Noise-Immune Cavity-Enhanced Optical Heterodyne Molecular Spectrometry: Modeling under Saturated Absorption

Patrick DUPRÉ

Laboratoire de Physico-Chimie de l'Atmosphère, Université du Littoral, Côte d'Opale Dunkerque, France

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## Outline

#### Introduction

- Background
- Absorption in a Resonant Cavity

### Modeling

- By Hands
- Fundamentals
- Equations



#### Simulations

- Para-Acetylene
- Ortho-Acetylene

## Summary

## Perspectives

- Noise-Immune Cavity-Enhanced Optical Heterodyne Molecular Spectrometry vs Cavity RingDown Spectroscopy
- Both Absorption-based Techniques (in High Finesse Cavity)
  - Heterodyne (Absorbance) vs. Response Time (Absorption) Detection
  - Continuous vs. Discontinuous Acquisition (LOD),
  - "Trichromatic" vs. "Monochromatic" EMF
  - Ultimate Sensitivity: Photon-Shot-Noise Limited
  - Lockings (Cavity, ...)
- "Weak" Absorption (Optically Thin Medium)
- Saturation: Lamb Dip, Crossover Resonances vs. Cross-Sideband Resonances, etc...

#### • Combining both advantages:

- Radio-Frequency Modulation: Noise Immunity and Photon Shot Noise Ultimate Limit
- High Finesse Cavity: i.e., Long Equivalent Absorption Length (> km), Narrow Bandwidth
- High EMF Trapped
  - NonLinear Couplings
- Phase Modulation of the Source





## Weak Absorption in a Resonant Cavity

Equivalent Absorption Length:

$$L_{eq} = \frac{2 \mathscr{F} L_{cav}}{\pi}$$

Cavity "Bandwidth":

 $\Delta_{\omega} = \frac{FSR}{\mathscr{F}}$ 

EM Field Enhancement at resonance ( $\mathscr{R}_{cav} \sim 0$ ):

$$\mathcal{I}_{cav} = \frac{\mathcal{F}}{\pi} \mathcal{I}_{in}$$

Photon noise:

$$\sqrt{\frac{2 e \Delta_{v}}{\eta \langle P \rangle}} = (\alpha L_{eq})_{min}$$

## Interaction with a TriChromatic EMF in Cavity



## Background

#### • "2-level System"

- Analytical Solution of the Liouville Equation in the Frequency Domain
- NonLinear Polarization in the Frequency Domain
- 2-step Iteration (Moderate Saturation)
- Multichromatic EMF (forth and back), Dirac or Lorentzian Profile
- Numerical Integration over the Doppler Shift (Boltzmann Distribution)
- Lamb Dips and Cross-Sideband Resonances
- Zeeman Structure and EMF Polarization ( $\sigma$  or  $\pi$ )
- Zeeman Sublevels with a unique Relaxation Time
- Transit Time Approximation
- Direct Absorption, "Absorption-like", "Dispersion-like"
- Comparison with O. Axner Group Data

#### • Approximated Model

• Zeeman Sub-transitions (Averaged "Saturation Coefficient")

#### • Not in the Model

- Strong Saturation
- N-level System

#### • Recovering Dipole Moment and Number Density

## **Solutions**

$$\begin{split} \Delta \mathscr{P}_{m_{u}m_{l}}^{(DC)}(\omega_{0},\delta) &= \mathscr{P}_{0} \, l_{eq} \begin{pmatrix} F_{u} & 1 & F_{l} \\ -m_{u} \; \Delta m \; m_{l} \end{pmatrix}^{2} \times \\ & \frac{J_{0}^{2}(\xi) \; d\alpha_{ul}^{(0)}(\omega_{0},\delta) + J_{1}^{2}(\xi) \left[ d\alpha_{ul}^{(0)}(\omega_{0} + \omega_{rf},\delta) + d\alpha_{ul}^{(0)}(\omega_{0} - \omega_{rf},\delta) \right]}{1 + 3 \begin{pmatrix} F_{u} & 1 & F_{l} \\ -m_{u} \; \Delta m \; m_{l} \end{pmatrix}^{2} S_{ul}(\omega_{0},\delta)} \\ d\mathscr{P}_{m_{u}m_{l}}^{(\cos)}(\omega_{0},\delta) &= J_{0}(\xi) J_{1}(\xi) \; \mathscr{P}_{0} \, l_{eq} \times \begin{pmatrix} F_{u} & 1 & F_{l} \\ -m_{u} \; \Delta m \; m_{l} \end{pmatrix}^{2} \times \\ & \frac{d\alpha_{ul}^{(0)}(\omega_{0} + \omega_{rf},\delta) - d\alpha_{ul}^{(0)}(\omega_{0} - \omega_{rf},\delta)}{1 + 3 \begin{pmatrix} F_{u} & 1 & F_{l} \\ -m_{u} \; \Delta m \; m_{l} \end{pmatrix}^{2} S_{ul}(\omega_{0},\delta)} \\ d\mathscr{P}_{m_{u}m_{l}}^{(\sin)}(\omega_{0},\delta) &= J_{0}(\xi) J_{1}(\xi) \; \mathscr{P}_{0} \, l_{eq} \times \begin{pmatrix} F_{u} & 1 & F_{l} \\ -m_{u} \; \Delta m \; m_{l} \end{pmatrix}^{2} \times \\ & \frac{2d\phi_{ul}^{(0)}(\omega_{0},\delta) - d\phi_{ul}^{(0)}(\omega_{0} + \omega_{rf},\delta) - d\phi_{ul}^{(0)}(\omega_{0} - \omega_{rf},\delta)}{1 + 3 \begin{pmatrix} F_{u} & 1 & F_{l} \\ -m_{u} \; \Delta m \; m_{l} \end{pmatrix}^{2} S_{ul}(\omega_{0},\delta)} \end{split}$$

## And (Numerical Sums and Integrals):

$$\Delta \mathscr{P}_{ul}^{(DC)}(\omega_0) = 3 \sum_{m_u} \sum_{m_l} \int \Delta \mathscr{P}_{m_u m_l}^{(DC)}(\omega_0, \delta)$$

$$\mathscr{P}_{ul}^{(\cos)}(\omega_0) = 3\sum_{m_u}\sum_{m_l}\int \mathrm{d}\mathscr{P}_{m_u m_l}^{(\cos)}(\omega_0,\delta)$$

$$\mathcal{P}_{ul}^{(\sin)}(\omega_0) = 3 \sum_{m_u} \sum_{m_l} \int d\mathcal{P}_{m_u m_l}^{(\sin)}(\omega_0, \delta)$$

$$\Omega_{ul}^2 = \frac{\mathscr{F}}{\pi} \left| \frac{E_0 \mu_{ul}}{\hbar} \right|^2$$

$$\mathbf{d}\mathcal{N}_0(\delta) = \mathcal{N}_0 \sqrt{\frac{\ln 2}{\pi}} \frac{1}{\Delta_D} e^{-\ln 2\left(\frac{\delta}{\Delta_D}\right)^2} \mathbf{d}\delta$$

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Saturated NICE-OHMS

$$3 \begin{pmatrix} F_u & 1 & F_l \\ -m_u & \Delta m & m_l \end{pmatrix}^2 S_{ul}(\omega_0, \delta) =$$

$$2\Omega_{ul}^2 \frac{\gamma_{ul} T_1}{dg_{rot_l}} \sum_{\kappa} \left[ \frac{1}{(\omega_{ul} - \omega_0 + \kappa \omega_{rf} + \delta)^2 + \gamma_{ul}^2} + \frac{1}{(\omega_{ul} - \omega_0 + \omega_{rf} - \delta)^2 + \gamma_{ul}^2} \right]$$

## Simulation C<sub>2</sub>H<sub>2</sub>: $R_0$ (I = 0): Direct Absorption ( $\sigma$ Polar.)

 $1_0^1 3_0^1 (\Sigma_u^+ \leftarrow \Sigma_g^+), \mu = 10.7 \text{ mD}, \omega_{rf} = 380 \text{ MHz}, \xi = 0.4 (1 \text{ mTorr}, w_0 = 0.45 \text{ mm}, \Gamma_{rel.} = 1 \text{ MHz})$ 



## Simulation C<sub>2</sub>H<sub>2</sub>: $R_0$ (I = 0): Absorption-like ( $\sigma$ Pol.)

 $1_0^1 3_0^1 (\Sigma_u^+ \leftarrow \Sigma_g^+), \mu = 10.7 \text{ mD}, \omega_{rf} = 380 \text{ MHz}, \xi = 0.4 (1 \text{ mTorr}, w_0 = 0.45 \text{ mm}, \Gamma_{rel.} = 1 \text{ MHz})$ 



## Simulation C<sub>2</sub>H<sub>2</sub>: $R_0$ (I = 0): Dispersion-like ( $\sigma$ Pol.)

 $1_0^1 3_0^1 (\Sigma_u^+ \leftarrow \Sigma_g^+), \mu = 10.7 \text{ mD}, \omega_{rf} = 380 \text{ MHz}, \xi = 0.4 (1 \text{ mTorr}, w_0 = 0.45 \text{ mm}, \Gamma_{rel.} = 1 \text{ MHz})$ 



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## Simulation C<sub>2</sub>H<sub>2</sub>: $P_{11}(I = 1)$ , $\mathcal{N} = 3.2 \times 10^{11} \,\mathrm{cm}^{-3}$ , $\sigma$ Pol.

 $1_0^1 3_0^1 (\Sigma_u^+ \leftarrow \Sigma_g^+), \mu = 10.75 \,\mathrm{mD}, \omega_{rf} = 380 \,\mathrm{MHz}, \xi = 0.36 (w_0 = 0.54 \,\mathrm{mm}, \Gamma_{rel.} = 0.5 \,\mathrm{MHz})$ 



### Summary

- Analytical Solutions + Doppler Shift Integral (Numerical) + Sum over the Sub-Transitions ("saturation coefficients" < 100)
- The Formalism (involving 2 counter propagating beams) reproduces the Lamb-dips and the Cross-Sideband Resonances: Position, Intensity and Width
- Can be applied to any value of the Phase Modulation Index ( $\xi$ )
- Spatial Beam Profiles: Longitudinal and Transverse (through additional Numerical Integration)
- Can be applied to non-monochromatic Sources
- Benchmarked on 2 transitions of  $C_2H_2$  ( $P_{10}$ , I = 0, and  $P_{11}$ , I = 1) in  $\sigma$ Polarization (very good agreement with O. Axner data)
- Weak sensitivity of the "Dispersion-like" signal with the Saturation (under the present experimental conditions)
- Optimum Sensitivity: Modulation Index, and Doppler vs. Sub-Doppler
- Approximation by mean "saturation coefficient" (excellent for "Dispersion" signal)
- Applicable to Saturated FMS (i.e., without cavity)

## **Perspectives**

- Simultaneous determination of the Number Density and of the Transition Strength
- The Strong Saturation Regime?
- Extension to N-level System?
- Application to Multichromatic EMF (OFC)
- Metrology
- Line Shape Analysis
- Etc..

# **Thank for Your Attention**

# Ove Axner Swedish French Institute



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