# A TOPOLOGICAL TECHNIQUE FOR ANALYSIS OF ACTIVE NETWORKS <br> John T. Barrows, Jr. 

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## ABSTRACT

This paper presents a generalization of Fuessner's method for the analysis of active networks. This method makes it possible to determine system functions of any network in a simple and straightforward manner.

In the usual topological analysis of networks it is necessary to list all possible trees or co-trees and the signs of their tree admittance products as well as determining the set of 2 -tree products and their signs for the network function in question. Usually two graphs, a current and a voltage graph, are used to implement this process when dependent sources are present in the network. The method presented here utilizes solely the original topology and the assigned branch current directions to determine these terms and their sign. Justification of the technique is given for the cases of one and two dependent sources present in a network with the summary of the method for each case. The analysis of a network with an arbitrary number of dependent sources is seen to follow directly. Finally an example of a network with two dependent sources is given which illustrates the method.

## INTRODUCTION

The early methods of Kirchhoff ${ }^{4,5}$, Maxwe $11^{9}$ and Feussner ${ }^{2,7,11}$ demonstrate the practicability and simplicity of topological techniques used in the determination of network functions of passive networks. This paper is concerned with the application of topological techniques to active (nonseparable) networks with the discussion in terms of branch impedances and admittance functions. Specifically, Feussner's method is extended to networks containing dependent voltage sources which depend on current in other branches of the network. By using elementary source transformations, both dependent current sources and dependent voltage sources with arbitrary dependencies can also be treated.

The discussion will be restricted to resistive networks for simplicity and uniformity but applies directly to general R, L, C networks. In fact, networks with mutual inductance can be treated if the coupling is represented by appropriate dependent sources ${ }^{9}$. Also, networks containing only one independent source will be considered since the Principle of Superposition can be applied. Each element will be considered as a separate branch except for voltage sources (both dependent and independent) in series with a passive element.

Since Feussner's method and the Method of Residual Networks after R. Seacat ${ }^{7}$ are basic to the discussion, a brief summary of these will be given prior to the main results.

## 1. Feussner's Method and the Method of Residual Networks

Feussner's method provides an easy means for finding the co-trees of a passive network and since the determinant of the mesh-impedance matrix of a network yields the sum of the co-tree products, this provides a simple way of calculating the denominator of the response functions of the network.

The current response in the $k^{0}$ th mesh of an mesh network containing a single voltage source in mesh 1 is given by

$$
\begin{equation*}
I_{k}=E_{1} \Delta_{1 k / \Delta} \tag{1}
\end{equation*}
$$

where $\Delta$ is the network determinant and $\Delta_{1 k}$ is the cofactor obtained from $\Delta$ by crossing out the first row and the $k$ 'th column.

## Calculation of the Denominator

If $R_{a}$ appears only in terms on the main diagonal of the system determinant (this can always be done by a suitable choice of mesh currents), $\Delta$ can be expressed as

$$
\Delta=R_{a} \Delta_{a}+\Delta_{a^{\prime}}
$$

where $\Delta_{a}$ is the cofactor obtained by crossing out the row and column containing $R_{a}$ and $\Delta_{a}$, is the original determinant with $R_{a}$ equal to zero. Upon inspection of $\Delta_{a}$ it is apparent that it is the determinant of a network $N_{a}$ which is identical to the original network except that $R_{a}$ has been replaced by an open circuit. Also $\Delta_{a^{\prime}}$ is the determinant of a network $N_{a^{\prime}}$ which is the same as the original network with $R_{a}$ replaced by a short circuit.

Now $N_{a}$ and $N_{a}$, can be considered as new networks and their determinants can be expanded similarly, say with respect to element $R_{b}$. Thus

$$
\begin{equation*}
\Delta=R_{a} R_{b} \Delta_{a b}+R_{a} \Delta_{a b^{\prime}}+R_{b} \Delta_{a^{\prime} b}+\Delta_{a^{\prime} b^{\prime}} \tag{2}
\end{equation*}
$$

This process can be continued until the original network has been reduced to a set of primitive networks whose system determinants can be found by inspection. These primitive networks, their system determinant, and the topological graph corresponding to the principle determinant are shown in Figure 1. From Figure 1 it is apparent that a graph which is connected but which has no circuits has a determinant equal to unity (see graph (a), $\Delta_{a}$; and graph (c), $\Delta_{a}$ ); a graph which contains a "double short" has a zero determinant (see graph (a), $\Delta_{a^{0}}$; and graph (b), $\Delta_{a^{0}}$ ); and a graph which is separated has a zero determinant (see graph (e), $\Delta_{a^{p} b^{p} c^{p} \text { ). The source can be }}$ shorted without affecting the determination of $\Delta$.

In this manner it is possible and quite simple to find all the terms of the denominator of an admittance function directly. This is a self-checking method which eliminates the tedious cancellation of terms in the solution of the system determinant by conventional techniques.

For passive networks containing inductors and capacitors the Method of Residual Networks due to R. Seacat ${ }^{7}$ provides a means of directly calculating the denominator (and numerator) in the form of a polynomial, and yields any coefficient of the polynomial essentially by inspection. (This fact can be used nicely in active networks to ascertain the presence and dependence of any degeneracy due to the active elements as well as indicating ordinary degeneracies due to all-capacitor and all-inductor loops and cut-sets). In this method the expansion is done in terms of all the inductors and capacitors of the network by replacing each of these elements by switches. Since the expansion is done in terms of all the reactive elements, residual networks containing only resistors are left and Feussner's primitive-graph correspondences are used to ascertain the complete coefficient. If a particular L-switch is


$$
\begin{aligned}
\Delta & =R_{o} \\
\Delta_{a} & =1 \\
\Delta_{a^{\prime}} & =0
\end{aligned}
$$

(a)

(b)


$$
\begin{aligned}
& \Delta=R_{\mathrm{a}}+R_{b} \\
& \Delta_{\mathrm{a}}=\Delta_{\mathrm{b}}=1 \\
& \Delta_{\mathrm{a}^{\prime}}=R_{b}
\end{aligned}
$$

(c)

$\Delta=R_{a} R_{b} R_{c}$
$\Delta=R_{c}\left(R_{b}+R_{c}\right)+R_{b} R_{c}$
$\Delta_{0}=R_{b} R_{c}$
$\Delta_{a}=R_{b}+R_{c} ; \Delta_{a^{\prime}}=R_{b} R_{c}$
$\Delta_{a^{\prime}}=0$

open ( $L$ is extremized to infinity), $L$ appears in a product of terms in a coefficient, and doesn ${ }^{\circ} t$ appear if its switch is closed; alternatively if a particular $C$-switch is closed (C is extremized to infinity) $C$ appears in a coefficient and doesn ${ }^{\prime} t$ appear if its switch is opened. To make this clear, suppose we have a network containing two capacitors $C_{1}$ and $C_{2}$ and one inductor $\mathrm{L}_{3}$. The denominator will be of the form

$$
\begin{align*}
& =\frac{1}{s^{2} C_{1} c_{2}}\left[s^{3}\left\{c_{1} c_{2} L_{3} \Delta_{1^{\prime} 2^{\prime} 3^{\prime}}\right\}+s^{2}\left\{c_{1} c_{2} \Delta_{1^{\prime} 2^{\prime} 3^{\prime}}\right.\right. \\
& \left.+c_{1^{\prime}} L_{3^{\prime}} \Delta_{1^{\prime} 23}+c_{2} L_{3} \Delta_{12^{\prime} 3^{\prime}}\right\}+s\left\{c_{1^{\prime}} \Delta_{1^{\prime} 23^{\prime}}\right. \\
& \left.\left.+c_{2} \Delta_{12^{\prime} 3^{\prime}}+L_{3} \Delta_{123^{\prime}}\right\}+\Delta_{123^{\prime}}\right] \tag{3}
\end{align*}
$$

where for example $\Delta_{1^{\prime} 2^{\circ} 3}$ is the determinant of the residual network which is obtained by shorting switches $C_{1}$ and $C_{2}$ and opening switch $L_{3}$.

The term $S^{2} C_{1} C_{2}$ which appears in the denominator of this expression will also appear in the same manner in the numerator of all admittance functions and therefore can be disregarded.

## Calculation of the Numerator

The underlying theory for the method of calculation of the numerator given here, which is due to Kirchhoff, will be deferred and taken up in the next section in conjunction with the analysis of active networks and only the technique is given at this point.

This calculation involves the concept of a path current which is defined as a current flowing from the source through the network in a single closed loop. The algebraic sum of all path currents through a particular branch equals that branch current. The number of path currents in a network
due to a single source will in general be greater than or equal to the number of independent mesh currents.

Since this concept is best explained by an example, let us calculate the admittance function $I_{2} / E_{1}$ of the network of Figure $2 a$ which can be written as

$$
I_{2} / E_{1}=\Delta_{12} / \Delta .
$$

In this simple example only one path current flows through the source and branch 2 , and in completing its path it flows through $R_{1}, R_{3}$ and $R_{2} . \Delta$ can be expanded in terms of these three elements:

$$
\begin{equation*}
\Delta=R_{1} R_{2} R_{3} \Delta(123)+R_{1^{2}} R_{2} \Delta\left(123^{\prime}\right)+\ldots+R_{3^{\Delta}}{\left(11^{\prime} 2^{\prime} 3\right)}+\Delta \Delta_{\left(1^{\prime} 2^{\prime} 3^{\prime}\right)} \tag{4}
\end{equation*}
$$

Similarly $\Delta_{12}$ can be expanded with respect to these same elements:

$$
\begin{align*}
\Delta_{12}= & R_{1} R_{2} R_{3} \Delta_{12(123)}+R_{1} R_{2} \Delta_{12\left(123^{\prime}\right)}+\cdots+ \\
& R_{3} \Delta_{12\left(1^{\prime} 2^{\prime} 3\right)}+\Delta_{12\left(1^{\prime} 2^{\prime} 3^{\prime}\right)} \tag{5}
\end{align*}
$$

By definition of this method of expansion, it is obvious that the only term which contributes to $\Delta_{12}$ is the term $\Delta_{12\left(1^{\prime} 2^{\prime} 3^{\prime}\right)}$ which is the residual determinant obtained from the original network by shorting $R_{1}, R_{2}$ and $R_{3}$. All of the other terms indicate that one or more of these three resistors have been opened and therefore won't allow this path current to flow. Therefore to topologically determine the terms in $\Delta_{12}$ arising from a path current $J$, open the terminals of the source and short all elements in path $J$ and determine the determinant of the resultant primitive network using Feussner's method (see Figure 2c).


Figure 2

Thus the expression for any branch current is obtained by algebraically summing the residual contributions of each path current flowing through the branch in question.

So far nothing has been said about the signs of these terms. For uniformity all path currents will be drawn through the source from the negative to the positive terminal. Branch current directions are assigned, where possible, to agree with the direction that a positive $E$ would drive a positive current through a particular branch. Now the contribution of a single path current through the branch in question will be positive if its direction agrees with the branch current direction, negative otherwise.

The determination of path currents using the method of residual networks for passive R, L, C networks follows directly with all the Lswitches and C-switches in a path under consideration necessarily being shorted and the determinant of the residual network calculated as above.

## 2. Extension to Active Networks

We now justify the techniques presented in Section $I$ and extend them to active networks. The network to be considered has b branches, $b_{1}, b_{2}, \ldots, b_{b}$, and $n$ nodes, $a_{1}, a_{2}, \ldots, a_{m}, a_{n}$ with node $a_{n}$ to be taken as datum and where $m=n-1$. Branch $b_{j}$ will be denoted by $j$ whenever the context is such that no confusion will arise.

We will utilize, in the analysis, the matrix equation ${ }^{11}$

(6)
where $\mathbb{E}$ is the $m \times 1$ node voltage matrix
II is the $\mathrm{b} \times 1$ branch current matrix
A is the $m \times b$ reduced incidence matrix
$\mathbb{R}$ is the $\mathrm{b} x \mathrm{~b}$ branch resistance matrix which for a passive network is diagonal; for an $R, L, C, M$ network with dependent sources $\mathbb{R}$ is replaced by $\mathbb{Z}$ which is not diagonal in general.

$$
\text { Denoting the }(m \times b)^{\text {th }} \text { order matrix by } \mathbb{D} \text { we have }
$$



We want to find, topologically, the network junction $I_{i / E_{b}}$ which can be written as

$$
\begin{equation*}
\frac{I_{i}}{E_{b_{1}}}=\frac{\Delta_{b_{1} i}^{o} \pm \sum_{j} \mu_{j} \Delta_{b_{1}} i_{j} d_{j} \pm \cdots \pm \mu_{1} \mu_{2} \ldots \mu_{s} \Delta_{b_{1} i_{1} d_{1}}::: s d_{s}}{\Delta^{o} \pm \sum_{j} \mu_{j} \Delta_{j d_{j}} \pm \sum_{j<k} \mu_{j} \mu_{k} \Delta_{j d_{j}} k d_{k} \cdots \pm \mu_{1} \cdot \mu_{s} \Delta_{1 d_{1}} \ldots s d_{s}} \tag{8}
\end{equation*}
$$

where

$$
\begin{aligned}
\Delta^{\circ}= & |\mathbb{D}| \text { with all } \mu^{\prime} s \text { in } \mathbb{D} \text { (and in the network) set to zero; } \\
\Delta_{b_{1} i}^{\circ}= & \left|\mathbb{D}_{b_{1} i}\right| \text { where row } b_{1} \text { and column } b_{i} \text { are crossed out in } \mathbb{D} \\
& \text { and all } \mu^{\prime} \text { s set to zero; } \\
\Delta_{j d_{j}}= & \left|\mathbb{D}_{j d_{j}}\right| \text { where row } b_{j} \text { and column } d_{j} \text { are crossed out in } \mathbb{D} \\
& \text { and all } \mu \text {-source (dependent source) voltages set to zero } \\
& \quad \text { except } \mu_{j} I_{d_{j}} ; \text { etc. }
\end{aligned}
$$

The ambiguity of signs as well as the signs of terms within each cofactor will be resolved in the sequel.
A. Determination of $\Delta^{\circ}$
$\Delta^{0}=||D|$ with all dependent sources set to zero, (i.e. shorted) leaving a passive resistive network. Therefore $\mathbb{R}$ is diagonal and $\mathbb{D}$ can be expanded by Laplace's rule to yield the sum of co-tree products of the network. The details of this expansion will be omitted here since the topological determination of $\Delta^{\circ}$ has been discussed and justified in section $I$.
B. Determination of $E_{1} \Delta_{b_{1}}^{\circ}$ and $\mu_{j} \Delta_{j d_{j}}$
$\Delta_{b_{1} i}^{o}=\left|\left|D_{b_{1} i}\right|\right.$ with all dependent sources set to zero in the network and therefore corresponds to the numerator of the network junction $I_{i / E_{1}}$ for a purely passive network. Without loss of generality it will be assumed that the desired network junction is $I_{2 / E}$ and the cofactor to be found is $\Delta_{12}$.

It is apparent that the analysis for $E_{1} \Delta_{12}$ applies equally well to the denominator term $\pm \mu_{1} \Delta_{12}$ where the independent source and all dependent sources
have been set to zero except $\mu_{1} I_{2}$. Thus in general, the analysis to follow applies to the term $\pm \mu_{j} \Delta_{j d_{j}}$ which arises from a single dependent source in branch $b_{j}$ depending on current in branch $d_{j}=b_{k}$ for some $k$.

As indicated above, a dependent source voltage will be written as $\mu_{j} I_{j}$ which means that $\mu_{j}$ effectively has units of resistance. It is assumed for simplicity that $R_{j}$ and $R_{d_{j}}$ are nonzero although the theory will hold if either or both are zero. Also the voltage $\mu_{j} I_{d}$ is to appear in branch $b_{j}$ preceded by a plus sign with the voltage reference signs to be adjusted accordingly.

We expand $\left|\left|D_{12}\right|\right.$ by Laplace's rule by first crossing out row $b_{1}$ and column $b_{2}$ of $\mathbb{D}$. Using the first m rows of $\mathbb{D}_{12}$ (i.e. all the rows of $A \mathbb{}$ ) we consider all minors $D_{\alpha}$ of order $m$ that can be found from the rows and columns of $A^{(2)}$ which is $A$ with column $b_{2}$ crossed out. There are $\binom{b}{b}$ possible minors. Since the rows $a_{1}, a_{2}, \ldots, a_{m}$ are the same for all such minors we denote a particular minor $D_{\alpha}$ by the ordered m-tuple of its column numbers as given in $\mid D$ taken in order of increasing magnitude. Therefore $D_{\alpha}=\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{m}\right)$ where $\alpha_{i}=b_{j}$ for some $j$ and is actually the $\left(m+\alpha_{i}\right)^{\text {th }}$ column of $\mid D$; also row $a_{i}$ of $D_{\alpha}$ is the $i^{\text {th }}$ row of $\mathbb{D}$.

For each minor $D_{\alpha}$ there is a corresponding algebraic complement $(-1) \eta^{\alpha} D^{\alpha}$ obtained from $\mid D$ by crossing out the rows and columns chosen in $D_{\alpha}$, where $\eta$ is the sum of the numbers of the rows and columns contained in $D_{\alpha}$. Therefore

$$
\begin{equation*}
\left|\mathrm{D}_{12}\right|=(-1)^{\mathrm{b}_{1}+\mathrm{b}_{2}} \Sigma(-1)^{\eta} \mathrm{D}_{\alpha} \mathrm{D}^{\alpha} \tag{9}
\end{equation*}
$$

Now

$$
D^{\alpha}=\left|\left[A^{T}\right]_{(1)}\right| \mathbb{R}^{\alpha} \mid
$$

where $\mathbb{R}^{\alpha}$ is obtained by crossing out row $b_{1}$, column $b_{2}$ and columns $\alpha_{1}, \alpha_{2}, \ldots, \alpha_{m}$ of $\mathbb{R}$. It is more convenient to work with

$$
D^{\alpha}=\left|\begin{array}{|l}
\mathbb{A}^{(1)} \\
\hdashline \mathbb{R}_{\alpha}
\end{array}\right|
$$

where $\mathbb{R}_{\alpha}=\left[\mathbb{R}^{\alpha}\right]^{T}$ and can be obtained from $\mathbb{R}$ by crossing out column $b_{1}$, row $b_{2}$, and rows $\alpha_{1}, \alpha_{2}, \ldots, \alpha_{m}$. Applying Laplace's rule to the expansion of $D^{\alpha}$ considering all minors $D_{\beta}$ of order $m$ obtained from $A l^{(1)}$ we have

$$
D^{\alpha}=\sum_{\beta}(-1)^{\nu} D_{\beta} R_{\alpha}^{\beta}
$$

where $D_{\beta}=\left(\beta_{1}, \beta_{2}, \ldots, \beta_{m}\right)$ and $R_{\alpha}^{\beta}$ is the minor of $\mathbb{R}$ formed by crossing out rows $b_{2}, \alpha_{1}, \alpha_{2}, \ldots, \alpha_{m}$ and columns $b_{1}, \beta_{1}, \beta_{2}, \ldots, \beta_{m}$. As before, $\nu$ is the sum of the numbers of the rows and columns contained in $D_{\beta}$. Therefore

$$
\begin{equation*}
\Delta_{12}=(-1)^{b_{1}+b_{2}} \Sigma(-1)^{\eta+\nu} D_{\alpha} D_{\beta} R_{\alpha}^{\beta} \tag{10}
\end{equation*}
$$

Since $D_{\alpha}$ and $D_{\beta}$ are $m^{\text {th }}$ order determinants taken from the reduced incidence matrix and $R_{\alpha}^{\beta}$ is not diagonal in general for arbitrary choices of $D_{\alpha}$ and $D_{\beta}$, some of the terms in this expansion will be zero. Considering just one term in this summation we have

Proposition 1. $\quad D_{\alpha} D_{\beta} R_{\alpha}^{\beta}$ is nonzero only if $D_{\alpha}$ and $D_{\beta}$ are $m^{\text {th }}$ order minors which both correspond to trees of the network.

Proposition 2. $D_{\alpha} D_{\beta} R_{\alpha}^{\beta}$ is nonzero if and only if $D_{\alpha}$ corresponds to a tree and contains column $b_{1}, D_{\beta}$ corresponds to a tree containing branch (column) $b_{2}$ and the remaining $m-1$ columns of $D_{\alpha}$ are identical to the last $m-1$ columns of $D_{\beta}$ in the same order.

Proof: Crossing out row $b_{1}$ and column $b_{2}$ of $\mathbb{D}$ leaves a submatrix $\mathbb{R}_{12}$ which is not diagonal and has only resistive terms for its nonzero entries. The rows and columns of $\mathbb{R}_{12}$ which are subsequently eliminated appear as columns of $D_{\beta}$ and $D_{\alpha}$, respectively. Thus crossing out row $b_{2}$ and column $b_{1}$ leaves a diagonal submatrix whose determinant is therefore nonzero. Similarly for every row eliminated in $\mathbb{R}_{12}$, the corresponding column must also be crossed out if the resulting minor is to remain diagonal and therefore nonzero.

> Q.E.D.

Letting $\alpha_{1}=b_{1}$ and $\beta_{1}=b_{2}$ we have

$$
\mathrm{D}_{\alpha}=\left(\mathrm{b}_{1}, \alpha_{2}, \ldots, \alpha_{m}\right)
$$

and

$$
D_{\beta}=\left(b_{2}, \beta_{2}, \ldots, \beta_{m}\right)
$$

where $\alpha_{i}=\beta_{i}$ for $i \neq 1$. Also $R_{\alpha}^{\beta}$ is a positive product of ( $b-m-1$ ) resistances which appear in the tree complements of $D_{\alpha}$ and $D_{\beta}$.

In the case of a $\mu$-source depending on its own branch current, i.e., $\mu_{1} I_{1}$, a term in the summation is nonzero if and only if $D_{\alpha}$ and $D_{\beta}$ both correspond to trees and all their columns are identical. (This case can be regarded as passive with total resistance $R_{1} \pm \mu_{1}$ in branch $b_{1}$.)

Now with these conditions and the fact that the original column numbers have been used to designate $D_{\alpha}$ we have

$$
\begin{align*}
\eta & =\left[a_{1}+a_{2}+\ldots a_{m}\right]+\left[\left(m+b_{1}\right)+\left(m+\alpha_{2}-1\right)+\ldots+\left(m+\alpha_{m}-\right.\right. \\
& =\frac{m(m+1)}{2}+m^{2}+b_{1}+\left(\alpha_{2}-1\right)+\left(\alpha_{3}-1\right)+\ldots+\left(\alpha_{m}-1\right) \tag{11}
\end{align*}
$$

since the column numbers of $\mathbb{D}_{12}$ have decreased by one after column $b_{2}$ because $b_{2}$ has been crossed out. Similarly

$$
\begin{align*}
v & =\left[a_{1}+a_{2}+\ldots a_{m}\right]+\left[\left(b_{2}-1\right)+\left(\beta_{2}-1\right)+\ldots+\left(\beta_{m}-1\right)\right] \\
& =\frac{m(m+1)}{2}+\left(b_{2}-1\right)+\left(\beta_{2}-1\right)+\cdots+\left(\beta_{m}-1\right) \tag{12}
\end{align*}
$$

since row $b_{1}$ has been crossed out and therefore the row numbers of $D_{\beta}^{T}$ are decreased by one after row $b_{1}$ from their original numbers in $\mathbb{D}$. Thus since $\alpha_{i}=\beta_{i}, i=2,3, \ldots, m$ we have

$$
\begin{aligned}
(-1)^{\eta+\nu} & =(-1)^{m(m+1)+m^{2}+b_{1}+\left(b_{2}-1\right)+2\left(\alpha_{2}-1\right)+2\left(\alpha_{3}-1\right)+\ldots 2\left(\alpha_{m}-1\right)} \\
& =(-1)^{m^{2}}(-1)^{b_{1}+b_{2}-1}
\end{aligned}
$$

It is apparent that the factor $(-1)^{\mathrm{m}^{2}}$ will appear in every term in both numerator and denominator of the network function and will therefore cancel. However, it is convenient for the rest of the discussion if $m$ is assumed an even number which means $(-1)^{m^{2}}=+1$. Therefore the $t$ erm becomes

$$
\begin{align*}
(-1)^{b_{1}+b_{2}(-1)^{\eta+\nu} D_{\alpha} D_{\beta} R_{\alpha}^{\beta}} & =(-1)^{b_{1}+b_{2}(-1)^{b_{1}+b_{2}-1} D_{\alpha} D_{\beta} R_{\alpha}^{\beta}} \\
& =-D_{\alpha} D_{\beta} R_{\alpha}^{\beta} \tag{13}
\end{align*}
$$

Now all that remains is to determine the sign of $D_{\alpha} D_{\beta}$. Since $D_{\alpha}= \pm 1$ and $D_{\beta}= \pm 1$ it is sufficient to determine the sign of $D_{\beta}$ with respect to $D_{\alpha}$.

Noting that $D_{\alpha}=\left(b_{1}, \alpha_{2}, \alpha_{3}, \ldots, \alpha_{m}\right)$ and $D_{\beta}=\left(b_{2}, \alpha_{2}, \alpha_{3}, \ldots, \alpha_{m}\right)$, when the two trees corresponding to these minors are coupled, a subnetwork is obtained which has exactly one loop containing both branch $b_{1}$ and branch $b_{2}$. Call this loop, $J$, and let its branches be labeled $b_{1}, b_{2}, \gamma_{1}, \gamma_{2}, \ldots, \gamma_{p}$ where $p \leq m-1$. The branches $\gamma_{1}, \gamma_{2}, \ldots, \gamma_{p}$ appear as columns in both $D_{\alpha}$ and $D_{\beta}$.

In $D_{\beta}$, add column $\gamma_{j}(j=1,2, \ldots, p)$ to column $b_{2}$ if branch currents $I_{b_{2}}$ and $I_{\gamma_{j}}$ have the same directions around the loop; subtract $\gamma_{j}$ from $b_{2}$ otherwise. The result of these operations which don't change the value of $D_{\beta}$ yields a column in position 1 which is plus or minus column $b_{1}$, the leading column of $D_{\alpha}$. This first column of $D_{\beta}$ equals minus $b_{1}$ if and only if $I_{b_{1}}$ is in the same direction as $I_{b_{2}}$ around the loop and in this case $D_{\beta}=-D_{\alpha}$ and $D_{\alpha} D_{\beta}=-1$; otherwise $D_{\alpha}=D_{\beta} \cdot D_{\beta}=+1$. Thus since $(-1)^{b_{1}+b_{2}}(-1)^{\eta+\nu}=-1$, equation 13 becomes $+R_{\alpha}^{\beta}$ if branch currents $I_{b_{1}}$ and $I_{b_{2}}$ have the same sense around loop $J$ and is $-R_{\alpha}^{\beta}$ otherwise.

Now the method given in Section 1 follows directly from the above and the terms of $E_{1} \Delta_{12}$ can be found topologically: open the terminals of $E_{1}$ and draw a path current through $E_{1}$ and $R_{2}$; set all elements in this path to zero (short-circuit them) and find the residual network terms by Feussner's method. The sign of these terms is plus if both $I_{1}$ and $I_{2}$ are in the same sense around the loop; otherwise the sign is minus. Since branch current directions are assigned in a natural way with respect to $E_{1}$ without loss of generality, the sign determination can be simplified. Draw the path current $J$ from - to + through $E_{1}$ traversing branch $b_{2}$; the sign of the terms is positive if the path current direction agrees with the direction of $I_{2}$; otherwise the sign is negative. This eliminates the comparing of directions of the two branch currents $I_{1}$ and $I_{2}$.

Drawing all possible path currents through the opened terminals of $E_{1}$ traversing branch $\mathrm{b}_{2}$ one at a time, finding the residual terms as above and adding these terms algebraically yields $E_{1} \Delta_{12}^{\circ}$.

By replacing $E_{1}$ with $\left(-\mu_{1}\right)$ we have the denominator term $-\mu_{1} \Delta_{12}$ due to a single dependent source $\mu_{1} I_{2}$. This replacement is valid since, by the
convention of assigning branch current directions, $+\mathrm{E}_{1}$ appears in position $d_{b_{1}} b_{2}$ of $\mid D$ in the determination of $I_{2}$ by Cramer's rule whereas a voltage rise across the dependent source in the direction of positive $I_{1}$ appears in this same position as $-\mu_{1}$.

Thus the ambiguity of sign is eliminated in equation 8 . That is, the sign preceding a cofactor in either numerator of denominator is + if the cofactor is multiplied by an even number of $\mu^{\prime} s$ and is minus otherwise.
C. Determination of $\left(-\mu_{j}\right)\left(-\mu_{k}\right) \Delta_{j d_{j}} k d_{k}$ and $E_{1}\left(-\mu_{k}\right) \Delta_{1 i k d_{k}}$

$$
\Delta_{j d_{j} k d_{k}}=\left|\left|D_{j d_{j} k d_{k}}\right|\right.
$$

where all but two dependent source voltages, $\mu_{j} I_{d}$ and $\mu_{k} I_{d_{k}}$, have been set to zero in the network. There are various configurations for these two $\mu$-sources:

1. both $\mu$-sources in distinct branches, each depending on distinct branch currents;
2. both $\mu$-sources in distinct branches, one depending on current in a distinct branch and the other depending on the branch current through either of the two $\mu$-sources;
3. both $\mu$-sources in distinct branches each depending on current in their own branches or on the current in each other's branch. This cofactor is zero if both $\mu$-sources are in the same branch since $j=k$ and both appear in the same row of $\mid D$. Also the cofactor is zero if both depend on current in the same branch since then $d_{j}=d_{k}$ and both $\mu_{j}$ and $\mu_{k}$ appear in the same column of $1 D$.

Case 1: Again, without loss of generality, we can relabel the branches and consider $\mu_{1} I_{3}$ and $\mu_{2} I_{4}$ and the cofactor $\Delta_{1324}$ where rows $b_{1}$ and $b_{2}$ and columns $b_{3}$ and $b_{4}$ have been eliminated from ID. Expanding $\Delta_{1324}$ by Laplace's rule, we have as before

$$
\begin{equation*}
\left(-\mu_{1}\right)\left(-\mu_{2}\right) \Delta_{1324}=\mu_{1} \mu_{2}(-1)^{b_{1}+b_{2}+b_{3}+b_{4} \sum_{\alpha, \beta}(-1)^{\eta+\nu} D_{\alpha} D_{\beta} R_{\alpha}^{\beta} . . . ~ . ~} \tag{14}
\end{equation*}
$$

The same reasoning used in Proposition 2 applies to this expansion also and we have

Proposition 3: $D_{\alpha} D_{\beta} R_{\alpha}^{\beta}$ is nonzero if and only if $D_{\alpha}$ corresponds to a tree and contains columns (branches) $b_{1}$ and $b_{2}$ and $D_{\beta}$ corresponds to a tree which contains branches $b_{3}$ and $b_{4}$, the other $m-2$ columns of $D_{\alpha}$ and $D_{\beta}$ being identical.
$R_{\alpha}^{\beta}$ now is a product of ( $b-m-2$ ) resistors which appear in the complement of $D_{\alpha}$ and $D_{\beta}$ and

$$
\begin{aligned}
& D_{\alpha}=\left(b_{1}, b_{2}, \alpha_{3}, \ldots, \alpha_{m}\right) \\
& D_{\beta}=\left(b_{3}, b_{4}, \beta_{3}, \ldots, \beta_{m}\right)
\end{aligned}
$$

where

$$
\begin{align*}
& \beta_{i}=\alpha_{i}=b_{j} \text { for some } j \neq 1,2,3,4 \text {. Also } \\
& \eta=\frac{m(m+1)}{2}+\left(m+b_{1}\right)+\left(m+b_{2}\right)+\left(m+\alpha_{3}-2\right)+\ldots+\left(m+\alpha_{m}-2\right) \tag{15}
\end{align*}
$$

and

$$
\begin{equation*}
v=\frac{m(m+1)}{2}+\left(b_{3}-2\right)+\left(b_{4}-2\right)+\left(\beta_{3}-2\right)+\cdots+\left(\beta_{m}-2\right) \tag{16}
\end{equation*}
$$

Therefore $(-1)^{\eta+\nu}=(-1)^{m^{2}}(-1)^{b_{1}+b_{2}+\left(b_{3}-2\right)+\left(b_{4}-2\right)}$ and since $m$ is assumed even, we have

$$
\begin{equation*}
(-1)^{b_{1}+b_{2}+b_{3}+b_{4}}(-1)^{\eta+\nu} D_{\alpha} D_{\beta} R_{\alpha}^{\beta}=+D_{\alpha} D_{\beta}^{R} R_{\alpha}^{\beta} \tag{17}
\end{equation*}
$$

We again consider the sign of $D_{\beta}$ with respect to $D_{\alpha}$ in the sign determination of $D_{\alpha} D_{\beta}$. Because of the characteristics of $D_{\alpha}$ and $D_{\beta}$, it is apparent that when their corresponding trees are coupled, a subnetwork containing two loops is obtained. Actually more than two loops may exist but there are only two with a single $\mu$-source in each. Both branches $b_{3}$ and $b_{4}$ may appear in the same loop with a single $\mu$-source but then only one of $b_{3}$ or $b_{4}$ can belong to the second loop. Branches $b_{1}$ and $b_{2}$ cannot both appear in two loops, and branches $b_{3}$ and $b_{4}$ cannot both appear in two loops. These last two possibilities have interesting interpretations when considered topologically as will be seen.

Therefore there are effectively two possibilities to consider for $D_{\alpha}$ and $D_{\beta}$ nonzero:
(a) one loop, $J_{1}$, contains branches $b_{1}$ and $b_{3}$; the second loop, $J_{2}$, contains $\mathrm{b}_{2}$ and $\mathrm{b}_{4}$. Call this the normal case.
(b) one loop, $J_{1}^{p}$, contains branches $b_{1}$ and $b_{4}$; the second loop, $J_{2}^{p}$, contains $b_{2}$ and $b_{3}$. Call this the permuted case.
(a) Normal case: Loop $J_{1}$ contains branches $b_{1}, b_{3}$ and $\delta_{1}, \delta_{2}, \ldots, \delta_{p}$, $\mathrm{p} \leq \mathrm{m}-2$; loop $\mathrm{J}_{2}$ contains branches $\mathrm{b}_{2}, \mathrm{~b}_{4}$ and $\gamma_{1}, \gamma_{2}, \ldots, \gamma_{q}, \mathrm{q} \leq m-2$ where $\delta_{i}=b_{j}$ for some $b_{j} \in\left\{b_{4}, \alpha_{3}, \alpha_{4}, \ldots, \alpha_{m}\right\}, i=1,2, \ldots, p$ and $\gamma_{i}=b_{k}$ for some $b_{k} \in\left\{b_{3}, \alpha_{3}, \alpha_{4}, \ldots, \alpha_{m}\right\} \quad i=1,2, \ldots, q$. It is possible that $\delta_{i}=b_{4}$ or $\gamma_{i}=b_{3}$, not both.

$$
\begin{aligned}
& D_{\alpha}=\left(b_{1}, b_{2}, \alpha_{3}, \ldots, \alpha_{m}\right) \\
& D_{\beta}=\left(b_{3}, b_{4}, \alpha_{3}, \ldots, \alpha_{m}\right)
\end{aligned}
$$

In $D_{\beta}$ add column $\delta_{i}(i=1,2, \ldots, p)$ to column $b_{3}$ if the branch currents $I_{b_{3}}$ and $I_{\delta_{i}}$ have the same sense in loop $J_{1}$; subtract it otherwise. The resulting column in the first position of $D_{\beta}$ is the same as column $b_{1}$, the first column of $D_{\alpha}$, if $I_{b_{3}}$ and $I_{b_{1}}$ have opposite senses in loop $J_{1}$ and is the negative of column $b_{1}$ otherwise. Similarly, adding or subtracting column $\gamma_{i}(i=1,2, \ldots, q)$ to column $b_{4}$ of $D_{\beta}$ yields plus column $b_{2}$ in this second column position of $D_{\beta}$ if $I_{b_{4}}$ and $I_{b_{2}}$ have opposite senses around loop $J_{2}$ and yields minus column $\mathrm{b}_{2}$ otherwise.

Thus if the branch current pair $I_{b_{1}}, I_{b_{3}}$ have the same direction around loop $J_{1}$ and the pair $I_{b_{2}}, I_{b_{4}}$ have the same direction around loop $J_{2}$ or if both pairs have opposite senses in the respective loops, then $D_{\alpha}=D_{\beta}$, $D_{\alpha} D_{\beta}=+1$ and from equation 17 the resulting term is $+R_{\alpha}^{\beta}$; otherwise $D_{\beta}=-D_{\alpha}$ and the term is $-R_{\alpha}^{\beta}$.

Applying the simplification discussed for the one $\mu$-source case, the topological determination follows directly from the above analysis: open the terminals of both $\mu$-sources and draw path current $J_{1}$ from - to + through $\mu_{1}$, going through $R_{3}$ and path current $J_{2}$ from - to + through $\mu_{2}$ and going through $\mathrm{R}_{4}$; set all elements in both paths to zero and calculate the residual terms using Feussner ${ }^{\circ}$ s method. The sign of these terms is positive if $J_{1}$ traverses $R_{3}$ in the same direction as $I_{b_{3}}$ and $J_{2}$ traverses $R_{4}$ in the same direction as $\mathrm{I}_{b_{4}}$ or if both $J_{1}$ and $J_{2}$ have opposite directions with respect to $I_{b_{3}}$ and $\mathrm{I}_{\mathrm{b}}$ respectively; otherwise the sign is negative. Drawing all possible path current pairs, $J_{1}$ and $J_{2}$, finding the residual terms corresponding to each pair and adding these set of terms algebraically yields terms in $\Delta_{1324^{\circ}}$

This procedure does not necessarily yield all the terms in $\Delta_{1324}$ as will be seen in the discussion of the permuted case. Therefore identify all terms obtained from this normal case as $\Delta_{1324^{\circ}}^{\mathrm{N}}$
(b) Permuted Case: When the trees of $D_{\alpha}$ and $D_{\beta}$ are coupled a subnetwork with (at least) two loops is obtained which may associate branches $b_{1}$ and $b_{4}$ with loop 1 and branches $b_{2}$ and $b_{3}$ with loop 2. Therefore terms in $\Delta_{13,24}$ can be thought of as arising either from the $\mu$-source pair $\mu_{1} I_{3}$, $\mu_{2} I_{4}$ or from the pair, $\mu_{1} I_{4}, \mu_{2} I_{3}$. In this case we write

$$
\begin{aligned}
& D_{\alpha}=\left(b_{1}, b_{2}, \alpha_{3}, \ldots, \alpha_{m}\right) \\
& D_{\beta}^{P}=\left(b_{4}, b_{3}, \alpha_{3}, \ldots, \alpha_{m}\right)
\end{aligned}
$$

where the first and second columns of $D_{\beta}$ have been permuted to yield $D_{\beta}^{P}$.
Now the procedure for comparing $D_{\beta}^{P}$ with $D_{\alpha}$ is the same as in the normal case and the resulting sign of $D_{\alpha} D_{\beta}^{P}$ is exactly the negative of the sign of $D_{\alpha} D_{\beta}$ determined in the normal case since some of the columns of $D_{\beta}$ have been permuted an odd number of times.

Thus the topological determination of these terms is effected by opening the terminals of both $\mu$-sources and drawing a path current $J_{1}{ }^{P}$ from - to + through $\mu_{1}$ traversing branch $b_{4}$ and a path current $J_{2}{ }^{P}$ from - to + through $\mu_{2}$ traversing $b_{3}$; the rest of the procedure is the same as for the normal case. The result is a set of signed terms which will be denoted as $\Delta_{1324}^{\mathrm{P}}$ and the effect of the (odd) permutation is accounted for by multiplying this "cofactor" by -1 .

An alternate way of describing the set of terms obtained in this case is

$$
-\Delta_{13,24}^{\mathrm{P}}=+\Delta_{14,23}^{\mathrm{N}}
$$

and $\Delta_{14,23}^{N}$ can be determined from the normal case analysis associated with the $\mu$-source pair of $\mu_{1} I_{4}$ and $\mu_{2} I_{3}$. (The sign of $(-1)^{\eta+\nu}$ changes if this idea is used and the two descriptions are seen to be equivalent.)

Thus a $\mu$-path current can be thought of as a loop containing the $\mu$-source and the branch from which it takes its dependence.

The sets of terms found in (a) and (b) constitute all of the terms in the cofactor. In the topological determination of $\Delta_{1324}$, there is a possibility that the same term appears in both $\Delta_{1324}^{N}$ and $\Delta_{1324}^{P}$ but these terms will always cancel since this case arises when $b_{3}$ and $b_{4}$ appear in both loops $J_{1}$ and $J_{2}$ and thus appear in both $J_{1}^{P}$ and $J_{2}^{P}$ also with relative current directions remaining the same. This agrees with the analysis of $\Delta_{13,24}$ by Laplace's rule since this term has $D_{\alpha} D_{\beta}=0$ and therefore doesn't appear in the expansion. Thus when both $b_{3}$ and $b_{4}$ appear in $J_{1}$ and in $J_{2}$ the resulting term can be neglected by inspection. The similar possibility of $b_{1}$ and $b_{2}$ appearing in two loops does not arise topologically if we agree that a path current due to one $\mu$-source cannot travel through another $\mu$-source (or E-source) because the terminals of the $\mu$-sources have been opened.

For the dependent source configuration of Case 2, we consider $\Delta_{12,23}$ or $\Delta_{13,22}\left(\right.$ recall $\Delta_{12,23}^{\mathrm{N}}=-\Delta_{13,22}^{\mathrm{P}}$ ). $\mathrm{D}_{\alpha}$ now must contain column $b_{1}, D_{\beta}$ must contain column $b_{3}$, and the remaining m-1 columns of $D_{\alpha}$ and $D_{\beta}$ must be identical. In Case 3, only $\mu_{1} I_{1}, \mu_{2} I_{2}$ need be considered which results in the cofactor $\Delta_{11,22}\left(\Delta_{12,21}^{N}=-\Delta_{11,22}^{\mathrm{P}}\right)$. In this case all columns of $D_{\alpha}$ and $D_{\beta}$ must be identical. The analysis for both Case 2 and Case 3 is essentially the same as for Case 1.

Now we can find topologically the network determinant of a network containing two dependent sources:

$$
\begin{equation*}
\Delta=\Delta^{\circ}-\mu_{1} \Delta_{13}-\mu_{2} \Delta_{24}+\mu_{1} \mu_{2}\left[\Delta_{1324}^{N}-\Delta_{1324}^{P}\right] \tag{18}
\end{equation*}
$$

where both $\mu_{1} \equiv 0$ and $\mu_{2} \equiv 0$ to find $\Delta^{\circ} ; \mu_{2} \equiv 0$ for the determination of $\mu_{1} \Delta_{13}$, etc.

As in the one $\mu$-source case, if $\left(-\mu_{1}\right)$ is replaced by $\left(+E_{1}\right)$ the above analysis applies directly to the terms in the numerator of the branch current $I_{3}$ which are multiplied by $\mu_{2}$. Thus the network function $I_{3 / E}$ with one dependent voltage source, $\mu_{2} I_{4}$, is given by

$$
\begin{equation*}
\frac{I_{3}}{E_{1}}=\frac{\Delta_{13}^{o}-\mu_{2} \Delta_{1324}}{\Delta^{o}-\mu_{2} \Delta_{24}} \tag{19}
\end{equation*}
$$

The general case of $s$ dependent voltage sources present in a network is treated in essentially the same manner as for two $\mu$-sources. Consider the topological determination of

$$
\begin{equation*}
\left(-\mu_{1}\right)\left(-\mu_{2}\right) \ldots\left(-\mu_{s}\right) \Delta_{1 d_{1}} 2 d_{2} 3 d_{3} \ldots s d_{s} \tag{20}
\end{equation*}
$$

where $\mu_{1} I_{d_{1}}, \mu_{2} I_{d_{2}}, \ldots, \mu_{s} I_{d}$ are the dependent source voltages. We will assume that $s \leq b-n+1$ since for $a$ network with $n$ nodes and $b$ branches, the network determinant consists of terms which are products of $b-n+1$ elements.

To find $\Delta_{1 d_{1}}, 2 d_{2}, \ldots, s d_{s}$ all $s$ ! permutations of the branches $d_{i}(i=1,2, \ldots, s)$ have to be considered. Since there are $\frac{s!}{2}$ even and $\frac{s!}{2}$ odd permutations, let $P_{i}$ denote an even permutation if $i$ is even; an odd permutation otherwise. Thus

$$
\begin{equation*}
\Delta_{1 d_{1} 2 d_{2} \ldots s d_{s}}=\sum_{i=0}^{s!-1}(-1)^{i} \Delta_{1 d_{1}}^{P_{i}} d_{2} \ldots s d_{s} \tag{21}
\end{equation*}
$$

where

$$
\Delta_{1 d_{1} \ldots s d_{s}}^{P_{o}}=\Delta_{1 d_{1} \ldots s d_{s}}^{N}
$$

$$
s!-1
$$

There are a total of $s!\sum_{i=0} N_{i}$ path current combinations that have to be checked where $N_{i}$ is the number of sets of path current possibilities for the $i^{\text {th }}$ permutation of $\mu$-source dependences.

The sign of a term within any of the permutation-cofactors is negative if an odd number of path currents traverse their respective permuted dependent branches in the opposite sense of the assumed branch current direction; otherwise it is positive.

By replacing $\left(-\mu_{1}\right)$ by $E_{1}$ we obtain the terms in the numerator function $I_{d_{1} / E_{1}}$ which are multiplied by $\mu_{2} \mu_{3} \ldots \mu_{s}$. Thus the complete topological determination of a network function $I_{i / E}$ with a $s-1$ dependent voltage sources present is at hand. The proof for the general case is necessarily rather involved and will be omitted here.

## CONCLUSION

The method given in this paper provides an easy means of analyzing network functions of arbitrary networks and facilitates the determination of the effect of any circuit element on the network. Specifically the effect of dependent sources can be ascertained immediately and independently without the necessity of calculating the entire network function.

Further investigation in this area would seem to lie in the direction of utilizing the insight and simplicity offered by this method in network synthesis.

The following example clarifies the arguments given above and illustrates the simplification effected by this method.

## Example



We wish to find the network function

$$
Y_{16}=\frac{I_{6}}{E_{1}}=\frac{\Delta_{16}^{\circ}-\mu_{2} \Delta_{1624}-\mu_{3} \Delta_{1635}+\mu_{2} \mu_{3} \Delta_{162435}}{\Delta^{o}-\mu_{2} \Delta_{24}-\mu_{3} \Delta_{35}+\mu_{2} \mu_{3} \Delta_{2435}}
$$

$\Delta^{\circ}:$
Expand by $\mathrm{R}_{1}$ and $\mathrm{R}_{2}$

$$
\Delta^{\circ}=R_{1} R_{2} \Delta_{(12)}+R_{1^{\prime}} \Delta_{(12 \prime)}+R_{2} \Delta_{\left(1^{\prime} 2\right)}+\Delta_{\left(1^{\prime} 2^{\prime}\right)}
$$



$$
\Delta_{(12)}=R_{3}+R_{4}+R_{5}+R_{6}
$$



$$
\Delta_{\left(12^{\prime}\right)}=\left(R_{3}+R_{4}\right)\left(R_{5}+R_{6}\right)
$$



$$
\Delta_{\left(1^{\prime} 2\right)}=\left(R_{3}+R_{5}\right)\left(R_{4}+R_{6}\right)
$$



$$
\Delta{\left(1{ }^{\prime} 2^{\prime}\right)}=R_{3}\left[R_{4}\left(R_{5}+R_{6}\right)+R_{5} R_{6}\right]+R_{4} R_{5} R_{6}
$$

Therefore

$$
\begin{gathered}
\Delta^{0}=R_{1} R_{2}\left(R_{3}+R_{4}+R_{5}+R_{6}\right)+R_{1}\left(R_{3}+R_{4}\right)\left(R_{5}+R_{6}\right)+R_{2}\left(R_{3}+R_{5}\right)\left(R_{4}+R_{6}\right) \\
+R_{3} R_{4}\left(R_{5}+R_{6}\right)+R_{5} R_{6}\left(R_{3}+R_{4}\right)
\end{gathered}
$$

$-\mu_{2} \Delta_{24}:$


$$
\Delta_{24}^{(1)}=-\left[R_{1}\left(R_{5}+R_{6}\right)+R_{5} R_{6}\right]
$$


$\Delta_{24}^{(2)}=-R_{3} R_{6}$

Therefore

$$
-\mu_{2} \Delta_{24}=\mu_{2}\left[R_{1}\left(R_{5}+R_{6}\right)+R_{6}\left(R_{3}+R_{5}\right)\right]
$$

$$
-\mu_{3} \Delta_{35}:
$$



$$
\Delta_{35}^{(1)}=-\left[R_{2}\left(R_{4}+R_{6}\right)+R_{4} R_{6}\right]
$$



$$
\Delta_{35}^{(2)}=-R_{1} R_{2}
$$

Therefore

$$
-\mu_{3} \Delta_{35}=+\mu_{3}\left[R_{2}\left(R_{1}+R_{4}\right)+R_{6}\left(R_{2}+R_{4}\right)\right]
$$

$\mu_{2}{ }^{\mu}{ }_{3}{ }_{2435}:$

$\Delta_{2435}^{N}(1)=+R_{6}$

$(-1)^{1} \Delta_{2435}^{P}=(-1)\left(-R_{1}\right)=R_{1}$

Therefore

$$
\mu_{2} \mu_{3} \Delta_{2435}=\mu_{2} \mu_{3}\left(R_{1}+R_{6}\right)
$$

$\Delta_{16}^{0}:$


$$
\Delta_{16}^{(2)}=+R_{4} R_{5}
$$

Therefore

$$
\Delta_{16}^{\circ}=R_{2}\left(R_{3}+R_{5}\right)+R_{5}\left(R_{3}+R_{4}\right)
$$

$$
-\mu_{2} \Delta_{1624}:
$$



$$
\Delta_{1624}^{\mathrm{N}}(1)=-\mathrm{R}_{5}
$$



Therefore

$$
-\mu_{2} \Delta_{1624}=\mu_{2}\left[R_{3}+R_{5}\right]
$$

$-\mu_{3} \Delta_{1635}:$


$$
\Delta_{1635}^{\mathrm{N}}(1)=-\mathrm{R}_{2}
$$


(-1) $\Delta_{1635}^{P}(1)=0$

$$
-\mu_{3} \Delta_{1635}=+\mu_{3} R_{2}
$$

$\mu_{2} \mu_{3}{ }^{\Delta}{ }_{162435}:$
By inspection $\quad \Delta_{162435}^{\mathrm{N}}=\Delta_{162435}^{\mathrm{P}_{4}}=\Delta_{162435}=0$


Therefore

$$
\mu_{2} \mu_{3} \Delta_{162435}=\mu_{2} \mu_{3}[1+1-1]=\mu_{2} \mu_{3}
$$

Finally

$$
\frac{I_{6}}{E_{1}}=\frac{R_{2}\left(R_{3}+R_{5}\right)+R_{5}\left(R_{3}+R_{4}\right)+\mu_{2}\left(R_{3}+R_{5}\right)+\mu_{3} R_{2}+\mu_{2} \mu_{3}}{\Delta^{0}+\mu_{2}\left[R_{1}\left(R_{5}+R_{6}\right)+R_{6}\left(R_{3}+R_{5}\right)\right]+\mu_{3}\left[R_{2}\left(R_{1}+R_{4}\right)+R_{6}\left(R_{2}+R_{4}\right)\right]+\mu_{2} \mu_{3}\left(R_{1}+R_{6}\right)}
$$

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