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Turing's Method for the Selberg Zeta-Function

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Abstract: In one of his final research papers, Alan Turing introduced a method to certify the completeness of a purported list of zeros of the Riemann zeta-function. In this paper we consider Turing's method in the analogous setting of Selberg zeta-functions and we demonstrate that it can be carried out rigorously in the prototypical case of the modular surface.

1. Introduction

In [28], Turing described and implemented a numerical procedure for verifying the Riemann Hypothesis (RH) up to a given height T in the critical strip. Turing's procedure was similar to earlier numerical investigations of RH by Gram [11], Backlund [2], Hutchinson [16], and Titchmarsh [27], in that they were all based on isolating zeros on the critical line by finding sign changes in the Hardy function $Z(t)$, and then confirming that no zeros had been missed. Turing's approach differs only in the latter step; where the earlier authors used ad hoc procedures that are valid only for small values of T , Turing introduced a method for certifying the completeness of a purported list of zeros of $Z(t)$ that is guaranteed to work (when the list is in fact complete). Turing's method has remained a small but essential ingredient in all subsequent verifications of RH and its many generalizations; see [5] and [6] for more on Turing's method and its historical background.

Meanwhile, researchers in the high energy physics community have since the early 1990s applied the same idea to certifying lists of zeros of Selberg zeta-functions for hyperbolic manifolds, albeit at a heuristic level (without explicit error estimates) and without attribution to Turing; see for instance [25], where it was used in one of the first investigations of large eigenvalues of the Laplacian for the modular group, $\mathrm{PSL}(2, \mathbb{Z})$. In this paper we show, much in the spirit of Turing's computations, that the method can be made rigorous in the case of the modular group.

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We begin by describing Turing’s method in greater detail, in the context of the Selberg zeta-function. Let $\mathbb{H} = \{z \in \mathbb{C} : \Im(z) > 0\}$ denote the hyperbolic plane, and let $\{f_j\}_{j=1}^\infty$ be a complete sequence of Hecke–Maass cuspforms on $\mathrm{PSL}(2, \mathbb{Z}) \backslash \mathbb{H}$, with Laplacian eigenvalues $\frac{1}{4} + r_j^2$ satisfying $0 < r_1 \leq r_2 \leq \dots$. Then the associated Selberg zeta-function $Z_{\mathrm{PSL}(2, \mathbb{Z})}(s)$ has zeros at $s = \frac{1}{2} \pm ir_j$.

Let $N(t) = \#\{j : r_j \leq t\}$ denote the counting function of zeros in the upper half plane, and suppose that we have accurately computed several r_j up to some height T , so we can construct a minorant $N^-(t)$ of the step function $N(t)$, for $t \leq T$. Suppose hypothetically that we miss an r_j at $T - H$ for some $H > 0$, so that $N(t) \geq N^-(t) + 1$ for $t \geq T - H$. Integrating this inequality, we get

$$\int_0^T N^-(t) dt + H \leq \int_0^T N(t) dt = \int_0^T (T - t) dN(t) = \mathrm{Tr}(h_0), \quad (1.1)$$

where $h_0(r) = \max(0, T - |r|)$ and $\mathrm{Tr}(h) = \sum_{j=1}^\infty h(r_j)$ denotes the trace of h over the cuspidal spectrum. Unfortunately, although h_0 has a trace in this sense, it is not suitable for applying the Selberg trace formula, but we can get around that by replacing h_0 by a majorant h_0^+ with Fourier transform of compact support. Thus, we have $\int_0^T N^-(t) dt + H \leq \mathrm{Tr}(h_0^+)$.

Applying the trace formula, the right-hand side will be the expected main term (described by Weyl’s law) plus the error that arises from truncating the support of the Fourier transform. If it happens that we have not actually missed any zeros and we know them precisely enough then we can expect $\int N^-(t) dt$ to be close to the main term, and in fact it may even exceed the main term sometimes. This gives us an upper bound $H \leq H^+$, i.e. we can provably show that there are no missing zeros up to height $T - H^+$. Moreover, using both upper and lower bounds for $\mathrm{Tr}(h_0)$ we can estimate $\int_{T_1}^{T_2} N(t) dt$, which would allow one to carry out the procedure using only the zeros in an interval around T .

This approach is guaranteed to work for large enough T because the error term $S(t)$ in Weyl’s law has mean value 0; precisely, it is known that $\frac{1}{T} \int_0^T S(t) dt = \mathcal{O}((\log T)^{-2})$ (see [14, Ch. 10, Thm. 2.29]). Consequently, one can expect to prove by this method that there are no missing zeros up to $T - \mathcal{O}(T(\log T)^{-2})$.

The remainder of the paper is devoted to obtaining such a bound with explicit (and practical!) constants. Our precise result is the following.

Theorem 1.1. *Define*

$$S(t) = N(t) - \left(\frac{t^2}{12} - \frac{2t}{\pi} \log \frac{t}{e\sqrt{\frac{\pi}{2}}} - \frac{131}{144} \right) \quad \text{and} \quad E(t) = \left(1 + \frac{6.59125}{\log t} \right) \left(\frac{\pi}{12 \log t} \right)^2.$$

Then for $T > 1$,

$$\frac{1}{T} \int_0^T S(t) dt \leq E(T). \quad (1.2)$$

To demonstrate the usefulness of this bound in practice, we applied it to the list of zeros $r_j \leq 178$, which are shown to 20 decimal place accuracy at [8]. This list was kindly provided to us by Andreas Strömbergsson, who computed the r_j using Hejhal’s algorithm [13] and certified them using the program from [10]. Using Theorem 1.1, we obtain the following:

Corollary 1.2. *The Selberg zeta-function for $\mathrm{PSL}(2, \mathbb{Z}) \backslash \mathbb{H}$ has exactly 2 184 zeros with imaginary part in $(0, 177.75]$. All of them are simple.*

Remarks.

- (1) Our method can also be used to show the lower bound

$$\frac{1}{T} \int_0^T S(t) dt \geq -(2 + o(1)) \left(\frac{\pi}{12 \log t} \right)^2,$$

though we do not prove that here. The factor of 2 difference between the upper and lower bounds is due to an asymmetry when approximating h_0 above and below by band-limited functions, as we make explicit in the next section.

- (2) Following Selberg, Hejhal proved the analogue of the estimate $\frac{1}{T} \int_0^T S(t) dt \ll (\log T)^{-2}$ for a general cofinite Fuchsian group, using the theory of the Selberg zeta-function; see [15, Ch. 2, Thm. 9.7] and [14, Ch. 10, Thm. 2.29]. Sarnak has suggested that this estimate could be obtained directly via the trace formula, and our work realizes that goal in the case of $\mathrm{PSL}(2, \mathbb{Z})$. Our method could be generalized to congruence subgroups, and we expect the implied constants that it produces to compare favorably to the Selberg–Hejhal method (which has not been made explicit, to our knowledge).

In the specific case of the modular group, the full asymptotic for $N(t)$ appearing in Theorem 1.1 was computed by Steil [25]. A detailed proof was given by Jorgenson, Smajlović, and Then [19], who also computed the lower-order terms of the asymptotic in the case of moonshine groups.

- (3) The leading order constant $(\pi/12)^2$ could in principle be divided by 4 by using the Kuznetsov formula and the method of Li and Sarnak [20]; however, it would be quite cumbersome to work out an explicit error term in that setting. A more practical means of achieving a factor of 2 savings in the asymptotic result would be to split the spectrum into even and odd parts and derive a bound for each separately; even there, however, the secondary error terms would be worse, and it would likely not result in a savings for T of any practical size.
- (4) The estimate (1.2) is substantially weaker than Turing's estimate in the context of the Riemann zeta-function, which is $\mathcal{O}(\frac{\log T}{T})$. One key reason for this is that the Selberg zeta-function has a much higher density of zeros at large height ($\sim \frac{1}{6}t$ vs. $\sim \frac{1}{2\pi} \log t$ for Riemann zeta), which in turn makes $S(t)$ noisier; see Figs. 1 and 2 for a comparison of the two over the range $t \in [0, 100]$. In fact it is only by virtue of the fact that the analogue of RH is known to hold in this context that we can prove that $\frac{1}{T} \int_0^T S(t) dt = o(1)$ as $T \rightarrow \infty$. The true rate of decay of $\frac{1}{T} \int_0^T S(t) dt$ is not known, but extensive numerics of Then [26] suggest that it should be $o(T^{-\frac{1}{2}})$.

The outline of the paper is as follows. In Sect. 1.1, we show how the optimization of our upper bound leads naturally to an extremal problem in Fourier analysis, and we sketch a solution that explains the leading-order constants that we can expect to achieve. In Sect. 2, we recall the Selberg trace formula for $\mathrm{PSL}(2, \mathbb{Z})$, which will be our main tool. We will then deduce an asymptotic for the main term of the trace formula applied to h_0 (Sect. 3) which we can use to estimate the average of $S(t)$ for various ranges of T (Sect. 4). Section 5 describes the computation to bound the constant B introduced

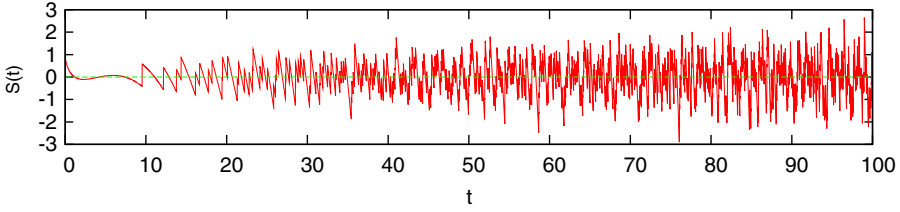


Fig. 1. $S(t)$ for the Selberg zeta-function

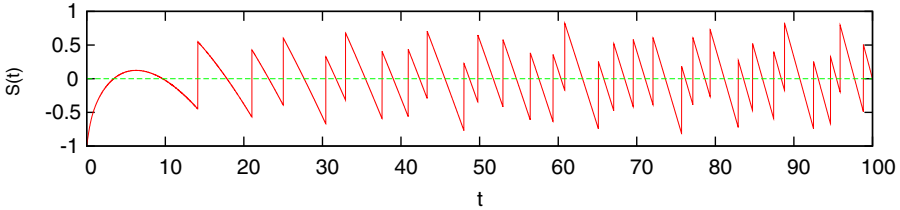


Fig. 2. $S(t)$ for the Riemann zeta-function

in Sect. 4, and finally, in Sect. 6, we verify Theorem 1.1 by considering T in the four ranges $(1, 100]$, $[100, 27\,400]$, $[27\,400, 10^6]$ and $[10^6, \infty)$. Appendix A contains some details of the techniques used to perform the necessary computations rigorously.

1.1. An extremal problem. The main error in our estimate comes from approximating h_0 by h_0^+ for the upper bound, and similarly by a minorant for the lower bound. To control these errors, we have two contrary objectives. First, we want h_0^+ to be a good approximation to h_0 , in order to minimize the contribution from the main terms of the trace formula, which for large T are essentially of the form $\frac{1}{12} \int_{\mathbb{R}} |r| h_0^+(r) dr$. Second, we want the support of the Fourier transform of h_0^+ to be as small as possible in order to control the contribution from the hyperbolic terms. Thus, we are led naturally to the following extremal problem.

Objective. Set $f(x) = \max(0, 1 - |x|)$. Given a large $\Delta > 0$, find functions f^\pm such that

- (1) the Fourier transforms \hat{f}^\pm are supported in $[-\Delta, \Delta]$;
- (2) $f^-(x) \leq f(x) \leq f^+(x) \quad \forall x \in \mathbb{R}$;
- (3) $\int_{\mathbb{R}} |x|(f(x) - f^-(x)) dx$ and $\int_{\mathbb{R}} |x|(f^+(x) - f(x)) dx$ are minimal.

In this section, we obtain the following result toward this objective.

Proposition 1.3. *There are functions f^\pm as above with*

$$\int_{\mathbb{R}} |x|(f(x) - f^-(x)) dx = \frac{1 + o(1)}{6\Delta^2} \quad \text{and} \quad \int_{\mathbb{R}} |x|(f^+(x) - f(x)) dx = \frac{1 + o(1)}{12\Delta^2} \quad (1.3)$$

as $\Delta \rightarrow \infty$.

Remark. It seems likely that this result is asymptotically best possible. However, we do not attempt to prove that assertion here.

Proof (sketch). Since the two cases are similar, we present the proof for f^+ only. Let $\varphi : \mathbb{R} \rightarrow \mathbb{R}$ be an even function with Fourier transform supported in $[-1, 1]$, $\hat{\varphi}(0) = 1$ and $\varphi(x) \ll \frac{1}{1+x^6}$. Set $\varphi_\Delta(x) = \Delta\varphi(\Delta x)$. Let $u(x) = \frac{1}{2}(1 + \operatorname{sgn} x)$ be the unit step function, and $v(x) = \max(0, x)$; these are approximated by the successive integrals $U(x) = \int_{-\infty}^x \varphi(y) dy$ and $V(x) = \int_{-\infty}^x U(y) dy$. Set $E(x) = U(x) - u(x)$ and $F(x) = V(x) - v(x)$. Then we have

$$\begin{aligned} f * \varphi_\Delta(x) - f(x) &= - \int_{\mathbb{R}} f(t) dE(\Delta(x - t)) = \int_{\mathbb{R}} f'(t) E(\Delta(x - t)) dt \\ &= \Delta^{-1} [F(\Delta(x - 1)) + F(\Delta(x + 1)) - 2F(\Delta x)]. \end{aligned} \tag{1.4}$$

Now suppose that $F(x) \geq 0$ for all x . Let F^+ be a majorant of F , with Fourier transform supported in $[-1, 1]$, such that $F^+(x) - F(x) = \mathcal{O}(x^{-4})$; such a function is easily constructed using the techniques of [29]. Then $f^+(x) = f * \varphi_\Delta(x) + 2\Delta^{-1}F^+(\Delta x)$ majorizes $f(x)$. Further, we have

$$\begin{aligned} \int_{\mathbb{R}} |x| (f^+(x) - f(x)) dx &= \Delta^{-3} \int_{\mathbb{R}} |x| [F(x - \Delta) + F(x + \Delta) + 2(F^+(x) - F(x))] dx \\ &= 2\Delta^{-2} \int_{\mathbb{R}} F(x) dx + 2\Delta^{-3} \int_{|x| > \Delta} (|x| - \Delta) F(x) dx \\ &\quad + 2\Delta^{-3} \int_{\mathbb{R}} |x| (F^+(x) - F(x)) dx \\ &= 2\Delta^{-2} \int_{\mathbb{R}} F(x) dx + \mathcal{O}(\Delta^{-3}). \end{aligned} \tag{1.5}$$

Clearly we want to minimize $\int_{\mathbb{R}} F(x) dx$. Equivalently, V should be a majorant of v of exponential type 2π such that $\int_{\mathbb{R}} (V(x) - v(x)) dx$ is minimal. The unique such function, found recently by Littmann [21], is $V_0(x) = \left(\frac{\cos \pi x}{\pi}\right)^2 [x\psi'(\frac{1}{2} - x) + 1]$, where $\psi(x) = \frac{\Gamma'}{\Gamma}(x)$; it satisfies $2 \int_{\mathbb{R}} (V_0(x) - v(x)) dx = \frac{1}{12}$. Unfortunately, we cannot simply set $V = V_0$ in our application, for then the integral over $|x| > \Delta$ in (1.5) would diverge. However, it is not hard to see that one can come arbitrarily close to the constant $\frac{1}{12}$ using an approximation argument (e.g., we can take V as defined in Proposition 6.4, with $X = 1$ and δ sufficiently small); we carry out the full details of this in Sect. 6.1. \square

2. The Trace Formula for $\operatorname{PSL}(2, \mathbb{Z})$

Our main tool will be the following version of the Selberg trace formula.

Proposition 2.1 (The Selberg trace formula for Maass forms on $\operatorname{PSL}(2, \mathbb{Z}) \backslash \mathbb{H}$). *Let $\{r_j\}_{j=1}^\infty$ be as in the introduction. Further, for a fixed $\delta > 0$, let $h(t)$ be an analytic function on $\{t \in \mathbb{C} : |\Im t| \leq \frac{1}{2} + \delta\}$ satisfying $|h(t)| \ll (1 + |\Re t|)^{-2-\delta}$ and $h(t) = h(-t)$. Define \hat{h} , the Fourier transform of h , via*

$$\hat{h}(t) = \int_{\mathbb{R}} h(r) \exp(-2\pi i r t) dr.$$

Then,

$$\mathrm{Tr}(h) = \sum_{j=1}^{\infty} h(r_j) = M(h) + R(h),$$

where

$$M(h) = (I + E + P)(h) - h(0), \quad R(h) = (D - C)(\hat{h}),$$

$$I(h) = \frac{1}{12} \int_{\mathbb{R}} r \tanh(\pi r) h(r) dr = -\frac{1}{24\pi} \int_{\mathbb{R}} \frac{\hat{h}'(t)}{\sinh \pi t} dt,$$

$$E(h) = \int_{\mathbb{R}} \frac{\frac{1}{8} + \frac{1}{3\sqrt{3}} \cosh\left(\frac{\pi r}{3}\right)}{\cosh(\pi r)} h(r) dr = \int_{\mathbb{R}} \left(\frac{1}{8 \cosh \pi t} + \frac{2 \cosh \pi t}{3 + 6 \cosh 2\pi t} \right) \hat{h}(t) dt,$$

$$\begin{aligned} P(h) &= \frac{1}{2\pi} \int_{\mathbb{R}} (\log 2\pi - 2\psi(1 + 2ir)) h(r) dr \\ &= \frac{\hat{h}(0)}{2\pi} \left(\log \frac{\pi}{2} + 2\gamma \right) - \frac{h(0)}{4} - \frac{1}{\pi} \int_0^{\infty} \log \left(4 \sinh \left(\frac{\pi t}{2} \right) \right) \hat{h}'(t) dt, \end{aligned}$$

$$\begin{aligned} D(\hat{h}) &= \frac{1}{\pi} \sum_{l=3}^{\infty} \frac{L(1, \chi_d)}{l} \prod_{p|l} \left(1 + (p - \chi_d(p)) \frac{(p^{\infty}, l) - 1}{p - 1} \right) \hat{h} \left(\frac{1}{\pi} \log \left(\frac{t + \sqrt{t^2 - 4}}{2} \right) \right) \\ &\quad + \sum_{n=1}^{\infty} \frac{\Lambda(n)}{n} \hat{h} \left(\frac{\log n}{\pi} \right) \end{aligned}$$

and

$$C(\hat{h}) = \int_{\mathbb{R}} \hat{h}(t) (\cosh \pi t - 1) dt.$$

In the definition of D , we write $t^2 - 4 = dl^2$, where d is a fundamental discriminant, $l > 0$, χ_d is the corresponding quadratic character, and $(p^{\infty}, l) = \gcd(p^{\infty}, l)$ is the largest power of p dividing l .

Remark. We refer to I , E , P , D , and C as the identity, elliptic, parabolic, discrete, and continuous terms, respectively. For “wide” test functions (those of the sort used to measure Weyl’s Law), one can think of M and R as the main term and remainder, respectively.

Although the terms of the trace formula are only defined for analytic test functions, both $\mathrm{Tr}(h)$ and $M(h)$ can be interpreted for any even, continuous function $h : \mathbb{R} \rightarrow \mathbb{C}$ satisfying $h(t) \ll (1 + |t|)^{-2-\delta}$. In turn, for any such h , we may define $R(h)$ by the equality $\mathrm{Tr}(h) = M(h) + R(h)$.

Proof. Theorem 3 of [9] gives the trace formula for $\Gamma_0^{\pm}(N)$ with character χ . Thus we need to set $N = 1$, take χ to be the trivial character of modulus $q(\chi) = 1$ and then sum over $\epsilon \in \{0, 1\}$ to account for both even and odd eigenfunctions. We also define χ_d and $t^2 - 4 = dl^2$ as in the statement of the theorem and $d(n)$ to be the usual divisor function. Now, considering the term

$$\frac{\prod_{p|N}(p+1)}{24} \int_{\mathbb{R}} r h(r) \tanh(r) dr,$$

we note that the product is empty so $\epsilon \in \{0, 1\}$ together contribute

$$\frac{1}{12} \int_{\mathbb{R}} r h(r) \tanh(r) dr = I(h).$$

Turning now to

$$-\frac{d(N)}{4} g(0) \left(\log(8N) + C_{\chi, \epsilon} \log \left(\frac{Nq(\chi)}{2\pi^2} \right) + \frac{1}{2} (C_{\chi, \epsilon} - 1) \log \gcd(N, 2) \right),$$

since $\chi(-1) = 1$ (or, in the language of [9], χ is "pure"), we have $C_{\chi, \epsilon} = 2, 0$ for $\epsilon = 0, 1$, respectively. Thus we get

$$-\frac{1}{2} \frac{\hat{h}(0)}{2\pi} (3 \log 2 - \log 2 - 2 \log \pi) = -\frac{\hat{h}(0)}{2\pi} \log \frac{2}{\pi} = \frac{1}{2\pi} \int_{\mathbb{R}} \log \left(\frac{\pi}{2} \right) h(r) dr.$$

The next line,

$$-\frac{d(N)}{4\pi} \int_{\mathbb{R}} h(r) \left[\psi \left(\frac{1}{2} + ir \right) + C_{\chi, \epsilon} \psi(1 + ir) \right] dr,$$

yields

$$-\frac{1}{2\pi} \int_{\mathbb{R}} h(r) [2\psi(1 + 2ir) - 2 \log 2] dr,$$

and combining this with the previous term we get

$$\frac{1}{2\pi} \int_{\mathbb{R}} [\log(2\pi) - 2\psi(1 + 2ir)] h(r) dr = P(h).$$

The final line contributes nothing as the first sum therein is empty and the first term of the penultimate line does not apply either as $\chi = 1$. The second term, i.e.

$$d(n) \sum_{n=1}^{\infty} \frac{\Lambda(n) \cdot \{\chi\}_{\epsilon}(n)}{n} g(2 \log n),$$

results in

$$\frac{1}{\pi} \sum_{n=1}^{\infty} \frac{\Lambda(n)}{n} \hat{h} \left(\frac{\log n}{\pi} \right). \tag{2.1}$$

This leaves us with the term spanning lines 2 and 3. [9, (2.60)] gives us

$$\langle \chi(\delta) \rangle_{\delta^2 - t\delta + n = 0} = 1$$

and the argument on page 142 of the same shows that

$$\sum_{f|l} \mathbf{h}^+(\tau[f]) \cdot [\tau[l]^l : \tau[f]^l] = \mathbf{h}^+(d) \prod_{p|l} \left(1 + (p - \chi_d(p)) \frac{(p^{\infty}, l) - 1}{p - 1} \right),$$

with $\mathbf{h}^+(d)$ denoting the narrow class number of $\mathbb{Q}(\sqrt{d})$. This leads us to

$$\sum_{\substack{t \in \mathbb{Z} \\ \sqrt{t^2-4N} \notin \mathbb{Q}}} \mathbf{h}^+(d) \prod_{p|l} \left(1 + (p - \chi_d(p)) \frac{(p^\infty, l) - 1}{p - 1} \right) A(t, 1),$$

where

$$A(t, n) = \begin{cases} \frac{\log \epsilon_1}{\sqrt{t^2-4n}} \cdot g \left(\log \frac{(|t| + \sqrt{t^2-4n})^2}{4} \right) & \text{if } t^2 - 4n > 0, \\ \frac{2}{|\tau[1]^1| \cdot \sqrt{4-t^2}} \cdot \int_{\mathbb{R}} \frac{\exp(-2r \cdot \arccos(|t|/2))}{1 - \exp(-2\pi r)} h(r) dr & \text{if } t^2 - 4n < 0. \end{cases}$$

Here ϵ_1 is the proper fundamental unit in $\mathbb{Z}[\frac{d+\sqrt{d}}{2}]$ and $|\tau[1]^1|$ is the size of its norm one unit group. Considering the sum for $|t| \geq 3$, where $t^2 - 4 > 0$, we get

$$\begin{aligned} & \frac{1}{\pi} \sum_{t=3}^{\infty} \mathbf{h}^+(d) \prod_{p|l} \left(1 + (p - \chi_d(p)) \frac{(p^\infty, l) - 1}{p - 1} \right) \frac{\log \epsilon_1}{\sqrt{t^2-4}} \hat{h} \left(\frac{1}{\pi} \log \left(\frac{t + \sqrt{t^2-4}}{2} \right) \right) \\ &= \frac{1}{\pi} \sum_{t=3}^{\infty} \frac{L(1, \chi_d)}{l} \prod_{p|l} \left(1 + (p - \chi_d(p)) \frac{(p^\infty, l) - 1}{p - 1} \right) \hat{h} \left(\frac{1}{\pi} \log \left(\frac{t + \sqrt{t^2-4}}{2} \right) \right), \end{aligned} \tag{2.2}$$

by Dirichlet's class number formula. We now combine (2.1) and (2.2) to get $D(\hat{h})$. Turning to the $t = 0$ term, we have $d = -4$, $l = 1$, $\mathbf{h}^+(-4) = 1$, $|\tau[1]^1| = 4$ and we get

$$\mathbf{h}^+(-4)A(0, 1) = \mathbf{h}^+(-4) \frac{2}{|\tau[1]^1| \cdot \sqrt{4}} \int_{\mathbb{R}} \frac{\exp(-\pi r)}{1 + \exp(-2\pi r)} h(r) dr = \int_{\mathbb{R}} \frac{1}{8 \cosh(\pi r)} h(r) dr.$$

Now we consider the $t = \pm 1$ terms ($d = -3$, $l = 1$, $\mathbf{h}^+(-3) = 1$, $|\tau[1]^1| = 6$), which contribute

$$\mathbf{h}^+(-3) \frac{4}{|\tau[1]^1| \cdot \sqrt{3}} \int_{\mathbb{R}} \frac{\exp\left(-\frac{2\pi r}{3}\right)}{1 + \exp(-2\pi r)} h(r) dr = \frac{1}{3\sqrt{3}} \int_{\mathbb{R}} \frac{\cosh\left(\frac{\pi r}{3}\right)}{\cosh(\pi r)} h(r) dr,$$

where we have used the fact that h is even. Thus the $A(\pm 1, 1)$ terms together with the $A(0, 1)$ term make up $E(h)$. Clearly the $t = 2$ term contributes nothing as $\sqrt{4-4} \in \mathbb{Q}$.

Finally, we have the contribution from the constant function with eigenvalue $\lambda = \frac{1}{4} + r^2 = 0$, which leads to

$$-h\left(\frac{i}{2}\right) = - \int_{\mathbb{R}} \hat{h}(t) \cosh \pi t dt,$$

since \hat{h} is even. We can now write

$$- \int_{\mathbb{R}} \hat{h}(t) \cosh \pi t dt = - \int_{\mathbb{R}} \hat{h}(t) [\cosh \pi t - 1] dt - h(0) = -C(\hat{h}) - h(0),$$

and we have the first form of the trace formula as stated.

We now use the results of [9, p. 141] (adjusted for our definition of the Fourier transform) to convert from h to \hat{h} , to whit:

$$\begin{aligned} h(0) &= \int_{\mathbb{R}} \hat{h}(t) dt, \\ \int_{\mathbb{R}} r h(r) \tanh(\pi r) dr &= -\frac{1}{2\pi} \int_{\mathbb{R}} \frac{\hat{h}'(t)}{\sinh(\pi t)} dt \text{ and} \\ \int_{\mathbb{R}} \frac{h(r)}{\cosh(\pi r)} dr &= \int_{\mathbb{R}} \frac{\hat{h}(t)}{\cosh(\pi t)} dt. \end{aligned}$$

Referring to [3, (12) on p. 31], we also have

$$\int_{\mathbb{R}} \frac{\cosh\left(\frac{\pi r}{3}\right)}{\cosh(\pi r)} h(r) dr = 2\sqrt{3} \int_{\mathbb{R}} \frac{\cosh \pi t}{1 + 2 \cosh 2\pi t} \hat{h}(t) dt.$$

Finally, we consider $P(h)$. Adjusting the identities in [9, p. 141] for our Fourier transform convention, we have

$$\begin{aligned} \frac{1}{2\pi} \int_{\mathbb{R}} \psi(1 + ir) h(r) dr &= -\gamma \frac{\hat{h}(0)}{2\pi} + \frac{1}{4\pi^2} \int_0^\infty \log(u) \hat{h}'(u/2\pi) du \\ &\quad + \frac{1}{8\pi} \int_{\mathbb{R}} \hat{h}(u/2\pi) du \\ &\quad + \frac{1}{4\pi^2} \int_0^\infty \log\left(\frac{\sinh(u/2)}{u/2}\right) \hat{h}'(u/2\pi) du \end{aligned}$$

and

$$\begin{aligned} \frac{1}{2\pi} \int_{\mathbb{R}} \psi\left(\frac{1}{2} + ir\right) h(r) dr &= -\gamma \frac{\hat{h}(0)}{2\pi} + \frac{1}{4\pi^2} \int_0^\infty \log(u) \hat{h}'(u/2\pi) du \\ &\quad - \frac{1}{4\pi^2} \int_0^\infty \log\left(\frac{\sinh(u/2)}{u/2}\right) \hat{h}'(u/2\pi) du \\ &\quad + \frac{1}{2\pi^2} \int_0^\infty \log\left(\frac{\sinh(u/4)}{u/4}\right) \hat{h}'(u/2\pi) du. \end{aligned}$$

Thus, starting from

$$\begin{aligned} P(h) &= \frac{1}{2\pi} \int_{\mathbb{R}} (\log(2\pi) - 2\psi(1 + 2ir)) h(r) dr \\ &= \frac{1}{2\pi} \int_{\mathbb{R}} (\log(2\pi) - \psi(1 + ir) - \psi(1/2 + ir) - 2 \log 2) h(r) dr, \end{aligned}$$

we obtain

$$\begin{aligned} P(h) &= \frac{\hat{h}(0)}{2\pi} \left(\log \frac{\pi}{2} + 2\gamma \right) - \frac{h(0)}{4} - 2 \int_0^\infty \log(4 \sinh(u/4)) \frac{\hat{h}'(u/2\pi)}{4\pi^2} du \\ &= \frac{\hat{h}(0)}{2\pi} \left(\log \frac{\pi}{2} + 2\gamma \right) - \frac{h(0)}{4} - \frac{1}{\pi} \int_0^\infty \log(4 \sinh(\pi t/2)) \hat{h}'(t) dt, \end{aligned}$$

as required. \square

3. An Asymptotic for $M(h_0)$

Recall that if we set $h_0(r) = \max(0, T - |r|)$ then $\text{Tr}(h_0) = \int_0^T N(t) dt$, where N is the counting function of Theorem 1.1. Motivated by this, we establish the following asymptotic upper bound for $M(h_0)$:

Proposition 3.1. *For $T \geq 4$, we have*

$$M(h_0) \leq \frac{1}{36}T^3 - \frac{\log(T)}{\pi}T^2 + \frac{3 + \log \frac{\pi}{2}}{2\pi}T^2 - \frac{131}{144}T + \frac{\log T}{24\pi} + C_0,$$

where

$$C_0 = \frac{\zeta(3)}{16\pi^3} - \frac{1}{4\pi^2} \left(2L(2, \chi_{-4}) + 3\sqrt{3}L(2, \chi_{-3}) \right) - \frac{1}{2\pi} \left(\zeta'(-1) - \frac{\log 2 + 1}{12} \right).$$

The proof proceeds by examining the contributions from the identity term, the elliptic terms, the parabolic terms and the constant eigenfunction in turn, as follows.

3.1. The identity term.

Lemma 3.2. *For $T \geq 0$, we have*

$$\frac{1}{12} \int_{-T}^T (T - |r|)r \tanh(\pi r) dr \leq \frac{T^3}{36} - \frac{T}{144} + \frac{\zeta(3)}{16\pi^3}.$$

Proof. We have

$$\begin{aligned} & \frac{1}{12} \int_{-T}^T (T - |r|)r \tanh(\pi r) dr \\ &= \frac{1}{6} \int_0^T (T - r)r dr + \frac{1}{6} \int_0^\infty (r - T)r(1 - \tanh(\pi r)) dr \\ & \quad - \frac{1}{6} \int_T^\infty (r - T)r(1 - \tanh(\pi r)) dr \\ &\leq \frac{T^3}{36} + \frac{1}{6} \int_0^\infty (r - T)r(1 - \tanh(\pi r)) dr = \frac{T^3}{36} - \frac{T}{144} + \frac{\zeta(3)}{16\pi^3}. \end{aligned}$$

□

3.2. The elliptic terms.

Lemma 3.3. *Similarly, for $T \geq 0$ we get*

$$\int_{-T}^T (T - |r|) \frac{\frac{1}{8} + \frac{\sqrt{3}}{9} \cosh(\pi r/3)}{\cosh(\pi r)} dr \leq \frac{25T}{72} - \frac{L(2, \chi_{-4})}{2\pi^2} - \frac{3\sqrt{3}L(2, \chi_{-3})}{4\pi^2} + E_e,$$

where

$$E_e = \frac{\exp(-\pi T)}{2\pi} \left(T + \frac{1}{\pi} + \frac{\sqrt{3}}{12} \left[e^{\pi T/3} \left(8T + \frac{12}{\pi} \right) + e^{-\pi T/3} \left(4T + \frac{3}{\pi} \right) \right] \right).$$

Proof. We have

$$\int_0^\infty \frac{dr}{\cosh(\pi r)} = \frac{1}{2}, \quad \int_0^\infty \frac{r dr}{\cosh(\pi r)} = \frac{2L(2, \chi_{-4})}{\pi^2},$$

$$\int_0^\infty \frac{\cosh(\pi r/3) dr}{\cosh(\pi r)} = \frac{1}{\sqrt{3}} \quad \text{and} \quad \int_0^\infty \frac{r \cosh(\pi r/3) dr}{\cosh(\pi r)} = \frac{27L(2, \chi_{-3})}{8\pi^2}.$$

Thus we can write

$$\begin{aligned} & \int_{-T}^T (T - |r|) \frac{\frac{1}{8} + \frac{\sqrt{3}}{9} \cosh(\pi r/3)}{\cosh(\pi r)} dr \\ &= \frac{T}{4} \int_0^T \frac{dr}{\cosh(\pi r)} - \frac{1}{4} \int_0^T \frac{r}{\cosh(\pi r)} dr + \frac{2\sqrt{3}T}{9} \int_0^T \frac{\cosh(\pi r/3)}{\cosh(\pi r)} dr \\ & \quad - \frac{2\sqrt{3}}{9} \int_0^T \frac{r \cosh(\pi r/3)}{\cosh(\pi r)} dr \\ &\leq \frac{T}{4} \int_0^\infty \frac{dr}{\cosh(\pi r)} - \frac{1}{4} \int_0^\infty \frac{r}{\cosh(\pi r)} dr + \frac{1}{4} \int_T^\infty \frac{r}{\cosh(\pi r)} dr \\ & \quad + \frac{2\sqrt{3}T}{9} \int_0^\infty \frac{\cosh(\pi r/3)}{\cosh(\pi r)} dr \\ & \quad - \frac{2\sqrt{3}}{9} \int_0^\infty \frac{r \cosh(\pi r/3)}{\cosh(\pi r)} dr + \frac{2\sqrt{3}}{9} \int_T^\infty \frac{r \cosh(\pi r/3)}{\cosh(\pi r)} dr \\ &\leq \frac{T}{8} - \frac{L(2, \chi_{-4})}{2\pi^2} + \frac{e^{-\pi T}}{2\pi} \left(T + \frac{1}{\pi} \right) + \frac{2T}{9} - \frac{3\sqrt{3}}{4\pi^2} L(2, \chi_{-3}) \\ & \quad + \frac{e^{-\pi T}}{2\pi} \cdot \frac{\sqrt{3}}{12} \left[e^{\pi T/3} \left(8T + \frac{12}{\pi} \right) + e^{-\pi T/3} \left(4T + \frac{3}{\pi} \right) \right]. \end{aligned}$$

□

3.3. *The parabolic terms.* We will need the following preparatory lemma.

Lemma 3.4. Write $\log \Gamma(z) = A(z) + R(z)$, where $A(z) = (z - \frac{1}{2}) \log z - z + \frac{1}{2} \log(2\pi) + 1/(12z)$. Then we have

$$C_{\frac{1}{2}} := \Re \int_{\frac{1}{2}}^{\frac{1}{2}+i\infty} R(z) dz = \frac{1}{2} \zeta'(-1) + \frac{\log 2}{12} + \frac{1}{48}$$

and

$$C_1 := \Re \int_1^{1+i\infty} R(z) dz = -\zeta'(-1) - \frac{1}{6}.$$

Proof.

$$\begin{aligned}
C_1 &= \Re \int_1^{1+i\infty} (\log \Gamma(z) - A(z)) dz \\
&= 2\Re \int_{\frac{1}{2}}^{\frac{1}{2}+i\infty} (\log \Gamma(2z) - A(2z)) dz \\
&= 2\Re \int_{\frac{1}{2}}^{\frac{1}{2}+i\infty} (\log \Gamma(z) + \log \Gamma(z + \frac{1}{2}) - \frac{1}{2} \log \pi + (2z - 1) \log 2 - A(2z)) dz \\
&= 2\Re \int_{\frac{1}{2}}^{\frac{1}{2}+i\infty} (A(z) + \log \Gamma(z + \frac{1}{2}) - \frac{1}{2} \log \pi + (2z - 1) \log 2 - A(2z)) dz + 2C_{\frac{1}{2}} \\
&= 2\Re \int_1^{1+i\infty} (A(z - \frac{1}{2}) - A(2z - 1) + \log \Gamma(z) - \frac{1}{2} \log \pi + (2z - 2) \log 2) dz + 2C_{\frac{1}{2}}.
\end{aligned}$$

This leads us to

$$\begin{aligned}
2C_{\frac{1}{2}} + C_1 &= 2\Re \int_1^{1+i\infty} (A(2z - 1) - A(z) - A(z - \frac{1}{2}) + \frac{1}{2} \log \pi + (2 - 2z) \log 2) dz \\
&= -2\Im \int_0^\infty (A(1 + it) - A(1 + 2it) + A(\frac{1}{2} + it) - \frac{1}{2} \log \pi + (1 + 2it) \log 2) dt,
\end{aligned}$$

which, by the miracle that is Maple™, gives

$$2C_{\frac{1}{2}} + C_1 = \frac{\log 2}{6} - \frac{1}{8}.$$

In addition, by [30, (A.14)] we have

$$\int_{\frac{1}{2}}^1 \log \Gamma(x) dx = \frac{\log \pi}{4} - \frac{1}{8} + \frac{7}{24} \log 2 + \frac{3}{2} \zeta'(-1),$$

so that

$$C_{\frac{1}{2}} = C_1 + \int_{\frac{1}{2}}^1 R(x) dx = C_1 + \int_{\frac{1}{2}}^1 (\log \Gamma(x) - A(x)) dx = C_1 + \frac{3}{16} + \frac{\log 2}{12} + \frac{3}{2} \zeta'(-1),$$

and we are done. \square

We can now handle the parabolic term as follows:

Lemma 3.5. *Let $T > 1$. Then*

$$\begin{aligned}
P(h_0) &= \frac{1}{2\pi} \int_{-T}^T (T - |r|)(\log(2\pi) - \psi(1 + 2ir)) dr \\
&\leq \frac{3 - 2 \log T + \log(\pi/2)}{2\pi} T^2 - \frac{T}{4} + \frac{1}{24\pi} + \frac{\log T}{24\pi} + \frac{\log 2}{24\pi} - \frac{\zeta'(-1)}{2\pi} \\
&\quad - \frac{241}{5760\pi} T^{-2} + \frac{17641}{161280\pi} T^{-4}.
\end{aligned}$$

Proof. We write

$$\begin{aligned}
 & \frac{1}{2\pi} \int_{-T}^T (T - |r|)(\log(2\pi) - \psi(1 + 2ir)) \, dr \\
 &= \frac{1}{2\pi} \int_{-T}^T (T - |r|)(\log(\pi/2) - \psi(1 + ir) - \psi(\frac{1}{2} + ir)) \, dr \\
 &= \log(\pi/2)T^2 - \frac{1}{2\pi} \int_{-T}^T (T - |r|)(\psi(1 + ir) + \psi(\frac{1}{2} + ir)) \, dr \\
 &= \log(\pi/2)T^2 + \frac{1}{\pi} \Re \int_1^{1+iT} \log \Gamma(z) \, dz + \frac{1}{\pi} \Re \int_{\frac{1}{2}}^{\frac{1}{2}+iT} \log \Gamma(z) \, dz.
 \end{aligned}$$

We now write $\log \Gamma(z) = A(z) + R(z)$ as in Lemma 3.4. Then we have

$$\frac{1}{\pi} \Re \int_{\sigma}^{\sigma+iT} \log \Gamma(z) \, dz = \frac{1}{\pi} \Re \left[\int_{\sigma}^{\sigma+iT} A(z) \, dz - \int_{\sigma+iT}^{\sigma+i\infty} R(z) \, dz + \int_{\sigma}^{\sigma+i\infty} R(z) \, dz \right].$$

The integrals involving A evaluate to

$$\begin{aligned}
 & \frac{1}{24\pi} \left[(1 - 6T^2) \log(T^2 + 1) + 36T^2 - (12T^2 + 1) \log \sqrt{4T^2 + 1} \right. \\
 & \quad \left. - 12T \arctan(T) + 12T^2 \log 2 \right],
 \end{aligned}$$

and since $T > 1$ we have

$$\begin{aligned}
 -\arctan(T) &\leq -\frac{\pi}{2} + \frac{1}{T} - \frac{1}{3T^3} + \frac{1}{5T^5}, \\
 \log(T^2 + 1) &\leq 2 \log T, \\
 -\log(T^2 + 1) &\leq -2 \log T - T^{-2} + \frac{1}{2}T^{-4}
 \end{aligned}$$

and

$$-\log \sqrt{4T^2 + 1} \leq -\log 2 - \log T - \frac{1}{8}T^{-2} + \frac{1}{32}T^{-4}$$

so we can maximise the contribution from A with

$$\frac{1}{\pi} \left[\left(\frac{3}{2} - \log T \right) T^2 - \frac{\pi}{4}T + \frac{3}{16} - \frac{\log 2}{24} + \frac{\log T}{24} - \frac{5}{128}T^{-2} + \frac{773}{7680}T^{-4} \right].$$

We can now use Lemma 3.4 to handle the integrals of R from σ to $\sigma + i\infty$ and then using [12, (4.1)] to bound the error in Stirling's approximation we have

$$R(z) \leq -\frac{1}{360z^3} + \frac{2}{315|z|^5}$$

so that

$$\begin{aligned}
 -\Re \left[\int_{1+iT}^{1+i\infty} R(z) \, dz + \int_{\frac{1}{2}+iT}^{\frac{1}{2}+i\infty} R(z) \, dz \right] &\leq -\Re \left[(1+iT)^{-2} + \left(\frac{1}{2} + iT \right)^{-2} \right] + \frac{1}{315}T^{-4} \\
 &\leq -\frac{1}{360}T^{-2} + \frac{1}{180}T^{-4} + \frac{1}{315}T^{-4}.
 \end{aligned}$$

□

3.4. The constant eigenfunction.

Lemma 3.6. *The contribution from the constant eigenfunction is trivially*

$$-h_0(0) = -T.$$

3.5. *Proof of Proposition 3.1.* We now combine Lemmas 3.2, 3.3, 3.5 and 3.6, observing that

$$-\frac{241}{5760\pi}T^{-2} + \frac{17641}{161280\pi}T^{-4} + E_e(T)$$

is negative for $T \geq 4$.

4. An Upper Bound on $\int S(t) dt$

We will now use the trace formula to derive an upper bound for $\int S(t) dt$. We start with some preparatory lemmas.

Lemma 4.1. *Let $\varphi : \mathbb{R} \rightarrow \mathbb{R}$ be a smooth, even function such that $x^2\varphi(x)$ is absolutely integrable and $\int_{\mathbb{R}} \varphi(x) dx = 1$. Define*

$$F(r) = \int_{-\infty}^r \int_{-\infty}^y \varphi(x) dx dy - \max(0, r).$$

Then F is even and absolutely integrable, with Fourier transform

$$\hat{F}(t) = \frac{1 - \hat{\varphi}(t)}{(2\pi t)^2}.$$

Proof. Define

$$U(y) = \int_{-\infty}^y \varphi(x) dx \quad \text{and} \quad u(y) = \frac{1 + \operatorname{sgn} y}{2},$$

so that

$$F(r) = \int_{-\infty}^r (U(y) - u(y)) dy.$$

Then $U(y) - u(y)$ is an odd function of y , and

$$|U(y) - u(y)| = \left| \int_{-\infty}^{-|y|} \varphi(x) dx \right| \leq y^{-2} \int_{-\infty}^0 x^2 |\varphi(x)| dx \ll y^{-2}.$$

Hence, for $r \leq 0$, we have

$$\int_{-\infty}^r (r - x)\varphi(x) dx = \int_{-\infty}^r U(y) dy = F(r),$$

by partial integration. Thus,

$$\begin{aligned} \int_{-\infty}^0 |F(r)| \, dr &\leq \int_{-\infty}^0 \int_{-\infty}^r (r-x)|\varphi(x)| \, dx \, dr = \int_{-\infty}^0 \int_x^0 (r-x)|\varphi(x)| \, dr \, dx \\ &= \frac{1}{2} \int_{-\infty}^0 x^2 |\varphi(x)| \, dx < \infty. \end{aligned}$$

Since $U - u$ is odd, F is even, and therefore absolutely integrable, by the above. Hence, F has a continuous Fourier transform.

Next, let $f : \mathbb{R} \rightarrow \mathbb{R}_{\geq 0}$ be a smooth, even function with mass 1 and support $[-1, 1]$. For $\varepsilon > 0$, define $f_\varepsilon(x) = \varepsilon^{-1} f(\varepsilon^{-1}x)$ and

$$F_\varepsilon(r) = \int_{-\infty}^r \int_{-\infty}^y (\varphi(x) - f_\varepsilon(x)) \, dx \, dy = \int_{-\infty}^r (r-x)(\varphi(x) - f_\varepsilon(x)) \, dx. \quad (4.1)$$

Then the difference

$$F_\varepsilon(r) - F(r) = \max(0, r) - \int_{-\infty}^r (r-x)f_\varepsilon(x) \, dx$$

is supported on $[-\varepsilon, \varepsilon]$ and bounded by ε . Therefore, \hat{F}_ε converges uniformly to \hat{F} as $\varepsilon \rightarrow 0^+$. Differentiating (4.1) twice, we have

$$F_\varepsilon''(r) = \varphi(r) - f_\varepsilon(r),$$

and taking the Fourier transform of both sides we get

$$-(2\pi t)^2 \hat{F}_\varepsilon(t) = \hat{\varphi}(t) - \hat{f}(\varepsilon t).$$

Taking $\varepsilon \rightarrow 0^+$ yields the desired identity. \square

Lemma 4.2. *Let $h_0(r) = \max(0, T - |r|)$ and φ, F be as in Lemma 4.1. Then*

$$(h_0 * \varphi - h_0)(r) = F(r - T) + F(r + T) + 2F(r).$$

Proof. By the convolution theorem, we have

$$\begin{aligned} (h_0 * \varphi - h_0)(r) &= \int_{\mathbb{R}} \hat{h}_0(t) \hat{\varphi}(t) \exp(2\pi i t r) \, dt - h_0(r) \\ &= \int_{\mathbb{R}} \hat{h}_0(t) \left[1 - (2\pi t)^2 \hat{F}(t) \right] \exp(2\pi i t r) \, dt - h_0(r) \\ &= - \int_{\mathbb{R}} \hat{h}_0(t) (2\pi t)^2 \hat{F}(t) \exp(2\pi i t r) \, dt. \end{aligned}$$

A direct computation shows that $\hat{h}_0(t) = T^2 \operatorname{sinc}^2(\pi T t)$, where

$$\operatorname{sinc} t := \begin{cases} \frac{\sin t}{t} & \text{if } t \neq 0, \\ 1 & \text{if } t = 0. \end{cases}$$

Thus,

$$\begin{aligned}
 (h_0 * \varphi - h_0)(r) &= - \int_{\mathbb{R}} 4 \sin^2(\pi t T) \hat{F}(t) \exp(2\pi i t r) dt \\
 &= \int_{\mathbb{R}} 2(\cos(2\pi t T) - 1) \hat{F}(t) \exp(2\pi i t r) dt \\
 &= F(r - T) + F(r + T) + 2F(r),
 \end{aligned}$$

as claimed. \square

Lemma 4.3. *Let φ be as in Lemma 4.1. Assume that $\hat{\varphi}$ has compact support, so that φ extends to an entire function, and set*

$$V(r) = \int_{-\infty}^r \int_{-\infty}^y \varphi(x) dx dy \quad \text{for } r \in \mathbb{C}.$$

Then, for any $T \in \mathbb{R}$, we have

$$C\left(\frac{\cos(2\pi T t)}{2(\pi t)^2} \hat{\varphi}(t)\right) = 2V(-T) - 2\Re V\left(\frac{i}{2} - T\right),$$

where $C(\cdot)$ is as defined in Proposition 2.1.

Proof. Define $f(T) := C\left(\frac{\cos(2\pi T t)}{2(\pi t)^2} \hat{\varphi}(t)\right)$. We first observe that $\frac{\cosh(\pi t) - 1}{2(\pi t)^2} \hat{\varphi}(t)$ is absolutely integrable, so by the Riemann–Lebesgue lemma, $\lim_{T \rightarrow \infty} f(T) = 0$. Now differentiating twice with respect to T gives us

$$\begin{aligned}
 f''(T) &= -2 \int_{\mathbb{R}} [\cosh(\pi t) - 1] \hat{\varphi}(t) \cos(2\pi T t) dt \\
 &= 2 \int_{\mathbb{R}} \hat{\varphi}(t) \cos(2\pi T t) dt - 2 \int_{\mathbb{R}} \cosh(\pi t) \hat{\varphi}(t) \cos(2\pi T t) dt \\
 &= \varphi(T) + \varphi(-T) - \frac{\varphi\left(\frac{i}{2} + T\right) + \varphi\left(\frac{i}{2} - T\right) + \varphi\left(-\frac{i}{2} + T\right) + \varphi\left(-\frac{i}{2} - T\right)}{2}.
 \end{aligned}$$

On the other hand, if we start from

$$g(T) := V(T) + V(-T) - \frac{V\left(\frac{i}{2} + T\right) + V\left(\frac{i}{2} - T\right) + V\left(-\frac{i}{2} + T\right) + V\left(-\frac{i}{2} - T\right)}{2}$$

and differentiate twice with respect to T , we get $g''(T) = f''(T)$. Furthermore, Lemma 4.1 shows that $V(r) - \max(0, r)$ is even, and we get

$$\begin{aligned}
 g(T) &= 2V(-T) + T - \Re V\left(\frac{i}{2} - T\right) - \Re V\left(\frac{i}{2} + T\right) \\
 &= 2V(-T) - 2\Re V\left(\frac{i}{2} - T\right),
 \end{aligned}$$

and this vanishes in the limit as $T \rightarrow \infty$. Therefore, $f = g$. \square

Proposition 4.4. *Let φ , F and V be as in Lemmas 4.1–4.3, and assume that $F(r) \geq 0$ for $r \in \mathbb{R}$. Let $\beta : \mathbb{R} \rightarrow \mathbb{R}$ be an even C^2 function of compact support with $\beta(0) = 1$, and let h_2 be the continuous function with Fourier transform*

$$\hat{h}_2(t) = \frac{1 - \beta(t)}{2(\pi t)^2}.$$

Define

$$B = R(h_2) + (D - C) \left(\frac{\beta(t)}{2(\pi t)^2} \right)$$

and

$$k(r) = \frac{r \tanh(\pi r)}{12} + \frac{\frac{1}{8} + \frac{1}{3\sqrt{3}} \cosh\left(\frac{\pi r}{3}\right)}{\cosh(\pi r)} + \frac{\log(2\pi) - 2\Re\psi(1 + 2ir)}{2\pi}.$$

Then for $T \geq 4$ we have

$$\begin{aligned} \int_0^T S(t) dt &\leq \int_{\mathbb{R}} [k(T+r) + k(T-r)] F(r) dr - D \left(\frac{\cos(2\pi Tt)}{2(\pi t)^2} \varphi(t) \right) + B + C_0 \\ &\quad + \frac{\log T}{24\pi} - 2\Re V \left(\frac{i}{2} - T \right), \end{aligned}$$

with the constant C_0 defined as in Proposition 3.1 above.

Proof. We wish to derive an upper bound for $\text{Tr}(h_0)$. As established by Proposition 3.1, we can handle $M(h_0)$, and we now proceed by considering $R(h_0) = R(h_1) + R(h_0 - h_1)$, where h_1 majorizes h_0 and has a trace we can actually compute.

By hypothesis and Lemma 4.1, F is even and non-negative, so if we set $h_1 = h_0 * \varphi + 2F$ then h_1 majorizes h_0 . Further, by Lemma 4.2, we have

$$(h_1 - h_0)(r) = F(r - T) + F(r + T).$$

Thus we can write

$$\begin{aligned} R(h_0 - h_1) &= M(h_1 - h_0) - \text{Tr}(h_1 - h_0) \leq M(h_1 - h_0) \\ &= \int_{\mathbb{R}} k(r) [F(r - T) + F(r + T)] dr - 2F(T). \end{aligned}$$

Returning to h_1 , by Lemma 4.1 we have

$$\hat{h}_1(t) = \hat{h}_0(t)\hat{\varphi}(t) + 2\hat{F}(t) = \frac{\sin^2(\pi Tt)}{(\pi t)^2} \hat{\varphi}(t) + \frac{1 - \hat{\varphi}(t)}{2(\pi t)^2} = \frac{1}{2(\pi t)^2} - \frac{\cos(2\pi Tt)}{2(\pi t)^2} \hat{\varphi}(t).$$

We are almost there, but we must cater for the $\frac{1}{2(\pi t)^2}$ term. To that end, we have

$$\begin{aligned} R(h_1) &= R(h_2) + R(h_1 - h_2) = R(h_2) + (D - C) \left[\frac{\beta(t)}{2(\pi t)^2} - \frac{\cos(2\pi Tt)}{2(\pi t)^2} \hat{\varphi}(t) \right] \\ &= (C - D) \left[\frac{\cos(2\pi Tt)}{2(\pi t)^2} \hat{\varphi}(t) \right] + B, \end{aligned}$$

where B is the constant defined in the statement of the theorem. Finally, Lemma 4.3 gives the formula for the continuous part. \square

4.1. *Bound for large T .* For relatively small T , we can engineer things so that computing the discrete term is tractable. For large T , we use the following:

Proposition 4.5. *Let the notation be as in Proposition 4.4 and assume that $\hat{\varphi}(t) \geq 0$ for $t \in \mathbb{R}$. Then for $T \geq 4$,*

$$\begin{aligned} \int_0^T S(t) dt &\leq \int_{\mathbb{R}} [k(T+r) + k(T-r) + 2k(r)] F(r) dr \\ &\quad - 2\Re \left(V\left(\frac{i}{2}\right) + V\left(\frac{i}{2} - T\right) \right) + 2B + C_0 + \frac{\log T}{24\pi}. \end{aligned} \quad (4.2)$$

Proof. We bound the contribution from the discrete part trivially, using the assumption that $\hat{\varphi}$ is non-negative:

$$\begin{aligned} -D \left(\frac{\cos(2\pi Tt)}{2(\pi t)^2} \hat{\varphi}(t) \right) &\leq D \left(\frac{\hat{\varphi}(t)}{2(\pi t)^2} \right) \\ &= R(h_2 - 2F) + D \left(\frac{\beta(t)}{2(\pi t)^2} \right) + C \left(\frac{\hat{\varphi}(t) - \beta(t)}{2(\pi t)^2} \right) \\ &= -2R(F) + B + \int_{\mathbb{R}} \frac{\cosh(\pi t) - 1}{2(\pi t)^2} \hat{\varphi}(t) dt \\ &\leq 2M(F) + B + \int_{\mathbb{R}} \frac{\cosh(\pi t) - 1}{2(\pi t)^2} \hat{\varphi}(t) dt. \end{aligned}$$

Lemma 4.3 with $T = 0$ gives

$$\int_{\mathbb{R}} \frac{\cosh(\pi t) - 1}{2(\pi t)^2} \hat{\varphi}(t) dt = 2V(0) - 2\Re V\left(\frac{i}{2}\right),$$

and we are done. \square

5. Bounding the Constant B

The first task is to derive a rigorous bound for the constant B , which we recall is defined via

$$B = R(h_2) + (D - C) \left(\frac{\beta(t)}{2(\pi t)^2} \right).$$

5.1. *A suitable form for β .* Let $a, b, c > 0$ and define

$$\hat{\beta}(r) = c \operatorname{sinc}(\pi ar)^8 (b^2 - r^2). \quad (5.1)$$

Choosing c to make $\beta(0) = 1$ now puts β in the required form, with support $[-4a, 4a]$. Also, we have

Lemma 5.1. *Let $\hat{\beta}(r)$ be as defined above. Write*

$$k = \frac{c}{10080\pi^2 a^8}.$$

Then for $t \in [0, a)$ we have

$$\begin{aligned} \beta(t) = & k[\pi^2 b^2(4832a^7 - 3360a^5 t^2 + 1120a^3 t^4 - 280at^6 + 70t^7) \\ & - 1680a^5 + 3360a^3 t^2 - 2100at^4 + 735t^5]; \end{aligned}$$

for $t \in [a, 2a)$ we have

$$\begin{aligned} \beta(t) = & k[\pi^2 b^2(4944a^7 - 784a^6 t - 1008a^5 t^2 - 3920a^4 t^3 + 5040a^3 t^4 \\ & - 2352a^2 t^5 + 504at^6 - 42t^7) \\ & - 504a^5 - 5880a^4 t + 15120a^3 t^2 - 11760a^2 t^3 + 3780at^4 - 441t^5]; \end{aligned}$$

for $t \in [2a, 3a)$ we have

$$\begin{aligned} \beta(t) = & -k[\pi^2 b^2(2224a^7 - 24304a^6 t + 38640a^5 t^2 - 27440a^4 t^3 + 10640a^3 t^4 \\ & - 2352a^2 t^5 + 280at^6 - 14t^7) \\ & + 19320a^5 - 41160a^4 t + 31920a^3 t^2 - 11760a^2 t^3 + 2100at^4 - 147t^5]; \end{aligned}$$

for $t \in [3a, 4a)$ we have

$$\begin{aligned} \beta(t) = & k[\pi^2 b^2(32768a^7 - 57344a^6 t + 43008a^5 t^2 - 17920a^4 t^3 + 4480a^3 t^4 \\ & - 672a^2 t^5 + 56at^6 - 2t^7) \\ & + 21504a^5 - 26880a^4 t + 13440a^3 t^2 - 3360a^2 t^3 + 420at^4 - 21t^5]; \end{aligned}$$

and for $t \in [4a, \infty)$ we have $\beta(t) = 0$.

Proof. This is a messy but straightforward application of known Fourier transforms. \square

A priori, there is no guarantee that Strömbergsson's list of r_j is complete. However, we will choose b so that there are no unknown $r_j \leq b$ (see Sect. 5.3 below). This makes $h_2(r_j)$ non-positive for any unknown r_j , so that

$$R(h_2) = \text{Tr}^*(h_2) + \text{Tr}^\dagger(h_2) - M(h_2)$$

where Tr^* is the trace over the known r_j and Tr^\dagger is the trace over the rest. Thus

$$B \leq \text{Tr}^*(h_2) - M(h_2) + (D - C) \left(\frac{\beta(t)}{2(\pi t)^2} \right).$$

5.2. *Procedure.* We aim to rigorously compute

$$\text{Tr}^*(h_2) - M(h_2) + (D - C) \left(\frac{\beta(t)}{2(\pi t)^2} \right).$$

5.2.1. *Computing $\text{Tr}^*(h_2)$* We have

$$h_2(r) = \int_{\mathbb{R}} \hat{h}_2(t) \cos(2\pi tr) dt = 2 \int_0^{4a} \frac{1 - \beta(t)}{2(\pi t)^2} \cos(2\pi tr) dt + 2 \int_{4a}^{\infty} \frac{\cos(2\pi tr)}{2(\pi t)^2} dt.$$

The first integral we compute numerically for each r_j in our database using Theorem A.1. The second integral becomes

$$\frac{2r}{\pi} \left[\text{Si}(8a\pi r) - \frac{\pi}{2} + \frac{\cos(8a\pi r)}{8a\pi r} \right],$$

which again we compute for each of our known r_j . The function Si above is the sine integral

$$\text{Si}(x) = \int_0^x \frac{\sin y}{y} dy.$$

5.2.2. *Computing $I(h_2)$* We have

$$I(h_2) = -\frac{1}{12\pi} \int_0^{\infty} \frac{\hat{h}'_2(t)}{\sinh(\pi t)} dt.$$

When t is small, computing the x_k for Theorem A.1 will produce an interval that straddles zero. To avoid this, we work instead with $\frac{\hat{h}'_2(t)}{t}$ and truncate the Taylor expansion of $\frac{t}{\sinh(\pi t)}$ after a few terms. The following lemma provides an error bound.

Lemma 5.2. *Let $N \geq 0$ and $|t| < \sqrt{(2N+4)(2N+5)}$. Then*

$$\left| \frac{\sinh t}{t} - \sum_{n=0}^N \frac{t^{2n}}{(2n+1)!} \right| \leq \frac{t^{2N+2}(2N+4)(2N+5)}{(2N+3)!((2N+4)(2N+5) - t^2)}.$$

Proof. We majorize the tail of the Taylor expansion with the obvious geometric series. \square

We can thus compute the integral over the intervals $[0, a]$, $[a, 2a]$, $[2a, 3a]$ and $[3a, 4a]$ with little difficulty. This leaves computing the integral over $[4a, \infty)$ which we will truncate at some $t_0 > 4a$ using the following lemma:

Lemma 5.3. *For $t_0 > 4a$ we have*

$$\left| \int_{t_0}^{\infty} \frac{\hat{h}'_2(t)}{\sinh(\pi t)} dt \right| \leq \left| \frac{1}{\pi^3 t_0^3} \log \tanh\left(\frac{\pi t_0}{2}\right) \right|.$$

Proof. With $t \geq t_0 > 4a$ we have

$$\hat{h}'_2(t_0) = -\frac{1}{\pi^2 t_0^3},$$

and $\log \tanh(t/2)$ is an anti-derivative of $1/\sinh(t)$. \square

Computing

$$\int_{4a}^{t_0} \frac{1}{2(\pi t)^2 \sinh(\pi t)} dt$$

via Theorem A.1 still requires some care because of the pole at $t = 0$. As discussed in Appendix A we can sidestep this by taking the integrals over $[4a\alpha^n, 4a\alpha^{n+1}]$ for $n = 0, \dots, \lceil \log_{\alpha}(t_0/4a) \rceil$ for some $\alpha \in (1, 3)$. We took $\alpha = 3 - \frac{1}{128}$ and $t_0 = 4a\alpha^4$.

5.2.3. *Computing $E(h_2)$* This is more straightforward. We need only the following auxiliary lemma.

Lemma 5.4. *Let $t_0 > 4a$. Then*

$$\int_{t_0}^{\infty} \left(\frac{1}{8 \cosh \pi t} + \frac{2 \cosh \pi t}{3 + 6 \cosh 2\pi t} \right) \hat{h}_2(t) dt \leq \frac{7 \exp(-\pi t_0)}{24\pi^3 t_0^2} dt.$$

Proof. We use the trivial estimate

$$\int_{t_0}^{\infty} \left(\frac{1}{8 \cosh \pi t} + \frac{2 \cosh \pi t}{3 + 6 \cosh 2\pi t} \right) \hat{h}_2(t) dt \leq \int_{t_0}^{\infty} \left(\frac{\exp(-\pi t)}{4} + \frac{\exp(-\pi t)}{3} \right) \frac{1}{2\pi^2 t_0^2} dt.$$

□

We again do four integrations covering $4a$ to $4a \left(3 - \frac{1}{128}\right)^4$ then use Lemma 5.4 to bound the tail.

5.2.4. *Computing $P(h_2)$* As a reminder, we have

$$P(h) = \frac{\hat{h}(0)}{2\pi} \left(\log \frac{\pi}{2} + 2\gamma \right) - \frac{h(0)}{4} - \frac{1}{\pi} \int_0^{\infty} \log \left(4 \sinh \left(\frac{\pi t}{2} \right) \right) \hat{h}'(t) dt.$$

Using Maple™ we get

$$\hat{h}_2(0) = c \frac{\pi^2 a^2 b^2 - 1}{6a^5 \pi^4}$$

where c is the constant in equation 5.1. To compute $h_2(0) = \int_{\mathbb{R}} \hat{h}_2(t) dt$ we estimate the integrals over $[0, a]$, $[a, 2a]$, $[2a, 3a]$ and $[3a, 4a]$ using Theorem A.1. We then have

$$\int_{4a}^{\infty} \hat{h}_2(t) dt = \frac{1}{8a\pi^2}.$$

To compute

$$\int_0^a \log(4 \sinh(\pi t/2)) \hat{h}'_2(t) dt$$

we must handle the singularity at $t = 0$. We compute

$$\int_0^a (\log(4 \sinh(\pi t/2)) - \log t) \hat{h}'_2(t) dt + \int_0^a (\log t) \hat{h}'_2(t) dt,$$

where the second integral can be computed analytically to yield

$$\int_0^a (\log t) \hat{h}'_2(t) dt = c \frac{(195 - 130\pi^2 a^2 b^2) \log a + 72\pi^2 a^2 b^2 - 115}{2880a^5 \pi^4}.$$

Once past $4a$ we compute the three further integrals

$$\sum_{n=0}^2 \int_{4a\alpha^n}^{4a\alpha^{n+1}} \frac{\log \left(4 \sinh \left(\frac{\pi t}{2} \right) \right)}{\pi^2 t^3} dt$$

and then bound the tail using the following lemma.

Lemma 5.5. *Let $t_0 > 0$. Then*

$$\left| \int_{t_0}^{\infty} \log\left(4 \sinh\left(\frac{\pi t}{2}\right)\right) \hat{h}'_2(t) dt + \left(\frac{\pi t_0}{2} + \log 2\right) \hat{h}_2(t_0) + \frac{\pi}{2} \int_{t_0}^{\infty} \hat{h}_2(t) dt \right| \leq \log(1 - \exp(-\pi t_0)) \hat{h}_2(t_0).$$

Proof. We use

$$\log(4 \sinh(t)) = \log 2 + t + \log(1 - \exp(-2t))$$

and integration by parts. \square

5.2.5. *Computing $D\left(\frac{\beta(t)}{2\pi^2 t^2}\right)$* This is straightforward other than the evaluations of $L(1, \chi_d)$. To this end we produced a database containing class numbers $h(d)$ and fundamental units (u and v such that $u^2 - dv^2 = \pm 4$ with $u, v \in \mathbb{Z}_{>0}$ and v minimal). The database covered each fundamental discriminant d such that $dI^2 = t^2 - 4$ with $t \in [3, 10^5]$. We used the PARI function `qfbclassno`¹ to compute the class numbers [4] and the PQA algorithm due to Lagrange [24] to compute the fundamental units. This then allows us to compute $L(1, \chi_d)$ for this range of d rapidly and rigorously using Dirichlet's class number formula

$$L(1, \chi_d) = \frac{2h(d) \log\left(\frac{u+v\sqrt{d}}{2}\right)}{\sqrt{d}}.$$

Since $\beta(t)$ is zero for $t \notin [-4a, 4a]$, we can take

$$a = \frac{1}{4\pi} \log\left(\frac{10^5 + \sqrt{10^{10} - 4}}{2}\right) = 0.916169\dots$$

without running out of pre-computed class group data. We also need

$$\frac{1}{\pi} \sum_{n=1}^{\infty} \frac{\Lambda(n)}{n} h\left(\frac{\log n}{\pi}\right),$$

so for this value of a we need to sum over primes and prime powers $\leq 99\,991$.

5.2.6. *Computing $C\left(\frac{\beta(t)}{2\pi^2 t^2}\right)$* This reduces to the finite integral

$$\int_{-4a}^{4a} \frac{(\cosh(\pi t) - 1)\beta(t)}{2\pi^2 t^2} dt.$$

The only issue here is the computation when t is small, when we resort to the series expansion of $t^{-2}(\cosh(t) - 1)$ using the following lemma.

Lemma 5.6. *Let $N \geq 0$ and $|t| < \sqrt{(2N+5)(2N+6)}$. Then*

$$\left| t^{-2}(\cosh(t) - 1) - \sum_{n=0}^N \frac{t^{2n}}{(2n+2)!} \right| \leq \frac{t^{2N+2}(2N+5)(2N+6)}{((2N+5)(2N+6) - t^2)(2N+4)!}.$$

Proof. We majorize the tail of the series with the obvious geometric series. \square

¹ `quadclassunit` is faster but the correctness of its results is conditional on GRH.

5.3. Results.

5.3.1. $b = \frac{\sqrt{6\pi^2-1}}{2}$ We will initially take $b = \frac{\sqrt{6\pi^2-1}}{2}$ since by [17, Theorem 11.4] there are no r_j below this. Setting $a = \frac{7505}{8192}$ and using no r_j (so $\text{Tr}^*(h_2) = 0$) we get

$$B \leq 0.272955804771976.$$

5.3.2. $b = 177.75$ Once we have verified Corollary 1.2 using the above bound for B (see Sect. 6.3 below), we can take $b = 177.75$ and run the computation again. This time we get

$$B \leq 0.2729558044747431.$$

The computation takes about twenty minutes on a single core, the time being dominated by computing the trace of the known $r_j \leq 177.75$.

5.4. *Improving the bound for B further.* Even though the bound on B derived above will more than suffice for our immediate needs, it is possible to improve on it further. This could be done by increasing the exponent 8 used in the definition of β (equation (5.1)). The true value of B is more like

$$B = 0.2729558044747424323066650413 \dots,$$

but we do not prove that here.

6. Verifying Theorem 1.1

Our verification proceeds in four stages. We will first prove the theorem for $T \in [100, 27\,400]$ via interval arithmetic and use this to verify Corollary 1.2. We will then rigorously verify that Theorem 1.1 holds for $T \in (1, 100]$ by direct computation. Next we will check, again using interval arithmetic, that the theorem holds for $T \in [27\,400, 10^6]$. Finally we will show that it holds for all $T \geq 10^6$.

First, however, we must define the function φ used in Proposition 4.4.

6.1. *A suitable form for $\hat{\varphi}$.* In [21], Littman shows that the unique best majorant of exponential type 2π to $\max(0, r)$ is

$$V_0(r) = \frac{\cos^2(\pi r)}{\pi^2} \left[r\psi' \left(\frac{1}{2} - r \right) + 1 \right].$$

This would suggest that we set $\varphi(r) = X\varphi_0(Xr)$, where $\varphi_0(r) = V_0''(r)$, with Fourier transform

$$\begin{aligned} \hat{\varphi}_0(t) &= \left[1 - \sum_{k=0}^{\infty} \frac{B_{k+2} \left(\frac{1}{2} \right)}{(k+1)!} \left(\frac{k+1}{k+2} - |t| \right) (2\pi i t)^{k+2} \right] \chi_{(-1,1)}(t) \\ &= \left(\frac{|t|}{\text{sinc}(\pi t)} + \frac{(1-|t|)\cos \pi t}{\text{sinc}^2(\pi t)} \right) \chi_{(-1,1)}(t), \end{aligned} \tag{6.1}$$

where $\chi_{(-1,1)}$ denotes the characteristic function of $(-1, 1)$. Unfortunately, this fails to be C^4 near $t = 0$, which makes the integral on the right-hand side of (4.2) diverge. We add the following correction term which will allow us to patch things up:

Lemma 6.1. *Let $X, \delta > 0$, and define*

$$\hat{\eta}_0(t) = \frac{\pi^2}{4 + \pi^2} \left[\frac{2\pi^2}{3}(1 - |t|)^3 + 4(1 - |t|)(1 - \cos(\pi t)) - \frac{8}{\pi} \sin(\pi |t|) \right] \chi_{(-1,1)}(t)$$

and

$$\hat{V}_1(t) = -\frac{\delta}{48\pi^2 X^3} \left[2\hat{\eta}_0\left(\frac{t}{\delta}\right) + \hat{\eta}_0\left(\frac{t+X}{\delta}\right) + \hat{\eta}_0\left(\frac{t-X}{\delta}\right) \right].$$

Then

$$V_1(r) = \frac{\cos^2(\pi X r)}{\pi^2 X} \left[-\frac{1}{12X^2 r^2} + \frac{2 \operatorname{sinc}^2(\pi \delta r) + \operatorname{sinc}^2\left(\pi \delta r + \frac{\pi}{2}\right) + \operatorname{sinc}^2\left(\pi \delta r - \frac{\pi}{2}\right)}{24\left(1 + \frac{4}{\pi^2}\right)X^2 r^2} \right].$$

Proof. We have

$$\eta_0(r) = \frac{1}{r^2} - \frac{2 \operatorname{sinc}^2(\pi r) + \operatorname{sinc}^2\left(\pi r + \frac{\pi}{2}\right) + \operatorname{sinc}^2\left(\pi r - \frac{\pi}{2}\right)}{2\left(1 + \frac{4}{\pi^2}\right)r^2},$$

and if we set

$$V_1(r) = -\frac{\delta^2 \cos^2(\pi X r)}{12\pi^2 X^3} \eta_0(\delta r) = -\frac{\delta^2 (\cos(2\pi X r) + 1)}{24\pi^2 X^3} \eta_0(\delta r)$$

then

$$\hat{V}_1(t) = -\frac{\delta}{48\pi^2 X^3} \left[2\hat{\eta}_0\left(\frac{t}{\delta}\right) + \hat{\eta}_0\left(\frac{t+X}{\delta}\right) + \hat{\eta}_0\left(\frac{t-X}{\delta}\right) \right],$$

as required. \square

We will need to bound the term involving the trigamma function. The following will suffice.

Lemma 6.2. *For $z > 0$ we have*

$$-z\psi'\left(\frac{1}{2} + z\right) + 1 - \frac{1}{12z^2} \in \left[-\frac{7}{120z^4}, 0 \right],$$

and for $z \in \mathbb{R} \setminus \{0\}$,

$$\left| -iz\psi'\left(\frac{1}{2} + iz\right) + 1 + \frac{1}{12z^2} \right| \leq \frac{112 + 105\pi}{3840z^4}.$$

Also, for $z = \sigma + it$ with $\sigma > 0$ and $t \in \mathbb{R}$, we have

$$\left| -z\psi'\left(\frac{1}{2} + z\right) + 1 - \frac{1}{12z^2} \right| \leq \frac{7(\sigma + |t|)}{120\sigma^5}.$$

Proof. We use the multiplication formula to write

$$\psi' \left(\frac{1}{2} + z \right) = 4\zeta(2, 2z) - \zeta(2, z)$$

for $-z \notin \mathbb{Z}_{\geq 0}$ where $\zeta(s, z)$ is the Hurwitz zeta-function. We now apply Euler–Maclaurin summation with

$$g(t) = 4(t + 2z)^{-2} - (t + z)^{-2}$$

to get

$$\sum_{n=0}^N g(n) + \int_N^\infty g(t) dt - \frac{g(N)}{2} - \sum_{k=1}^K \frac{B_{2k} g^{(2k-1)}(N)}{(2k)!} - \int_N^\infty \frac{B_{2K}(\{t\})}{(2K)!} g^{(2K)}(t) dt.$$

We now take $N = 0$ and $K = 2$. In the case $z > 0$ the $k = 2$ term is $-\frac{7}{240z^5}$ and the final integral is trivially less in absolute terms than that. In the second case, the $k = 2$ term is bounded by $\frac{7}{240|z|^5}$ and the final integral contributes less than

$$\frac{B_4}{24} \int_0^\infty \left(\frac{480}{(t^2 + 4z^2)^3} - \frac{120}{(t^2 + z^2)^3} \right) dt = \frac{7\pi}{256|z|^5}.$$

In the final case, the $k = 2$ term yields $\frac{7}{240|z|^5}$. The integral can be bounded by $\frac{7}{240\sigma^5}$ and the result follows. \square

We bound the sinc terms using [7, Lemma A.2]:

Lemma 6.3. *Let $r, \delta > 0$ with $\delta^2 r^2 \geq \frac{1}{12}$. Then*

$$2 \operatorname{sinc}^2(\pi \delta r) + \operatorname{sinc}^2\left(\pi \delta r + \frac{\pi}{2}\right) + \operatorname{sinc}^2\left(\pi \delta r - \frac{\pi}{2}\right) \geq \frac{2}{\pi^2 \delta^2 r^2}.$$

We can now combine the above to get:

Proposition 6.4. *Let $\delta > 0$, $X \geq \delta \sqrt{\frac{7}{20} \left(1 + \frac{4}{\pi^2}\right)}$. Take $\hat{\varphi}_0$ and $\hat{\eta}_0$ to be as defined in equation (6.1) and Lemma 6.1, respectively. Now define*

$$\hat{\varphi}(t) = \hat{\varphi}_0\left(\frac{t}{X}\right) + \frac{\delta t^2}{12X^3} \left[2\hat{\eta}_0\left(\frac{t}{\delta}\right) + \hat{\eta}_0\left(\frac{t+X}{\delta}\right) + \hat{\eta}_0\left(\frac{t-X}{\delta}\right) \right].$$

Then $\hat{\varphi}$ meets all of the hypotheses of Propositions 4.4 and 4.5. In particular,

$$V(r) = \frac{\cos^2(\pi Xr)}{\pi^2 X} \left[Xr \psi' \left(\frac{1}{2} - Xr \right) + 1 - \frac{1}{12X^2 r^2} + \frac{2 \operatorname{sinc}^2(\pi \delta r) + \operatorname{sinc}^2(\pi \delta r + \frac{\pi}{2}) + \operatorname{sinc}^2(\pi \delta r - \frac{\pi}{2})}{24(1 + \frac{4}{\pi^2})X^2 r^2} \right] \quad (6.2)$$

is non-negative for $r \in \mathbb{R}$.

Proof. First we have the identity

$$\cos^2(\pi z) \left[\psi' \left(\frac{1}{2} + z \right) + \psi' \left(\frac{1}{2} - z \right) \right] = \pi^2$$

which implies $V(z) = V(-z) + z$. Thus to establish positivity, we need only consider

$$\frac{\cos^2(\pi Xr)}{\pi^2 X} \left[-Xr \psi' \left(\frac{1}{2} + Xr \right) + 1 - \frac{1}{12X^2r^2} + \frac{2 \operatorname{sinc}^2(\pi \delta r) + \operatorname{sinc}^2(\pi \delta r + \frac{\pi}{2}) + \operatorname{sinc}^2(\pi \delta r - \frac{\pi}{2})}{24(1 + \frac{4}{\pi^2})X^2r^2} \right]$$

for $r > 0$. Now by Lemmas 6.2 and 6.3 we get, for $\delta^2 r^2 \geq \frac{1}{12}$,

$$V(-r) \geq \frac{\cos^2(\pi Xr)}{\pi^2 X} \left[\frac{1}{12(\pi^2 + 4)X^2\delta^2r^4} - \frac{7}{120X^4r^4} \right],$$

which is non-negative providing $X \geq \delta \sqrt{\frac{7}{20}(\pi^2 + 4)}$, as claimed. For $\delta^2 r^2 < \frac{1}{12}$ we consider

$$-(Xr)^3 \psi' \left(\frac{1}{2} + Xr \right) + (Xr)^2 + \frac{2 \operatorname{sinc}^2(\pi \delta r) + \operatorname{sinc}^2(\pi \delta r + \frac{\pi}{2}) + \operatorname{sinc}^2(\pi \delta r - \frac{\pi}{2})}{24(1 + \frac{4}{\pi^2})} - \frac{1}{12}.$$

Since $-t^3 \psi' \left(\frac{1}{2} + t \right) + t^2$ is an increasing function of $t > 0$, it suffices to consider $X = c\delta$, where $c = \sqrt{\frac{7}{20}(\pi^2 + 4)}$. Dividing by $(Xr)^2$ and writing $\delta r = t$, we obtain

$$-ct \psi' \left(\frac{1}{2} + ct \right) + 1 + \frac{f(t) - f(0)}{24(1 + \frac{4}{\pi^2})c^2 t^2},$$

where $f(t) = 2 \operatorname{sinc}^2(\pi t) + \operatorname{sinc}^2(\pi t + \frac{\pi}{2}) + \operatorname{sinc}^2(\pi t - \frac{\pi}{2})$.

By the Lagrange form of the error in Taylor's theorem, $(f(t) - f(0))/t^2 = \frac{1}{2} f''(u)$ for some $u \in [0, t]$. We verify via interval arithmetic that $\frac{1}{2} f''(t) > -5$ for $t \in [0, 1/\sqrt{12}]$ and that

$$-ct \psi' \left(\frac{1}{2} + ct \right) + 1 - \frac{5}{24(1 + \frac{4}{\pi^2})c^2} > 0$$

over the same interval. \square

6.2. *Verifying Theorem 1.1 for $T \in [100, 27\,400]$.* For T in this range, we proceed computationally. We chose $\delta = 0.1$ and $X = 2.55$ and kept them fixed throughout. If we computed all the terms of Proposition 4.4 using interval arithmetic, we would be forced to use a relatively narrow width, and the cost of computing

$$\int_{\mathbb{R}} [k(T+r) + k(T-r)]F(r) dr$$

using Theorem A.1 would be prohibitive. Fortunately, the part which determines the maximum width we can get away with is the $\cos(2\pi tT)$ term within $D(\cdot)$. The rest, including the expensive integrals, are much less susceptible, and we can use a width of between 2 and 128 depending on the size of T .

Having computed everything else with a relatively coarse interval, we then compute $D(\cdot)$ many times with a much narrower interval (typically about 2^{-8}). Even here, most of the terms making up $D(\cdot)$ do not depend on T so they can be pre-computed once and stored.

The entire computation coded in C++ using interval arithmetic takes less than an hour on a 16-core node of Bluecrystal Phase III [1]. Despite the foregoing, this time was still dominated by the rigorous computation of

$$\int_{\mathbb{R}} [k(T+r) + k(T-r)]F(r) dr.$$

Lemma 6.5. *The computation shows that Theorem 1.1 holds for $T \in [100, 27\,400]$.*

6.3. *Verifying Corollary 1.2.* We now aim to partially validate our database. Writing $\bar{N}(t) = \frac{t^2}{12} - \frac{2t}{\pi} \log \frac{t}{e\sqrt{\frac{\pi}{2}}} - \frac{131}{144}$, we compute using interval arithmetic

$$\int_0^{178} \bar{N}(t) dt \in (121\,643.023\,932, 121\,643.023\,933).$$

Using Theorem 1.1, we compute

$$\int_0^{178} S(t) dt < 0.398\,780$$

so we have

$$\int_0^{178} N(t) dt < 121\,643.422\,713.$$

Also, using our (possibly incomplete) database of r_j we can compute

$$\int_0^{178} N^-(t) dt,$$

where N^- is a minorant of N . Specifically, we have

$$\int_0^{178} N^-(t) dt > 121\,643.206\,595.$$

Now if an r_j were missing anywhere in $(0, 178 - (121\,643.422\,713 - 121\,643.206\,595)) = (0, 177.783\,882)$ then our revised value for $\int N^-(t) dt$ would exceed our upper bound for $\int N(t) dt$, a contradiction.

6.4. *Verifying Theorem 1.1 for $T \in (1, 100]$.* Once we are satisfied that our database of zeros is complete, we can use it to confirm Theorem 1.1 for small T computationally.

Lemma 6.6. *Theorem 1.1 holds for $T \in (1, 100]$.*

Proof. The lemma holds trivially for $T \in (1, r_1)$ where $r_1 = 9.533\dots$ is the location of the first zero as $\int_0^T N(t) dt = 0$ and $\int_0^T \bar{N}(t) dt + T \cdot E(T)$ is positive. Letting N^+ be the majorant of N given by our list of r_j , we check that

$$\int_0^T (N^+(t) - \bar{N}(t)) dt - T \cdot E(T) < 0$$

at each r_j and divide the interval between the r_j 's into small enough sub-intervals so that the inequality works there too. In all, we checked a little under 500 000 intervals covering $(r_1, 100]$ in a little under 10 seconds to verify the theorem. The nearest miss in absolute terms was around $T = 20.6862978$, where

$$\int_0^T (N(t) - \bar{N}(t)) dt - T \cdot E(T) = -9.8826792\dots \times 10^{-8}.$$

□

6.5. *Verifying Theorem 1.1 for $T \in [27\,400, 10^6]$.* Once T is large enough, we can dispense with $D(\cdot)$ and appeal to Proposition 4.5 instead. We start with $T \in [27\,400, 27\,402]$ and check that Theorem 1.1 holds, and then move on to the next interval for T . After every 20 iterations, we try to double the width of the interval and continue. If at any point we fail (presumably because our interval for T grew too wide too quickly) we halve the width of the interval and repeat. Coded in C++ using interval arithmetic, the computation takes less than a couple of hours on a single core (and we could have parallelized it trivially). By the end, the width of the interval had increased to 16 384 and we have:

Lemma 6.7. *Theorem 1.1 holds for $T \in [27\,400, 10^6]$.*

6.6. *Verifying Theorem 1.1 for $T \geq 10^6$.* We now look at each of the contributions to our bound for $\int S(t) dt$ coming from Proposition 4.5. We will need several preparatory lemmas. We assume the notation and hypotheses of Propositions 4.4 and 6.4, and we fix $\delta = 0.842$ throughout.

Lemma 6.8. *We have*

$$\hat{F}(0) = \frac{3(\pi^2 + 4)X - 2\pi^2\delta}{72X^3(\pi^2 + 4)}.$$

Proof. We have

$$\hat{F}(0) = \lim_{t \rightarrow 0} \frac{1 - \hat{\varphi}(t)}{(2\pi t)^2}.$$

Then taking t positive and sufficiently close to zero eliminates the $\hat{\varphi}((t \pm X)/\delta)$ terms and we have

$$\hat{\varphi}(t) = \hat{\varphi}_0\left(\frac{t}{X}\right) + \frac{\delta t^2}{6X^3} \hat{\eta}_0\left(\frac{t}{\delta}\right).$$

Now we have

$$\hat{\varphi}_0\left(\frac{t}{X}\right) = 1 - \frac{1}{6} \frac{\pi^2}{X^2} t^2 + \mathcal{O}(t^3) \quad \text{and} \quad \frac{\delta t^2}{6X^3} \hat{\eta}_0\left(\frac{t}{\delta}\right) = \frac{\pi^4 \delta}{9(\pi^2 + 4)X^3} t^2 + \mathcal{O}(t^3),$$

so that

$$1 - \hat{\varphi}(t) = \frac{\pi^2(3(\pi^2 + 4)X - 2\pi^2\delta)}{18(\pi^2 + 4)X^3} t^2 + \mathcal{O}(t^3).$$

The result follows on taking $t \rightarrow 0$. \square

Lemma 6.9. For $r \in \mathbb{R} \setminus \{0\}$, we have

$$F(r) \leq \frac{1}{400X^3r^4}.$$

Proof. We use Lemma 6.2 and the estimate

$$2 \operatorname{sinc}^2(\pi\delta r) + \operatorname{sinc}^2(\pi\delta r + \pi/2) + \operatorname{sinc}^2(\pi\delta r - \pi/2) \leq \frac{24(\pi^2 + 4)}{400r^2}.$$

\square

Lemma 6.10. We have

$$k(r) \leq \begin{cases} \frac{4}{5} & \text{if } |r| < 2, \\ \frac{|r|}{12} & \text{if } |r| \geq 2. \end{cases}$$

Proof. Since k is an even function of r , we need only consider $r \geq 0$. We first check that $k(r) \leq \frac{4}{5}$ for $r \in [0, 2]$ by interval arithmetic. Then since

$$\frac{r \tanh(\pi r)}{12} \leq \frac{r}{12}$$

for all $r \geq 0$, we need only check that at $r = 2$

$$\frac{\frac{1}{8} + \frac{1}{3\sqrt{3}} \cosh\left(\frac{\pi r}{3}\right)}{\cosh(\pi r)} + \frac{\log 2\pi - 2\Re\psi(1 + 2ir)}{2\pi} < 0$$

and observe that the $\Re\psi(1 + 2ir)$ term is an increasing function of r whereas the cosh term is decreasing. \square

Lemma 6.11. For $T \geq 10$ we have

$$\int_{\mathbb{R}} [k(T+r) + k(T-r)] F(r) dr \leq \frac{T}{144X^2} - \frac{T\pi^2\delta}{216X^3(\pi^2 + 4)} + \frac{5T^4 + 15T^3 + 258T^2 + 20T + 264}{18000X^3(T-2)^3(T+2)^3}.$$

Proof. Since F is even, we write

$$\int_{\mathbb{R}} [k(T+r) + k(T-r)]F(r) \, dr = 2 \int_{\mathbb{R}} k(T+r)F(r) \, dr$$

and then split the integral into three pieces:

$$\begin{aligned} 2 \int_{-\infty}^{-T-2} k(T+r)F(r) \, dr &\leq 2 \int_{-\infty}^{-T-2} \frac{|T+r|}{12} \frac{dr}{400X^3r^4} \\ &= \frac{1}{3600X^3} \left(\frac{1}{4}(T+2)^{-2} + (T+2)^{-3} \right), \\ 2 \int_{-T-2}^{-T+2} k(T+r)F(r) \, dr &\leq 2 \int_{-T-2}^{-T+2} \frac{4}{5} \cdot \frac{dr}{400X^3r^4} = \frac{2}{375X^3} \frac{3T^2+4}{(T-2)^3(T+2)^3}, \end{aligned}$$

and

$$\begin{aligned} 2 \int_{-T+2}^{\infty} k(T+r)F(r) \, dr &\leq 2 \int_{-T+2}^{\infty} \frac{T+r}{12} F(r) \, dr \\ &= 2 \int_{-T+2}^{\infty} \frac{T}{12} F(r) \, dr + 2 \int_{-T+2}^{\infty} \frac{r}{12} F(r) \, dr \\ &\leq \frac{T}{6} \hat{F}(0) + 2 \int_{T-2}^{\infty} \frac{r}{12} \frac{dr}{400X^3r^4} \\ &= \frac{3(\pi^2+4)X - 2\pi^2\delta}{72X^3(\pi^2+4)} \frac{T}{6} + \frac{1}{4800X^3} (T-2)^{-2}. \end{aligned}$$

□

Lemma 6.12. *We have*

$$2 \int_{\mathbb{R}} k(r)F(r) \, dr \leq \frac{1}{15X^2}.$$

Proof. We have

$$\begin{aligned} 2 \int_{\mathbb{R}} k(r)F(r) \, dr &\leq 2 \int_{-2}^2 \frac{4}{5} F(r) \, dr + 4 \int_2^{\infty} \frac{r}{12} \frac{dr}{400X^3r^4} \\ &\leq \frac{8}{5} \hat{F}(0) + \frac{1}{9600X^3} = \frac{1}{15X^2} - \frac{2\pi^2\delta}{45(\pi^2+4)X^3} + \frac{1}{9600X^3} \\ &\leq \frac{1}{15X^2}. \end{aligned}$$

□

Lemma 6.13. *We have*

$$-2\Re V\left(\frac{i}{2}\right) \leq \cosh^2\left(\frac{\pi X}{2}\right) \left[0.09752X^{-3} + 0.3731X^{-5}\right].$$

Proof. We have

$$\frac{\cos^2\left(\frac{\pi Xi}{2}\right)}{\pi^2 X} = \frac{\cosh^2\left(\frac{\pi X}{2}\right)}{\pi^2 X}$$

and (from Lemma 6.2)

$$\left| -\frac{i}{2} X \psi' \left(\frac{1}{2} + \frac{iX}{2} \right) + 1 - \frac{1}{12X^2(i/2)^2} \right| \leq \frac{112 + 105\pi}{240X^4}.$$

Setting $\delta = 0.842$, the terms of (6.2) involving δ can be computed directly, and this yields the claimed inequality. \square

Lemma 6.14. *We have*

$$|2\Re V(i/2 - T)| \leq \frac{\cosh^2(\pi X/2)}{X^5 T^4} (0.07914 + 0.01183T^{-1}).$$

Proof. For $x, y \in \mathbb{R}$ we have

$$|\cos(x + iy)| \leq \sqrt{2} \cosh(y) \quad \text{and} \quad |\sin(x + iy)| \leq \sqrt{2} \cosh(y),$$

so that

$$\left| \cos^2 \left(\pi X \left(\frac{i}{2} - T \right) \right) \right| \leq 2 \cosh^2 \left(\frac{\pi X}{2} \right).$$

We also have

$$\begin{aligned} & \left| 2 \operatorname{sinc}^2 \left(\pi \delta \left(\frac{i}{2} - T \right) \right) + \operatorname{sinc}^2 \left(\pi \delta \left(\frac{i}{2} - T + \frac{1}{2} \right) \right) + \operatorname{sinc}^2 \left(\pi \delta \left(\frac{i}{2} - T - \frac{1}{2} \right) \right) \right| \\ & \leq \frac{8 \cosh^2 \left(\frac{\pi \delta}{2} \right)}{\pi^2 \delta^2 T^2} \end{aligned}$$

and Lemma 6.2 gives us

$$\left| -\left(XT - \frac{iX}{2} \right) \psi' \left(\frac{1}{2} + XT - \frac{iX}{2} \right) + 1 - \frac{1}{12 \left(XT - \frac{iX}{2} \right)^2} \right| \leq \frac{7}{120X^4 T^4} + \frac{7}{240X^4 T^5}.$$

The result follows on setting $\delta = 0.842$. \square

We can now establish our bound for $T \geq 10^6$.

Lemma 6.15. *Let $S(t)$ and $E(T)$ be as defined in Theorem 1.1. Then for $T \geq 10^6$ we have*

$$\frac{1}{T} \int_0^T S(t) dt - E(T) \leq 0.$$

Proof.

$$\begin{aligned}
\int_0^T S(t) dt &\leq \int_{\mathbb{R}} [k(T+r) + k(T-r) + 2k(r)] F(r) dr - 2\Re \left(V\left(\frac{i}{2}\right) + V\left(\frac{i}{2} - T\right) \right) \\
&\quad + 2B + C_0 + \frac{\log T}{24\pi} \\
&\leq \frac{T}{144X^2} - \frac{T\pi^2\delta}{216(\pi^2+4)X^3} + \frac{5T^4 + 15T^3 + 258T^2 + 20T + 264}{18000X^3(T-2)^3(T+2)^3} + \frac{1}{15X^2} \\
&\quad + 2B + C_0 + \frac{\log T}{24\pi} \\
&\quad + \cosh^2\left(\frac{\pi X}{2}\right) [0.09752X^{-3} + 0.3731X^{-5}] \\
&\quad + \frac{\cosh^2(\pi X/2)}{X^5T^4} [0.07914 + 0.01183T^{-1}].
\end{aligned}$$

Assuming that $T \geq 10^6$, this is at most

$$\begin{aligned}
&\frac{T}{144X^2} - \frac{T\pi^2\delta}{216(\pi^2+4)X^3} + \frac{1}{15X^2} + 2B + C_0 + \frac{\log T}{24\pi} \\
&\quad + \cosh^2\left(\frac{\pi X}{2}\right) [0.09753X^{-3} + 0.3732X^{-5}].
\end{aligned}$$

Now set $X = \frac{1}{\pi} \log(T/5)$. Then, since $X \geq \frac{1}{\pi} \log(200000)$, we have

$$\begin{aligned}
&\frac{1}{15X^2} + 2B + C_0 + \frac{\pi X + \log 5}{24\pi} \\
&\quad + \cosh^2\left(\frac{\pi X}{2}\right) [0.09753X^{-3} + 0.3732X^{-5}] \\
&\leq 0.0308 \exp(\pi X) X^{-3} = 0.00616T X^{-3}.
\end{aligned}$$

Since $\pi^2\delta/216(\pi^2+4) > 0.00277$, we get

$$\frac{1}{T} \int_0^T S(t) dt \leq \frac{\pi^2}{144 \log^2(T/5)} + \frac{0.1052}{\log^3(T/5)}.$$

Note that

$$\begin{aligned}
&\frac{144}{\pi^2} \log^3(T/5) \log^3(T) \left(\frac{\pi^2}{144 \log^2(T/5)} + \frac{0.1052}{\log^3(T/5)} - E(T) \right) \\
&= \log(T/5) \log^3(T) - \log(T) \log^3(T/5) + \frac{144}{\pi^2} 0.1052 \log^3(T) - 6.59125 \log^3(T/5) \\
&= 2 \log(5) \log(T/5) \log(T) \log(T/\sqrt{5}) + \frac{144}{\pi^2} 0.1052 \log^3(T) - 6.59125 \log^3(T/5).
\end{aligned}$$

This is a cubic polynomial in $\log T$, and it is straightforward to see that it is negative for $T \geq 10^6$. \square

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Appendix A: Rigorous Quadrature

Part of the proof of Theorem 1.1 will involve computational verification of certain statements involving integrals. We will control rounding errors by employing interval arithmetic² to obtain rigorous results and we will rely on the following theorem to perform rigorous quadrature:

Theorem A.1 (Molin). *Let f be a holomorphic function on $D(0, 2) = \{z \in \mathbb{C} : |z| \leq 2\}$. Then for all $n \geq 1$ we have*

$$\left| \int_{-1}^1 f(x) dx - \sum_{k=-n}^n a_k f(x_k) \right| \leq \exp\left(4 - \frac{5n}{\log(5n)}\right) \sup_{D(0,2)} |f|,$$

where $h = \frac{\log(5n)}{n}$, $a_k = \frac{h \cosh(kh)}{\cosh(\sinh(kh))^2}$ and $x_k = \tanh(\sinh(kh))$.

Proof. See [23, Theorem 1.1], which is derived from [22, Theorem 2.10]. \square

Under the conditions of the above theorem, the supremum of $|f(z)|$ with $z \in D(0, 2)$ lies somewhere on the boundary $|z| = 2$. If we are unable (or unwilling) to bound this supremum analytically, we can “cheat” and divide $[0, 1]$ into sufficiently small intervals θ_n , compute using interval arithmetic each

$$|f(2 \exp(2\pi i \theta_n))|$$

and take the maximum.

To apply this theorem in practice, consider estimating

$$\int_{t_0}^{t_1} f(t) dt$$

where $t_1 > t_0$. We can rescale this to

$$\frac{t_1 - t_0}{2} \int_{-1}^1 f\left(\frac{2w - t_1 - t_0}{t_1 - t_0}\right) dw$$

and appeal to Theorem A.1 directly providing the rescaled function has no poles inside the circle $|z| = 2$. The most common awkward situation is where f does have a single pole on the real line at $t = \rho$ and $t_1 > t_0 > \rho$. In this case, we are still in the clear so long as $t_1 < 3t_0 - 2\rho$ which may force us to perform the integral piecewise to cover the required interval.

² We use Johansson’s ARB package [18].

References

1. ACRC, BlueCrystal Phase 3 User Guide (2015)
2. Backlund, R.J.: Über die Nullstellen der Riemannschen Zetafunktion. *Acta Math.* **41**(1), 345–375 (1916)
3. Bateman, H.: *Tables of Integral Transforms*, vol. 1. McGraw-Hill, New York (1954)
4. Batut, C., Belabas, K., Bernardi, D., Cohen, H., Olivier, M.: *User's Guide to PARI-GP* (2000)
5. Booker, A.R.: Turing and the Riemann hypothesis. *Not. Am. Math. Soc.* **53**(10), 1208–1211 (2006)
6. Booker, A.R.: *Turing and the Primes, the Once and Future Turing*. Cambridge University Press, Cambridge, pp. 34–52 (2016)
7. Booker, A.R., Hiary, G.A., Keating, J.P.: Detecting squarefree numbers. *Duke Math. J.* **164**(2), 235–275 (2015)
8. Booker, A.R., Platt, D.J.: The First 2184 Zeros of the Selberg Zeta-Function for $\mathrm{PSL}(2, \mathbb{Z}) \backslash \mathbb{H}$. <https://people.maths.bris.ac.uk/~madjp/Datasets/Selberg.txt>
9. Booker, A.R., Strömbergsson, A.: Numerical computations with the trace formula and the Selberg eigenvalue conjecture. *Journal für die reine und angewandte Mathematik (Crelles J.)* **2007**(607), 113–161 (2007)
10. Booker, A.R., Strömbergsson, A., Venkatesh, A.: Effective computation of Maass cusp forms. *Int. Math. Res. Not. Article ID 71281*, 34 (2006)
11. Gram, J.P.: Note sur les zéros de la fonction $\xi(s)$ de Riemann. *Acta Math.* **27**(1), 289–304 (1903)
12. Hare, D.E.G.: Computing the principal branch of log-Gamma. *J. Algorithms* **25**(2), 221–236 (1997)
13. Hejhal, D.A.: On eigenfunctions of the Laplacian for Hecke triangle groups. In: *Emerging Applications of Number Theory* Minneapolis (1996), IMA Volumes in Mathematics and Its Applications, vol. 109, pp. 291–315. Springer, New York (1999)
14. Hejhal, D.A.: *The Selberg Trace Formula for $\mathrm{PSL}(2, \mathbb{R})$* , vol. 2. Springer, Berlin (2006)
15. Hejhal, D.A.: *The Selberg Trace Formula for $\mathrm{PSL}(2, \mathbb{R})$* , vol. 1. Springer, Berlin (2006)
16. Hutchinson, J.I.: On the roots of the Riemann zeta function. *Trans. Am. Math. Soc.* **27**(1), 49–60 (1925)
17. Iwaniec, H.: *Spectral Methods of Automorphic Forms*, vol. 53. American Mathematical Society Providence (2002)
18. Johansson, F.: Arb: a C library for ball arithmetic. *ACM Commun. Comput. Algebra* **47**(4), 166–169 (2013)
19. Jorgenson, J., Smajlović, L., Then, H.: On the distribution of eigenvalues of Maass forms on certain moonshine groups. *Math. Comp.* **83**(290), 3039–3070 (2014)
20. Li, X., Sarnak, P.: Number Variance for $\mathrm{SL}(2, \mathbb{Z}) \backslash \mathbb{H}$ (2004). <http://publications.ias.edu/sites/default/files/SarnakNumberPaper04.pdf>
21. Littmann, F.: One-sided approximation by entire functions. *J. Approx. Theory* **141**(1), 1–7 (2006)
22. Molin, P.: *Intégration numérique et calculs de fonctions L*, Ph.D. thesis, L'Université Bordeaux I (2010)
23. Molin, P.: *L'intégration numérique par la méthode double-exponentielle*, HAL (2010); [arXiv:hal-00491561v2](https://arxiv.org/abs/1004.1561v2)
24. Niven, I., Zuckerman, H.S., Montgomery, H.L.: *An Introduction to the Theory of Numbers*. Wiley, New York (2008)
25. Steil, G.: Eigenvalues of the Laplacian and of the Hecke Operators for $\mathrm{PSL}(2, \mathbb{Z})$, Tech. Report DESY-94-028, P00022028 (1994)
26. Then, H.: Large Sets of Consecutive Maass Forms and Fluctuations in the Weyl Remainder (2012)
27. Titchmarsh, E.C.: The zeros of the Riemann zeta-function. *Proc. R. Soc. Lond.* **151**, 234–255 (1935)
28. Turing, A.M.: Some calculations of the Riemann zeta-function. *Proc. Lond. Math. Soc.* (3) **3**, 99–117 (1953)
29. Vaaler, J.D.: Some extremal functions in Fourier analysis. *Bull. Am. Math. Soc. (N.S.)* **12**(2), 183–216 (1985)
30. Voros, A.: Spectral functions, special functions and the Selberg zeta function. *Comm. Math. Phys.* **110**(3), 439–465 (1987)

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