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## Studies on some parameters in heavily doped $Al_xGa_{1-x}As/GaAs$ double heterostructure

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Abstract : The variation of several parameters within heavily doped Al<sub>x</sub>Ga<sub>1x</sub>A/GaAs double heterostructure has been investigated theoretically taking into account band gap narrowing and carrier degeneracy as heavy doping effects. The results of the computational analyses are shown graphically.

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The minority carrier lifetime in III-V semiconductors like GaAs and  $AI_xGa_{1-x}As$  is an important parameter for a number of applications. To produce electrical and optical devices like high electron mobility transistors (HEMTs) [1], double heterostructure lasers [2] and high-efficiency devices [3],  $AI_xGa_{1-x}As/GaAs$  heterostructure is commonly used. Time resolved photo-luminescence [4] technique is used to measure minority carrier lifetime in some III. V semiconducting materials. Minority carrier lifetime is useful for minority-carrier devices like light emitting diodes (LEDs), photovoltaic cells, bipolar transistors. Various methods are used widely to determine the minority carrier lifetime in photovoltaic devices [5–7]. In III–V compounds, time-resolved photoluminescence decay method is useful for the measurement of the minority carrier lifetime. Carrier lifetime in silicon is measured by using pulse optical excitation and photo-conductivity decay technique. Heavy doping effects play important role in producing high emitter efficiency in bipolar transistors and high open-circuit voltage in solar cells, and have influence on lifetime in band-to-band processes. Experimentally measured lifetimes [4,8] of heavily doped GaAs and quaternary alloy are reported earlier to study Auger effects.

In this presentation, assuming uniform distribution of SRH (Shockley-Read-Hall) defects within specified region, the variation of bulk minority carrier lifetime with carrier

density and photoluminescence lifetime with interface recombination velocity in heavily doped  $AI_xGa_{1-x}As/GaAs$  double heterostructure have been investigated. Band gap narrowing and carrier degeneracy are considered as heavy doping effects. As band-to-band Auger recombination effect is a must, the variation of quantum efficiency and photoluminescence lifetime with nominal current density are studied for the same hetero-structure. Spatial variation of minority carrier density has been ignored owing to lower value of diffusion transit-time compared to the minority carrier lifetime. The results so obtained by computational analyses are shown graphically. For  $AI_xGa_{1-x}As/GaAs$  double heterostructure, photoluminescence [9] lifetime, under low interface recombination velocity, approaches the bulk minority carrier lifetime. The photoluminescence lifetime  $\tau_{PL}$  is given by [10]

$$\frac{1}{\tau_{\text{PL}}} = \frac{1}{\tau_{R}} + \frac{1}{\tau_{\text{SRH}}} + \frac{1}{\tau_{S}}$$
(1)  
$$\frac{1}{\tau_{S}} = \frac{2S}{d}.$$

where

 $\tau_R$  is the radiative lifetime;  $\tau_{SRII}$ , the Shockley-Read-Hall lifetime;  $\tau_S$ , the surface lifetime; S, the interface recombination velocity, and d is the active layer thickness.

The relationship among minority carrier diffusivity (D), decay time (t) and mobility ( $\mu$ ) can be written as

$$D = \frac{d^2}{2t} \text{ and } D = \frac{KT}{q}\mu, \qquad (2)$$

where K is the Boltzmann constant; q, the electronic charge, and T is the absolute temperature. For *n*-type semiconductor with donor concentration  $N_D$ , the electron mobility  $\mu$  is given by [11]

$$\mu_e(N_D) = \frac{\mu_o}{1 + (N_D/N_{\text{eff}})\alpha} + \mu_{\text{min}}.$$
(3)

 $\mu_o$  is the difference between the expected maximum and minimum mobilities;  $\mu_{\min}$ , the expected minimum mobility value;  $N_{\text{eff}}$ , a reference concentration, and  $\alpha$  is an exponential factor that controls the slope around  $N_D = N_{\text{eff}}$ . Again the radiative and SRH lifetime [12] are given by

$$\tau_R = \frac{\Delta n}{Bnp}, \quad \tau_{\text{SRH}} = \frac{\Delta n}{np} \left[ \tau_p (n + n_i) + \tau_n (p + n_i) \right], \quad (4)$$

where  $\Delta n$  is the excess photogenerated carrier; *B*, the radiative recombination coefficient;  $n_i$ , the intrinsic carrier concentration and  $\tau_p$  and  $\tau_n$  are the hole and electron lifetimes, respectively. Under heavy doping condition, electron (*n*) and hole (*p*) concentrations can be expressed as [13]

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$$n = n_{i} F_{1/2}(\eta_{n}) \exp(-\eta_{n}) \exp \frac{A\Delta E_{g}}{KT} \exp \frac{E_{fn} - E_{i}}{KT},$$
 (5a)

$$p = n_{r} F_{1/2}(\eta_{p}) \exp(-\eta_{p}) \exp\frac{(1-A)\Delta E_{g}}{KT} \exp\frac{E_{r} - E_{fp}}{KT},$$
 (5b)

where  $F_{1/2}$  ( $\eta_n$ ) is the Fermi-Dirac integral of order 1/2;  $\eta$ , the reduced Fermi energy; A, the asymmetry factor:  $\Delta E_g$ , the bandgap narrowing:  $E_i$ , the intrinsic Fermi energy, and  $E_{fn}$  and  $E_{fp}$  are the quasi-Fermi energies of electron and hole, respectively. At equilibrium,  $E_{fn} = E_{fp}$ .

Moreover, the interface recombination velocity (S) is related to recombination current density  $(J_s)$  as

$$S = \frac{J_s}{q\Delta n},\tag{6}$$

$$J_{s} = q \Big[ R_{\rm SRII} + R_{\rm Aug} \Big], \qquad (7)$$

where  $R_{SRH}$  is the interface recombination due to SRH process [14,15] and  $R_{Aug}$  is the net Auger recombination. These are expressed as

$$R_{\rm SRH} = \frac{np - n_i^2}{\tau_n(p + n_i) + \tau_p(n + n_i)},$$
(8)

$$R_{Aug} = r \left( \frac{np - n_i^2}{n_i^2} \right) (n + p), \qquad (9)$$

r is the Auger recombination coefficient.

Thus, from eqs. (1) - (9), one obtains

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$$\frac{1}{\tau_{PL}} = \frac{n_i}{\Delta n} \left[ \left[ F_{1/2}(\eta_n) F_{1/2}(\eta_p) \exp \frac{\Delta E_g}{KT} \exp \left\{ -(\eta_n + \eta_p) \right\} \right] \right]$$

$$\times \left[ n_i B + \left[ \tau_n \left\{ F_{1/2}(\eta_p) \exp \left( -\eta_p \right) \exp \frac{(1 - A)\Delta E_g}{KT} \right] \right] \right]$$

$$\times \exp \left( \frac{E_i - E_{fp}}{KT} \right) + 1 \right] + \tau_p \left\{ F_{1/2}(\eta_n) \exp \left( -\eta_n \right) \right\}$$

$$\times \exp \left( \frac{A\Delta E_g}{KT} \exp \left( \frac{E_{fn} - E_i}{KT} \right) + 1 \right\} - 1 \right]^{-1} + \left( \frac{2q}{tKT} \right)^{1/2}$$

$$\times \left\{ \frac{\mu_o}{1 + (N_D/N_{eff})\alpha} + \mu_{mun} \right\}^{-1/2} \left[ \left[ \tau_n \left\{ F_{1/2}(\eta_p) \right\} \right]^{-1/2} \right] \right] = 0$$

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$$\times \exp\left(-\eta_{p}\right) \exp\left(\frac{(1-A)\Delta E_{g}}{KT} \exp\left(\frac{E_{i}-E_{fp}}{KT}\right) + 1\right)$$

$$+ \tau_{p} \left\{F_{1/2}(\eta_{n}) \exp\left(-\eta_{n}\right) \exp\left(\frac{A\Delta E_{g}}{KT} + \exp\left(\frac{E_{fn}-E_{i}}{KT}\right) + 1\right\}\right]^{1}$$

$$+ r\left\{F_{1/2}\left(\eta_{n}\right) \exp\left(-\eta_{n}\right) \exp\left(\frac{A\Delta E_{g}}{KT}\right)$$

$$\times \exp\left(\frac{E_{fn}-E_{i}}{KT}\right) + F_{1/2}(\eta_{p}) \exp\left(-\eta_{p}\right) \exp\left(\frac{(1-A)\Delta E_{g}}{KT}\right)$$

$$\times \exp\left(\frac{E_{i}-E_{fp}}{KT}\right)\right\}\left[\left[F_{1/2}(\eta_{n})F_{1/2}(\eta_{p})\exp\left(\frac{\Delta E_{g}}{KT}\right) + \frac{\Delta E_{g}}{KT}\right]$$

$$\times \exp\left\{-(\eta_{n}+\eta_{p})-1\right]\right].$$

$$(10)$$

The computational analyses of eq. (10) have been carried out for a fixed composition under different circumstances. Figure 1 represents the variation of bulk minority carrier lifetime of





Figure 2. Photoluminescence lifetime versus interface recombination velocity of an  $Al_{0.3}Ga_{0.7}$ . As/GaAs/Al<sub>0.3</sub>Ga<sub>0.7</sub>As DH device ( $d = 4 \mu m$ ).

heavily doped  $Al_{0.3}Ga_{0.7}As/GaAs/Al_{0.3}Ga_{0.7}As$  DH ( $d = 4 \mu m$ ) with carrier density under 20 mw focussed power. The nature of variation reveals that bulk minority carrier lifetime decreases in a nonuniform manner with the increase of concentrations under the focussed

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power for a given cross-sectional area and thickness of a double heterostructure. Numerical computations of equation (10) are worked out considering  $\Delta E_g = 10.23$  (N/10<sup>18</sup>)<sup>1/3</sup> + 13.12 (N/10<sup>18</sup>)<sup>1/4</sup> + 2.93 (N/10<sup>18</sup>)<sup>1/2</sup> meV [16], T = 300 K,  $B \simeq 10^{-10}$  cm<sup>3</sup> s<sup>-1</sup>. The values of  $\eta_p$ ,  $\eta_n$ ,  $F_{1/2}$  ( $\eta_p$ ) and  $F_{1/2}$  ( $\eta_n$ ) are chosen suitably [17] for the given carrier densities. In numerical analysis, change of carrier lifetime with carrier densities is considered.



Figure 3. Variation of quantum efficiency with nominal current density of an Al  $_{0.3}$ Ga $_{0.7}$ As/GaAs/Al $_{0.3}$ -Ga $_{0.7}$ As/DII device ( $d = 2 \mu m$ ).

Figure 4. Variation of photoluminescence lifeting with nominal current density of an Al<sub>0.3</sub>Ga<sub>0.7</sub>As DH device  $(d = 2 \mu m)$ .

The values of minority carrier lifetimes are chosen within  $5 \times 10^{-7} - 10^{-7}$  s when majori carrier concentration varies in the range  $10^{18} - 10^{20}$  cm<sup>-3</sup>. Again, the values of majority carrie lifetimes are chosen within  $10^{-8} - 10^{-12}$  s when minority carrier concentration varies in the range  $(1-8) \times 10^{16}$  cm<sup>-3</sup>. These data are used to study the change of photoluminescen lifetime with interface recombination velocity of an A1<sub>0.3</sub>Ga<sub>0.7</sub>As/GaAs/Al<sub>0.3</sub>Ga<sub>0.7</sub>As D ( $d = 4 \mu$ m), which is depicted in Figure 2. Figure 3 represents the change of quantum efficiency with nominal current density for the same semiconducting compound with the same composition and  $d = 2 \mu$ m. The variation of photoluminescence lifetime with nominal current density of the same DH (Al<sub>0.3</sub>Ga<sub>0.7</sub>As/GaAs/Al<sub>0.3</sub>Ga<sub>0.7</sub>As) with  $d = 2 \mu$ m is shown Figure 4.

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