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Citation	Proceedings of Science (2016), 251
Issue Date	2016-07-15
URL	http://hdl.handle.net/2433/233619
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Type	Conference Paper
Textversion	publisher

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Study of entropy production in Yang-Mills theory with use of Husimi function[*]

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Understanding the thermalization process in a pure quantum system is a challenge in theoretical physics. In this work, we explore possible thermalization mechanism in Yang-Mills(YM) theory by using a positive semi-definite quantum distribution function called a Husimi function which is given by a coarse graining of the Wigner function within the minimum uncertainty. Then entropy is defined in terms of the Husimi function, which is called the Husimi-Wehrl(HW) entropy. We propose two numerical methods to calculate the HW entropy. We find that it is feasible to apply the semi-classical approximation with the use of classical YM equation. It should be noted that the semi-classical approximation is valid in the systems of physical interest including the early stage of heavy-ion collisions. Using a product ansatz for the Husimi function, which is confirmed to reproduce the HW entropy within 20% error (overestimate) for a few-body quantum system, we succeed in a numerical evaluation of HW entropy of YM fields and show that it surely has a finite value and increases in time.

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[*]Report No. : YITP-15-93, KUNS-2592

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1. Introduction

Motivation.— The early thrmalization in relativistic heavy ion collisions suggested by the hydrodynamical simulation is a long-standing problem (see review [1]). One possible way to tackle the problem is to explore the thermalization process by directly calculating entropy that is defined in terms of a quantum distribution function. In the relevant stage of the collisions, a strong gluon field is dominant and its time evolution is described well by classical Yang-Mills equations. In the previous works[2, 3], Kolmogolov-Sinai entropy, an entropy production rate, was found to be positive in this stage and it was pointed out that initial fluctuations play an important role. In this work, we calculate entropy itself in the semi-classical approximation including initial quantum fluctuations.

Because the heavy ion collision occurs in an isolated system, the time-evolution is unitary and von Neumann entropy does not grow in time. Thus we expect that some kinds of coarse-graining can be responsible for entropy production. Partial trace over the subspace of Hilbert space is one of the way to perform coarse-graining. Entropy calculated with the partial-traced density matrix is nothing but the entanglement entropy. We here propose another way of coarse-graining in phase space.

Quantum distribution function and Husimi-Wehrl entropy.— In the n-dimensional quantum mechanical system, a Wigner function is given by the Wigner transform of the density matrix[4];

$$f_W(\vec{p}, \vec{q}; t) = \int d\vec{\eta} \exp(-i\vec{p} \cdot \vec{\eta}/\hbar) \langle \vec{q} + \vec{\eta}/2 | \hat{\rho} | \vec{q} - \vec{\eta}/2 \rangle. \tag{1.1}$$

The expectation value of an observable $\hat{A}$ is given in the phase space integral of observables with a weight of the Winger function; $\langle A(t) \rangle = \int \frac{d\vec{p}d\vec{q}}{(2\pi\hbar)^n} f_W(\vec{p}, \vec{q}; t) A_W(\vec{p}, \vec{q}; t)$ where $A_W(\vec{p}, \vec{q}; t)$ is the Wigner transform of the observable $\hat{A}$. The Wigner function, however, is not positive definite and cannot be treated as a practical quantum distribution function. In order to overcome this shortcoming, we consider Gaussian smearing of the Wigner function;

$$f_H(\vec{q}, \vec{p}; t) = \int \frac{d\vec{p'}d\vec{q'}}{(\pi\hbar)^n} \exp[-\frac{1}{\Delta\hbar}(\vec{p} - \vec{p'})^2 - \frac{\Delta}{\hbar}(\vec{q} - \vec{q'})^2] f_W(\vec{q'}, \vec{p'}; t). \tag{1.2}$$

This is known as the Husimi function[5], which is semi-positive definite and is regarded as a quantum distribution function. Thus we may define the Wehrl entropy in terms of the Husimi function;

$$S_{HW}(t) = -\int \frac{d\vec{p}d\vec{q}}{(2\pi\hbar)^n} f_H(\vec{q}, \vec{p}; t) \log f_H(\vec{q}, \vec{p}; t). \tag{1.3}$$

We call this entropy Husimi-Wehrl(HW) entropy[6].

In the following section, we propose two numerical methods to calculate the time evolution of the HW entropy in the semi-classical approximation. We first apply these methods to quantum mechanical systems with a few degrees of freedom and examine their characteristics and usefulness, which is reported in Ref.[7]. Next we apply them to Yang-Mills theory.

2. The formulation of time evolution and two numerical methods

Time evolution in semi-classical approximation.— We first give the equation of motion for the Wigner function derived from the von Neumann equation;

$$\frac{\partial f_W}{\partial t} = \sum_i^n \frac{\partial H}{\partial q_i} \frac{\partial f_W}{\partial p_i} - \sum_i^n \frac{\partial H}{\partial p_i} \frac{\partial f_W}{\partial q_i} + \mathcal{O}(\hbar^2). \tag{2.1}$$

In the semi-classical approximation, we find that the Wigner function is constant along the classical trajectory; $\dot{p}_i = \partial H/\partial q_i, \dot{q}_i = -\partial H/\partial p_i$. Thus we can obtain the time evolution of the Wigner function in the semi-classical approximation by solving the classical equation of motion. Based on this property, we propose the following two methods to perform an integration needed for the HW entropy.

Numerical methods.— One of the problem to evaluate the HW entropy is that we need integration over large dimensional space, especially in field theories. We propose two Monte-Carlo based numerical methods. The first one is a direct Monte-Carlo integral, referred to as the two step Monte-Carlo(tsMC) method. The HW entropy is given in terms of the Wigner function;

$$S_{HW}(t) = -\int \underbrace{\frac{d\vec{p}d\vec{q}}{(\pi\hbar)^n}}_{A_I} \underbrace{\exp[-\frac{1}{\Delta\hbar}\vec{p}^2 - \frac{\Delta}{\hbar}\vec{q}^2]}_{A_W} \int \underbrace{\frac{d\vec{p'}d\vec{q'}}{(2\pi\hbar)^n}}_{B_I} \underbrace{f_W(\vec{q'},\vec{p'};t)}_{B_W}$$

$$\times \log \int \underbrace{\frac{d\vec{p''}d\vec{q''}}{(\pi\hbar)^n}}_{C_I} \underbrace{\exp[-\frac{1}{\Delta\hbar}(\vec{p}+\vec{p'}-\vec{p''})^2 - \frac{\Delta}{\hbar}(\vec{q}+\vec{q'}-\vec{q''})^2]}_{C_W} f_W(\vec{q''},\vec{p''};t). \tag{2.2}$$

$A_{I(W)}, B_{I(W)}$ and $C_{I(W)}$ are defined by the underlined integrals(functions) in Eq. (2.2). We evaluate the integrals, A_I, B_I and C_I, by Monte Carlo method. For the importance sampling, we generate MC samples according to the weight functions, A_W, B_W and C_W. To prepare the MC samples with the weight of the Wigner function, we generate the samples at $t = 0$ according to the initial distribution, and obtain the corresponding phase space sample points at time t by solving the classical equation of motion. Then the number of the MC samples of integrals, $A_{I(W)}$ and $B_{I(W)}$ $(C_{I(W)})$, out of (in) the logarithmic in Eq. (2.2) is described by $N_{\text{in(out)}}$. In this work we assume that the Wigner function is a semi-positive definite function to apply the MC method.

The second one is a test particle(TP) method. We prepare N_{TP} samples of the (test) particle in phase space so that the sample density is proportional to the initial distribution; $f_W(\vec{q},\vec{p};t) = \frac{(2\pi\hbar)^n}{N_{\text{TP}}} \sum_i^{N_{\text{TP}}} \delta^{(n)}(\vec{p}-\vec{p}^i(t))\delta^{(n)}(\vec{q}-\vec{q}^i(t))$, where N_{TP} is the total number of the test particles. Then the Wigner function is approximately given by the sum of delta functions. We substitute it into Eq. (2.2) and obtain

$$S_{HW}(t) = -\int \underbrace{\frac{d\vec{p}d\vec{q}}{(\pi\hbar)^n}}_{D_I} \underbrace{\exp[-\frac{1}{\Delta\hbar}\vec{p}^2 - \frac{\Delta}{\hbar}\vec{q}^2]}_{D_W} \frac{1}{N} \sum_i^{N_{\text{TP}}}$$

$$\times \log \frac{2^{N_{\text{TP}}}}{N} \sum_j^{N_{\text{TP}}} \exp[-\frac{1}{\Delta\hbar}(\vec{p}+\vec{p}^i(t)-\vec{p}^j(t))^2 - \frac{\Delta}{\hbar}(\vec{q}+\vec{q}^i(t)-\vec{q}^j(t))^2]. \tag{2.3}$$

3

$D_{I(W)}$ is defined by the underlined integral(function) in Eq. (2.3).We evaluate the integral D_I by the Monte Carlo method with the weight D_W. The number of the MC samples is described by N_{MC}.

We have examined these two methods in some two-dimensional quantum mechanical systems with quartic interaction. We find that two methods give consistent results as long as the sampling number is large enough [7].

One of the interesting findings is that when the sampling number is not large enough tsMC tends to overestimate the entropy, while the test particle method tends to underestimate the entropy. The actual HW entropy should be between the two results, because, as the number of samples becomes larger, the results of tsMC(TP) method approach to the converged value from above(below).

3. Extention to Yang-Mills field theory

Classical Yang-Mills field theory.— Now let us discuss a Yang-Mills field. We work in the temporal gauge $A_0 = 0$, then the Hamiltonian in non-compact formalism is given by

$$H = \frac{1}{2} \sum_{x,a,i} E_i^a(x)^2 + \frac{1}{4} \sum_{x,a,i,j} F_{ij}^a(x)^2, \tag{3.1}$$

with $F_{ij}^a = \partial_i A_j^a(x) - \partial_j A_i^a(x) + \sum_{b,c} f^{abc} A_i^b(x) A_j^c(x)$.

If we assign (A, E) to the phase space coordinate, it is straightforward to apply the numerical methods developed in quantum mechanics.

Product ansatz.— In the system with a large degrees of freedom, an extremely large number of samples is necessary for the precision calculation of the HW entropy even when we adopt the present MC methods. To avoid the difficulty, we here propose the product ansatz to evaluate the HW entropy of the Yang-Mills field. We assume that the Husimi function is written as a product of that for one degree of freedom, $f_H(q,p;t) = \prod_i^D h_i(q_i, p_i; t)$. Under this ansatz, we can obtain the HW entropy as a sum of the HW entropy for one
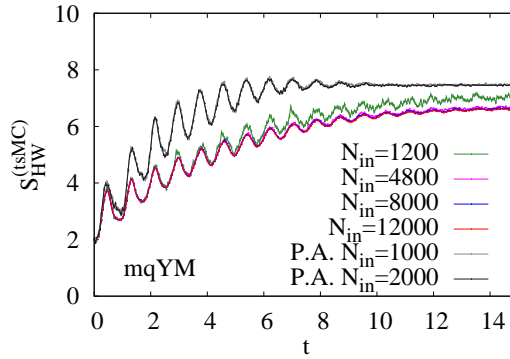

Figure 1: The variation in the product ansatz in the two-dimensional quantum model with tsMC.
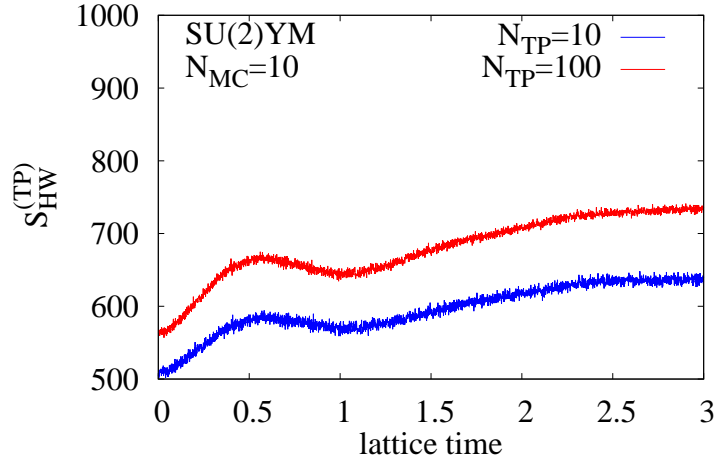
Figure 2: Time evolution of Husimi-Wehrl entropy for the SU(2) Yang-Mills field theory on 4^3 lattice in TP method
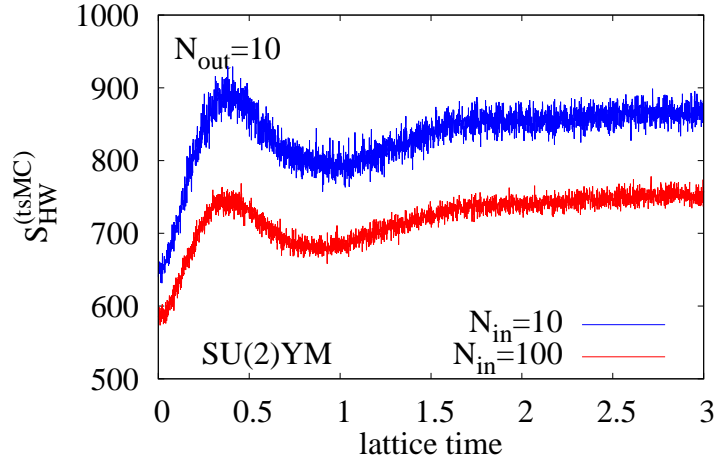

Figure 3: Time evolution of Husimi-Wehrl entropy for SU(2) Yang-Mills field theory on 4^3 lattice in tsMC method

degree of freedom;

$$S_{HW} \simeq -\sum_i^D \int \frac{dq_i dp_i}{2\pi\hbar} h_i(q_i, p_i; t) \log h_i(q_i, p_i; t). \qquad (3.2)$$

Confirmation in quantum mechanical model.— Figure 1 shows that the product ansatz is found to give the HW entropy consistent with the precisely evaluated value within around 10 % error in the case of the two-dimensional quantum mechanical problem.

Results in Yang-Mills field theory.— We apply the product ansatz to the color SU(2) Yang-Mills field on 4^3 and 6^3 lattices.

Figures 2 and 3 show the time dependence of the HW entropy in SU(2) Yang-Mills theory on the 4^3 lattice with the TP and tsMC methods, respectively. We find that the
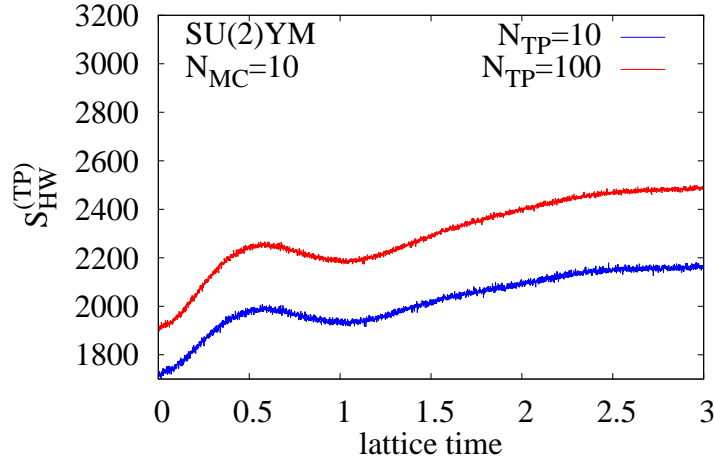
Figure 4: Time evolution of Husimi-Wehrl entropy for SU(2) Yang-Mills field theory on 6^3 lattice in TP method
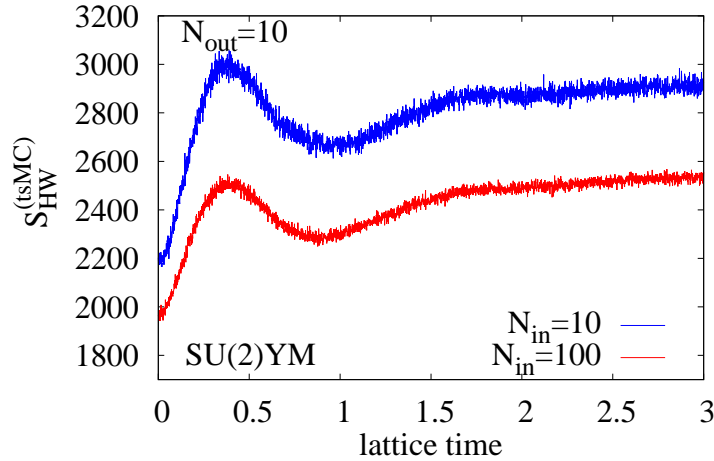

Figure 5: Time evolution of Husimi-Wehrl entropy for SU(2) Yang-Mills field theory on 6^3 lattice in tsMC method

HW entropy increases as a function of time in both methods, which implies thermal entropy production. As in the quantum mechanics cases, the HW entropy in the tsMC(TP) method decreases(increases) with increasing Monte-Carlo samples(test particles). The entropy value at t=3 is almost the same in the two methods with a large number of MC samples(test particles).

Figures 4 and 5 are the results in SU(2) Yang-Mills theory on 6^3 lattice with the TP and tsMC methods, respectively. The qualitative behavior is the same as that in 4^3 lattice. The HW entropy is created also on the larger lattice.

4. Summary

In order to explore the thermalization process in the early stage of the relativistic heavy ion collisions, we have proposed to use a positive semi-definite quantum distribution function called the Husimi function, which is given by a Gaussian smearing of the Wigner function within the minimum uncertainty: Since the semiclassical approximation is valid in the system of our interest and the classical Yang-Mills(YM) dynamics is known to show a chaotic behaior which can be a significant origin of the entropy production, the use of the distribution function is useful because it allows a classical-quantum correspondence. Then a kind of entropy can be defined in terms of the Husimi function, which is called the Husimi-Wehrl(HW) entropy. We have shown that it is feasible to apply the semiclassical approximation to calculate the HW entropy with the use of the classical YM equation. We have proposed two numerical methods to calculate the HW entropy on the basis of the Monte-Carlo method, which we call two step Monte-Carlo(tsMC) and test particle(TP) methods, respectively: It is found that the combination of these two methods gives reliable results for the entropy production for two-dimensional quantum mechanical systems. We have also introduced a product ansatz for calculating the HW entropy for a system with large number of degrees of freedom. We have confirmed that the product ansatz roughly reproduce the HW entropy with about 10-20% overestimate for a few-body quantum system. We have calculated the HW entropy of (quantum) Yang-Mills field theory with the tsMC and TP methods coupled with product ansatz, and found that the HW entropy increases in time. This implies that thermal entropy can be surely created in the early stage of the relativistic heavy-ion collisions, which is triggered by the initial fluctuations.

Acknowledgments

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