

Local amplification of Rayleigh waves in the continental United States observed on the USArray

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Abstract

We develop a method based on ratios of amplitudes measured at adjacent stations to determine local amplification of surface waves across an array of seismic stations. We isolate the effects of local structure from those of the earthquake and propagation by systematic averaging of ratios corresponding to many sources. We apply the method to data recorded on the USArray for the years 2006–2011 and determine amplification factors at each station of the array for Rayleigh waves at periods between 35 s and 125 s. Local amplification factors are spatially coherent and display variations of $\pm 10\%$ at a period of 125 s and greater variations at shorter periods. Maps of local amplification exhibit spatial correlation with topography and geologic structures in the western and central United States. At long periods, the observed amplification factors correlate well with predictions from a regional crust and mantle model of North America. At short periods, correlations are weaker, suggesting that the local amplification factors can be useful for constraining shallow structure better.

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1 **1. Introduction**

2 Surface waves provide one of the principal constraints on structure in the
3 uppermost part of the Earth. Both phase and amplitude, the two primary
4 surface-wave observables, contain useful information about structure in the
5 crust and upper mantle. Several different factors contribute to affect the
6 amplitudes of surface waves, including source, path, and receiver effects (e.g.,
7 Selby and Woodhouse, 2000; Dalton and Ekström, 2006b). In addition to this
8 complexity, stations are sometimes poorly calibrated in amplitude. Because
9 of this, amplitude data are used less frequently than phase travel time data
10 in tomography studies. In recent decades, significant work has been done
11 using measurements of surface-wave phase travel times to constrain both
12 two-dimensional (e.g., Zhang and Tanimoto, 1991; Trampert and Woodhouse,
13 1995; Ekström et al., 1997; Trampert and Woodhouse, 2003; Ekström, 2011;
14 Lin and Ritzwoller, 2011) and three-dimensional (e.g., Masters et al., 1996;
15 Boschi and Ekström, 2002; Shapiro and Ritzwoller, 2002; Kustowski et al.,
16 2008; Ritsema et al., 2011) velocity structure in the Earth. Although some
17 studies model both surface-wave phase and amplitude using two-plane and
18 multi-plane wave methods (e.g., Yang and Forsyth, 2006; Pollitz and Snoke,
19 2010), less work has been done to constrain Earth structure using amplitude
20 data. Surface-wave amplitudes are, however, a potentially rich source of
21 information about both elastic and anelastic structure of the crust and upper
22 mantle.

23 Surface-wave amplitudes contain information about propagation effects,

24 including both attenuation and elastic focusing (Selby and Woodhouse, 2000).
25 The majority of surface-wave amplitude studies have focused on constraining
26 anelastic structure of the crust and upper mantle. Previous studies that have
27 developed global models of surface-wave attenuation include Durek et al.
28 (1993), Selby and Woodhouse (2002), Gung and Romanowicz (2004), and
29 Dalton and Ekström (2006b). The use of amplitude measurements to infer
30 elastic structure has been less common (e.g., Dalton and Ekström, 2006a).

31 In addition to propagation effects, recorded surface-wave amplitudes in-
32 clude an amplification factor that depends on the instrument response and
33 Earth structure local to the station. Ideally, the instrument response is
34 known perfectly, and the receiver factor can be attributed entirely to Earth
35 structure. Unfortunately, the absolute instrument gain is difficult to measure
36 and verify, and is therefore associated with significant uncertainty. Errors in
37 the seismometer gain of several percent or more are common. For example,
38 Ekström et al. (2006) investigated the calibration of stations in the Global
39 Seismographic Network and discovered constant and time-dependent calibra-
40 tion errors of the order of 10% at a significant fraction of the stations. At
41 this level, calibration errors may be the dominant contribution to the receiver
42 factor. Interpretation of receiver-amplitude factors in terms of Earth struc-
43 ture is therefore difficult and depends critically on the quality of the station
44 calibration.

45 The Earth structure contribution to the receiver-amplitude factor is a
46 local effect that depends on the elastic structure beneath the station (e.g.,
47 Wang and Dahlen, 1994; Ferreira and Woodhouse, 2007b). Relatively few
48 studies have investigated this site effect in part because of the difficulties in

49 making accurate measurements of surface-wave amplitude. However, some
50 recent work has demonstrated how local elastic structure can be inferred from
51 surface-wave amplitudes. For example, Ferreira and Woodhouse (2007a),
52 Tanimoto and Rivera (2008), Yano et al. (2009), and Lin et al. (2012a)
53 showed that the ratio between vertical and horizontal Rayleigh wave am-
54 plitudes can be used to determine radial shear-velocity structure beneath
55 stations. Recently, Lin et al. (2012b) used measurements of surface-wave
56 phase and amplitude recorded on the USArray to derive maps of local ampli-
57 fication across the western United States using Helmholtz tomography. Local
58 surface-wave amplification in the continental United States is also the focus
59 of this paper.

60 Surface-wave amplitudes can provide a complementary constraint to phase
61 velocity on elastic structure in the crust and upper mantle. Sensitivity of am-
62 plification and phase velocity to perturbations in v_P , v_S , and density, ρ , is
63 shown in the radial sensitivity kernels in Figure 1. The phase velocity ker-
64 nels for velocity perturbations are always positive, while the amplification
65 kernels change sign with depth. Peak sensitivity of surface-wave amplifica-
66 tion to velocity perturbations is slightly shallower than the peak of phase
67 velocity sensitivity. In addition, surface-wave amplitudes are more sensitive
68 than phase velocity to shallow crustal structure. Although the density sen-
69 sitivity is relatively small compared to the velocity sensitivity, amplification
70 is more sensitive than phase velocity to density perturbations. Because of
71 these different sensitivities, combining observations of surface-wave ampli-
72 tude with measurements of surface-wave phase could help to refine current
73 elastic models of the Earth.

74 In this paper, we develop further a two-station method to derive local
75 amplification factors at seismic stations (Eddy and Ekström, 2011). We
76 apply it to data recorded on the USArray and derive maps of local amplifi-
77 cation across the footprint of the array. We assess the quality of calibration
78 of USArray stations. Additionally, we make quantitative comparisons with
79 predictions of local amplification made with mantle and crust models.

80 2. Theory

81 For a given angular frequency, ω , surface-wave seismograms can be writ-
82 ten as a function of amplitude and phase (e.g., Tromp and Dahlen, 1992,
83 1993);

$$u(\omega) = A(\omega) \exp[i\Phi(\omega)], \quad (1)$$

84 where $u(\omega)$ denotes the recorded displacement at the station, and $A(\omega)$ and
85 $\Phi(\omega)$ are the amplitude and phase, respectively. The amplitude of the seis-
86 mogram can be considered to be a product of four separate effects;

$$A(\omega) = A_S(\omega)A_R(\omega)A_F(\omega)A_Q(\omega), \quad (2)$$

87 where $A_S(\omega)$ is the effect of the source, $A_R(\omega)$ is the effect of the receiver,
88 $A_F(\omega)$ is the effect of geometric spreading and focusing, and $A_Q(\omega)$ is the ef-
89 fect of attenuation (e.g., Dalton and Ekström, 2006b). Source effects include
90 both earthquake source parameters and Earth structure local to the source.
91 Receiver effects include both the instrument response and Earth structure
92 local to the receiver. Because there are four contributions to each amplitude
93 signal, there is inherent difficulty in isolating the effect of each.

94 In this study, we form ratios of measured signal amplitudes recorded at
 95 adjacent stations to isolate the component of the surface-wave-amplitude
 96 signal local to the receiver. Consider the ratio of surface-wave amplitudes
 97 from one earthquake recorded at two stations denoted i and j (Note that
 98 although from here on we drop the frequency dependence in the equations,
 99 all quantities still depend on the frequency, ω .);

$$\frac{A_i}{A_j} = \frac{A_{S,i}A_{R,i}A_{F,i}A_{Q,i}}{A_{S,j}A_{R,j}A_{F,j}A_{Q,j}} \quad (3)$$

100 For a given earthquake and station pair, all amplitude effects will contribute
 101 to the observed ratio. For nearby stations, the contribution to the ratio of the
 102 effects associated with the source and propagation is small, since the stations
 103 sample adjacent points of a coherent surface-wave wave field. Because take-
 104 off angles and ray paths for waves recorded at two adjacent stations will be
 105 similar for a given earthquake, the ratios of all effects except the receiver can
 106 be expected to vary around 1.0. In contrast, local effects associated with
 107 receiver structure or instrument response will contribute a factor to the ratio
 108 that is consistent for all earthquakes recorded on the two stations.

109 To isolate the receiver effect, we first construct a datum, d_{ij}^k , from loga-
 110 rithmic amplitude ratios;

$$d_{ij}^k = \ln(A_i/A_j) = \ln(A_i) - \ln(A_j), \quad (4)$$

111 where A_i and A_j are the individual surface-wave amplitude measurements at
 112 a given frequency for each earthquake, k , that was recorded on two neigh-
 113 boring stations. For station pairs on which many earthquakes have been

114 recorded, we then derive an average datum;

$$\bar{d}_{ij} = \frac{1}{N_E} \sum_{k=1}^{N_E} d_{ij}^k, \quad (5)$$

115 where N_E is the total number of earthquakes recorded on both stations in
116 the pair.

117 We wish to attribute the average datum, \bar{d}_{ij} , to the difference between
118 the local receiver effect at each station in the pair, $\ln(A_{R,i}) - \ln(A_{R,j})$. The
119 cancellation of source and path effects in the averaging should work best for
120 stations that are separated by a small distance, for which the surface waves
121 generated by a single earthquake have traveled along nearly identical paths.
122 Ideally, the station separation is a fraction of the wavelength of the wave and
123 of any local complexity in the wavefield.

124 In the surface-wave ray-theoretical framework, the receiver-amplitude fac-
125 tor, A_R , is a function of the local radial elastic structure and has two con-
126 tributing factors (e.g., Tromp and Dahlen, 1992; Wang and Dahlen, 1994).
127 First, the surface amplitudes of the displacement eigenfunction vary depend-
128 ing on the local structure. A second, typically smaller, effect is related to
129 the local speed of propagation. Given a radial profile taken from a three-
130 dimensional Earth model, a predicted receiver-amplification factor can be
131 calculated for a specific location;

$$A_R = \frac{D}{D_0} \sqrt{\frac{U_0}{U}} \quad (6)$$

132 where D and D_0 are the displacement eigenfunctions for the receiver location
133 in the three-dimensional model and a reference model, respectively, U is the
134 group velocity of the model, and U_0 is the group velocity of the reference

135 model. As pointed out by Ferreira and Woodhouse (2007b), specific expres-
136 sions for the receiver-amplification factor vary depending on the convention
137 used in the normalization of the surface-wave eigenfunctions. Here we follow
138 the normal-mode convention (e.g., Gilbert and Dziewoński, 1975; Ferreira
139 and Woodhouse, 2007b).

140 **3. Data and Analysis**

141 The data used in this study are recordings on the USArray of earthquakes
142 with $M_W > 5.5$ occurring between January 2006 and December 2011. Ampli-
143 tudes are measured for minor-arc arrivals of Rayleigh waves between periods
144 of 25 s and 125 s from 2172 earthquakes recorded on a total of 1384 stations.
145 Cumulative geographic station coverage through 2011 is over more than half
146 of the continental United States.

147 Measurements of amplitude anomalies are made using the technique of
148 Ekström et al. (1997). In this method, the model surface wave is first ex-
149 pressed as a function of amplitude and phase using the earthquake location
150 and source geometry taken from the Global Centroid Moment Tensor (CMT)
151 project (Dziewoński et al., 1981; Ekström et al., 2012) and using the excita-
152 tion and propagation characteristics calculated for the Preliminary Reference
153 Earth Model (PREM) (Dziewoński and Anderson, 1981). A misfit function,
154 representing the difference between the modeled and observed waveforms,
155 is then minimized in an iterative process in which phase and amplitude are
156 varied to best fit the waveforms. The amplitude anomaly is the variation
157 in amplitude relative to the spherical Earth synthetic prediction necessary
158 to match the modeled waveform to the observation at each period. In prac-

159 tice, this means that the amplitude ratios used in the analysis (equation 3)
160 have been corrected for the small differences in source radiation pattern and
161 propagation effects on a reference spherical Earth as well as for instrument-
162 response functions.

163 For this analysis, we consider Rayleigh wave amplitudes measured on
164 the vertical component. We select high-quality amplitude-anomaly measure-
165 ments as our data. We only include measurements from earthquakes that are
166 at a distance greater than 15° away from the recording station. Using these
167 single-station data, we form the logarithmic amplitude ratios for station pairs
168 in the array and average the ratios for each pair as in equation 5.

169 A selection based on the maximum distance between stations in each pair
170 is needed to ensure that the source, focusing, and attenuation effects will all
171 be nearly equal for amplitudes of surface waves from one earthquake recorded
172 at two different stations. The smaller the maximum distance between sta-
173 tions in each pair, the more likely it is that this assumption holds. At the
174 same time, a small maximum distance between stations excludes a large por-
175 tion of the raw amplitude dataset. A distance of two degrees was chosen as
176 a compromise that minimizes the differences of source, focusing, and atten-
177 uation effects and retains a large amount of the data. For the USArray, this
178 typically leads to about 20–30 station neighbors.

179 To suppress the source and path effects in the average logarithmic ampli-
180 tude ratios, a sufficiently large number of recorded events must be included.
181 After experimentation, we selected ten as the minimum number of earth-
182 quakes for the analysis. The station pairs that are excluded from our selec-
183 tion as a consequence of this criterion are primarily located on the eastern

184 edge of the array, for which there is the least amount of data.

185 Table 1 provides information about the data used in this study. At long
186 periods (≥ 50 s), significantly more data were used, reflecting the larger num-
187 ber of high-quality observations available. For each station-pair observation,
188 \bar{d}_{ij} , we calculate an associated uncertainty, $\bar{\sigma}_{ij}$, by $\bar{\sigma}_{ij} = \sigma_{ij}/\sqrt{N_E}$, where σ_{ij}
189 is the standard deviation of the observations and N_E is the total number of
190 observations for each station pair. Since the observations are derived from
191 multiple differences, the covariance will be significant for observations on sta-
192 tion pairs which have one station in common, and $\bar{\sigma}_{ij}$ will not be a complete
193 characterization of the data uncertainties. However, the $\bar{\sigma}_{ij}$ values provide
194 a good relative measure of the uncertainties associated with different pairs.
195 The mean of these station-pair uncertainties for each period is reported in
196 Table 1. The increasing uncertainty with decreasing period reflects the larger
197 scatter in the short-period measurements. We find that the large scatter and
198 smaller number of observations available at periods shorter than 35 s lead to
199 unstable results. We therefore limit the further analysis to periods 35 s and
200 longer to ensure that the path-dependent effects on surface-wave amplitudes
201 are averaged out.

202 The average logarithmic amplitude ratio for each station pair is taken as
203 the datum in a least-squares inversion for local station amplification factors
204 by minimizing χ^2 in the following equation;

$$\chi^2 = \sum_{ij} \frac{1}{\bar{\sigma}_{ij}^2} [(\ln(A_{R,i}) - \ln(A_{R,j})) - \bar{d}_{ij}]^2, \quad (7)$$

205 where $A_{R,i}$ and $A_{R,j}$ are the station amplification factors that are inverted
206 for and the datum, \bar{d}_{ij} , and uncertainty, $\bar{\sigma}_{ij}$, are defined as above. Only
207 those stations that are linked together by observations are included in the

208 inversion. Absolute amplification factors cannot be resolved because the
209 data are derived from amplitude differences at each period, and we apply
210 the constraint that the logarithmic station amplification factors must sum to
211 a value of 0 across the array for each period using the method of Lagrange
212 multipliers. With this constraint, the inverse problem is overdetermined and
213 no additional regularization is necessary.

214 **4. Results**

215 Local amplification factors are derived by inversion for Rayleigh waves
216 at discrete periods between 35 s and 125 s. Figure 2 shows the range of
217 local Rayleigh wave amplification variations at different periods. Local am-
218 plification factors are distributed around 1.0, with the spread of variations
219 increasing with decreasing period. The observed variation in surface-wave
220 amplitude due to effects local to the receiver reaches $\pm 10\%$ at 125 s and
221 exceeds $\pm 20\%$ at 35 s.

222 The average logarithmic amplitude ratios are explained well by the de-
223 rived amplitude factors. Figure 3 shows the variance reduction of the local
224 amplification factors at all periods considered in this study. Variance reduc-
225 tion is highest for long-period Rayleigh waves, reaching values up to 95%.
226 The variance reduction is slightly lower at short periods, which could be
227 due to several different reasons, including the smaller number of high-quality
228 measurements and higher variability in the short-period amplitudes due to
229 focusing and deformation of the wavefront at a wavelength less than or on
230 the order of the station spacing of the USArray (70 km).

231 To assess whether the local amplification factors explain the observations

232 at the level of our estimated uncertainties, we calculate the goodness-of-fit
233 parameter χ^2/M (Figure 3), where M is the number of degrees of freedom,
234 here the difference between the number of station-pair observations and the
235 number of stations at each period. We find goodness-of-fit values in the range
236 0.5–0.7. We attribute the apparent over-fitting (i.e., $\chi^2/M < 1.0$) of the data
237 to our incomplete consideration of covariance in the station-pair data.

238 Figure 4 shows maps of the station amplification factors for Rayleigh
239 waves at periods of 35 s, 50 s, 75 s, and 125 s. The amplification factors are
240 relative and can be considered to be local amplification (for values > 1.0) or
241 local deamplification (for values < 1.0) of an incoming wavefield. The anoma-
242 lies exhibit spatial correlation with topography and geologic structures. For
243 example, at the shorter periods, very large amplification (+15%) is seen in
244 Colorado and in the Snake River Plain. Extreme deamplification (-15%) is
245 seen along the Gulf of Mexico.

246 A striking feature of the maps is the spatial coherence of anomalies at the
247 level of a few percent. We investigate the length over which the anomalies
248 are correlated by first calculating the absolute value of the differences in lo-
249 cal amplification for pairs of stations at varying inter-station distances. For
250 0.5-degree bins, we then find the average absolute difference in amplification
251 between stations in each distance range. Figure 5 shows the average ampli-
252 fication difference as a function of inter-station distance for Rayleigh wave
253 amplification factors at the different periods. The distance over which the
254 amplification factors are well correlated is short (< 2 degrees) at 35 s and in-
255 creases at longer periods. Due to finite-frequency effects, this length reflects
256 not only an average scale of structure, but also the averaging of amplification

257 over a period-dependent length. The minimum difference, observed at short
258 inter-station distances for all periods, can be considered a measure of the av-
259 erage quality of calibration of the stations. Based on this analysis, we infer
260 that the relative calibration of USArray stations is very good, with average
261 errors less than 2–4%.

262 Because of the spatial coherence of anomalies, it is possible to identify
263 outlier stations that have problems with instrument response and calibration.
264 Many of these problem stations are visible by inspection of the amplification-
265 factor maps, often as values that are significantly lower than those for the
266 stations surrounding them. One example is TA station N02C-TA, located at
267 40.8°N and 123.3°W. This station has a factor of two difference in gain with
268 respect to neighboring stations, and appears as a distinct negative anomaly
269 in the otherwise smooth local Rayleigh wave amplification maps.

270 We define outliers at each period to be stations that have an amplification
271 that is $> 2\sigma$ different from at least five neighboring stations, where σ
272 is the standard deviation of the observed amplification factors at that period
273 (Figure 2). Here, neighboring stations are defined to be those stations that
274 lie less than 1 degree away from the station of interest. Table 2 lists the
275 outliers for local Rayleigh wave amplification determined using this criteria at
276 different periods. At long periods, a larger number of stations are identified as
277 outliers because there is less variability in the amplification. The anomalous
278 amplification factors of outlier stations are isolated and do not affect the
279 observed amplification at nearby stations. This indicates that we can also
280 resolve variations in amplification that are due to structural anomalies local
281 to each station.

282 5. Discussion

283 Observed local Rayleigh wave amplification factors contain anomalies that
284 are spatially coherent and show good correlation with topography and ge-
285 ologic structures (Figure 4). We attribute the local Rayleigh wave amplifi-
286 cation factors to variations in local elastic structure beneath each station.
287 We observe that the range in variation of the amplification factors (10–20%,
288 depending on the period) is always larger than our inferred estimate of the
289 quality of calibration of USArray stations (Figure 5). For well-calibrated
290 arrays, the effect of local structure on surface-wave amplitudes is larger than
291 effects due to errors in instrument response.

292 Maps of local amplification factors show similarities to surface-wave phase
293 velocity maps in the western United States (e.g., Lin and Ritzwoller, 2011).
294 Locations of high amplification coincide with regions of extremely slow phase
295 velocity in the Rocky Mountains and in the region of the Snake River Plain
296 and Yellowstone Hotspot. Local amplification is also correlated with crustal
297 thickness (Gilbert, 2012). Stations located in regions of thicker crust tend to
298 have larger amplification factors, such as the Sierra Nevada Mountains, the
299 Rocky Mountains in Colorado, and the Snake River Plain and Yellowstone
300 Hotspot. A location with thinner crust, the Columbia Plateau, has lower
301 amplification factors relative to surrounding stations. Even at long periods,
302 there are clear spatial correlations with surface geologic features.

303 The sensitivity kernels in Figure 1 illustrate the difficulties of directly as-
304 sociating the amplitude anomalies with radial elastic structure. In contrast
305 with phase-velocity kernels, which are dominantly positive, so that a slow v_P
306 or v_S intrinsic anomaly at any depth results in a slow phase velocity, ampli-

307 tude kernels change sign with depth. Thus, a shallow, slow v_S anomaly will
308 lead to deamplification while an anomaly located at greater depth will lead
309 to an amplification. Additionally, at shallow depth the v_P and v_S sensitivi-
310 ties are of opposite sign. This effect may explain the extreme deamplification
311 observed along the Gulf of Mexico, an area with a thick layer of low-velocity
312 sediments with high v_P/v_S (Laske and Masters, 1997; Bassin et al., 2000).

313 In order to investigate the extent to which our amplification observations
314 are qualitatively and quantitatively consistent with existing knowledge of the
315 elastic structure of the crust and mantle beneath North America, we make
316 corresponding predictions based on two models. First, we predict surface-
317 wave amplitudes at each station using a model consisting of PREM core
318 and mantle (Dziewoński and Anderson, 1981) and crustal structure from
319 CRUST2.0, a crustal model with 2° by 2° resolution (Bassin et al., 2000).
320 For each station, a radial profile is taken from that location in the model
321 and normal modes are calculated. The predicted vertical amplification fac-
322 tor is then calculated as in equation 6, where D denotes the vertical dis-
323 placement eigenfunction for PREM and CRUST2.0 at that location and the
324 reference model is PREM. Second, to assess the importance of mantle het-
325 erogeneity, we also predict surface-wave amplitudes at each station using the
326 three-dimensional model ND08 (Nettles and Dziewoński, 2008), a radially
327 anisotropic shear-velocity model of the mantle beneath North America, em-
328 bedded in a lower resolution global model. The model incorporates crustal
329 structure from CRUST2.0. Local amplification is predicted in the same man-
330 ner using radial profiles from this model. The effect of density heterogeneity
331 is included in the predictions of amplification; density in the crust is taken

332 from CRUST2.0 and density in the mantle is taken from PREM.

333 Figure 6 shows the correlations between the observed amplification and
334 local amplification predicted by these two models at periods between 35 s
335 and 125 s. At short periods (≤ 40 s), predictions made by the two models are
336 equally well correlated with observed amplification factors, reflecting the fact
337 that they both incorporate and are dominated by the effects of CRUST2.0.
338 At periods longer than 40 s, elastic structure in the mantle begins to dom-
339 inate the receiver-amplitude signal, leading to the large separation of the
340 correlations between the observations and predictions from the two mod-
341 els. Correlations for predictions made by PREM with CRUST2.0 are low
342 or slightly anticorrelated at long periods, indicating that lateral variations
343 in crustal structure alone are not sufficient to predict accurate surface-wave
344 amplitudes.

345 The correlation between observed amplification factors and local ampli-
346 fication predicted by ND08 is large at longer periods with a maximum at
347 100 s ($R = 0.64$). Figure 7 shows maps of the local amplification predicted
348 by model ND08 at periods of 35 s, 50 s, 75 s, and 125 s. The predicted
349 range of amplification is $\pm 10\%$ at 125 s and exceeds $\pm 30\%$ at 35 s. These
350 predicted ranges agree well with the observed range in local amplification
351 (Figure 2). Comparison between the predictions made by these two models
352 demonstrates that models containing laterally varying mantle structure are
353 necessary for accurate predictions of surface-wave amplitudes, despite the
354 strong sensitivity of amplitudes to shallow crustal structure (Figure 1).

355 The agreement found between the observations and the predictions pro-
356 vides corroboration that our method of isolating the local amplification effect

357 is working. We also find good agreement with the recent study of Lin et al.
358 (2012b), who measured Rayleigh wave amplification using Helmholtz tomog-
359 raphy. In their study, Lin et al. (2012b) corrected for the effects of focusing
360 and defocusing by computing the curvature of the phase travel time. The
361 visual agreement between our 50 s map of amplification and their 60 s map
362 (the closest period available) is very good. At 60 s, the correlation of the
363 two sets of amplification factors is high ($R = 0.67$). The good agreement
364 indicates that our observed amplification factors are not strongly biased by
365 unmodeled elastic focusing effects, even though we do not explicitly correct
366 for this effect on amplitudes. We believe that any existing bias due to fo-
367 cusing effects is largest at the shortest periods and negligible at the longest
368 periods considered in this study.

369 There are clear differences between the amplification predicted by model
370 ND08 and that observed from measurements of surface-wave amplitude, de-
371 spite the good agreement at long periods. This suggests that surface-wave
372 amplification factors can help constrain elastic models of the crust and up-
373 per mantle, potentially as an additional data set to invert for radial structure
374 beneath the region in which they were measured. In particular, surface-wave
375 amplification is more sensitive than phase velocity to structure in the crust
376 (Figure 1), indicating that local amplification could improve constraints on
377 shallow structure.

378 We have isolated and investigated local amplification separately from
379 other effects on surface-wave amplitude, including both attenuation and elas-
380 tic focusing. As a result, the derived local amplification factors could also be
381 used as correction factors to constrain these path-dependent effects better.

382 For example, correction of local amplification observations in an inversion for
383 attenuation would reduce the risk that local elastic structure be erroneously
384 mapped into anelastic structure.

385 **6. Conclusions**

386 We have developed a technique to derive local surface-wave amplifica-
387 tion across an array of seismic stations and applied it to data recorded on
388 the USArray. Observed Rayleigh wave amplification factors are spatially co-
389 herent and exhibit good correlation with topography and geologic features,
390 indicating that they are strongly controlled by elastic structure local to each
391 USArray station. Local amplification varies by $\pm 10\%$ at a period of 125 s,
392 demonstrating that local elastic structure has a significant effect on observed
393 surface-wave amplitudes even at long periods. Predictions of local ampli-
394 fication factors show a similar range to the observations but exhibit some
395 differences in pattern, indicating that the surface-wave amplification factors
396 can be used to refine current elastic models of the crust and upper mantle.

397 The quality of amplitude calibration of USArray stations is very good, and
398 variations of surface-wave amplification due to the effect of local structure can
399 be resolved at the level of a few percent. This local amplification effect should
400 not be ignored when studying surface-wave amplitudes on well-calibrated
401 seismic arrays. We are making local amplification factors derived in this
402 study available online.

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Period (sec)	N_S	N_P	\bar{N}_E	$\bar{\sigma}$
25	1065	12930	20.46	0.076
30	1065	12930	20.46	0.061
35	1232	17080	34.00	0.047
40	1232	17080	34.00	0.040
45	1232	17080	34.00	0.035
50	1303	18956	58.07	0.030
75	1303	18956	58.07	0.015
100	1303	18956	58.07	0.011
125	1300	18911	57.52	0.013

Table 1: Data used in this study. Number of stations, N_S , number of station pairs, N_P , average number of observations per station pair, \bar{N}_E , and average station-pair uncertainty, $\bar{\sigma}$.

Station	Latitude	Longitude	35 s	50 s	75 s	125 s
343A-TA	31.284°N	91.617°W	–	0.77	0.78	1.09
H17A-TA	44.395°N	110.576°W	–	1.24	1.51	1.11
IBP-CI	32.661°N	116.093°W	1.01	1.08	1.17	1.20
ISA-CI	35.663°N	118.474°W	0.86	0.77	0.72	0.63
J17A-TA	43.363°N	110.712°W	0.87	0.85	0.84	0.80
MSO-US	46.829°N	113.941°W	0.70	0.67	0.67	0.65
N02C-TA	40.822°N	123.306°W	0.62	0.57	0.57	0.55
OSI-CI	34.614°N	118.724°W	0.95	0.90	0.91	0.87
S43A-TA	37.572°N	90.075°W	0.54	0.53	0.47	0.50
SAO-BK	36.764°N	121.447°W	0.82	0.89	0.91	0.90
T41A-TA	37.044°N	91.764°W	0.86	0.93	0.83	0.80
X41A-TA	34.495°N	92.514°W	0.97	1.17	1.10	0.86
YBH-BK	41.732°N	122.710°W	1.12	1.04	0.94	0.83

Table 2: Outlier stations, latitudes, longitudes, and local amplification factors found for Rayleigh waves at periods of 35 s, 50 s, 75 s, and 125 s. Bold entries are amplification factors which are $> 2\sigma$ from the values of at least 5 neighboring stations, where σ is the standard deviation of the amplification factors at each period. Blank entries indicate that the station was not included in the inversion at that period.

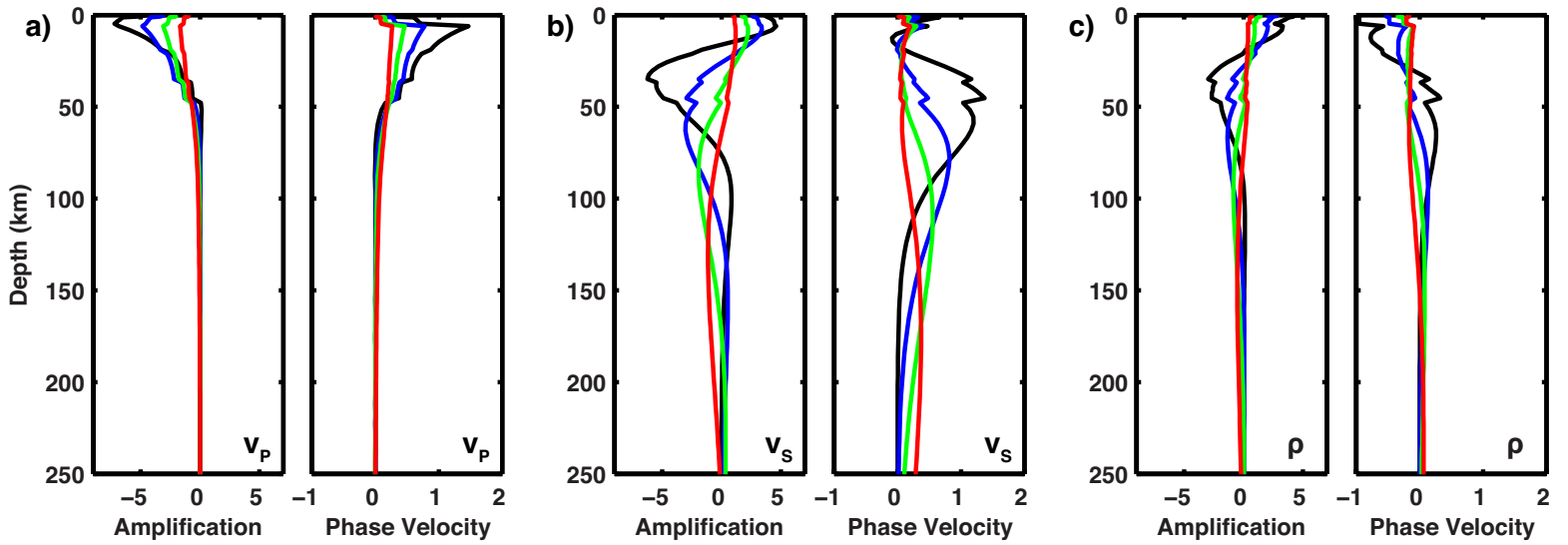


Figure 1: a) v_P , b) v_S , and c) density, ρ , sensitivity kernels for Rayleigh wave amplification (left) and phase velocity (right) at periods of 35 s (black), 50 s (blue), 75 s (green), and 125 s (red), calculated using a one-dimensional profile from ND08 (Nettles and Dziewoński, 2008) and CRUST2.0 (Bassin et al., 2000) located at 35°N and 99°W . Sensitivity kernels are the relative amplitude or phase velocity perturbation due to a 1% velocity or density increase in a 1 km thick layer. Note that the units of amplification and phase velocity are $\times 10^{-5} \%$ and the amplification and phase velocity kernels are always plotted on the same scale.

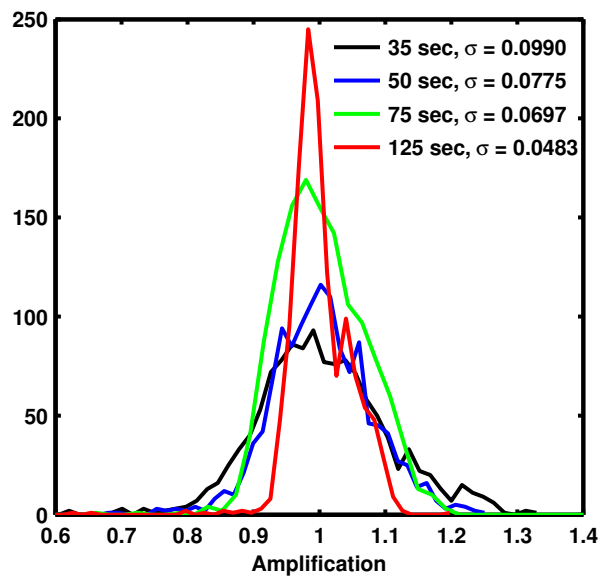


Figure 2: Histogram showing the distribution of local Rayleigh wave amplification factors at periods of 35 s, 50 s, 75 s, and 125 s and the standard deviation, σ , at each period. Observed variations in local amplification reach $\pm 10\%$ at 125 s. The range of variation in local amplification is larger at short periods.

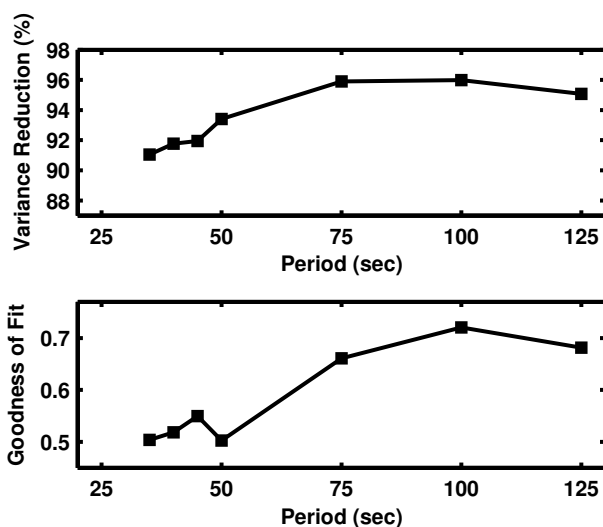


Figure 3: Variance reduction (top) and goodness-of-fit (bottom) for Rayleigh wave amplification factors at periods between 35 s and 125 s. Variance reduction is high for all periods. The goodness-of-fit parameter is equal to χ^2/M , where M is the number of degrees of freedom. We attribute over-fitting (i.e., $\chi^2/M < 1.0$) of the data to overestimated uncertainties.

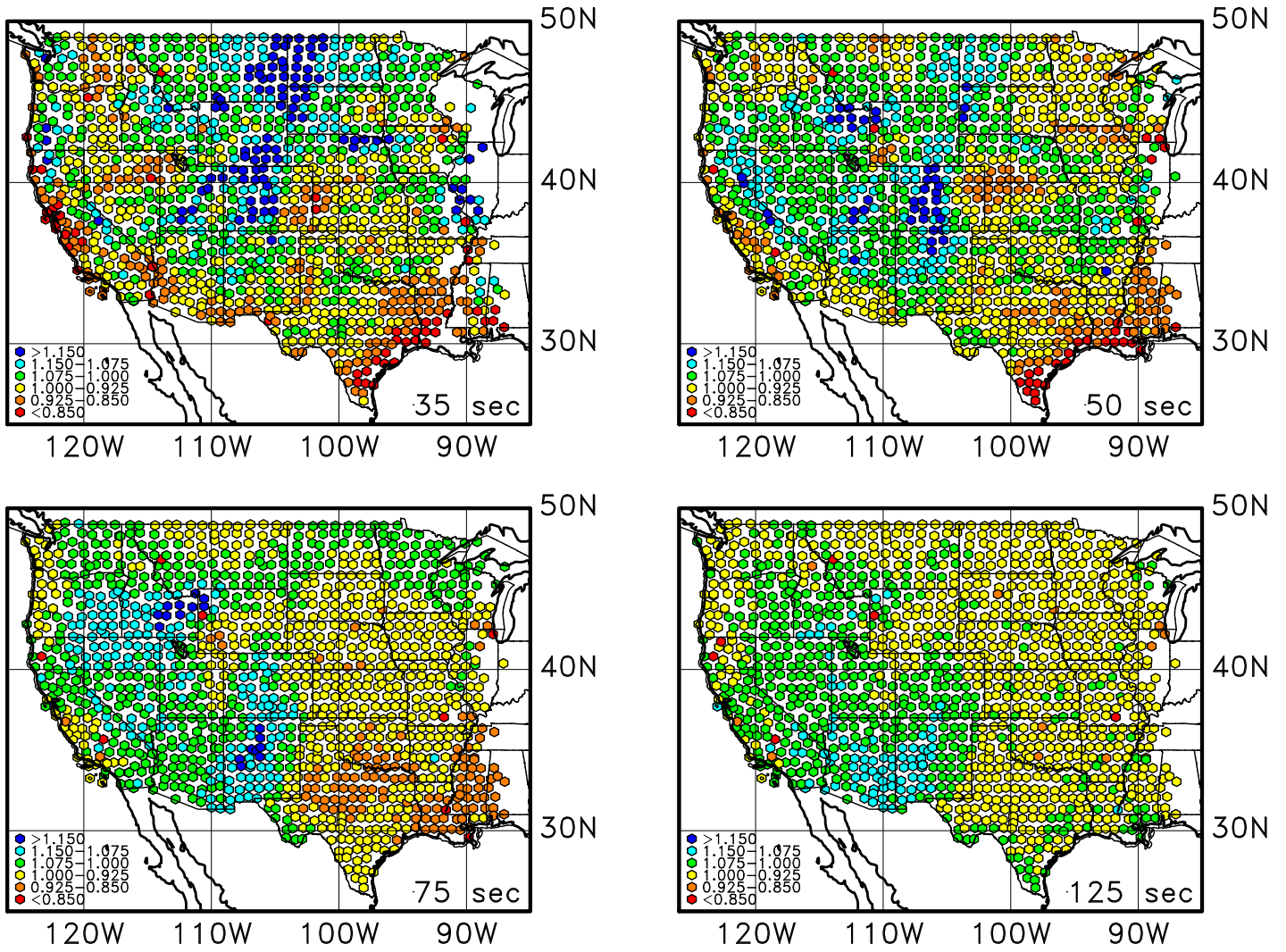


Figure 4: Observed local Rayleigh wave amplification factors at periods of 35 s, 50 s, 75 s, and 125 s. Each symbol corresponds to one USArray station and the color represents the derived amplification factor.

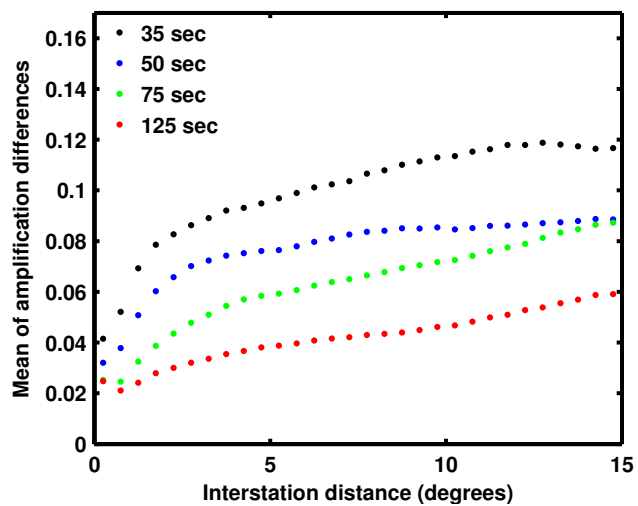


Figure 5: Average absolute differences between local Rayleigh wave amplification factors at stations as a function of inter-station distance at periods of 35 s, 50 s, 75 s, and 125 s. Observed amplification factors are best correlated at short inter-station distances. The minimum difference (2–4%) is an estimate of the average quality of calibration of the stations in the array.

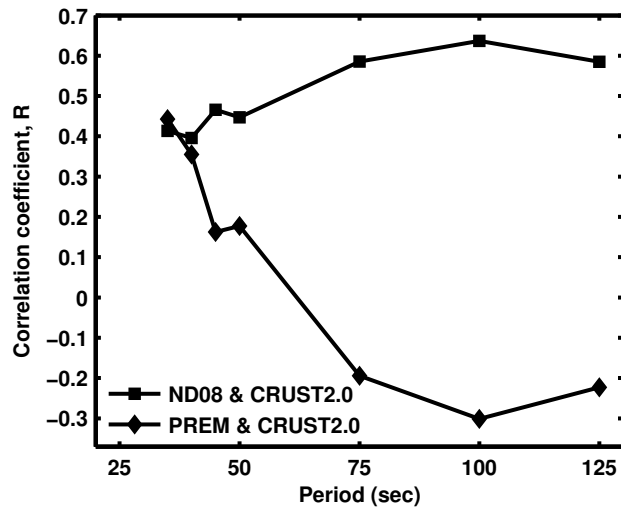


Figure 6: Correlation coefficients, R , between observed local Rayleigh wave amplification factors and predictions of local amplification factors made using models ND08 with CRUST2.0 and PREM with CRUST2.0 at periods between 35 s and 125 s. Predictions made at long periods from model ND08 are best correlated with observed amplification factors.

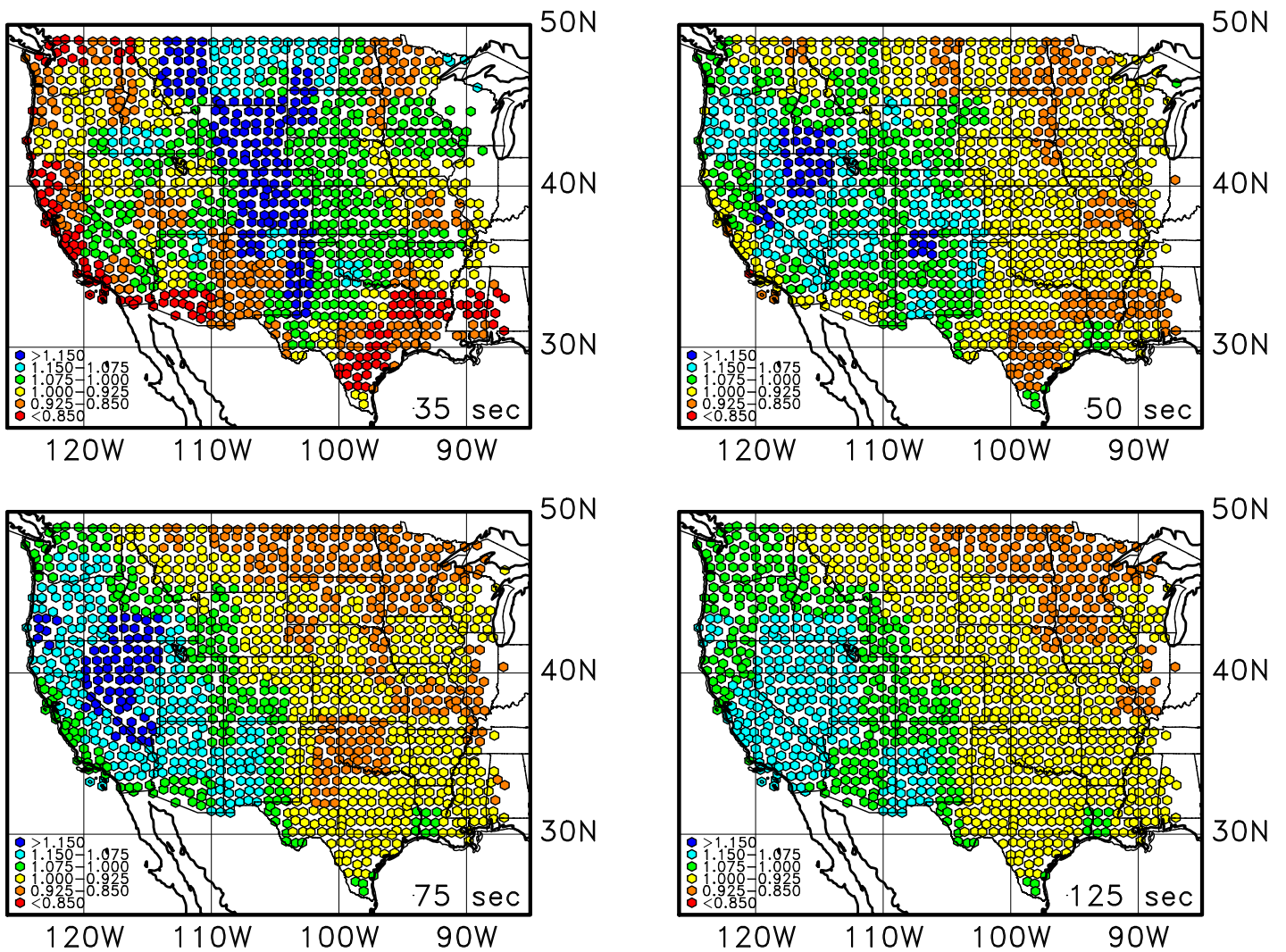


Figure 7: Predicted local Rayleigh wave amplification factors at periods of 35 s, 50 s, 75 s, and 125 s made using model ND08 with CRUST2.0. Each symbol corresponds to one USArray station and the color represents the predicted local amplification factor. Amplification factors are normalized to the mean value at each period for better comparison with the maps of observed amplification and the color scale is the same as in Figure 4.