-[amino(polyethylene glycol)-2000] (DSPE-PEG(2000)-amine) were purchased from Avanti Polar Lipids (Alabaster, AL). 1,2-distearoyl-*sn*-glycero-3-phosphoethanolamine-N-[methoxy(polyethylene glycol)5000] (DSPE-PEG(5000)) was obtained from NOF America Corporation (White Plains, NY). 1,4,7,10-Tetraazacyclododecane-1,4,7,10-tetraacetic acid mono(*N*-hydroxysuccinimide ester) (DOTA-NHS) was purchased from

Macrocyclics (Dallas, TX) and dissolved in N,N-dimethylformamide (DMF; Sigma-Aldrich) prior to use. Gadolinium (III) chloride ( $\text{GdCl}_3$ ) was purchased from Sigma Aldrich and dissolved in 0.2 M, pH 5.6 acetate buffer (VWR, Radnor, PA).

#### 4.2.2. Microbubble Generation and Size Isolation

Microbubbles with primary amine lipid groups were formulated using a lipid suspension of 90% DSPE and 10% DSPE-PEG(5000) at 2 mg/mL in 100 mL PBS (pH 7.2, 0.15 M NaCl, 0.2 M phosphate). Microbubbles with primary amine PEG groups were formulated using a lipid suspension of 90% DSPC and 10% DSPE-PEG(2000)-amine at 2 mg/mL in 100 mL PBS. The lipid DSPC was selected to prevent headgroup conjugation. Each solution was degassed by applying house vacuum with constant stirring. The DSPE/DSPE-PEG(5000) solution was then preheated to 6 °C above the main phase transition temperature ( $T_m = 74$  °C) of DSPE (Cevc and Marsh 1985). Similarly, the DSPC/DSPE-PEG(2000)-amine solution was then preheated to 5 °C above the main phase transition temperature ( $T_m = 55$  °C) of DSPC.

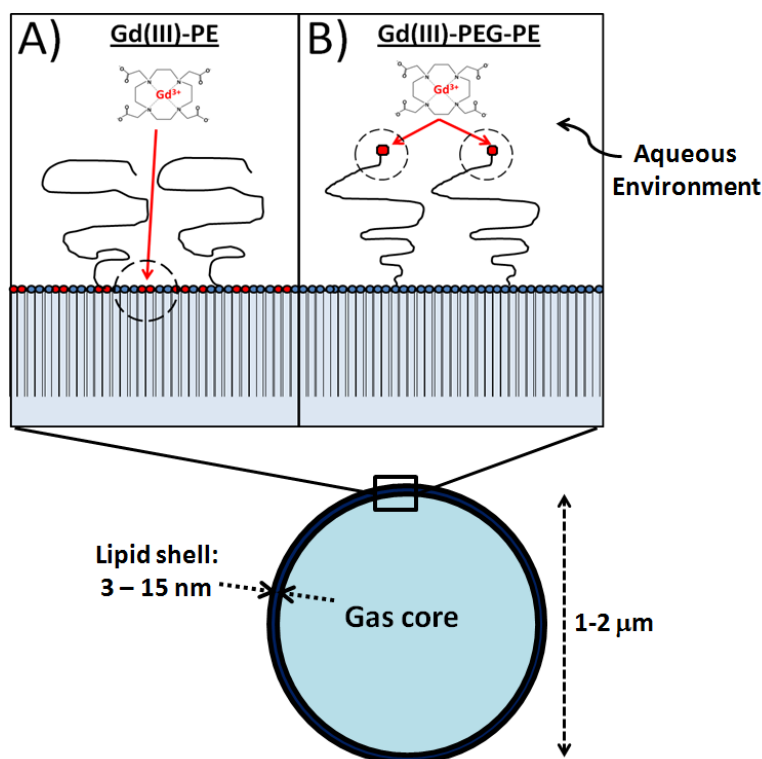
The lipid mixture was sonicated with a 20-kHz probe (model 250A, Branson Ultrasonics; Danbury, CT) at low power (3 W) to disperse the lipid aggregates into small, unilamellar liposomes. PFB gas was introduced by flowing it over the surface of the lipid suspension. Higher power sonication (33 W) was applied to the suspension for about 10 seconds at the gas-liquid interface to generate microbubbles. Size distribution of the microbubbles was determined by laser light obscuration and scattering (Accusizer 780A, NICOMP Particle Sizing Systems, Santa Barbara, CA). During measurements, 2  $\mu\text{L}$  samples of each microbubble suspension were

diluted into a flask containing 30 mL of distilled water under gentle mixing. All samples were measured in triplicate and analyzed for both number- and volume-weighted size distributions. The size-selected microbubbles used in this study were 1-2  $\mu\text{m}$  in size since this size range provided lower buoyancy and, consequently, improved measurement consistency during NMR data collection as compared to the 4-5  $\mu\text{m}$  bubbles. The 1-2  $\mu\text{m}$  microbubble size distribution was refined using methods described in chapter 2.

#### **4.2.3. Synthesis of Headgroup-labeled (Gd(III)-PE) and PEG-labeled (Gd(III)-PEG-PE) Microbubbles**

Gd(III)-bound microbubbles shells were fabricated using the post-labeling technique (Figure 4.1) described in chapter 3. To summarize briefly, during headgroup labeled Gd(III)-microbubble (Gd(III)-PE) formulation, the macrocyclic ligand DOTA was conjugated via an NHS reaction at room temperature to the primary amine located on the headgroup of DSPE. Excess DOTA-NHS ions were removed by several cycles of washing/centrifugation (5 minute, 300 RCF) using pH 5.6 Acetate buffer. The lanthanide Gd(III) was chelated to the DOTA group bound to the lipid headgroup. The chelation reaction was carried out at 50 °C for 2 hours. Excess Gd(III) ions were removed by several cycles of washing/centrifugation (5 minute, 300 RCF) using pH 5.6 Acetate buffer followed by several cycles of washing/centrifugation using pH 7.4 PBS.





**Figure 4.1: Schematic of Gd(III)-DOTA conjugation to microbubble shell. A) Gd(III)-PE. B) Gd(III)-PEG-PE.**

Similarly, PEG-labeled Gd(III) microbubbles (Gd(III)-PEG-PE) were fabricated using the same post-labeling technique. DOTA was conjugated to the primary amine group located on the distal end of DSPE-PEG(2000) amine. Excess DOTA-NHS ions were removed by several cycles of washing/centrifugation (5 minute, 300 RCF) using pH 5.6 Acetate buffer. The lanthanide Gd(III) was chelated to the DOTA group bound to the PEG-PE group. The chelation reaction was carried out at 45 °C for 2 hours. Excess Gd(III) ions were removed by several cycles of washing/centrifugation (5 minute, 300 RCF) using pH 5.6 Acetate buffer followed by several cycles of washing/centrifugation using pH 7.4 PBS.

#### 4.2.4. Characterization of Gd(III)-bound Microbubbles and Cavitation Fragments

Microbubble size distribution and concentration after chelation reaction were determined by using the Accusizer. The concentration of Gd(III) bound to the microbubble shell was determined by inductively coupled mass spectroscopy (ICP-MS, Agilent technologies, Santa Clara, CA). Destruction/fragmentation of Gd(III)-PE and Gd(III)-PEG-PE microbubble suspensions was accomplished by simultaneous bath sonication and heating to 80 °C and 60 °C for 10 mins respectively. A Malvern ZetaSizer (Worcestershire, United Kingdom) was used to determine the size of liposomal fragments.

#### 4.2.5. NMR Characterization of Intact and Fragmented Gd(III)-bound Microbubbles

The effect of the Gd(III)-bound microbubbles on the  $T_1$  and  $T_2^*$  relaxation times was determined using NMR. Intact Gd(III)-PE or Gd(III)-PEG-PE microbubbles suspension at a concentration of at least  $3 \times 10^{10} \text{ mL}^{-1}$  were mixed with saline in four different volume ratios (25, 50, 75 and 100 %) in a 2 mm outer diameter NMR compatible capillary tube (Wilmad-LabGlass, Vineland, NJ), which was flame-sealed. The average starting Gd(III) concentration of Gd(III)-PE and Gd(III)-PEG-PE samples was 0.45 mM and 0.1 mM respectively. Intact and fragmented 1-2  $\mu\text{m}$  microbubbles without Gd(III) binding were used as controls. End-over-end mixing, performed five times before each measurement, homogenized the microbubble suspensions and further mitigated the effects of bubbles rising due to their buoyancy. A 1.5 T vertical NMR spectroscopy system was used to measure  $T_1$  and  $T_2^*$  relaxation times.  $T_1$  relaxation times were determined using an inversion recovery sequence with recovery times ranging from as low as 1

ms up to 30 s.  $T_2^*$  relaxation times were determined using a Carr-Purcell-Meiboom-Gill (CPMG) sequence with a 4 ms intercho time, and a repetition time of at least 5  $T_1$  for each sample. The  $T_1$  and  $T_2^*$  times induced by the intact microbubbles were measured first. Next, the sealed capillary containing the intact microbubbles was heated and sonicated to produce a solution of Gd(III)-bound lipid fragments. Fragmentation of Gd(III)-PE and Gd(III)-PEG-PE microbubbles were accomplished by heating to 80 °C and 60 °C for 10 minutes respectively. Finally, the  $T_1$  and  $T_2^*$  times induced by the fragmented microbubbles were determined. All  $T_1$  and  $T_2^*$  measurements were repeated in triplicate. Molar relaxivities ( $r_1$  and  $r_2^*$ ) were calculated from slopes of regression lines of the plot of  $R_1$  and  $R_2^*$  versus Gd(III) concentration. A two-tailed unpaired student t-test was used to determine the significance between  $r_1$  and  $r_2^*$  of fragmented versus intact samples assuming a Gaussian distribution.

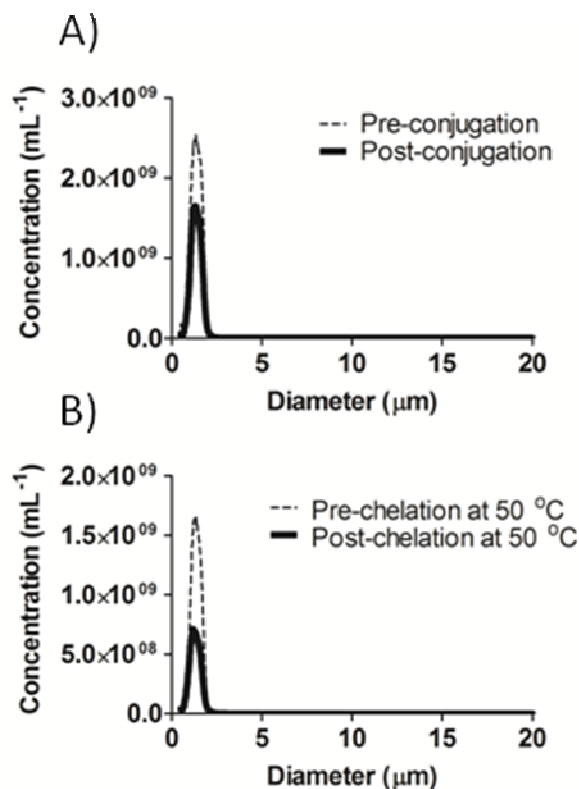
### **4.3. Results and Discussions**

#### **4.3.1. Preparation of Gd(III)-bound Microbubbles**

##### **4.3.1.1. Headgroup-labeled Microbubbles (Gd(III)-PE)**

The size isolation technique produced 1-2  $\mu\text{m}$  diameter microbubbles at a concentration of at least  $1 \times 10^{11} \text{ mL}^{-1}$ . Figure 4.2a shows the size distribution of microbubbles before and after DOTA conjugation to the amine-containing lipid (DSPE) on the monolayer via an NHS-coupling reaction at room temperature. The average microbubble concentration decreased by approximately 41 % after DOTA conjugation. Figure 4.2b shows the size distribution of microbubbles before and after Gd(III) chelation to DOTA bound to the DSPE monolayer at 50 °C. The average microbubble concentration decreased by 54 % after the 2-hour chelation

reaction. The mean and median number diameter of microbubbles changed by less than 1% before and after DOTA conjugation and Gd(III) chelation. This indicated that the post-labeling strategy to available PE groups reduced microbubble concentration but minimally affected 1-2 micron size monodispersity. The average Gd(III) ions per Gd(III)-PE microbubble determined by ICP-MS was  $7.3 \times 10^6 \pm 2.3 \times 10^6$  ions/microbubble ( $1.3 \times 10^6 \pm 4.1 \times 10^5$  ions/ $\mu\text{m}^2$ ). Assuming spherical microbubbles with an average molecular area of  $0.4 \text{ nm}^2$ , there was on average ~64 % of Gd(III)-DOTA binding to available PE groups. Dynamic light scattering measurements indicated that fragmentation of the Gd(III)-PE microbubbles by heating and bath sonication resulted in liposomes of size  $217 \pm 3 \text{ nm}$  in diameter.

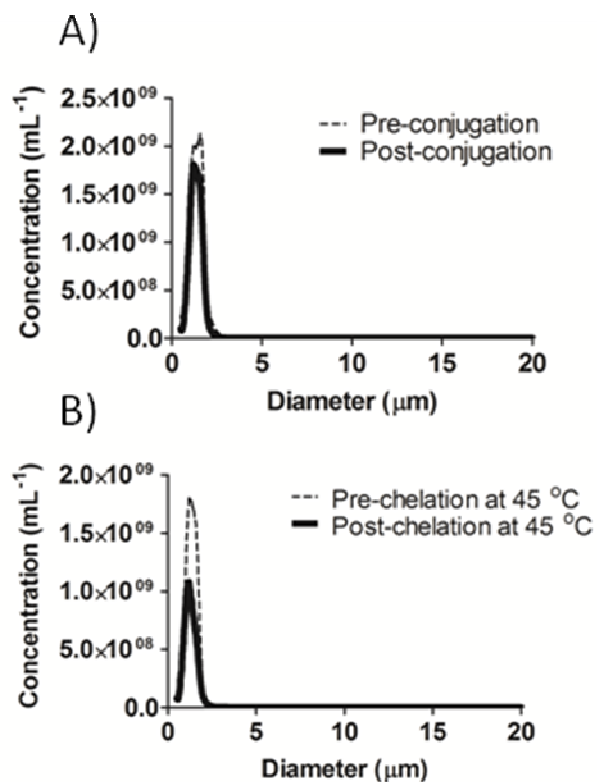


**Figure 4.2. A) Number-weighted size distributions of DSPE microbubbles before and after conjugation with DOTA-NHS. B) Number-weighted size distributions of lipid-bound DOTA-microbubbles before and after Gd(III) chelation at 50 °C.**

#### **4.3.1.2. PEG-labeled Microbubbles (Gd(III)-PEG-PE)**

As before, the size isolation technique produced 1-2 μm diameter microbubbles at a concentration of at least  $1 \times 10^{11} \text{ mL}^{-1}$ . Figure 4.3a shows the size distribution of microbubbles before and after DOTA conjugation to the distal end of the amine-containing PEG-PE on the monolayer via an NHS-coupling reaction at room temperature. The average concentration of 1-2 μm microbubbles decreased by approximately 27 % after DOTA-conjugation. Figure 4.3b

shows the size distribution of microbubbles before and after Gd(III) chelation to the DOTA-bound PEG-PE on the monolayer at 50 °C. The average microbubble concentration decreased by 41 % after the 2-hour chelation. The mean and median number diameter of microbubbles changed by less than 1 % before and after DOTA conjugation and Gd(III) chelation. This indicated that the post-labeling strategy to available PEG-PE groups reduced microbubble concentration but minimally affected 1-2 micron size monodispersity. The average Gd(III) ions per Gd(III)-PEG-PE microbubble determined by ICP-MS was  $6.3 \times 10^5 \pm 2.1 \times 10^5$  ions/microbubble ( $1.4 \times 10^5 \pm 6.7 \times 10^4$  ions/ $\mu\text{m}^2$ ). Again, assuming spherical microbubbles with an average molecular area of  $0.4 \text{ nm}^2$ , there was on average ~53 % of Gd(III) binding to available reactive PEG-PE groups. Dynamic light scattering measurements indicated that fragmentation of Gd(III)-PEG-PE microbubbles by heating and bath sonication resulted in bimodal lipid fragments of  $84 \pm 18 \text{ nm}$  and  $340 \pm 61 \text{ nm}$  in diameter. This bimodality suggests that Gd(III)-PEG-PE microbubbles formed a mixture of lipid vesicles such as micelles and liposomes after fragmentation.



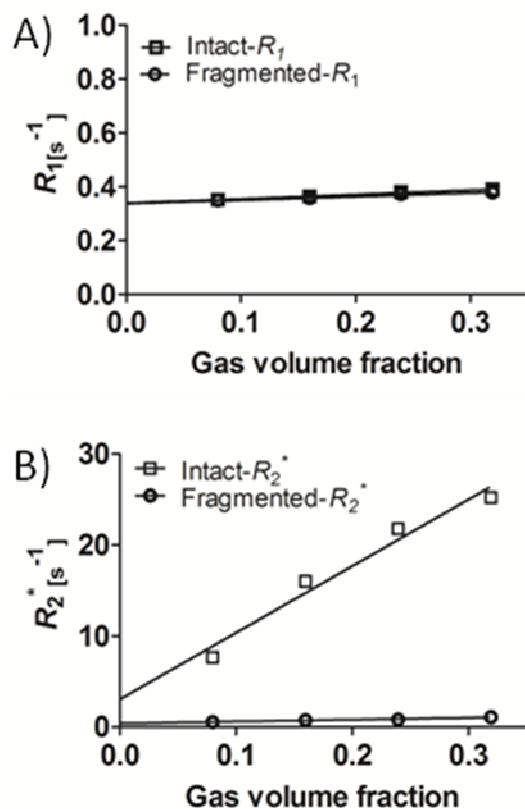
**Figure 4.3.** A) Number-weighted size distributions of DSPE/DSPE-PEG 2000 amine microbubbles before and after conjugation with DOTA-NHS. B) Number-weighted size distributions of PEG-bound DOTA-microbubbles before and after Gd(III) chelation at 45 °C.

#### 4.3.2. NMR Characterization of Intact and Fragmented Gd(III)-Bound Microbubbles

##### 4.3.2.1. Control Microbubbles (No Gd(III))

Figure 4.4a and 4.4b show plots of the longitudinal ( $R_1 = 1/T_1$ ) and transverse ( $R_2^* = 1/T_2^*$ ) relaxation rates of the control microbubble versus gas volume fraction.  $r_2^*$  decreased from  $73 \pm$

9.8 to  $2.1 \pm 0.31 \text{ s}^{-1}$  volume-fraction<sup>-1</sup> during microbubble fragmentation. There was no appreciable change in  $r_l$ .



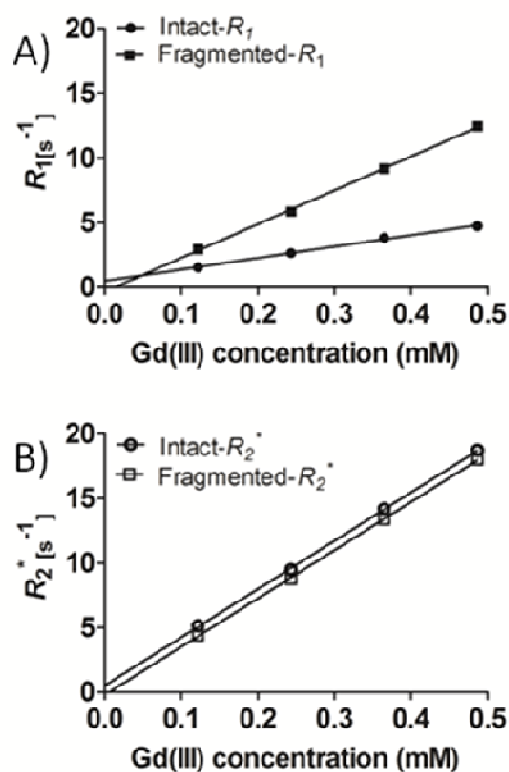
**Figure 4.4: Relaxation rate of control microbubbles (No Gd(III)) versus gas volume fraction. A) Longitudinal relaxation rate ( $R_1$ ). B) Transverse relaxation rate ( $R_2^*$ ).**

#### 4.3.2.2. Relaxation Rate of Intact and Fragmented Gd(III)-PE Microbubbles

Figure 4.5 shows representative plots of  $R_l$  and  $R_2^*$  versus Gd(III) concentration for Gd(III)-PE microbubbles.  $R_l$  and  $R_2^*$  increased in a dose-dependent manner for both the intact and fragmented microbubbles. Table 4.1 summarizes the calculated  $r_l$  and  $r_2^*$  values of the intact



and fragmented Gd(III)-PE microbubbles for 3 independent experiments. On average the value of  $r_1$  increased by a factor of 2.4 after microbubble fragmentation ( $p < 0.05$ ). However,  $r_2^*$  remained relatively unchanged after microbubble fragmentation ( $p = 0.9$ ). Therefore, fragmentation of Gd(III)-PE microbubbles enhanced  $R_1$ , but only slightly changed  $R_2^*$ .



**Figure 4.5: Relaxation rates of intact and fragmented Gd(III)-PE microbubbles versus Gd(III) concentration. A) Longitudinal relaxation rate ( $R_1$ ). B) Transverse relaxation rate ( $R_2^*$ ). Figure shown is a representative data set.**

**Table 4.1: Relaxivities ( $r_1$  and  $r_2^*$ ) of intact and fragmented Gd(III)-PE microbubbles.**

Gd(III)-PE	Gd(III) ions/ Microbubble	Intact	Fragmented	Intact	Fragmented
		$r_1$ (mM <sup>-1</sup> s <sup>-1</sup> )	$r_1$ (mM <sup>-1</sup> s <sup>-1</sup> )	$r_2^*$ (mM <sup>-1</sup> s <sup>-1</sup> )	$r_2^*$ (mM <sup>-1</sup> s <sup>-1</sup> )
Trial 1	$8.6 \times 10^6$	$6.6 \pm 0.19$	$16 \pm 0.39$	$23 \pm 1.4$	$22 \pm 2.5$
Trial 2	$8.7 \times 10^6$	$8.8 \pm 0.36$	$26 \pm 0.55$	$37 \pm 0.40$	$37 \pm 0.33$
Trial 3	$4.7 \times 10^6$	$7.3 \pm 8.1$	$14 \pm 13$	$43 \pm 4.4$	$41 \pm 1.6$

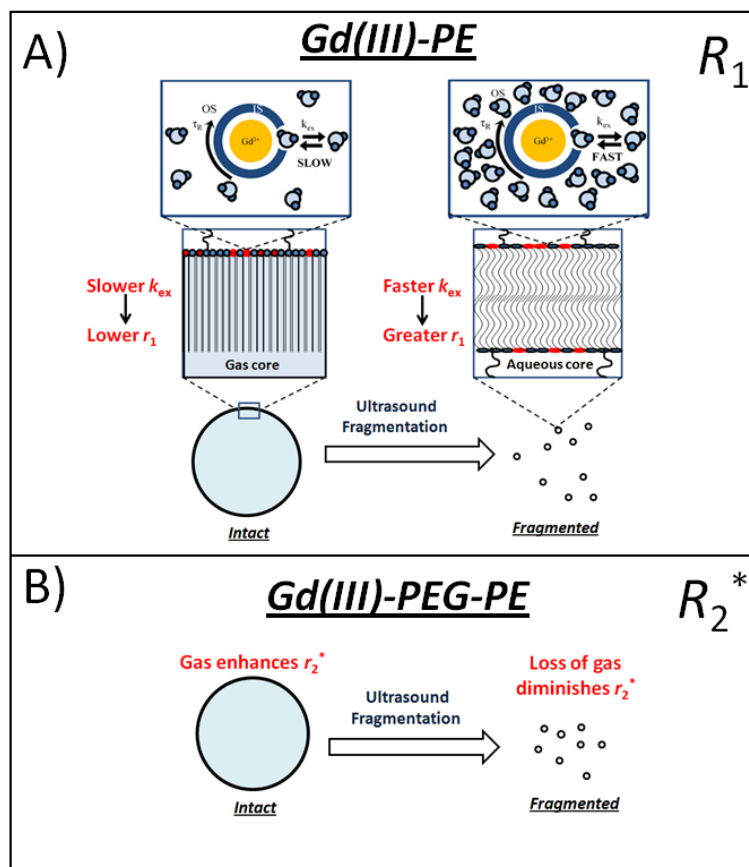
Similar to the  $T_1$ -weighted MRI analysis of 4-5 micron sized Gd(III)-PE microbubbles in chapter 3, fragmentation of the 1-2 micron sized Gd(III)-PE microbubbles resulted in an increase in  $r_1$ . From the control microbubble experiments, it was evident that there is virtually no contribution to  $T_1$  relaxation from either the lipid bilayer or the perfluorobutane gas core. These results support the hypothesis in chapter 3 of an enhanced exchange rate of protons with Gd(III) in the relaxed bilayer versus the condensed monolayer.

We suggest that the enhancement in proton exchange occurred as the lipid headgroup area increased with the transformation from a monolayer to bilayer configuration (figure 4.6A). Israelachvili provided a simple relation for the optimal headgroup area ( $a_o$ ) in terms of a balance between the attractive and repulsive intermolecular forces between lipid molecules (Israelachvili 2011):

$$a_o = \sqrt{\frac{K}{\gamma}} \quad (4.2)$$

where  $K$  is a constant accounting for steric repulsion in the interfacial region and  $\gamma$  is the interfacial tension. Based on this simple approximation, the equilibrium headgroup area is expected to increase ~70% upon going from a gas-water interface ( $\gamma_{monolayer} \sim 73$  mN/m) to a bilayer membrane ( $\gamma_{bilayer} \sim 25$  mN/m), assuming that repulsive forces remain unchanged. Such

an increase in headgroup area may allow more water molecules to permeate into the headgroup region and interact with the chelated Gd(III) ions.



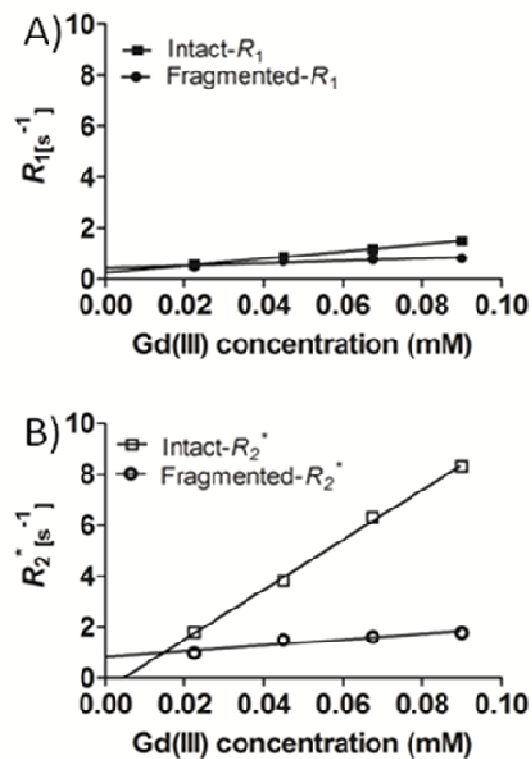
**Figure 4.6. Schematic demonstrating hypothetical mechanisms governing changes in a)  $r_1$  of Gd(III)-PE microbubbles and b)  $r_2^*$  of Gd(III)-PEG-PE microbubbles.**

The trend in  $T_2^*$  is different (Figure 4.5B, table 4.1) than  $T_1$ , as  $r_2^*$  of the intact and fragmented Gd(III)-PE microbubbles were similar in magnitude, especially for the first two trials. From observation of the control microbubbles without Gd(III), it is apparent that the intact microbubbles provide a ~35-fold greater increase of  $R_2^*$  than the lipid fragments. This result indicates that the gas core of the Gd(III)-PE microbubbles induces an independent enhancement of the  $R_2^*$ . This result also corroborates previous reports of an improvement in  $R_2^*$  using gas-

filled microbubbles (without Gd(III))(Chow, Chan et al. ; Alexander, McCreery et al. 1996; Wong, Huang et al. 2004; Cheung, Chow et al. 2009; Yang, Li et al. 2009; Liu, Lammers et al. 2011). Our results demonstrate that this gas susceptibility behavior affects  $R_2^*$ , but not  $R_1$ .  $R_1$  is only sensitive to relaxation processes at the proton Larmor frequency (64 MHz at 1.5 T) (Abragam 1961; Slichter 1990; Levitt 2008), while the  $R_2^*$  signal is sensitive to inhomogeneities in the magnetic field such as those present at gas-liquid interfaces. Therefore, we propose that the increase in  $R_2^*$  due to intact Gd(III)-PE microbubbles stems from a combination of a weak augmentation induced by Gd(III) in the highly-packed lipid monolayer and a susceptibility enhancement from the gas core. After microbubble fragmentation, the gas-susceptibility enhancement dissipates, but the remnant lipid fragments (containing no gas) still enhance  $R_2^*$  by virtue of enhanced proton exchange with Gd(III) in the relaxed lipid bilayer. Observation of the sample from trial 3, further supports this explanation, as a slightly higher gas-fraction to Gd(III) ratio, ultimately resulted in intact Gd(III)-PE microbubbles with a slightly larger  $r_2^*$  value than the lipid fragments. It may be worthwhile to explore the effects of changing the gas composition of the Gd(III)-PE microbubbles on  $r_2^*$ . The magnetic susceptibility constant of gases varies depending on whether the gas is diamagnetic or paramagnetic, and although the magnetic susceptibility of PFB gas has not been reported, the values for diamagnetic gases like Nitrogen and paramagnetic gases like Oxygen are known (Schenck 1996).

#### 4.3.2.3. Relaxation Rate of Intact and Fragmented Gd(III)-PEG-PE Microbubbles

Figure 4.7 shows representative plots of  $R_1$  and  $R_2^*$  versus Gd(III) concentration for Gd(III)-PEG-PE microbubbles. Again,  $R_1$  and  $R_2^*$  increased in a dose-dependent manner for both the intact and fragmented microbubbles. However, the trendlines changed significantly upon fragmentation in comparison to Gd(III)-PE microbubbles. Table 4.2 summarizes the calculated values of  $r_1$  and  $r_2^*$  of the intact and fragmented Gd(III)-PEG-PE microbubbles for 3 independent experiments. The calculated  $r_1$  values of the Gd(III)-PEG-PE microbubbles decreased on average by a factor of 2.1 after microbubble fragmentation ( $p = 0.26$ ). Similarly,  $r_2^*$  decreased on average by a factor of 8 after fragmentation ( $p < 0.05$ ). Moreover, the decrease in magnitude of  $r_2^*$  was about 2 times greater than that of  $r_1$ . Evidently, fragmentation of Gd(III)-PEG-PE microbubbles resulted in an overall weakening of both  $R_1$  and  $R_2^*$  enhancement.



**Figure 4.7: Relaxation rates of intact and fragmented Gd(III)-PEG-PE microbubbles versus Gd(III) concentration. A) Longitudinal relaxation rate ( $R_1$ ). B) Transverse relaxation rate ( $R_2^*$ ). Figure shown is a representative data set.**

**Table 4.2: Relaxivities ( $r_1$  and  $r_2^*$ ) of intact and fragmented Gd(III)-PEG-PE microbubbles.**

Gd(III)-PEG-PE	Gd(III) ions/ Microbubble	Intact $r_1$ (mM <sup>-1</sup> s <sup>-1</sup> )	Fragmented $r_1$ (mM <sup>-1</sup> s <sup>-1</sup> )	Intact $r_2^*$ (mM <sup>-1</sup> s <sup>-1</sup> )	Fragmented $r_2^*$ (mM <sup>-1</sup> s <sup>-1</sup> )
Trial 1	$5.0 \times 10^5$	$18 \pm 0.030$	$14 \pm 0.10$	$190 \pm 1.3$	$36 \pm 1.2$
Trial 2	$8.7 \times 10^5$	$14 \pm 0.33$	$4.6 \pm 1.3$	$98 \pm 2.9$	$11 \pm 2.9$
Trial 3	$5.1 \times 10^5$	$7.8 \pm 9.1$	$3.7 \pm 6.3$	$210 \pm 25$	$21 \pm 3.3$

We sought to further explore the proton exchange rate mechanism, proposed earlier as the reason for the increase in  $r_1$  for Gd(III)-PE microbubbles, by characterizing the MR signal induced from a construct consisting of Gd(III) bound to the PEG brush of the lipid microbubble. PEG is a highly flexible and dynamic molecule under these conditions (Kuhl, Leckband et al. 1994; Kuhl, Majewski et al. 1999; Moore and Kuhl 2006), and therefore the relaxation rate of terminal Gd(III) groups is expected to be relatively insensitive to changes in packing of the underlying lipids. Contrary to Gd(III)-PE microbubble behavior, both the  $r_1$  and  $r_2^*$  relaxivity of the Gd(III)-PEG-PE microbubbles dropped in magnitude after microbubble fragmentation.

The decrease in  $r_1$  suggests that the underlying mechanism governing changes in  $r_1$  after fragmentation of the Gd(III)-PEG-PE microbubbles is different from that governing the changes in  $r_1$  of Gd(III)-PE microbubbles. With Gd(III) bound to the PEG, the proton exchange mechanism proposed to govern changes in relaxivity of the Gd(III)-PE microbubbles would be expected to play a minimal role in the Gd(III)-PEG-PE microbubbles since the effective PEG molecular area and water access to Gd(III) would remain the same after microbubble fragmentation. One possible explanation for the decrease in  $r_1$  of Gd(III)-PEG-PE microbubbles is that differences in rotational correlation times between the intact and fragmented complexes dominates  $r_1$  behavior. Relaxation theory predicts that optimization of one or more parameters governing the inner sphere relaxivity of Gd(III) such as hydration number, electronic relaxation, proton exchange rate and rotational correlation time, makes the other parameters more limitative (Toth, Helm et al. 2002). Thus, assuming that the proton exchange mechanism and all other relaxation parameters remains unchanged, then the differences in rotational correlation times or between the microbubbles and lipid fragments could dictate the differences in relaxivity from fragmentation of Gd(III)-PEG-PE microbubbles. The microbubble as the larger sized complex

possibly tumbles at a slower rate, leading to larger  $r_1$  as compared to the smaller-sized lipid fragments that tumble at a faster rate. However, further experimentation will be needed to validate this hypothesis.

The changes in  $r_2^*$  can be attributed primarily to the strong effect the gas creates on the transverse magnetization signal. As explained above for the Gd(III)-PE microbubbles, the gas core of the intact Gd(III)-PEG-PE microbubble induces an independent enhancement of  $R_2^*$ . This is supported by the observation that fragmentation of the Gd(III)-PEG-PE microbubbles correlated with a significant decrease of  $R_2^*$ . However, unlike the Gd(III)-PE sample, the signal loss from microbubble fragmentation was not compensated by the presence of the remnant Gd(III)-PEG-PE lipid fragments. Overall, the lowering of  $r_1$  and  $r_2^*$  values from fragmenting Gd(III)-PEG-PE microbubbles seems to support the hypothesis proposed for the changes in  $r_1$  in the Gd(III)-PE microbubbles. Ultimately, it appears the effect of Gd(III)-bound microbubble cavitation on the MR signal depends on the placement of the Gd(III) (lipid labeled versus PEG labeled) on the shell and the presence of the gas core (for  $T_2^*$ ). For example, one may monitor FUS treatment by monitoring cavitation of lipid headgroup-labeled Gd(III)-microbubbles via changes in  $T_1$ , or via changes in  $T_2^*$  in the case of PEG-labeled Gd(III)-microbubbles.

#### 4.4. Conclusions

The paramagnetic MRI contrast agent, Gd(III)-DOTA, was conjugated to two different groups on the lipid monolayer shell of the microbubble: the PE lipid headgroup region or the distal region of the PEG brush. NMR analysis revealed that the microbubble gas core



specifically enhanced  $R_2^*$  in a dose-dependent manner, but not  $R_1$ . The increase in  $r_1$  after fragmentation of Gd(III)-PE microbubbles suggested that the corresponding expansion of the lipid headgroup area increased proton exchange with the Gd(III)-bound lipid headgroup and enhanced  $R_1$ . However, the decrease in  $r_1$  from fragmenting Gd(III)-PEG-PE microbubbles indicated that a different relaxation mechanism was at play, which was minimally affected by changes in lipid headgroup area. Therefore, the effect of Gd(III)-bound microbubble cavitation on the NMR signal depends on the location of Gd(III) on the lipid shell (PEG-labeled or lipid headgroup-labeled) and the presence of the gas core. These results show how NMR may be used to provide information on lipid headgroup packing, and they may hold implications for detecting and monitoring microbubble-assisted MRIGFUS.

## **Chapter 5. Opening the Blood-Brain Barrier with Gd(III)-bound Lipid-stabilized Microbubbles**

### **5.1. Introduction**

One main bottleneck for the treatment of diseases found in the central nervous system is lack of effective delivery of larger drug compounds to the brain. This restriction is primarily imposed by action of the brain's semi-permeable microvasculature known as the blood-brain barrier (BBB). The BBB generally excludes compounds greater than 400 Da, and thereby prevents neuropharmaceutical agents, such as inhibitors and antibodies, from reaching their desired target (Pardridge 2005). Focused ultrasound (FUS) with an intravenous injection of lipid-stabilized, gas-filled microbubbles has been demonstrated to induce the transient non-invasive and localized opening of the BBB (Hynynen, McDannold et al. 2005; Choi, Pernot et al. 2007). Subsequent studies have demonstrated a dependence of the acoustic pressure threshold needed to induce BBB opening on the size range of the microbubbles (Choi, Feshitan et al. 2010; Tung, Vlachos et al. 2011; Vlachos, Tung et al. 2011). Microbubbles of size class 4-5 and 6-8 micron in diameter induced BBB opening at a peak rarefactional driving pressure of 0.45 MPa, while 1-2 micron sized microbubbles required a driving pressure of 0.6 MPa. Additional studies demonstrated the feasibility to open the BBB in non-human primates (Marquet, Tung et al. 2011).

For these applications, it would be advantageous to use a microbubble formulation detectable by magnetic resonance imaging (MRI), which could be used to measure microbubble concentration, image cavitation events and determine the biodistribution of microbubble shell debris during or after FUS. For example, MRI-detectable microbubbles may be useful to

monitor FUS treatment, similar to how changes in temperature are used to monitor treatment in real-time during MRI-guided FUS therapy applications (Cline, Schenck et al. 1992; Cline, Hynynen et al. 1995; Hynynen, Freund et al. 1996). Alternatively, MRI-detectable microbubbles could be used to track the deposition profile of microbubble shell material (a potential surrogate for an attached drug) by scanning the sonicated region after FUS treatment.

In chapter 3 of this thesis, an MRI-detectable microbubble formulation was developed method by labeling the MRI-contrast agent Gd(III)-DOTA to the lipid shell of size-selected gas-filled microbubbles. Gd(III)-DOTA was conjugated to the primary amine on the headgroup of phosphatidylethanolamine (PE). The effect of Gd(III)-bound microbubble cavitation on the MR signal (at 9.4 T) was determined *in vitro* by comparing both longitudinal ( $r_1$ ) and transverse ( $r_2$ ) relaxivity of 4-5 micron-sized gas core-containing Gd(III)-bound microbubbles to those of microbubbles that were fragmented into lipid shell fragments. Analysis revealed that both  $r_1$  and  $r_2$  increased after the fragmentation of intact Gd(III)-bound microbubbles. Additionally, the Gd(III)-bound microbubbles were also found to be echogenic *in vivo* during high-frequency ultrasound imaging of the mouse kidney. The echogenicity indicated that conjugation of the MRI contrast agent to the microbubble lipid shell did not affect its behavior as an ultrasound contrast agent.

In another study, Liao et al (Liao, Liu et al. 2012) developed albumin microbubbles containing a shell surface-conjugated with Gd(III)-DTPA to serve as dual mode contrast agents for ultrasound and MR imaging. FUS in combination with these Gd(III)-labeled albumin microbubbles were used to disrupt the BBB and induce intracerebral hemorrhaging at a pressure range of 0.98 to 1.35 MPa. Additionally, the capability to image the regions of intracerebral

hemorrhaging was demonstrated using MRI in combination with the Gd(III)-labeled microbubbles. However, it is desirable to open the BBB at a safer pressure threshold of 0.45 MPa using the Gd(III)-labeled microbubbles, and then image the sonicated region of the brain with MRI for the presence of Gd(III)-labeled lipid vesicles. Furthermore, lipid-stabilized microbubbles are preferred for this purpose due to their higher ultrasound compliance in comparison to the albumin-shelled counterparts (Sirsi and Borden 2009).

In this chapter, we detail the opening of the BBB using FUS at 1.5 MHz and 0.4 MPa and 4-5 micron lipid-stabilized microbubbles labeled with Gd(III)-DOTA, and the feasibility to use MRI to image the sonicated region for presence of Gd(III)-labeled lipid vesicles after FUS treatment. The microbubbles in this study were surface-conjugated with Gd(III)-DOTA since the ligand DOTA is known to provide a stronger chemical bond with Gd(III) than DTPA, which reduces the potential for the release of free toxic Gd(III) ions (Sherry, Caravan et al. 2009).

## **5.2. Materials and methods**

### **Sonication Protocol**

All animal procedures were approved by the Columbia University Institutional Animal Care and Use Committee. The sonication protocol was adapted from the previously reported techniques (Choi, Pernot et al. 2007; Choi, Feshitan et al. 2010; Vlachos, Tung et al. 2011). A FUS transducer (frequency: 1.5 MHz) was confocally aligned with a pulse-echodiagnostic transducer (frequency: 7.5 MHz). A water-filled cone was mounted onto the transducer and attached to a positioning system (Velmex Inc., Lachine, QC, CAN). The FUS transducer was connected to a

matching circuit and driven by a function generator (Agilent, Palo Alto, CA) and a 50-dB power amplifier (ENI Inc., Rochester, NY). The diagnostic transducer was driven by a pulser-receiver system (Panametrics, Waltham, MA) connected to a digitizer (Gage Applied Technologies, Inc., Lachine, QC, CAN). Pressure measurements of the FUS transducer were made with a needle hydrophone in a water tank. Each mouse ( $n = 2$ , Strain C57BL/6 sex: male) was anesthetized with isoflurane, placed prone, and immobilized by a stereotaxic apparatus. The mouse hair was removed, ultrasound-coupling gel was applied on the skin, and a water-filled container sealed at the bottom was placed on the head. The transducers were submerged in the water and their foci were positioned to overlap with the right hippocampus of the brain. The left hippocampus was not targeted and was used as the control. The 4-5 micron lipid-stabilized microbubbles labeled with Gd(III)-DOTA were fabricated using the post-labeling technique previously described in chapter 3 and injected intravenously (IV) through the tail vein at a concentration of  $10^7 \text{ mL}^{-1}$  using a 30 G needle. The right hippocampal region of the brain was sonicated 1 min after the microbubble injection using pulsed FUS (burst rate: 10 Hz; burst cycles: 100; duty cycle: 0.067%; frequency: 1.5 MHz) at a peak-rarefactional pressure of 0.45 MPa for a duration of 60 s. Previous studies (Choi, Feshitan et al. 2010; Tung, Vlachos et al. 2011) have shown that the safe acoustic pressure for microbubble-mediated BBB opening lies between 0.30 and 0.45 MPa.

### **MRI protocol**

All the mice were imaged in a vertical bore 9.4 T MRI system (DRX400, Bruker Biospin, Billerica, MA) following previously developed protocols (Choi, Feshitan et al. 2010; Tung, Vlachos et al. 2011; Vlachos, Tung et al. 2011). Isoflurane gas (1–2%) was used to anesthetize

the mouse during the MRI procedure. A two-dimensional FLASH  $T_1$ -weighted sequence (TR/TE = 230/3.3 ms; flip angle:  $70^\circ$ ; NEX = 18; scan time: 9 min 56 s, matrix size: 256 x 192; spatial resolution:  $86 \times 86 \mu\text{m}^2$ ; slice thickness: 500  $\mu\text{m}$  with no interslice gap) was utilized to image the sonicated hippocampus (Vlachos, Tung et al. 2010; Vlachos, Tung et al. 2011). The scanning sequence was initially applied ~30 minutes after sonication of the right hippocampus in the presence of the Gd(III)-bound microbubbles. The presence of Gd(III) on the lipid shell of microbubbles was qualitatively confirmed by MR imaging of a centrifuge tube containing a 100  $\mu\text{L}$  aliquot of Gd(III) microbubble that were fragmented into constitutive lipid fragments using a combination of ultrasound and heating (data not shown). A second scan was applied ~50 minutes (~90 mins after sonication) after intraperitoneal injection of 0.30 ml BBB-impermeable, MRI contrast agent, Omniscan (GE Healthcare, Princeton, NJ), which allowed for sufficient time for Omniscan to diffuse into the sonicated hippocampus (Choi, Pernot et al. 2007; Choi, Feshitan et al. 2010). Omniscan (Gd(III)-DTPA-based) is used to confirm BBB opening because it reduces the longitudinal relaxation time ( $T_1$ ) when released in the extravascular extracellular space, thus augmenting the local  $T_1$ -weighted contrast.

### 5.3. Results

Figure 5.1B shows the image of the mouse brain after sonication of the right hippocampus in the presence of the Gd(III)-labeled microbubbles. As is evident from the image, there was no observable increase in  $T_1$ -weighted contrast in the right (sonicated) hippocampus after FUS sonication in presence of the Gd(III)-labeled microbubbles. Figure 5.1C shows the image of the mouse brain after sonication of the right hippocampus in the presence of the Gd(III)-labeled

microbubbles followed by intraperitoneal injection of Omniscan. As is evident from the image, there was a significant increase in  $T_1$ -contrast in the sonicated region, which provided confirmation that the BBB was opened. There was no observable increase in contrast in the control hippocampus in all cases.



**Figure 5.1. MRI images of the mouse brain: a) Unperturbed. b) ~40 mins after sonication of the right hippocampus in the presence of Gd(III)-labeled microbubbles. c) At least 100 mins after sonication in the presence of Gd(III)-labeled microbubbles and ~60 mins after injection of Omniscan.**

#### **5.4. Discussion**

FUS in combination with 4-5 micron sized Gd(III)-DOTA-labeled microbubbles was used to open the BBB in the right hippocampus of the mouse brain at a pressure threshold of 0.45 MPa. The BBB was confirmed to be opened because of the  $T_1$ -contrast enhancement in the sonicated

region after the injection of Omniscan. This enhancement is attributed to stem from combination of the intraperitoneal administration of Omniscan and time difference (50 mins) between its injection and the MRI scan. The intraperitoneal injection method is known to allow for the slow uptake and diffusion of the contrast media to the sonicated area, which improves the probability that it is still present at the region of interest during MRI scanning (Choi, Pernot et al. 2007). This result indicated that the conjugation of the MRI contrast agent to the lipid shell minimally affected its capability to open the BBB at the safe pressure threshold used in previous studies (Choi, Feshitan et al. 2010; Tung, Vlachos et al. 2011; Vlachos, Tung et al. 2011).

MRI was used to image the sonicated region for the presence of Gd(III)-labeled vesicles and to confirm BBB opening. In chapter 3 and 4 of this thesis, the fragmentation of the lipid headgroup labeled Gd(III) microbubbles into constitutive lipid vesicles resulted in an increase in longitudinal relaxivity ( $r_1$ ) *in vitro*. As a result, the presence of Gd(III)-bound lipid shell fragments at the sonicated hippocampus was expected to produce an increase in the local  $T_1$ -weighted MRI contrast; however, there was no observable increase in  $T_1$  signal. One hypothesis for the lack of  $T_1$ -contrast enhancement *in vivo* may be due to lengthy time difference (~30 mins) between sonication in the presence of the intravenously administered Gd(III)-labeled microbubbles and commencement of MRI scanning. This time difference possibly resulted in clearance of the Gd(III)-labeled vesicles from the sonicated regions by various mechanisms such as blood flow. Alternatively, the concentration of the Gd(III) lipid shell that may have been present at the scanning region may not have been high enough to produce a noticeable increase in  $T_1$ -contrast above the baseline level.

One potential method to improve the chances of imaging the presence of Gd(III)-labeled lipid vesicles at the region of interest is to minimize the timing between FUS application and



imaging using real-time MRI guidance similar to that adapted in MRI-guided FUS therapy devices (Cline, Schenck et al. 1992; Cline, Hynynen et al. 1995; Hynynen, Freund et al. 1996; Jolesz and McDannold 2008). Alternatively, the Gd(III)-labeled microbubble shell may be designed to incorporate a targeting peptide that can specifically bind to the BBB receptor and reduce the possibility of clearance from the region of interest. These and other considerations may potentially improve the ability to track the fate of the Gd(III)-labeled microbubbles after FUS treatment *in vivo*.

## 5.5. Conclusion

FUS in combination with intravenously injected 4-5 micron lipid-stabilized microbubbles labeled with Gd(III)-DOTA successfully induced BBB opening at a safe pressure threshold. The increase in  $T_1$ -weighted MRI contrast after injection of Omniscan was used to confirm BBB opening at the sonicated region. The threshold corresponded to that used to open the BBB using 4-5 micron and 6-8 micron sized microbubbles without Gd(III)-DOTA. However, there was no observable increase in  $T_1$ -weighted contrast in the sonicated region after FUS application in the presence of the Gd(III)-labeled microbubbles. In the future, a potential strategy to increase likelihood of imaging the Gd(III)-labeled shell after FUS sonication may require minimizing the time difference between FUS treatment and MRI scanning.

## **Chapter 6: Concluding Remarks and Future Directions**

### **6.1. Accomplishment of Specific Aims**

#### **6.1.1 Microbubble Size Isolation by differential centrifugation**

This study detailed a methodology to generate lipid-coated, perfluorobutane-filled microbubbles and isolate their size fractions based on migration in a centrifugal field. Polydispersity of the freshly sonicated suspension was characterized by particle sizing and counting through light obscuration/scattering and electrical impedance sensing, fluorescence and bright-field microscopy and flow cytometry. We found that the initial microbubble size distribution was polydisperse. Smaller microbubbles were more abundant. Differential centrifugation was used to isolate the 1–2 and 4–5  $\mu\text{m}$  diameter fractions. Isolated microbubbles were demonstrated to be stable over two days. After two weeks, however, more dilute suspensions (<1 vol%) were susceptible to Ostwald ripening. For example, 4–5  $\mu\text{m}$  microbubbles disintegrated into 1–2  $\mu\text{m}$  microbubbles. This latter observation indicated the existence of an optimally stable diameter in the 1–2  $\mu\text{m}$  range for these lipid-coated microbubbles. Overall, differential centrifugation provided a rapid and robust means for size selection and reduced polydispersity of lipid-coated microbubbles.

#### **6.1.2. Theranostic Gd(III)-Lipid Microbubbles for MRI-guided Focused Ultrasound Surgery**

This study described a technique to synthesize a biomaterial consisting of Gd(III) ions chelated to lipid-coated, size-selected microbubbles for utility in both magnetic resonance and ultrasound imaging. The macrocyclic ligand DOTA-NHS was bound to phosphatidylethanolamine (PE) headgroups on the lipid shell of pre-synthesized microbubbles.

Gd(III) was then chelated to DOTA on the microbubble shell. The reaction temperature of 50 °C was optimized to increase the rate of Gd(III) chelation while maintaining microbubble stability. The Gd(III)-bound microbubbles were found to be echogenic *in vivo* during high-frequency ultrasound imaging of the mouse kidney. The Gd(III)-bound microbubbles also were characterized by magnetic resonance imaging (MRI) at 9.4 T by a spin-echo technique. This initial analysis with MRI revealed that Gd(III) relaxivity increased significantly, and in a concentration-dependent manner, after microbubble fragmentation into non-gas-containing lipid bilayer remnants. We attributed this behavior to an increase in water proton exchange with the Gd(III)-labeled lipid fragments caused by an increase in the lipid headgroup area that accompanied the lipid shell monolayer to bilayer transition.

### **6.1.3. Magnetic Resonance Properties of Gd(III)-Bound Lipid-Coated Microbubbles and their Cavitation Fragments**

In this study, we sought to explore the mechanism governing the changes in  $r_1$  and  $r_2^*$  after fragmentation of microbubbles consisting of Gd(III) labeled to two different groups on the lipid monolayer shell: the phosphatidylethanolamine (PE) lipid headgroup region or the distal region of the polyethylene-glycol (PEG) brush. Nuclear magnetic resonance (NMR) analysis at 1.5 T of the lipid headgroup labeled Gd(III)-microbubbles revealed that  $r_1$  increased on average by a factor of 2.4 after microbubble fragmentation, while  $r_2^*$  was unchanged. Analysis of PEG-labeled Gd(III)-microbubbles revealed that  $r_1$  and  $r_2^*$  decreased on average by a factor of 2.1 and 8 respectively, after microbubble fragmentation. Further analysis revealed that the microbubble gas core enhanced the transverse MR signal ( $T_2^*$ ) in a dose-dependent manner but minimally affected the longitudinal ( $T_1$ ) signal. These results illustrate how the effect of lipid-stabilized

Gd(III)-microbubble cavitation on the MR signal is dependent on the location of Gd(III) on the lipid shell (lipid headgroup-labeled or PEG-labeled) and the presence of the gas.

#### **6.1.4. Opening the Blood-brain Barrier with Gd(III)-bound Lipid-stabilized Microbubbles**

The aim of this study was to demonstrate the feasibility to safely and non-invasively open the BBB with Gd(III)-DOTA-labeled lipid-stabilized microbubbles and to determine the potential of imaging the sonicated region for the presence of lipid shell fragments. The Gd(III)-bound microbubbles successfully induced BBB opening at a pressure threshold of 0.45 MPa. This threshold corresponded to that used to open the BBB using 4-5 micron and 6-8 micron sized microbubbles without Gd(III)-DOTA. The increase in  $T_1$ -weighted MRI contrast after injection of Omniscan was used to confirm BBB opening at the sonicated region. There was no observable increase in  $T_1$ -weighted contrast in the sonicated region after FUS application in the presence of the Gd(III)-bound microbubbles.

## **6.2. Impact on the Field**

The differential centrifugation technique developed in this thesis has been useful in producing the different size classes of microbubbles needed to test the effect of microbubble monodispersity in CEUS studies. For instance, Sirsi *et al.* (Sirsi, Feshitan *et al.*) demonstrated that 4-5 and 6-8 micron diameter offer increased acoustic signal and circulation persistence *in vivo* in comparison to polydisperse and 1-2 micron sized microbubbles, which allowed for much longer imaging sessions during high-frequency ultrasound imaging. Streeter *et al.* (Streeter, Gessner *et al.* 2010) also demonstrated that larger-sized microbubbles, generated using the

differential centrifugation methodology, provided improved CEUS diagnostics of rat tumors in comparison to polydisperse sizes.

These size-selected microbubbles have also been extensively studied and adapted in applications involving opening the BBB with FUS. Choi *et al.* (Choi, Feshitan *et al.*) was the first to demonstrate the dependence of microbubble size on BBB opening with FUS. The 4-5 micron sized microbubbles were demonstrated to induce BBB opening in mice at a lower acoustic threshold than polydisperse and smaller-sized microbubbles. Marquet *et al.* (Marquet, Tung *et al.* 2011) later demonstrated the first successful transcranial BBB opening in non-human primates using FUS and 4-5 micron size-selected microbubbles. Subsequently, Vlachos *et al.* (Vlachos, Tung *et al.* 2011) studied the permeability of BBB opening in the hippocampus of mice after the application of FUS at different acoustic pressures and microbubble sizes. Using dynamic contrast-enhanced MRI, the volume of BBB opening was determined to be proportional to both acoustic pressure and microbubble diameter. Tung *et al.* (Tung, Vlachos *et al.* 2011) later explored the mechanism of BBB opening by comparing the efficacy of 1-2, 4-5 and 6-8 micron size-selected classes. It was demonstrated that BBB opens with nonlinear bubble oscillation when the bubble size is closer to the capillary diameter and with inertial cavitation when it is much smaller than the capillary diameter. The volume of opening was demonstrated to increase with both pressure and microbubble diameter. Finally, Samiotaki *et al.* (Samiotaki, Vlachos *et al.* 2012) demonstrated the dependence of the reversibility of BBB opening in mice on the peak-rarefactional pressure and microbubble size using contrast-enhanced MRI.

The MR behavior of the size-selected lipid-stabilized Gd(III)-bound microbubbles may ultimately prove to be beneficial in the biomedical field. One proposed application for these constructs is to provide an alternative means of monitoring FUS surgery i.e. by monitoring

microbubble fragmentation state with MRI as opposed to the conventional method of MR thermometry. For instance, owing to the observed differences in longitudinal relaxivity between the intact and fragmented constructs,  $T_1$ -weighted MRI along with the lipid-headgroup labeled (Gd(III)-PE) microbubbles can potentially be used to monitor microbubble fragmentation state during FUS therapy. Alternatively, due to the observed differences in transverse relaxivity of intact and fragmented constructs, the PEG labeled Gd(III)-microbubbles and microbubbles without Gd(III) may provide the option to use the  $T_2^*$ -weighted MRI to monitor microbubble fragmentation state.

### 6.3. Future Directions

The next step in adapting the size-selection methodology is the development of continuous size selection process as opposed to the batch system used in this thesis. One possibility is the design of a continuous-flow centrifugal separator that consists of a process input (containing the polydisperse suspension), a process output (containing the size-separated microbubbles) and a recycle stream for added efficiency. This approach may reduce microbubble separation time and minimize the error associated with operator handling that stems from the batch separation method.

Additional studies on the parameters affecting the relaxivity of the Gd(III)-bound lipid microbubbles should also be explored in the future. This includes determining the dependence of varying microbubble shell parameters and gas composition on the relaxivity of Gd(III)-bound construct. For instance, microbubbles with a paramagnetic gas composition such as Oxygen may have different transverse relaxivity from one with a diamagnetic gas composition, such as

Nitrogen. Moreover, the effects of changing external parameters such as the magnetic field strength and local temperature should also be considered in subsequent analysis.

The next step in understanding the behavior of Gd(III)-bound lipid microbubbles *in vivo* is to adapt techniques that improve the chances of tracking the Gd(III)-bound shell. In the BBB opening study performed in chapter 5, the main drawback in detecting the shell fragments at the sonicated hippocampus was the delay between scanning with MRI and FUS sonications. This limitation stemmed from the inherent differences in accessibility of both equipments during experimentation. In order to improve the chances of imaging the shell fragments, it is preferable to utilize an integrated MRI and FUS device similar to the Exablate 2000 (InSightec, Ltd), which would reduce the time difference between imaging and FUS sonications.

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