# Convergence to fractional kinetics for random walks associated with unbounded conductances 

Martin T. Barlow • Jiří Černý

Received: 15 January 2009 / Revised: 17 November 2009 / Published online: 2 December 2009
© Springer-Verlag 2009


#### Abstract

We consider a random walk among unbounded random conductances whose distribution has infinite expectation and polynomial tail. We prove that the scaling limit of this process is a Fractional-Kinetics process-that is the time change of a $d$-dimensional Brownian motion by the inverse of an independent $\alpha$-stable subordinator. We further show that the same process appears in the scaling limit of the non-symmetric Bouchaud's trap model.


Mathematics Subject Classification (2000) $\quad 60 \mathrm{~F} 17 \cdot 60 \mathrm{~K} 37 \cdot 82 \mathrm{C} 41$

## 1 Introduction and main results

In this paper we establish a quenched non-Gaussian scaling limit theorem for two processes in a random environment in $\mathbb{Z}^{d}$ (with $d \geq 3$ ); these are the simple random walk among random conductances (often called the Random Conductance Model), and the non-symmetric Bouchaud's trap model on $\mathbb{Z}^{d}$. We will show that if the distribution of the environment is sufficiently heavy-tailed (and regular) then both these models, suitably normalised, converge to the Fractional Kinetics process, which is the non-Markovian, self-similar continuous process defined as the time change of a standard $d$-dimensional Brownian motion by the inverse of a stable subordinator.

[^0]Since the Fractional Kinetics process is sub-diffusive, that is its mean-square increment increases sub-linearly with time, our results prove that these models have anomalous diffusion. In [1] an analogous scaling limit theorem was established for the much simpler symmetric Bouchaud's trap model.

Before discussing our results we describe the models more precisely. We begin by defining continuous time random walks associated with a family of (non-random) conductances on $\mathbb{Z}^{d}$. Let $E^{d}$ be the set of all nearest-neighbour edges in $\mathbb{Z}^{d}$, and let $\mu_{e}, e \in E^{d}$, be strictly positive. Write $x \sim y$ if $x, y$ are neighbours in $\mathbb{Z}^{d}$. Set

$$
\begin{align*}
\mu_{x} & =\sum_{e \ni x} \mu_{e} \quad \text { for } x \in \mathbb{Z}^{d},  \tag{1.1}\\
p_{x y} & =\mu_{(x y)} / \mu_{x} \quad \text { if } x \sim y ; \tag{1.2}
\end{align*}
$$

here $(x y) \in E^{d}$ is the edge connecting $x$ and $y$. We consider random walks which jump from $x$ to a neighbour $y$ according to the transition probabilities $p_{x y}$. Let $v$ be a measure on $\mathbb{Z}^{d}$, and write $v_{x}=v(\{x\}), x \in \mathbb{Z}^{d}$. We study the continuous-time Markov chain on $\mathbb{Z}^{d}$ with transition rate from $x$ to $y$ given by $\mu_{x y} / v_{x}$. This random walk is reversible, and $v_{x}$ is its reversible measure.

There are two choices of $v$ which will concern us. If we take $v_{x}=\mu_{x}$, then the random walk has transition rates $p_{x y}$. We use $X=(X(t), t \geq 0)$ and $P_{x}^{\mu}$ to denote this Markov chain and its law on $D^{d}:=D\left([0, \infty), \mathbb{R}^{d}\right)$, and call this the constant-speed random walk (CSRW) in the configuration of conductances $\boldsymbol{\mu}=\left\{\mu_{e}: e \in E^{d}\right\}$. The term 'constant-speed' refers to the fact that the total jump rate out of any point $x$ is independent of $x$ and is equal to one.

The second walk, which we call the variable-speed random walk (VSRW) among the random conductances $\boldsymbol{\mu}$, is given by taking $v_{x} \equiv 1$. We use $Y=(Y(t), t \geq 0)$ and (with a slight abuse of notation) $P_{x}^{\mu}$ to denote this process on $D^{d}$. Since the CSRW and VSRW have the same jump probabilities, they are time-changes of each other, and this will play an important role in our proofs.

To obtain a random walk in random environment we take $\mu_{e}$ to be random. The first model we study is obtained by taking $\Omega=(0, \infty)^{E^{d}}$ to be the set of configurations of conductances, and let $\mathbb{P}$ be the product measure on $\Omega$ under which the canonical coordinates $\mu_{e}, e \in E^{d}$, are i.i.d. random variables. This gives us the Random Conductance Model (RCM).

We are interested in obtaining a $\mathbb{P}$-a.s. limit for the CSRW $X$ under the law $P_{0}^{\mu}$. In order to state our principal result we need to introduce the limiting Fractional-Kinetics (FK) process.

Definition 1.1 Let $\mathrm{BM}_{d}(t)$ be a standard $d$-dimensional Brownian motion started at 0 , and let $V_{\alpha}$ be an $\alpha$-stable subordinator independent of $\mathrm{BM}_{d}$, which is determined by $\mathbb{E}\left[e^{-\lambda V_{\alpha}(t)}\right]=e^{-t \lambda^{\alpha}}$. Let $V_{\alpha}^{-1}(s):=\inf \left\{t: V_{\alpha}(t)>s\right\}$ be the generalised rightcontinuous inverse of $V_{\alpha}(t)$. We define the fractional-kinetics process $\mathrm{FK}_{d, \alpha}$ by

$$
\begin{equation*}
\mathrm{FK}_{d, \alpha}(s)=\mathrm{BM}_{d}\left(V_{\alpha}^{-1}(s)\right), \quad s \in[0, \infty) \tag{1.3}
\end{equation*}
$$

The FK process is a non-Markovian process which is $\gamma$-Hölder continuous for all $\gamma<\alpha / 2$. It is self-similar: $\mathrm{FK}_{d, \alpha}(\cdot) \stackrel{\text { law }}{=} \lambda^{-\alpha / 2} \mathrm{FK}_{d, \alpha}(\lambda \cdot), \lambda>0$, and the density of its fixed time distribution $p(t, x)$ satisfies the fractional-kinetics equation

$$
\begin{equation*}
\frac{\partial^{\alpha}}{\partial t^{\alpha}} p(t, x)=\frac{1}{2} \Delta p(t, x)+\delta_{0}(x) \frac{t^{-\alpha}}{\Gamma(1-\alpha)} \tag{1.4}
\end{equation*}
$$

This process is well known in the physics literature. See the broad survey by Zaslavsky [2], the recent book [3] about the relevance of this process for chaotic deterministic systems, and also [4-8] for more on this class of processes and further references.

Our main result is the following quenched functional limit theorem for the CSRW in a heavy-tailed environment.

Theorem 1.2 Assume that $d \geq 3$,

$$
\begin{equation*}
\mathbb{P}\left[\mu_{e} \geq u\right]=C_{1} u^{-\alpha}(1+o(1)), \quad u \rightarrow \infty \tag{1.5}
\end{equation*}
$$

for some $\alpha \in(0,1), C_{1} \in(0, \infty)$, and that $\mathbb{P}\left[\mu_{e}>\underline{c}\right]=1$ for some $\underline{c} \in(0, \infty)$. Let

$$
\begin{equation*}
X_{n}(t)=n^{-1} X\left(t n^{2 / \alpha}\right), \quad t \in[0, \infty), n \in \mathbb{N} \tag{1.6}
\end{equation*}
$$

be the rescaled CSRW. Then there exists a constant $\mathcal{C} \in(0, \infty)$ such that $\mathbb{P}$-a.s., under $P_{0}^{\mu}$, the sequence of processes $X_{n}$ converges in law to a multiple of the fractionalkinetics process $\mathcal{C} \mathrm{FK}_{d, \alpha}$ on $D^{d}$ equipped with the topology of the uniform convergence on compact subsets of $[0, \infty)$.

Theorem 1.2 contrasts with a long line of Gaussian functional central limit theorems for both the CSRW, the VSRW and also for the discrete-time walk corresponding to the CSRW. Let

$$
\begin{equation*}
X_{n}^{\prime}(\cdot)=n^{-1} X\left(n^{2} \cdot\right), \quad Y_{n}(\cdot)=n^{-1} Y\left(n^{2} \cdot\right) \tag{1.7}
\end{equation*}
$$

In the 1980 s it was proved that, provided $\mathbb{E} \mu_{e}<\infty$, the processes $X_{n}^{\prime}$ and $Y_{n}$ converge to a multiple of a standard Brownian motion, $\sigma \mathrm{BM}_{d}$, in law under semi-direct product measure $\mathbb{P} \times P_{0}^{\mu}$, with the possibility that $\sigma=0$ in some cases-see [9-11].

These 'annealed' or averaged invariance principles were greatly improved in [12], where, under the ellipticity assumption $\mathbb{P}\left[\mu_{e} \in(a, b)\right]=1$ for some $0<a \leq b<\infty$, it was shown that for all $d \geq 1$, the rescaled discrete-time walk converges to $\sigma \mathrm{BM}_{d}$ with $\sigma>0$, almost surely with respect to $\mathbb{P}$. The ellipticity assumption was relaxed to boundedness from above (that is $\mathbb{P}\left[\mu_{e} \in[0, b]\right]=1$ ) in $[13,14]$, with the same non-trivial Brownian limit. (These papers impose the condition $\mathbb{P}\left(\mu_{e}=0\right)<p_{c}(d)$, where $p_{c}(d)$ is the critical probability for bond percolation in $\mathbb{Z}^{d}$, in the cases when $\mathbb{P}\left(\mu_{e}=0\right)>0$. Note also the papers [15-17] for the percolation case $\mu_{e} \in\{0,1\}$.) When the conductances are bounded from above, the CSRW, the VSRW, and the discrete-time walk with jump probabilities $p_{x y}$ are related by time changes that are asymptotically linear. Hence there is little difference in the asymptotic behaviour of these processes.

This is no longer the case for unbounded conductances. This situation was studied in [18] where it was proved that, under the assumption that $\mathbb{P}\left[\mu_{e}>a\right]=1$, both $X_{n}^{\prime}$ and $Y_{n}$ converge $\mathbb{P}$-a.s. to multiples of Brownian motion: $\sigma_{X} \mathrm{BM}_{d}$ and $\sigma_{Y} \mathrm{BM}_{d}$, respectively. While $\sigma_{Y}>0$ always, it was shown that the constant for the CSRW satisfies $\sigma_{X}>0$ if and only if $\mathbb{E}\left[\mu_{e}\right]<\infty$. The results of [18] therefore only give that $X_{n}^{\prime}$ converges to 0 ; Theorem 1.2 above identifies the right scaling and gives a non-trivial limiting process.

We now describe our second random environment, Bouchaud's trap model. This was introduced in the physics literature [19,20], on the complete graph, to explain some strange dynamical properties of complex disordered systems, in particular aging. The version on $\mathbb{Z}^{d}$ was first studied (using physics arguments) in [21].

It is defined as follows. Let $\tilde{\Omega}=\mathbb{R}^{\mathbb{Z}^{d}}$ and let $\tilde{\mathbb{P}}$ be a product measure on $\tilde{\Omega}$ under which the canonical coordinates $\mathcal{E}_{x}, x \in \mathbb{Z}^{d}$, interpreted as energies, are i.i.d. variables. We define $\tau_{x}=e^{\beta \mathcal{E}_{x}}$ to be the non-normalised Gibbs measure on $\mathbb{Z}^{d}$, and write $\boldsymbol{\tau}=\left(\tau_{x}: x \in \mathbb{Z}^{d}\right)$. Let $a \in[0,1]$ be a parameter, and define random conductances $\tilde{\mu}_{e}$ by

$$
\begin{equation*}
\tilde{\mu}_{(x y)}=\tau_{x}^{a} \tau_{y}^{a} \quad \text { if } x \sim y, \tag{1.8}
\end{equation*}
$$

Let $v_{x}=\tau_{x}$; then Bouchaud's trap model (BTM) is the continuous-time Markov chain on $\mathbb{Z}^{d}$ whose transition rates $w_{x y}$ are given by

$$
\begin{equation*}
w_{x y}=\frac{\tilde{\mu}_{(x y)}}{\tau_{x}}=\tau_{x}^{a-1} \tau_{y}^{a}=e^{-\beta\left((1-a) \mathcal{E}_{x}-a \mathcal{E}_{y}\right)}, \quad \text { for } x \sim y . \tag{1.9}
\end{equation*}
$$

We use $\tilde{X}=(\tilde{X}(t), t \geq 0)$ to denote the BTM, and we write $\tilde{P}_{x}^{\tau}$ for the law of $\tilde{X}$ on $D^{d}$. Note that for any $a \in[0,1]$ the Gibbs measure $\tau_{x}$ is reversible for the BTM. If $a=0$, then $\tilde{\mu}_{e}=1$ for all $e$, and the BTM is a time change of the simple random walk on $\mathbb{Z}^{d}$. This case is sometimes called symmetric BTM, while non-symmetric refers to the general case $a \neq 0$.

We can define the VSRW associated with the conductances $\tilde{\mu}_{e}$ in the same way as before; this process has jump rates given by (1.8) and counting measure as its reversible measure. We write $\tilde{Y}=(\tilde{Y}(t), t \geq 0)$ for this process, and $\tilde{P}_{x}^{\tau}$ for its law on $D^{d}$.

The following theorem, analogical to Theorem 1.2, is our main result on the BTM.
Theorem 1.3 Let $d \geq 3, a \in[0,1]$, and suppose that

$$
\begin{equation*}
\tilde{\mathbb{P}}\left[\tau_{x} \geq u\right]=C_{1} u^{-\alpha}(1+o(1)), \quad u \rightarrow \infty \tag{1.10}
\end{equation*}
$$

for some $\alpha \in(0,1), C_{1} \in(0, \infty)$, and $\tilde{\mathbb{P}}\left[\tau_{x}>\underline{c}\right]=1$ for some $\underline{c} \in(0, \infty)$. Let

$$
\begin{equation*}
\tilde{X}_{n}(t)=n^{-1} \tilde{X}\left(t n^{2 / \alpha}\right), \quad t \in[0, \infty) \tag{1.11}
\end{equation*}
$$

Then there exists a constant $\tilde{\mathcal{C}} \in(0, \infty)$ such that $\tilde{\mathbb{P}}$-a.s., under $\tilde{P}_{0}^{\tau}$, the sequence of processes $\tilde{X}_{n}$ converges in law to a multiple of the fractional-kinetics process $\tilde{\mathcal{C}} \mathrm{FK}_{d, \alpha}$ on $D^{d}$.

Let us recall previous rigorous results on the BTM. The first papers on this model concentrated on its aging behaviour. Scaling limit statements, if present, were used as technical tools. Moreover, with the exception of [22], only the symmetric BTM is explored in these papers. This include [23] considering the BTM on a $n$-dimensional hypercube, and [24] where aging of the BTM on $\mathbb{Z}$ is proved. In [24] it is also shown that the scaling limit (in the sense of finite-dimensional distributions) of the one-dimensional symmetric BTM is a singular diffusion in a random environment, and so quite different from the FK process. The result of [24] was extended to the non-symmetric case in [22]. It was proved there that the scaling limit is independent of $a$, giving an indication that this parameter has a little influence on the asymptotic behaviour of the BTM. The techniques of [24] and [22] used strongly the fact that $d=1$.

Aging for the symmetric BTM on $\mathbb{Z}^{d}$ was shown in [25] $(d=2)$ and [26] $(d \geq 3)$. The scaling limit approach to studying the BTM was put forward in [27]. In [1], a theorem analogous to our Theorem 1.3 was shown for $d \geq 2$, but in the symmetric case only. Theorem 1.3 confirms the fact that the influence of $a$ on the asymptotic behaviour of the BTM is small.

There is another natural Markov chain which has the Gibbs measure $\tau_{x}$ as the reversible measure, namely the Metropolis chain. Our techniques can be used to prove the same result as Theorem 1.3 for this chain too.

The case when $d \geq 3$ and $\alpha=1$ is treated (for the RCM) in [28]. In this case the process $n^{-1} X_{n^{2}(\log n) t}$ converges to a multiple of Brownian motion.

The assumptions of Theorems 1.2 and 1.3 are not optimal. The assumptions (1.5), (1.10) on the tail of $\mathbb{P}$ and $\tilde{\mathbb{P}}$ could be replaced by a weaker condition $\mathbb{P}\left[\mu_{e} \geq u\right]=$ $u^{-\alpha} L(u)$, where $L(u)$ is a function slowly varying at infinity. We decided to use (1.5) and (1.10) only to avoid unnecessary technical complications. It is clear that some tail regularity of $\mathbb{P}$ and $\tilde{\mathbb{P}}$ is necessary if we are to obtain convergence to the $\mathrm{FK}_{d, \alpha}$ process.

We also believe that both Theorems 1.2 and 1.3 hold in $d=2$, with an appropriate change of the scaling (cf. [1]). It should be possible to adapt the approach of the present paper to that case, but since the underlying random walks on $\mathbb{Z}^{2}$ are recurrent, slightly more care is needed. In particular, in this paper we use the fact that the random diagonal Green function $g_{B(x, R)}^{\mu}(x, x)$ of VSRW killed on the exit from a large ball $B(x, R)$ is essentially local function, that is for all $R \geq r=O\left(\log ^{c} n\right)$ and for all $x \in B(0, K n)$ it can be approximated by the Green function $g_{B(x, r)}^{\mu}(x, x)$ in the smaller ball $B(x, r)$ (cf. Lemmas 3.5 and 6.2). This implies that the diagonal Green functions at sufficiently distant points are essentially independent. In dimension $d=2$, a different control of the diagonal Green function in the balls is necessary.

It is also likely that the assumptions in the Theorems 1.2 and 1.3 that $\mu_{e}$ and $\tau_{x}$ are bounded from below, could be avoided. However, this would require combining and extending the arguments of [13,14] and [18] to cover the general RCM with positive conductances, and this is out of the scope of this paper.

The methods of our paper will likely apply when instead of the nearest-neighbour edges we consider edges of finite range. On the other hand, our methods are not strong enough if the independent collections $\boldsymbol{\mu}$ and $\boldsymbol{\tau}$ are replaced by collections with a finite range of dependence. In particular, Lemma 6.6 does not hold in the dependent situation.

The method that we use to prove both main theorems is based on the coarse-graining procedure developed in $[1,25]$ to control the symmetric BTM. In those papers, however, the proofs exploit heavily the fact that the symmetric BTM is the time change of the simple random walk. They use the local limit theorems and precise estimates on hitting probabilities and Green functions that are available for the simple random walk. In addition, they use the fact that all these asymptotic results hold uniformly in the starting position.

In this paper, however, the CSRW and the (non-symmetric) BTM are not time changes of the simple random walk. The process that corresponds to the simple random walk, that is the process with the flat reversible measure, is the VSRW. We do not have such precise estimates for the VSRW as for the simple random walk-and it is clear that any estimates that do hold cannot do so uniformly. In fact, for the RCM, we have available only the quenched FCLT and Gaussian heat kernel bounds for the VSRW proved in [18], and results (such as Harnack inequalities) that follow from them. Although these do yield a local limit theorem (see [18, Theorem 5.13]), this local limit theorem is obtained for at most a finite number of starting points simultaneously. (The ergodic theorems used to prove the FCLT for the VSRW do not give any information on the rate of convergence and its dependence on the starting position.) This proves to be a significant obstacle to using the coarse graining procedure of [1,25].

In order to overcome this difficulty, we had to improve considerably the original coarse graining. The coarse graining used in this paper requires as input only the FCLT for the process started at the origin and Gaussian heat-kernel estimates. It is therefore much more robust than the original one. These inputs are known for the VSRW for the RCM, and can be easily obtained for the VSRW $\tilde{Y}$ associated with the BTM by checking that the BTM satisfies the conditions of [18, Theorem 6.1].

We close the introduction by a short discussion related to the sub-diffusivity of the limiting FK process. Anomalous diffusions has been studied extensively in the physics literature since 1960s (see, e.g., [6] for an extensive list of references). It has been recognised that there are essentially two mechanisms leading to it. First, the presence of obstacles at a broad range of size scales can slow down the process. Mathematically this leads to studying processes on fractal sets that are nowadays quite well understood [29,30]. Second, the distribution of some local characteristic of the environment can be very broad and the process can be trapped at places where this characteristic is exceptionally large. To our knowledge, mathematical studies of this second mechanism are sparse and address to some extent trivial situations: either processes on the one-dimensional lattice $\mathbb{Z}$ [22,24], or processes that can be expressed as a time change of a simple random walk on $\mathbb{Z}^{d}, d \geq 1$, ( $[31,32]$ where the time change is independent of the trajectory and [1] where the time change depends on the environment). Our result is thus an important step in broadening our understanding of this second mechanism.

## 2 Time changes and other results

We have stated in Introduction that the CSRW and the VSRW are related by a time change. Since, as in [1], this time change will be an important subject of our study we
define it precisely here. We further state two auxiliary theorems, one for the VSRW and one for the BTM, that will be used to prove Theorems 1.2 and 1.3.

Consider the VSRW $Y$ and define the clock process $S:[0, \infty) \rightarrow[0, \infty)$ by

$$
\begin{equation*}
S(t)=\int_{0}^{t} \mu_{Y(t)} \mathrm{d} t \tag{2.1}
\end{equation*}
$$

Since $\mu_{x} \geq 2 d \underline{c}$ (see below (1.5) for definition of $\underline{c}$ ), $S$ is strictly increasing and its inverse $S^{-1}$ is well defined. It is easy to see that the process $X(t)=Y\left(S^{-1}(t)\right)$ has the same law as the CSRW. Hence, $Y$ and $X$ can be constructed on the same probability space and we always suppose that they are.

While the clock process $S$ is continuous, its scaling limit is not, as we will see soon. It follows that the clock process does not converge in the usual Skorokhod $J_{1}$-topology, and that the right topology to consider on $D^{1}$ is the weaker Skorokhod $M_{1}$-topology. To distinguish which topology we use, we let $D^{d}(U)$, resp. $D^{d}\left(M_{1}\right)$, stand for the space $D^{d}$ equipped with the topology of uniform convergence on compacts, resp. with the $M_{1}$-topology.

Theorem 1.2 is a consequence of the following more detailed result, which gives the joint convergence of the clock process and of the position of the embedded VSRW $Y$ to constant multiples of an independent stable subordinator $V_{\alpha}$ and Brownian motion $\mathrm{BM}_{d}$.

Theorem 2.1 Let, for $t \geq 0, n \in \mathbb{N}$,

$$
\begin{equation*}
Y_{n}(t)=n^{-1} Y\left(n^{2} t\right) \text { and } S_{n}(t)=n^{-2 / \alpha} S\left(n^{2} t\right) \tag{2.2}
\end{equation*}
$$

Under the assumptions of Theorem 1.2, there exist constants $\mathcal{C}_{Y}, \mathcal{C}_{S} \in(0, \infty)$ such that $\mathbb{P}$-a.s., under $P_{0}^{\mu}$, the joint distribution of $\left(S_{n}, Y_{n}\right)$ converges to the distribution of $\left(\mathcal{C}_{S} V_{\alpha}, \mathcal{C}_{Y} \mathrm{BM}_{d}\right)$ weakly on the space $D^{1}\left(M_{1}\right) \times D^{d}(U)$.

Theorem 1.3 is a consequence of an analogous statement. Recall that $\tilde{Y}$ is the VSRW associated with the conductances $\tilde{\mu}_{e}$. Defining the clock process $\tilde{S}$ analogously to (2.1),

$$
\begin{equation*}
\tilde{S}(t)=\int_{0}^{t} \tau_{\tilde{Y}(t)} \mathrm{d} t \tag{2.3}
\end{equation*}
$$

is easy to see that the BTM can be written as $\tilde{X}(t)=\tilde{Y}\left(\tilde{S}^{-1}(t)\right)$.
Theorem 2.2 Under the hypotheses of Theorem 1.3, the conclusions of Theorem 2.1 hold after the replacement of $S, Y, \mathbb{P}, P_{0}^{\mu}$ by $\tilde{S}, \tilde{Y}, \tilde{\mathbb{P}}$ and $\tilde{P}_{0}^{\tau}$.

The rest of the paper is organised as follows. We prove Theorems 1.2 and 2.1 for the CSRW in Sects. 3-8. In Sect. 3 we recall the results of [18] on the VSRW and prove some preliminary facts. In Sect. 4 we obtain some estimates on the Green function of the VSRW that replace the precise estimates for the simple random walk used in
[1,25]. In the next three sections we study the contribution to $S_{n}(t)$ of different sizes of conductances. In Sect. 5 we show that edges with conductance less than $\varepsilon n^{2 / \alpha}$ make little contribution. The main work is in Sect. 6, where we study edges with conductance between $\varepsilon n^{2 / \alpha}$ an $\varepsilon^{-1} n^{2 / \alpha}$. Here we introduce the improved coarse-graining procedure used to control the time that the CSRW spends in the vicinity of these edges. Section 7 treats the remaining corrections: edges with conductivity greater than $\varepsilon^{-1} n^{2 / \alpha}$, or closely spaced edges with conductivity greater than $\varepsilon n^{2 / \alpha}$. Finally in Sect. 8, Theorem 2.1 and then Theorem 1.2 are proved.

Theorems 1.3 and 2.2 for the BTM can be proved analogously to those for the RCM, and so we do not give a detailed proof. In Sect. 9 we prove that this process satisfies the conditions of [18, Theorem 6.1], so that Gaussian heat kernels and the FCLT hold for the VSRW $\tilde{Y}$ associated with the BTM. We then indicate the places where the proofs of Sects. 3-8 need to be modified, and give more details only at places where a different argument is necessary.

## 3 Preliminaries

We begin by introducing some further notation. Let $B(x, R)$ be the Euclidean ball centred at $x$ of radius $R$ and let $Q(x, R)$ be a cube centred at $x$ with side length $R$ whose edges are parallel to the coordinate axes. Both balls and cubes can be understood either as subsets of $\mathbb{R}^{d}, \mathbb{Z}^{d}$ or of $E^{d}$ (an edge is in $B(x, R)$ if both its vertices are), depending on the context. For $A \subset \mathbb{Z}^{d}$ we write $\partial A=\left\{y \notin A \exists x \in A,(x y) \in E^{d}\right\}$ and $\bar{A}=A \cup \partial A$. For $A, B \subset \mathbb{Z}^{d}$ we set $d(A, B)=\inf \{|x-y|: x \in A, y \in B\}$, where $|x-y|$ stands for the Euclidean distance of $x$ and $y$.

Sometimes, we identify a set of edges $G \subset E^{d}$ with the set of their vertices. It allows us to write, e.g., $x \in G$, meaning that $x \in \mathbb{Z}^{d}$ is a vertex of an edge in $G$. For a set $A \subset \mathbb{Z}^{d}$ we write $B(A, R)=\bigcup_{x \in A} B(x, R)$.

We use the convention that all large values appearing in the proofs are rounded above to the closest integer, if necessary. It allows us to write that, e.g., $\varepsilon n \mathbb{Z}^{d} \subset \mathbb{Z}^{d}$ for $\varepsilon \in(0,1)$ and $n$ large.

The following quenched FCLT for the VSRW in the RCM case was recently obtained in [18].

Proposition 3.1 (Theorem 1.1 of [18]) Let $Y_{n}$ be as in (2.2) and $d \geq 2$. Then there exists $\mathcal{C}_{Y} \in(0, \infty)$ such that $\mathbb{P}$-a.s., under $P_{0}^{\mu}$, the sequence $Y_{n}$ converges in law on $D^{d}(U)$ to a multiple of a standard d-dimensional Brownian motion, $\mathcal{C}_{Y} \mathrm{BM}_{d}$.

Let $Z$ be a process or path in $\mathbb{R}^{d}$. For $A \subset \mathbb{R}^{d}$ we define its hitting and exit times $\sigma_{A}(Z)=\inf \{t \geq 0: Z(t) \in A\}, \tau_{A}(Z)=\inf \{t \geq 0: Z(t) \notin A\}$. When $Z$ is the VSRW $Y$ we will usually write $\sigma_{A}$ and $\tau_{A}$ for $\sigma_{A}(Y)$ and $\tau_{A}(Y)$. We write $g_{A}^{\mu}(x, y)$ for the Green function of $Y$ killed on exiting from $A$ :

$$
\begin{equation*}
g_{A}^{\mu}(x, y)=E_{x}^{\mu} \int_{0}^{\tau_{A}} \mathbf{1}\left\{Y_{s}=y\right\} \mathrm{d} s \tag{3.1}
\end{equation*}
$$

If $A=\mathbb{Z}^{d}$, we omit it from the notation, and we sometimes also omit $\mu$ and write $g(x, y)=g_{\mathbb{Z}^{d}}^{\mu}(x, y)$.

Recall the definition of $p_{x y}$ in (1.2). We say that a function $h$ is harmonic on $A \subset \mathbb{Z}^{d}$ if $h$ is defined on $\bar{A}$ and

$$
\begin{equation*}
h(x)=\sum_{y \sim x} p_{x y} h(y) \text { for all } x \in A . \tag{3.2}
\end{equation*}
$$

We say that the elliptic Harnack inequality (EHI) with constant $C_{H}$ holds in a ball $B(x, r)$ if whenever $h \geq 0$ is harmonic in $B(x, r)$, then

$$
\begin{equation*}
h(y) \leq C_{H} h\left(y^{\prime}\right) \quad \text { for all } y, y^{\prime} \in B(x, r / 2) \tag{3.3}
\end{equation*}
$$

From [18] we obtain the following-see Theorem 1.3 for the bounds on $g^{\mu}$ and Theorem 4.7 for the Harnack inequality. (Note that the EHI follows immediately from the parabolic Harnack inequality.)

Proposition 3.2 Let $d \geq 3$. There exists a collection of random variables $\left(V_{x}, x \in\right.$ $\mathbb{Z}^{d}$ ) on $\Omega$ and positive constants $c_{1}, c_{2}, C_{H}, \eta=1 / 3$ with the following properties.
(a) For all $x \in \mathbb{Z}^{d}$,

$$
\begin{equation*}
\mathbb{P}\left[V_{x} \geq n\right] \leq c_{2} \exp \left(-c_{1} n^{\eta}\right) \tag{3.4}
\end{equation*}
$$

(b) If $x, y \in \mathbb{Z}^{d}$,

$$
\begin{equation*}
\frac{c_{1}}{|x-y|^{d-2}} \leq g^{\mu}(x, y) \leq \frac{c_{2}}{|x-y|^{d-2}} \quad \text { if }|x-y| \geq V_{x} \wedge V_{y} \tag{3.5}
\end{equation*}
$$

(c) If $R \geq V_{x}$, then $E H I$ holds with constant $C_{H}$ for $B(x, R)$.
(d) Let $C_{0}=\Gamma\left(\frac{d}{2}-1\right) /\left(2 \pi^{d / 2} \mathcal{C}_{Y}^{2}\right)$. For any $\varepsilon>0$ there exists a random variable $M_{\varepsilon}$ on $\Omega$ with $\mathbb{P}\left(M_{\varepsilon}<\infty\right)=1$ such that

$$
\begin{equation*}
\frac{(1-\varepsilon) C_{0}}{|x|^{d-2}} \leq g^{\mu}(0, x) \leq \frac{(1+\varepsilon) C_{0}}{|x|^{d-2}} \text { for }|x|>M_{\varepsilon} \tag{3.6}
\end{equation*}
$$

For all $f \in \ell^{2}\left(\mathbb{Z}^{d}, \mu_{x}\right)$ we set

$$
\begin{equation*}
\mathcal{E}^{\mu}(f)=\frac{1}{2} \sum_{x, y \in \mathbb{Z}^{d}} \mu_{x y}(f(y)-f(x))^{2}, \tag{3.7}
\end{equation*}
$$

and we define the conductance between two disjoint subsets $A, B$ of $\mathbb{Z}^{d}$ as

$$
\begin{equation*}
C_{\mathrm{eff}}[A, B]=\inf \left\{\mathcal{E}^{\mu}(f):\left.f\right|_{A}=0,\left.f\right|_{B}=1\right\} \tag{3.8}
\end{equation*}
$$

We use $C_{\text {eff }}[A, \infty]$ to denote $\lim _{n \rightarrow \infty} C_{\text {eff }}\left[A, B(A, n)^{c}\right]$, and $R_{\text {eff }}[A, B]=$ $C_{\text {eff }}[A, B]^{-1}$ to denote the resistance between $A$ and $B$. We recall the well-known fact

$$
\begin{equation*}
g_{A}(x, y) \leq g_{A}(x, x)=R_{\mathrm{eff}}\left[x, A^{c}\right], \quad \forall x, y \in A \tag{3.9}
\end{equation*}
$$

We now prove several auxiliary lemmas. Let $K \geq 1$ be a fixed large number.
Lemma 3.3 There exist $c_{3}>0$ and $N_{0}=N_{0}(\boldsymbol{\mu})$ with $\mathbb{P}\left(N_{0}<\infty\right)=1$ such that on $\left\{\boldsymbol{\mu}: n \geq N_{0}(\boldsymbol{\mu})\right\}$,

$$
\begin{equation*}
\sup \left\{V_{x}: x \in B(0, K n)\right\} \leq c_{3}(\log n)^{1 / \eta}=: b_{1}(n) . \tag{3.10}
\end{equation*}
$$

Proof Using (3.4), the probability that (3.10) fails is at most $C(K n)^{d} e^{\left(-c_{1} c_{3}^{\eta} \log n\right)}$, which is summable if $c_{1} c_{3}^{\eta}$ is large enough.

Lemma 3.4 (a) There exists a constant $c_{4}$ such that the Green function satisfies

$$
\begin{equation*}
g^{\mu}(x, y) \leq c_{4} \text { for all } x, y \in \mathbb{Z}^{d} \tag{3.11}
\end{equation*}
$$

(b) Further, for any $x \in \mathbb{Z}^{d}, g^{\mu}(x, x) \geq c_{1} V_{x}^{2-d}$. In particular, $\mathbb{P}$-a.s. for all but finitely many $n \in \mathbb{N}$,

$$
\begin{equation*}
\inf _{x \in B(0, K n)} g^{\mu}(x, x) \geq c_{1} b_{1}(n)^{2-d}=c_{1} c_{3}^{2-d}(\log n)^{(2-d) / \eta} . \tag{3.12}
\end{equation*}
$$

Proof (a) is immediate from (3.9) and the comparison of the random conductances network $\boldsymbol{\mu}$ with a homogeneous network where every conductance equals $\underline{c}$ (see (1.5)). The first claim in (b) follows from (3.5), and the fact that $g^{\mu}(x, \cdot)$ attains its maximum at $x$. Using (3.10) then gives (3.12).

We need estimates on resistance to boundaries of large balls.
Lemma 3.5 There exists a constant $c_{5}$ such that for every $\varepsilon \in(0,1), \mathbb{P}$-a.s. for all but finitely many $n \in \mathbb{N}$,

$$
\begin{equation*}
R_{\mathrm{eff}}[x, \infty] \geq R_{\mathrm{eff}}\left[x, B\left(x, b_{2}(n, \varepsilon)\right)^{c}\right] \geq(1-\varepsilon) R_{\mathrm{eff}}[x, \infty] \quad \forall x \in B(0, K n) \tag{3.13}
\end{equation*}
$$

where $b_{2}(n, \varepsilon)=c_{5} \varepsilon^{\frac{1}{2-d}} b_{1}(n)$.
Proof The first inequality in (3.13) is obvious. To prove the second, set $r=b_{2}(n, \varepsilon)$. Observe that the function $f_{x}(y):=g(x, y)-g_{B(x, r)}(x, y)$ is non-negative and harmonic on $B(x, r)$. By (3.5) and Lemma 3.3, $f_{x}(y) \leq c_{2} r^{2-d}$ for $x \in B(0, K n)$ and $y \in \partial B(x, r), \mathbb{P}$-a.s. for all $n$ large. Hence, by the maximum principle, $f_{x}(y) \leq c_{2} r^{2-d}$ on $B(x, r)$. On the other hand, $\mathbb{P}$-a.s., $g(x, x) \geq c(\log n)^{(2-d) / \eta}$. Combining these two bounds, (3.9), and taking $c_{5}$ large enough, the second inequality follows easily.

We now fix the constant $C_{H}$ in EHI to be as in Proposition 3.2(c). Let $N_{0}$ be as in Lemma 3.3. We will require the Harnack inequality to hold for many balls simultaneously.

Lemma 3.6 (a) If $n \geq N_{0}, x \in B(0, K n)$ and $r \geq b_{1}(n)$ then EHI holds for $B(x, r)$. (b) Let $n \geq N_{0}$. If $z \in B(0, K n), r \geq b_{1}(n), m \geq 1$, and $h \geq 0$ is harmonic on $B\left(z, 2^{m} r\right)$, then, writing $\rho=\left(2 C_{H}\right)^{-1}$,

$$
\begin{equation*}
\frac{\sup _{B(z, r)} h}{\inf _{B(z, r)} h} \leq 1+(1-\rho)^{m-1} C_{H} \tag{3.14}
\end{equation*}
$$

Proof (a) This is immediate from Proposition 3.2(b) and Lemma 3.3.
(b) We consider the chain of balls $B_{i}=B\left(z, 2^{i} r\right), i=0, \ldots, m$. By (a) EHI holds for each $B_{i}$. Write $\operatorname{Osc}(f, A)=\sup _{A} f-\inf _{A} f$. A standard argument (see e.g. [33, Proposition 3.2]) gives that

$$
\begin{equation*}
\operatorname{Osc}\left(h, B_{i-1}\right) \leq(1-\rho) \operatorname{Osc}\left(h, B_{i}\right), \quad 1 \leq i \leq m \tag{3.15}
\end{equation*}
$$

Further, using EHI in $B_{m}$,

$$
\begin{equation*}
\operatorname{Osc}\left(h, B_{m-1}\right) \leq \sup _{B_{m-1}} h \leq C_{H} \inf _{B_{m-1}} h \leq C_{H} \inf _{B_{0}} h . \tag{3.16}
\end{equation*}
$$

So,

$$
\begin{align*}
\sup _{B_{0}} h-\inf _{B_{0}} h & =\operatorname{Osc}\left(h, B_{0}\right) \leq(1-\rho)^{m-1} \operatorname{Osc}\left(h, B_{m-1}\right) \\
& \leq C_{H}(1-\rho)^{m-1} \inf _{B_{0}} h . \tag{3.17}
\end{align*}
$$

This finishes the proof.

## 4 Behaviour of Green functions

In this section we prove the estimates on Green functions that we use to replace the standard asymptotic formulae on the Green function of the simple random walk used in [1,25]. Before we present our results, the following approximate calculations on a simplified model may prove useful.

Let $d \geq 3$, and consider $\mathbb{Z}^{d}$ with $\mu_{e} \equiv 1$, except in a ball $B(x, r)$, where $\mu_{e} \equiv$ $A \gg 1$. Write $\tilde{g}(y, z)$ for the Greens function for this graph. If $|x-y|>2 r$ then $\tilde{g}(y, y)=R_{\text {eff }}(y, \infty)=O(1)$, while $\tilde{g}(x, x) \simeq R_{\text {eff }}\left(B(x, r)^{c}, \infty\right) \simeq c r^{2-d}$. Since $|x-y|>2 r$ the anomalous region around $x$ has little effect on the hitting probability of $y$ starting from $z \in B(x, 3 r / 2)^{c}$, and so $P_{x}\left(\sigma_{y}<\infty\right) \simeq|x-y|^{2-d}$. Thus we have

$$
\begin{equation*}
\tilde{g}(x, y)=P_{x}\left(\sigma_{y}<\infty\right) \tilde{g}(y, y) \simeq c|x-y|^{2-d}, \quad \text { if }|x-y|>2 r . \tag{4.1}
\end{equation*}
$$

Thus the anomalous region in $B(x, r)$ has little effect on $\tilde{g}(x, \cdot)$ outside $B(x, 2 r)$.

Proposition 4.1 Let $0<2 \varepsilon_{o}<\varepsilon_{g} \leq \frac{1}{2}$, and $\delta>0$. Then there exists $\varepsilon_{b} \in\left(0, \varepsilon_{o}\right)$ with $\lim _{\delta \rightarrow 0} \varepsilon_{b}=0$ such that if $n \geq N_{0}, \varepsilon_{b} n \geq b_{1}(n), x \in B\left(0,\left(K-\varepsilon_{g}\right) n\right)$ and $A \subset \mathbb{Z}^{d}$ with $B\left(x, \varepsilon_{g} n\right) \subset A$, then

$$
\begin{array}{cl}
\frac{\sup _{y \in B\left(z, \varepsilon_{b} n\right)} g_{A}^{\mu}(x, y)}{\inf _{y \in B\left(z, \varepsilon_{b} n\right)} g_{A}^{\mu}(x, y)} \leq 1+\delta \quad \text { whenever } \varepsilon_{o} n \leq|x-z| \leq\left(\varepsilon_{g}-\varepsilon_{o}\right) n, \\
\frac{\sup _{x^{\prime} \in B\left(x, \varepsilon_{b} n\right)} g_{A}^{\mu}\left(x^{\prime}, y\right)}{\inf _{x^{\prime} \in B\left(x, \varepsilon_{b} n\right)} g_{A}^{\mu}\left(x^{\prime}, y\right)} \leq 1+\delta \quad \text { whenever } \varepsilon_{o} n<|y-x|<\varepsilon_{g} n . \tag{4.3}
\end{array}
$$

Proof Choose $m$ so that $C_{H}(1-\rho)^{m-1}<\delta$, and write $h=g_{A}^{\mu}(x, \cdot)$. Fix $\varepsilon_{b} \leq 2^{-m} \varepsilon_{o}$. If $z$ is as in (4.2), then $z \in B(0, K n)$ and $B\left(z, 2^{m} \varepsilon_{b} n\right) \subset B\left(z, \varepsilon_{o} n\right) \subset A \backslash\{x\}$, and therefore $h$ is harmonic on $B\left(z, 2^{m} \varepsilon_{b} n\right)$. Hence, for $n$ satisfying the assumptions of the proposition, using Lemma 3.6,

$$
\begin{equation*}
\frac{\sup _{B\left(z, \varepsilon_{b} n\right)} h}{\inf _{B\left(z, \varepsilon_{b} n\right)} h} \leq 1+(1-\rho)^{m-1} C_{H} \leq 1+\delta \tag{4.4}
\end{equation*}
$$

This proves (4.2). The proof of (4.3) is similar.
Write $g_{A}^{*}(x, y)$ for the Green function of $W:=\mathcal{C}_{Y} \mathrm{BM}_{d}$ killed on exit from $A \subset \mathbb{R}^{d}$ (see Proposition 3.1 for $\mathcal{C}_{Y}$ ). We have

$$
\begin{equation*}
g^{*}(x, y)=C_{0}|x-y|^{2-d}, \quad g_{B(x, r)}^{*}(x, y)=C_{0}\left(|x-y|^{2-d}-r^{2-d}\right) \tag{4.5}
\end{equation*}
$$

where the constant $C_{0}$ is as in Proposition 3.2. Using the FCLT we now show that an estimate similar to the second part of (4.5) holds for $g_{B(n x, n r)}^{\mu}(n x, n y)$, provided that $x, y \in \mathbb{R}^{d}$ and $r$ are fixed, and $y$ is not too close either to $x$ or to $B(x, r)^{c}$.

Lemma 4.2 Let $0<3 \varepsilon_{o}<\varepsilon_{g}<\frac{1}{2}, \delta>0$. Let $r \in\left(\varepsilon_{g}-\varepsilon_{o}, \varepsilon_{g}+\varepsilon_{o}\right)$ and let $x, y \in \mathbb{R}^{d}, x \in B\left(0, K-\varepsilon_{g}\right), y \in B\left(x, r-\varepsilon_{o}\right) \backslash B\left(x, \varepsilon_{o}\right)$. Then there exists $c_{0}(\delta)$ with $\lim _{\delta \rightarrow 0} c_{0}(\delta)=0$ such that if

$$
\begin{equation*}
A_{n}^{r}(x, y)=\left\{\mu: 1-c_{0}(\delta) \leq \frac{g_{B(n x, r n)}^{\mu}(n x, n y)}{C_{0} n^{2-d}\left(|x-y|^{2-d}-r^{2-d}\right)} \leq 1+c_{0}(\delta)\right\} \tag{4.6}
\end{equation*}
$$

then $A_{n}^{r}(x, y)$ holds for all sufficiently large $n, \mathbb{P}$-a.s.
Proof We use the notation $c_{i}(\delta)$ to denote functions of $\delta$ such that $\lim _{\delta \rightarrow 0} c_{i}(\delta)=0$. Let $\varepsilon_{b}<\varepsilon_{o}$ be as in Proposition 4.1. For any path $\gamma \in D^{d}$ write $\sigma_{1}(\gamma)=\sigma_{B\left(x, \varepsilon_{b}\right)}(\gamma)$, and $\tau_{1}(\gamma)=\inf \left\{t \geq \sigma_{1}: \gamma \notin B\right\}$ where $B=B(x, r)$, and let

$$
\begin{equation*}
F(\gamma)=\int_{\sigma_{1}(\gamma)}^{\tau_{1}(\gamma)} \mathbf{1}\left\{\gamma(s) \in B\left(y, \varepsilon_{b}\right)\right\} \mathrm{d} s \tag{4.7}
\end{equation*}
$$

By the FCLT (Proposition 3.1) for $Y$, we have,

$$
\begin{equation*}
\lim _{n \rightarrow \infty} E_{0}^{\mu} F\left(Y_{n}\right)=E_{0} F(W) \tag{4.8}
\end{equation*}
$$

The right hand side of (4.8) is

$$
\begin{equation*}
E_{0} \mathbf{1}\left\{\sigma_{1}(W)<\infty\right\} \int_{B\left(y, \varepsilon_{b}\right)} g_{B}^{*}\left(W_{\sigma_{1}}, z\right) \mathrm{d} z . \tag{4.9}
\end{equation*}
$$

Using the uniform results for $g_{B}^{*}$ analogous to Proposition 4.1, and writing $v_{d}=$ $|B(0,1)|$ for the volume of the ball in $\mathbb{R}^{d}$, it follows that

$$
\begin{equation*}
E_{0} F(W) \geq\left(1-c_{1}(\delta)\right) P_{0}\left[\sigma_{1}(W)<\infty\right] v_{d} \varepsilon_{b}^{d} g_{B}^{*}(x, y) \tag{4.10}
\end{equation*}
$$

Similarly, we have, writing $\sigma_{n}=\sigma_{B\left(n x, n \varepsilon_{b}\right)}(Y), n B=B(n x, n r)$,

$$
\begin{equation*}
E_{0}^{\mu} F\left(Y_{n}\right)=E_{0}^{\mu}\left[n^{-2} \mathbf{1}\left\{\sigma_{n}<\infty\right\} \sum_{z \in B\left(n y, n \varepsilon_{b}\right)} g_{n B}^{\mu}\left(Y_{\sigma_{n}}, z\right)\right] \tag{4.11}
\end{equation*}
$$

where the additional factor $n^{-2}$ comes from the scaling of $Y$ to $Y_{n}$. Using Proposition 4.1 (with $\varepsilon_{o} / 2$ on place of $\varepsilon_{o}$ and $\varepsilon_{g}-\varepsilon_{o}$ on place of $\varepsilon_{g}$ ) we can choose $\varepsilon_{b}$ small such that, for $n$ sufficiently large and $z \in B\left(n y, n \varepsilon_{b}\right)$

$$
\begin{equation*}
g_{n B}^{\mu}\left(Y_{\sigma_{n}}, z\right) \leq(1+\delta) g_{n B}^{\mu}(n x, z) \leq(1+\delta)^{2} g_{n B}^{\mu}(n x, n y), \tag{4.12}
\end{equation*}
$$

with a similar lower bound. Therefore,

$$
\begin{equation*}
E_{0}^{\mu} F\left(Y_{n}\right) \leq\left(1+c_{2}(\delta)\right) P_{0}^{\mu}\left[\sigma_{n}<\infty\right] n^{d-2} v_{d} \varepsilon_{b}^{d} g_{n B}^{\mu}(n x, n y) . \tag{4.13}
\end{equation*}
$$

The functional central limit theorem also implies that

$$
\begin{equation*}
P_{0}^{\mu}\left[\sigma_{n}<\infty\right] \xrightarrow{n \rightarrow \infty} P_{0}\left[\sigma_{1}(W)<\infty\right] . \tag{4.14}
\end{equation*}
$$

Combining this with (4.8), (4.10) and (4.13), we obtain that

$$
\begin{equation*}
n^{d-2} g_{n B}^{\mu}(n x, n y) \geq\left(1-c_{0}(\delta)\right) g_{B}^{*}(x, y) \tag{4.15}
\end{equation*}
$$

provided $n$ is sufficiently large. The upper bound on $g_{n B}^{\mu}(n x, n y)$ is proved in the same way. From these bounds and (4.5) we deduce that $A_{n}^{r}$ holds for all large $n$.

Finally, we show that the last lemma holds uniformly over $B(0, K n)$.

Proposition 4.3 Let $0<3 \varepsilon_{o}<\varepsilon_{g}<\frac{1}{2}$, and $\delta>0$. Let $r \in\left(\varepsilon_{g}, \varepsilon_{g}+\varepsilon_{o} / 2\right)$. Then there exists $c_{1}(\delta)$ with $\lim _{\delta \rightarrow 0} c_{1}(\delta)=0$ such that, for all but finitely many $n, \mathbb{P}$-a.s., for all $x \in B\left(0,\left(K-\varepsilon_{g}\right) n\right), y \in B\left(x,\left(\varepsilon_{g}-\varepsilon_{o}\right) n\right) \backslash B\left(x, \varepsilon_{o} n\right)$,

$$
\begin{equation*}
1-c_{1}(\delta) \leq \frac{g_{B(x, r n)}^{\mu}(x, y)}{C_{0}\left(|x-y|^{2-d}-(n r)^{2-d}\right)} \leq 1+c_{1}(\delta) \tag{4.16}
\end{equation*}
$$

Proof Let $\varepsilon_{b}$ be as in the proof Lemma 4.2. Let $\left\{x_{1}, \ldots, x_{m}\right\} \subset B\left(0, K-\varepsilon_{g}\right)$ be a finite collection of points chosen so that $B\left(0, K-\varepsilon_{g}\right) \subset \cup_{i} B\left(x_{i}, \varepsilon_{b}\right)$. Write $J=\{(i, j)$ : $\left.r-\varepsilon_{o} / 2 \geq\left|x_{i}-x_{j}\right| \geq \varepsilon_{o} / 2\right\}$. Lemma 4.2 implies that there exists $N^{\prime}=N^{\prime}(\boldsymbol{\mu}) \geq N_{0}$ such that $A_{n}^{r+\varepsilon_{b}}\left(x_{i}, x_{j}\right)$ holds for all $(i, j) \in J$ and $n \geq N^{\prime}$.

Now let $n \geq N^{\prime}$ and let $x, y$ be as required. Then there exists $(i, j) \in J$ such that $\left|x-n x_{i}\right|<n \varepsilon_{b},\left|y-n x_{j}\right|<n \varepsilon_{b}$. By Proposition 4.1, since $A_{n}^{r+\varepsilon_{b}}\left(x_{i}, x_{j}\right)$ holds,

$$
\begin{align*}
g_{B(x, n r)}^{\mu}(x, y) & \leq(1+\delta)^{2} g_{B(x, n r)}^{\mu}\left(n x_{i}, n x_{j}\right) \\
& \leq(1+\delta)^{2} g_{B\left(n x_{i}, n\left(r+\varepsilon_{b}\right)\right)}^{\mu}\left(n x_{i}, n x_{j}\right) \\
& \leq(1+\delta)^{2} C_{0} n^{2-d}\left(\left|x_{i}-x_{j}\right|^{2-d}-\left(r+\varepsilon_{b}\right)^{2-d}\right) \\
& \leq(1+c(\delta)) C_{0}\left(|x-y|^{2-d}-(n r)^{2-d}\right) . \tag{4.17}
\end{align*}
$$

The lower bound in (4.16) is proved in the same way.

## 5 Edges with small conductance

The next three sections extend the methods of $[1,25]$ to the CSRW. We begin with one simplification, and replace the hypothesis (1.5) by

$$
\begin{equation*}
\mathbb{P}\left[\mu_{e} \geq u\right]=u^{-\alpha}(1+o(1)), \quad u \rightarrow \infty \tag{5.1}
\end{equation*}
$$

that is, we assume $C_{1}=1$ in (1.5). An easy rescaling argument recovers the general case.

We aim at proving that the rescaled clock process $S_{n}$ (see Theorem 2.1) converges to a stable subordinator. Since the stable subordinator at time $T$ is well approximated by a large but finite number of its largest jumps before $T$, we will, in the next section, control the (suitably coarse-grained) large jumps of $S_{n}$ and prove that their distribution converge to the distribution of the large jumps of the subordinator. As we will see these large jumps are due to the visits of the VSRW to edges with conductance of order $n^{2 / \alpha}$.

Before studying the large jumps we prove in this section that the contribution of less conducting edges to the clock process can be neglected. More precisely, we show that the contribution of the edges with conductances smaller than $\varepsilon n^{2 / \alpha}$ to the rescaled clock process $S_{n}$ at time $T$ a is very likely to be smaller than $\delta T$, for suitably chosen $\varepsilon, \delta$. This is the content of the next proposition.

## Proposition 5.1 Let

$$
\begin{equation*}
\tilde{T}_{n}(0, \varepsilon)=\left\{x \in \mathbb{Z}^{d}: \mu_{x y} \leq \varepsilon n^{2 / \alpha} \forall y \sim x\right\} . \tag{5.2}
\end{equation*}
$$

Then, for every $\delta>0$ there exists $\varepsilon$ such that for all $T>0, \mathbb{P}$-a.s. for all but finitely many $n$,

$$
\begin{equation*}
P_{0}^{\mu}\left[T^{-1} n^{-2 / \alpha} \int_{0}^{T n^{2}} \mu_{Y(t)} \mathbf{1}\left\{Y(t) \in \tilde{T}_{n}(0, \varepsilon)\right\} \mathrm{d} t \geq \delta\right] \leq \delta . \tag{5.3}
\end{equation*}
$$

Proof Let $\mathbb{B}=B(0, K n)$, and write $\overline{\mathbb{B}}$ for the set of edges with at least one vertex in $\mathbb{B}$. We first insert the trivial term $\mathbf{1}\{Y(t) \in \mathbb{B}\}+\mathbf{1}\{Y(t) \notin \mathbb{B}\}$ in the integral in (5.3), and write the resulting integral as $I_{\mathbb{B}}(T)+I_{\mathbb{B}}(T)$. By the FCLT (Proposition 3.1), for each $T>0$ it is possible to choose $K=K(T)$ large such that, $\mathbb{P}$-a.s., $P_{0}^{\mu}\left[I_{\mathbb{B}^{c}}(T)>0\right]<\delta / 2$.

To bound the contribution of the term $I_{\mathbb{B}}(T)$ we show that for every $\delta$ there exists $\varepsilon$ such that, $\mathbb{P}$-a.s. for large $n$,

$$
\begin{equation*}
E_{0}^{\mu}\left[\int_{0}^{T n^{2}} \mu_{Y(t)} \mathbf{1}\left\{Y(t) \in \tilde{T}_{n}(0, \varepsilon) \cap \mathbb{B}\right\} \mathrm{d} t\right] \leq \frac{1}{2} T n^{2 / \alpha} \delta^{2} . \tag{5.4}
\end{equation*}
$$

The claim (5.3) then follows using the Chebyshev inequality.
To show (5.4), we set $i_{\text {max }}=\min \left\{i \in \mathbb{N}: 2^{-i} \varepsilon n^{2 / \alpha} \leq \underline{c}\right\}=O(\log n)$, and

$$
\begin{equation*}
H_{n}(i)=\left\{e \in \overline{\mathbb{B}}: \mu_{e} \in \varepsilon n^{2 / \alpha}\left(2^{-i}, 2^{-i+1}\right]\right\}, \quad i \in\left\{1, \ldots, i_{\max }\right\} . \tag{5.5}
\end{equation*}
$$

Using Proposition 3.2, Lemmas 3.3, 3.4, and the notation $g^{\mu}(0,(x y))=g^{\mu}(0, x)+$ $g^{\mu}(0, y)$, the left-hand side of (5.4) can bounded from above by

$$
\begin{equation*}
\sum_{(x y) \in E^{d} \cap \overline{\mathbb{B}}} \mu_{(x y)} g^{\mu}(0,(x y)) \mathbf{1}\left\{\mu_{(x y)} \leq \varepsilon n^{2 / \alpha}\right\} \leq \sum_{i=1}^{i_{\max }} 2^{-i+1} \varepsilon n^{2 / \alpha} \sum_{e \in H_{n}(i)} \bar{g}_{n}(0, e), \tag{5.6}
\end{equation*}
$$

where

$$
\bar{g}_{n}(0, e)= \begin{cases}c_{4} & \text { if } d(0, e) \leq b_{1}(n)  \tag{5.7}\\ c_{2} d(0, e)^{2-d} & \text { otherwise }\end{cases}
$$

Let $p_{n, i}=\mathbb{P}\left[\mu_{e} \in \varepsilon n^{2 / \alpha}\left(2^{-i}, 2^{-i+1}\right]\right] \leq c \varepsilon^{-\alpha} 2^{i \alpha} n^{-2}, \lambda_{n}>0$, and let $C \geq 1$. For fixed $i$, using the i.i.d. property of the environment, we get

$$
\begin{align*}
& \mathbb{P}\left[2^{-i+1} \varepsilon n^{2 / \alpha} \sum_{e \in H_{n}(i)} \bar{g}_{n}(0, e) \geq C \varepsilon^{1-\alpha} 2^{i(\alpha-1)} \operatorname{Tn}^{2 / \alpha}\right] \\
& \quad \leq e^{-\lambda_{n} c C \varepsilon^{-\alpha} 2^{i \alpha} T} \prod_{e \in \overline{\mathbb{B}}}\left(1+p_{n, i}\left(e^{\lambda_{n} \bar{g}_{n}(0, e)}-1\right)\right) . \tag{5.8}
\end{align*}
$$

The logarithm of the product above is bounded by

$$
\begin{equation*}
\log \prod_{e \in \mathbb{B}}\left(1+p_{n, i}\left(e^{\lambda_{n} \bar{g}_{n}(0, e)}-1\right)\right) \leq \sum_{e \in \mathbb{B}} c \varepsilon^{-\alpha} 2^{i \alpha} n^{-2}\left(e^{\lambda_{n} \bar{g}_{n}(0, e)}-1\right) \tag{5.9}
\end{equation*}
$$

Taking $\lambda_{n}=(\log n) / c_{4}$ and dividing the sum according to $d(0, e)$ being smaller or larger than $b_{1}(n)$, we get

$$
\begin{equation*}
\sum_{e: d(0, e) \leq b_{1}(n)} c \varepsilon^{-\alpha} 2^{i \alpha} n^{-2}\left(e^{\lambda_{n} \bar{g}_{n}(0, e)}-1\right) \leq c 2^{i \alpha} n^{-2} b_{1}(n)^{d} \varepsilon^{-\alpha} n \xrightarrow{n \rightarrow \infty} 0 \tag{5.10}
\end{equation*}
$$

and, using $\lambda_{n} b_{1}(n)^{2-d}=O\left((\log n)^{1+2 \eta^{-1}(2-d)}\right) \rightarrow 0$,

$$
\begin{align*}
\sum_{e \in \mathbb{B}: d(0, e)>b_{1}(n)} c \varepsilon^{-\alpha} 2^{i \alpha} n^{-2}\left(e^{\lambda_{n} \bar{g}_{n}(0, e)}-1\right) & \leq \sum_{\substack{e \in \overline{\mathbb{B}}: d(0, e)>b_{1}(n)}} c \varepsilon^{-\alpha} 2^{i \alpha} n^{-2} \lambda_{n} d(0, e)^{2-d} \\
& \leq c^{\prime} 2^{i \alpha} \varepsilon^{-\alpha} T K^{2} \lambda_{n} \tag{5.11}
\end{align*}
$$

Hence, (5.8) is smaller than $\exp \left(\varepsilon^{-\alpha} 2^{i \alpha} T \lambda_{n}\left(-c C+c^{\prime} K^{2}\right)\right)$. Let $c_{5}=\left(c C-c^{\prime} K^{2}\right)$ $\log 2$, and choose $C$ large enough so that $T \varepsilon^{-\alpha} c_{5} / c_{4} \geq 2$. The probability that (5.4) fails is thus bounded by

$$
\begin{aligned}
& \mathbb{P}\left[\bigcup_{i=1}^{i_{\max }}\left\{c 2^{-i} \varepsilon n^{2 / \alpha} \sum_{e \in H_{n}(i)} g^{\mu}(0, e) \geq C \varepsilon^{1-\alpha} 2^{i(\alpha-1)} T n^{2 / \alpha}\right\}\right] \\
& \quad \leq \sum_{i=1}^{i_{\max }} \exp \left\{\varepsilon^{-\alpha} 2^{i \alpha} T \lambda_{n}\left(-c C+c^{\prime} K^{2}\right)\right\} \\
& \quad \leq \sum_{i=1}^{\infty} \exp \left\{-c_{5} \varepsilon^{-\alpha} i T \lambda_{n}\right\} \leq c n^{-c_{5} T \varepsilon^{-\alpha} / c_{4}} .
\end{aligned}
$$

Therefore, for all large $n$ (5.6) is $\mathbb{P}$-a.s. smaller than $\sum_{i=1}^{i_{\text {max }}} C \varepsilon^{1-\alpha} 2^{i(\alpha-1)} T n^{2 / \alpha}$, which is smaller than $T n^{2 / \alpha} \delta^{2} / 2$ if $\varepsilon$ is small enough. This completes the proof of (5.4) and thus of Proposition 5.1.

## 6 Coarse graining

In this section we fix $\varepsilon_{s} \in(0,1)$ and we use a coarse-graining inspired by [25] and [1] to control the contribution of edges with conductances between $\varepsilon_{s} n^{2 / \alpha}$ and $\varepsilon_{s}^{-1} n^{2 / \alpha}$ (we call these deep edges) to the rescaled clock process. We will show that these edges create jumps of $S_{n}$ in the limit and that the distribution of these jumps converges to the distribution of the jumps of a stable subordinator. Since the construction is quite technical we find it useful to give a short verbal description here.

Note that caution is necessary when speaking about the jumps, since the clock process $S_{n}$ is continuous by definition. It has, however, some very steep pieces at instants when the VSRW visits a deep edge. Moreover, the visits to deep edges occur in clusters, since after visiting a deep edge, the VSRW has a reasonable chance of returning there soon. It is thus suitable to identify the jump of $S_{n}$ with the total contribution to $S_{n}$ of one cluster of visits to one deep edge.

In the coarse-graining construction we observe the VSRW only before the exit from a large ball $B(0, K n)$. The VSRW spends a time of order $K^{2} n^{2}$ in this ball. We show that before exiting this ball, only finitely many (depending on $\varepsilon_{s}$ and $K$ ) clusters of visits occur and that they are well separated. Moreover, to prove the convergence to the stable subordinator, we need to know that the times that the VSRW spends to in going from one cluster to another are asymptotically independent and exponentially distributed. To this end we cut the trajectory of the VSRW into small pieces of spatial size $\varepsilon_{g} n$ (and temporal size $\sim \varepsilon_{g}^{2} n^{2}$ ). We show that the probability that any such piece contains a cluster of visits to a deep edge is proportional to $\varepsilon_{g}^{2}$ if $\varepsilon_{g}$ is small. Further, we show that (asymptotically as $n \rightarrow \infty$ ) no more than one cluster occurs during each piece of the path before exiting $B(0, K n)$. We further control the distribution of the contribution of one cluster to $S_{n}$. All these results are contained in Proposition 6.7 below.

The rough strategy is thus quite similar to [1,25]. There are, however, several important technical differences, which we would like to point out.

First, in previous papers the process was cut into pieces using balls of mesoscopopic radius, i.e., of radius much smaller than $n$. Due to the imprecision of our estimates on the Green function, we are forced to use macroscopic ball-like sets of radius $\varepsilon_{g} n$ here. These introduces some additional, mainly technical, difficulties.

Second, since, unlike the simple random walk, the VSRW is a process in a nonuniform environment, some additional random variables appear in the argument. As an example consider the total time spent by the CSRW in a site $x$ given $X(0)=x$. This time has mean $\mu_{x} g^{\mu}(x, x)$, where $g^{\mu}(x, x)$ is the usual Green function of the VSRW $Y$. The same was of course also true in the context of [25], except that there $Y$ was a simple random walk, and so the Green function $g(x, x), x \in \mathbb{Z}^{d}$, was constant. Hence, we need, in addition, to deal with the random variables of $g^{\mu}(x, x)=C_{\text {eff }}[x, \infty]^{-1}$, $x \in \mathbb{Z}^{d}$, which are not independent. To recover, at least partially, the independence we use Lemma 3.5 to approximate the diagonal Green functions by conductances to balls of size $O\left((\log n)^{1 / \eta}\right)$-see Lemma 6.2 below.

We now start the construction. Let $\mathbb{B}=B(0, K n)$, where $K$ is as in the previous section. Recall that we identify a set of edges $G \subset E^{d}$ with the set of their vertices.

We define

$$
\begin{align*}
& \tilde{E}_{n}(u, w)=\left\{e \in E^{d}: \mu_{e} \in[u, w) n^{2 / \alpha}\right\} \\
& \tilde{T}_{n}(u, w)=\left\{x \in \mathbb{Z}^{d}: x \in \tilde{E}_{n}(u, w), x \notin \tilde{E}_{n}(w, \infty)\right\} \tag{6.1}
\end{align*}
$$

For the arguments of this section we need that the deep edges that we observe are well separated from each other. Therefore we introduce a set of bad sites, that is the set of sites in $\tilde{T}_{n}(u, w)$ which are close to more than one edge in $\tilde{E}_{n}(u, \infty)$. We write

$$
\begin{array}{r}
\mathcal{B}_{n}(u, w)=\left\{x \in \tilde{T}_{n}(u, w): \exists y \text { s.t. }(x y) \in \tilde{E}_{n}(u, w),\right. \\
\left.B(\{x, y\}, v) \cap \tilde{T}_{n}\left(n^{-\iota}, \infty\right) \neq\{x, y\}\right\}, \tag{6.2}
\end{array}
$$

where $v=n^{\omega}, \omega \in(0,1 / d)$, is a mesoscopic scale, and $\iota \in(0,2 / \alpha)$ an arbitrary fixed small constant. We show in the next section that the bad sites are never hit before exiting $\mathbb{B}$, with a large probability. Finally, we set

$$
\begin{align*}
T_{n}(u, w) & =\tilde{T}_{n}(u, w) \backslash \mathcal{B}_{n}(u, w), \\
E_{n}(u, w) & =\left\{e=(x y) \in E^{d}: x, y \in T_{n}(u, w)\right\} . \tag{6.3}
\end{align*}
$$

We begin with an easy bound on the distance between the origin and $\tilde{E}_{n}(u, \infty)$.
Lemma 6.1 For any $u>0, \mathbb{P}$-a.s., for all but finitely many $n$,

$$
\begin{equation*}
B(0, v) \cap \tilde{E}_{n}(u, \infty)=\emptyset \tag{6.4}
\end{equation*}
$$

Proof For any edge $e$
$\mathbb{P}\left[e \in E_{n}(u, \infty)\right] \leq \mathbb{P}\left[e \in \tilde{E}_{n}(u, \infty)\right]=\mathbb{P}\left[\mu>n^{2 / \alpha} u\right]=n^{-2} u^{-\alpha}(1+o(1))$.

So

$$
\begin{equation*}
\mathbb{P}\left[B(0, v) \cap \tilde{E}_{n}(u, \infty) \neq \emptyset\right] \leq c n^{\omega d-2} u^{-\alpha}, \tag{6.6}
\end{equation*}
$$

and using Borel-Cantelli completes the proof.
We start by introducing the random variables that we use to approximate the diagonal Green function. We fix $\varepsilon_{c}>0$ and set $\bar{b}(n)=b_{2}\left(n, \varepsilon_{c} / 2\right)$ (see Lemma 3.5). For any $e \in E^{d}$ and $z \in \mathbb{Z}^{d}$ we set

$$
\begin{align*}
& \gamma_{n}(e)=C_{\mathrm{eff}}\left[e, B(e, \bar{b}(n))^{c}\right],  \tag{6.7}\\
& \gamma_{n}(z)=C_{\mathrm{eff}}\left[z, B(z, \bar{b}(n)+1)^{c}\right] . \tag{6.8}
\end{align*}
$$

We write $F_{C}$ for the law of $C_{\text {eff }}[e, \infty]$.

Lemma 6.2 (i) For all $e, \gamma_{n}(e)$ is independent of $\mu_{e}$.
(ii) If $z \in e$ then $\gamma_{n}(z) \leq \gamma_{n}(e)$.
(iii) $\mathbb{P}$-a.s. for all but finitely many $n$, for all $e \in E_{n}(u, v) \cap \mathbb{B}$ and $z \in e$,

$$
\begin{equation*}
\left(1+\varepsilon_{c}\right) \gamma_{n}(z) \geq \gamma_{n}(e) \tag{6.9}
\end{equation*}
$$

(iv) For any $k>0$ there exists $n_{0}$ such that for all $n \geq n_{0}$

$$
\begin{equation*}
\mathbb{E}\left[\gamma_{n}(z)^{k}\right] \leq \mathbb{E}\left[\gamma_{n}(e)^{k}\right]<\infty . \tag{6.10}
\end{equation*}
$$

(v) The law of $\gamma_{n}(e)$ converges weakly as $n \rightarrow \infty$ to $F_{C}$.

Proof The claims (i), (ii) and (v) are obvious. To prove (iii) let $h(x)=P_{x}^{\mu}\left[\sigma_{z}<\right.$ $\left.\tau_{B(z, \bar{b}(n))}\right]$, so that $C_{\mathrm{eff}}\left[z, B(z, \bar{b}(n))^{c}\right]=\mathcal{E}(h, h)$. Let $y \in e, y \neq z$ be the second vertex of $e$. Since $e \in E_{n}(u, v)$ we know that $\mu_{e} \geq u n^{2 / \alpha}$ and that the conductances of all edges attached to $y$ different from $e$ are at most $n^{2 / \alpha} n^{-\iota}$. Therefore, $h(y) \geq 1-c n^{-\iota}$. Let $\bar{h}(\cdot)=h(y)^{-1} h(\cdot) \wedge 1$. Then $\bar{h}=0$ on $B(z, \bar{b}(n))^{c} \supset B(e, \bar{b}(n))^{c}$ and $\overline{\bar{h}}=1$ on $e$. Hence,

$$
\begin{equation*}
C_{\mathrm{eff}}\left[z, B(z, \bar{b}(n))^{c}\right]=\mathcal{E}(h, h) \geq h(y)^{2} \mathcal{E}(\bar{h}, \bar{h}) \geq\left(1-c n^{-\iota}\right)^{2} \gamma_{n}(e) . \tag{6.11}
\end{equation*}
$$

and taking $n$ large enough gives (6.9).
To prove (iv) we set, without lost of generality, $z=0$. We define $k^{\prime}=\alpha^{-1}(k+1)$ and take $n_{0}$ such that $\bar{b}\left(n_{0}\right)>k^{\prime}$. Consider a new electric network, where all edges (xy) such that $\|x\|_{\infty}=\|y\|_{\infty}$ are short-cut, that is their conductance is set to be infinite. Recall the notation $Q(x, r)$ for cubes centre $x$ and side $r$ introduced at the start of Sect. 3. If $C^{\prime}[\cdot, \cdot]$ denotes the conductance in this new network we have, for $n \geq n_{0}$,

$$
\begin{equation*}
\mathbb{P}\left[\gamma_{n}(z) \geq u\right] \leq \mathbb{P}\left[C_{\mathrm{eff}}\left[z, Q\left(z, k^{\prime}\right)^{c}\right] \geq u\right] \leq \mathbb{P}\left[C^{\prime}\left[z, Q\left(z, k^{\prime}\right)^{c}\right] \geq u\right] \tag{6.12}
\end{equation*}
$$

Let $L_{i}=\left\{(x y) \in E^{d}:\|x\|_{\infty}=i,\|y\|_{\infty}=i-1\right\}$. Then

$$
\begin{equation*}
C^{\prime}\left[z, Q\left(z, k^{\prime}\right)^{c}\right]=\left\{\sum_{i=1}^{k^{\prime}}\left(\sum_{e \in L_{i}} \mu_{e}\right)^{-1}\right\}^{-1}=:\left(\sum_{i=1}^{k^{\prime}}\left(C_{L_{i}}\right)^{-1}\right)^{-1} \tag{6.13}
\end{equation*}
$$

Hence,

$$
\begin{align*}
& \mathbb{P}\left[C^{\prime}\left[z, Q\left(z, k^{\prime}\right)^{c}\right] \geq u\right] \leq \mathbb{P}\left[\min \left\{C_{L_{i}}: i=1, \ldots, k^{\prime}\right\} \geq u\right] \\
& =\prod_{i=1}^{k^{\prime}} \mathbb{P}\left[C_{L_{i}} \geq u\right] \leq \prod_{i=1}^{k^{\prime}} c(d, i) u^{-\alpha} \leq c\left(d, k^{\prime}\right) u^{-\alpha k^{\prime}} \leq c u^{-k-1} . \tag{6.14}
\end{align*}
$$

This proves that $\mathbb{E}\left[\gamma_{n}(z)^{k}\right]<\infty$. If $z$ is replaced by $e$, the proof is analogous.

We will split the sets $E_{n}(u, v)$ according to the value of $\gamma_{n}(e)$. To this end we define

$$
\begin{align*}
E_{n}\left(u, v, w, w^{\prime}\right) & =\left\{e \in E_{n}(u, w): \gamma_{n}(e) \in\left[v, w^{\prime}\right)\right\} \\
T_{n}\left(u, v, w, w^{\prime}\right) & =\mathbb{Z}^{d} \cap E_{n}\left(u, v, w, w^{\prime}\right) \tag{6.15}
\end{align*}
$$

We need the next estimate on the size of these sets and their spatial distribution.
Lemma 6.3 Let $u, v, w, w^{\prime}>0$ and $\delta, \varepsilon_{b}>0$ be fixed. Then $\mathbb{P}$ - $a . s$. for all but finitely many $n$,

$$
\begin{align*}
& n^{d} \varepsilon_{b}^{d} d(1-\delta) p_{n}\left(u, v, w, w^{\prime}\right) \leq \inf _{x \in \varepsilon_{b} n \mathbb{Z}^{d} \cap \mathbb{B}}\left|Q\left(x, \varepsilon_{b} n\right) \cap E_{n}\left(u, v, w, w^{\prime}\right)\right| \\
& \quad \leq \sup _{x \in \varepsilon_{b} n \mathbb{Z}^{d} \cap \mathbb{B}}\left|Q\left(x, \varepsilon_{b} n\right) \cap E_{n}\left(u, v, w, w^{\prime}\right)\right| \leq n^{d} \varepsilon_{b}^{d} d(1+\delta) p_{n}\left(u, v, w, w^{\prime}\right), \tag{6.16}
\end{align*}
$$

where $p_{n}\left(u, v, w, w^{\prime}\right)=n^{-2}\left(u^{-\alpha}-w^{-\alpha}\right) F_{C}\left(\left[v, w^{\prime}\right)\right)$.
Proof The number of points in $\varepsilon_{b} n \mathbb{Z} \cap \mathbb{B}$ is bounded uniformly in $n$. It is hence sufficient to show that

$$
\begin{equation*}
\sum_{n=1}^{\infty} \mathbb{P}\left[\frac{\left|Q\left(0, \varepsilon_{b} n\right) \cap E_{n}\left(u, v, w, w^{\prime}\right)\right|}{n^{d} \varepsilon_{b}^{d} d p_{n}\left(u, v, w, w^{\prime}\right)} \notin(1-\delta, 1+\delta)\right]<\infty . \tag{6.17}
\end{equation*}
$$

The lemma will then follow from the Borel-Cantelli lemma and the translation invariance of $\mu$.

To prove (6.17) we set $A_{n}=E^{d} \cap Q(0,3 v)$ and for $e \in A_{n}$ we define $G_{n}(e)=$ $Q\left(0, \varepsilon_{b} n\right) \cap\left(e+3 v \mathbb{Z}^{d}\right)$. Hence

$$
\begin{equation*}
\left|A_{n}\right|=d(3 v)^{d}(1+o(1)) \quad \text { and } \quad\left|G_{n}(e)\right|=\left(n \varepsilon_{b} / 3 v\right)^{d}(1+o(1)) . \tag{6.18}
\end{equation*}
$$

The series (6.17) is bounded by

$$
\begin{equation*}
\sum_{n=1}^{\infty} \sum_{e \in A_{n}} \mathbb{P}\left[\frac{\left|G_{n}(e) \cap E_{n}\left(u, v, w, w^{\prime}\right)\right|}{\left|G_{n}(e)\right| p_{n}\left(u, v, w, w^{\prime}\right)} \notin(1-\delta, 1+\delta)\right] . \tag{6.19}
\end{equation*}
$$

From the definitions of $E_{n}\left(u, v, w, w^{\prime}\right)$ and $\gamma_{n}(e)$, and the independence of $\mu_{e}$ and $\gamma_{n}(e)$ given by Lemma 6.2(i) it follows that $\left\{\mathbf{1}_{e^{\prime} \in E_{n}\left(u, v, w, w^{\prime}\right)}: e^{\prime} \in G_{n}(e)\right\}$ is an i.i.d. family of Bernoulli random variables with success probability

$$
\begin{equation*}
\mathbb{P}\left[\mu_{e} \in n^{2 / \alpha}[u, w)\right] \mathbb{P}\left[\mu_{e^{\prime}} \leq n^{-\iota} n^{2 / \alpha} \forall e^{\prime} \in B(e, v) \backslash e, \gamma_{n}(e) \in\left[v, w^{\prime}\right)\right] . \tag{6.20}
\end{equation*}
$$

The first probability in (6.20) is $n^{-2}\left(u^{-\alpha}-w^{-\alpha}\right)(1+o(1))$. To control the second term in (6.20) note that if $\iota$ is sufficiently small, then $\mathbb{P}\left[\mu_{e^{\prime}} \leq n^{-\iota} n^{2 / \alpha} \forall e^{\prime} \in B(e, v) \backslash e\right]$ $\xrightarrow{n \rightarrow \infty} 1$. Therefore it follows from Lemma 6.2(v) that the second probability in (6.20) converges to $F_{C}\left(\left[v, w^{\prime}\right)\right)$. Thus (6.20) equals $p_{n}\left(u, v, w, w^{\prime}\right)(1+o(1))$.

It follows that $\left|G_{n}(e) \cap E_{n}\left(u, v, w, w^{\prime}\right)\right|$ has binomial distribution with parameters $\left|G_{n}(e)\right|$ and $p_{n}\left(u, v, w, w^{\prime}\right)(1+o(1))$. Standard estimates on the binomial distribution yield that for any $\delta>0$ there is $c>0$ such that all summands in (6.19) are bounded by $\exp \left(-c\left|G_{n}(e)\right| p_{n}\left(u, v, w, w^{\prime}\right)\right)$. Since $\left|A_{n}\right|=O\left(v^{d}\right)$, the series (6.17) is summable and the proof is finished.

The papers $[1,25]$ used a family of (large but mesoscopic) Euclidean balls to coarse grain the trajectory. They then used the fact that the probability that the simple random walk exits a large ball $B(x, R)$ at a given site $x^{\prime} \in \partial B(x, R)$ can be bounded from above by $c R^{1-d}$. This fact was used to control the probability that the simple random walk does not exit a large ball too close to a site in $T_{n}\left(\varepsilon_{s}, \infty\right)$.

It is not clear that bounds of this type hold for the VSRW $Y_{t}$, since large values of $\mu_{e}$ for edges close to a boundary point $x^{\prime}$ might have a substantial effect on the probability of $Y$ exiting at $x^{\prime}$. We therefore replace balls by ball-like sets whose boundary avoids $\tilde{T}_{n}(\varepsilon, \infty)$ and $\tilde{E}_{n}(\varepsilon, \infty)$. This construction will then allow us to apply the coarse-graining without requiring such precise information on the exit distribution of $Y$ from large balls.

Lemma 6.4 Let $\varepsilon_{s}>0$ and $K \geq 1$. Then $\mathbb{P}$-a.s. for all but finitely many $n$, for all $x \in \mathbb{Z}^{d}$ and $r \in(0, K n)$ there exist sets $\mathfrak{B}_{n}(x, r)$ which satisfy
(i) $\mathfrak{B}_{n}(x, r)$ is simply connected in $\mathbb{Z}^{d}$.
(ii) For all $x \in B(0, K n), B(x, r) \subset \mathfrak{B}_{n}(x, r) \subset B(x, r+3 \nu \log n)$.
(iii) For all $x \in B(0, K n), \partial \mathfrak{B}_{n}(x, r) \cap \bigcup_{e \in \tilde{E}_{n}\left(\varepsilon_{s}, \infty\right)} B(e, v)=\emptyset$.

Proof If $x \notin B(0, K n)$ we just take $\mathfrak{B}_{n}(x, r)=B(x, r)$.
Let $\mathbb{B}=B(0, K n)$ and let $\mathcal{O}=\bigcup_{e \in \tilde{E}_{n}\left(\varepsilon_{s}, \infty\right)} B(e, v) \cap \mathbb{Z}^{d}$, and let $\mathcal{C}(z)$ be the $\mathbb{Z}^{d}$ connected component of $\mathcal{O}$ containing $z$. Finally, let $K(z)=\mid\left\{(x y) \in \tilde{E}_{n}\left(\varepsilon_{s}, \infty\right)\right.$ : $\{x, y\} \subset \mathcal{C}(z)\} \mid$. Using (6.6) we have

$$
\begin{align*}
& \mathbb{P}[\operatorname{diam} \mathcal{C}(z) \geq 2 k v] \leq \mathbb{P}[K(z) \geq k] \\
& \quad \leq\left(c v^{d} n^{-2}\right)\left(2 c v^{d} n^{-2}\right) \ldots\left(k c v^{d} n^{-2}\right) \leq c k!\left(n^{-2} v^{d}\right)^{k} \tag{6.21}
\end{align*}
$$

There are, $\mathbb{P}$-a.s., for all but finitely many $n$, at most $c n^{-2} n^{d}$ sites in $\tilde{T}_{n}\left(\varepsilon_{s}, \infty\right) \cap 2 \mathbb{B}$. Hence,

$$
\begin{equation*}
\mathbb{P}\left[\max _{z \in 2 \mathbb{B}} \operatorname{diam} \mathcal{C}(z) \geq 2 v \log n\right] \leq c n^{d-2}\lceil\log n\rceil!\left(n^{-2} v^{d}\right)^{\lceil\log n\rceil} \tag{6.22}
\end{equation*}
$$

which is summable if $v \leq(n / \log n)^{2 / d}$. The Borel-Cantelli lemma then implies that (for all $n$ large) the largest component of $\mathcal{O} \cap 2 \mathbb{B}$ has diameter smaller than $2 v \log n$. The claim of the lemma then follows easily, by taking $\mathfrak{B}_{n}(x, r)$ to be the union of $B(x, r)$ and every connected component of $\mathcal{O}$ which intersects $B(x, r)$.

This lemma does not uniquely specify the family $\left(\mathfrak{B}_{n}(x, r)\right)$. In what follows we will take $\mathfrak{B}_{n}(x, r)$ to be the set given by the procedure in the proof of Lemma 6.4.

We finally have all the ingredients that we need to start the coarse-graining construction. Let $\varepsilon_{g}>0, t_{n}(0)=0, y_{n}(0)=0$ and for $i \geq 1$ let

$$
\begin{align*}
t_{n}(i) & =\inf \left\{t>t_{n}(i-1): Y(t) \notin \mathfrak{B}_{n}\left(y_{n}(i-1), \varepsilon_{g} n\right)\right\},  \tag{6.23}\\
y_{n}(i) & =Y\left(t_{n}(i)\right) .
\end{align*}
$$

We denote by $Y[i]=\left\{Y(s): t_{n}(i) \leq s<t_{n}(i+1)\right\}$ the $i$ th a piece of the trajectory of $Y$.

Observe that, by the definition of $\mathfrak{B}_{n}(x, r), d\left(y_{n}(i), \tilde{E}_{n}\left(\varepsilon_{s}, \infty\right)\right) \geq v$ for all $i \geq 1$. The same is $\mathbb{P}$-a.s. true also for $0=y_{n}(0)$, by Lemma 6.1.

The functional central limit theorem (Proposition 3.1) and the construction of the sets $\mathfrak{B}_{n}(x, r)$ imply the following lemma.

Lemma 6.5 Let $\mathrm{BM}_{d}$ be the standard Brownian motion in $\mathbb{R}^{d}$ and let $W(\cdot)=\mathcal{C}_{Y}$ $\mathrm{BM}_{d}(\cdot)$ (see Proposition 3.1). Define $t_{\infty}(0)=0$ and, for $i \geq 1, t_{\infty}(i)=\inf \{t \geq$ $\left.t_{\infty}(i-1): W \notin B\left(W\left(t_{\infty}(i-1)\right), \varepsilon_{g}\right)\right\}$. Then, $\mathbb{P}$-a.s., under $P_{0}^{\mu}$, the law of the sequence ( $n^{-2} t_{n}(i): i \in \mathbb{N}$ ) converges as $n \rightarrow \infty$ to the law of $\left(t_{\infty}(i): i \in \mathbb{N}\right)$. In particular, for every $s>0$ and $\delta>0$ there exists $\varepsilon_{g}>0$ such that $\mathbb{P}$-a.s. for all but finitely many $n$,

$$
\begin{equation*}
P_{0}^{\mu}\left[n^{-2} t_{n}\left(\left\lfloor d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s\right\rfloor\right) \notin(s-\delta, s+\delta)\right]<\delta . \tag{6.24}
\end{equation*}
$$

Proof The first claim is a direct consequence of the functional limit theorem. From standard properties of Brownian motion, we have that $\left(t_{\infty}(i)-t_{\infty}(i-1): i \geq 1\right)$ is an i.i.d. sequence, the expectation of $t_{\infty}(1)-t_{\infty}(0)=t_{\infty}(1)$ equals $\varepsilon_{g}^{2} \mathcal{C}_{Y}^{-2} d^{-1}$, and its variance is finite. Hence, by the weak law of large numbers (for $\varepsilon_{g} \rightarrow 0$ ), there exists $\varepsilon_{g}$ small such that $\mathbb{P}\left[t_{\infty}\left(d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s\right) \notin(s-\delta, s+\delta)\right] \leq \delta / 2$. Since, as $n \rightarrow \infty$, $\mathbb{P}$-a.s., $n^{-2} t_{n}\left(d \mathcal{C}_{Y} \varepsilon_{g}^{-2} s\right)$ converges in law to $t_{\infty}\left(d \mathcal{C}_{Y} \varepsilon_{g}^{-2} s\right)$, the proof is finished.

We define

$$
\begin{equation*}
s_{n}\left(i ; u, v, w, w^{\prime}\right)=n^{-2 / \alpha} \int_{t_{n}(i)}^{t_{n}(i+1)} \mu_{Y(t)} \mathbf{1}\left\{Y(t) \in T_{n}\left(u, v, w, w^{\prime}\right)\right\} \mathrm{d} t \tag{6.25}
\end{equation*}
$$

this is the increment of the clock process between times $t_{n}(i)$ and $t_{n}(i+1)$ caused by sites in $T_{n}\left(u, v, w, w^{\prime}\right)$.

We now wish to calculate the distribution of these random variables, and begin with the following lemma.

Lemma 6.6 Let $z \in \mathbb{B}_{0}=B\left(0,\left(K-\varepsilon_{g}\right) n\right)$, $e=(x y) \in E_{n}\left(\varepsilon_{s}\right) \cap \mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)$ be such that $\mu_{e}=u n^{2 / \alpha}, \gamma_{n}(e)=v$. Then, $\mathbb{P}$-a.s., the distribution of

$$
\begin{equation*}
n^{-2 / \alpha} \int_{0}^{\tau_{\mathfrak{B}}\left(z, g_{g} n\right)} 1\{Y(t) \in\{x, y\}\} \mu_{Y(t)} \mathrm{d} t \tag{6.26}
\end{equation*}
$$

under both $P_{x}^{\mu}$ and $P_{y}^{\mu}$ converges as $n \rightarrow \infty$ to the exponential distribution with mean $2 u / v$.
Proof Since, $e \in E_{n}\left(\varepsilon_{s}\right)$ (and thus all edges incident to $e$ have conductance smaller than $\left.n^{2 / \alpha} n^{-\imath}\right)$, it is easy to see that $\mu_{x}, \mu_{y}=\mu_{e}(1+o(1))$. Moreover, Lemma 6.2 implies that $\gamma_{n}(x), \gamma_{n}(y)=\gamma_{n}(e)\left(1+O\left(\varepsilon_{c}\right)\right)$. Let $N_{x}, N_{y}$ be the number of visits to $x$ and $y$ by $Y$ before exiting $\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)$ and let $T_{x}, T_{y}$ be the total time spend there, so that $T_{x}=\int_{0}^{\tau_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}} \mathbf{1}\{Y(t)=x\} \mathrm{d} t$. It is well known fact that under $P_{x}^{\mu}, T_{x}$ has exponential distribution with mean $R_{\text {eff }}\left[x, \partial \mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)\right]$. By the definition of $\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)$, $d\left(\{x, y\}, \partial \mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)\right) \geq v$. Therefore, $\gamma_{n}(x)^{-1} \leq R_{\text {eff }}\left[x, \partial \mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)\right] \leq \gamma_{n}(x)^{-1}$ $\left(1+\varepsilon_{c}\right)$. Hence, the mean of $T_{x}$ is equal to $1 / \gamma_{n}(x)\left(1+O\left(\varepsilon_{c}\right)\right)=v^{-1}\left(1+O\left(\varepsilon_{c}\right)\right)$, and thus the random variable $N_{x}$ has geometrical distribution with mean $\mu_{x} v^{-1}\left(1+O\left(\varepsilon_{c}\right)\right)$. Moreover, since $p_{x y}=1-O\left(n^{-l}\right)$, it follows that with probability converging to one, $N_{y} \geq\left(1-o\left(n^{-\iota / 2}\right)\right) N_{x}$. Similar argument gives also $N_{x} \geq\left(1-o\left(n^{-\iota / 2}\right)\right) N_{y}$ with a large probability. Therefore, $T_{x}+T_{y}$ is asymptotically exponentially distributed with mean $2 v^{-1}\left(1+O\left(\varepsilon_{c}\right)\right)$. Taking $\varepsilon_{c}$ arbitrarily small, the distribution of (6.26), which is $n^{-2 / \alpha}\left(T_{x} \mu_{x}+T_{y} \mu_{y}\right)$ under $P_{x}^{\mu}$, converges to the exponential distribution with mean $2 u / v$ 。

The following proposition, which is the main result of this section, gives the distribution of $s_{n}\left(i, \varepsilon_{s}, \varepsilon_{s}, \varepsilon_{s}^{-1}, \varepsilon_{s}^{-1}\right)$.
Proposition 6.7 Let $T, \varepsilon_{s}, \varepsilon_{g}>0$. Define $s_{n}(i)=s_{n}\left(i, \varepsilon_{s}, \varepsilon_{s}, \varepsilon_{s}^{-1}, \varepsilon_{s}^{-1}\right)$. Then, $\mathbb{P}$ a.s., under $P_{0}^{\mu}$, the sequence $\left(s_{n}(i), i \in\left\{1, \ldots, \varepsilon_{g}^{-2} T\right\}\right)$ converges as $n \rightarrow \infty$ to an i.i.d. sequence $\left(s_{\infty}(i): i \in\left\{1, \ldots, \varepsilon_{g}^{-2} T\right\}\right)$. Moreover, as $\varepsilon_{g} \rightarrow 0$,

$$
\begin{align*}
& P_{0}^{\mu}\left[s_{\infty}(i)=0\right]=1-c_{\varepsilon_{s}} \varepsilon_{g}^{2}+o\left(\varepsilon_{g}^{2}\right),  \tag{6.27}\\
& P_{0}^{\mu}\left[s_{\infty}(i) \in A\right]=\varepsilon_{g}^{2} v_{\varepsilon_{s}}(A)+o\left(\varepsilon_{g}^{2}\right), \quad A \subset(0, \infty), \tag{6.28}
\end{align*}
$$

where

$$
\begin{equation*}
c_{\varepsilon_{s}}=\mathcal{C}_{Y}^{-2} \int_{\varepsilon_{s}}^{\varepsilon_{s}^{-1}} \int_{\varepsilon_{s}}^{\varepsilon_{s}^{-1}} \alpha v u^{-\alpha-1} \mathrm{~d} u F_{C}(\mathrm{~d} v) \tag{6.29}
\end{equation*}
$$

and

$$
\begin{equation*}
\nu_{\varepsilon_{s}}(\mathrm{~d} x)=\mathcal{C}_{Y}^{-2} \int_{\varepsilon_{s}}^{\varepsilon_{s}^{-1}} \int_{\varepsilon_{s}}^{\varepsilon_{s}^{-1}} \frac{v}{2 u} \exp \left\{-\frac{x v}{2 u}\right\} \alpha v u^{-\alpha-1} \mathrm{~d} u F_{C}(\mathrm{~d} v) \mathrm{d} x \tag{6.30}
\end{equation*}
$$

For the reader's convenience, before starting the proof we summarise the role of various small and large parameters appearing there and their dependence. The parameter $\varepsilon_{s}$ in the definition of a deep edge is kept constant in this section: it will be chosen to be small in Sect. 8. The same holds for $\varepsilon_{g}$, which determines the scale of the coarse-graining. In the proof we will choose $\delta>0$ which will eventually tend to zero. All others $\varepsilon$ 's appearing in the proof are chosen so that our approximations have a
multiplicative error $1+O(\delta)$. These include $\varepsilon_{o}$, which gives the size of the zones close to the centre and the border of $\mathcal{B}_{n}\left(x, \varepsilon_{g} n\right)$ where we do not have precise estimates on the Green function (see Proposition 4.3), $\varepsilon_{b}$, the parameter giving the size of boxes where we apply the homogeneity estimates of Lemma 6.3 , and $\varepsilon_{c}$, which controls the approximation of the diagonal Green function by $\gamma_{n}(z)$. These three $\varepsilon$ 's are mutually independent. Finally $K$ is chosen so, that the VSRW does not exit $\mathbb{B}$ before $t_{n}\left(\varepsilon_{g}^{-2} T\right)$ with a large probability depending on $\varepsilon_{g}$.

Proof We first prove (6.27) and (6.28) for a fixed $i \leq \varepsilon_{g}^{-2} T$. We comment on the asymptotic independence of $s_{n}(i)$ 's at the end of the proof.

We choose $K$ as function of $\varepsilon_{g}$, such that $P_{0}^{\mu}\left[\tau_{\mathbb{B}_{0}} \leq t_{n}\left(\varepsilon_{g}^{-2} T\right)\right]=o\left(\varepsilon_{g}^{2}\right)$, where $\mathbb{B}_{0}=B\left(0,\left(K-\varepsilon_{g}\right) n\right)$. We write, as usual, $\mathbb{B}=B(0, K n)$. Let $E_{n}\left(\varepsilon_{s}\right)=E_{n}\left(\varepsilon_{s}, \varepsilon_{s}, \varepsilon_{s}^{-1}\right.$, $\varepsilon_{s}^{-1}$ ), and let

$$
\begin{equation*}
W_{n}(i)=\sum_{e \in E_{n}\left(\varepsilon_{s}\right)} P_{0}^{\mu}[Y[i] \cap e \neq \emptyset] . \tag{6.31}
\end{equation*}
$$

We first show that, $\mathbb{P}$-a.s.,

$$
\begin{equation*}
\lim _{n \rightarrow \infty} W_{n}(i)=c_{\varepsilon_{s}} \varepsilon_{g}^{2}+o\left(\varepsilon_{g}^{2}\right) \tag{6.32}
\end{equation*}
$$

For any edge $e \in E_{n}\left(\varepsilon_{s}\right)$ we choose one of its vertices as its 'representative'. Let $H_{n}\left(\varepsilon_{s}\right)=H_{n}\left(\varepsilon_{s}, \varepsilon_{s}, \varepsilon_{s}^{-1}, \varepsilon_{s}^{-1}\right)$ stands for the set of these representatives. Since, due to the definition of $E_{n}\left(\varepsilon_{s}\right)$, all edges incident to $E_{n}\left(\varepsilon_{s}\right)$ have conductances smaller than $n^{2 / \alpha} n^{-\iota}$, the event that $Y[i]$ intersects an edge $e \in E_{n}\left(\varepsilon_{s}\right)$ but not its representative has probability $O\left(n^{-\iota}\right)$. Therefore,

$$
\begin{equation*}
W_{n}(i)=\left(1+O\left(n^{-\iota}\right)\right) \sum_{x \in H_{n}\left(\varepsilon_{s}\right)} P_{0}^{\mu}[x \in Y[i]] . \tag{6.33}
\end{equation*}
$$

Let $\varepsilon_{m}>0$ small and let $\varepsilon_{s}=u_{0}<u_{1}<\cdots<u_{q}=\varepsilon_{s}^{-1}$ be such that $u_{i}-u_{i-1} \in$ $\left(\varepsilon_{m}, 2 \varepsilon_{m}\right), 1 \leq i \leq q$. Using the notation $H_{n}(j, k)=H_{n}\left(u_{j}, u_{k}, u_{j+1}, u_{k+1}\right)$, we get

$$
\begin{equation*}
W_{n}(i)=\left(1+O\left(n^{-\iota}\right)\right) \sum_{j=0}^{q-1} \sum_{k=0}^{q-1} \sum_{x \in H_{n}(j, k)} P_{0}^{\mu}[x \in Y[i]] . \tag{6.34}
\end{equation*}
$$

Decomposing on the value $y_{n}(i)$ of the starting point of $Y[i]$, the interior sum equals

$$
\begin{align*}
\sum_{x \in H_{n}(j, k)} P_{0}^{\mu}[x \in Y[i]] & =\sum_{z \in \mathbb{Z}^{d}} P_{0}^{\mu}\left[y_{n}(i)=z\right] \sum_{x \in H_{n}(j, k)} P_{z}^{\mu}\left[\sigma_{x}<\tau_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}\right] \\
& =\sum_{z \in \mathbb{B}_{0}} P_{0}^{\mu}\left[y_{n}(i)=z\right] \sum_{x \in H_{n}(j, k)} \frac{g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(z, x)}{g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(x, x)}+o\left(\varepsilon_{g}^{2}\right), \tag{6.35}
\end{align*}
$$

where the error results from restricting the summation to $\mathbb{B}_{0}$. By (3.9), Lemmas 3.5, 6.2 , and the definition of $\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)$, for $x \in H_{n}(k, j) \cap \mathbb{B}$,

$$
\begin{aligned}
u_{k}\left(1+2 \varepsilon_{m} \varepsilon_{s}^{-1}\right) & \geq u_{k+1} \geq \gamma_{n}(x) \geq g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(x, x)^{-1} \\
& \geq\left(1-\varepsilon_{c}\right) \gamma_{n}(x) \geq\left(1-\varepsilon_{c}\right) u_{k} .
\end{aligned}
$$

Hence, up to an multiplicative error $1+O\left(\varepsilon_{m} / \varepsilon_{s}\right)+O\left(\varepsilon_{c}\right)$, we can replace the Green function in the denominator of (6.35) by $u_{k}^{-1}$.

To apply Proposition 4.3, we choose $\delta>0$ and $\varepsilon_{o} \in\left(0, \varepsilon_{g} / 100\right)$, say. Then, for this $\delta, \varepsilon_{o}$ and $K, \varepsilon_{g}$ as above, we fix $\varepsilon_{b} \in\left(0, \varepsilon_{o} / 2\right)$ as in Lemma 4.1, and we split the sum in (6.35) again,

$$
\begin{align*}
& o\left(\varepsilon_{g}^{2}\right)+u_{k}\left(1+O\left(\varepsilon_{c}, \varepsilon_{m} / \varepsilon_{s}\right)\right) \sum_{z \in \mathbb{B}_{0}} P_{0}^{\mu}\left[y_{n}(i)=z\right] \\
& \quad \times\left\{\sum_{\substack{y \in \varepsilon_{b} n \mathbb{Z} \\
|z-y|<2 \varepsilon_{o} n}}+\sum_{\substack{y \in \varepsilon_{b} n \mathbb{Z} \\
2 \varepsilon_{o} n \leq|z-y| \leq\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n}}+\sum_{\substack{y \in \varepsilon_{b} n \mathbb{Z} \\
|z-y|>\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n}}\right\} \\
& \quad \sum_{x \in H_{n}(j, k) \cap Q\left(y, \varepsilon_{b} n\right)} g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(z, x) . \tag{6.36}
\end{align*}
$$

We first estimate the second sum over $y$, since the other two sums contribute to the error only. By Lemma 6.3, any square $Q\left(y, \varepsilon_{b} n\right)$ contains $d \varepsilon_{b}^{d} n^{d} p_{n}\left(u_{j}, u_{k}, u_{j+1}, u_{k+1}\right)(1+$ $O(\delta))$ points in $H_{n}(j, k)$. On the other hand, since $|y-x|>2 \varepsilon_{o} n$, by Proposition 4.1 the Green function $g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(z, x)$ is almost constant in this square and can be approximated by $g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(z, y)(1+O(\delta))$. Hence, the second sum over $y$ in (6.36) equals

$$
\begin{align*}
& d \varepsilon_{b}^{d} n^{d} p_{n}\left(u_{j}, u_{k}, u_{j+1}, u_{k+1}\right) \sum_{\substack{y \in \varepsilon_{b} n \mathbb{Z}^{d} \\
|z-y| \in\left[2 \varepsilon_{o} n,\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n\right]}} g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(z, y)(1+O(\delta)) \\
& =d p_{n}\left(u_{j}, u_{k}, u_{j+1}, u_{k+1}\right) \sum_{\substack{y \in \varepsilon_{n} n \mathbb{Z}^{d} \\
|z-y| \in\left[2 \varepsilon_{o} n,\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n\right]}} \sum_{x \in Q\left(y, n \varepsilon_{b} n\right)} g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}(z, x)(1+O(\delta)), \tag{6.37}
\end{align*}
$$

where we once more used the regularity of the Green function. Using Proposition 4.3, the sumation over $y$ and $x$ in (6.37) can be estimated, and equals

$$
\begin{equation*}
n^{2} \frac{\varepsilon_{g}^{2}}{d \mathcal{C}_{Y}}\left(1+O\left(\varepsilon_{o} / \varepsilon_{g}\right)+c_{1}(\delta)\right) \tag{6.38}
\end{equation*}
$$

Inserting these estimates back into (6.34), and replacing the summation over $j$ and $k$ by an integration, we get that the contribution of the second sum over $y$ in (6.36) to
(6.34) equals $\varepsilon_{\varepsilon_{s}} \varepsilon_{g}^{2}\left(1+O\left(\delta, c_{1}(\delta), \varepsilon_{m} / \varepsilon_{s}, \varepsilon_{c}\right)\right)+o\left(\varepsilon_{g}^{2}\right)$. This can be made arbitrarily close to the right-hand side of (6.32) by choosing $\delta, \varepsilon_{o}, \varepsilon_{m}$ and $\varepsilon_{c}$ small.

It remains to show that the other two summations over $y$ in (6.36) contribute to the error term only. For $|y-z| \geq\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n$, observe first that

$$
\begin{equation*}
\left|H_{n}\left(\varepsilon_{s}\right) \cap \mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right) \cap \bigcup_{\substack{\in \varepsilon_{b} n \mathbb{Z} \\|z-y|>\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n}} Q\left(y, \varepsilon_{b} n\right)\right| \leq C n^{d} \varepsilon_{o} \varepsilon_{g}^{d-1} n^{-2} \varepsilon_{s}^{-\alpha} . \tag{6.39}
\end{equation*}
$$

For such $y$, using the global upper bound on the Green function (3.11),

$$
\begin{equation*}
g_{\mathfrak{B}_{n}\left(z, \varepsilon_{g} n\right)}^{\mu}(z, x) \leq g^{\mu}(z, x) \leq C\left(\varepsilon_{g}-3 \varepsilon_{o}\right)^{2-d} n^{2-d} . \tag{6.40}
\end{equation*}
$$

Inserting these two estimates into (6.36), it is easy to see that the contribution of $y$ with $|y-z|>\left(\varepsilon_{g}-2 \varepsilon_{o}\right) n$ is bounded from above by $C\left(\varepsilon_{g}-3 \varepsilon_{o}\right)^{2-d} \varepsilon_{g}^{d-1} \varepsilon_{s}^{-\alpha} \varepsilon_{o}$, which can be made arbitrarily small by choosing $\varepsilon_{o}$ small.

It is slightly more delicate to bound the contribution of $|y-z|<2 \varepsilon_{o} n$. We use similar argument as in [25]. We need to improve our homogeneity estimates (Lemma 6.3) first: Let $j_{\max }$ be the smallest integer satisfying $2^{i} v \geq 2 \varepsilon_{o} n$, that is $j_{\max }=O(\log n)$. Then there exists $K>0$ such that $\mathbb{P}$-a.s. for all large $n$, all $j \in\left\{0, \ldots, j_{\max }\right\}$ and all $x \in \mathbb{B}$,

$$
\begin{equation*}
\left|Q\left(x, 2^{j} v\right) \cap H_{n}\left(\varepsilon_{s}\right)\right| \leq K\left(\log n \vee 2^{j d} v^{d} n^{-2} \varepsilon_{s}^{-\alpha}\right) . \tag{6.41}
\end{equation*}
$$

Indeed, by definition $\left|Q\left(x, 2^{j} v\right) \cap H_{n}\left(\varepsilon_{s}\right)\right| \leq\left|Q\left(x, 2^{j} v\right) \cap \tilde{T}_{n}\left(\varepsilon_{s}, \varepsilon_{s}, \varepsilon_{s}^{-1}, \varepsilon_{s}^{-1}\right)\right|$. Moreover, if $x_{1}, x_{2}$ be such that $\left|x_{1}-x_{2}\right| \geq 2 \bar{b}$, then the events that $x_{1}$, resp. $x_{2}$, is in $T_{n}\left(\varepsilon_{s}, \varepsilon_{s}, \varepsilon_{s}^{-1}, \varepsilon_{s}^{-1}\right)$ are independent. Hence, the probability of the complement of (6.41) is bounded, using a similar decomposition to summations of i.i.d. random variables as in Lemma 6.3, by

$$
\begin{align*}
& c \sum_{x \in \mathbb{B}} \sum_{j=0}^{j_{\max }} d \bar{b}^{d} e^{-K \lambda\left(\log n \vee 2^{j d} v^{d} n^{-2} \varepsilon_{s}^{-\alpha}\right) / d \bar{b}^{d}}\left(1+c \varepsilon_{s}^{-\alpha}\left(e^{\lambda}-1\right) n^{-2}\right)^{2^{j d} v^{d} / \bar{b}^{d}} \\
& \quad \leq c n^{d} j_{\max } \bar{b}^{d} n^{-K^{\prime}}, \tag{6.42}
\end{align*}
$$

where $K^{\prime}$ can be made arbitrarily large by choosing $K$ large. This proves (6.41).
Let $E=\{-1,0,1\}^{d} \backslash\{(0,0,0)\}$ and let $\mathcal{O}_{j}$ be the union of $3^{d}-1$ cubes of size $2^{j} v$ centred at $y_{n}(i)+2^{j} v E$,

$$
\begin{equation*}
\mathcal{O}_{j}=\bigcup_{x \in E} Q\left(y_{n}(i)+x 2^{j} v, 2^{j} v\right) \tag{6.43}
\end{equation*}
$$

We cover $B\left(y_{n}(i), 2 \varepsilon_{o} n\right) \backslash B\left(y_{n}(i), \nu\right)$ by $\bigcup_{j=0}^{j_{\max }} \mathcal{O}_{j}$. It is not necessary to cover the interior ball $B\left(y_{n}(i), \nu\right)$, since there are not any sites from $H_{n}\left(\varepsilon_{s}\right)$ by the definition
of sets $\mathfrak{B}_{n}(x, r)$. The contribution of the sum over $y<2 \varepsilon_{o} n$ in (6.36) to (6.34) is bounded from above, using the global upper estimate on the Green function (3.11) only, by

$$
\begin{align*}
& C \varepsilon_{s}^{-1} \sum_{j=0}^{j_{\max }}\left(\log n \vee 2^{j d} v^{d} n^{-2} \varepsilon_{s}^{-\alpha}\right)\left(2^{j} v\right)^{2-d} \\
& \quad \leq C \varepsilon_{s}^{-1} \sum_{j=0}^{j_{\max }}\left\{\left(2^{j} v\right)^{2-d} \log n+2^{2 j} v^{2} n^{-2} \varepsilon_{s}^{-\alpha}\right\}, \tag{6.44}
\end{align*}
$$

where the $\varepsilon_{s}^{-1}$ before the summation comes from the bound on the Green function in the denominator of (6.35). The last display is bounded by $C \varepsilon_{s}^{-1}\left(v^{2-d} \log n+\varepsilon_{o}^{2} \varepsilon_{s}^{-\alpha}\right)$, which can be made arbitrarily small by choosing $n$ large and $\varepsilon_{o}$ small enough. This finishes the proof of (6.32).

We now show that the probability that two different edges from $E_{n}\left(\varepsilon_{s}\right)$ are visited during one part of the trajectory is small: that is $\mathbb{P}$-a.s.,

$$
\begin{equation*}
\lim _{n \rightarrow \infty} P_{0}^{\mu}\left[\exists i \leq \varepsilon_{g}^{-2} T, Y[i] \cap E_{n}\left(\varepsilon_{s}\right) \geq 2\right]=o\left(\varepsilon_{g}^{2}\right) \tag{6.45}
\end{equation*}
$$

(6.32) implies that the probability that $Y[i]$ hits at least one edge in $E_{n}\left(\varepsilon_{s}\right)$ is $O\left(\varepsilon_{g}^{2}\right)$. Given that one such edge is visited, one can prove, using the just the global upper bound on the Green function, that the probability that $Y[i]$ visits a second such edge is $O\left(\varepsilon_{g}^{2}\right)$ again. From this (6.45) follows immediately.

Fix $u, v>\varepsilon_{s}$. By the same reasoning as for (6.32) and (6.45), $\mathbb{P}$-a.s.,

$$
\begin{align*}
& \lim _{n \rightarrow \infty} P_{0}^{\mu}\left[\left|Y[i] \cap E_{n}\left(u, v, \varepsilon_{s}^{-1}, \varepsilon_{s}^{-1}\right)\right|=1\right] \\
& =\frac{\varepsilon_{g}^{2}}{\mathcal{C}_{Y}^{2}} \int_{v}^{\varepsilon_{s}^{-1}} \int_{u}^{\varepsilon_{s}^{-1}} \alpha v^{\prime} u^{\prime-\alpha-1} \mathrm{~d} u^{\prime} F_{C}\left(\mathrm{~d} v^{\prime}\right)+o\left(\varepsilon_{g}^{2}\right) . \tag{6.46}
\end{align*}
$$

Combining this formula with Lemma 6.6, we obtain (6.28). (6.27) is then a consequence of (6.32), (6.45) and Bonferroni's inequalities.

To show the asymptotic independence of $s_{n}(i)$ 's, it is sufficient to inspect more carefully the above argument. If we replace the distribution $P_{0}^{\mu}[\cdot]$ by $P_{0}^{\mu}\left[\cdot \mid s_{n}(0) \in\right.$ $\left.A_{0}, \ldots, s_{n}(i-1) \in A_{n-1}\right]$, the only object that changes in the above computations is the distribution of $y_{n}(i)$. However, the only property of this distribution we used is $\sum_{z \in \mathbb{Z}^{d}} P_{0}^{\mu}\left[y_{n}(i)=z\right]=1$, which remains valid when we consider the conditional distribution. Hence, given $s_{n}(0), \ldots, s_{n}(i-1)$, the distribution of $s_{n}(i)$ satisfies (6.27), which implies the asymptotic independence.

As a consequence of Proposition 6.7 we get the following lemma, which we need to show the asymptotic independence of the clock process and the trajectory. We define

$$
\begin{equation*}
r_{n}(i)=n^{-1}\left(y_{n}(i)-y_{n}(i-1)\right) . \tag{6.47}
\end{equation*}
$$

Lemma 6.8 For all $\varepsilon_{g}, \varepsilon_{s}>0, \ell \in \mathbb{N}, \lambda_{1}, \ldots, \lambda_{\ell}>0, \xi_{1}, \ldots, \xi_{\ell} \in \mathbb{R}^{n}$, and $i_{1}<\cdots<i_{\ell} \leq T \varepsilon_{g}^{-2}, \mathbb{P}$-a.s.,

$$
\begin{align*}
& \lim _{n \rightarrow \infty} E_{0}^{\mu}\left[\exp \left\{-\sum_{j=1}^{\ell}\left[\lambda_{j} s_{n}\left(i_{j}\right)+\xi_{j} \cdot r_{n}\left(i_{j}\right)\right]\right\}\right] \\
& =\prod_{j=1}^{\ell}\left[1+\varepsilon_{g}^{2}\left(\frac{\left|\xi_{j}\right|^{2}}{2 d}-c_{\varepsilon_{s}}+G\left(\lambda_{j}\right)\right)+o\left(\varepsilon_{g}^{2}\right)\right], \tag{6.48}
\end{align*}
$$

where

$$
\begin{equation*}
G(\lambda)=G^{\varepsilon_{s}}(\lambda)=\int_{\varepsilon_{s}}^{\varepsilon_{s}^{-1}} \int_{\varepsilon_{s}}^{\varepsilon_{s}^{-1}} \frac{\alpha v^{2} u^{-\alpha-1}}{v+2 u \lambda} \mathrm{~d} u F_{C}(\mathrm{~d} v) . \tag{6.49}
\end{equation*}
$$

Proof In Proposition 6.7 we have proved that $s_{n}(i)$ 's converge as $n \rightarrow \infty$ weakly to an i.i.d. sequence. The same is true for $r_{n}(i)$ 's. Moreover, the same reasoning as at the end of the proof of Proposition 6.7 can be used to show that $s_{n}(i)$ is asymptotically independent of $r_{n}(j)$ for $i \neq j$. This implies the product structure of (6.48). It remains to compute the joint Laplace transform of one pair $\left(s_{n}(i), r_{n}(i)\right)$,

$$
\begin{align*}
& E_{0}^{\mu}\left[e^{-\lambda s_{n}(i)-\xi \cdot r_{n}(i)}\right] \\
& \quad=E_{0}^{\mu}\left[e^{-\xi \cdot r_{n}(i)} \mathbf{1}\left\{s_{n}(i)=0\right\}\right]+E_{0}^{\mu}\left[e^{-\lambda s_{n}(i)-\xi \cdot r_{n}(i)} \mathbf{1}\left\{s_{n}(i) \neq 0\right\}\right] \\
& \quad=E_{0}^{\mu}\left[e^{-\xi \cdot r_{n}(i)} \mathbf{1}\left\{s_{n}(i)=0\right\}\right]+E_{0}^{\mu}\left[e^{-\lambda s_{n}(i)} \mathbf{1}\left\{s_{n}(i) \neq 0\right\}\right] \mathcal{R}(n), \tag{6.50}
\end{align*}
$$

where, since $\left|r_{n}(i)\right| \leq 2 \varepsilon_{g}$, the error term satisfies $e^{-2 \varepsilon_{g}|\xi|} \leq \mathcal{R}(n) \leq e^{2 \varepsilon_{g}|\xi|}$, and thus $\mathcal{R}(n)=1+O\left(\varepsilon_{g}\right)$. The first term in (6.50) can be written as

$$
\begin{aligned}
& E_{0}^{\mu}\left[e^{-\xi \cdot r_{n}(i)}\right]+E_{0}^{\mu}\left[e^{-\xi \cdot r_{n}(i)} \mathbf{1}\left\{s_{n}(i) \neq 0\right\}\right] \\
& \quad=1+\frac{\varepsilon_{g}^{2}|\xi|^{2}}{2 d}+o\left(\varepsilon_{g}^{2}\right)-\mathcal{R}^{\prime}(n) P_{0}^{\mu}\left[s_{n}(i) \neq 0\right],
\end{aligned}
$$

where we used the fact that, by the functional central limit theorem, the distribution of $r_{n}(i)$ converges to the uniform distribution on the sphere with radius $\varepsilon_{g}$. The error term $\mathcal{R}^{\prime}(n)$ has the same asymptotics as $\mathcal{R}(n)$. The second term in (6.50) can be computed using Proposition 6.7. As $n \rightarrow \infty$ it converges to

$$
\begin{equation*}
\varepsilon_{g}^{2} \int_{0}^{\infty} e^{-\lambda x} \nu_{\varepsilon_{s}}(\mathrm{~d} x)=\varepsilon_{g}^{2} G^{\varepsilon_{s}}(\lambda) \tag{6.51}
\end{equation*}
$$

The sum of the last two displays is exactly equal to the square bracket on the right-hand side of (6.48).

## 7 Remaining corrections

In this section we prove that the contribution to the clock process $S_{n}$ of those edges that were not considered in the previous two sections is zero with a high probability. We should control edges that have either $\gamma_{n}(e) \geq \varepsilon_{s}^{-1}$, or $\mu_{e} \geq n^{2 / \alpha} \varepsilon_{s}^{-1}$, or are in the set $\mathcal{B}_{n}\left(\varepsilon_{s}, \varepsilon_{s}^{-1}\right)$ (see (6.2)).

The following lemma treats the first case.
Lemma 7.1 For every $\delta, m>0$ there exists $\varepsilon_{s}>0$ such that, $\mathbb{P}$-a.s., for all but finitely many $n$,

$$
\begin{equation*}
P_{0}^{\mu}\left[\sigma_{\tilde{E}_{n}\left(\varepsilon_{s}, \varepsilon_{s}^{-1}, \infty, \infty\right)}<\tau_{\mathbb{B}}\right] \leq \delta \tag{7.1}
\end{equation*}
$$

Proof By Lemma 6.2(iv), $F_{C}([u, \infty)) \leq \mathbb{P}[\gamma(e) \geq u] \leq c u^{-4}$. This implies that for a large constant $C, \mathbb{P}$-a.s., $\max _{e \in \mathbb{B}} \gamma_{e} \leq C n^{d / 4} \log n$. Let $i_{\text {min }}$ be the largest integer such that $2^{i_{\min }} \leq \varepsilon_{s}^{-1}$ and let $i_{\max }$ be the smallest integer such that $2^{i_{\max }+1} \geq C n^{d / 4} \log n$. By very similar arguments as for (6.41), using the inequality $F_{C}\left(\left[2^{i}, 2^{i+1}\right)\right) \leq c 2^{-4 i}$, one can easily check that, $\mathbb{P}$-a.s., for large $n$, all $x \in \mathbb{B}$ and all $i \in\left\{i_{\min }, \ldots, i_{\max }\right\}$

$$
\begin{equation*}
\left|Q\left(x, \varepsilon_{b} n\right) \cap \tilde{E}\left(\varepsilon_{s}, 2^{i}, \infty, 2^{i+1}\right)\right| \leq C \varepsilon_{b}^{d}\left(\varepsilon_{s}^{-\alpha} n^{d-2} 2^{-4 i}+\log n\right) \tag{7.2}
\end{equation*}
$$

Then arguments analogous to those leading to (6.32) yield that

$$
\begin{align*}
& \sum_{i=i_{\min }}^{i_{\max }} P_{0}^{\mu}\left[\sigma_{\tilde{E}_{n}\left(\varepsilon_{s}, 2^{i}, \infty, 2^{i+1}\right)}<\tau_{\mathbb{B}}\right] \leq \sum_{i=i_{\min }}^{i_{\max }} c n^{2-d} 2^{i} C\left(\varepsilon_{s}^{-\alpha} n^{d-2} 2^{-4 i}+\log n\right) \\
& \leq c \varepsilon_{s}^{3-\alpha}+c^{\prime} n^{2-d+(d / 4)}(\log n)^{2} \tag{7.3}
\end{align*}
$$

which can be made arbitrarily small by choosing $\varepsilon_{s}$ small and $n$ large.
The remaining corrections are treated in the next lemma. Its proof is almost analogous to the previous one (actually simpler), and we do not present it here.

Lemma 7.2 For every $\delta, m>0$ there exists $\varepsilon_{s}>0$ such that, $\mathbb{P}$-a.s., for all but finitely many $n$,

$$
\begin{equation*}
P_{0}^{\mu}\left[\sigma_{\tilde{E}_{n}\left(\varepsilon_{s}^{-1}, \varepsilon_{s}, \infty, \varepsilon_{s}^{-1}\right)}<\tau_{\mathbb{B}}\right] \leq \delta \quad \text { and } \quad P_{0}^{\mu}\left[\sigma_{\mathcal{B}_{n}\left(\varepsilon_{s}, \varepsilon_{s}^{-1}\right)}<\tau_{\mathbb{B}}\right] \leq \delta . \tag{7.4}
\end{equation*}
$$

## 8 Proof of the main theorems

Proof of Theorem 2.1 We first show that $S_{n}(t)$ converges as $n \rightarrow \infty$ to $\mathcal{C}_{S} V$ in the $M_{1}$-topology. To check the convergence of finite dimensional distribution we fix $\ell \in \mathbb{N}$
and $0<s_{1}<\cdots<s_{\ell}<\infty$. By Lemma 6.5, for every $\delta>0$ it is possible to choose $\varepsilon_{g}$ small enough, such that except on a set of probability smaller than $\delta$, for all $j \leq \ell$

$$
\begin{equation*}
t_{n}\left(\left\lfloor(1-\delta) d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s_{j}\right\rfloor\right) \leq n^{2} s_{n}(j) \leq t_{n}\left(\left\lfloor(1+\delta) d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s_{j}\right\rfloor\right) \tag{8.1}
\end{equation*}
$$

Further, it follows from Proposition 5.1 and Lemmas 7.1 and 7.2 that by choosing $\varepsilon_{s}$ small enough, except on a set of probability smaller than $\delta$, for all $i \leq T \varepsilon_{g}^{-2}$ and $n$ large enough,

$$
\begin{equation*}
0 \leq S_{n}\left(t_{n}(i)\right)-\sum_{j=1}^{i} s_{n}(i) \leq \delta . \tag{8.2}
\end{equation*}
$$

Hence, for $u_{1}, \ldots, u_{\ell} \in[0, \infty)$,

$$
\begin{align*}
& \lim _{n \rightarrow \infty} P_{0}^{\mu}\left[S_{n}\left(s_{j}\right) \geq u_{j} \forall j \leq \ell\right]=\lim _{\varepsilon_{s}, \varepsilon_{g} \rightarrow 0} \lim _{n \rightarrow \infty} P_{0}^{\mu}\left[S_{n}\left(s_{j}\right) \geq u_{j} \forall j \leq \ell\right] \\
& \leq \lim _{\varepsilon_{s}, \varepsilon_{g} \rightarrow 0} \lim _{n \rightarrow \infty} P_{0}^{\mu}\left[S_{n}\left(n^{-2} t_{n}\left(\left\lfloor(1+\delta) d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s_{j}\right\rfloor\right)\right) \geq u_{j} \forall j \leq \ell\right] \\
& \leq \lim _{\varepsilon_{s}, \varepsilon_{g} \rightarrow 0} \lim _{n \rightarrow \infty} P_{0}^{\mu}\left[\sum_{i=1}^{\left\lfloor(1+\delta) d C_{Y}^{2} \varepsilon_{g}^{-2} s_{j}\right\rfloor} s_{n}(i) \leq u_{j} \forall j \leq \ell\right] . \tag{8.3}
\end{align*}
$$

The lower bound can be obtained analogously by changing $\delta$ to $-\delta$. The distribution of the sum of $s_{n}(i)$ is given in Proposition 6.7. Sending first $n \rightarrow \infty$ and then $\varepsilon_{g} \rightarrow 0$ we get that the distribution of the process $\sum_{i=1}^{\left\lfloor\cdot \varepsilon_{g}^{-2}\right\rfloor} s_{n}(i)$ converges to a compound Poisson process with intensity measure $\nu_{\varepsilon_{s}}$. As $\varepsilon_{s} \rightarrow 0, \nu_{\varepsilon_{s}}$ converges to the measure

$$
\begin{equation*}
x^{-1-\alpha} \alpha 2^{\alpha} \Gamma(\alpha+1) \mathcal{K}_{C} \mathrm{~d} x, \tag{8.4}
\end{equation*}
$$

where $\mathcal{K}_{C}=\int_{0}^{\infty} v^{1-\alpha} F_{C}(\mathrm{~d} v)$. Therefore, as $n \rightarrow \infty$ and then $\varepsilon_{s}, \varepsilon_{g} \rightarrow 0$, the sum $\sum_{i=1}^{s \varepsilon_{g}^{-2}} s_{n}(i)$ converges (in the sense of finite-dimensional distributions) to a stable subordinator $V^{\prime}$ whose log-Laplace transform at time $s, \log \mathbb{E}\left[e^{-\lambda V^{\prime}(s)}\right]$, is given by

$$
\begin{equation*}
-s \mathcal{K}_{C} \int_{0}^{\infty}\left(1-\varepsilon^{-\lambda x}\right) x^{-1-\alpha} \alpha 2^{\alpha} \Gamma(\alpha+1)=-s \mathcal{K}_{C} 2^{\alpha} \pi \alpha \csc (\pi \alpha) \lambda^{\alpha} \tag{8.5}
\end{equation*}
$$

Putting this together with the estimate (8.3), we obtain that $S_{n}(\cdot)$ converges, in the sense of the finite-dimensional distributions, to $\mathcal{C}_{S} V_{\alpha}(\cdot)$ where

$$
\begin{equation*}
\mathcal{C}_{S}=\left(d \mathcal{C}_{Y}^{2} \mathcal{K}_{C} 2^{\alpha} \pi \alpha \csc (\pi \alpha)\right)^{1 / \alpha} \tag{8.6}
\end{equation*}
$$

To check the asymptotic independence of $Y_{n}$ and $S_{n}$ we use Lemma 6.8. As follows from the previous discussion $Y_{n}(s)$ and $S_{n}(s)$ are well approximated by $\sum_{i=0}^{d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s} r_{n}(i)$ and $\sum_{i=0}^{d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s} s_{n}(i)$. The joint Laplace transform of these two sums

$$
\begin{align*}
& E_{0}^{\mu}\left[\exp \left\{-\sum_{i=0}^{d \mathcal{C}_{Y}^{2} \varepsilon_{g}^{-2} s}\left[\xi \cdot r_{n}(i)+\lambda s_{n}(i)\right]\right\}\right] \xrightarrow{n \rightarrow \infty} \\
& \left(1+\varepsilon_{g}^{2}\left(\frac{|\xi|^{2}}{2 d}-c_{\varepsilon_{s}}+G(\lambda)\right)+o\left(\varepsilon_{g}^{2}\right)\right)^{d C_{Y} \varepsilon_{g}^{-2} s} \tag{8.7}
\end{align*}
$$

which converges, as $\varepsilon_{g} \rightarrow 0$, to

$$
\begin{equation*}
\exp \left\{d \mathcal{C}_{Y} s\left(\frac{|\xi|^{2}}{2 d}-c_{\varepsilon_{s}}+G(\lambda)\right)\right\} \tag{8.8}
\end{equation*}
$$

The same calculation applies for the higher dimensional distributions, which implies that the sums of $s_{n}$ and $r_{n}$ (as processes) are independent in the limit.

We further prove that the sequence $S_{n}$ is tight on $D^{1}\left(M_{1}\right)$, that is on the space $D=D^{1}=D([0, \infty), \mathbb{R})$ equipped with the Skorokhod $M_{1}$-topology. We recall the criterion for the tightness on on this space (see, e.g., [34, Theorem 12.12.3])

Lemma 8.1 (a) The sequence of probability measures $\left\{P_{n}\right\}$ on $D([0, T], \mathbb{R})$ is tight in the $M_{1}$-topology if
(i) For each positive $\varepsilon$ there exists $c$ such that

$$
\begin{equation*}
P_{n}\left[f:\|f\|_{\infty}>c\right] \leq \varepsilon, \quad n \geq 1 \tag{8.9}
\end{equation*}
$$

(ii) For each $\varepsilon>0$ and $\eta>0$, there exist a $\delta, 0<\delta<T$, and an integer $n_{0}$ such that

$$
\begin{equation*}
P_{n}\left[f: w_{f}(\delta) \geq \eta\right] \leq \varepsilon, \quad n \geq n_{0} \tag{8.10}
\end{equation*}
$$

and

$$
\begin{equation*}
P_{n}\left[f: v_{f}(0, \delta) \geq \eta\right] \leq \varepsilon \text { and } P_{n}\left[f: v_{f}(T, \delta) \geq \eta\right] \leq \varepsilon, \quad n \geq n_{0} \tag{8.11}
\end{equation*}
$$

Here, $w_{f}(\delta)$ and $v_{f}(t, \delta)$ stands for

$$
\begin{align*}
w_{f}(\delta) & =\sup \left\{\inf _{\alpha \in[0,1]}\left|f(t)-\alpha f\left(t^{\prime}\right)-(1-\alpha) f\left(t^{\prime \prime}\right)\right|: t^{\prime} \leq t \leq t^{\prime \prime} \leq T, t^{\prime \prime}-t^{\prime} \leq \delta\right\}, \\
v_{f}(t, \delta) & =\sup \left\{\left|f\left(t^{\prime}\right)-f\left(t^{\prime \prime}\right)\right|: t^{\prime}, t^{\prime \prime} \in[0, T] \cap(t-\delta, t+\delta)\right\} . \tag{8.12}
\end{align*}
$$

(b) The sequence of probability measures $\left\{P_{n}\right\}$ on $D^{1}=D([0, \infty), \mathbb{R})$ is tight in the $M_{1}$-topology iffor every $T>0$ its natural projection to $D([0, T], \infty)$ is tight.

Returning back to $S_{n}$, we note that since $S_{n}$ are increasing, the condition (i) of Lemma 8.1 is equivalent to the tightness of $S_{n}(T)$ which can be easily checked from the convergence of finite dimensional distributions. In order to check condition (ii) of Lemma 8.1 remark that for increasing functions the oscillation function $w_{f}(\delta)$ is equal to zero. So checking (ii) reduces to controlling the boundary oscillations $v_{S_{n}}(0, \delta)$ and $v_{S_{n}}(T, \delta)$. For the first quantity (using again the monotonicity of $S_{n}$ ) this amounts to checking that for any $\varepsilon, \eta>0$ there is $\delta$ such that $\mathbb{P}\left[S_{n}(\delta) \geq \eta\right]<\varepsilon$, which follows again from the convergence of the finite-dimensional distributions. The reasoning for $v_{S_{n}}(T, \delta)$ is analogous. The sequence of distributions of $S_{n}$ is thus $M_{1}$-tight on $D([0, T], \mathbb{R})$ for all $T>0$, and therefore, by Lemma 8.1 , on $D^{1}\left(M_{1}\right)$.

The tightness of $Y_{n}$ in $D^{d}(U)$ follows from Proposition 3.1. The tightness of both components implies the tightness of the pair $\left(S_{n}, Y_{n}\right)$ in the product topology on $D^{1}\left(M_{1}\right) \times D^{d}(U)$.

Proof of Theorem 1.2 This follows from the description of $X_{n}$ as $X_{n}(\cdot)=Y_{n}\left(S_{n}^{-1}(\cdot)\right)$. Let $D_{u, \uparrow}$ denote the subset of $D^{1}$ consisting of unbounded increasing functions. By Corollary 13.6.4 of [34] the inverse map from $D_{u, \uparrow}\left(M_{1}\right)$ to $D_{u, \uparrow}(U)$ is continuous at strictly increasing functions. Since the stable subordinator $V_{\alpha}$ (the limit of $S_{n}$ in $\left.\left(D_{u, \uparrow}, M_{1}\right)\right)$ is a.s. strictly increasing, the distribution of $S_{n}^{-1}$ converges to the distribution of $V_{\alpha}^{-1}$ weakly on $D_{u, \uparrow}(U)$ and the limit is a.s. continuous. It is easy to check that the composition $(f, g) \mapsto f \circ g$ as a mapping from $D^{d}(U) \times D_{u, \uparrow}(U)$ to $D^{d}(U)$ is continuous on $C^{d} \times C$ (here $C$ is the space of continuous function). The weak convergence of $X_{n}$ on $D^{d}(U)$ then follows.

## 9 The proof for the BTM

In this Section we discuss the proofs of Theorems 1.3 and 2.2 for the BTM. More precisely, we will explain how the proofs from [18], and from Sects. 3 to 8 should be modified for the BTM.

Let $\tilde{Y}$ be the VSRW associated with the conductances $\tilde{\mu}_{x y}=\tau_{x}^{a} \tau_{y}^{a}$, and let

$$
\begin{equation*}
\tilde{g}_{A}^{\tau}(x, y)=\tilde{E}_{x}^{\tau} \int_{0}^{\tau_{A}} \mathbf{1}\left\{\tilde{Y}_{s}=y\right\} \mathrm{d} t \tag{9.1}
\end{equation*}
$$

be the associated Greens function. The first step is to obtain the FCLT (Proposition 3.1) and the Greens function bounds (Proposition 3.2) for $\tilde{Y}$ and $\tilde{g}_{A}^{\tau}(x, y)$.

Note that the collection $\left(\tilde{\mu}_{e}\right)_{e \in E^{d}}$ is not i.i.d., but that it is still a stationary ergodic process which is bounded below. We can therefore use [18, Theorem 6.1], provided we verify the condition (9.4) on the metric $\tilde{d}(x, y)$ given there. This metric is defined as follows. For all edges $(x y) \in \mathbb{E}^{d}$, let

$$
\begin{equation*}
t(x y)=\min \left\{\tilde{\mu}_{x y}^{-1 / 2}, C\right\} \tag{9.2}
\end{equation*}
$$

for some $C>0$. The new metric $\tilde{d}$ is given by the first-passage percolation distance,

$$
\begin{equation*}
\tilde{d}(x, y)=\inf \left\{\sum_{i=1}^{n} t\left(x_{i-1} x_{i}\right)\right\}, \tag{9.3}
\end{equation*}
$$

where the infimum is taken over all nearest-neighbours paths, $x=x_{0} \sim x_{1} \sim \ldots \sim$ $x_{n}=y$, connecting $x$ and $y$. In the case of the VSRW for the RCM [18] uses firstpassage percolation results from [35] to show that the metric $\tilde{d}$ is equivalent with the Euclidean one, with a high probability. Since these results are only available when $t_{e}$ are i.i.d., we need an equivalent of Lemma 4.2 of [18] to prove the necessary estimates on $\tilde{g}^{\tau}$.

Lemma 9.1 Let $\tilde{B}(x, r)=\{y: \tilde{d}(x, y) \leq r\}$ be the balls in the $\tilde{d}$-metric. Then there exists $c_{1}, \ldots c_{4} \in(0, \infty)$ such that

$$
\begin{equation*}
\tilde{\mathbb{P}}\left[\tilde{B}\left(0, c_{1} r\right) \subset B(0, r) \subset \tilde{B}\left(0, c_{2} r\right)\right] \geq 1-c_{3} e^{-c_{4} r} \tag{9.4}
\end{equation*}
$$

Consequently, the conclusions of $\left[18\right.$, Theorem 6.1] hold for $\tilde{Y}$ and $g_{A}^{\tau}(x, y)$, and in particular Propositions 3.1, 3.2 hold for $\tilde{Y}$ and $\tilde{g}_{A}^{\tau}(x, y)$.

Proof As $t(x y)$ is bounded from above by $C$, there exists $c_{2}$ such that the right hand inclusion always holds. To show the left one define, for some small $\delta>0$ which will be fixed later,

$$
\begin{equation*}
\bar{t}(x y)=\delta \mathbf{1}\{\tilde{t}(x y) \geq \delta\} \tag{9.5}
\end{equation*}
$$

and $\bar{d}(x, y)$ analogously to $\tilde{d}(x, y)$. Obviously $\tilde{d}(x, y) \geq \bar{d}(x, y)$ and thus $\tilde{B}(x, y) \subset$ $\bar{B}(x, y)$, where $\bar{B}$ is ball in the $\bar{d}$ metric. We have $\tilde{\mathbb{P}}(\bar{t}(x y)=\delta)=\psi(\delta)$ with $\psi(\delta) \rightarrow 1$ as $\delta \rightarrow 0$. Moreover, $\bar{t}(x y)$ and $\bar{t}\left(x^{\prime} y^{\prime}\right)$ are independent if $\{x, y\}$ and $\left\{x^{\prime}, y^{\prime}\right\}$ are disjoint. We can now use coupling with independent percolation as in [36]. According to it, it is possible to construct an i.i.d. family $\left(t^{\prime}(e): e \in E^{d}\right)$ of Bernoulli random variables on $\{0, \delta\}$ on the same probability space as $\bar{t}(e)$ such that $\bar{t}(e) \geq t^{\prime}(e)$, $\tilde{\mathbb{P}}$-a.s., and $\tilde{\mathbb{P}}\left[t^{\prime}(e)=\delta\right]=\phi(\psi(\delta))$ and $\phi(u) \rightarrow 1$ as $u \rightarrow 1$. We now fix $\delta$ such that $\phi(\psi(\delta))$ is larger that the percolation threshold on $\mathbb{Z}^{d}$. Since $t^{\prime}(e)$ are independent, by the same argument as in $[18], \mathbb{P}\left[B^{\prime}\left(0, c_{1}^{\prime} r\right) \subset B(0, r)\right] \geq 1-c_{3} e^{-c_{4} r}$, where $B^{\prime}$ is the ball corresponding to the metric $t^{\prime}$. Since $t^{\prime}(e) \leq \bar{t}(e)$ and thus $B^{\prime}(0, r) \supset \bar{B}(0, r)$, the proof is finished.

Finally, we mention the modifications that are necessary in Sects. 3-8. The proof for the BTM is actually simpler, since the 'sites with large equilibrium measure' do not come in pairs as in the CSRW, but are typically isolated.

The only changes in Sect. 3 are the obvious replacement of $\boldsymbol{\mu}$ by $\boldsymbol{\tau}$ and adding tildas everywhere. Section 5 can also be easily adapted after replacing the definition (5.2) by

$$
\begin{equation*}
\tilde{T}_{n}(0, \varepsilon)=\left\{x \in \mathbb{Z}^{d}: \tau_{x} \leq \varepsilon n^{2 / \alpha}\right\} \tag{9.6}
\end{equation*}
$$

Similarly, in Sect. 6 one should define

$$
\begin{align*}
\tilde{T}_{n}(u, w) & =\left\{x \in \mathbb{Z}^{d}: \tau_{x} \in[u, w) n^{2 / \alpha}\right\}, \\
\mathcal{B}_{n}(u, w) & =\left\{x \in \tilde{T}_{n}(u, w): B(x, v) \cap \tilde{T}_{n}\left(n^{-\iota}, \infty\right)>1\right\},  \tag{9.7}\\
T_{n}\left(u, v, w, w^{\prime}\right) & =\left\{x \in T_{n}(u, w): \gamma_{n}(x) \in\left[v, w^{\prime}\right)\right\} .
\end{align*}
$$

It is not necessary to define $E(u, w)$ and $E\left(u, v, w, w^{\prime}\right)$. In Lemma 6.2 only $\gamma_{n}(z)$ should be considered. In the proof of this lemma, in (6.14), the minimum and products should be taken only over even integers to recover the independence. The remaining parts of Sect. 6 are essentially unchanged. Note, however, that the values of the constants in (6.29) and (6.30) will change. Similar changes apply also for Sects. 7 and 8.

Acknowledgments The authors thank G. Ben Arous and T. Kumagai for valuable discussions.

## References

1. Ben Arous, G., Černý, J.: Scaling limit for trap models on $\mathbb{Z}^{d}$. Ann. Probab. 35(6), 2356-2384 (2007)
2. Zaslavsky, G.M.: Chaos, fractional kinetics, and anomalous transport. Phys. Rep. 371, 461-580 (2002)
3. Zaslavsky, G.M.: Hamiltonian Chaos and Fractional Dynamics. Oxford University Press, Oxford (2005)
4. Gorenflo, R., Mainardi, F.: Fractional diffusion processes: probability distributions and continuous time random walk. In: Rangarajan, G., Ding, M. (eds.) Processes with Long Range Correlations, Lecture Notes in Physics, no. 621, pp. 148-166. Springer, Berlin (2003)
5. Hilfer, R. (ed.): Applications of Fractional Calculus in Physics. World Scientific Publishing Co. Inc., River Edge (2000)
6. Metzler, R., Klafter, J.: The random walk's guide to anomalous diffusion: a fractional dynamics approach. Phys. Rep. 339(1), 77 (2000)
7. Montroll, E.W., Shlesinger, M.F.: On the wonderful world of random walks. In: Nonequilibrium Phenomena, II. Stud. Stat. Mech. XI, pp. 1-121. North-Holland, Amsterdam (1984)
8. Shlesinger, M.F., Zaslavsky, G.M., Klafter, J.: Strange kinetics. Nature 363, 31-37 (1993)
9. Kipnis, C., Varadhan, S.R.S.: Central limit theorem for additive functionals of reversible Markov processes and applications to simple exclusions. Commun. Math. Phys. 104(1), 1-19 (1986)
10. Kozlov, S.M.: The averaging method and walks in inhomogeneous environments. Uspekhi Mat. Nauk 40(2(242)), 61-120, 238 (1985)
11. De Masi, A., Ferrari, P.A., Goldstein, S., Wick, W.D.: An invariance principle for reversible Markov processes. Applications to random motions in random environments. J. Stat. Phys. 55(3-4), 787855 (1989)
12. Sidoravicius, V., Sznitman, A.-S.: Quenched invariance principles for walks on clusters of percolation or among random conductances. Probab. Theory Relat. Fields 129(2), 219-244 (2004)
13. Biskup, M., Prescott, T.M.: Functional CLT for random walk among bounded random conductances. Electron. J. Probab. 12(49), 1323-1348 (electronic, 2007)
14. Mathieu, P.: Quenched invariance principles for random walks with random conductances. J. Stat. Phys. 130(5), 1025-1046 (2008)
15. Barlow, M.T.: Random walks on supercritical percolation clusters. Ann. Probab. 32(4), 30243084 (2004)
16. Berger, N., Biskup, M.: Quenched invariance principle for simple random walk on percolation clusters. Probab. Theory Relat. Fields 137(1-2), 83-120 (2007)
17. Mathieu, P., Piatnitski, A.: Quenched invariance principles for random walks on percolation clusters. Proc. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci. 463(2085), 2287-2307 (2007)
18. Barlow, M.T., Deuschel, J.-D.: Invariance principle for the random conductance model with unbounded conductances. Ann. Probab. (2009, to appear)
19. Bouchaud, J.-P.: Weak ergodicity breaking and aging in disordered systems. J. Phys. I (France) 2, 1705-1713 (1992)
20. Bouchaud, J.-P., Dean, D.S.: Aging on Parisi's tree. J. Phys. I (France) 5, 265 (1995)
21. Rinn, B., Maass, P., Bouchaud, J.-P.: Hopping in the glass configuration space: subaging and generalized scaling laws. Phys. Rev. B 64, 104417 (2001)
22. Ben Arous, G., Černý, J.: Bouchaud's model exhibits two aging regimes in dimension one. Ann. Appl. Probab. 15(2), 1161-1192 (2005)
23. Ben Arous, G., Bovier, A., Gayrard, V.: Glauber dynamics of the random energy model. I. Metastable motion on the extreme states. Commun. Math. Phys. 235(3), 379-425 (2003)
24. Fontes, L.R.G., Isopi, M., Newman, C.M.: Random walks with strongly inhomogeneous rates and singular diffusions: convergence, localization and aging in one dimension. Ann. Probab. 30(2), 579604 (2002)
25. Ben Arous, G., Černý, J., Mountford, T.: Aging in two-dimensional Bouchaud's model. Probab. Theory Relat. Fields 134(1), 1-43 (2006)
26. Cerný, J.: On Two Properties of Strongly Disordered Systems, Aging and Critical Path Analysis. Ph.D. thesis, EPF Lausanne (2003)
27. Ben Arous, G., Černý, J.: The arcsine law as a universal aging scheme for trap models. Commun. Pure Appl. Math. 61(3), 289-329 (2008)
28. Barlow, M.T., Zheng, X.: The random conductance model with Cauchy tails. Ann. Appl. Probab. (2009, to appear)
29. Barlow, M.T.: Aspects of first passage percolation. In: École d'été de Probabilités de Saint-Flour, XXV-1995. Lecture Notes in Mathematics, vol. 1690, pp. 1-121. Springer, Berlin (1998)
30. Kigami, K.: Diffusions on Fractals. Cambridge University Press, Cambridge (2001)
31. Saichev, A.I., Zaslavsky, G.M.: Fractional kinetic equations: solutions and applications. Chaos 7(4), 753-764 (1997)
32. Meerschaert, M.M., Scheffler, H.-P.: Limit theorems for continuous-time random walks with infinite mean waiting times. J. Appl. Probab. 41(3), 623-638 (2004)
33. Barlow, M.T., Hambly, B.M.: Parabolic Harnack inequality and local limit theorem for percolation clusters. Electron. J. Probab. 14, 1-27 (2009)
34. Whitt, W.: Stochastic-process limits: an introduction to stochastic-process limits and their application to queues. In: Springer Series in Operations Research. Springer, New York (2002)
35. Kesten, H.: Aspects of first passage percolation. In: École d'été de Probabilités de Saint-Flour, XIV1984, Lecture Notes in Mathematics, vol. 1180, pp. 125-264. Springer, Berlin (1986)
36. Liggett, T.M., Schonmann, R.H., Stacey, A.M.: Domination by product measures. Ann. Probab. 25(1), 71-95 (1997)

[^0]:    Research partially supported by NSERC (Canada).
    M. T. Barlow

    Department of Mathematics, University of British Columbia, Vancouver, BC, V6T 1Z2, Canada
    e-mail: barlow@math.ubc.ca
    J. Černý ( $\boxtimes$ )

    Department of Mathematics, ETH Zürich, Rämistr. 101, 8092 Zurich, Switzerland
    e-mail: jiri.cerny@math.ethz.ch

