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Theoretical study of 2 × 2 planar array of equilateral triangular patch microstrip antenna on ferrite substrate

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att This paper presents theoretical investigations of the radiation properties of a 2×2 element planar array of equilateral triangular patch and thin a built up on a magnetized lerite substrate $N_{1,ma}$ ($^{5}\alpha_{ma}F_{1,ma}$, $^{0}\alpha_{ma}$), in the presence of plasma medium. The far zone EM-mode and Eadvision Lds are derived using vector wave function techniques and pattern multiplication approaches. The results are obtained both in plasma a did in fice space. Computation of maximum field intensity half power beam width (HPRW), full null beam width (FNRW) and other action a parameters such as radiation conductance, directivity and quality factor has been made for the plasma and the free space medium atts of this theoretical analysis suggest that due to ferrite substrate this planar array antenna can be operated at lower frequency range and there e_{max} or derive of such antenna when designed on N(G oFeO, ferrite substrate

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ecol past, microstrip antennas built up on ferrite substrate

TOS Territe materials have a significant amount of TOPV at microwave frequencies. This anisotropy gets kedby applying external dic magnetic field in ferrite matehand brings about non-reciprocal behavior in them.

Martice ently, with the availability of low loss commercial towave, ferrite substrates have been used in microstrip tana applications to provide loading for antenna size when Microstrip antennas on pure dielectric substrate have a cutenvicely analysed in these years due to their advantioner bulky antennas. The band width of such antennas in the order of 1 percent and base not tried at lower ultra high frequency (UHF) due to the vize considerations.

 \mathfrak{M}_{eh} antennas when mounted on aerospace vehicles en- \mathfrak{M}_{ef} plasma medium during their travel in space. It has been

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pointed out that the radiation properties of microstrip antennas in plasma medium are modified to a great extent due to the generation of electroacoustic waves in addition to electromagnetic waves [1-5]

This paper describes the radiation characteristics of a 2×2 planar array of equilateral triangular patch microstrip antenna built up on a ferrite substrate Ni_{1.062}Co_{0.02}Fe_{1.948}O₄. Design requirement and substrate characteristics considered for this analysis are listed in Table 1.

2. Radiation field expressions

The geometry and coordinate system of planar array antenna is shown in Figure 1. It consists of four identical triangular microstrip patch elements of arm length *a*, on a territe substrate $N_{1,062}Co_{0.02}Fe_{1.948}O_4$ of thickness *h* and substrate permittivity \mathcal{F}_r . The array elements which are positioned along *x*-axis are separated by a distance d_y and those along *y*-axis are separated by a distance d_y . Each patch can be excited by a microstrip transmission line connected to the edge or by a coaxial line from the back at the plane $\phi = 0$. We have considered the patch as

Table 1. Design requirement and substrate characteristics of ferrite substrate $N_{1.062}Co_{0.02}Fc_{1.948}O_4$

Design frequency (f)	1.0 GHz
Relative permittivily (e,)	14 78
Dieletric loss tangent (tan δ_i)	0 0005
Magnetic loss tangent $(\tan \delta_m)$	0.005
Applied d.c. magnetic bias field (H_{μ})	796 × 104 Amp/m
Saturation magnetization $(\mu_0 M_s)$	0.03 Testa
Patch dimension (a)	0.022 meter
Substrate thickness (h)	0.0016 meter
Gyromagnetic ratio (γ)	1 76 × 10 ¹¹ rad /sec Tesla
Separation between the array element along x_i and y_i axis respectively $(d_y = d_y)$	λ/2
Progressive phase excitation difference along x and x-axis respectively $(\beta_x - \beta_y)$	π/2

a cavity which acts as a disc resonator. Among the various modes that may be excited in such disc resonator, we have considered TM_{nm} mode with respect to z-axis. Here, m and n are the mode numbers associated with λ_2 and γ_2 -directions respectively.



Figure 1. Configuration and coordinate system of 2×2 element planal array equilateral friangular patch array microstrip antenna on ferrite substrate

The E component of field inside the cavity for dominant mode is given as

$$E_{z} = A_{1,0-1} \Big[2 \cos \left(2\pi x / \sqrt{3a} + 2\pi / 3 \right) \cos \left(2\pi y / 3a \right) \\ + \cos \left(4\pi y / 3a \right)$$
(1)

A ferrite slab may support various guided modes and these propagating modes get affected in the presence of applied magnetic bias field. A strong bias field may be obtained easily along the plane of antenna by placing two pole pieces of magnet In close contact with substrate. When propagalisingelectromagnetic waves take place along this $bias_{1bl}$ infinite ferrite having two orthogonal feeds, two $plar_{a}$ modes namely left hand circularly polarized (LHCP) mode right hand circularly polarized (RHCP) mode are generate

The magnetic properties of ferrite substrate attect hat, modes, therefore the effective permeability for tights circularly polarization and left hand circularly polarization be given as [6-8]

$$\begin{aligned} \boldsymbol{\mu}_{effR} &= \boldsymbol{\mu}_{o} \; \frac{1 + \boldsymbol{\omega}_{m}}{\boldsymbol{\omega}_{o} - \boldsymbol{\omega}} \\ \boldsymbol{\mu}_{effI} &= \boldsymbol{\mu}_{o} \; \frac{1 + \boldsymbol{\omega}_{m}}{\boldsymbol{\omega}_{o} + \boldsymbol{\omega}} , \end{aligned}$$

where ω_0 and ω_m are the precession and forced preptrequencies, respectively and may be defined as

$$\omega_0 = \mu \gamma H_0, \omega_m = \mu \gamma H$$
, and $\omega = 2\pi f$,

The expression of the resonant frequency of constequilateral triangular microstrip antenna on ferrite subset TM_{10} mode is given as

$$f_{\mu} = \frac{2 \epsilon \sqrt{\mu_{\mu}}}{3 a \sqrt{\epsilon_{\mu} \mu_{\mu}}}$$

Other field components are obtained by solving Maw equations. By image theory, the ground plane may be up by an image of the top conductor. The magnetic currents exist along the edges of triangular conductor and m. evaluated from

$$M = 2(E \times n),$$

where n is a unit vector normal to the aperture

The total far zone fields of planar array antenna of expressed by the fields of a single element positioned origin multiplied by a factor which is referred to as the factor. This method is widely known as pattern multipliapproach. Thus, using the expression of single de equilateral triangular patch microstrip antenna and negl the coupling between the elements, the total field of the array antenna can be expressed as [9-12]

E (total) = E (single element placed at the origin) × array factor (AF)

As the entire array is taken as uniform, the normalized of the array factor (AF) is obtained and may be written as

$$AF_n(\theta,\phi) = \frac{1}{M} \times \frac{\sin(M\psi_{\chi}/2)}{\sin(\psi_{\chi}/2)} \times \frac{1}{N} \times \frac{\sin(N\psi_{\chi}/2)}{\sin(\psi_{\chi}/2)}$$

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 $F_{\rm off}$ the case of 2 × 2 element planar array of equilateral mular patch microstrip antenna on territe substrate, the EM to mode fields are given as

FM mode

$$F_{vi} = i\eta_0 \omega \left[-F_v \sin \phi + F_v \cos \phi \right] \times \frac{1}{4}$$

$$\times \frac{\sin(\psi_v)}{\sin(\psi_v/2)} \times \frac{\sin(\psi_v)}{\sin(\psi_v/2)}$$
(8)

$$I_{x,y} = i\eta_0 \omega \left[F_y \cos\theta \cos\phi + F_y \cos\theta \sin\phi \right] \times \frac{1}{4}$$

$$\frac{\sin(\psi_{\chi})}{\sin(\psi_{\chi}/2)} \times \frac{\sin(\psi_{\chi})}{\sin(\psi_{\chi}/2)}$$
(9)

$$I_{1} = 2h\beta_{p}\omega_{p,r}^{2}/3a\omega\epsilon_{0}\left(\omega^{2} - \omega_{p}^{2}\right) \times \exp\left(i\beta_{p}r\right)/r$$

$$\leq \exp\left(-i\beta_{p}a\sin\theta\cos\phi/\sqrt{3}\right)$$

$$= \frac{\sin\left\{\beta_{p}/d_{x}\sin\theta\cos\phi+\beta_{x}\right\}}{\sin\left\{05(\beta_{p}/d_{x}\sin\theta\cos\phi+\beta_{x})\right\}}$$

$$= \frac{\sin\left\{\beta_{p}/d_{x}\sin\theta\cos\phi+\beta_{x}\right\}}{\sin\left\{05(\beta_{p}/d_{x}\sin\theta\cos\phi+\beta_{x})\right\}} \times \left[E_{m} + E_{px}\right], (10)$$

 $\Psi_{\chi} = \beta_{\chi} d_{\chi} \sin \theta \cos \phi + \beta_{\chi}$ and $\Psi_{\chi} = \beta_{e} d_{\chi} \sin \theta$ $\sin + \beta_{\chi}$, M = elements placed along the x-axis andis respectively, β_{χ} , $\beta_{\chi} = \text{progressive phase excitation}$ $\cos \text{along x- and y-directions respectively, } E_{\theta t}, E_{\phi t} =$ proments of total electric field vectors for EM-mode, $Ept = \frac{1}{2}$ electric field vector for P-mode, $F_{\chi}, F_{\chi} = \text{vector electric}$ maks for y-component and y-component, respectively. $E_{\mu\nu}, E_{\mu\nu}$ = P-mode electric field vectors for x-component and ycomponent respectively. β_{ν} = phase propagation constant for EM- mode given by $2\pi A/\lambda_0$, β_{μ} = phase propagation constant for P- mode given by $\beta_{\nu}c/\nu$, c = velocity of light, ν = root mean square thermal velocity of electron, A = plasma frequency parameter given by $(1-\omega_{\mu}^2/\omega^2)^{1/2}$, ω_{μ} = angular plasma frequency, ω = angular source frequency, η_0 = free space impedance equal to 120 π ohms

3. Field patterns

The expression for total field pattern $R(\theta, \phi)$ is obtained as

$$R(\theta, \phi) = E_{\theta_l} \Big|_{\theta_l}^{\infty} + \Big| E_{\phi_l} \tag{11}$$

Then, the radiation field patterns in the E-plane ($\phi = 0$) and H-plane ($\phi = \pi/2$) are given as

$$R_{e}(\theta,\phi) = \left|E_{\theta t}\right|^{2} + \left|E_{\phi t}\right|^{2} = \eta_{0}^{2}\omega^{2}\left(\left|F_{s}\right|^{2} + \left|F_{s}\right|^{2}\cos^{2}\theta\right),$$
(E-plane) (12)

$$\mathbf{R}_{h}(\boldsymbol{\theta},\boldsymbol{\phi}) = \left| \boldsymbol{E}_{\boldsymbol{\theta} t} \right|^{2} + \left| \boldsymbol{E}_{\boldsymbol{\phi} t} \right|^{2} = \boldsymbol{\eta}_{0}^{-} \boldsymbol{\omega}^{2} \left(\left| \boldsymbol{F}_{x} \right| + \left| \boldsymbol{F}_{x} \right|^{-} \cos^{2} \boldsymbol{\theta} \right)$$
(13)

The values of R_{ν} and R_{h} are calculated for a case taking $f = 1 \text{ GHz}, a = 2.2 \text{ cm}, \epsilon_{\nu} = 14.78$, Applied d c magnetic bias field $(H_{\mu}) = 7.96 \times 10^4 \text{ Amp/m}$, Saturation magnetization $(\mu_0 M_{\chi}) = 0.03 \text{ Tesla}$, $d_{\chi} = d_{\chi} = \lambda/2$ and the phase difference $\beta_{\chi} = \beta_{\chi} = \pi/2$. The results are plotted in Figures 2 and 3 for two different planes ($\phi = 0$ and $\phi = \pi/2$) for $A = 1.0 \nu c$ in free space and in Figures 4 and 5 for two different planes ($\phi = 0$ and $\phi = \pi/2$) A = 0.5 νc in plasma medium for TM₁₀ mode of excitation Pattern characteristics of planar ariay are computed in Table 2. Comparison between radiation properties of the planar array antenna on dielectric and ferrite substrate are shown in Table 3.

Table 2. Pattern characteristics of the 2 × 2 clement planar array of triangular microstrip patch antenna on ferrite substrate

St No	Pattern	$\phi = 0$ plane $\Lambda = 1.0$		$\phi = 0$ plane A = 0.5		$\phi - \pi/2 \text{plane}$ $A = 1 \ 0$		$\phi = \pi/2 \text{plane}$ $A = 0.5$	
	characteristics _								
		RHCP	LHCP	RHCP	LHCP	RHCP	LHCP	RHCP	LHC P
I	Half-power beam width (HPBW)	60°	80"	140°	150°	60"	80"	140°	150°
2	Full null beam width (FNBW)	120°	140°	180°	280°	120°	140"	180°	280"
1	Direction of max	90°	90°	180°	180°	90°	90°	180"	180°



Figure 2. B plane radiation patterns of 2×2 element planar array and single-element equilateral triangular patch microstrip antennas on ferrite substrate. For A = 1.0 (the space)



Figure 3. H-plane radiation patterns of 2×2 element planar array and single-element equilateral triangular patch microstrip antennas on ferrite substrate for A = 1.0 (the space)



Figure 4. E-plane radiation patterns of 2×2 element planar array and single-element equilateral triangular patch microstrip antennas on ferrite substrate for A = 0.5 (plasma)

4. Other antenna parameters

4.1 Radiation conductance

The total power radiated can be $cal_{U_{0}}$ by employing Poynting's theorem and κ_{p} as .

$$P_{i} = (\frac{1}{2}) (\frac{1}{2}) \left\{ \operatorname{Re} \iint s \left(E \times H^{*} \right)_{d} \right\}$$

The factor $\frac{1}{2}$ is due to the fact that the p is radiated through the upper half space and S is the total spherical surface area

$$P_{i} = \left(\frac{1}{4}\right) \left\{ \operatorname{Re} \iint \operatorname{s} \left(E_{\theta} H_{\phi}^{*} - E_{\phi} H_{u} \right) \right\}$$

Thus, the expression for radiated $p_{\rm the}$ EM-modes is obtained using the relation

$$P_{\epsilon} = (A/4\eta_0) \int_0^{2\pi} \int_0^{\pi} \left\{ \left| E_{\theta_1} \right| \right\}$$

$\times r^2 \sin\theta \, d\theta d\phi$

Here, $A = (1 - \omega_p^2 / \omega^2)^{1/2}$ and ω^2 the plasma to source frequency rate

The radiation conductance of an anten EM-mode can be defined as

$$G_e \doteq 2 P_e / V_0^2$$

4.2 Directivity

The directivity of an antenna is defined, ratio of the maxium radiation intensity per unit solid angle) to the average radiintensity It can be expressed as

$$D_r = 4 \pi \left\{ \max\left(\left| E_{\theta t} \right|^2 + \left| E_{\phi t} \right|^2 \right) \right\}$$
$$= \left\{ \int_0^{2\pi} \int_0^{\pi} \left| E_{\theta t} \right|^2 + \left| E_{\phi t} \right|^2 r^2 \sin\theta \ d\theta \ d\phi \right\}$$

4 3 Quality factor

A parameter specifying frequency selected of a resonant circuit is the quality factor Q, can be defined as the ratio between er stored in the system and the energy loa

The total Q of a microstrip radiating of comprises contributions due to the rad Q_r , conductor loss Q_i , and dielectric log quality factors, so

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	Patch dimensions (meters)	Resonant Irequency (GHz)	Radiation conduction (mho)	Quality factor	Directive gain (dB)
Pure-dielectric substrate	0 022	5 97	0 001	7 756	6 518
$\varepsilon_r \simeq 2.32.$					
h = 0.0016					
Fernite substrate, LHCP	0 0 2 2	12	1 116×10 4	252 15	75
$L_1 = 14.78$					
h = 0.0016					
Ferrite substrate RHCP	0 0 2 2	09	7 356×10 °	259 09	6 6 2 5
$\iota_{,} = 14.78$					
h ≕0.0016					

Table 3. Comparison between radiation properties of the 2×2 element planar array of triangular microstrip patch antenna on pure dielectric substrate and territe substrate in TM_{10} mode

$$\mu_{Q} = 1/Q_{t} + 1/Q_{t} + 1/Q_{d}, \qquad (17)$$

$$\mu_{L} = \omega_{L} W_{t}/P_{t},$$

 $(v - \omega W_i P_i = (\pi f \mu \sigma)^{W} h$

$$\psi = \omega W_t / P_d = 1/\tan\delta$$

lete W is the energy stored in the antenna element, P_i and weilloss factors due to the conductors and dielectric,

carely σ is the conductivity of the conductors

nergy stored in the triangular radiating element is given

$$= (\varepsilon h/2) \iint \left| E_{\varepsilon}(xv) \right|^2 dx dv \tag{18}$$

onclusion

radiation characteristics of 2×2 planar array of equilateral galar patch microstrip antenna on ferrite substrate in plasma





as well as in free space have been studied by considering presence of bias magnetic field in the direction of propagation of e m waves The results of the array geometry are compared with those of single-element of equilateral triangular patch microstrip antenna Figures 2 and 3 illustrate the radiation patterns of 2 × 2 element planar array of equilateral triangular microstrip antenna for two planes ($\phi = 0$ and $\phi = \pi/2$) in free space. It is observed that the beam splits up into major and minor lobes and corresponding position of maximum radiation is shifted. The radiation patterns of a single element antenna for both RHCP and LHCP waves contain only one major lobe of considerably wide beam width, while the array geometry for both RHCP and LHCP waves produces a directive beam with a narrow beam width. The radiation field patterns of this array antenna indicate that beam width with LHCP waves is larger than that with RHCP waves. Figures 4 and 5 illustrate the radiation patterns of 2×2 element planar array of equilateral triangular microstrip antenna for two planes ($\phi=0$ and $\phi=\pi/2$) in plasma medium (A=0.5) From these figures, it is observed that the shape of the field patterns has been

changed and it redistributes the field intensity The computed values of pattern characteristics of the planar array are given in Table 2 It is clear from this table that half power beam width of the patterns in free space is relatively small in comparison to plasma medium Thus it reveals that the patterns are more directive in free-space than in plasma medium From Tables 4 and 5, it is clear that the effect of plasma on radiation conduction of single element as well as planar array of equilateral triangular microstrip antenna in RHCP mode is significantly weaker than the LHCP mode The directivity of 2×2 element planar array of equilateral triangular microstrip antenna with LHCP waves is little higher than

Table 4. Antenna parameters for single element and 2×2 planar array equilateral triangular patch microstrip antenna = 0 ferrite substrate with RHCP waves

Si No	Plasma parameter (A)	G, tor single clement (mho)	G _c for 2×2 planar array (mho)	D _i for single element (dB)	D, for 2 ×2 planar artay	Q, foi single element	Q, for 2 planar array
1	10	1.588×10 ⁻⁵	7 356×10-5	4 751	6 625	444 418	259 80
2	0.9	1.428×10-1	6 819×10-1	4 755	6 4 1 4	453 355	270 26
3	08	1 268×10 `	6 152 ×10 5	4 758	6 244	462 635	284 45
4	07	1 109×10 5	5 379×10	4 761	6 108	472 281	302.93
5	06	9.468×10	4 537×10 °	4 764	5 99	482 319	325 98
6	0 5	7 911×10 -6	3 673×10 ⁻¹	4 766	5 885	492 774	153.60
7	04	4 743×10 -6	2 829×10 ⁻⁵	4 768	5 74	503 678	385 474
8	03	4 743×10 °	2 073×10 5	4 769	5 523	515 061	421 119
9	0 2	3-161×10 ⁻⁶	1 309×10 1	4 77	5 248	526 961	460 223
0	0.1	1 58×10 °	6 179×10-6	4 771	4 931	539 417	503-30-
	0.0	0.0	0.0	4 771	4 772	-	-

Table 5 Antenna parameters: for single element and 2×2 planar array equilateral triangular patch microstrip antenna e ferrite substrate with LHCP waves

Si No	Plasma parameter (A)	G ₂ for single clement (mho)	G ₁ tor 2×2 planar array (mho)	D _i for single clement (dB)	D _i for 2 ×2 planar array	Q, for single clement	Q, tor 2x planar array
1	10	2 811×10 `	1 116×10 4	4 7 3 6	7 5	449 784	252 11
2	09	2 526×10 ⁻⁵	1 071×10 ⁴	4 742	7 151	462 144	258 296
٦	0.8	2 242×10-5	1.013×10.4	4 748	6 767	475 146	266.621
4	07	1 959×10 ¹	9 287×10 1	4 754	6 475	488 847	279 728
5	U G	1678×10 `	8 142×10 ⁻¹	4 758	6 24	503 311	299 766
6	0 1	1 397×10 ⁻⁵	6 751×10 -5	4 762	6-063	518 612	328 351
7	04	1 117×10 ⁻⁵	5 237×10 '	4 765	5 921	534 831	366 356
я	0 1	8 37×10 -6	3 741×10 5	4 768	5 7 3 7	552 062	413 675
9	0 2	5 578×10-4	2 364×10 -5	4 77	5 435	570 413	469 492
10	01	2 788×10-6	1.143×10 ⁻¹	4 771	5 024	590 006	533 852
11	0.0	0.0	0.0	4 771	4 77	-	-

that with RHCP waves for all the plasma frequency and free space. The total quality factor of planar array antenna with LHCP waves is lower than that with RHCP waves.

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