

Muon Spin Rotation Study of the Intercalated Graphite Superconductor CaC₆ at Low Temperatures

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Abstract Muon spin rotation (μ SR) experiments were performed on the intercalated graphite CaC₆ in the normal and superconducting state down to 20 mK. In addition, AC magnetization measurements were carried out resulting in an anisotropic upper critical field H_{c2} , from which the coherence lengths $\xi_{ab}(0) = 36.3(1.5)$ nm and $\xi_c(0) = 4.3(7)$ nm were estimated. The anisotropy parameter $\gamma_H = H_{c2}^{ab}/H_{c2}^c$ increases monotonically with decreasing temperature. A single isotropic s -wave description of superconductivity cannot account for this behavior. From magnetic field dependent μ SR experiments, the absolute value of the in-plane magnetic penetration depth $\lambda_{ab}(0) = 78(3)$ nm was determined. The temperature dependence of the superfluid density $\rho_s(T)$ is slightly better described by a two-gap than a single-gap model.

Keywords Superconducting materials · Penetration depth · Muon spin rotation and relaxation

1 Introduction

The field of graphite intercalation compounds (GICs) gained attention after the discovery of the superconductor CaC₆

with a rather high value of the superconducting transition temperature $T_c \simeq 11.5$ K [1]. Superconductivity in GICs was first reported in the potassium-graphite compound KC₈ with $T_c \simeq 0.14$ K [2]. Until the discovery of the superconductor CaC₆, the highest T_c was observed in KTi_{1.5}C₅ with $T_c \simeq 2.7$ K, synthesized under ambient pressure [3]. However, high pressure synthesis was found to increase T_c up to 5 K in metastable compounds such as NaC₂ and KC₃ [4, 5]. According to experimental and theoretical work, CaC₆ can be described as a classical BCS superconductor with a single isotropic gap $\Delta_0 \simeq 1.7$ meV [6, 7]. The upper critical field H_{c2} shows a remarkable anisotropy with zero temperature values $\mu_0 H_{c2}^c(0) \simeq 0.3$ T and $\mu_0 H_{c2}^{ab}(0) \simeq 1.9$ T, for the external field H parallel to the c -axis or parallel to the ab -plane, respectively [8]. A recent ARPES study indicated the existence of a possible second superconducting gap [9] with a small zero temperature value $\Delta_{0,2} \simeq 0.2(2)$ meV. Tunneling experiments [10] gave strong indications for the existence of a superconducting anisotropic s -wave gap in CaC₆, providing different zero temperature gap values for changing current injection directions, $\Delta_0^c \simeq 1.7$ meV and $\Delta_0^{ab} \simeq 1.44$ meV. The muon-spin rotation (μ SR) technique is a powerful method to characterize the superconducting gap symmetry in superconductors [11]. μ SR measurements down to very low temperatures may allow to access a possible second small gap, which should manifest itself in the low temperature superfluid density in terms of an inflection point [11]. A recent μ SR study [12] supports a single gap isotropic s -wave description of superconductivity in CaC₆, although the low temperature region was not studied. In this work, we present extended measurements down to 20 mK and investigate possible order parameter symmetries (single-gap isotropic s -wave, single-gap anisotropic s -wave, two-gap isotropic s -wave) by means of μ SR experiments.

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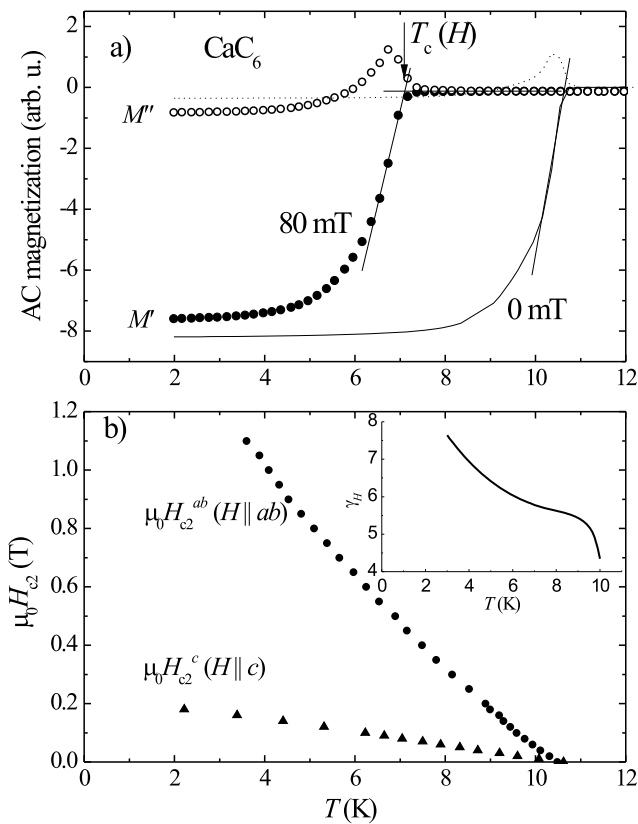


Fig. 1 (a) AC magnetization curves of CaC₆ taken at $\mu_0 H = 80$ mT and 0 mT. M' and M'' denote the real and the imaginary part of the AC magnetization, respectively. The *arrow* indicates the onset criterion to determine $T_c(H)$ (crossing of the two lines). (b) Temperature dependence of the upper critical fields H_{c2}^c and H_{c2}^{ab} for the external magnetic field $H\parallel c$ and $H\parallel ab$, respectively, as determined by AC magnetization. The *inset* shows the calculated upper critical field anisotropy $\gamma_H = H_{c2}^{ab}/H_{c2}^c$ as a function of temperature

2 Experimental Results

CaC₆ samples were prepared from Highly Oriented Pyrolytic Graphite (HOPG, Structure Probe Inc., SPI-2 Grade) and calcium metal (Sigma Aldrich, 99.99% purity) by means of the molten alloy method. The detailed preparation procedure is described elsewhere [13]. In order to reduce the melting point of the calcium, lithium (Sigma Aldrich, 99.9% purity) was added. The atomic ratio 3:1 of calcium and lithium was measured in a stainless steel reactor under argon atmosphere. The molten alloy was cleaned in the glove-box. After immersing the graphite pieces, the reactor was tightly sealed. Intercalation at 350 °C was carried out in an industrial furnace for 30 days. The resulting sample's color is gold and the surface possesses a shiny metallic lustre. The final sample is a superconductor with an onset temperature $T_c \simeq 11.2$ K, in reasonable agreement with earlier studies [1, 6, 7].

To characterize our sample, we carried out detailed AC magnetization measurements presented in Fig. 1. The AC

component of the applied magnetic field was $\mu_0 H = 1$ mT at $\nu = 1$ kHz. Figure 1a shows typical AC magnetization curves at 80 mT (real $M'(T)$ and imaginary $M''(T)$ part), and the straight lines indicate the onset criterion to determine T_c . The temperature dependence of the upper critical field components $H_{c2}^c (H\parallel c)$ and $H_{c2}^{ab} (H\parallel ab)$ are displayed in Fig. 1b. The values of the H_{c2}^c and H_{c2}^{ab} were determined from AC magnetization using the procedure illustrated in Fig. 1a. The upper critical field is anisotropic and follows the temperature dependence as reported in [8] and [14]. In the framework of the anisotropic Ginzburg–Landau theory, the upper critical field for $H\parallel c$ and $H\parallel ab$ is given by

$$\begin{aligned} \mu_0 H_{c2}^{ab}(0) &= \frac{\Phi_0}{2\pi \xi_{ab}(0) \xi_c(0)}, \\ \mu_0 H_{c2}^c(0) &= \frac{\Phi_0}{2\pi \xi_{ab}^2(0)}, \end{aligned} \quad (1)$$

where $\Phi_0 = h/2e = 2.07 \times 10^{-15}$ Tm² is the flux quantum, and $\xi_{ab,c}$ are the corresponding coherence length components at zero temperature. With the linearly extrapolated values of the upper critical field at zero temperature, $\mu_0 H_{c2}^{ab}(0) = 2.1(3)$ T and $\mu_0 H_{c2}^c(0) = 0.25(2)$ T, and in the equation above we obtain $\xi_{ab}(0) = 36.3(1.5)$ nm and $\xi_c(0) = 4.3(7)$ nm. These values are comparable to earlier reports such as resistivity measurements ($\xi_{ab}(0) = 29.0$ nm and $\xi_c(0) = 5.7$ nm) [8], or susceptibility studies ($\xi_{ab}(0) = 35$ nm and $\xi_c(0) = 13$ nm) [15]. Remarkably, the upper critical field H_{c2}^{ab} shows a positive curvature in the studied temperature region which was not observed in previous investigations. A similar positive curvature was observed near T_c in MgB₂ [16] where it was explained by a two-gap model. However, we should keep in mind that other explanations are also possible. Complementary experiments, such as studies of the superfluid density, as performed in this work, are necessary to draw definite conclusions. The temperature dependence of the upper critical field anisotropy $\gamma_H = H_{c2}^{ab}/H_{c2}^c$ is shown in the inset of Fig. 1b. We interpolated the measured upper critical field values with a third order polynomial and determined the ratio γ_H . The anisotropy parameter γ_H increases monotonically with decreasing temperature, showing a similar temperature dependence as with MgB₂ [17]. Note that a single gap isotropic *s*-wave description of superconductivity cannot account for this behavior of $\gamma_H(T)$.

μ SR experiments were carried out at the π M3 beamline using the GPS (General Purpose Spectrometer) and the LTF (Low Temperature Facility) spectrometers at the Paul Scherrer Institute PSI Villigen, Switzerland. Six pieces of CaC₆ forming an area of $\simeq 10 \times 14$ mm² were used in the experiment. In the intercalation process, the Ca atoms penetrate from the side of the graphite sample and diffuse along the *ab*-plane. Reducing the sample size in the *ab*-direction favors Ca diffusion, and consequently leads to a higher sample

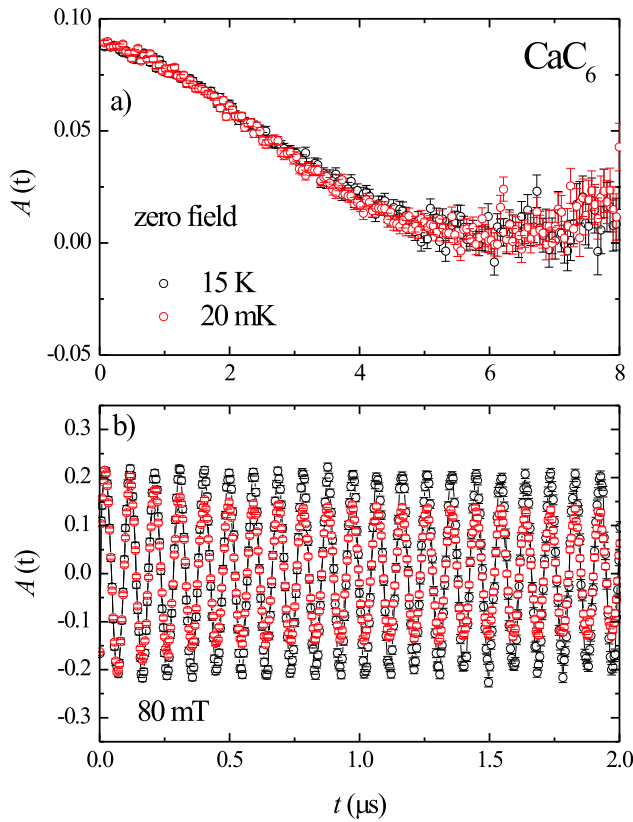


Fig. 2 (Color online) Measured μ SR asymmetries $A(t)$ in CaC_6 . (a) Asymmetry in zero magnetic field (ZF). (b) Asymmetry in 80 mT transverse field (TF). The black data points were taken at 15 K (well above T_c) and the red ones at 20 mK

quality than just using one large piece. For the LTF experiments, the samples were glued on a silver plate with the help of Apiezon N grease and covered by silver coated polyester foil. In contrast, for the GPS experiments the samples were supported between the two arms of a specially designed fork by means of kapton tape. This technique allows only the incoming muons that stop in the sample to be counted and, therefore, the background signal is reduced. The magnetic field ($\mu_0 H = 80$ mT) was applied along the crystallographic c -direction at 15 K well above T_c . Then the sample was cooled down to the lowest temperature (20 mK for the LTF spectrometer), and μ SR spectra were recorded with increasing temperature. The same procedure was applied in the GPS spectrometer, down to the lowest temperature of ≈ 1.6 K available in this device. The initial spin polarization of the implanted muons is perpendicular to the applied magnetic field H . The μ SR experiments were performed at $H > H_{c1}^c$, i.e., the sample was in the vortex state. Magnetic field dependent measurements were performed in the GPS spectrometer. At each magnetic field H , the field was set at a sample-temperature of 15 K (well above T_c), and a μ SR time spectrum was recorded, then it was cooled down to 1.6 K, and another μ SR time spectrum was taken. Before

and after the μ SR experiments, the same piece of CaC_6 was measured in the AC susceptometer, where no difference was observed by comparing the two measurements.

In the vortex state of a type II superconductor, the muons probe the inhomogeneous local magnetic field distribution $P(B)$ due to the vortex lattice [18]. From the second moment $\langle \Delta B^2 \rangle$ of $P(B)$, the in-plane magnetic penetration depth λ_{ab} can be extracted according to the relation $\langle \Delta B^2 \rangle = C \lambda_{ab}^{-2}$ (C is a field dependent quantity) from which the superfluid density $\rho_s \propto \lambda_{ab}^{-2}$ can be extracted [19]. Thus, μ SR provides a direct method to measure the superfluid density in the mixed state of a type II superconductor. Figure 2a shows the μ SR asymmetries $A(t)$ recorded in LTF in zero magnetic field above T_c and at 20 mK. The two time spectra overlap, indicating the same magnetic ground state at both temperatures. The time dependence of the asymmetry can be described by the static Kubo–Toyabe formula, showing only the presence of randomly distributed nuclear magnetic moments, but no realization of magnetism. Figure 2b shows the μ SR asymmetries taken at 15 K and at 20 mK in a field of $\mu_0 H = 80$ mT. A clear damping of the μ SR signal at 20 mK is visible, which can be ascribed to the presence of the vortex lattice. For a detailed description of μ SR studies of the vortex state in type II superconductors, see, e.g., [18]. The μ SR time spectra were fitted to the following expression:

$$A^{\text{TF}}(t) = A_{\text{SC}} e^{(\sigma_{\text{SC}} t)^2 / 2} \cos(\gamma_{\mu} B_{\text{SC}} t + \phi) + A_{\text{bg}} e^{(\sigma_{\text{bg}} t)^2 / 2} \cos(\gamma_{\mu} B_{\text{bg}} t + \phi) + A_{\text{Ag}} \cos(\gamma_{\mu} B_{\text{Ag}} t + \phi) \tag{2}$$

Here, the indices SC, bg and Ag denote the sample (superconductor), the background arising from possible non-superconducting parts of the sample and the nonrelaxing silver background, respectively. A denotes the initial asymmetry, σ is the Gaussian relaxation rate, $\gamma_{\mu} = 2\pi \cdot 135.5342$ MHz/T is the muon gyromagnetic ratio, B is the average magnetic field at the μ^+ stopping site, and ϕ is the initial phase of the muon-spin ensemble. A set of μ SR data were fitted simultaneously with A_{SC} , A_{bg} , A_{Ag} , B_{bg} , B_{Ag} , σ_{bg} , and ϕ as common parameters, and B_{SC} and σ_{SC} as free parameters for each temperature. The same expression was used to analyse the data recorded with the GPS spectrometer (here no silver background signal was present).

The analysis of the field dependence of the measured μ SR relaxation rate $\sigma_{\text{SC}} \propto \langle \Delta B^2 \rangle$ allows to estimate the absolute value of the in-plane magnetic penetration depth λ_{ab} , assuming that CaC_6 can be described as a single isotropic s -wave gap superconductor. The measured values of σ_{SC} are plotted in Fig. 3. To analyze our data, we used the formula

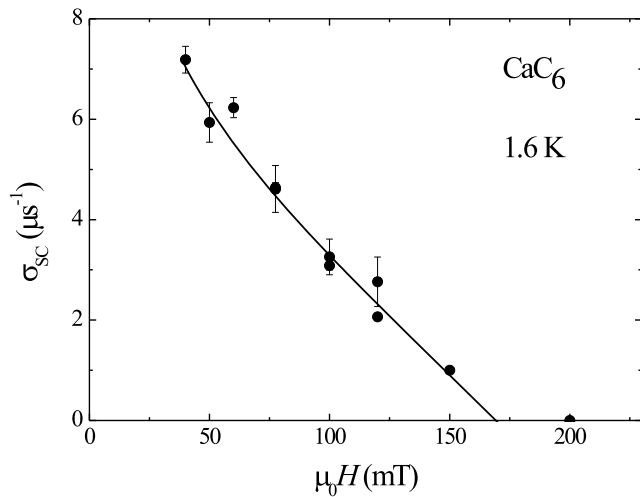


Fig. 3 Field dependence of the Gaussian relaxation rate σ_{SC} in CaC_6 measured at $T = 1.6$ K. The *solid line* corresponds to the calculated curve using (3) (see text)

developed by Brandt [19]:

$$\sigma_{SC} \simeq 0.172\gamma\mu \frac{\Phi_0}{2\pi} \left(1 - \frac{B}{B_{c2}^c}\right) \times \left[1 + 1.21 \left(1 - \sqrt{\frac{B}{B_{c2}^c}}\right)^3\right] \lambda_{ab}^{-2}. \quad (3)$$

Although the formula was derived for high κ type II superconductors ($\kappa \geq 5$), it was suggested to work also for low values of κ at magnetic fields $\sqrt{B/B_{c2}^c} \geq 0.5$ [12]. From the measured μ SR relaxation rates σ_{SC} and (3), one obtains $\lambda_{ab}(0) = 78(3)$ nm and $B_{c2}^c(0) = 170(5)$ mT. The solid line in Fig. 3 represents the corresponding fit (the point at 200 mT was not included in the fit). To estimate the Ginzburg–Landau parameter κ , we use our values of $\xi_{ab}(0) = 36.3(1.5)$ nm and $\lambda_{ab}(0) = 78(3)$ nm giving $\kappa \approx 2.1$. It is worth to mention that the value of B_{c2}^c is smaller than that estimated from the AC magnetization measurements ($B_{c2}^c(0) \simeq 250$ mT).

It is generally assumed that the superfluid density ρ_S is related to the in-plane magnetic penetration depth λ_{ab} by the simple relation $\rho_S \propto \lambda_{ab}^{-2}$. However, for magnetic fields close to B_{c2} this proportionality must be corrected because the order parameter $\psi(r)$ is reduced due to the overlapping vortices. In this case, the spatial average of the superfluid density is given by [12]:

$$\rho_S(T) \simeq \langle |\psi(r)|^2 \rangle \lambda^{-2}(T) \simeq (1 - B/B_{c2}^c(T)) \lambda_{ab}^{-2}(T), \quad (4)$$

where $\langle \dots \rangle$ means the spatial average. Values of λ_{ab}^{-2} were calculated using (3) and for the values of B_{c2}^c we used the results of our AC magnetization measurements (see Fig. 1b). The temperature dependence of the superfluid density $\rho_S(T)$

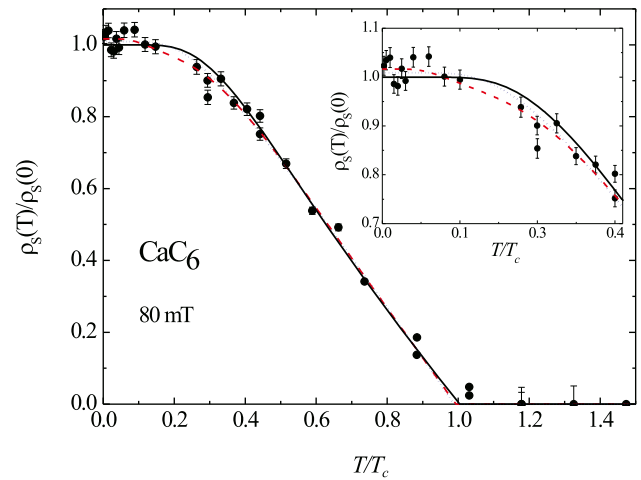


Fig. 4 (Color online) Temperature dependence of the normalized superfluid density $\rho_S(T)/\rho_S(0)$ in CaC_6 measured at 80 mT. The *inset* shows the low temperature region. The *black line* corresponds to a single-gap isotropic *s*-wave, the *blue, dotted line* to a single-gap anisotropic *s*-wave, and the *red, dashed line* to a two-gap isotropic *s*-wave description of $\rho_S(T)$. For a detailed explanation, see text

determined at 80 mT is plotted in Fig. 4. We analyzed $\rho_S(T)$ using three different models: (i) single-gap isotropic *s*-wave (ii) single-gap anisotropic *s*-wave, and (iii) two-gap isotropic *s*-wave. To calculate the temperature dependence of the magnetic penetration depth within the local approximation ($\lambda \gg \xi$), we applied the following formula [20, 21]:

$$\rho_S(T) = \rho_S(0) \times \left(1 + \frac{1}{\pi} \int_0^{2\pi} \int_{\Delta(T,\phi)}^{\infty} \left(\frac{\partial f}{\partial E}\right) \frac{E dE d\phi}{\sqrt{E^2 - \Delta(T,\phi)^2}}\right), \quad (5)$$

where $\rho_S(0)$ is the zero temperature value of the superfluid density, $f = [1 + \exp((E - E_F)/k_B T)]^{-1}$ is the Fermi function, ϕ is the angle along the Fermi surface, and $\Delta(T, \phi) = \Delta_0 h(T/T_c) g(\phi)$. The temperature dependence of the gap is expressed by $h(T/T_c) = \tanh(1.82[1.018(T_c/T - 1)]^{0.51})$ [22]. For the isotropic *s*-wave gap (i) $g^s(\phi) = 1$ and for the anisotropic *s*-wave gap (ii) $g^{sAn}(\phi) = (1 + a \cos 4\phi)/(1 + a)$, where a denotes the anisotropy of the gap. The two-gap fit was calculated using the so called α -model and assuming that the total superfluid density is the sum of the two components [21, 22]:

$$\rho_S(T) = \rho_S(0) \left(\omega \frac{\rho(T, \Delta_{0,1})}{\rho(0, \Delta_{0,1})} + (1 - \omega) \frac{\rho(T, \Delta_{0,2})}{\rho(0, \Delta_{0,2})} \right). \quad (6)$$

Here, $\Delta_{0,1}$ and $\Delta_{0,2}$ are the zero temperature values of the larger and the smaller gap, and ω ($0 \leq \omega \leq 1$) is a weighting factor representing the relative contribution of the larger gap to ρ_S . The results obtained for the different models are plotted in Fig. 4. The black line represents the result for

the single s -wave gap model (i) with $\rho_s(0) = 96.5(5) \mu\text{m}^{-2}$, $\Delta_0 = 0.79(1) \text{ meV}$, $T_c = 6.80(2) \text{ K}$, and $2\Delta_0/k_B T_c = 2.70(5)$, describing CaC_6 as weak coupling BCS superconductor. The reduction of T_c agrees well with our AC magnetization measurements for 80 mT. The blue, dotted line represents the anisotropic s -wave gap analysis (ii) with $\Delta_0 = 0.81(1) \text{ meV}$ and $a \simeq 0.37$ for the gap anisotropy. These values do not agree with the reported ones from tunneling experiments ($\Delta_0^c \simeq 1.7 \text{ meV}$ and $\Delta_0^{ab} \simeq 1.44 \text{ meV}$) [10]. The two-gap analysis (iii) (red, dashed line) yields $\Delta_{0,1} = 0.85(2) \text{ meV}$ for the larger gap and $\Delta_{0,2} = 0.23(1) \text{ meV}$ for the smaller one. Only $\sim 9.4(3)\%$ of the superfluid density is associated with the smaller gap. The smaller gap agrees with $\Delta_{0,2} = 0.2(2) \text{ meV}$ reported in [9]. Our larger gap value is significantly smaller ($\Delta_{0,1} = 1.9(2) \text{ meV}$). Although the best agreement is found for the two-gap scenario, no final conclusion can be drawn from this analysis. All three scenarios describe the observed temperature dependence of the superfluid density. For a direct comparison with [12], we omitted the values of $\rho_s(T)$ below 2 K. In this case, the single s -wave gap fit yields $\rho_s(0) = 91.6(9) \mu\text{m}^{-2}$, $\Delta_0 = 0.84(1) \text{ meV}$, $T_c = 6.75(2) \text{ K}$, and $2\Delta_0/k_B T_c = 2.90(5)$. The values presented in [12] are: $\rho_s(0) = 85.1(3) \mu\text{m}^{-2}$, $\Delta_0 = 0.868(5) \text{ meV}$, and $2\Delta_0/k_B T_c = 3.6(1)$. The higher applied magnetic field (120 mT in [12]) may account for the differences. Here, one should point out that our low temperature data are essential to determine the zero temperature values of ρ_s and Δ .

3 Conclusions

In conclusion, we carried out extensive AC magnetization measurements to map the temperature dependence of the upper critical field H_{c2}^{ab} and H_{c2}^c in CaC_6 and to evaluate the coherence length. The values of $\xi_{ab}(0) = 36.3(1.5) \text{ nm}$ and $\xi_c(0) = 4.3(7) \text{ nm}$ are in good agreement with those reported previously [8, 15]. We found that the upper critical field anisotropy $\gamma_H = H_{c2}^{ab}/H_{c2}^c$ increases with decreasing temperature. From magnetic field dependent μSR experiments, the absolute value of the in-plane magnetic penetration depth was determined to be $\lambda_{ab} \simeq 78(3) \text{ nm}$, in agreement with previously reported values [7, 12]. Furthermore, low temperature μSR experiments were performed in order to map the whole temperature dependence of the superfluid density $\rho_s(T)$. We analyzed the temperature dependence of ρ_s with three different models: (i) single-gap isotropic s -wave, (ii) single-gap anisotropic s -wave, and (iii) two-gap isotropic s -wave. All models describe the measured μSR data almost equally well, although a slightly better agreement was achieved using the two-gap model.

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