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Observation of successive phase transitions in ferroelectric liquid crystal (CS-1013) from dielectric measurements and a pseudospin type model for studying ferro and antiferroelectric behaviour

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Abstract From dielectric measurements, we have shown the appearance of successive phase transitions in a ferroelectric liquid crystal CS-1013 Between crystalline and isotropic phases, we observed successive phase transitions ν_{12} Cr $\sim 35^{\circ}$ SmC^{*} = 63° C = SmA = 70° C=N* 80° C=lso Both Goldstone mode (GM) and Soft mode (SM) are observed in the SmC* phase. A theoretical pseudospin type model has been suggested for studying ferro and antiferroelectric phase transitions in liquid crystals.

keywords Liquid crystal, ferroelectricity, Goldstone and soft mode

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Phase transition in crystalline ferroelectric material is a vastly studied field of interest [1]. But phase transition in ferroelectric hquid crystal is a most recent discovery when 4-(1-methylhepty loxycatbonyl) phenyl 4'-octylbiphenyl-4-carboxylate (MHPOBC) was found to be the first ferroelectric liquid crystal [2] After this discovery, many chiral compounds with mesophases viz. ferroelectric (FE)-SmC*, antiferroelectric (AFE)- SmC_{4}^{*} , ferrielectric (FEE)-SmC_v, Sm_aC^{*} etc. have been discovered by electro optic [3, 4], switching [2, 5, 6], pyroelectric [7] and other measuring techniques. However, none of these techniques is self-sufficient to provide detailed information of the ferroelectric phase transition, molecular dynamics etc. Temperature and frequency-dependent dielectric permittivity and loss data provide important information about the successive phase transitions and molecular dynamics in liquid crystals.

In the present communication, we report the different phase transitions occurring between the crystalline and the isotropic phases of a typical ferroelectric liquid crystal FLC viz. CS-1013 using dielectric spectroscopic technique. The sample was prepared by Chisso Petrochemical Corporation, Japan. Dielectric constants (both real and the imaginary parts) were measured by a computer controlled HP 4192A impedance analyzer having measuring frequency range between 5Hz and 13MHz. The temperature was controlled (with an accuracy of ± 0.1°C) by an Eurotherm temperature controller. The temperature range of our investigation was between 29-82°C. Our experimental set-up is similar to that described elsewhere [8, 9]. Sputtered gold plated electrodes were used to measure the dielectric constant of the FLC layer of thickness ~ 15 µm. The FLC filled cell was prepared by capillary action method in an isotropic phase. Stabilized low frequency ac field (~ 1.67 kV/cm) was used for better molecular alignment in this FLC sample. The different phases of SC-1013 studied by other methods [10] are shown below

Some important parameters of the liquid crystal under investigation are shown in Table 1.

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 Table 1. Some important parameters of the ferroelectric liquid crystal

 CS-1013

Spontaneous polarization (Ps)	25°C	15 0nCcm ⁻²
Tilt angle (θ)	25"C	26 deg
Helical pitch in N* phase	- T _N ''	38 µm
Helical pitch in SmC* phase	25°C	6 µm
Response time (τ)	25"C, E = ± 10V/μm	739 µsec
Rotational Viscosity $(\eta_0^{(2)})$	25°C	1282 mPa s
Optical anisotropy (An)	25°C, $\lambda = 546$ nm	0 12

1) T_N means N*-SmA just above transition temperature 2) $\eta_0 = \eta / \sin^2 \theta$

Figure 1 shows the temperature dependent dielectric loss (\mathcal{E}'') factor at two different fixed frequencies (1 and 10 kHz). Different phase transitions described by eq. (1) are clearly



Figure 1. Temperature-dependent dielectric loss (ϵ ") of ferroelectric CS10-13 showing different phases



Figure 2. Temperature-dependent dielectric permittivity (real part ε') of ferroelectric CS-1013 at two different temperatures.

observed from this Figure. The SmA-SmC* phase transition appears to be first order in nature. Temperature dependence of dielectric permittivity (ε') shown in Figure 2 also indicates the presence of different phases in SC-1013.

It is well known that in planar alignment, the dielectru permittivity due to Goldstone mode [8, 11] is dominated in SmC4 phase which appears due to the phase fluctuations in tilt angle of the molecules. The other electric permittivity contributions are due to soft mode [12] (caused by amplitude fluctuation in the angle), domain mode [13] and unwinding mode [14] (earlier theoretically predicted). These modes are generally observed by suppressing the Goldstone mode in SmC* phase. In the present sample, the dielectric permittivity due to Goldstone mode is dominated in the SmC* phase. However, soft mode is also observed near the SmA-SmC* transition region, as observed from Figure 3. It is interesting to note that one also observes (from Figure 3) a new molecular mode around 70kHz This new mode is considered to be due to the re-orientation of the long helical axis. Similar curves have also been obtained at different temperatures and frequencies showing several interesting features in this FLC which would be published elsewhere



Log f (Hz)

Figure 3. Frequency dependence of the imaginary part of dielectile constant (ϵ'') in ferroelectric SmC* phase at different fixed temperatures

We find that the crystalline to ferroelectric (SmC*) phase transition can be well described by a pseudospin Hamiltonian of the form

$$H = -\mu E \Sigma \sigma_i^z - U \Sigma (\sigma_i^z)^2 - B \Sigma (\sigma_i^z)^2 - \Sigma J_{ii} \sigma_i^z \sigma_j^z.$$

Here, the pseudospin variable σ (with $\sigma_{\lambda}, \sigma_{\nu}$, and σ components) is associated with the long axis of the molecule μ is the electric field (E)-dependent dipole moment, U is the magnitude of the uniaxiality parameter representing single molecular anisotropy in the z direction of polarization (a measure of the magnitude of pitch), B is the magnitude of the biaxiality parameter representing the molecular anisotropy transverse to the z-direction (a measure of the tilt angle θ) and J_{ij} (>0 for ferroelectric transition) is the effective exchange interaction between the dipoles (pseudospin). J is also a measure of the interaction strength between B and U (*i.e.* pitch and θ) at sites *i* and *j*. All the parameters U, B, and J undergo thermodynamic fluctuations along with tilt angle and pitch) and are responsible for different mesophase transitions in the liquid crystals. Though this model Hamiltonian [eq. (2)] is similar to that of Devil's staircase model used by Yamashita and Miyazima [15], our derivation of the model and physical significance of the model parameters are juite different from those used by them. Detailed calculations of the transition temperature, spontaneous polarization *etc.* using this model [eq. (2)] will be published elsewhere.

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