# Kinematics and dynamics motion planning by polar piecewise interpolation and geometric considerations 

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#### Abstract

The importance of numerical methods in science and engineering[1] was long recognised and considered a fundamental factor in improving productivity and reducing production costs. The ability to model flexible systems and describe their trajectories [5] involves usually the study of nonlinear coupled partial differential equations. Since their exact solutions are not normally feasible in practice, computational methods [8] can be considered.


Keywords: Numerical methods, motion planning, kinematics, dynamics.

The development of such computationally efficient numerical algorithms and methods for solving PDE and handling of complicated geometries [13] require the use of implicit (trajectory and forces computed along the motion) and explicit (trajectory computed before execution) motion planning methods [14]. Implicit approaches based on potential field methods consider a potential

[^0]function in the C-space to asymptotically generate an attractive goal configuration. The potential function can be defined [8] by
\[

$$
\begin{equation*}
U(q)=\frac{1}{2} \zeta\left\|q-q_{\text {Goal }}\right\| \tag{1}
\end{equation*}
$$

\]

where $\zeta$ is a scale factor. Explicit methods [8] can be separate in continuous methods namely kinematic or dynamic motion planning [16] or discrete methods. The kinematic approach often referred as geometry of motion is essential the initial step in the computation of the forces related to the dynamical model (dynamical equations of the system) and their effect on the motion. Kinematic motion planning [8] formulated as a variational approach consist in generating a smooth trajectory $\delta:[0,1] \rightarrow \mathbb{R}$ by minimizing a functional

$$
\begin{equation*}
J(\delta)=\int_{0}^{1}<L\left(\delta, \frac{d \delta}{d t}\right), L\left(\delta, \frac{d \delta}{d t}\right)>d t, \quad \text { and } \tag{2}
\end{equation*}
$$

where $t$ is the integration time, $<\cdot \cdot \cdot>$ is the Riemannian metric and $L$ the initial conditions [8]. Dynamic motion planning [8] can be defined as a mathematical minimization of a cost function [16] by

$$
\begin{equation*}
\min \left(J_{c}\right) \text { where } J_{c}=\Phi\left(\left(y \mid\left(t_{f}\right)\right)+\int_{t_{0}}^{t_{f}} L(y, u) d t\right. \tag{3}
\end{equation*}
$$

subject to initial conditions $y\left(t_{0}\right)=x_{0}$ and $\dot{y}=f(y, u)$. Parameterizations of joint trajectories transform the dynamic motion planning problem into a parameter optimization [8] problem of a family of smooth curves generated by parametric and/or geometric continuous splines, uniform cubic $B$ spline with parametric and geometric continuity, or by piecewise interpolating curves [4,6,7,9].

Since the motion along a set of control point can be eventually expressed as a combination of rotations about a fixed point, that is, rotations about $O x, O y$ and $O z$ axes that passes through the origin, the use of polar and/or Cartesian piecewise interpolating curves $[2,11]$ can be then related with the robot dynamics $[2,10,12]$ through its geometry [2]. The set of the considered rotations can be represented as a 3D manifold, that is, the real projective 3 -space denoted by $R P^{3}$. Depending on the geometry of a robotic manipulator arm, the fixed point representing the center of rotation, e.g., base of the manipulator, can be related to either polar zenithal gnomic, polar zenithal stereographic or polar zenithal orthographic projections, by the use of the time dependent generalized coordinates expressed as spherical or cylindrical coordinates. To interpolate between some a priori defined control points, a piecewise
polar interpolant that approximate the polar trajectory can be expressed as a Hermite-type function [2,7] defined by

$$
\begin{equation*}
r(\theta)=\sum_{k=0}^{q} c_{k}^{i}\left(\theta-\theta_{i}\right)^{k} \tag{4}
\end{equation*}
$$

where $q=3$ (cubic approximation) is the order of the polynomial, $c_{0}^{i}=r_{i}$, $c_{1}^{i}=\dot{r}_{i}, c_{2}^{i}=\frac{1}{h_{i}}\left[\left(2 \dot{r}_{i}+\dot{r}_{i+1}\right)+3 \Delta r_{i}\right], c_{3}^{i}=\frac{1}{h_{i}^{2}}\left[\dot{r}_{i}+\dot{r}_{i+1}-2 \Delta r_{i}\right], r_{i}=r\left(\theta_{i}\right)$, $r_{i+1}=r\left(\theta_{i+1}\right), h_{i}=\theta_{i+1}-\theta_{i}, \Delta r_{i}=\frac{r_{i+1}-r_{i}}{h_{i}}$, and the derivatives at the endpoints are calculated using $\dot{r}\left(\theta_{i}\right)=\frac{d r\left(\theta_{i}\right)}{d \theta}=\dot{r}_{i}$ and $\dot{r}\left(\theta_{i+1}\right)=\frac{d r\left(\theta_{i+1}\right)}{d \theta}=\dot{r}_{i+1}$ respectively.

Considering the case of polar zenithal orthographic projection and related cylindrical coordinates (Fig. 1.c), trajectories representing the system dynamics can be generated using a combination of polar piecewise interpolants that approximate the polar trajectory, and Cartesian piecewise interpolants to approximate the trajectory height. Trajectory height is computed in the projective plane obtained by unfolding the cylinder with generators passing through the polar trajectory of the motion.

Considering the case of polar gnomic or stereographic projection and related spherical coordinates (Fig. 1.a and Fig. 1.b), trajectories representing the system dynamics can be generated by a combination of polar piecewise interpolants over a "spherical" domain relatively similar with [6] by

$$
\begin{equation*}
R(\lambda, \phi)=S_{0}(\psi)+S_{1}(\psi) \chi+S_{2}(\psi) \chi^{2}+S_{3}(\psi) \chi^{3} \tag{5}
\end{equation*}
$$

where $\lambda$ and $\phi$ represents the generalised coordinates related to the system model, $S(\psi)=c_{0,0}+c_{1,0} \psi+c_{2,0} \psi^{2}+c_{3,0} \psi^{3}$, and and $\psi$ and $\chi$ are defined as in [6].

Therefore, adequate geometric considerations - such as manipulators geometry and related coordinate systems - used in the motion planning studies backed up by polar piecewise cubic interpolants allows easy calculation of kinematics variable and handling of the dynamical equations of motion [12] through the use of Lagrange's equations defined by

$$
\begin{equation*}
\frac{d}{d t}\left(\frac{\partial T}{\partial \dot{q}_{j}}\right)-\frac{\partial T}{\partial q_{j}}=Q_{j} \tag{6}
\end{equation*}
$$

The kinetic energy $T=$ and the acting contact and body forces $Q_{i}$ can be expressed as in [12] by

$$
\begin{equation*}
T_{j}=\frac{1}{2} m_{j} \mathbf{v}_{C_{j}} \cdot \mathbf{v}_{C_{j}}+\frac{1}{2} \bar{I}_{j} \omega_{j} \cdot \omega_{j} \tag{7}
\end{equation*}
$$

and

$$
\begin{equation*}
Q_{j}=\frac{\partial \omega}{\partial \dot{q}_{j}} \cdot \mathbf{T}+\frac{\partial \mathbf{v}}{\partial \dot{q}_{j}} \cdot \mathbf{R} \tag{8}
\end{equation*}
$$

where $v$ is the velocity, $\omega$ is the angular velocity, $\mathbf{T}$ is the applied torque, and $\mathbf{R}$ is the applied force.


Figure 1

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