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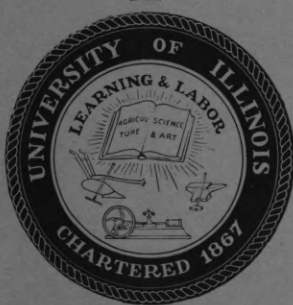
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# Coordinated Science Laboratory



UNIVERSITY OF ILLINOIS - URBANA, ILLINOIS

# TUNNEL DIODE DETECTOR

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REPORT R-138

MAY, 1962

COORDINATED SCIENCE LABORATORY  
UNIVERSITY OF ILLINOIS  
URBANA, ILLINOIS

Contract DA-36-039-SC-85122  
D/A Sub-Task 3-99-01-002

The research reported in this document was made possible by support extended to the University of Illinois, Coordinated Science Laboratory, jointly by the Department of the Army (Signal Corps and Ordnance Corps), Department of the Navy (Office of Naval Research), and the Department of the Air Force (Office of Scientific Research, Air Research and Development Command) under Signal Corps Contract DA-36-039-SC-85122.

## ABSTRACT

By use of a tunnel diode, it is possible to design a detector which is capable of amplifying a detected signal. Suppose input signal is an amplitude modulated signal whose modulating signal is  $e_s \sin \omega_i t$  and the degree of modulation is very small. From the experimental results, the insertion gain

$$g_s = \left| \frac{e_o}{e_o'} \right|$$

is larger than 14 where  $e_o$  is the peak value of output signal corresponding to  $e_s$  with a tunnel diode being active and  $e_o'$  is the peak value of output signal corresponding to  $e_s$  with the tunnel diode being shorted.



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## 1. INTRODUCTION

The use of only the negative resistance region of V-I characteristic of a tunnel diode, amplifiers<sup>1-8</sup> and oscillators<sup>9-12</sup> can be designed. By using only the positive regions, tunnel diodes can be used as bistable devices.<sup>13-16</sup> If both positive and negative resistance regions are used, a tunnel diode can be used as a detector which is capable of amplifying detected signal.<sup>17-18</sup>

This paper shows the principle of tunnel diode detectors and analytic and experimental results of such a detector with commercially available tunnel diodes.

## 2. PRINCIPLE OF TUNNEL DIODE DETECTOR

There are two typical circuits for tunnel diode amplifiers.<sup>3</sup> One is the series amplifier and the other is the parallel amplifier as shown in Fig. 1 a and b respectively.

Since the amplification by tunnel diodes is possible only because of the negative resistance of the diode, the biasing point must be set on the middle of the negative resistance region so that in each half cycle of input signal, the operating point A swings inside the negative resistance region as shown in Fig. 2a.

The idea of our detector circuit is to set the biasing point at the end of the negative characteristic, B or C, as shown in Figure 2b, so that this circuit acts as amplifier at a half cycle of input signal and acts as attenuator at the other half cycle.

In the paper the circuit (a) in Fig. 1 is used to perform detection by the proper setting of biasing voltage. However, it is also possible to use the circuit (b) in Fig. 1 to use as a detector. An equivalent circuit for such a detector is shown in Fig. 3.

## 3. GAIN OF DETECTORS

Suppose the bias point of the diode in the circuit in Fig. 3 is set at the point B in Fig. 2. Then the gain of the circuit can approximately be calculated by considering the characteristic of the diode as the inverse of capital letter V as shown in Fig. 4a. Furthermore by proper transformation, the characteristic of the diode can be assumed to be the one shown in Fig. 4b and the circuit in Fig. 3 becomes the circuit in Fig. 5.

When the operating point is in the negative resistance region, the following equations from the circuit in Fig. 4b can be obtained:

$$\left. \begin{aligned} v_d^+ &= -iR_t - L_t \frac{di}{dt} + v_{in} \\ i^+ &= C_d \frac{dv_d}{dt} - \frac{v_d}{R_d} \end{aligned} \right\} \quad (1)$$

for  $v_d \geq 0$  where

$$v_{in} = V_i \sin \omega t \quad (2)$$

and

$$R_t = R'_g + R_L \quad (3)$$

From these,

$$\frac{d^2 v_d^+}{dt^2} + \left( \frac{R_t}{L_t} - \frac{1}{R_d C_d} \right) \frac{dv_d^+}{dt} + \frac{R_d - R_t}{L_t C_d R_d} v_d^+ = \frac{V_i \sin \omega t}{L_t C_d} \quad (4)$$

The solution  $v_d^+$  of Eq. 4 is

$$v_d^+ = Ae^{\alpha_1 t} + Be^{\alpha_2 t} + \frac{(V_i/L_t C_d) \sin(\omega t - \beta)}{\sqrt{(b - \omega^2)^2 + \omega^2 a^2}} \quad \text{for } v_d \geq 0 \quad (5)$$

where A, B are arbitrary constants

$$\alpha_1, \alpha_2 = \frac{1}{2} (-a \pm \sqrt{a^2 - 4b}) \quad (6)$$

$$a = \frac{R_t}{L_t} - \frac{1}{R_d C_d} \quad (7)$$



$$b = \frac{R_d - R_t}{L_t C_d R_d} \quad (8)$$

and

$$\beta = \tan^{-1} \frac{\omega a}{b - \omega^2} \quad (9)$$

From the Eq.'s 9 and 12,

$$i^+ = (C_d \alpha_1 - \frac{1}{R_d}) A e^{\alpha_1 t} + (C_d \alpha_2 - \frac{1}{R_d}) B e^{\alpha_2 t} \quad (10)$$

$$+ \frac{\frac{V_i}{L_t C_d}}{\sqrt{(b - \omega^2)^2 + \omega^2 a^2}} \left\{ \omega \cos(\omega t - \beta) - \frac{1}{R_d} \sin(\omega t - \beta) \right\}$$

and output voltage is obtained by  $i^+ \times R_L$ .

$$v_{out}^+ = i^+ \times R_L \quad (11)$$

When the operating point is in the positive resistance region, the following equations can be obtained by replacing  $-R_d$  by  $R_d$  in Eq.'s 2 and 11.

Therefore,

$$v_d = C \cdot e^{\alpha_1' t} + D \cdot e^{\alpha_2' t} + \frac{\frac{V_i}{L_t C_d} \sin(\omega t - \beta')}{\sqrt{(b' - \omega^2)^2 + \omega^2 a'^2}} \quad \text{for } v_d \leq 0 \quad (12)$$

where C and D are arbitrary constants

$$\alpha_1', \alpha_2' = \frac{1}{2} (-a' \pm \sqrt{a'^2 - 4b'}) \quad (13)$$

$$a' = \frac{R_t}{L_t} + \frac{1}{R_d C_d} \quad (14)$$

$$b' = \frac{R_d R_t}{L_t C_d R_d} \quad (15)$$

and

$$\beta' = \frac{a' \omega}{b' - \omega^2} \quad (16)$$

Similarly

$$i^- = \left( C_d \alpha_1' + \frac{1}{R_d} \right) C \cdot e^{\alpha_1' t} + \left( C_d \alpha_2' + \frac{1}{R_d} \right) D e^{\alpha_2' t} + \frac{V_i}{L_t C_d} \frac{1}{\sqrt{(b' - \omega^2)^2 + \omega^2 a'^2}} \left\{ \omega \cos(\omega t - \beta') + \frac{1}{R_d} \sin(\omega t - \beta') \right\} \quad (17)$$

and

$$v_{out}^- = i^- \times R_L \quad (18)$$

In order to make this circuit stable, it is necessary that the real part of roots  $\alpha$ ,  $\alpha$ ,  $\alpha'$ , and  $\alpha$  should be negative and in that case, the transient terms in Eq.'s 10 and 17 will vanish. Suppose A, B, C and D are not infinite. Then,

$$v_{out}^+ \approx \frac{\frac{V_i R_L}{L_t C_d} \sqrt{\omega^2 + (1/R_d^2)}}{\sqrt{(b - \omega^2)^2 + \omega^2 a^2}} \sin(\omega t - \beta - \gamma) \quad \text{as } t \rightarrow \infty \quad (19)$$

where

$$\gamma = \tan^{-1} \omega \cdot R_d \sqrt{\omega^2 + \frac{1}{R_d^2}} \quad (20)$$

Also

$$v_{\text{out}}^- \approx \frac{(V_i R_L / L_t C_d) \sqrt{\omega^2 + \frac{1}{R_d^2}} \sin(\omega t - \beta' + \gamma)}{\sqrt{(b' - \omega^2)^2 + \omega^2 a'^2}} \quad \text{as } t \rightarrow \infty \quad (21)$$

From the necessary condition that the circuit should be stable,

$$a = (R_t / L_t) - (1 / R_d C_d) = \epsilon > 0 \quad (22)$$

and from the condition that the circuit should be monostable,

$$b = (R_d - R_t / L_t C_d R_d) \geq 0 \quad (23)$$

For the circuit constants which satisfy Eq.'s 22 and 23,  $a'$  and  $b'$  become automatically

$$a' = (R_t / L_t) + (1 / R_d C_d) > 0 \quad (24)$$

$$b' = (R_d + R_t / L_t C_d R_d) > 0 \quad (25)$$

Then, the real part of  $\alpha_1'$  and  $\alpha_2'$  can be negative and, therefore, this circuit is stable in the positive resistance region.



By expressing  $a'$  and  $b'$  by the terms of  $a$  and  $b$ , the equations

$$a' = a + \frac{2}{R_d C_d} \quad (26)$$

$$b' = -b + \frac{2}{L_t C_d} \quad (27)$$

can be obtained. By letter  $b = \omega^2$ , Eq.'s 7, 15, and 27 give

$$\therefore \frac{1}{L_t C_o} = \left( \omega^2 + \frac{a}{R_d C_o} + \frac{1}{R_d^2 C_o^2} \right) \approx \left( \omega^2 + \frac{1}{R_d^2 C_o^2} \right) \quad (28)$$

for small  $a$ . Hence,

$$\left| \frac{v_{out}^-}{v_{out}^+} \right| \approx \frac{\omega \cdot \alpha R_d C_d}{2 \sqrt{\left( \omega^2 + \frac{1}{R_d^2 C_d^2} \right)}} \quad (29)$$

From Eq.'s 14 and 15 the following equations can be obtained

$$\omega^2 = (1/L_t C_d) - (a/C_d R_d) - (1/C_d R_d)^2 \quad (30)$$

Since

$$(a/C_d R_d) \ll 1/(C_d R_d)^2 \quad (31)$$

$$\omega^2 \approx (1/L_t C_d) - (1/(C_d R_d)^2) \quad (32)$$

Substitution of this equation into Eq. 29 gives

$$\left| \frac{v_{out}^-}{v_{out}^+} \right| \cong (aC_o/2) \sqrt{R_d^2 - (L_t/C_d)} \quad (33)$$

It is clear from Eq.'s 23 and 32 that as

$$R_d^2 - L_t/C_d \geq 0 \quad (34)$$

Eq. 33 shows that the ratio will decrease by decreasing  $C_d$ ,  $R_d$  or by increasing  $L_t$  with fixed  $a$ , which means that better output waveform of a tunnel diode detector can be obtained either by decreasing  $C_o$  or  $R_d$  or by increasing  $L_t$  under the fixed  $a$ . However, it is known that a maximum  $L_t$  by the stability point of view. When the tunnel diode is shorted, the steady state term of the current can be expressed as

$$i_s = \frac{\ddot{V}_1 \sin(\omega t - \beta\delta)}{\sqrt{R_t^2 + \omega^2 L_c^2}} \quad (35)$$

where

$$\delta = \tan^{-1} \frac{\omega L_c}{R_t} \quad (36)$$

Hence

$$v_{s \text{ out}} = \frac{V_1 R_L \sin(\omega t - \delta)}{\sqrt{R_t^2 + \omega^2 L_c^2}} \quad (37)$$

Therefore, the ratio of  $|v_{out}^+|$  and  $|v_{s \text{ out}}|$  which is an insertion voltage gain in detection circuit becomes

$$\left| \frac{v_{\text{out}}^+}{v_{\text{s out}}} \right| = \frac{\sqrt{R_t^2 + \omega^2 L_c^2} \cdot \sqrt{\omega^2 + (1/R_d)^2}}{L_t C_d \sqrt{(b-\omega^2)^2 + \omega^2 a^2}} \quad (38)$$

usually  $L_c = 0$ , hence

$$\left| \frac{v_{\text{out}}^+}{v_{\text{s out}}} \right| \approx \frac{R_t \cdot \sqrt{\omega^2 + \frac{1}{R_d^2}}}{L_t C_d \sqrt{(b-\omega^2)^2 + \omega^2 a^2}} \quad (39)$$

If  $R_L = R'_g$

$$\left| \frac{v_{\text{out}}^+}{v_{\text{s out}}} \right| \approx 2 \cdot \left| \frac{v_{\text{out}}^+}{v_{\text{in}}} \right| \quad (40)$$



## 4. NUMERICAL ANALYSIS OF TUNNEL DIODE DETECTORS

The previous section shows the conditions by which the maximum detection ratio can be obtained from tunnel diode detectors by assuming the characteristic of the diode as an inverted capital letter  $v$ . Since a tunnel diode does not have such a characteristic the conditions given in previous section may not be proper conditions in order to design special purpose tunnel diode detectors which will be seen later. In order to design high gain tunnel diode detectors, it is necessary to investigate the relationship between the parameters in the circuit and output signals under the actual tunnel diode characteristics which can be done either by experimental set up or by numerical analysis.

In the experiment, it is difficult to adjust the values of circuit constants like  $C_d$ ,  $L_t$ ,  $R_t$  and so on as they are desired. However, when the circuits are analyzed by the computer, it is possible to control them. Furthermore, it is also possible to analyze the circuit with the ideal elements.

From the circuit in Fig. 3, the following equations can be obtained:

$$L_t \frac{di}{dt} + R_t i + v_d = v_i \sin \omega t + E \quad (41)$$

$$C_d \frac{dv_d}{dt} = i - i_d(v_d) \quad (42)$$

Combining above equations, Eq. 43 can be obtained.

$$\begin{aligned} \frac{d^2 v_d}{dt^2} + \frac{R_t}{L_t} \frac{dv_d}{dt} + \frac{v_d}{L_t C_d} + \frac{1}{C_d} \frac{di_d(v_d)}{dt} \\ + \frac{R_t}{L_t C_d} i_d(v_d) - v_i \sin \omega t + E \end{aligned} \quad (43)$$

By the use of a digital computer, Eq. 43 can be solved with parameters  $R_t$ ,  $L_t$ ,  $C_d$ ,  $V_i$ ,  $\omega$  and  $E$  and characteristic  $i_d(v_d)$  of a tunnel diode as shown in Fig. 6.

$-R_d$  at P is approximately equal to  $100 \Omega$ . The junction capacitor  $C_o$  of such a tunnel diode is approximately  $C_d = 6 \times 10^{-12}$  f (which is estimated by considering present A<sub>s</sub> tunnel diode). The frequency  $\omega$  of the input signal  $V_i \sin \omega t + E$  is set to  $2\pi \times 10^8$  cps for convenience. From Eq.'s 7 and 8, the following equation can be obtained:

$$R_t = \frac{(R_d C_d a + 1) R_d}{\omega^2 (C_d R_d)^2 + a R_d C_d + 1} \quad (42)$$

and

$$L_t = \frac{R_t R_d C_d}{(a R_d C_d + 1)} \quad (43)$$

For convenience,  $R_{t_o}$  is defined as

$$R_{t_o} = \frac{R_d}{\omega^2 C_d^2 R_d^2 + 1} \quad (44)$$

Also  $L_{t_o}$  is defined as

$$L_{t_o} = R_t R_d C_d \quad (45)$$

Notice that  $R_t \approx R_{t_o}$  and  $L_t = L_{t_o}$  if  $a \ll \frac{1}{R_d C_d}$ . Also  $R_t > R_{t_o}$  and

$L_t > L_{t_o}$  for  $a > 0$  with fixed  $\omega$ . Then the values  $R_d = 100 \Omega$ ,  $C_d = 6 \times 10^{-12}$  f,

and  $\omega = 2 \times 10^8$  cps,  $R_{t_0}$  and  $L_{t_0}$  are approximately equal to 87.6 ohms and  $5.25 \times 10^{-8}$  h respectively.

The following are the results obtained by the numerical calculation using a digital computer:

(i)  $R_t$  v.s.  $v_{out}^+$

By varying  $R_t$ , the curve in Fig. 7 is obtained.  $L_t$  is also varied in order to make an invariant.

Notice that  $R_t$  gives maximum  $v_{out}^+$  by Eq.'s 19 and 20. From the curve in Fig. 7 it is clear that  $R_t > R_{t_0}$  gives larger  $v_{out}^+$  which means that the conditions obtained by linealized characteristics of tunnel diodes are not applicable for the designing of a tunnel diode and detector.

The reasons are as follows: (1) because the negative resistance region is finite, the magnitude of the output signal under the situation that the operating point would be inside of the negative region is bounded which makes finite magnitude at output signal; (2) since the magnitude of the negative resistance near the peak current is larger than that near the middle point of the negative resistance region, with the small signal input and with the biasing point at the peak current region, larger  $R_t$  should give the larger output voltage than  $R_{t_0}$  which is almost equal to  $R_d$ .

(ii) Biasing Point vs.  $V_{out}$ .

When the biasing points of the tunnel diode in Fig. 2 are varied from C to A to B as shown in Fig. 8, the wave form of output voltage  $V_{out}$  will be changed as in Fig. 9.

(1) When the biasing point is at B ( $V_B = .06V$  in Fig. 8) the positive side of the output voltage  $v_{out}^+$  is a little larger than that when  $V_B = .05V$ . However, the negative side of output voltage  $v_{out}^-$  is also larger than that when  $V_B = .05 V$ .

Hence, as far as detection is concerned biasing at the peak current region would be better than biasing at the point which is away from the peak current region such as C or B in Fig. 8.

(iii) Input voltage vs. Output voltage

The curves in Fig. 10 show that  $v_{in}$  vs.  $v_{out}$  with different values of parameters. Every curve in the figure shows the sudden increase of gain which



can be seen in the experimental results shown in the next section. After sudden increase the operating point will reach to the valley region of the characteristic curve in Fig. 8 which makes the detector circuit into saturating states (the output voltage will not increase by increasing the input voltage). Notice that even though  $a < 0$ , the circuit will not be unstable with small input signal which contradicts the results in the previous section.

(iv)  $L_t$  vs.  $V_{out}$

With  $R_t = 114 \Omega$  and  $C_o = 6$  PF, the curve in Fig. 11 is obtained. Even though  $R_t > R_d$ , the circuit will act as a detector with small input signals. It is important to notice that the increase of  $L_t$  will increase the output voltage  $v_{out}$ . However with  $V_i = 10$  mv, the end of the right hand side of the curve is approaching the saturation point of the circuit. Also,  $L_t$  must be small enough to avoid self-oscillation.

## 5. EXPERIMENTAL RESULTS

The calculation of Eq. 4 by a digital computer for a large number of combinations of different values of parameters in the tunnel diode detector needs a large amount of time. Also, in order to verify that the results from the calculation by a computer are applicable to design a tunnel diode detector, the following experiment is carried out. The circuit scheme is shown in Figure 12.

(i) Miniature variable resistors were used as series resistors and by adjusting the value of these resistors, we investigated the relations between  $v_{out}$  and the values of the resistors  $R'_g$  and  $R_L$ .

In order to avoid any influence on the results by stray capacitances and conductances in the circuit to be measured, the carrier frequency of the amplitude modulated input signal is limited to 3 mc. The tunnel diode which is used in the circuit is GE 1N2939 which has about 1 ma peak current. The V-I characteristic of the diode which is traced by the curve traces shown in Fig. 13 is shown in Fig. 14.

According to the manual of the tunnel diode, the value of  $C_d$  and  $L_d$  of 1N2939 are 5.0 to 1.5  $\mu$ f and 6 nh respectively

The D. C. bias voltage of the tunnel diode is measured by the oscilloscope with differential amplifier.

(ii) For the bias voltage the output voltage of the power transistor which gives the smooth varying D.C. voltage by controlling the base current is used.

In order to obtain the output wave form as a function of the magnitude of the input signal, it is convenient to make  $V_i$  of input signal  $V_i \sin 2\pi ft$  not as a constant but as a linearly increasing time varying signal. Hence the sawtooth wave with period of about 1 m sec is used which is easily obtained from the sawtooth wave output of the oscilloscope.

Table I shows the output waveform  $v_{out}$  corresponding to amplitude modulated wave with modulating signal being a sawtooth input waveform and carrier frequency of .5 mc for different values of  $R_t$  and biasing voltage  $V_B$ .

The symbol  $v_{s out}$  is the output voltage when the tunnel diode is shorted.  $v_{s out}$  of the circuit under experiment is as shown in Fig. 15.

If there is only one waveform in a picture in Table I, then the waveform indicates  $v_{out}$  and  $v_{s out}$  are equal to that in Fig. 15. If there are two waveforms in a picture, one whose slope is an isosceles triangle is  $v_{s out}$  and the other is  $v_{out}$ .

The unit of the vertical axis in every picture in the table and Fig. 15 is 5 mc except two pictures corresponding to  $R_t = 160$  ohms and the biasing voltages of 80 mv and 165 mv (each of which has unit of the vertical axis 10 mv). The maximum peak to peak voltage of  $v_{s out}$  in Fig. 15 in most of the pictures in Table I is, therefore, 15 mv except in the three pictures corresponding to  $R_t = 136$  ohms and the biasing voltages of 100 mv, 120 mv and 140 mv (each of which has peak to peak  $v_{s out}$  voltage of 5 mv).

From V-I characteristic of a tunnel diode in Fig. 14, the voltage  $v_d$  at the peak amount is about 70 mv. The biasing point at about 80 mv are where the point on the curve is a little away from the peak current toward to the valley current. The biasing voltages from 120 mv to 140 mv are where the point on the curve in Fig. 14 is about the middle on the negative resistance region at which the magnitude of the negative resistance is the minimum. The biasing voltage of 165 mv is the point at which the magnitude of the negative resistance begins to increase rapidly. Hence, at the biasing voltage between 100 mv and 140 mv in Table I, the circuit in Fig. 12 acts as a detector. The minimum negative resistance  $-R_d$  of the tunnel diode in the circuit in Fig. 12 is between -135 ohms and -138 ohms. Hence there is the case which  $R_t = 160 \gg R_d$ . It is impossible to obtain pictures for which the biasing point is between 100 mv and 140 mv. The reason is simply that the circuit is unstable at these biasing voltages. Notice that the output signal increases suddenly in the pictures for the biasing voltages of 70 mv and 80 mv ( $R_t = 160$  ohms).

For convenience, the symbol  $g_{s_i}$  is defined as

$$g_{s_i} = \frac{P_i}{q} \quad (i = 1, 2) \quad (46)$$

where  $P_1$  is the slope of the upper envelope and  $P_2$  is the slope of the lower envelope at time  $t$  of the output waveform  $v_{out}$  in a picture in Table I in which the horizontal axis indicates time  $t$ ,  $q$  is the slope of one side of envelope of  $v_{s out}$  corresponding to the picture. Then it is clear from Table I that the



$\max(g_{s_1}, g_{s_2})$  is larger if  $R_t$  is larger when the biasing voltage  $V_B$  is one of 70 mv, 80 mv, and 165 mv.

Notice that  $g_{s_i}$  is a function of  $v_{s_{out}}$ .

Suppose  $v_{s_{out}}$  is as shown in Fig. 16a which indicates that the input signal to a tunnel diode is an amplitude modulated signal with very small degree of modulation. Also suppose  $v_{out}$  is as shown in Fig. 16 b. Then  $g_{s_i}$  is almost equal to  $e_o/e_s$  with very small  $e_s$ . Also the following equation will hold;

$$v_{out} = k + g_{s_i} v_{s_{out}} \quad (47)$$

where  $k$  and  $g_s$  are functions at  $v_{s_{out}}$ . Hence  $k$  and  $g_{s_i}$  give the relationship between  $v_{s_{out}}$  and  $v_{out}$  which indicates the quality of detection, that is, for a small input signal  $e_s$ , larger  $g_{s_i}$  gives larger amplification (which means larger  $e_o$ ). Notice that for small  $e_s$ ,  $k$  is not important at all. However, for a larger  $e_s$ , constant  $g_{s_i}$  and  $k$  will be the best for a linear amplification of  $e_s$ .

Fig. 17a shows that  $g_{s_i}$  varies rather rapidly. Also it indicates that for higher frequency, maximum  $g_{s_i}$  does not change much for changing  $v_{s_{out}}$  near maximum  $g_{s_i}$  when the carrier frequency becomes higher.

From Eq. 19 it can be seen that there exists a critical frequency  $\omega_c$  of carrier frequency where  $\omega_c^2 = b$  in Eq. 19, such that with increasing carrier frequency  $\omega$  above  $\omega_c$ , the output  $v_{out}^+$  decreases. But increasing carrier frequency  $\omega$  which is still below  $\omega_c$  increases the output  $V_{out}^+$ . However, the results in Fig. 17a show that the increasing carrier frequency  $\omega$  decreases  $v_{out}^+$  and  $g_{s_i}$ . Even  $\omega$  is much less than  $\omega_c$ .

Fig. 17a also shows the  $g_{s_2}$  of  $v_{out}$  corresponding to the negative half-cycle at the input signal by which the detection by the tested circuit is excellent. Fig. 17b shows the  $k$  of  $v_{out}$  corresponding to the positive and the negative half-cycle of the input signal respectively.

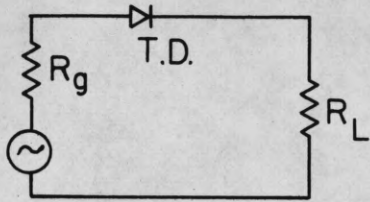
With carrier frequency  $\omega = .5$  mc, total series resistance  $R_t = 136$  ohms, small signal gain  $g_s$  vs. input signal is shown in Fig. 18. From this it is learned that there are two regions on V-I characteristics of a tunnel diode such that by setting biasing point in one of these two regions, the circuit becomes a detection with reasonably high gain obtainable. These two are regions near points B and C in Fig. 2b. Notice that at the region near point C in Fig. 2b, the output waveform  $v_{out}$  has  $g_s > 2$  for half-cycle at which maximum  $(g_{s_1}, g_{s_2})$  for the other half-cycle is obtained. On the other hand, at region near point B in Fig. 2b,  $g_s < 1$  for a half-cycle at which maximum  $(g_{s_1}, g_{s_2})$  for the other half-cycle is obtained.

Fig. 19 shows how changing total series resistance  $R_t$  influences  $g_s$  vs.  $v_{s out}$ . Notice that when  $R_t > r_{d_1}$   $\max(g_{s_1})$  is large which is predictable from the pictures in Table 2 in which there is a sudden increase of  $v_{out}$ . For a comparison, the output waveforms of a carrier frequency by experiment and by the computation (by a digital computer) are shown in Fig. 20 a and b.

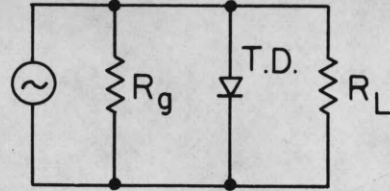
## 6. CONCLUSIONS

It is not difficult to design a tunnel diode detector with small signal gain  $g_s$  at about 15. Hence if the input signal is an amplitude modulated wave with small percentage of modulation, a tunnel diode detector is a very practical circuit which has advantages over an ordinal diode detector. There are many problems related to the detector which are left for future work. These problems are indicated in each section of this paper.



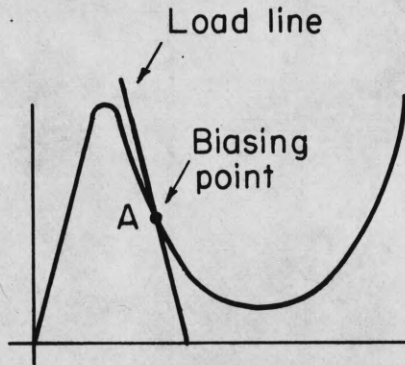


(a) Series amplifier

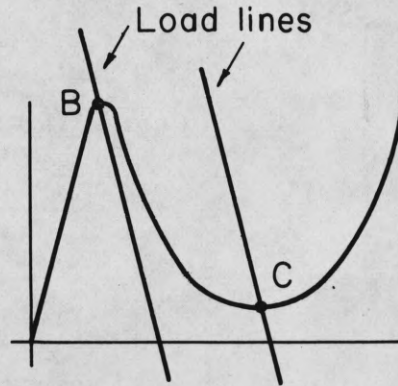


(b) Parallel amplifier

Fig. 1



(a) Biasing point for usual amplifier



(b) The biasing points B or C for detector

Fig. 2

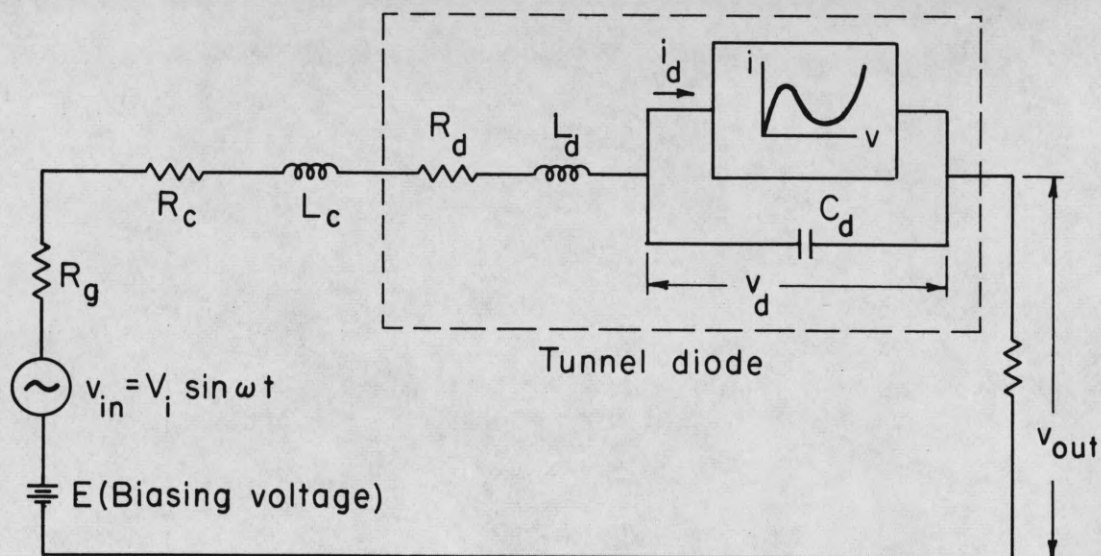


Fig. 3. Equivalent circuit for a tunnel diode detector

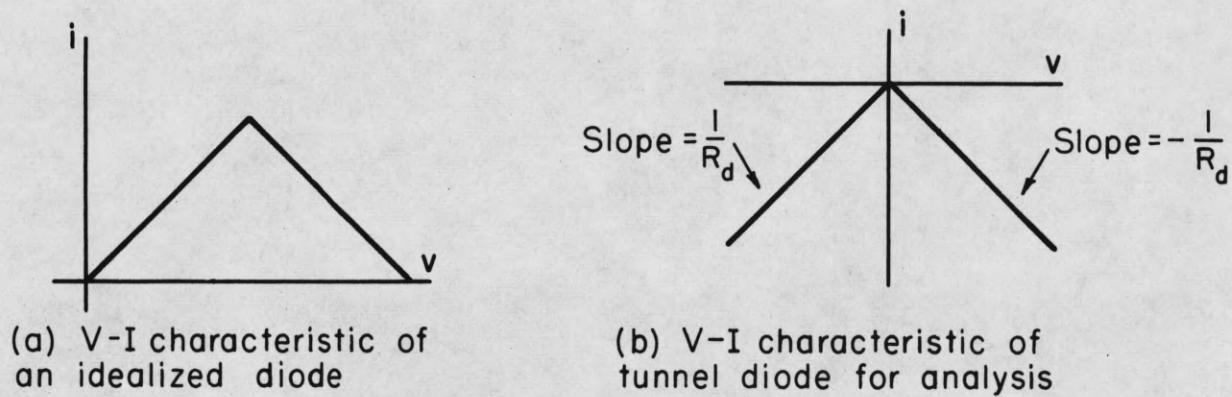


Fig. 4. Simplified V-I characteristic of a tunnel diode

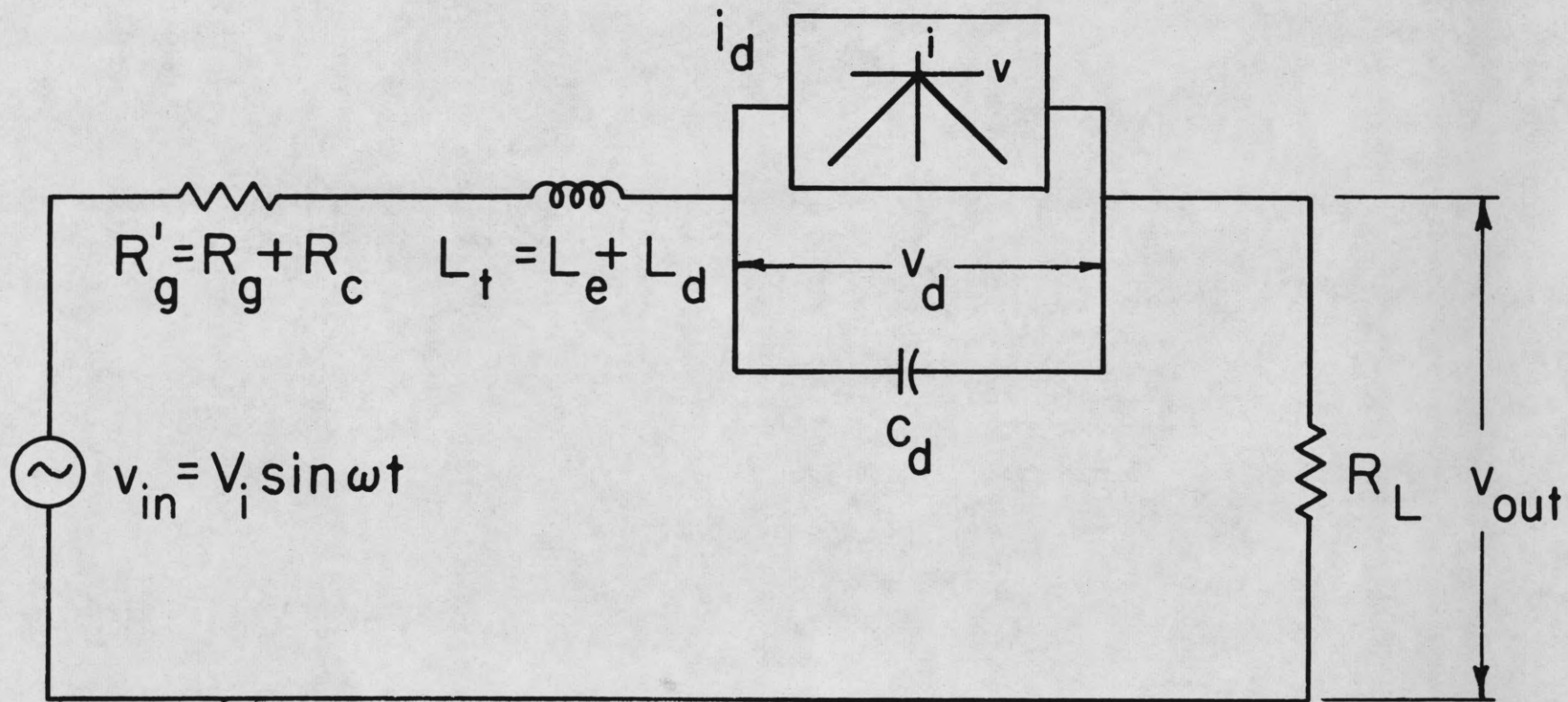


Fig. 5. Equivalent circuit for tunnel diode detector with specified B-I characteristic.



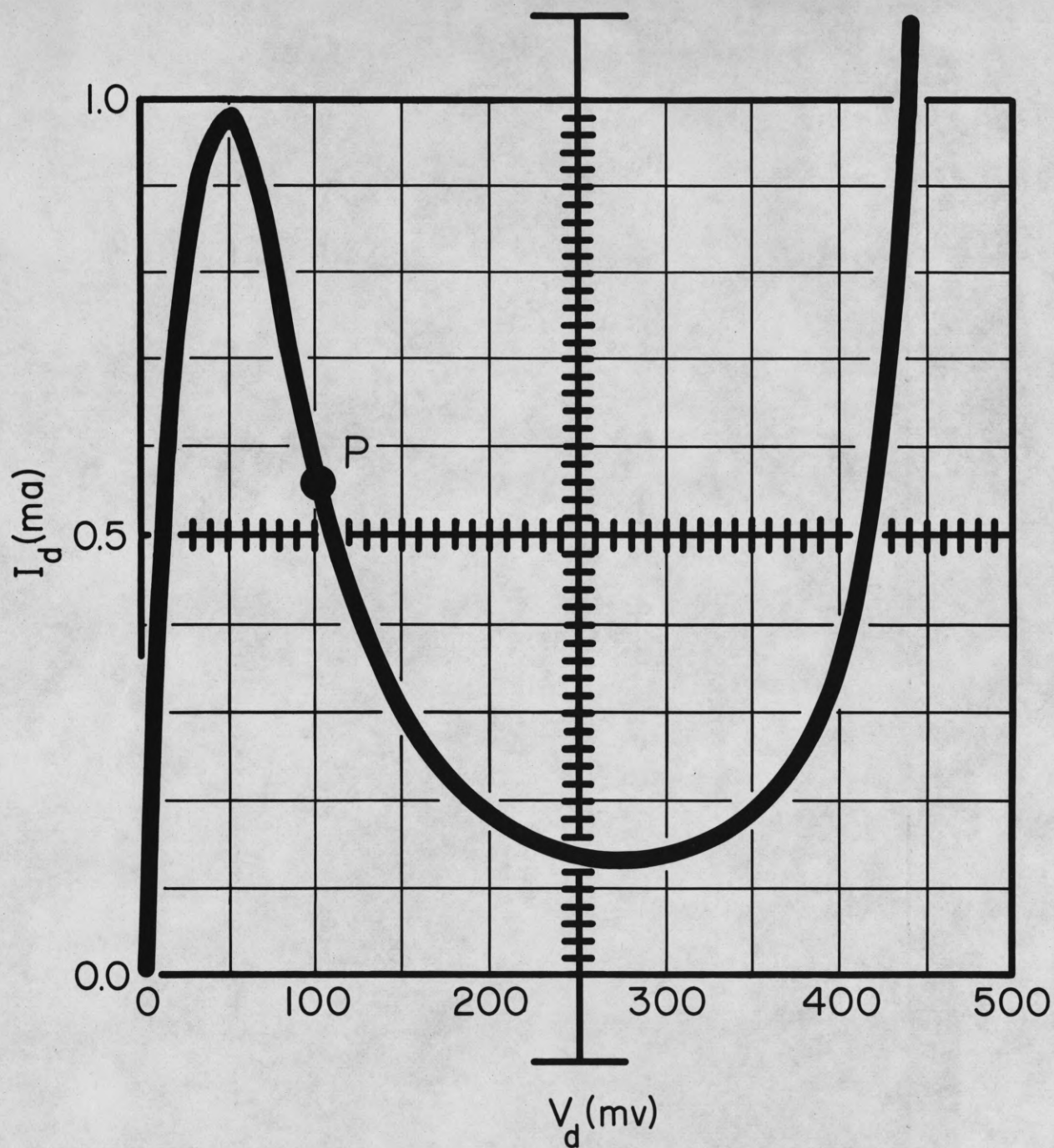


Fig. 6 V-I characteristic of a tunnel diode

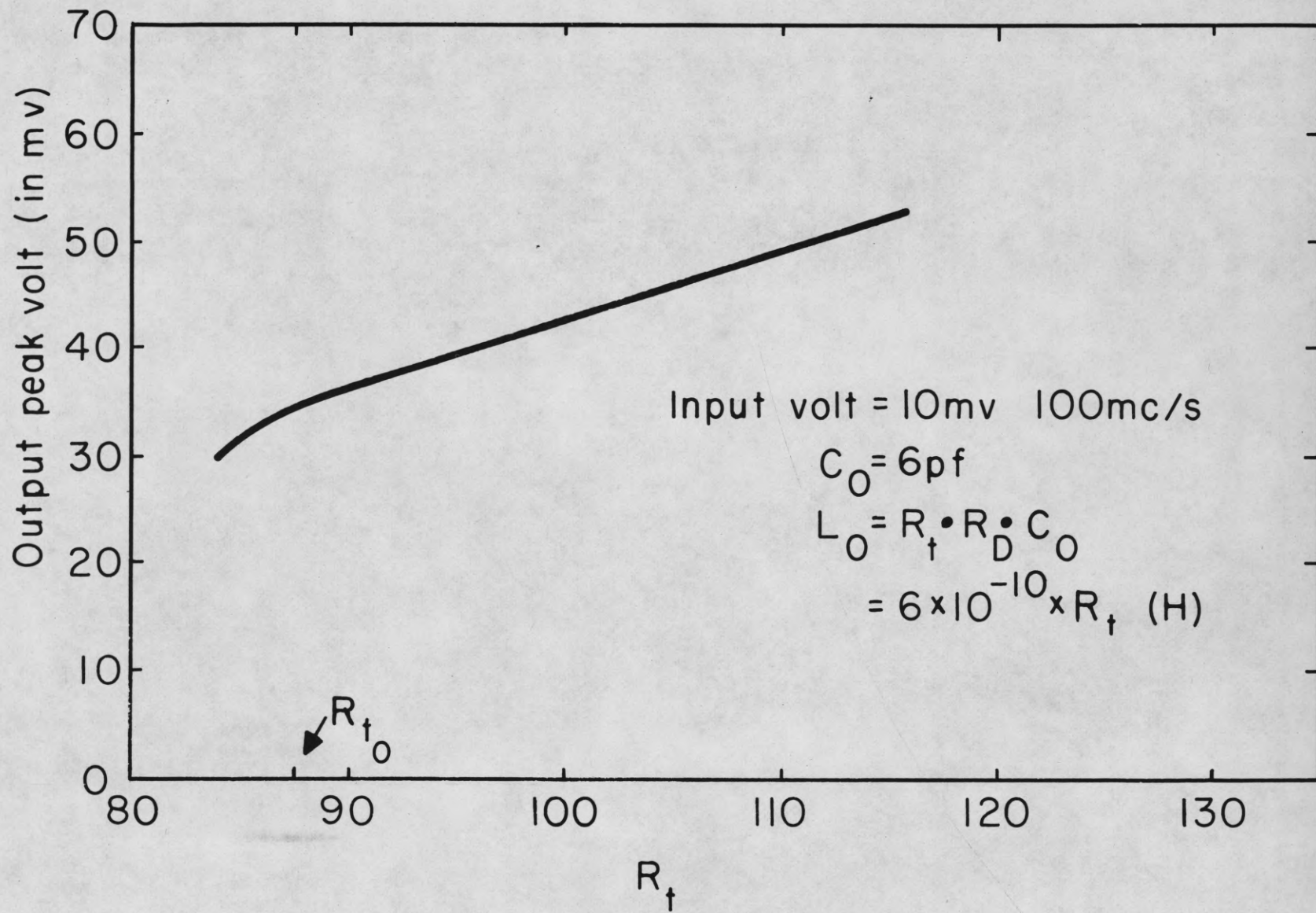


Fig. 7. Peak  $v_{out}$  v.s.  $R_t$

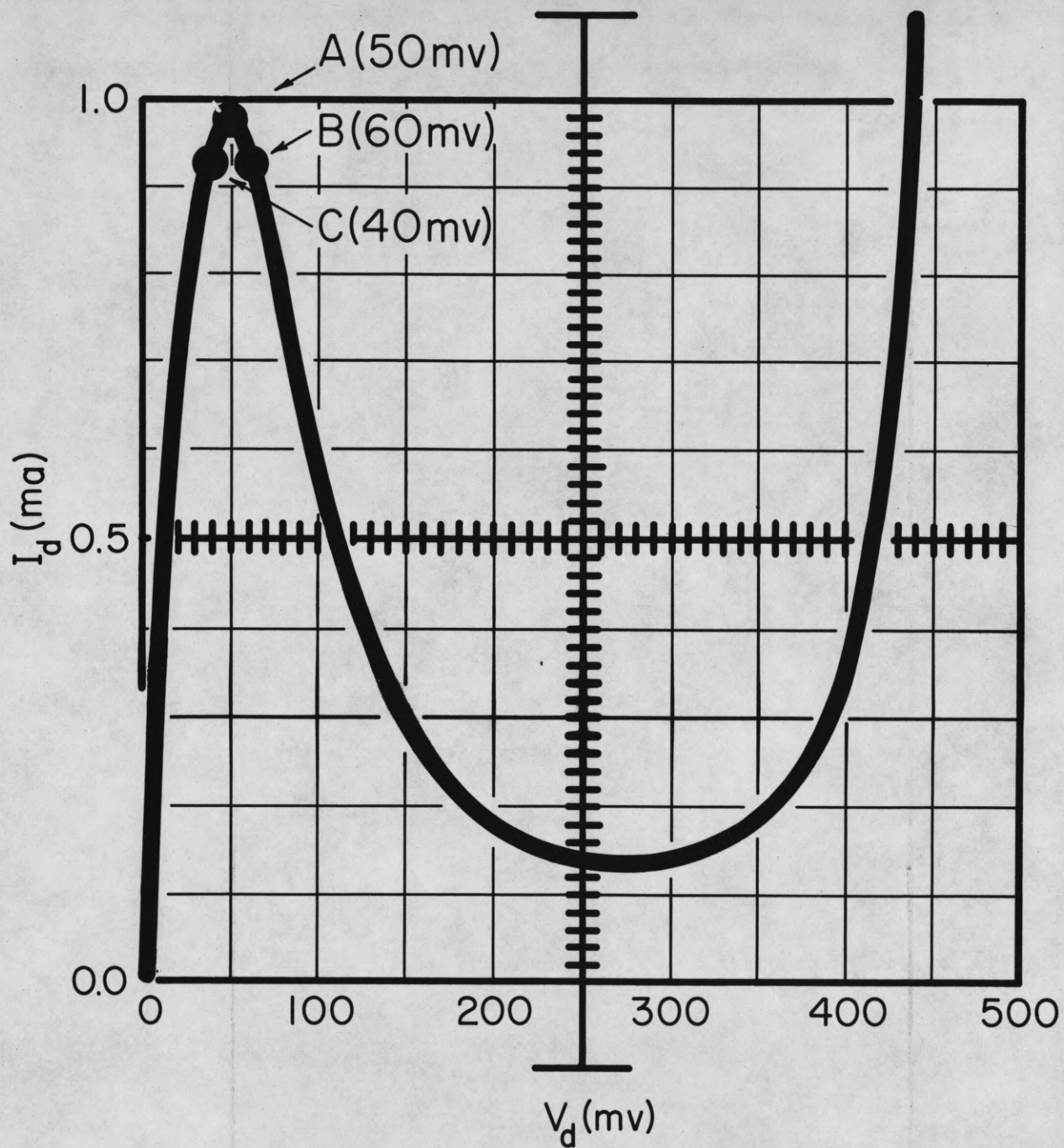


Fig. 8. Biasing Points A, B, and C.



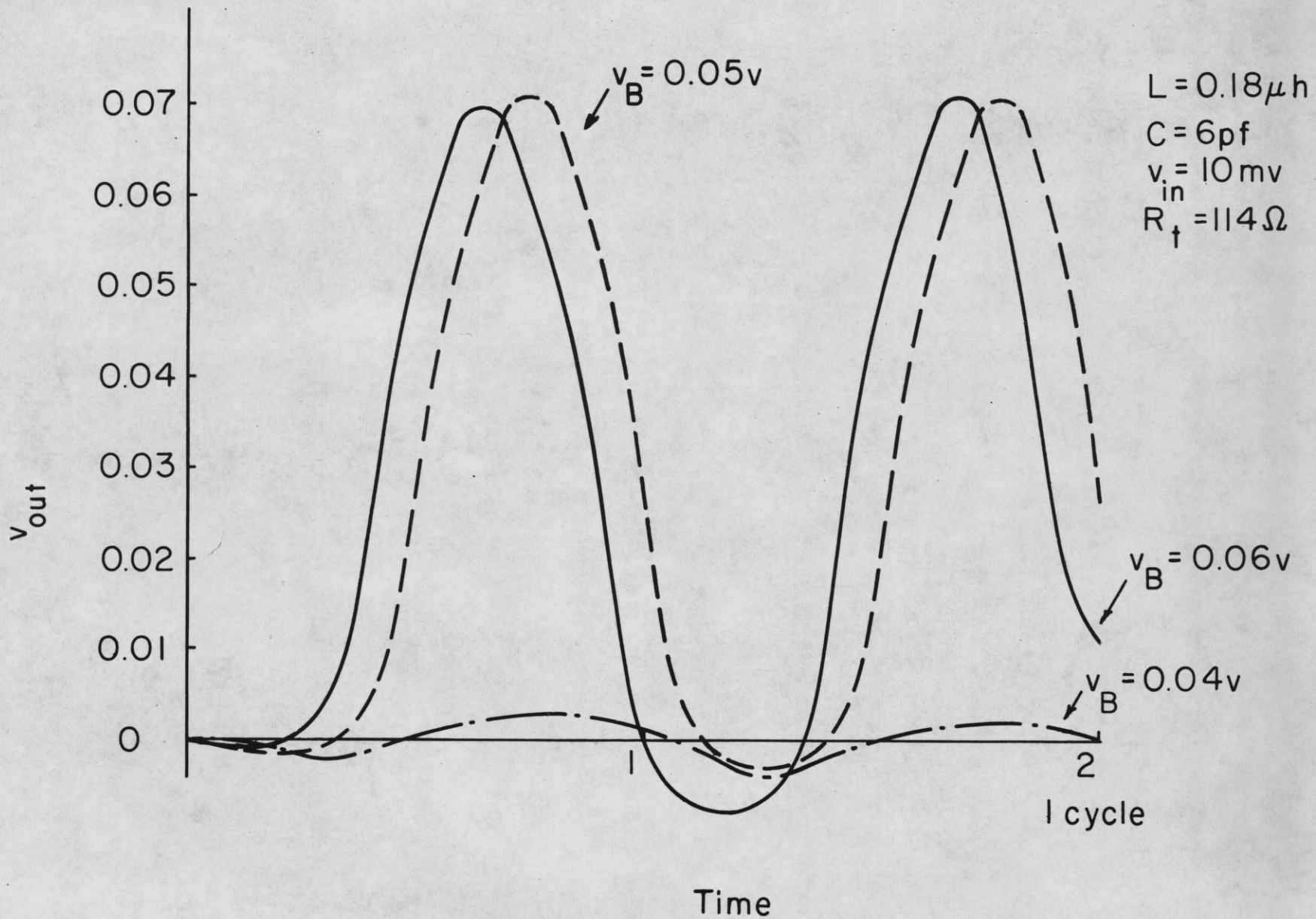


Fig. 9. Variation of  $v_{out}$  by different biasing points

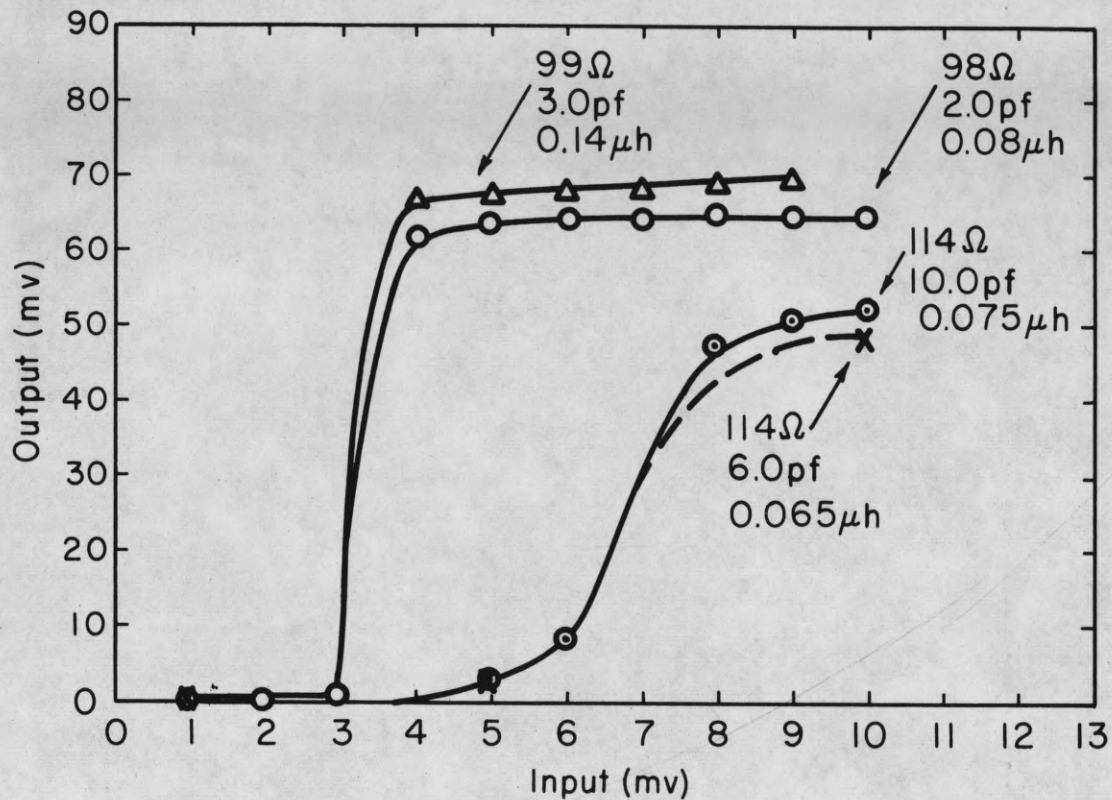


Fig. 10.  $v_{in}$  v.s.  $v_{out}$

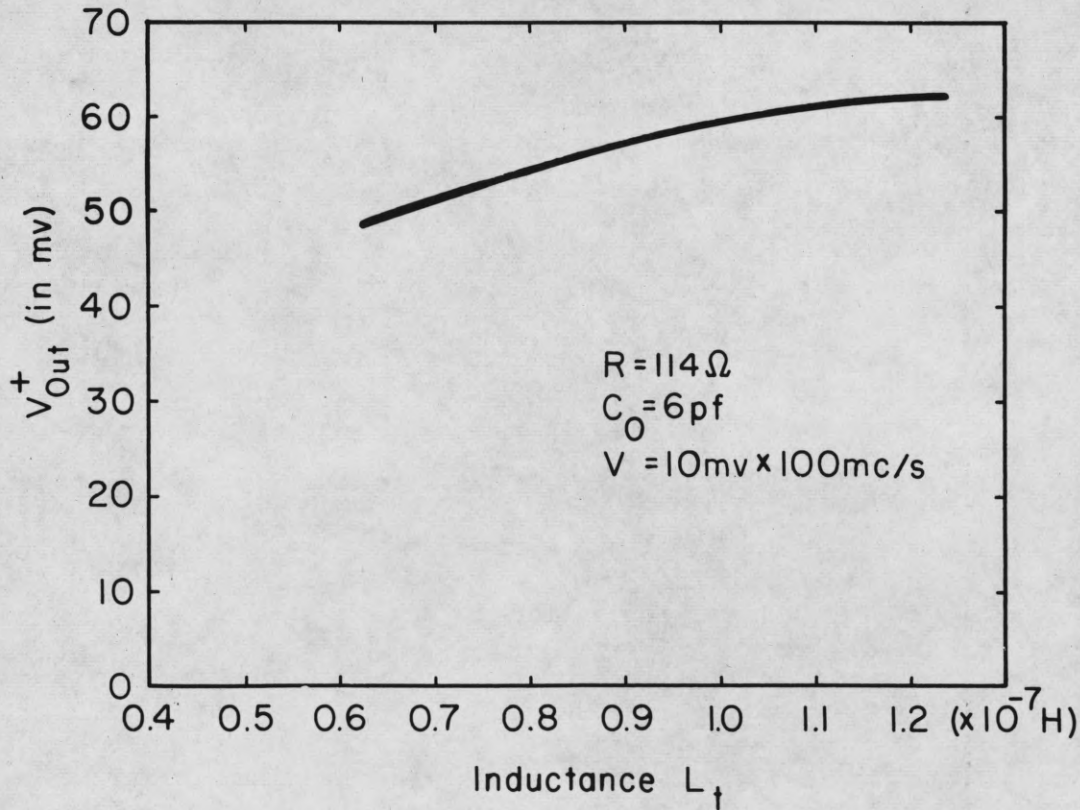


Fig. 11.  $v_{out}$  v.s.  $L_t$

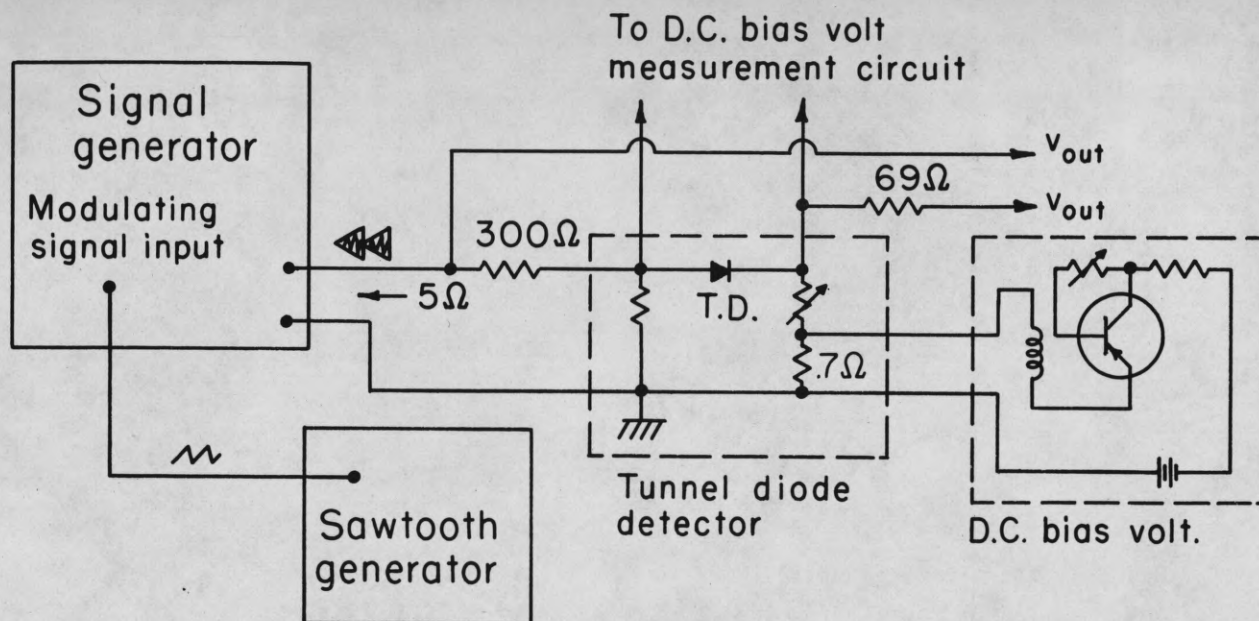


Fig. 12. Circuit scheme for testing a tunnel diode detector.

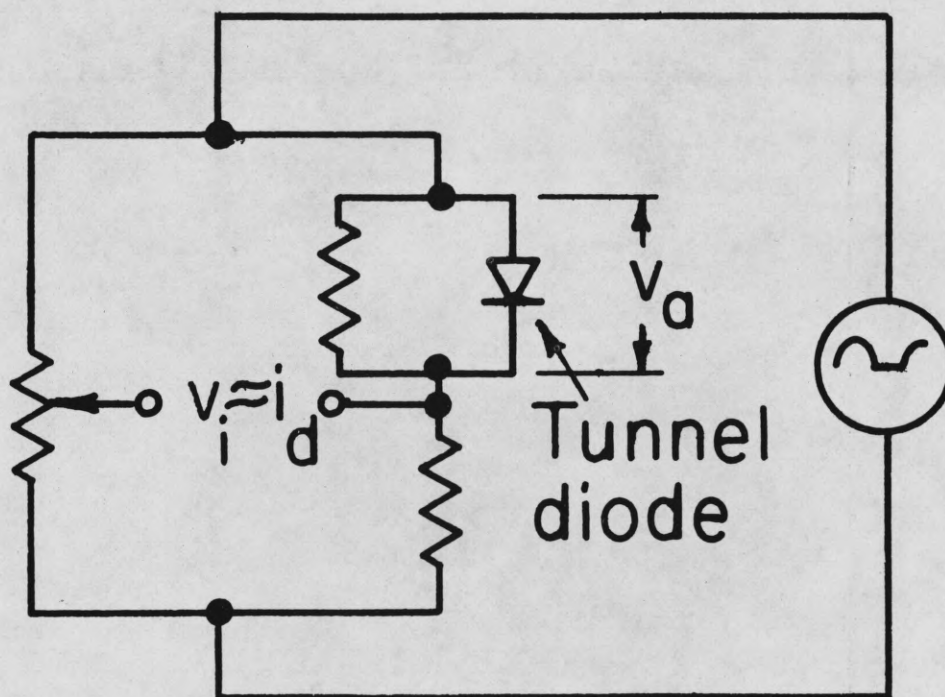


Fig. 13. Tunnel diode curve traces.



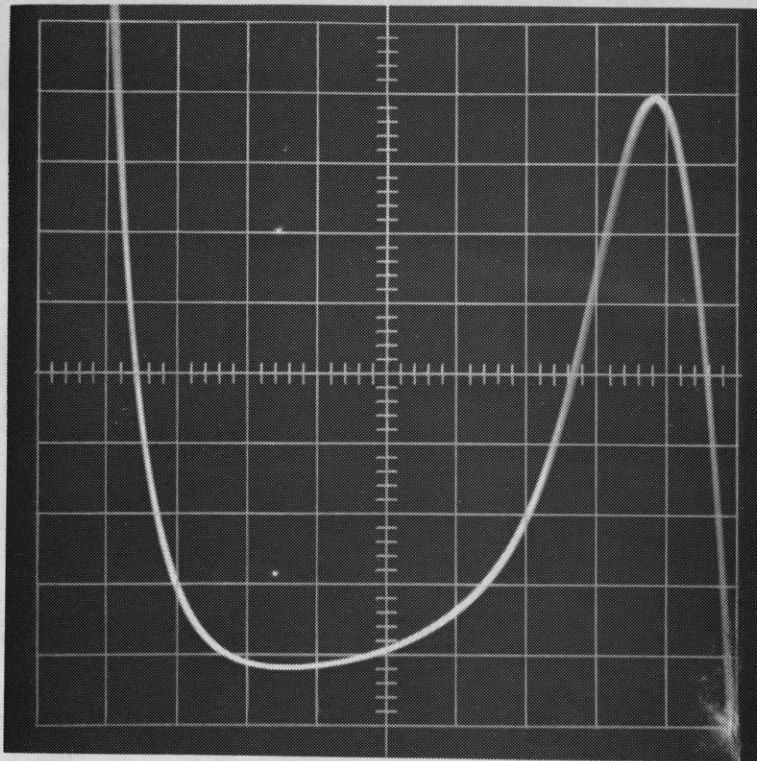


Fig. 14. V-I characteristic of 1N2939

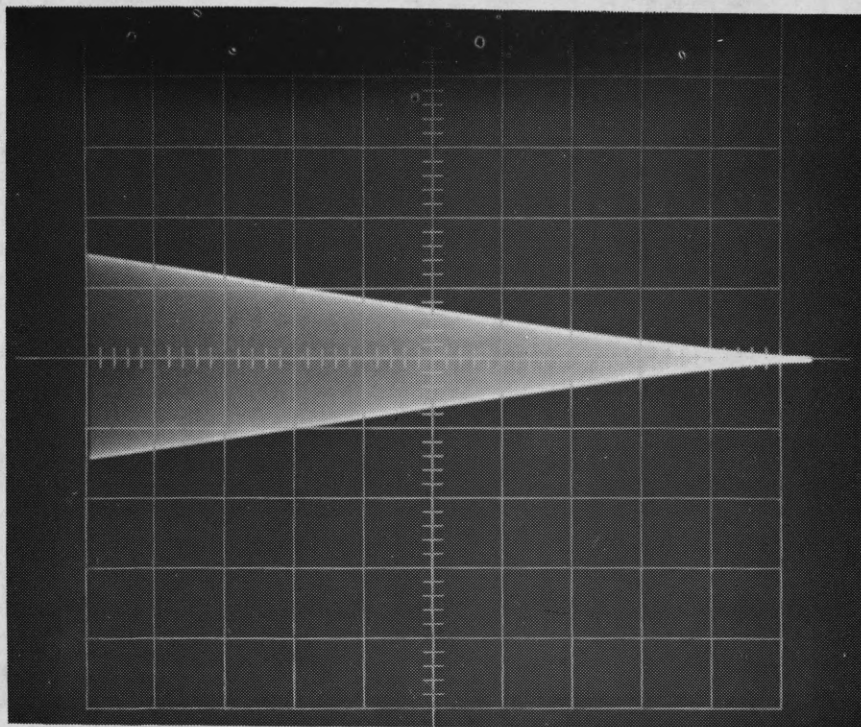
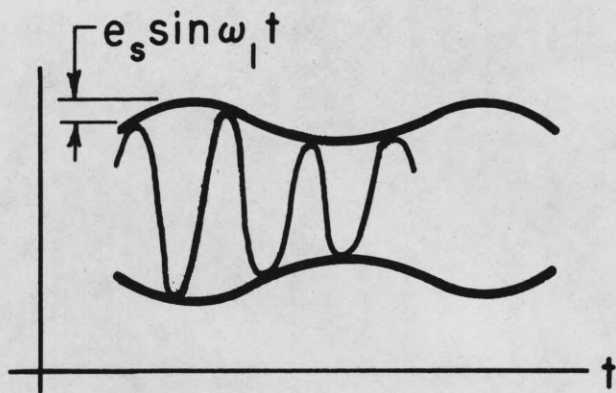


Fig. 15. Waveform of  $v_{in}(t)$

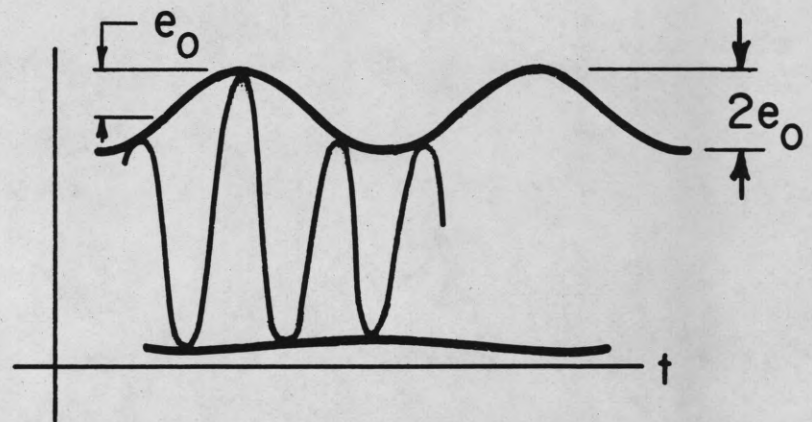
Table 1  
 (  $R_T = R_g' + R_L$  and  $R_g' \cong R_L$ ,  $V_B$ : Bias Voltage )

$R_T$ (ohm) $V_B$ (mv)	160	140	136	130	120
70					
80	 Vert. 10mv/Unit				
100			 Max. $V_{s_{out}}$ : 5mv p-p		
120			 Max. $V_{s_{out}}$ : 5mv p-p		
140			 Max. $V_{s_{out}}$ : 5mv p-p		
165	 Vert. 10 mv/Unit				





(a)  $v_{s_{out}}$



(b)  $v_{out}$

Fig. 16. Definition of  $e_s$  and  $e_o$ .



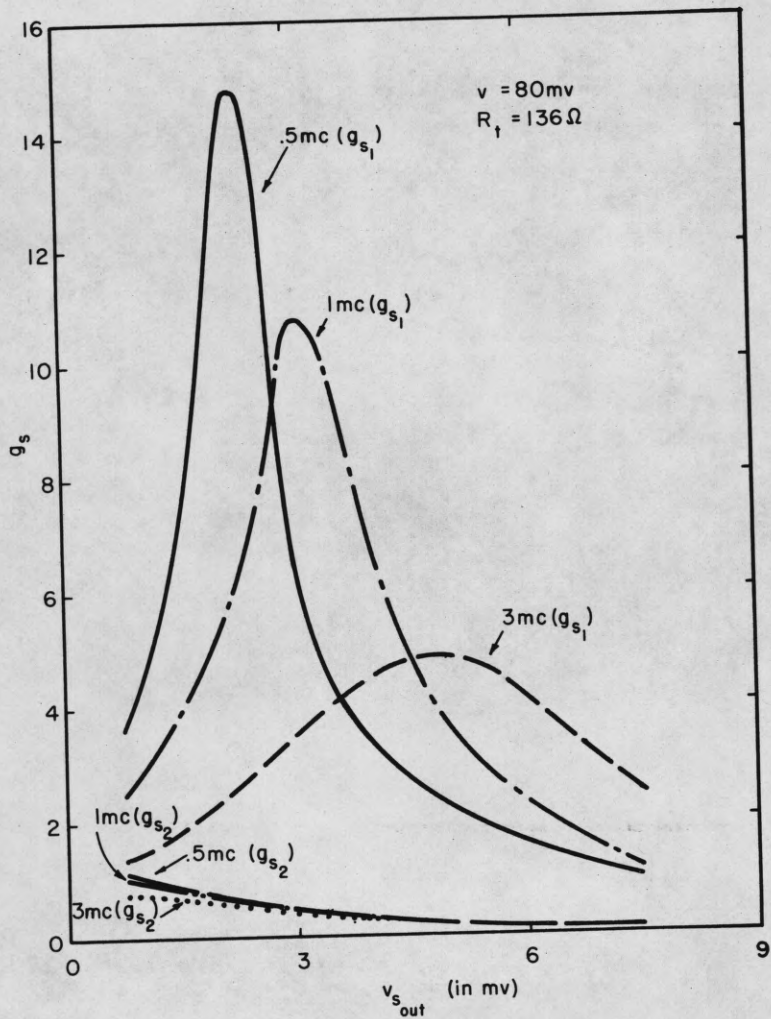


Fig. 17a. Small signal gain.

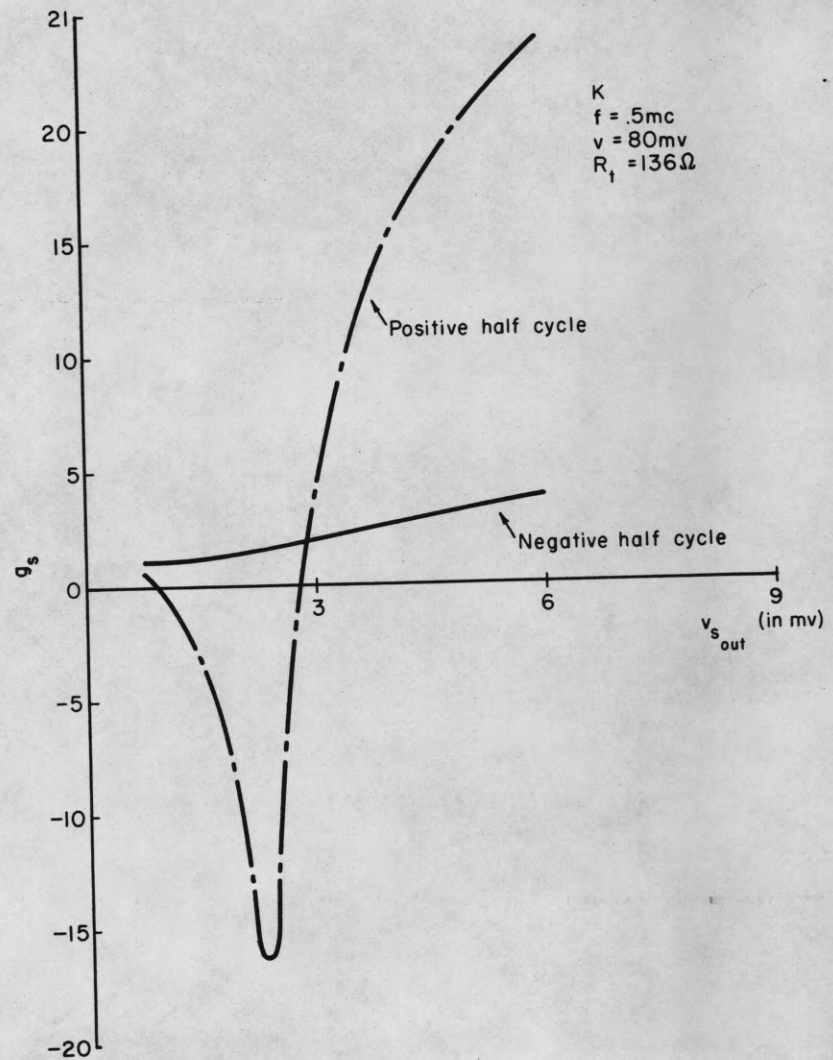


Fig. 17b. Gains of positive and negative half cycles.

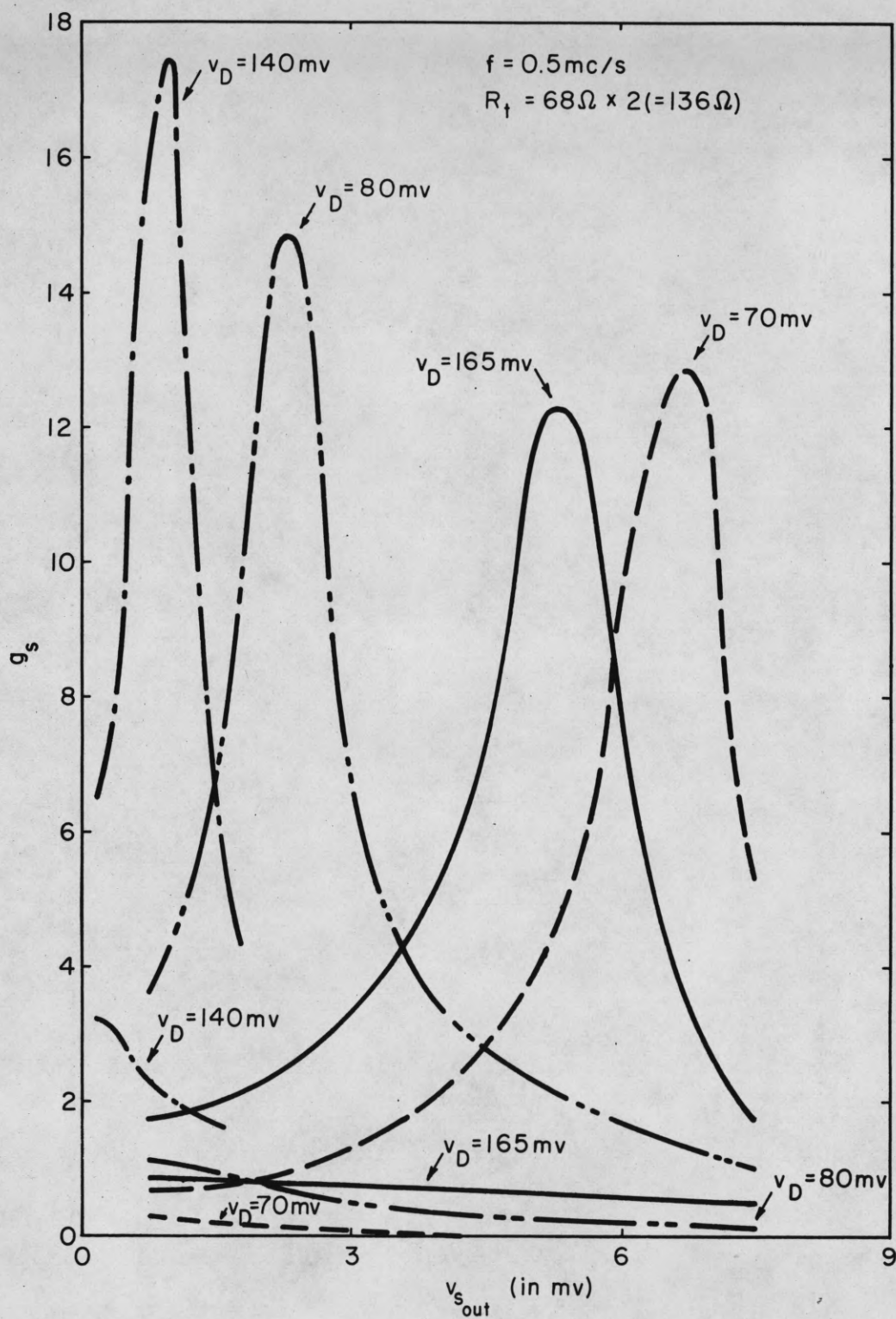


Fig. 18. Small signal gains with different  $V_D$  (voltage across the diode)

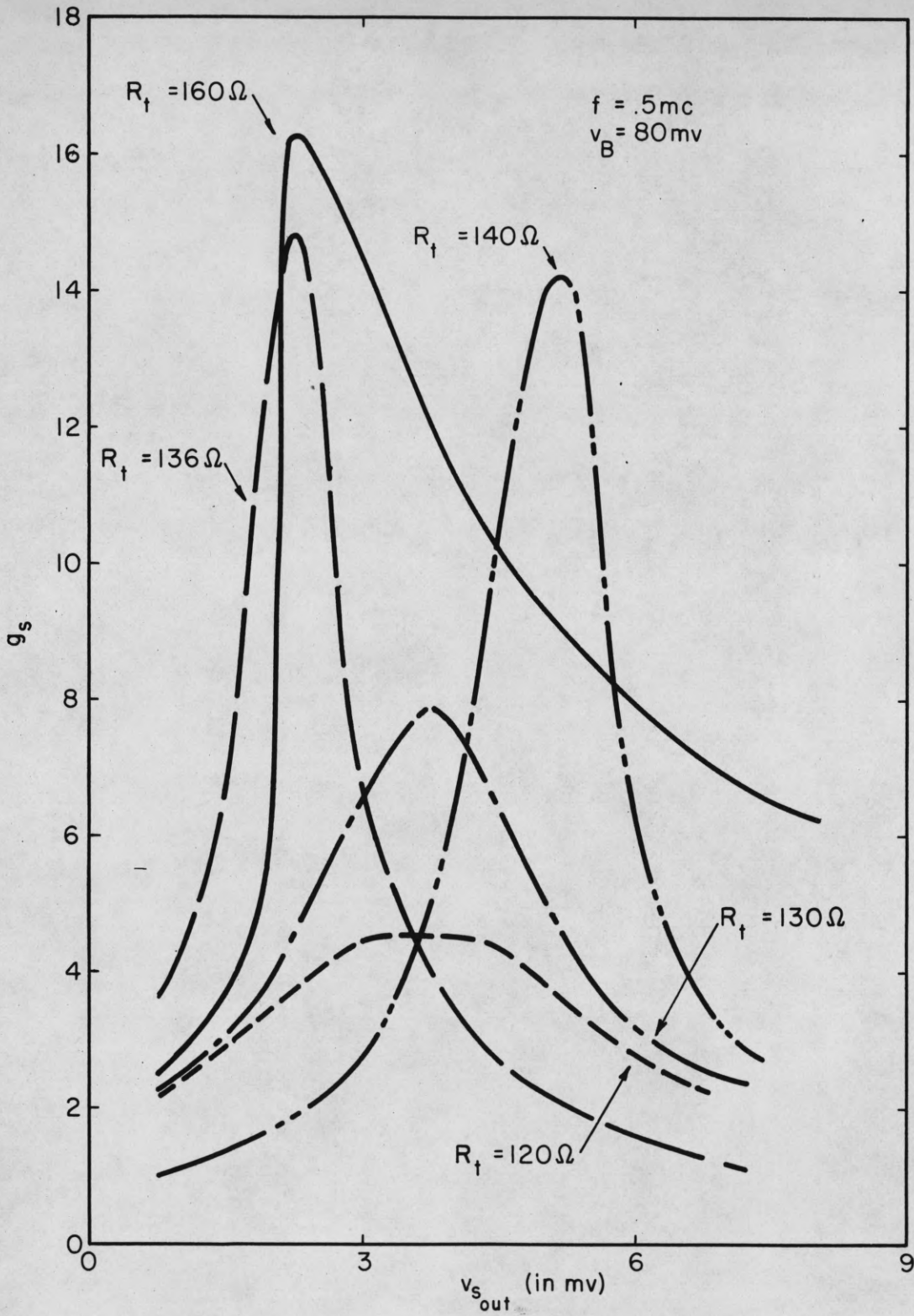


Fig. 19. Small signal gains with different  $R_t$ .



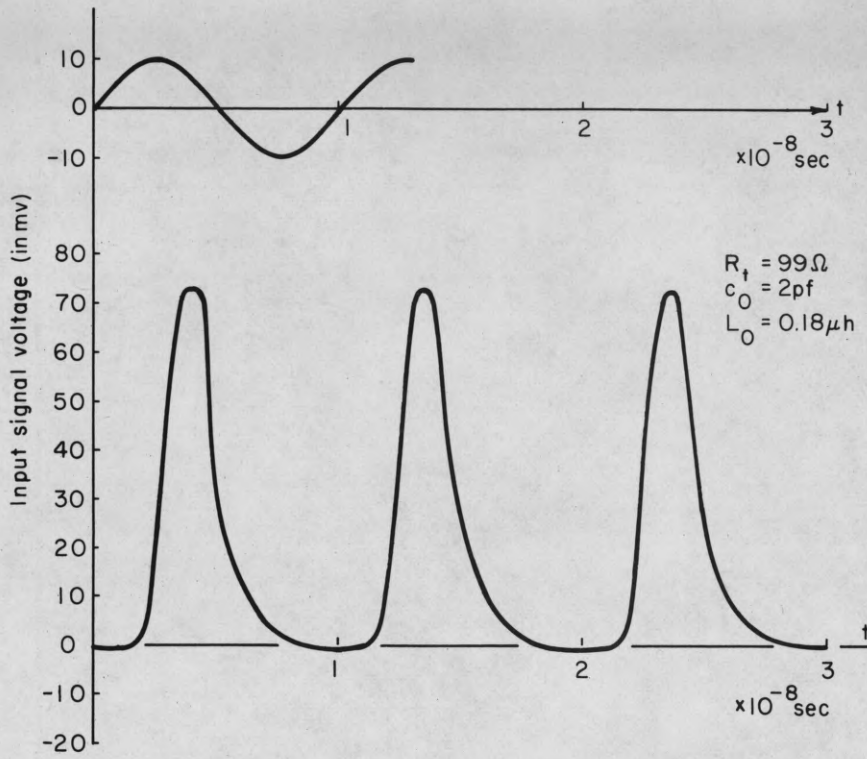


Fig. 21a. Waveform obtained by calculation.

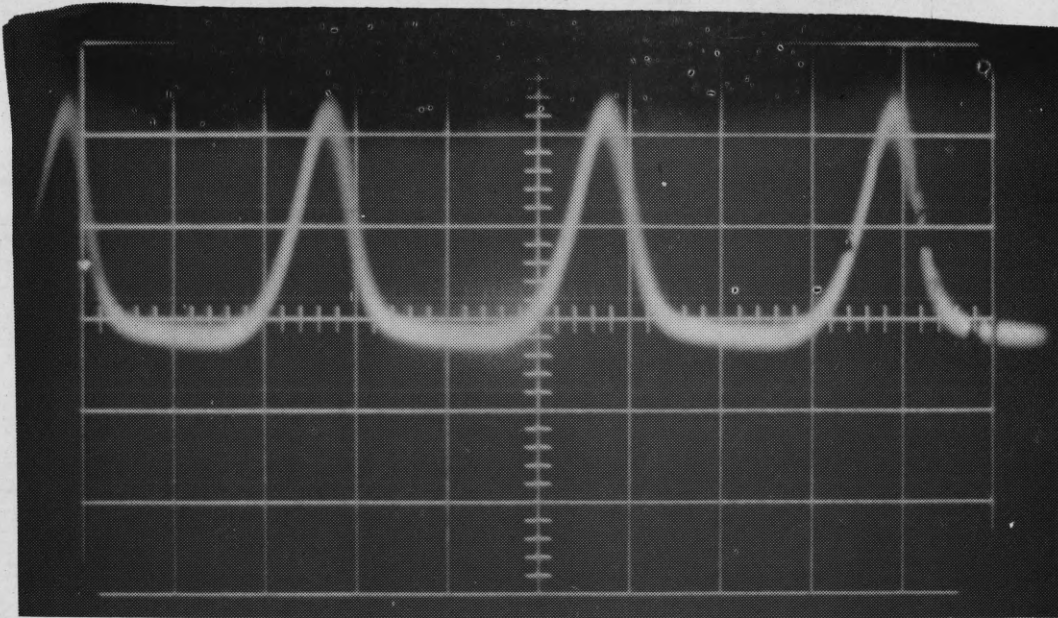


Fig. 21b. Waveform obtained by experiment.

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