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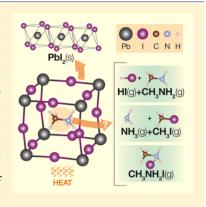
Perspective

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1 Thermodynamics and the Intrinsic Stability of Lead Halide 2 Perovskites CH₃NH₃PbX₃

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ABSTRACT: The role of thermodynamics in assessing the intrinsic instability of the CH₃NH₃PbX₃ perovskites (X = Cl₂Br₂I) is outlined on the basis of the available experimental information. Possible decomposition/degradation pathways driven by the inherent instability of the material are considered. The decomposition to precursors $CH_3NH_3X(s)$ and $PbX_2(s)$ is first analyzed, pointing out the importance of both the enthalpic and the entropic factor, the latter playing a stabilizing role making the stability higher than often asserted. For CH3NH3PbI3, the disagreement between the available calorimetric results makes the stability prediction uncertain. Subsequently, the gasreleasing decomposition paths are discussed, with emphasis on the discrepant results presently available, probably reflecting the predominance of thermodynamic or kinetic control. The competition between the formation of $NH_3(g) + CH_3X(g)$, $CH_3NH_2(g) +$ HX(g) or CH₃NH₃X(g) is analyzed, in comparison with the thermal decomposition of methylammonium halides. In view of the scarce and inconclusive thermodynamic studies to-date available, the need for further experimental data is emphasized.



t is no exaggeration to say that the largest part of the current research efforts on lead halide-based and similar perovskite 21 materials is directed toward the search for higher stability 22 needed in photovoltaic applications. In the last five years, many 23 authors expressed the concern that the low stability under the 24 action of a number of external agents, including other device 25 components, could be the main Achilles' heel of this class of 26 light harvester materials, very attractive in other respects, 27 whose prototype is the well-known methylammonium lead 28 iodide, CH₃NH₃PbI₃. As a consequence, a wealth of 29 strategies were put in place to improve the material and 30 device stability, based on chemical modifications^{7–9} protection 31 layers, 10 and encapsulation. 11

Interaction with water/moisture has been soon identified as 33 a major drawback. 12 Other external agents which many 34 researchers focused on are oxygen and UV/visible radiation. 35 The study of these degradation processes was typically 36 performed by following the change of properties/performances 37 of the material/device, while the interaction takes place or ex 38 post. To elucidate the progress and mechanism of degradation, 39 a "microscopic" approach is most often applied, based on 40 techniques such as XRD, UV-vis and NIR spectroscopy, 41 fluorescence, microscopy, XPS, etc. 14-17 Theoretical calcu-42 lations may be of help in identifying mechanistic details. 18 The 43 effect of temperature on the perovskite stability was usually 44 studied in conjunction with that of such chemical and physical 45 agents. Incidentally, it should be mentioned that in some cases 46 the diagnostic means used to study degradation (e.g., X-ray 47 irradiation, electron currents) can themselves play a role in 48 degradation phenomena. 19,20

Comparatively little work has been carried out on the 50 intrinsic (in)stability of these materials in itself, both in vacuo 51 and under inert atmosphere, as a function of temperature. 21-24 In particular, very few studies are available based on a 52 macroscopic thermodynamic approach.^{25–27} Moreover, the 53 available experimental and computational results are often 54 discrepant. The goal of this Perspective is to discuss the issue 55 of the intrinsic stability of CH3NH3PbX3 materials from a 56 thermodynamic point of view in light of the information 57 currently available, pointing out the persistent uncertainties 58 and inconsistencies, which make urgent further experimental 59 efforts. Special focus is done on the decomposition processes 60 leading to the release of gaseous products.

Overall, we would like to emphasize in this paper the 62 contribution that classic macroscopic thermodynamics can (or 63 cannot) provide to the stability issues raised from the 64 application of these materials to energy conversion technolo- 65 gies. It is important to underline that the thermodynamic 66 characterization of a material as such is a crucial prerequisite to 67 undertake reliable thermodynamic predictions and simulations 68 of its behavior in various chemical and physical environments. 69 Furthermore, ascertaining the intrinsic (in)stability of a 70 material is essential for practical applications, because if the 71 material is found to be inherently unstable, any protection 72 strategy may be undermined. Regrettably, from an analysis of 73 the literature trends, one has the clear impression that 74 experimental thermodynamic studies cannot keep pace with 75 the wealth of new material modifications that are proposed at 76 an ever-increasing rate to overcome instability issues.

In the broadest sense, investigating the thermodynamic 78 stability of a material means to wonder if, under a given set of 79 external conditions, the material will remain unchanged or it 80

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81 will undergo some kind of chemical/physical transformation, 82 ultimately driven by entropy production. The number of 83 possible transformation pathways of the system is, in general, 84 high and difficult to predict. In a stricter and more practical 85 sense, one usually evaluates the thermodynamic driving force ₈₆ or affinity, $\Delta_{\rm r}G = \sum_i \nu_i \mu_i$, for one or a few among the many 87 possible decomposition/degradation reactions, where ν_i and μ_i 88 are, respectively, the stoichiometric coefficients (taken as 89 negative for the left-hand reactants of the chemical/physical 90 degradation process) and the chemical potentials of all the 91 species involved. If the affinity is found to be negative under 92 the conditions of interest (for example, for given pressure and 93 temperature), the selected decomposition/degradation path is 94 thermodynamically favored. Otherwise, the material is stable as 95 far as that path is considered, and indeed its formation from 96 right-hand products is thermodynamically favored. In this 97 approach, the final products of the process are to be known or 98 an hypothesis has to be done. In the case of CH₃NH₃PbX₃ 99 compounds, the decomposition to the synthesis precursors was 100 most often considered in theoretical evaluations (Figure 1):

$$CH_3NH_3PbX_3(s) = CH_3NH_3X(s) + PbX2(s)$$
 (1)

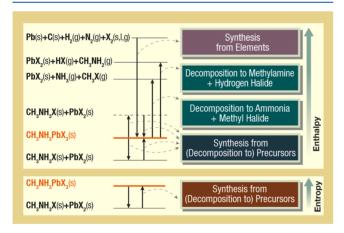


Figure 1. Enthalpy and entropy level scheme for possible formation/decomposition processes of CH₃NH₃PbX₃ perovskites. Decomposition to precursors can be both endothermic and exothermic. The minor process leading to the formation of CH₃NH₃X(g) (see Figure 3) is not shown. Enthalpy levels are not to scale.

Although considering this reaction is probably the most natural choice in assessing the formability of $CH_3NH_3PbX_3$ compounds, it should be noted that, while $PbI_2(s)$ was much often reported as the main decomposition product under various conditions, apparently no experimental study reported the CH_3NH_3X solids among the observed products.

A huge number of theoretical calculations were performed to evaluate the energy/enthalpy change of the above reaction by

the DFT approach. However, the results are significantly 110 dependent on the chosen functional, the best performance 111 being usually obtained with the PBEsol one, with the 112 additional inclusion of spin—orbit contributions. 28,29 A fairly 113 rich selection of results is reported in Table 1. As for 114 t1 experiments, regrettably, only two direct determinations of 115 $\Delta_{\rm r}H^{\circ}(1)$ are available in the literature, both obtained by 116 solution calorimetry at T=298 K, using DMSO 25 or aqueous 117 HCl^{26} as a solvent. Note that, although at 298 K CH₃NH₃PbI₃ 118 is stable in the tetragonal form, measurements of ref 25 were 119 performed on the (metastable) cubic high temperature phase. 120 A third experimental value for reaction 1 can be derived from 121 the vapor pressure measurements carried out by effusion-based 122 techniques on the equilibrium (reaction 8) discussed in the 123 next section. 27

A compilation of all the results is reported in Table 1. 125 Somehow surprisingly, the two calorimetric determinations are 126 not in agreement, showing a discrepancy definitely outside the 127 claimed experimental uncertainties. For CH3NH3PbI3, in 128 particular, the very negative value of $\Delta_r H^{\circ}(1)$ found in ref 129 26., which led those authors to claim the instability of the 130 compound, was not confirmed by the subsequent measure- 131 ments in DMSO,²⁵ making difficult any conclusive prediction 132 on the spontaneous direction of reaction (1) at room 133 temperature for X = I. Also the stability trend from Cl to Br 134 to I is not in agreement between the two studies (Figure 2). 135 f2 The values of ref 26 suggest the Cl > Br > I stability trend, 136 which is consistent with the Goldschmidt's tolerance factor 137 traditionally used to rationalize the stability of perovskite 138 phases. However, tensimetric results²⁷ agree well with the 139 Ivanov's data,²⁵ which is a nice occurrence in view of the 140 completely different experimental approach used.

While the enthalpic term is often the most important factor 142 p driving the thermodynamic direction of an isothermal chemical 143

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reaction, it may well be that the entropic factor comes into play 144 and affects the spontaneity evolution in a decisive manner. 145 This is especially true for reactions involving gaseous phases, 146 aggregation/disaggregation processes, high temperatures, and 147 when enthalpic effects are small, which is the case of reactions 148 (1). Regrettably, in the last few decades the number of papers 149 presenting heat capacity and absolute entropy measurements is 150 decreasing, making thermodynamic evaluations more difficult 151 and less accurate. It is thus a really lucky occurrence that 152 almost 30 years ago, when lead halide perovskites were far 153 from bursting on the scientific scene, Suga and co-workers 154 published low temperature heat capacity data for the 155 CH₃NH₃PbX₃ crystal phases. Absolute entropies of the 156 other compounds involved in reaction 1 are known with the 157 exception of CH₃NH₃Br, whose entropy was estimated by us 158 by the empirical Volume-Based-Thermodynamic approach.

Table 1. Energetic and Thermodynamic Properties of the Decomposition Reactions of $CH_3NH_3PbX_3$ Perovskites to Solid Precursors: $CH_3NH_3PbX_3(s) = CH_3NH_3X(s) + PbX_2(s)^a$

	theory (DFT)		thermodynamic experiments					
X	ΔE^{b}	ref.	$\Delta H_{298 ext{K}}^{\circ}$	method and ref	S _{298K} ° €	$\Delta S_{298K}^{\circ}^{d}$	$\Delta G_{298 ext{K}}^{\circ}$	stable with respect to precursors at 298 K, 1 bar^e
Cl (cubic)	12	30	9.03 ± 1.68	solution calorim. in HCl ²⁶	313.37	- 38.77	11.6	YES
	68	31	4.45 ± 0.34	solution calorim. in DMSO ²⁵			16.0	YES
	0.39-3.9	32	2.8 ± 7.8	vapor pressure (KEML, KEMS) ²⁷			11.6	YES
Br (cubic)	12	30	-6.69 ± 1.41	solution calorim. in HCl ²⁶	349.29	- 39.9	5.2	YES
	24	31	6.78 ± 0.97	solution calorim. in DMSO ²⁵			18.7	YES
	1.4-4.1	32	3.3 ± 8.7	vapor pressure (KEML, KEMS) ²⁷			15.2	YES
I (tetragonal)	2.7 - 3.4	22		solution calorim. in HCl ²⁶	374.15	-39.6	-22.7	NO
	9.6	33	- 34.50 ± 1.01					
	2.2	29						
	4.8	34						
	-8.7	35	-1.913 ± 1.12^{f}	solution calorim. in DMSO ²⁵			9.9	YES
	5.8	36						
	9.6	31						
	-5.8/-6.1	32						
	0.39	37	0.39 ± 9.7	vapor pressure (KEML,			11.4	YES
	2.4	38		KEMS) ²⁷				
	3.9	39						
I (cubic) ^g	-4.8	28	-4.493 ± 1.12	solution calorim. in DMSO ²⁵	383.85 ^c	-49.3	10.2	YES
	-1.9	30						
	26	40						
	-11/-12	32						

^aEnergies and enthalpies are in kJ/mol, entropies in J/K mol. ^bFor an accurate comparison of the theoretical ΔE s with the thermochemical ΔH_{298K}° small effects due to zero-point energy, finite temperature and standard pressure should be considered. The corrections due to the heat content difference ($H_{298~K}-H_{0K}$) and the $P^{\circ}\Delta V$ terms are of the order of few kJ/mol and J/mol, respectively. ^cAbsolute entropies of CH₃NH₃PbX₃ are from ref 41. The value for the high temperature cubic phase of CH₃NH₃PbI₃ was estimated by adding the tetragonal-cubic transition entropy (9.7 J/K mol, measured at 330 K⁴¹) to S_{298K}° of the tetragonal phase. ^dEntropies of CH₃NH₃X(s) and PbX₂(s) are from the compilation of ref 25. except for CH₃NH₃Br, whose S_{298K}° was estimated by us as 148.2 J/K mol by a volume-based-thermodynamics approach. ⁴² ^eBased on the sign of ΔG_{298K}° which for reactions (1) corresponds to the driving force $\Delta_r G$ at 298 K and 1 bar. ^fEvaluated by adding the t-c transition enthalpy (2.58 kJ/mol, measured at 330 K⁴¹) to the ΔH_{298K}° value measured for the cubic phase. ²⁵ ^gThe tetragonal to cubic transition enthalpy of CH₃NH₃PbI₃ is 2.58 kJ/mol, measured at 330 K. ⁴¹

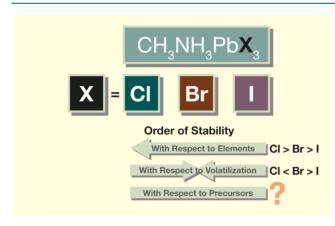


Figure 2. Stability order of $CH_3NH_3PbX_3$ for $X = Cl_3Br_3I$ with respect to different processes.

160 The so-derived entropy changes $\Delta_r S^{\circ}(1)$ at 298 K, also 161 reported in Table 1, are significantly negative (Figure 1) and 162 therefore give a large positive contribution to the decom-163 position Gibbs energy change, in most cases larger than the 164 enthalpic term, so stabilizing the perovskite phases with respect

to precursors. Furthermore, since
$$\left(\frac{\partial \Delta_{\rm r} G^{\circ}}{\partial T}\right)_{p=p^{\circ}} = -\Delta_{\rm r} S^{\circ}$$
,

negative entropy changes cause $\Delta_r G^\circ$ of reaction 1 to increase 166 with T, and the decomposition to become more and more 167 disfavored with increasing temperature. It is interesting to note 168 that the largely negative $\Delta_r S^\circ(1)$ values are related to the high 169 absolute entropies of the CH₃NH₃PbX₃ phases, in turn due to 170 the high entropy changes associated with the phase transitions 171 in the crystal.

Since reaction 1 involves pure solid phases, $\Delta_r G^\circ$ is 173 practically coincident with the thermodynamic driving force 174 for decomposition, $\Delta_r G$, provided that pressure is not too high. 175 The last column of Table 1 indicates the stability of perovskite 176 compounds with respect to precursors of reaction 1. On the 177 basis of the available information, it is concluded that for X=178 Cl and Br decomposition is thermodynamically disfavored at 179 298 K, whereas for X=I the results are conflicting: depending 180 on the selected enthalpy change, the negative entropic term 181 can compensate it or not. On the basis of the agreement with 182 tensimetric values, it seems reasonable to recommend 183 provisionally the calorimetric results of Ivanov et al., 25 which 184 provide the following expressions for the $\Delta_r G^\circ$ of decom-185 position reaction 1:

$$\Delta_{\rm r}G^{\circ}({\rm CH_3NH_3PbCl_3})$$
 (1) = (4.45 + 38.7710⁻³T) kJ/mol (2) ₁₈₇

$$\Delta_{\rm r} G^{\circ}({\rm CH_3NH_3PbBr_3}) (1) = (6.78 + 39.9010^{-3}T) \,{\rm kJ/mol}$$
 (3) ₁₈₈

$$\Delta_{\rm r} G^{\circ}({\rm CH_3NH_3PbI_3})$$
 (1) = (- 1.91 + 39.6010⁻³T) kJ/mol (4)

190 valid for the room temperature phases (cubic phase for 191 CH₃NH₃PbCl₃ and CH₃NH₃PbBr₃ and tetragonal for 192 CH₃NH₃PbI₃) in a reasonably large temperature range (the 193 temperature dependence of $\Delta_r H^{\circ}$ and $\Delta_r S^{\circ}$ is not accounted 194 for in the above equations), including temperatures of practical 195 interest for photovoltaic devices (typically in the range 30-80 196 °C). According to these expressions, the pervoskite phases are 197 stable at room temperature, and their stability increases at 198 higher temperature. In principle, by putting $\Delta_r G$ equal to zero, 199 the lowest temperature at which the perovskite phases are 200 stable can be roughly estimated. These temperatures 201 correspond to the invariant temperatures (at $P = P^{\circ}$) where 202 perovskites would coexist in equilibrium with their precursors. 203 Simple linear extrapolation of eqs 2-4 gives temperatures 204 extremely low or even negative, indicating that no decom-205 position occurs in practice (obviously, both CH₃NH₃PbX₃ and 206 the decomposition products would undergo transitions to the 207 low-temperature phases at the corresponding transition 208 temperatures). It is then clear from this analysis, that 209 perovskite phases are "high temperature compounds" in the 210 pseudobinary PbX₂-CH₃NH₃X phase diagram, and reaction 1 211 should not be regarded as a "thermal" decompositions at all. 212 However, if the $\Delta_r H^{\circ}$ value of ref 26 is used for the iodide 213 phase, a minimum temperature of stability of 870 K is 214 estimated by rough extrapolation (even higher than the 215 melting point of PbI₂), meaning that the material is completely 216 unsuitable for photovoltaic applications (unless strong kinetic 217 hindrance comes into play).

As mentioned, ascertaining the stability of a material referring to only one or several processes is not a conclusive proof of its absolute intrinsic stability. For example, in order to identify the most favored among a number of possible decomposition pathways of $CH_3NH_3PbI_3$, a convex hull approach was used, leading to select the decomposition to $NH_4I + PbI_2 + CH_2$ as the most stable path. In this connection, it is interesting to note that the cleavage of the $NH_4I + NH_4I + NH_$

Another process that is sometimes taken as representative of the intrinsic stability of a material is the opposite of the 131 formation from elements, Pb + 1.5 X_2 + 3 H_2 + 0.5 N_2 + C = 132 $CH_3NH_3PbX_3$, with all species in their reference phase (at the 133 temperature of interest) and in the standard state (see Figure 134 1). Consistent with eqs 2–4, the following expressions can 135 be derived for $\Delta_f G^\circ$ of the $CH_3NH_3PbX_3$ compounds, valid for 136 the room temperature phases in a reasonably large temperature 137 range:

$$\Delta_{\rm f} G^{\circ}({\rm CH_3NH_3PbCl_3}) = (-662.2 + 579.110^{-3}T) \,{\rm kJ/mol}$$
²³⁸ (5)

$$\Delta_{\rm f} G^{\circ}({\rm CH_3NH_3PbBr_3}) = (-543.1 + 437.010^{-3}T) \,{\rm kJ/mol}$$

$$\Delta_{\rm f} G^{\circ}({\rm CH_3NH_3PbI_3}) = (-371.6 + 358.110^{-3}T) \, {\rm kJ/mol} \end{cases} \label{eq:deltafine}$$

These equations indicate that $\Delta_f G^{\circ}$ is strongly negative at temperatures close to room temperature and show the expected trend of stability with respect to elements, Cl > Br (Figure 2). Equations 5–7 can be combined with the

corresponding expressions for other substances (H_2O , PbO, 245 PbCO₃, HI, CH₃NH₂, etc.) to estimate the standard Gibbs 246 energy change of chemical reactions potentially involved in the 247 intrinsic or extrinsic (e.g., due to water or oxygen) degradation 248 of the materials as well in synthesis and annealing processes. 249 Note that for several important reactions (for example, those 250 forming perovskite hydrate phases such as $(CH_3NH_3)_4PbI_6$. 251 $2H_2O$ under exposure to moisture) Δ_rG° cannot be evaluated 252 owing to the lack of relevant thermodynamic data.

In view of the above discussion, the decomposition 254 processes (1) should not be a worrisome decomposition 255 path as far as methylammonium lead chloride and bromide 256 perovskites are concerned, whereas the discrepancy between 257 calorimetric data makes a conclusive assessment for iodide 258 difficult. However, other decomposition channels could be at 259 work under operative conditions. In particular, gas-releasing 260 decomposition processes are of crucial importance for 261 investigating the stability of CH3NH3PbX3 and similar 262 materials. Grazing-incidence wide-angle X-ray diffraction 263 measurements have shown recently 46 that the dramatic heat- 264 induced performance decrease of encapsulated perovskite- 265 based devices is due to surface modifications related to the 266 intercalation of thermally decomposed methylammonium 267 fragments into PbI₂ planes. Since encapsulation is expected 268 to prevent interaction with external agents, this is a clear 269 evidence of intrinsically driven degradation related to the loss 270 of volatile fragments, as already suggested in previous 271 papers. 22,24 Furthermore, the type of gas that these materials 272 tend to lose under heating may be of interest for real devices 273 because gaseous products could interact with the sealing 274 materials and with the other components of the cell. Finally, 275 gas-phase releasing degradation is very important under 276 conditions where the system is allowed to vaporize, such as 277 postsynthesis annealing⁴⁷ and synthesis by vapor deposition ²⁷⁸ techniques.⁴⁸ In spite of this, the direct experimental study of 279 the released gaseous species has been the subject of relatively 280 few studies that, unfortunately, presented problematic 281 results. 27,49,50

In this connection, the following gas-releasing processes are 283 worth considering:

$$CH_3NH_3PbX_3(s) = PbX_2(s) + HX(g) + CH_3NH_2(g)$$
(8) 285

$$CH_3NH_3PbX_3(s) = PbX_2(s) + CH_3X(g) + NH_3(g)$$
 (9) ₂₈₆

$$CH_3NH_3PbX_3(s) = PbX_2(s) + CH_3NH_3X(g)$$
 (10) ₂₈₇

In all three cases, solid lead dihalide is formed, as invariably 288 observed by means of solid state techniques, 27 and the 289 $\rm CH_3NH_3X$ portion of the perovskite phase is lost, either as 290 undissociated methylammonium halide or in the form of 291 smaller molecules. In principle, gas-phase dissociation can 292 occur by formation of HX or $\rm CH_3X$, depending on whether the 293 methyl group or the proton associates with the halide ion. 294 Theoretical analyses indicate as an energetically favored path 295 to this kind of degradation the creation of HI vacancies and the 296 subsequent combination of the amine fragment with Pb atoms, 297 which disintegrates the inorganic framework. 51

That the heat-induced degradation of CH₃NH₃PbI₃ perov- 299 skites proceeds through mass loss has been shown by a number 300 of thermogravimetric (TGA) measurements and has been also 301 inferred by solid state techniques, for example by measuring 302

303 the evolution of I/Pb and N/Pb ratios by photoelectron 304 spectroscopy.²⁰

Most investigations of the degradation of $CH_3NH_3PbX_3$ to 306 volatile species were carried out by classic TGA, where the

Most investigations of the degradation of CH₃NH₃PbX₃ to volatile species were carried out by classic TGA, where the mass loss rate is recorded as a function of temperature, usually under an inert dynamic atmosphere

307 mass loss rate is recorded as a function of temperature, usually 308 under an inert dynamic atmosphere. A number of early TGA measurements $^{23,52-54}$ led to the conclusion that thermal 310 decomposition of CH₃NH₃PbX₃ compounds begins at rather 311 high temperature, above 200-250 °C, giving support to the 312 view that the intrinsic thermal stability was not a major 313 drawback in real applications where temperature does not 314 exceed 80-90 °C. Various authors also showed that mass loss 315 starts at lower temperatures for chloride and mixed iodide-316 chloride compounds compared to bromide and iodide. ^{26,54} In 317 the same years some papers based on XRD and other solid-318 state techniques were less optimiztic^{22,55} pointing out that also 319 at temperatures as low as 85-100 °C appreciable decom-320 position occurred even under an inert atmosphere, especially 321 under prolonged heating. Actually, the dynamic nature of TGA 322 experiments, especially if the used scan rate is not very low, 323 could be insufficient for assessing the long-term stability of a 324 photovoltaic material, which is required to work for a long time 325 at temperatures lower than the decomposition temperatures 326 detectable by TGA curves. Furthermore, the TGA method 327 does not permit distinguishing among processes 8−10, and the 328 attribution of the chemical nature of the gaseous phase has 329 been basically speculative or indirect in the literature, with 330 reaction 8 generally preferred on the basis of chemical wisdom 331 (acid-base interaction between the organic cation and the 332 halide anion). Two-step TGA curves observed by some 333 authors for $CH_3NH_3PbI_3$ suggested the sequential loss of 334 HI(g) and $CH_3NH_2(g)$, 23 supporting this hypothesis.

To the best of our knowledge, the first reports on the direct detection of the gas phase released by the CH₃NH₃PbX₃ perovskites appeared in 2016, ^{27,49,54} although FTIR spectros-338 copy experiments were previously reported for the vaporization of dimethylformamide-CH₃NH₃PbCl₂I_{3-x} solutions. ⁵⁶

Extensive vaporization experiments were first reported based on the classic Knudsen Effusion Mass Loss (KEML) and Knudsen Effusion Mass Spectrometry (KEMS) techniques. KEMS measurements were carried out on all three CH₃NH₃PbX₃ compounds in the overall temperature range found in TGA measurements. Mass spectra showed the large dominance of peaks attributable to HX(g) and CH₃NH₂(g), previously proposed in ref 54 based on Temperature-sponding to undissociated CH₃NH₃X(g) were also detected (it standard countries and countries should be pointed out that electron impact mass spectrometry could cause the undissociated CH₃NH₃X(g) species to break into smaller fragment ions.). The thermodynamic analysis of

partial pressure data allowed decomposition enthalpies to be 355 derived for reaction 8 and, thereafter, formation enthalpies of 356 CH₃NH₃PbX₃ perovskites to also be evaluated. The latter were 357 subsequently found to agree well with calorimetric results (see 358 Table 1). 25,50 Shortly after, another study 49 was published 359 where a very different behavior was observed by TGA-Mass 360 Spectrometry experiments for CH3NH3PbI3 and for its 361 precursor halide, CH3NH3I, which were found to release 362 only $NH_3(g)$ and $CH_3I(g)$. Although the largest part of TGA- 363 MS experiments were carried out at much higher temperatures 364 (300-420 °C) than Knudsen measurements, authors obtained 365 some indication that the same process would also take place at 366 temperatures as low as 80 °C, close to those of interest for 367 photovoltaic applications. Interestingly, the findings of ref 49 368 confirmed in part those previously obtained by FTIR.⁵⁶ FTIR ³⁶⁹ spectra of the gas phase recorded at 265 °C indicated the 370 decomposition of solid CH3NH3I to NH3(g) and CH3I(g), in 371 contrast with CH₃NH₃Cl and CH₃NH₃PbCl₃, which was 372 observed to release HCl(g) and CH₃NH₂(g).

In order to shed some light on this scanty and discrepant 374 experimental information, a thermodynamic analysis is useful. 375

In principle, reactions 8–10 may occur simultaneously, 376 giving a three-phase monovariant equilibrium with seven 377 components, two of them thermodynamically independent. 378 While the absolute values of the partial pressures depend on 379 the properties of the two condensed phases, the thermody- 380 namic competition between the three reactions is basically due 381 to the different stability of the gaseous species.

Using relations 5–7 for the Gibbs energy of formation of the 383 perovskite phases and the corresponding well-established 384 expressions for the products of reactions 8 and 9, the results 385 reported in Table 2 are derived for the corresponding standard 386 t2

Table 2. Standard Gibbs Energy Changes ($\Delta_r G^{\circ}$, in kJ/mol) for the Decomposition Reactions of CH₃NH₃PbX₃ and CH₃NH₃X with Release of Gaseous Products^a

X	$CH_3NH_3PbX_3(s) = PbX_2(s) + HX(g) + CH_3NH_2(g)$ (8)	$CH_3NH_3PbX_3(s) = PbX_2(s) + CH_3X(g) + NH_3(g)$ (9)
Cl	$187.9 - 252.3 \times 10^{-3} T$	$174.9 - 249.6 \times 10^{-3} T$
Br	$206.9 - 253.3 \times 10^{-3} T$	$183.3 - 250.3 \times 10^{-3} T$
I^{b}	$200.2 - 250.1 \times 10^{-3} T$	$164.7 - 247.2 \times 10^{-3} T$
	CH3NH3X(s) = HX(g) + CH3NH2(g)	$CH_3NH_3X(s) = CH_3X(g) + NH_3(g)$
Cl	$183.5 - 291.1 \times 10^{-3} T$	$170.5 - 288.3 \times 10^{-3} T$
Br	$200.1 - 293.3 \times 10^{-3} T$	$176.6 - 290.3 \times 10^{-3} T$
I	$204.7 - 289.7 \times 10^{-3} T$	$169.2 - 286.8 \times 10^{-3} T$

^aThese expressions can be applied in a reasonably large temperature range near to 298 K. ^bCubic phase of CH₃NH₃PbI₃.

Gibbs energy changes. Data in Table 2 indicate that, at 387 variance with the trend of formation enthalpies, bromide is the 388 most stable with respect to vaporization decomposition 389 (Figure 2) for both processes 8 and 9. However, the stability 390 trend is Br > I > Cl and Br > Cl > I, respectively, for the two 391 processes. Note that only in the case of reaction 8 is the 392 stability order in agreement with TGA experiments (see 393 above). For the sake of comparison, in the same table, the 394 $\Delta_r G^\circ$'s are also reported for the corresponding decomposition 395 reactions of methylammonium halides, CH₃NH₃X. It is 396 interesting to note that $\Delta_r G^\circ$ values for the latter are lower 397 (i.e., the decomposition pressures are higher for a given 398 temperature) than those of the corresponding perovskites. A 399

400 large part of this effect is due to the higher entropy changes for 401 the decomposition of CH₂NH₂X compounds.

A similar analysis for process 10 is made difficult by the lack 403 of thermodynamic data for the species CH₃NH₃X(g), which is 404 not surprising if we consider that, even for the simpler 405 ammonium halide species, NH₄X(g), experimental data are 406 more uncertain than one might believe. For example, the 407 dissociation degree of $NH_4Cl(g)$ to $NH_3(g)$ and HCl(g) has 408 been the subject of a longstanding debate, with experimental 409 results ranging from complete to very limited dissociation and 410 theoretical studies claiming for a fair stability of the 411 undissociated hydrogen-bond linked species. 57 Although the 412 CH₃NH₂ + HCl potential energy surface has been the subject 413 of a number of theoretical studies aimed at clarifying the 414 mechanism of the methyl exchange in solution phase (the so-415 called Menshutkin S_N2 reaction, CH₃Cl + NH₃ = 416 CH₃NH₃+Cl⁻), apparently a stable structure for the 417 CH₃NH₃Cl complex in the gas phase was reported only 418 recently by Patterson⁵⁸ based on DFT calculations. In order to 419 evaluate the relative importance of reaction 10 for the iodide 420 perovskite under thermodynamic conditions, similar calcu-421 lations were done for the CH₃NH₃I(g) species, ⁵⁹ that allowed 422 us to derive an approximate estimate of $\Delta_r G^{\circ}(10)$ as

$$\Delta_{\rm r} G^{\circ}(10) = (143.8 - 111.310 - 3T) \, \text{kJ/mol} \quad (X = I)$$
⁴²³

424 (cubic phase of $CH_3NH_3PbI_3$). From the equations in Table 2 425 and eq 11, the partial pressures of all the gaseous species in 426 equilibrium with $CH_3NH_3PbI_3$ and PbI_2 were estimated. The 427 so-derived *total* pressures produced from reactions 8–10 are 428 reported in Figure 3 along with the total pressures measured by 429 the classic Knudsen effusion method $^{2.7,60}$ (the only exper-430 imental value available to date). This plot shows clearly that, 431 under the hypothesis of thermodynamic equilibrium, the 432 decomposition channel of $CH_3NH_3PbI_3$ leading to $NH_3 + 433$ CH_3I (process 9) should be by far the most important,

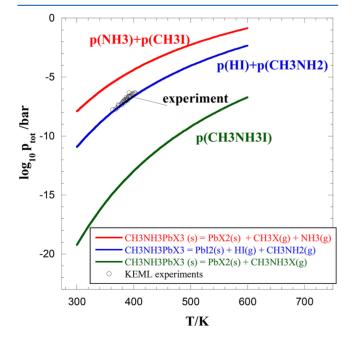


Figure 3. Total pressures produced by decomposition processes 8-10 evaluated from calorimetric data, compared to the results of KEML experiments. 27,60

especially at moderate temperatures, followed by that one 434 releasing CH₃NH₂ + HI and eventually by the evaporation to 435 the undissociated species, which is negligible at any temper- 436 ature of interest. However, the decomposition pressure due to 437 the loss of NH₃ + CH₃I should be 2–3 orders of magnitude 438 higher than that experimentally observed under effusion 439 conditions, which instead agrees well with the occurrence of 440 process 8. It should be pointed out that our thermodynamic 441 prediction is entirely based on calorimetric results (eqs 5–7). 442 Being such data are completely independent from tensimetric 443 measurements, the agreement with the latter is very 444 satisfactory. It should also be pointed out that our calculation 445 for process 10 is based on computationally estimated 446 thermodynamic data.

Overall, from this analysis the conclusion can be drawn that 448 (i) process 9 has a much larger thermodynamic driving force, 449 but (ii) under effusion conditions, which, in principle, should 450 allow an approach toward thermodynamic equilibrium, the 451 occurrence of this process is kinetically hindered, and the 452 measured pressures suggests that process 8 takes place instead. 453 The kinetic limitation of reactions 9 seems quite plausible 454 because the breaking of a strong C-N bond (330 kJ/mol at 455 298 K) is required, instead of the hydrogen-bond breaking 456 involved in process 8. This view is supported by DFT 457 calculations, 49 which predict a remarkable activation barrier for 458 reaction 9, and by the comparison with the thermal 459 decomposition behavior of simple alkylammonium halides 460 (see below). It should be noted that the very good agreement 461 between calculated and experimental data indicates that, as far 462 as process 8 is concerned, thermodynamic equilibrium is fully 463 attained under Knudsen conditions, giving confidence in the 464 thermodynamic properties derived from tensimetric measure- 465 ments, provided that data are analyzed on the basis of the 466 proper reaction.

In view of the above analysis, the aforementioned 468 experimental results on the decomposition of CH₃NH₃PbX₃ 469 (especially in regard to CH₂NH₃PbI₃) are somehow puzzling. 470 Apparently, measurements under Knudsen conditions, which 471 are supposed to favor the attainment of thermodynamic 472 equilibrium, gave evidence for the occurrence of the 473 thermodynamically disfavored process 8, whereas "open-pan" 474 TGA-MS and IR experiments provided evidence for the 475 occurrence of the thermodynamic pathway. Since the latter 476 experiments were mostly performed in a much higher 477 temperature range where kinetic hindrance might be over- 478 come, the question arises whether different temperatures alone 479 can account for the experimental findings or other effects are to 480 be invoked, such as vacuum versus dynamic inert atmosphere, 481 heating scan versus thermal equilibration, presence of a heated 482 transfer line, etc.

More recent KEMS experiments⁵⁰ where the competition 484 between processes 8 and 9 was studied by measuring the 485 $p(\mathrm{HI})/p(\mathrm{CH_3I})$ ratio under different conditions, provided 486 some additional information. The competition between the 487 two processes is basically driven by the following homoge- 488 neous pressure-independent gaseous equilibrium: 50 489

$$HX(g) + CH_3NH_2(g) = CH_3X(g) + NH_3(g)$$
 (12) ₄₉₀

Since the gaseous species come in 1:1 molar ratio from the 491 $CH_3NH_3PbX_3$ solid, at equilibrium $p(HX) = p(CH_3NH_2) \equiv 492$ p(8) and $p(CH_3X) = p(NH_3) \equiv p(9)$, with

$$\frac{p(8)}{p(9)} = \exp\left(\frac{\Delta_r G^{\circ}(12)}{2RT}\right) \tag{13}$$

In view of the small entropy change of reaction 12, the 496 partial pressure ratio 13 is ruled by the enthalpic factor. Since 497 ammonia is much more thermally stable than methylamine 498 ($\Delta_f H_{298}^\circ = -45.94$ and -22.5 kJ/mol, respectively), at the 499 temperatures of interest equilibrium 12 is shifted toward the 500 right, regardless of the nature of X. Furthermore, since HI is 501 thermally unstable ($\Delta_f H_{298}^\circ = +26.5$ kJ/mol) compared to 502 CH₃I ($\Delta_f H_{298}^\circ = 14.4$ kJ/mol), the decomposition channel 8 is 503 especially disfavored for iodide. The reverse holds for chloride 504 ($\Delta_f H_{298}^\circ = -92.3$ and -81.9 kJ/mol for HCl and CH₃Cl, 505 respectively), whereas the thermal stabilities of CH₃Br and 506 HBr are practically equal ($\Delta_f H_{298}^\circ = -36.3$ and -36.4 kJ/mol). 507 Being exothermal, reaction 12 tends to shift to left at higher 508 temperatures, although right-hand products remain strongly 509 favored at any temperature of interest. In Figure 4, the p(HI)/

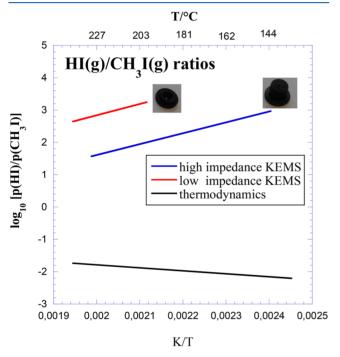


Figure 4. HI(g)/CH₃I(g) partial pressures ratios illustrating the competition between the decomposition processes 8 and 9. Red and blue lines refer to KEMS measurements performed, respectively, with low flow impedance (ordinary effusion cap) and high flow impedance (chimney-like cap).⁵⁰

510 p(CH₃I) pressure ratio calculated by eq 13 is reported, along 511 with the corresponding ratios measured by KEMS in a higher 512 temperature range compared to ref 27, using two different 513 effusion caps. 50 The red line refers to experiment carried out 514 with an ordinary effusion hole (1 mm diameter, negligible 515 thickness), the blue one was derived using a sort of "chimney" 516 orifice, with 0.5 mm in diameter and a 6.5 mm long channel 517 making the effusion rate lower (see the cap pictures in Figure 518 4). The higher impedance to effusion flow is expected to favor 519 approaching the heterogeneous equilibrium. Indeed, Figure 4 520 clearly shows that (i) higher temperatures do favor the 521 formation of CH₃I (g) versus HI(g), contrarily to thermody-522 namic predictions (black line with negative slope in Figure 4), 523 and (ii) high impedance effusion conditions have the same

effect, decreasing the HI/CH₃I ratio by more than 1 order of 524 magnitude. A schematic picture of these findings is reported in 525 Figure 5. Nevertheless, under all the explored conditions, 526 fS



Figure 5. Graphic representation of the experimental conditions leading to kinetically controlled (process 8) or thermodynamically controlled (process 9) decomposition pathways. Higher temperatures and/or closer-to-equilibrium conditions favor the thermodynamic control. ⁵⁰

process 8 remains the dominant decomposition pathway. 527 Therefore, according to KEMS results, even at temperature as 528 high as 250 °C, process 9 cannot come out by the extent 529 predicted by equilibrium thermodynamics.

In an attempt to rationalize the observed decomposition 531 behavior of $CH_3NH_3PbX_3$ perovskites, a possible benchmark is 532 given by the thermal decomposition of simple mono- di-, tri-, 533 and tetraalkylammonium halides, which have been the subject 534 of quite a high number of studies. 61–63 On the basis of TGA 535 experiments, Błażejowski and co-workers concluded that 536 compounds of general formula $R_pNH_{4-p}X$ decompose under 537 heating according to the following processes: 538

$$R_p NH_{4-p} X(s) = R_p NH_{3-p}(g) + HX(g)$$
 $p = 1,2,3$ (14) 539

$$R_p NH_{4-p} X(s) = R_{p-1} N(g) + RX(g) \quad p = 4$$
 (15) 540

In other words, mono-, di- and trialkyl ammonium halides 541 decompose, releasing the corresponding amine and hydrogen 542 halide rather than by the alternative process releasing the alkyl 543 halide and the less substituted amine:

$$R_p NH_{4-p} X(s) = R_{p-1} NH_{4-p}(g) + RX(g)$$
 (16) ₅₄₅

whereas tetralkylammonium halides, which do not contain 546 hydrogen atoms, decompose by loss of alkyl halide and 547 trialkylamine. The occurrence of the latter process was also 548 confirmed by direct FTIR spectroscopy experiments.⁶³ 549

As far as thermodynamic equilibrium is concerned, the \$50 above-reported discussion on the competition between \$51 processes 8 and 9 could be extended to processes 14 and 16. \$52 Actually, the pressure ratio (process 12) in equilibrium with \$53 the $CH_3NH_3PbX_3 + PbX_2$ mixture is the same as that in \$54 equilibrium with the methylammonium halide solids \$55 CH_3NH_3X , and it is ruled by the relative stability of $NH_3(g)$ \$56 versus $CH_3NH_2(g)$ and HX(g) versus $CH_3X(g)$, which would \$57 strongly favor process 16 (see above). Other cases can be more \$58 complex, since the relative stability of HX(g) and RX(g) \$59

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560 depends markedly on X and that of $R_pNH_{3-p}(g)$ versus 561 $R_{p-1}NH_{4-p}(g)$ for p=2,3 is not obvious.

The kinetic hindrance of the C-N bond-breaking process 563 involved in decomposition reactions 15 and 16 is supported by 564 the fact that quaternary amines, where process 14 cannot take 565 place, are observed to decompose at much higher temperatures 566 than less substituted compounds. 63 Furthermore, the enthalpy 567 changes of the thermal decomposition of tetralkylammonium 568 halides as measured by TGA have been found to be much 569 higher than those measured by DSC (for instance, 319.7 kJ/ 570 mol versus 186.7 kJ/mol for (CH₃)₄NI⁶³), which correspond 571 to the actual energy required to convert the solid into gaseous 572 products (i.e., the thermodynamic decomposition enthalpy). In 573 TGA measurements, where the enthalpy change is obtained by 574 the temperature dependence of the mass loss rate, the 575 activation barrier is instead measured. DFT calculations fully 576 support this view. For example, an activation energy of 239 kJ/ 577 mol was calculated recently 64 for the thermodynamically 578 favored decomposition of dioctylammonium chloride to 579 dioctylamine +1-chlorooctane, whereas no activation barrier 580 is found for the alternative process leading to trioctylammo-581 nium chloride and HCl. While the outlined frame of the 582 thermal behavior of alkylammonium halide seems fairly well-583 established, nonetheless the direct experimental evidence of 584 process 14 for mono-, di-, and trisubstituted compounds seems 585 scarce, and indeed recent perovskite-related experiments seem 586 to call it into question. 49,5

In conclusion, the intrinsic stability of CH₃NH₃PbX₃ 588 perovskites can be analyzed in the light of a classical 589 thermodynamic analysis relying on the limited experimental 590 information available to date. Decomposition reactions to solid 591 precursors PbX₂(s) and CH₃NH₃X(s) are shown to be 592 thermodynamically disfavored for X = Cl, Br, in great part 593 because of the large negative decomposition entropies. For X = 594 I, the serious discrepancy between the calorimetric determi-595 nations does not allow one to draw definitive conclusions on 596 the thermodynamic driving force of this decomposition 597 pathway, although the entropic contribution certainly also 598 plays a large stabilizing role in this case. Decomposition 599 reactions are most likely to occur by the release of gaseous 600 products, a process that, according to recent experimental 601 findings, may play an important role even in encapsulated 602 devices. However, the identification of the molecular species 603 lost by the perovskite structure is still uncertain, and the 604 limited experimental information is not conclusive. Decom-605 position to NH₃(g) and CH₃X(g) is largely favored from a 606 thermodynamic point of view, but it seems to suffer from a 607 severe kinetic limitation related to the breaking of the strong 608 C-N bond in the organic cation, as previously observed in the 609 thermal decomposition of alkylammonium halides. Indeed, 610 TGA-MS measurements support the occurrence of this 611 decomposition channel at high temperature. However, effusion 612 experiments indicate, in spite of thermodynamic driving forces, 613 the release of HX(g) and CH₃NH₂(g), rather than NH₃(g) 614 and CH₃X(g), at temperatures much lower than the 615 decomposition temperatures detected by TGA experiments. 616 The thermodynamic pathway becomes more important at 617 higher temperature and under closer-to-equilibrium effusion 618 conditions. The release of undissociated CH₃NH₃X(g) 619 molecules seems thermodynamically disfavored, although 620 accurate thermodynamic data for these species are lacking. 621 While thermodynamics prove to be very useful in rationalizing 622 the degradation behavior of perovskite materials and

corresponding precursors, careful attention has to paid to 623 kinetic effects. However, the scarcity and the uncertainty of the 624 experimental data currently available is a serious limit to 625 thermodynamic predictions, which would greatly benefit from 626 increased research efforts aimed at determining accurate 627 thermodynamic information by independent techniques 628 under various conditions. It is desirable that, in the next 629 years, chemical thermodynamics research will give a greater 630 contribution to assess the stability and the suitability of 631 perovskite materials for photovoltaics and, more generally, to 632 help the development of advanced materials for energy 633 applications. 65

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