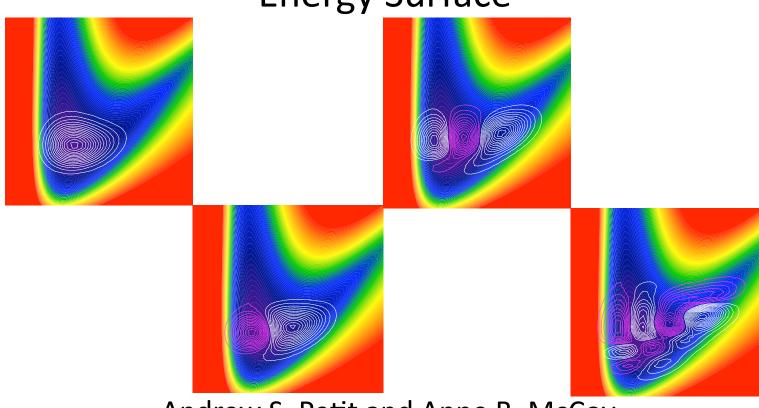
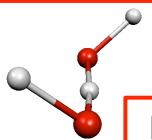
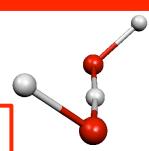
Progress Towards the Accurate Calculation of Anharmonic Vibrational States of Fluxional Molecules and Clusters Without a Potential Energy Surface

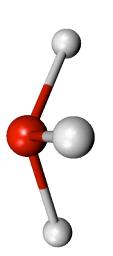


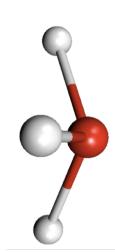
Andrew S. Petit and Anne B. McCoy
The Ohio State University

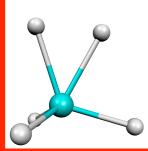


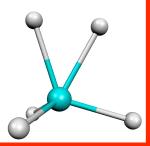


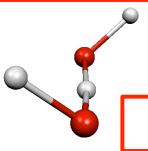
Harmonic analysis pathologically fails

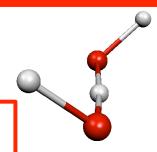




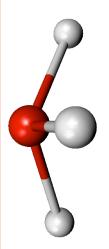


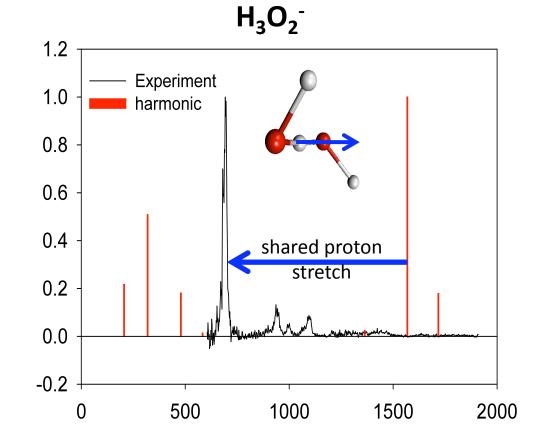




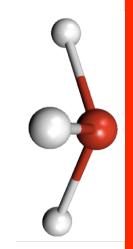


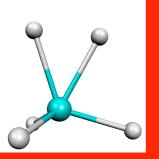
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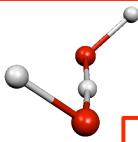




Photon energy

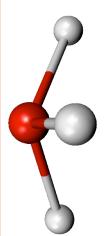




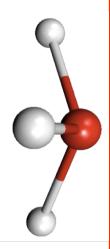


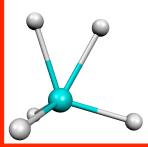


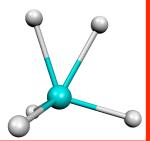
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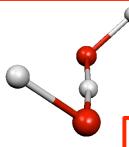


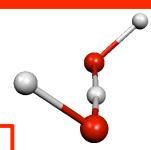
An accurate and general treatment of these highly fluxional systems requires either a global PES or a grid-based sampling of configuration space



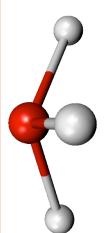




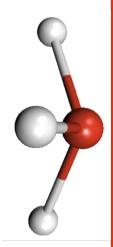




Harmonic analysis pathologically fails

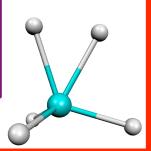


An accurate and general treatment of these highly fluxional systems requires either a global PES or a grid-based sampling of configuration space





Large basis set expansions required to fully capture the anharmoncity of these systems





Monte Carlo Sample Configuration Space

Use importance sampling Monte Carlo

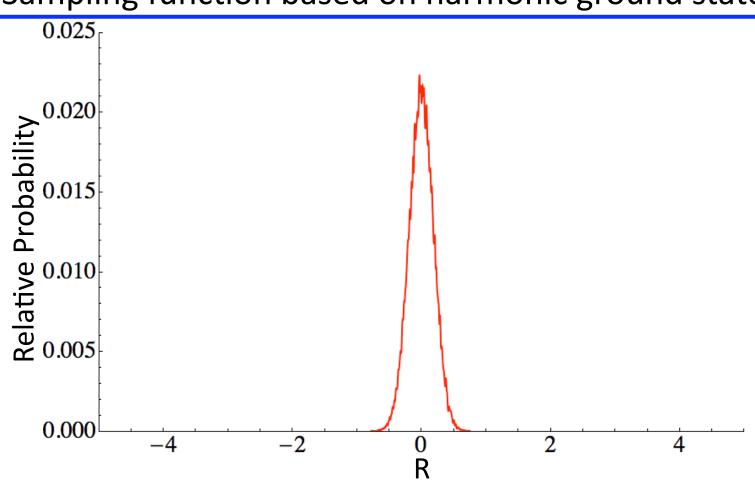
Use importance sampling Monte Carlo

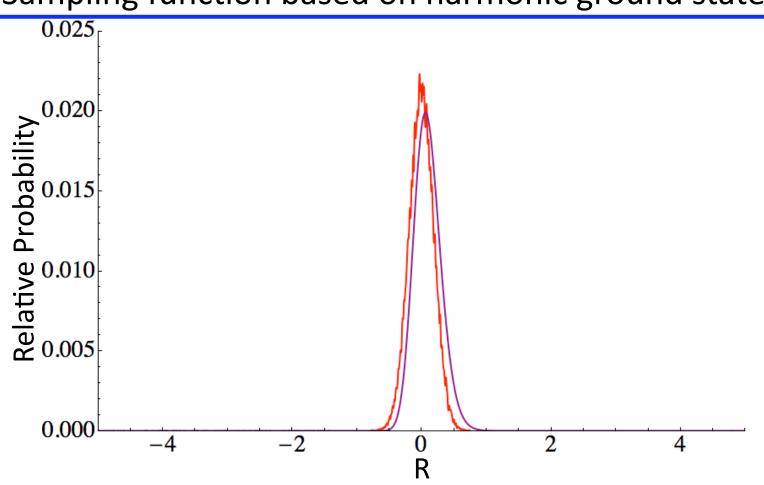
$$\int_{V} f(\vec{x}) dV = \int_{V} \frac{f(\vec{x})}{g(\vec{x})} g(\vec{x}) dV \approx \frac{V}{N} \sum_{i=1}^{N} \frac{f(\vec{x}_i)}{g(\vec{x}_i)}$$

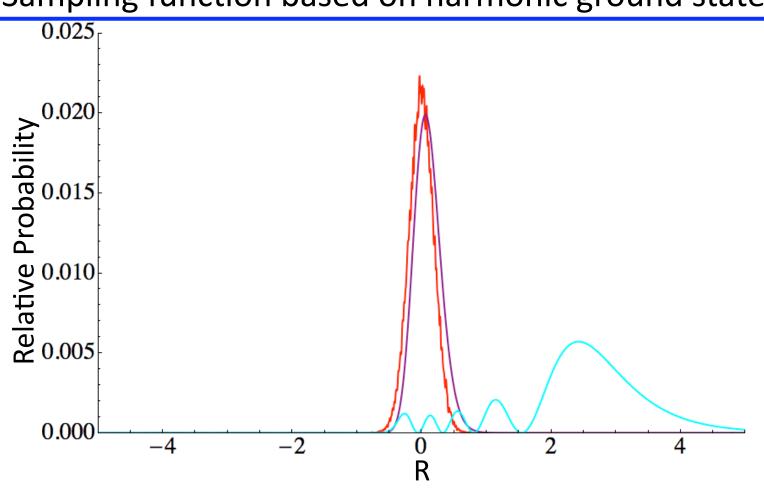
Use importance sampling Monte Carlo

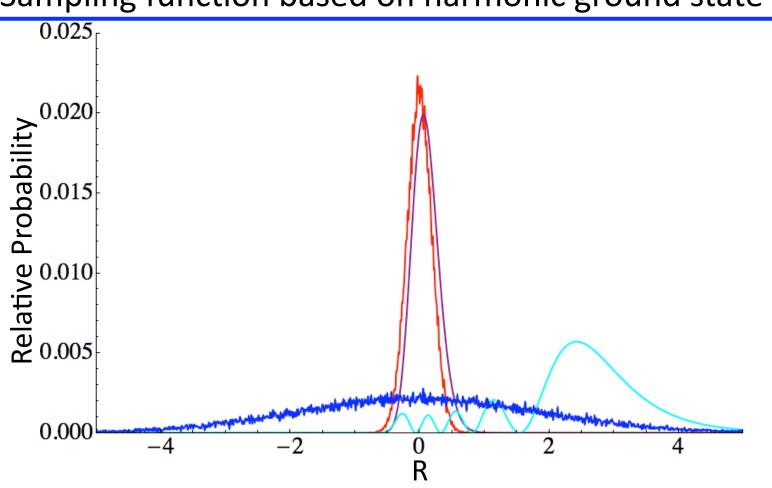
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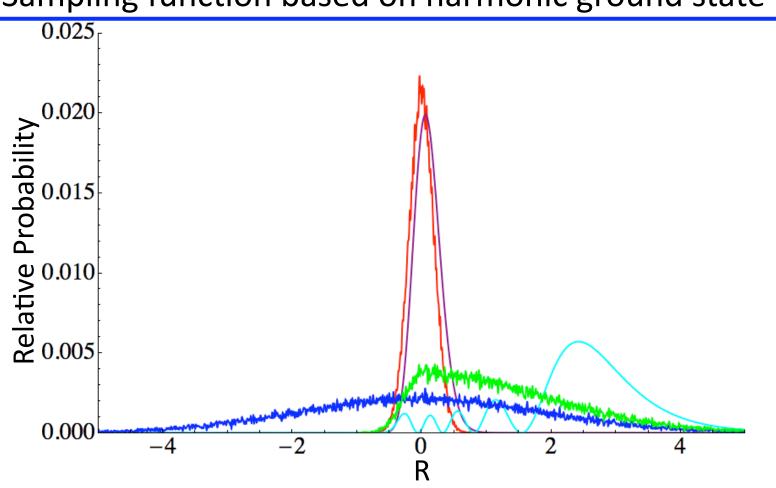
Normalized probability distribution that is peaked, ideally, where the function of interest, f, is peaked.





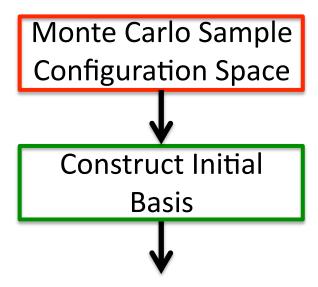




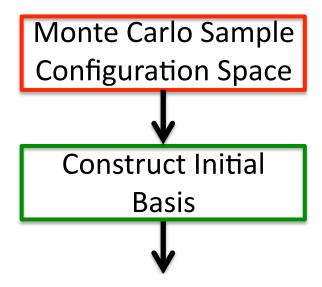


Monte Carlo Sample
Configuration Space

Construct Initial
Basis

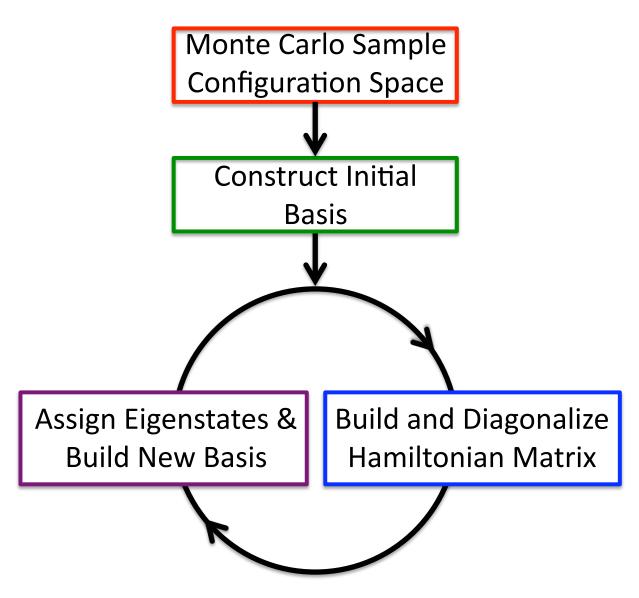


Build and Diagonalize Hamiltonian Matrix



Assign Eigenstates & Build New Basis

Build and Diagonalize Hamiltonian Matrix



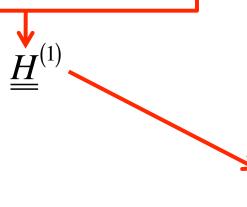


Harmonic oscillator basis

Harmonic oscillator basis

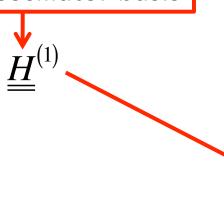


Harmonic oscillator basis



 $\left\{ ilde{\Psi}^{\left(1
ight) }
ight\} ,\left\{ E^{\left(1
ight) }
ight\}$

Harmonic oscillator basis

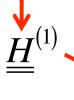


 $\left\{ ilde{\Psi}^{\left(1
ight)}
ight\} ,\left\{ E^{\left(1
ight)}
ight\}$

Eigenstates used to construct new basis

$$\max\left(\left|\left\langle \tilde{\Psi}_{k}^{(1)}\left|\hat{r}\right|\Psi_{m-1}^{(1)}\right\rangle\right|^{2}\right) \Longrightarrow \tilde{\Psi}_{m}^{(1)}$$

Harmonic oscillator basis



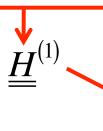
$$\left\{ \Psi^{\left(2\right) }
ight\}$$

$$\left\{ ilde{\Psi}^{\left(1
ight)}
ight\} \!\!,\! \left\{ E^{\left(1
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$$\max\left(\left|\left\langle \tilde{\Psi}_{k}^{(1)} \left| \hat{r} \right| \Psi_{m-1}^{(1)} \right\rangle\right|^{2}\right) \Longrightarrow \tilde{\Psi}_{m}^{(1)}$$

Harmonic oscillator basis



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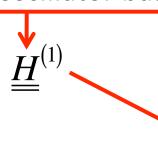
Define an active space:

$$\left\{ \Psi_{m}^{(2)} \right\}_{active} = \left\{ \tilde{\Psi}_{m}^{(1)} \right\}_{active}$$

Eigenstates used to construct new basis

$$\max\left(\left|\left\langle \tilde{\Psi}_{k}^{(1)}\left|\hat{r}\right|\Psi_{m-1}^{(1)}\right\rangle\right|^{2}\right) \Longrightarrow \tilde{\Psi}_{m}^{(1)}$$

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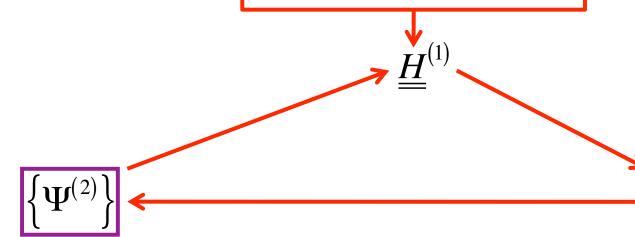
$$\left\{ \Psi_{m}^{(2)} \right\}_{active} = \left\{ \tilde{\Psi}_{m}^{(1)} \right\}_{active}$$

$$\overline{\Psi}_{m \notin active}^{(2)} = \sum_{k=1}^{m} \kappa_{m,m-k}^{(2)} r^k \Psi_{m-k}^{(2)}$$

Eigenstates used to construct new basis

$$\max\left(\left|\left\langle \tilde{\Psi}_{k}^{(1)} \left| \hat{r} \right| \Psi_{m-1}^{(1)} \right\rangle\right|^{2}\right) \Longrightarrow \tilde{\Psi}_{m}^{(1)}$$

Harmonic oscillator basis



 $\left\{ ilde{\Psi}^{\left(1
ight)}
ight\},\left\{E^{\left(1
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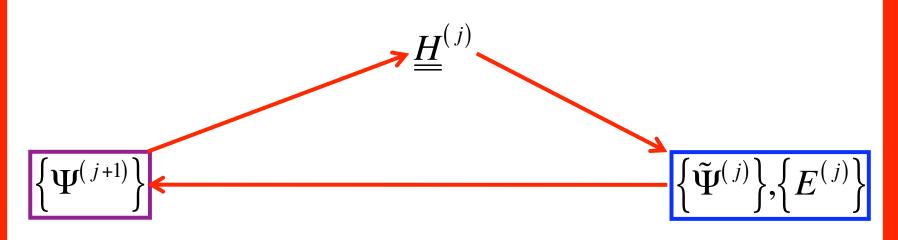
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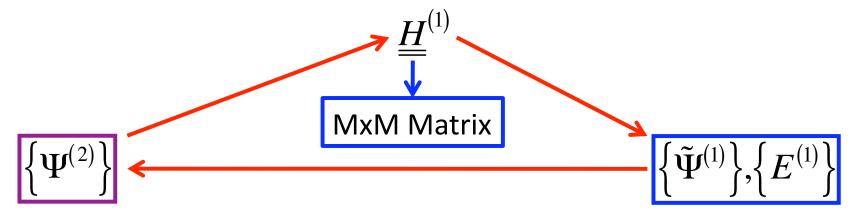
Define an active space:

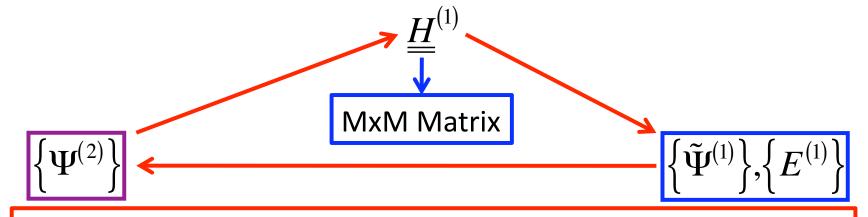
$$\left\{ \Psi_{m}^{(j+1)} \right\}_{active} = \left\{ \tilde{\Psi}_{m}^{(j)} \right\}_{active}$$

$$\overline{\Psi}_{m \notin active}^{(j+1)} = \sum_{k=1}^{m} \kappa_{m,m-k}^{(j+1)} r^k \Psi_{m-k}^{(j+1)}$$

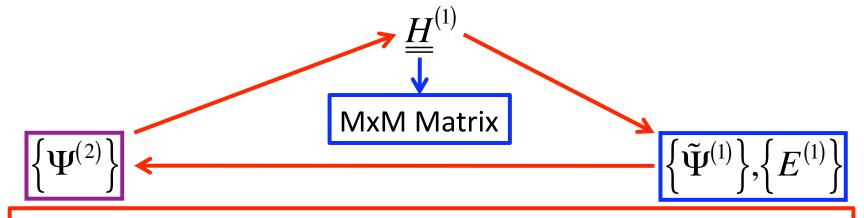
Eigenstates used to construct new basis

$$\max\left(\left|\left\langle \tilde{\Psi}_{k}^{(j)}\left|\hat{r}\right|\Psi_{m-1}^{(j)}\right\rangle\right|^{2}\right) \Longrightarrow \tilde{\Psi}_{m}^{(j)}$$



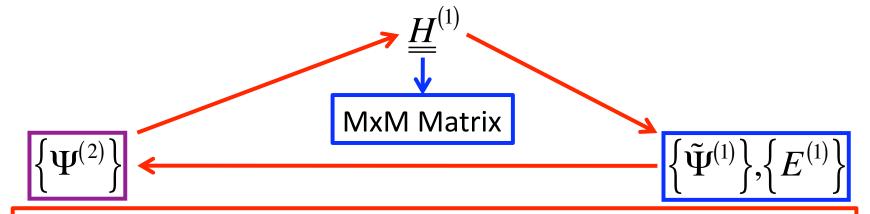


Suppose active space only contains ground state



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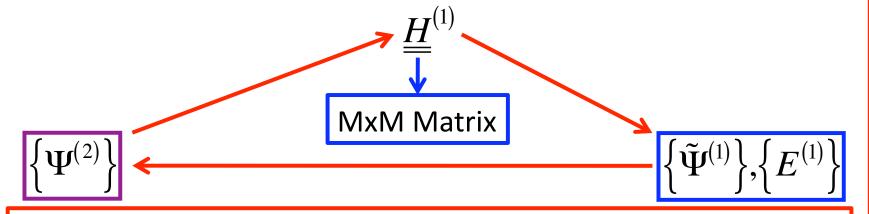
$$\overline{\Psi}_{1}^{(1)} = ra_{0}^{(1)}e^{-\frac{\alpha r^{2}}{2}}$$



Suppose active space only contains ground state

$$\overline{\Psi}_{1}^{(1)} = ra_{0}^{(1)}e^{-\frac{\alpha r^{2}}{2}}$$

$$\overline{\Psi}_0^{(2)} = \left(a_0^{(2)} + a_1^{(2)}r + a_2^{(2)}r^2 + \dots + a_{M-1}^{(2)}r^{M-1}\right)e^{-\frac{\alpha r^2}{2}}$$

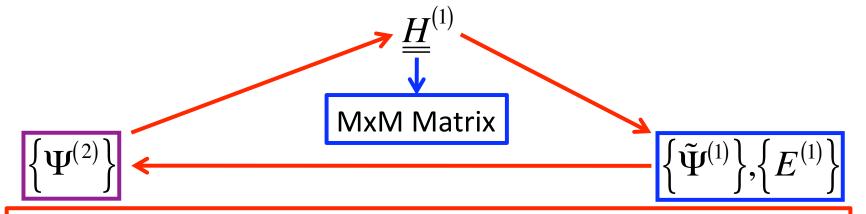


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$$\overline{\Psi}_{1}^{(2)} = r\overline{\Psi}_{0}^{(2)} = \left(a_{0}^{(2)}r + a_{1}^{(2)}r^{2} + a_{2}^{(2)}r^{3} + \dots + a_{M-1}^{(2)}r^{M}\right)e^{-\frac{\alpha r^{2}}{2}}$$

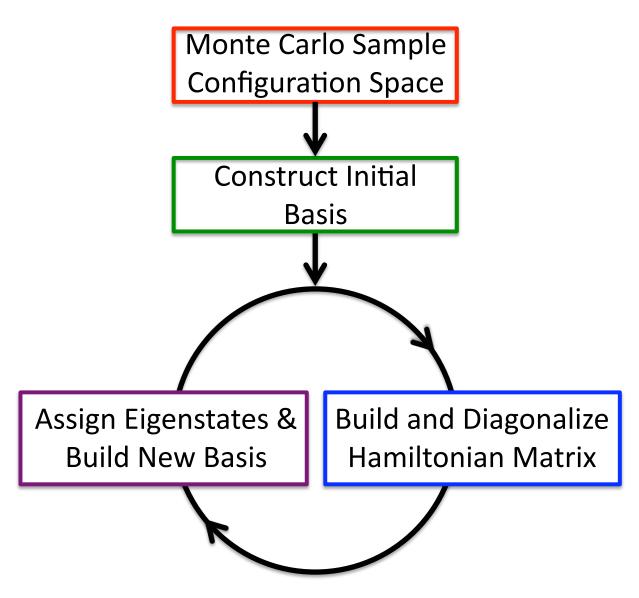


Suppose active space only contains ground state

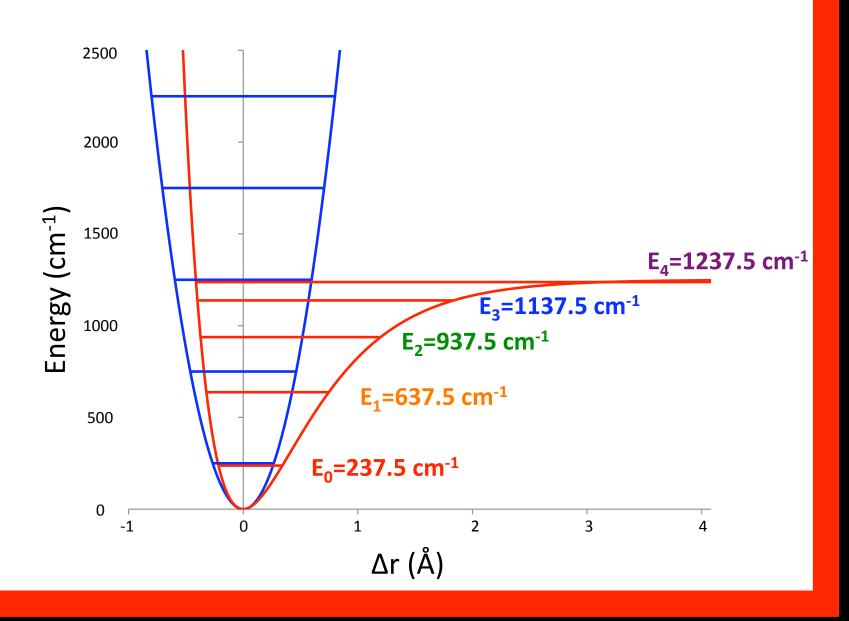
$$\overline{\Psi}_0^{(2)} = \left(a_0^{(2)} + a_1^{(2)}r + a_2^{(2)}r^2 + \dots + \overline{a_{M-1}^{(2)}r^{M-1}}\right)e^{-\frac{\alpha r^2}{2}}$$

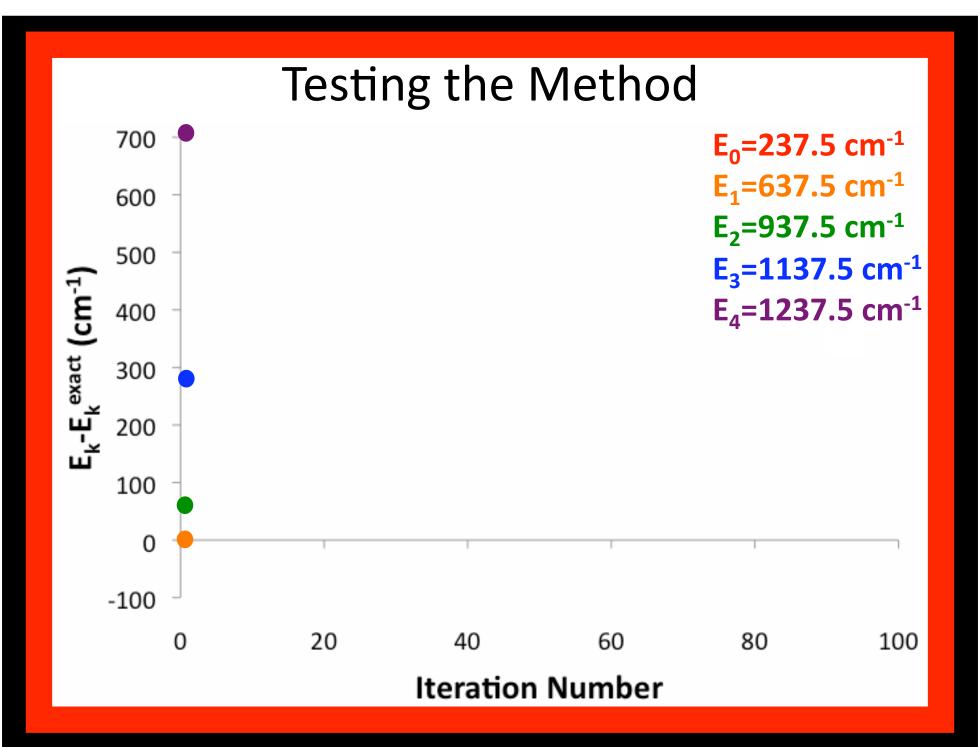
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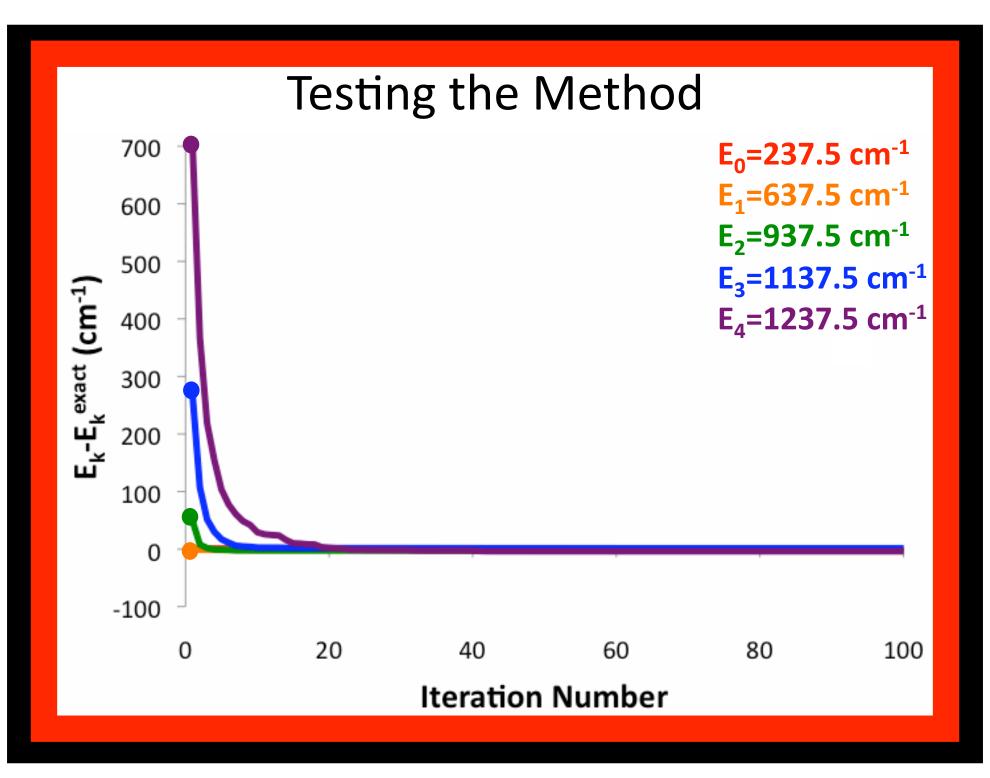
Effective size of basis grows each iteration



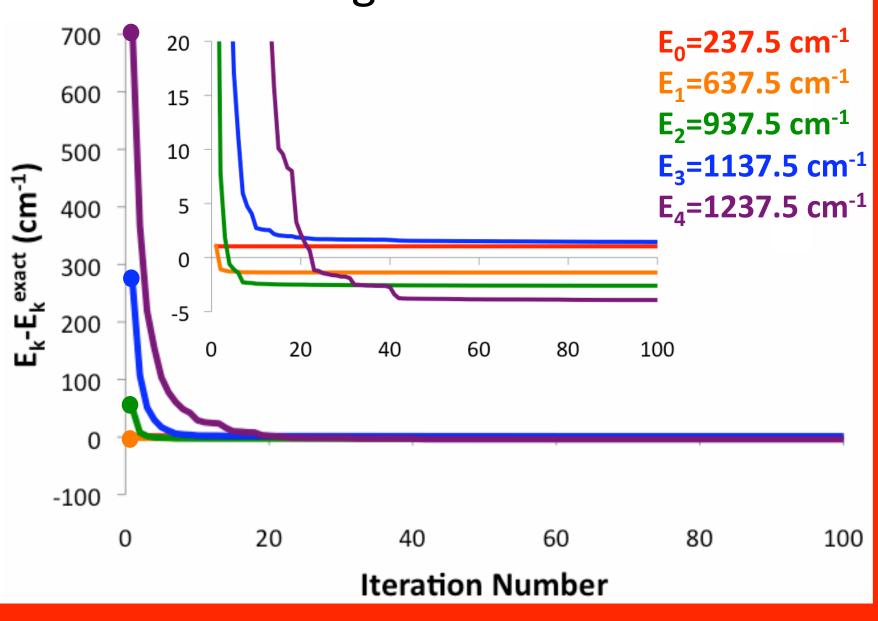
Testing the Method











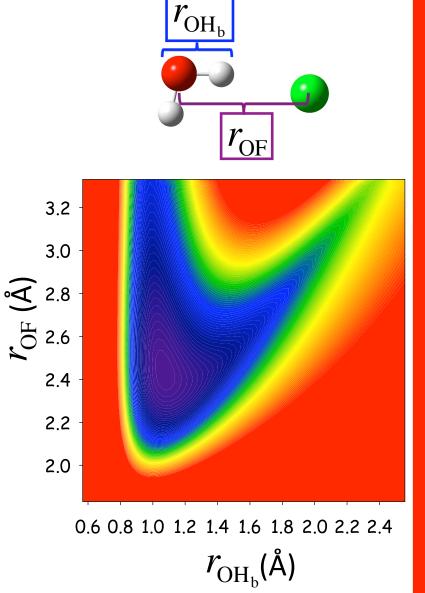
Moving on to 2D

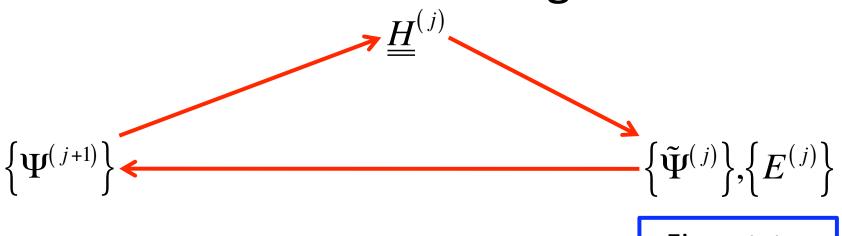
2D PES and G matrix elements constructed by S. Horvath from a bicubic spline interpolation of a grid of points on which MP2/aug-cc-pVTZ calculations were performed.

Highly anharmonic system with a strong coupling observed between $r_{\rm OH_h}$ and $r_{\rm OF}$.

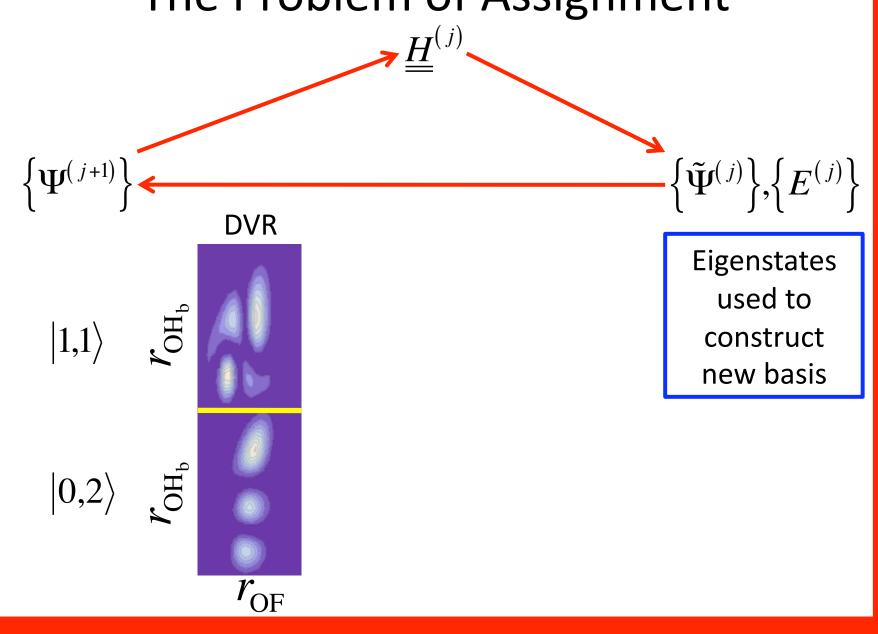
Our goal is to accurately capture all states with up to 3 total quanta of excitation.

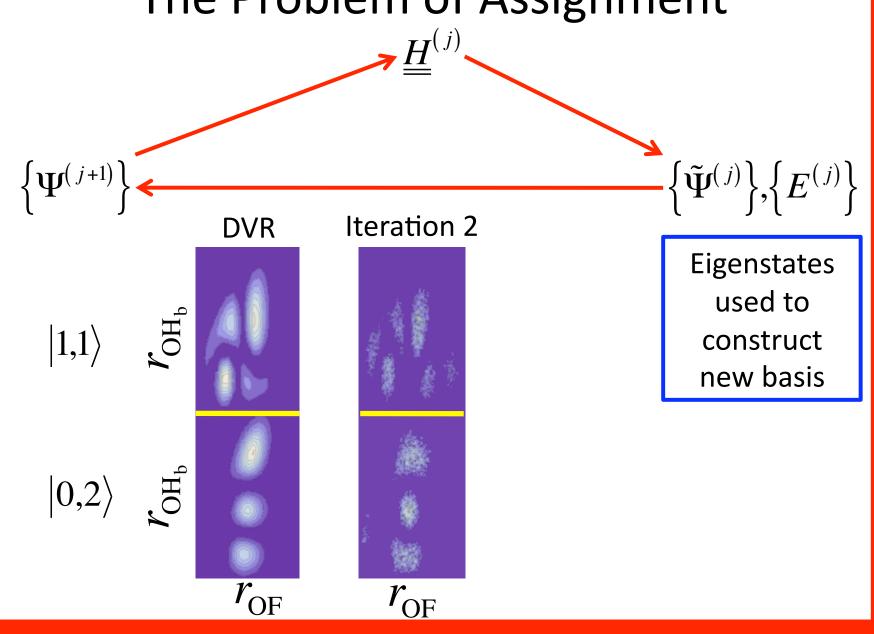
Horvath et al. J. Phys. Chem. A **2008**, 112, 12337-12344.

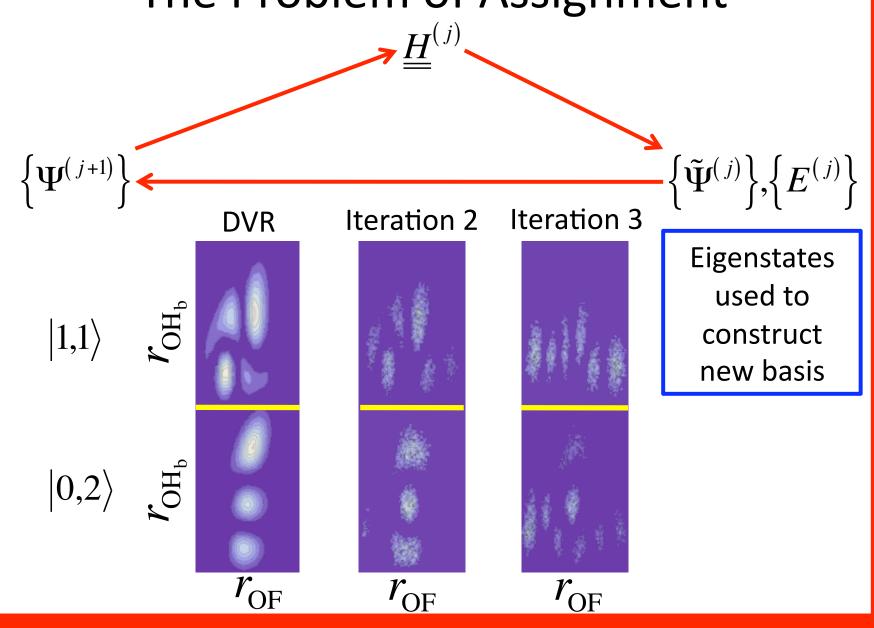




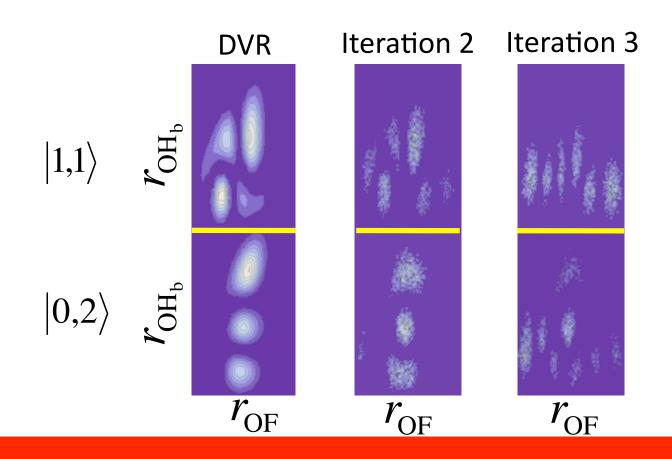
Eigenstates used to construct new basis







Identity of problem states dependent on parameters of algorithm AND exact distribution of Monte Carlo sampling points



$$\Psi_{m,n}^{(j)}$$
 Basis function

$$\phi_k^{(j)}$$
 Un-assigned eigenstate

$$\Psi_{m,n}^{(j)}$$
 Basis function

$$oldsymbol{\phi}_k^{(j)}$$
 Un-assigned eigenstate

$$\left|\left\langle oldsymbol{\phi}_{k}^{\left(j
ight)}\left|\Psi_{m,n}^{\left(j
ight)}
ight
angle
ight|^{2}$$

$$\Psi_{m,n}^{(j)}$$
 Basis function

 $oldsymbol{\phi}_k^{(j)}$ Un-assigned eigenstate

$$\left|\left\langle oldsymbol{\phi}_{k}^{\left(j
ight)}\left|\Psi_{m,n}^{\left(j
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ight|^{2}$$

$$\left|\left\langle \boldsymbol{\phi}_{k}^{(j)} \left| \hat{\boldsymbol{r}}_{OF} \right| \boldsymbol{\Psi}_{m-1,n}^{(j)} \right\rangle \right|^{2} + \left|\left\langle \boldsymbol{\phi}_{k}^{(j)} \left| \hat{\boldsymbol{r}}_{OH_{b}} \right| \boldsymbol{\Psi}_{m,n-1}^{(j)} \right\rangle \right|^{2}$$

 $\Psi_{m,n}^{(j)}$ Basis function

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Compare moments of the two sets of wavefunctions:

$$\begin{cases}
\left\langle \left(r_{OF} - \left\langle r_{OF} \right\rangle_{k}\right)^{2} \right\rangle_{k} \\
\left\langle \left(r_{OF} - \left\langle r_{OF} \right\rangle_{k}\right) \left(r_{OH_{b}} - \left\langle r_{OH_{b}} \right\rangle_{k}\right) \right\rangle \\
\left\langle \left(r_{OH_{b}} - \left\langle r_{OH_{b}} \right\rangle_{k}\right)^{2} \right\rangle_{k}
\end{cases}$$

$$F_{k,l;m,n}^{(j)} = \frac{1}{\sqrt{1 + \left(\frac{E_{k,l}^{(j)} - E_{m,n}^{(j)}}{H_{k,l;m,n}^{(j)}}\right)^2}}$$

Perform calculation using an assignment scheme based on

$$\left|\left\langle oldsymbol{\phi}_{k}^{\left(j
ight)}\left|oldsymbol{\Psi}_{m,n}^{\left(j
ight)}
ight
angle
ight|^{2}$$

Perform calculation using an assignment scheme based on

$$\left|\left\langle oldsymbol{\phi}_{k}^{\left(\,j
ight)}\left|\Psi_{m,n}^{\left(\,j
ight)}
ight.
ight|^{2}$$

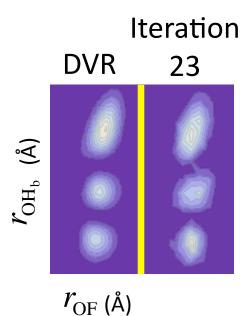
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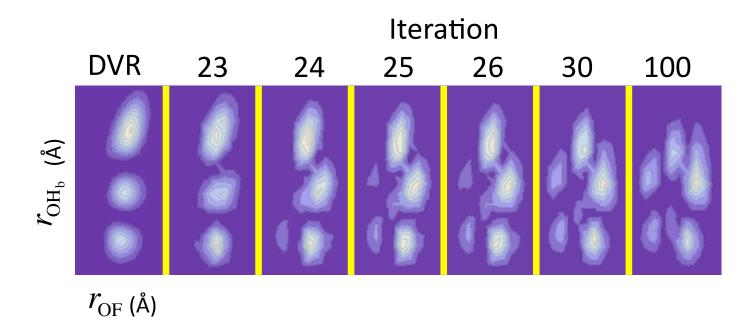
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ight)}\left|oldsymbol{\Psi}_{m,n}^{\left(\,j
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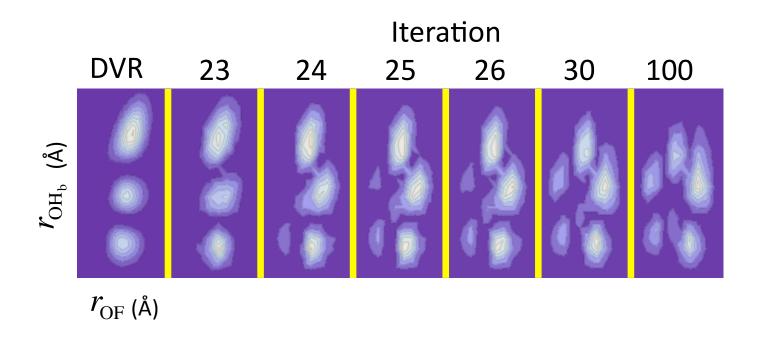


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ight)}\left|oldsymbol{\Psi}_{m,n}^{\left(\,j
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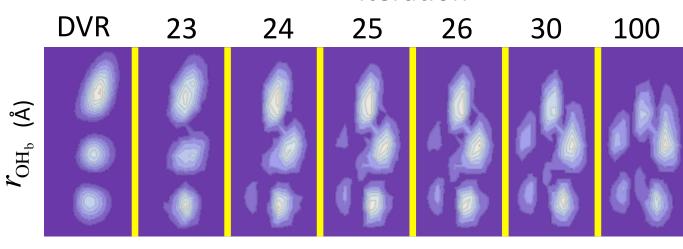
Contamination occurs under the radar of overlaps, transition moments, moment distributions, etc.



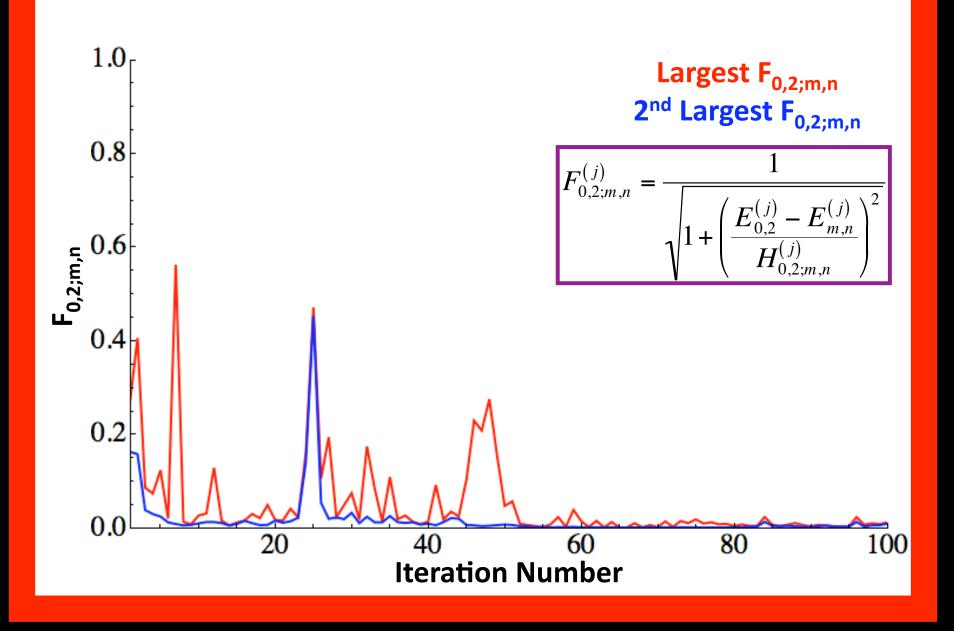
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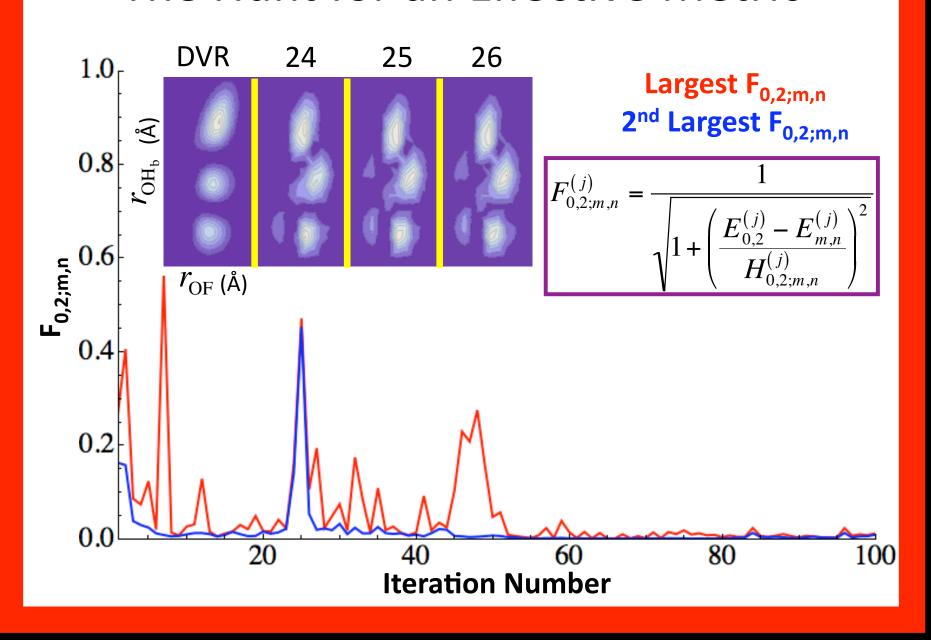
$$F_{k,l;m,n}^{(j)} = \frac{1}{\sqrt{1 + \left(\frac{E_{k,l}^{(j)} - E_{m,n}^{(j)}}{H_{k,l;m,n}^{(j)}}\right)^2}}$$

Iteration



 $r_{
m OF}$ (Å)





Developed a methodology that uses importance sampling Monte Carlo and and evolving basis to intelligently sample configuration space and minimize size of basis required

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Promising developments in tackling the problem of multidimensional eigenstate assigment

Complete the optimization of the multi-dimensional algorithm's treatment of spurious resonances

Further multi-dimensional benchmarking

Developed a methodology that uses importance sampling Monte Carlo and and evolving basis to intelligently sample configuration space and minimize size of basis required

The method has been shown to be able to accurately describe a highly anharmonic Morse oscillator

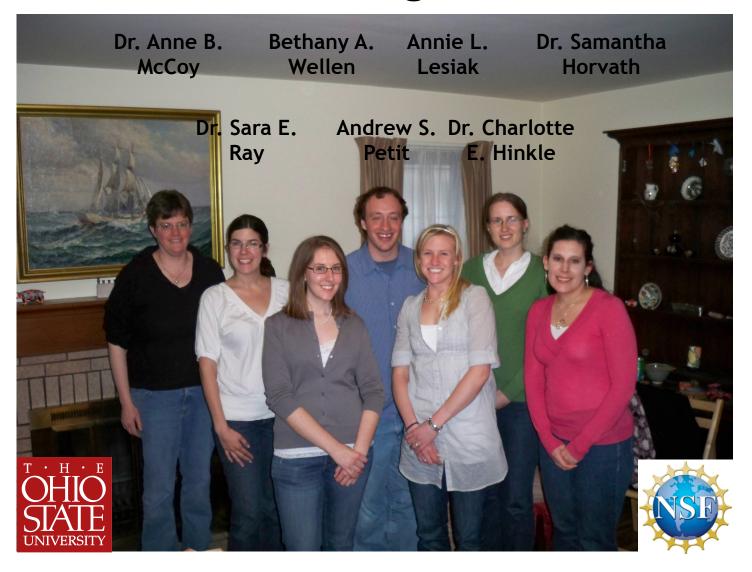
Promising developments in tackling the problem of multidimensional eigenstate assigment

Complete the optimization of the multi-dimensional algorithm's treatment of spurious resonances

Further multi-dimensional benchmarking

Couple algorithm with ab initio electronic structure calculations

Acknowledgements



The Initial Basis (2D Example) $\Psi_{0,0}^{(1)} \propto e^{-\frac{\alpha_1 r^2}{2}} e^{-\frac{\alpha_2 R^2}{2}}$

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Basis functions change each iteration

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$$\overline{\Psi}_{m,n}^{(1)} = \sum_{j=1}^{m} \kappa_{m,n;m-j,n}^{(1)} r^{j} \Psi_{m-j,n}^{(1)} + \sum_{j=1}^{m} \kappa_{m,n;m,n-j}^{(1)} R^{j} \Psi_{m,n-j}^{(1)}$$

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Phase factors

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Phase factors

Will be made orthogonal to all previously built states and normalized

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$$\overline{\Psi}_{3,0}^{(1)} = \kappa_{3,0;2,0}^{(1)} r \Psi_{2,0}^{(1)} + \kappa_{3,0;1,0}^{(1)} r^2 \Psi_{1,0}^{(1)} + \kappa_{3,0;0,0}^{(1)} r^3 \Psi_{0,0}^{(1)}$$

$$\overline{\Psi}_{m,n}^{(1)} = \sum_{j=1}^{m} \kappa_{m,n;m-j,n}^{(1)} r^{j} \Psi_{m-j,n}^{(1)} + \sum_{j=1}^{n} \kappa_{m,n;m,n-j}^{(1)} R^{j} \Psi_{m,n-j}^{(1)}$$

$$\overline{\Psi}_{m,n}^{(1)} = \sum_{j=1}^{m} \kappa_{m,n;m-j,n}^{(1)} r^{j} \Psi_{m-j,n}^{(1)} + \sum_{j=1}^{n} \kappa_{m,n;m,n-j}^{(1)} R^{j} \Psi_{m,n-j}^{(1)}$$

$$\overline{\Psi}_{2,1}^{(1)} = \kappa_{2,1;2,0}^{(1)} R \Psi_{2,0}^{(1)} + \kappa_{2,1;1,1}^{(1)} r \Psi_{1,1}^{(1)} + \kappa_{2,1;0,1}^{(1)} r^2 \Psi_{0,1}^{(1)}$$