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Remarks on simple interpolation between Jordanian twists

Stjepan Meljanac^{*1}, Daniel Meljanac^{†1}, Anna Pachol^{‡2} and Danijel Pikutić^{§1}

¹Division of Theoretical Physics, Ruđer Bošković Institute, Bijenička c.54,
HR-10002 Zagreb, Croatia

²Queen Mary, University of London, School of Mathematics, Mile End Rd.,
London E1 4NS, UK.

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Abstract

In this paper, we propose a simple generalization of the locally r -symmetric Jordanian twist, resulting in the one-parameter family of Jordanian twists. All the proposed twists differ by the coboundary twists and produce the same Jordanian deformation of the corresponding Lie algebra. They all provide the κ -Minkowski spacetime commutation relations. Constructions from noncommutative coordinates to the star product and coproduct, and from the star product to the coproduct and the twist are presented. The corresponding twist in the Hopf algebroid approach is given. Our results are presented symbolically by a diagram relating all of the possible constructions.

1 Introduction

In the Hopf algebras framework it is known that the given Hopf algebra $\mathcal{H}(\mu, \Delta, \epsilon, S)$ can be deformed by using the (Drinfeld) twist $\mathcal{F} \in \mathcal{H} \otimes \mathcal{H}$ [1] which allows to deform the coproduct and an antipode map in such a way that the compatibility conditions are still satisfied and one gets a new Hopf algebra with deformed maps. One special example of such twist is the so-called Jordanian twist $\mathcal{F}_0 = \exp(\ln(1 + \alpha E) \otimes H)$ with α as the deformation parameter ¹. It was firstly constructed by O. Ogievetsky in [2] while studying the moduli space of Hopf structures on the Borel subalgebra of $sl(2) \ni \{H, E | [H, E] = E\}$. This twist provides the simplest example of the triangular deformation of $sl(2)$ algebra with the quantum R-matrix given by $R = \tilde{\mathcal{F}}\mathcal{F}^{-1} = 1 \otimes 1 + r + \mathcal{O}(\alpha^2)$, where $\tilde{\mathcal{F}} = \tau \circ \mathcal{F}$ is the flipped twist (with the flip map: $\tau : c \otimes d \rightarrow d \otimes c$) and r is denoting the classical r -matrix.

*meljanac@irb.hr

†Daniel.Meljanac@irb.hr

‡a.pachol@qmul.ac.uk

§dpikutic@irb.hr

¹To distinguish different Jordanian twists we introduce a notation with sub-indices 0, 1, or 1/2 which will become clear in section 2.

One can notice that \mathcal{F}_0 is not r -symmetric [3] in the sense that the term of the expansion in the deformation parameter α , at the first order, is not given by the full classical r -matrix r , i.e. is not of the form: $\mathcal{F} = 1 \otimes 1 + \frac{1}{2}r + \mathcal{O}(\alpha^2)$. However it can be symmetrized by the so called coboundary twist $\mathcal{F}_\omega = \omega^{-1} \otimes \omega^{-1} \Delta(\omega)$ as:

$$\mathcal{F}_r^{(\omega)} := \omega^{-1} \otimes \omega^{-1} \mathcal{F}_0 \Delta(\omega),$$

where $\omega = \sqrt{\mu[(1 \otimes S)\mathcal{F}_0]}$. The change of the twist by the coboundary twist \mathcal{F}_ω does not provide a new Hopf algebra. It provides a new presentation in the co-algebraic sector only, i.e. new form of the coproduct $\tilde{\Delta} = \mathcal{F}_\omega \Delta \mathcal{F}_\omega^{-1}$.

One of the symmetrized versions (i.e. ‘‘locally r -symmetric’’) of the Jordanian twist $\tilde{\mathcal{F}} = \tau \circ \mathcal{F}_0$ known in the literature is due to V. N. Tolstoy [3] which appeared in the context of studying quantum deformations of relativistic symmetries. The locally r -symmetric version of the Jordanian twist calculated therein has the form:

$$\mathcal{F}_{1/2} = e^{\frac{\alpha}{2}(HE \otimes 1 + 1 \otimes HE)} e^{H \otimes \ln(1 + \alpha E)} e^{-\frac{\alpha}{2}(HE \otimes 1 + H \otimes E + E \otimes H + 1 \otimes HE)} \quad (1)$$

Here $\omega = \exp(\frac{1}{2}\alpha HE)$ was used to construct the coboundary twist. There exists another version of symmetrized \mathcal{F}_0 and it is due to Giaquinto and Zhang [4], but it will not be the object of our study here.

The Jordanian twist investigated, firstly in [2] and then in [3], reappeared in the context of the so-called κ -Minkowski noncommutative spacetime [5, 6, 7], which is an algebra of coordinate functions equipped in the noncommutative star-product leading to the following commutation relations $[\hat{x}^\mu, \hat{x}^\nu] = \frac{i}{\kappa}(v^\mu \hat{x}^\nu - v^\nu \hat{x}^\mu)$ where the deformation parameter is $\frac{1}{\kappa}$ and v^μ is the vector on Minkowski spacetime $\mathcal{M}_{1,n-1}$ in n -dimensions such that $v^2 \in \{-1, 0, 1\}$. In physical applications κ is usually interpreted as the Planck mass or Quantum Gravity scale. The natural quantum symmetry of this noncommutative spacetime is the κ -Poincaré quantum group [8] and it constitutes one of the examples of deformed relativistic spacetime symmetries. In this paper we are interested in the relation with the symmetry of the κ -Minkowski spacetime therefore we shall work with the generators of relativistic symmetries: dilatation D and momenta p_α (instead of the generators H and E of the $sl(2)$ algebra) satisfying the same commutation relation, i.e. $[p_\alpha, D] = p_\alpha$ ². The dilatation generator D is included in the minimal extension of the relativistic spacetime symmetry, the so-called Poincaré-Weyl symmetry: $\{M_{\mu\nu}, p_\mu, D\}$ as well as in the conformal algebra $\{M_{\mu\nu}, p_\mu, D, K_\mu\}$, therefore the Jordanian twists have the support in both of these algebras and can be used in their deformations [5], [9], respectively. However for the presentation of our results it is enough to consider the Lie algebra g generated only by the dilatation and momenta operators satisfying $[p_\alpha, D] = p_\alpha$ and $[p_\mu, p_\nu] = 0$. Recently Jordanian twists also have been considered in application to gravitational theory. For example, in [11], the Jordanian twist was used to construct the noncommutative differential calculus providing metrics as the solutions of non-vacuum Einstein equations, including cosmological constant or spatial curvature cases. Additionally, it can be shown [10] that the star product, reproducing the κ -Minkowski Lie algebra, obtained by suitably reducing the so-called Wick-Voros star product, is in fact the star product from the Jordanian twist \mathcal{F}_0 . Jordanian deformations also have become popular in the context of applications in AdS/CFT correspondence [12].

In this paper we are interested in the generalization of the symmetrized version of Jordanian twist $\mathcal{F}_{1/2}$ from [3]. We introduce the real parameter characterizing a whole family of Jordanian twists

²The correspondence with the generators of $sl(2)$ algebra is $H \rightarrow -D$ and $E \rightarrow p_\alpha$.

interpolating between the original Jordanian twist $\mathcal{F}_0 = \exp(-\ln(1 - a^\alpha p_\alpha) \otimes D)$ and Jordanian twist $\mathcal{F}_1 = \tau \circ \mathcal{F}_0|_{-a^\alpha} = \exp(-D \otimes \ln(1 + a^\alpha p_\alpha))$, where $a^\alpha = \frac{1}{\kappa} v^\alpha$ (in section 2). They both lead to the κ -Minkowski spacetime commutation relations. The aim of this paper is to present different methods to obtain the same family of twists³, developed in the framework of noncommutative spacetimes.

In section 2, a generalization of the r-symmetric twist from [3] is given. Corresponding deformed Hopf algebra symmetry, star products and differential realizations for noncommutative coordinates are presented. This family of Jordanian twists is also constructed as one exponential formula. However, in general, the twist might not always be known (e.g. in examples coming from deformation quantization framework).

In section 3, we provide an example of a method used to construct the twist. Starting from realizations of noncommutative coordinates, obtained in section 2, we construct the corresponding star product and coproduct. Also, from the star product and the coproduct, presented in section 2 (and coinciding with those in section 3), inverses of corresponding twists are obtained. They differ from the ones presented in section 2 by the right ideal, but indeed give the same deformed Hopf algebra, star product and realization of noncommutative coordinates. At the end of section 3, twists in Hopf algebroid approach, in terms of momenta and noncommutative coordinates, are given, showing that the used techniques appear within this more general framework. In section 4, concluding remarks are presented.

1.1 Notation and formalism

Here we will present the notation and recall some standard formulas related to the twist formalism which will be necessary for the following sections of the paper.

The Lie algebra g generated by the dilatation and momenta operators is defined by the commutation relations:

$$[p_\mu, D] = p_\mu, \quad [p_\mu, p_\nu] = 0.$$

The differential representation of the generators is the following: $D = x^\alpha \partial_\alpha = x \cdot \partial$ and $p_\mu = -i\partial_\mu$. We will also use the short notation for the momenta p_μ contracted⁴ with the vector a^ν as $A = ia^\nu \partial_\nu = -a^\nu p_\nu = -a \cdot p$. We do not specify signature of the metric, therefore the metric may be of any signature.

The generator D can be also rewritten in Heisenberg realization as $D = ix^\alpha p_\alpha = ix \cdot p$, in terms of Heisenberg algebra H , i.e.:

$$[x^\mu, x^\nu] = 0, \quad [p_\mu, x^\nu] = -i\delta_\mu^\nu, \quad [p_\mu, p_\nu] = 0. \quad (2)$$

The deformation of the Hopf algebra $U(g)(\mu, \Delta_0, \epsilon, S_0)$ of the universal enveloping algebra of $g = \{p_\alpha, D : [p_\alpha, D] = p_\alpha, [p_\mu, p_\nu] = 0\}$ given by the twist element $\mathcal{F} \in U(g) \left[\left[\frac{1}{\kappa} \right] \right] \otimes U(g) \left[\left[\frac{1}{\kappa} \right] \right]$ into $U^\mathcal{F}(g)(\mu, \Delta, \epsilon, S)$ is provided by the deformation of the coproduct and antipode maps as follows: $\Delta h = \mathcal{F} \Delta_0 h \mathcal{F}^{-1}$, $S(h) = [\mu((1 \otimes S)\mathcal{F})] S_0(h) [\mu((S \otimes 1)\mathcal{F}^{-1})]$, where $h \in g$.

The algebra of coordinates \mathcal{A} with multiplication map $m : \mathcal{A} \otimes \mathcal{A} \rightarrow \mathcal{A}$, i.e. the spacetime algebra is the Hopf module algebra with the action $U(g) \otimes \mathcal{A} \rightarrow \mathcal{A}$ of the Hopf algebra $U(g)$ on the module algebra $\mathcal{A} \ni f, g$ such that $h \triangleright (m(f \otimes g)) = m[\Delta h(\triangleright \otimes \triangleright)(f \otimes g)]$, where $h \in g$ and the

³up to the right ideal

⁴Note that here we are using the relativistic notation: $a^\nu b_\nu = a \cdot b$ (summation over ν index is assumed).

action \triangleright is defined by

$$p_\mu \triangleright f(x) = -i\partial_\mu \triangleright f(x) = -i\frac{\partial f(x)}{\partial x^\mu} \text{ and } D \triangleright f(x) = x^\alpha \frac{\partial f(x)}{\partial x^\alpha} \quad (3)$$

and the module acts on itself as $x^\mu \triangleright f(x) = x^\mu f(x)$.

The algebra of functions \mathcal{A} becomes noncommutative during the twist deformation once the usual multiplication is replaced by the star-product between the functions i.e.:

$$f \star g = m_\star(f \otimes g) = m[\mathcal{F}^{-1}(\triangleright \otimes \triangleright)(f \otimes g)] \quad (4)$$

for $f, g \in \mathcal{A}$. The star product is associative (due to the fact that the twist \mathcal{F} satisfies cocycle condition).

2 Simple one-parameter family of Jordanian twists

Let us introduce the following one-parameter family of twists $\mathcal{F}_u \in U(\mathfrak{g})[[\frac{1}{\kappa}]] \otimes U(\mathfrak{g})[[\frac{1}{\kappa}]]$:

$$\mathcal{F}_u = \exp(-u(DA \otimes 1 + 1 \otimes DA)) \exp(-\ln(1+A) \otimes D) \exp(\Delta_0(uDA)), \quad u \in \mathbb{R}, \quad (5)$$

The twists \mathcal{F}_u can be obtained from \mathcal{F}_0 by the transformation with the coboundary twist with the element $\omega = e^{uDA}$. Note that the deformation parameter $1/\kappa$ is included in A for the purpose of simplified notation. The twists \mathcal{F}_u are Drinfeld twists [1] $\forall u \in \mathbb{R}$, they satisfy normalization and cocycle condition⁵ and generalize the construction by V. N. Tolstoy [3], (cf. (1)). The correspondence of this family of the Jordanian twists (5) and the r-symmetric twist $\mathcal{F}_{1/2}$ is given by taking $u = \frac{1}{2}$ in (5). For $u = 0$, twist (5) simplifies to $\mathcal{F}_0 = \exp(-\ln(1+A) \otimes D)$ and for $u = 1$, simplifies to the twist $\mathcal{F}_1 = \tau \circ \mathcal{F}_0|_{-a} = \tau \mathcal{F}_0|_{-a} \tau = \exp(-D \otimes \ln(1-A))$.

The inverse of the above family of twists is

$$\mathcal{F}_u^{-1} = \exp(-\Delta_0(uDA)) \exp(\ln(1+A) \otimes D) \exp(u(DA \otimes 1 + 1 \otimes DA)), \quad u \in \mathbb{R}. \quad (6)$$

Now we can use the standard formulae (from section 1.1) for describing the deformation of the Hopf algebra maps of $U^{\mathcal{F}}(\mathfrak{g})(\mu, \Delta, \epsilon, S)$. The coproducts, star products and realizations in the rest of the paper depend on the parameter u , but for the sake of simplicity, the u -dependence will be omitted.

Deformed Hopf algebra

Coproducts Δp_μ , corresponding to the above twist, for any parameter $u \in \mathbb{R}$ are

$$\Delta p_\mu = \mathcal{F}_u \Delta_0 p_\mu \mathcal{F}_u^{-1} = \frac{p_\mu \otimes (1-uA) + (1+(1-u)A) \otimes p_\mu}{1 \otimes 1 + u(1-u)A \otimes A} \quad (7)$$

$$\Delta D = \mathcal{F}_u \Delta_0 D \mathcal{F}_u^{-1} = \left(D \otimes \frac{1}{1-uA} + \frac{1}{1+(1-u)A} \otimes D \right) (1 \otimes 1 + u(1-u)A \otimes A) \quad (8)$$

⁵It is based on the fact that the change of the twist by the coboundary twist \mathcal{F}_ω influences the change of the coproduct by $\tilde{\Delta} = \mathcal{F}_\omega \Delta \mathcal{F}_\omega^{-1}$ which is isomorphic to the coproduct Δ . Also adding the real valued parameter u will not influence the cocyclicity.

The coproduct is coassociative. The corresponding antipode is given by

$$S(p_\mu) = \frac{-p_\mu}{1 + (1 - 2u)A} \quad (9)$$

$$S(D) = -D - (1 - u)AD + uDA - \frac{u(1 - u)^2 A^2}{(1 + (1 - u)A)(1 - uA)} \quad (10)$$

The antipode is antihomomorphism. The corresponding counit is trivial, i.e. $\epsilon(p_\mu) = 0, \epsilon(D) = 0$ and $\epsilon(1) = 1$.

Interesting special cases are $u = 0$ (cf.[10]) and $u = 1$ (cf.[5]).

- For $u = 0$ we get:

$$\Delta p_\mu = p_\mu \otimes 1 + (1 + A) \otimes p_\mu \quad (11)$$

$$\Delta D = D \otimes 1 + \frac{1}{1 + A} \otimes D \quad (12)$$

$$S(p_\mu) = \frac{-p_\mu}{1 + A} \quad (13)$$

$$S(D) = -(1 + A)D \quad (14)$$

- For $u = 1$:

$$\Delta p_\mu = p_\mu \otimes (1 - A) + 1 \otimes p_\mu \quad (15)$$

$$\Delta D = D \otimes \frac{1}{1 - A} + 1 \otimes D \quad (16)$$

$$S(p_\mu) = \frac{-p_\mu}{1 - A} \quad (17)$$

$$S(D) = -D(1 - A) \quad (18)$$

Star product

The inverse of the twist \mathcal{F}_u^{-1} also provides the star product between the functions, as indicated in equation (4) This star product is associative (due to the fact that the twist \mathcal{F}_u (5) satisfies cocycle condition).

When we choose our functions to be exponential functions $e^{ik \cdot x}$ and $e^{iq \cdot x}$, then we define new function $\mathcal{D}_\mu(k, q)$:

$$e^{ik \cdot x} \star e^{iq \cdot x} = m [\mathcal{F}^{-1}(\triangleright \otimes \triangleright)](e^{ik \cdot x} \otimes e^{iq \cdot x}) = e^{i\mathcal{D}_\mu(k, q)x^\mu}, \quad (19)$$

where $k, q \in \mathcal{M}_{1, n-1}$ - n -dimensional Minkowski spacetime.

One can calculate explicitly that in the case of twist \mathcal{F}_u (5) the function $\mathcal{D}_\mu(k, q)$ is given by

$$\mathcal{D}_\mu(k, q) = \frac{k_\mu(1 + u(a \cdot q)) + (1 - (1 - u)(a \cdot k))q_\mu}{1 + u(1 - u)(a \cdot k)(a \cdot q)}. \quad (20)$$

Note that the function $\mathcal{D}_\mu(k, q)$ can be seen as rewriting the coproduct Δp_μ without using the tensor product notation (denoting left and right leg by k and q respectively). Therefore the relation between the coproduct Δp_μ and the function $\mathcal{D}_\mu(k, q)$ is given by

$$\Delta p_\mu = \mathcal{D}_\mu(p \otimes 1, 1 \otimes p), \quad (21)$$

hence Δp_μ uniquely determines $\mathcal{D}_\mu(k, q)$. In the case $u = 0$, $\mathcal{D}_\mu(k, q) = k_\mu + (1 - a \cdot k)q_\mu$, while in the case $u = 1$, $\mathcal{D}_\mu(k, q) = k_\mu(1 + a \cdot q) + q_\mu$. Deformed addition of momenta is given by $(k \oplus q)_\mu = \mathcal{D}_\mu(k, q)$.

Coordinates

Noncommutative coordinates \hat{x}^μ , corresponding to the twist \mathcal{F}_u (5), are given by

$$\begin{aligned}\hat{x}^\mu &= m [\mathcal{F}_u^{-1}(\triangleright \otimes 1)(x_\mu \otimes 1)] = x^\mu(1 - uA) + ia^\mu(1 - u)D(1 - uA) \\ &= (x^\mu + ia^\mu(1 - u)D)(1 - uA)\end{aligned}\quad (22)$$

Alternatively, we notice that they can also be obtained from the coproducts⁶

$$\begin{aligned}\hat{x}^\mu &= x^\mu + ix^\alpha m [(\Delta - \Delta_0)p_\alpha(\triangleright \otimes 1)(x^\mu \otimes 1)] \\ &= (x^\mu + ia^\mu(1 - u)D)(1 - uA).\end{aligned}\quad (23)$$

The noncommutative coordinates \hat{x}^μ satisfy

$$\begin{aligned}[\hat{x}^\mu, \hat{x}^\nu] &= i(a^\mu \hat{x}^\nu - a^\nu \hat{x}^\mu), \\ [p_\mu, \hat{x}^\nu] &= (-i\delta_\mu^\nu + ia^\nu(1 - u)p_\mu)(1 - uA).\end{aligned}\quad (24)$$

In the case $u = 0$, $\hat{x}^\mu = x^\mu + ia^\mu D$, while in the case $u = 1$, $\hat{x}^\mu = x^\mu(1 - A)$.

Note that we can define another set of the noncommutative coordinates \hat{y}^μ coming from the flipped version of (5) as

$$\begin{aligned}\hat{y}^\mu &= m [\tilde{\mathcal{F}}_u^{-1}(\triangleright \otimes 1)(x^\mu \otimes 1)] \\ &= x^\mu + ix^\alpha m [(\tilde{\Delta} - \Delta_0)p_\alpha(\triangleright \otimes 1)(x^\mu \otimes 1)] \\ &= (x^\mu - ia^\mu u D)(1 + (1 - u)A),\end{aligned}\quad (25)$$

where $\tilde{\mathcal{F}}_u = \tau \circ \mathcal{F}_u = \tau \mathcal{F}_u \tau$ and $\tilde{\Delta} = \tau \circ \Delta = \tau \Delta \tau$. Generators \hat{y}_μ define a dual coordinate algebra in the sense:

$$[\hat{y}^\mu, \hat{y}^\nu] = -i(a^\mu \hat{y}^\nu - a^\nu \hat{y}^\mu) \quad \text{and} \quad [\hat{x}^\mu, \hat{y}^\nu] = 0 \quad (26)$$

i.e. they obey κ -Minkowski commutation relations with $a^\mu \rightarrow -a^\mu$. The commutation relation $[p_\mu, \hat{y}^\nu]$ follows from the realization (25). In the case $u = 0$, $\hat{y}^\mu = x^\mu(1 + A)$, while in the case $u = 1$, $\hat{y}^\mu = x^\mu - ia^\mu D$.

One exponent formula for a family of Jordanian twists (5)

Above family of twists \mathcal{F}_u given by (5) can be written as

$$\mathcal{F}_u = \exp \left((D \otimes uA)\theta - ((1 - u)A \otimes D)\tilde{\theta}|_{(-a)} \right), \quad (27)$$

where

$$\theta = \sum_{n=0}^{\infty} \sum_{\substack{k, l=0 \\ k+l=n}}^n c_{k, l} A^k \otimes A^l \times \begin{cases} (1 - u)^k u^l & \text{for } u \in \mathbb{R} \setminus \{0, 1\} \\ \delta_{l0} & \text{for } u = 0 \\ \delta_{k0} & \text{for } u = 1 \end{cases} \quad (28)$$

⁶Later on we introduce (43) from which this relation follows naturally.

$$\tilde{\theta} = \tau \circ \theta = \tau \theta \tau, \quad (29)$$

where $c_{k,l} \in \mathbb{R}$ and $c_{0,0} = 1$. The first three terms in the above expansion in $\frac{1}{\kappa}$ are:

$$(\ln \mathcal{F}_u)_1 = D \otimes uA - (1-u)A \otimes D, \quad (30)$$

$$(\ln \mathcal{F}_u)_2 = \frac{1}{2} [D \otimes uA + (1-u)A \otimes D] [1 \otimes uA + (1-u)A \otimes 1], \quad (31)$$

$$(\ln \mathcal{F}_u)_3 = D \otimes uA \left[\frac{1}{3} (1 \otimes u^2 A^2) + \frac{1}{6} (1-u)A \otimes uA - \frac{1}{6} (1-u)^2 A^2 \otimes 1 \right] \\ - (1-u)A \otimes D \left[\frac{1}{3} (1-u)^2 A^2 \otimes 1 + \frac{1}{6} (1-u)A \otimes uA - \frac{1}{6} (1 \otimes u^2 A^2) \right]. \quad (32)$$

In special cases $u = 0$ and $u = 1$, \mathcal{F}_u reduces to

$$\mathcal{F}_0 = e^{-\ln(1+A) \otimes D}, \quad \mathcal{F}_1 = e^{-D \otimes \ln(1-A)}, \quad (33)$$

respectively.

The corresponding quantum R-matrix is $\mathcal{R} = \tilde{\mathcal{F}}_u \mathcal{F}_u^{-1} = 1 \otimes 1 + r + \mathcal{O}(\frac{1}{\kappa^2})$, where $r = A \otimes D - D \otimes A$, $\forall u \in \mathbb{R}$.

3 From realization to star product and twist

Realizations of noncommutative coordinates \hat{x}^μ can be generally expressed in terms of Heisenberg algebra H (2), generated by x^μ and p_μ . If \hat{x}^μ generate a Lie algebra, there exists universal formula for \hat{x}^μ , related to Weyl ordering [17].⁷ Starting from the realization (22) (see for example [18])

$$\hat{x}^\mu = x^\alpha \varphi_\alpha^\mu(p) = (x^\mu + ia^\mu(1-u)D)(1-uA), \quad (34)$$

one can reconstruct the star product using the following method. Recalling the action introduced in section 2 (3) we can explain how the exponent function of the noncommutative coordinates acts on the usual exponential function, via the realization (22) of \hat{x} .

$$e^{ik \cdot \hat{x}} \triangleright e^{iq \cdot x} = e^{ik \cdot x^\alpha \varphi_\alpha^\mu(p)} \triangleright e^{iq \cdot x} = e^{i\mathcal{P}(k,q) \cdot x}, \quad (35)$$

Here we introduced another set of functions, denoted by $\mathcal{P}_\mu(k, q)$ which satisfy the following differential equations [19, 20, 21, 22]

$$\frac{d\mathcal{P}_\mu(\lambda k, q)}{d\lambda} = \varphi_\mu^\alpha(\mathcal{P}(\lambda k, q))k_\alpha. \quad (36)$$

Note that the equation involves the same function $\varphi_\mu^\alpha(p)$ from the realization of the noncommutative coordinates. The boundary conditions are $\mathcal{P}_\mu(0, q) = q_\mu$. The solution is

$$\mathcal{P}_\mu(k, q) = \frac{K_\mu(k)(1 + u(a \cdot q)) + (1 - (1-u)(a \cdot K(k)))q_\mu}{1 + u(1-u)(a \cdot K(k))(a \cdot q)}, \quad (37)$$

⁷For a more general case, see [22].

where

$$K_\mu(k) = \mathcal{P}_\mu(k, 0) = k_\mu \frac{e^{a \cdot k} - 1}{a \cdot k} \frac{1}{(1-u)e^{a \cdot k} + u}. \quad (38)$$

The inverse function of $K_\mu(k)$, defined as $K_\mu(K^{-1}(k)) = K_\mu^{-1}(K(k)) = k_\mu$, is given by

$$K_\mu^{-1}(k) = k_\mu \frac{1}{a \cdot k} \ln \left(\frac{1 + u(a \cdot k)}{1 - (1-u)a \cdot k} \right). \quad (39)$$

There exists a relation between the function $\mathcal{P}_\mu(k, q)$ and introduced before function $\mathcal{D}_\mu(k, q)$ as $\mathcal{D}(k, q) = \mathcal{P}(K^{-1}(k), q)$. Thanks to this, we can rewrite the corresponding star product of exponential functions as

$$e^{ik \cdot x} \star e^{iq \cdot x} = e^{iK^{-1}(k) \cdot \hat{x}} \triangleright e^{iq \cdot x} = e^{i\mathcal{P}(K^{-1}(k), q) \cdot x} = e^{i\mathcal{D}(k, q) \cdot x}. \quad (40)$$

The function $\mathcal{D}_\mu(k, q)$ from equation (40) coincides with equation (20).

From the function $\mathcal{D}_\mu(k, q)$, the coproduct Δp_μ can be recovered using (21) and antipodes $S(p_\mu)$ follow analogously. Alternatively, it is possible to find the coproduct Δp_μ as

$$\Delta p_\mu = e^{iK_\alpha^{-1}(p) \otimes \text{ad}_{\hat{z}^\alpha}} (1 \otimes p_\mu), \quad (41)$$

From star product and coproduct to twist

From equations (21) and (40) it follows [13, 14, 21]:

$$\begin{aligned} e^{ik \cdot x} \star e^{iq \cdot x} &= \lim_{\substack{y \rightarrow x \\ z \rightarrow x}} \left(e^{ix \cdot [\mathcal{D}(-i\partial^y, -i\partial^z) + \partial^y + \partial^z]} (e^{ik \cdot y} e^{iq \cdot z}) \right) \\ &= m \left[: e^{i((1-t) \otimes x^\alpha + tx^\alpha \otimes 1)(\Delta - \Delta_0)p_\alpha} : (\triangleright \otimes \triangleright) (e^{ik \cdot x} \otimes e^{iq \cdot x}) \right], \quad \forall t \in \mathbb{R}. \end{aligned} \quad (42)$$

Note that on the right hand side of the above equation the terms multiplied by t cancel each other due to the normal ordering and the multiplication map m .

From this identity and equations (4), (19), it follows:

$$\mathcal{F}^{-1} =: e^{i((1-t) \otimes x^\alpha + tx^\alpha \otimes 1)(\Delta - \Delta_0)p_\alpha} : + \mathcal{I}_0, \quad (43)$$

where $\mathcal{I}_0 \subset H \otimes H$ is the right ideal of the coordinate algebra \mathcal{A} , defined by

$$m[\mathcal{I}_0(\triangleright \otimes \triangleright)(\mathcal{A} \otimes \mathcal{A})] = 0 \quad (44)$$

For the case $u = 0$, with $t = 0$, twist \mathcal{F}_0^{-1} is given by

$$\mathcal{F}_0^{-1} =: e^{A \otimes D} := e^{\ln(1+A) \otimes D}, \quad (45)$$

while for the case $u = 1$, with $t = 1$, twist \mathcal{F}_1^{-1} is given by

$$\mathcal{F}_1^{-1} =: e^{-D \otimes A} := e^{D \otimes \ln(1-A)}. \quad (46)$$

Calculation of \mathcal{F}^{-1} in equation (43) in the form $\mathcal{F}^{-1} = e^{-f}$ for linear realizations of \hat{x}^μ is presented in [21] and for Abelian twists in [13].

Twist in the Hopf algebroid approach

Deformed phase spaces of Lie algebra type are presented and studied in [15, 23, 26, 27]. Twists in Hopf algebroid approach were studied and constructed in [14, 16, 24, 25].

Generally, in Hopf algebroid approach, the twist is given by

$$\mathcal{F}^{-1} = e^{-ip_\alpha \otimes x^\alpha} e^{iK_\gamma^{-1}(p) \otimes \hat{x}^\beta} + \mathcal{I}_0 = e^{-ip_\alpha \otimes x^\alpha} e^{iK_\gamma^{-1}(p) \otimes x^\beta \varphi_\beta^\gamma(p)} + \mathcal{I}_0 \quad (47)$$

This is a generalization of the result presented in [16]. The full Hopf algebroid analysis will be presented elsewhere.

4 Concluding remarks

Inspired by the recent interest in Jordanian deformations, in this paper we focused on the simple generalization of the locally r -symmetric Jordanian twist. By introduction of one real valued parameter u we obtained the family of Jordanian twists which provides interpolation between the original Jordanian twist (for $u = 0$) and its flipped version (for $u = 1$, up to minus sign in the deformation parameter). All of the proposed twists provide the κ -Minkowski spacetime and have the support in the Poincaré-Weyl or conformal algebras as deformed symmetries of this noncommutative spacetime. Another important issue considered here was the presentation of different methods relating the twist, deformed coproducts with the star-product and the realizations for the noncommutative coordinates.

It is important to note that we can present our results symbolically by the following diagram:

$$\begin{array}{ccc} & \hat{x}^\mu = x^\alpha \varphi_\alpha^\mu(p) & \\ & \uparrow \downarrow & \\ & \mathcal{F} & \\ & \downarrow \uparrow & \\ e^{ik \cdot x} \star e^{iq \cdot x} & \longleftrightarrow & \Delta p_\mu \end{array}$$

In section 2, starting from the twist (5), we have found deformed Hopf algebra symmetry (7)-(10), star products (19), (20) and corresponding realizations for noncommutative coordinates (22). Relation between coproduct Δp_μ and star products $e^{ik \cdot x} \star e^{iq \cdot x}$ (19), (20) are given in (21). Noncommutative coordinates \hat{x}^μ are also obtained from the coproduct Δp_μ in equation (23) and also from the related star product $e^{ik \cdot x} \star e^{iq \cdot x} = e^{i\mathcal{D}(k,q) \cdot x}$. In section 3, starting from realizations of noncommutative coordinates (34), we have constructed corresponding star products (20), (40) and coproducts. Alternatively, coproduct Δp_μ is obtained from noncommutative coordinates (41). Also, from the star product and the coproduct in section 2, see also equations (40), (42), inverses of the corresponding twist (43) are obtained. Twists in the Hopf algebroid approach, in terms of momenta and noncommutative coordinates, are given in equation (47). Note that the inverse twist \mathcal{F}^{-1} is not uniquely determined from the star product, coproduct Δp_μ and noncommutative coordinates \hat{x}^μ and it is determined up to a right ideal \mathcal{I}_0 .

More general realizations are of the type $\hat{x}^\mu = x^\alpha \varphi_\alpha^\mu(p) + \chi^\mu(p)$, but we restricted our considerations to the case $\chi^\mu(p) = 0$ and the above diagram corresponds to $\chi^\mu(p) = 0$. We point out that all our results are exact. There are few possible extensions of these results, leading to new insights of the deformation quantization. However, they require an extension to the Hopf algebroid

framework. Drinfeld twists can be lifted to Hopf algebroids [28] from the Hopf algebra, but it is worth to note that not all the twists obtained from the coproduct are the Hopf algebra twists. The issues of the twist reconstruction within Hopf algebroids and the generalization of the other version of symmetrized Jordanian twist presented in [4] will be thoroughly discussed in another paper.

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