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Dynamic topography of passive continental margins and their hinterlands since the Cretaceous

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Dynamic topography of passive continental margins and their hinterlands since the Cretaceous

Abstract

Even though it is well accepted that the Earth's surface topography has been affected by mantle-convection induced dynamic topography, its magnitude and time-dependence remain controversial. The dynamic influence to topographic change along continental margins is particularly difficult to unravel, because their stratigraphic record is dominated by tectonic subsidence caused by rifting. We follow a three-fold approach to estimate dynamic topographic change along passive margins based on a set of seven global mantle convection models. We first demonstrate that a geodynamic forward model that includes adiabatic and viscous heating in addition to internal heating from radiogenic sources, and a mantle viscosity profile with a gradual increase in viscosity below the mantle transition zone, provides a greatly improved match to the spectral range of residual topography end-members as compared with previous models at very long wavelengths (spherical degrees 2-3). We then combine global sea level estimates with predicted surface dynamic topography to evaluate the match between predicted continental flooding patterns and published paleo-coastlines by comparing predicted versus geologically reconstructed land fractions and spatial overlaps of flooded regions for individual continents since 140 Ma. Modelled versus geologically reconstructed land fractions match within 10% for most models, and the spatial overlaps of inundated regions are mostly between 85% and 100% for the Cenozoic, dropping to about 75-100% in the Cretaceous. Regions that have been strongly affected by mantle plumes are generally not captured well in our models, as plumes are suppressed in most of them, and our models with dynamically evolving plumes do not replicate the location and timing of observed plume products. We categorise the evolution of modelled dynamic topography in both continental interiors and along passive margins using cluster analysis to investigate how clusters of similar dynamic topography time series are distributed spatially. A subdivision of four clusters is found to best reveal end-members of dynamic topography evolution along passive margins and their hinterlands, differentiating topographic stability, long-term pronounced subsidence, initial stability over a dynamic high followed by moderate subsidence and regions that are relatively proximal to subduction zones with varied dynamic topography histories. Along passive continental margins the most commonly observed process is a gradual motion from dynamic highs towards lows during the fragmentation of Pangea, reflecting the location of many passive margins now over slabs sinking in the lower mantle. Our best-fit model results in up to 500 (± 150) m of total dynamic subsidence of continental interiors while along passive margins the maximum predicted dynamic topographic change over 140 million years is about 350 (± 150) m of subsidence. Models with plumes exhibit clusters of transient passive margin uplift of about 200 \pm 200 m, but are mainly characterised by long-term subsidence of up to 400 m. The good overall match between predicted dynamic topography to geologically mapped paleo-coastlines makes a convincing case that mantle-driven topographic change is a critical component of relative sea level change, and indeed the main driving force for generating the observed geometries and timings of large-scale continental inundation through time.

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1 **Dynamic topography of passive continental margins and their hinterlands**
2 **since the Cretaceous**

3

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22

23 **Abstract**

24 Even though it is well accepted that the Earth's surface topography has been affected by
25 mantle-convection induced dynamic topography, its magnitude and time-dependence
26 remain controversial. The dynamic influence to topographic change along continental
27 margins is particularly difficult to unravel, because their stratigraphic record is dominated by
28 tectonic subsidence caused by rifting. We follow a three-fold approach to estimate dynamic
29 topographic change along passive margins based on a set of seven global mantle convection
30 models. We first demonstrate that a geodynamic forward model that includes adiabatic and
31 viscous heating in addition to internal heating from radiogenic sources, and a mantle
32 viscosity profile with a gradual increase in viscosity below the mantle transition zone,
33 provides a greatly improved match to the spectral range of residual topography end-
34 members as compared with previous models at very long wavelengths (spherical degrees 2-
35 3). We then combine global sea level estimates with predicted surface dynamic topography
36 to evaluate the match between predicted continental flooding patterns and published paleo-
37 coastlines by comparing predicted versus geologically reconstructed land fractions and
38 spatial overlaps of flooded regions for individual continents since 140 Ma. Modelled versus
39 geologically reconstructed land fractions match within 10% for most models, and the spatial
40 overlaps of inundated regions are mostly between 85% and 100% for the Cenozoic, dropping
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42 plumes are generally not captured well in our models, as plumes are suppressed in most of
43 them, and our models with dynamically evolving plumes do not replicate the location and
44 timing of observed plume products. We categorise the evolution of modelled dynamic
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46 investigate how clusters of similar dynamic topography time series are distributed spatially.

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48 evolution along passive margins and their hinterlands, differentiating topographic stability,
49 long-term pronounced subsidence, initial stability over a dynamic high followed by moderate
50 subsidence and regions that are relatively proximal to subduction zones with varied dynamic
51 topography histories. Along passive continental margins the most commonly observed
52 process is a gradual motion from dynamic highs towards lows during the fragmentation of
53 Pangea, reflecting the location of many passive margins now over slabs sinking in the lower
54 mantle. Our best-fit model results in up to 500 (± 150) m of total dynamic subsidence of
55 continental interiors while along passive margins the maximum predicted dynamic
56 topographic change over 140 million years is about 350 (± 150) m of subsidence. Models
57 with plumes exhibit clusters of transient passive margin uplift of about 200 ± 200 m, but are
58 mainly characterised by long-term subsidence of up to 400 m. The good overall match
59 between predicted dynamic topography to geologically mapped paleo-coastlines makes a
60 convincing case that mantle-driven topographic change is a critical component of relative
61 sea level change, and indeed the main driving force for generating the observed geometries
62 and timings of large-scale continental inundation through time.

63

64 **Introduction**

65 The vertical motions and water depth of passive margins are dominated by the intensity of
66 lithospheric thinning, and sediment accumulation through time (Kirschner et al., 2010). A
67 number of mechanisms have been suggested to account for additional, anomalous vertical
68 motions of passive margins, and many lines of evidence suggest that there is no single
69 mechanism that can account for all observed subsidence and uplift anomalies in this context.
70 Changes in intraplate stresses have been widely inferred to cause flexure, either uplift or

71 subsidence, and inversion along passive margins (Cloetingh, 1988; Lowell, 1995). However,
72 there are few published models for intraplate stress variations through geological time that
73 could be used to predict their effect on basins and margins, and inversion of faults
74 underlying passive margins is also relatively localised. Recently, Yamato et al. (2013)
75 proposed that major changes in mantle convection regimes can induce margin compression
76 and uplift, while Schiffer and Nielsen (2016) investigated the effect of plumes on margin
77 uplift and changes in lithospheric stress in the North Atlantic. Japsen et al. (2012) favoured
78 lithospheric-scale folding at craton boundaries as a universal explanation for anomalous
79 margin uplift, but observations supporting this idea are limited to only a few regions. Braun
80 (2010) reviewed the expressions of mantle dynamic surface topography on continental
81 interiors globally, without specifically investigating the effect on continental margins; this
82 partly reflects that dynamic topographic changes affecting continents are more readily
83 observed in continental interiors far away from geologically recent plate deformation.

84

85 Recently, the potential influence of mantle-driven dynamic topography at present-day was
86 analysed in a number of different ways, comparing a variety of observations and
87 assumptions to derive residual, non-isostatic topography with geodynamic model
88 predictions (Hoggard et al., 2016). However, present-day estimates of residual topography
89 alone do not provide insights into dynamic topography affecting continents and their
90 margins through time, as dynamic topography by its nature is time variable. In response to
91 the need to understand the long-term effect of plate-mantle interaction on passive margins,
92 different mantle convection approaches have been developed (e.g. Gurnis, 1993). A widely-
93 used approach for modelling these geodynamic processes in the recent geological past is the
94 inversion of tomographically imaged mantle structure together with other observations to

95 model surface dynamic topography (see Flament et al., 2013, for a review), as recently
96 applied, for example, to the last three million years to the eastern margin of the United
97 States by (Moucha et al., 2008). However, this approach is not useful for modelling these
98 processes through deep geological time, because the current mantle structure does not
99 provide sufficient information to model plate mantle interaction since the breakup of the
100 supercontinent Pangea, in the course of which most current passive margins formed. This
101 issue was recently evaluated in detail by Rowley et al. (2013), who confirmed that
102 retrodictions of mantle flow into deep geological time are possible only if the plate velocity
103 field is used as an additional constraint. It is well established that a plate motion model is
104 needed to model plate-mantle interactions in deep geological time (see recent discussions of
105 the opportunities and limitations of this approach by Colli et al. (2015), as fully dynamic
106 mantle convection models are not yet able to reproduce the evolution of the plate-mantle
107 system realistically; however, recently developed sequential data assimilation methods
108 (Bocher et al., 2016; Colli et al., 2015), currently only tested in 2D simulations, and adjoint
109 methods (Li et al., 2017), hold the promise of more physically realistic plate-mantle models
110 to evaluate the effect of mantle dynamics on surface topography.

111

112 In order to overcome some of these current methodological limits, a geodynamic forward
113 modelling approach with time-dependent slab assimilation constrained by a global tectonic
114 model has been developed (Bower et al., 2015). This method has been applied previously to
115 investigate the role of mantle convection in driving large-scale (> 1000 km wavelength)
116 anomalous subsidence or uplift of passive margins in a number of regions including the east
117 coast of North America (Flament et al., 2013), the South Atlantic domain (Flament et al.,

118 2014), the Arctic (Shephard et al., 2014), Southeast Asia (Zahirovic et al., 2016), northern
119 Africa (Barnett - Moore et al., 2017) and the east Australian margin (Müller et al., 2016).
120
121 However, this approach has not yet been evaluated in the context of a global analysis of the
122 influence of large-scale mantle flow on the subsidence and uplift of passive margins and
123 their hinterlands through time. This partly reflects the difficulties in comparing the output of
124 global mantle dynamic flow models with detailed local observations from continental
125 margins, either based on present-day residual topography analysis (Hoggard et al., 2016) or
126 stratigraphic data from wells providing estimates of tectonic subsidence or uplift anomalies
127 through time (e.g. Xie et al., 2006). Currently observed residual topography may reflect
128 several processes other than large-scale mantle convection, including lithospheric thickness
129 and/or density anomalies (e.g. Xie et al., 2006), as well as asthenospheric temperature
130 anomalies and small-scale convection (Pedersen et al., 2016), complicating its interpretation.
131 Similarly backstripped tectonic subsidence derived from individual wells may be influenced
132 by local tectonic reactivation and faulting processes (e.g. Colli et al., 2014; Hoggard et al.,
133 2016), in addition to subsidence following rift-related lithospheric thinning, partly obscuring
134 dynamic topography signals from large-scale mantle convection ((e.g. Johan and Kleinspehn,
135 2000) for a discussion of the interaction of these signals in the South China Sea). Because of
136 the difficulties in assessing the effect of deep mantle convection-driven dynamic topography
137 for passive continental margins directly, we follow a three-fold approach here. We first
138 analyse the power spectra of a set of seven alternative mantle convection models in the
139 context of the spectra of residual topography end-members to establish their relationship at
140 long wavelengths, and then combine global sea level estimates (Fig. 1) with predicted
141 surface dynamic topography (Figs 2-4) to evaluate modelled continental inundation as

142 compared with published paleo-coastlines through time. Subsequently, we use an
143 established cluster analysis approach to investigate the time-dependent dynamic
144 topography evolution predicted for continental interiors and passive margins by our models.

145

146 **Methods**

147 ***Plate reconstructions***

148 We model global mantle flow based on the subduction and plate motion histories predicted
149 by topologically-evolving plate boundaries from three alternative plate reconstructions
150 (Gurnis et al., 2012; Müller et al., 2016; Seton et al., 2012). We use the reconstruction by
151 Müller et al. (2016) as reference, because it includes many recent improvements to our
152 model of regional relative plate motions and plate boundary evolution, including revised
153 maps of the evolution of the age-area distribution of the ocean floor through time, providing
154 improved constraints for the age and thus thickness of subducting oceanic lithosphere
155 through time, a key constraint for assimilating subducting lithosphere into mantle
156 convection models. We also use the reconstruction by Seton et al. (2012) as it has been the
157 reference plate tectonic model for post-Pangea geodynamic modelling for the last few years,
158 providing an opportunity to evaluate the different choices for absolute plate motion models
159 that were made by Seton et al. (2012) versus Müller et al. (2016), especially considering that
160 in the latter model episodes of large global RMS plate velocity, net rotation, and trench
161 migration were minimised to reduce potential artefacts in forward geodynamic models. The
162 tectonic reconstruction by Gurnis et al. (2012), an earlier version of the tectonic
163 reconstruction by Seton et al. (2012) for the last 140 million years, is used to evaluate the
164 predictions of the mantle flow models by Spasojevic and Gurnis (2012) for the last 90 million
165 years.

166

167 ***Geodynamic models***

168 Our calculations begin at 230 Ma in all models with the exception of that by Spasojevic and
169 Gurnis (2012), but we only analyse mantle evolution from the Early Cretaceous (140 Ma)
170 since it takes at least 50 million years for the models to reach an equilibrium from the initial
171 condition (Flament et al., 2014), and because published digital paleo-coastline maps are
172 available only for the period after 140 Ma. The earlier period of forward integration is
173 avoided in Model M1, a hybrid model (Spasojevic and Gurnis, 2012), as the initial global
174 mantle temperature field at present-day is estimated through a combination of seismic
175 tomographic inversions of surface and body waves using model S2ORTS (Ritsema et al.,
176 2004) in the lower mantle and one based on Benioff zone seismicity for the upper mantle
177 seismicity. This temperature field is integrated backward using the SBI (simple backward
178 integration) method of Liu and Gurnis (2008) back to the Late Cretaceous by reversing the
179 direction of gravity and plate motions. A hybrid paleo-buoyancy field is generated by
180 merging the backward-advected mantle temperature field with synthetic subducted slabs
181 assimilated into the model based on the location of subduction zones, the age of the
182 subducted lithosphere and relations among subduction zone parameters (Spasojevic and
183 Gurnis, 2012). In all other models analysed here viscous mantle flow is driven in forward
184 models by thermal convection with plate velocities applied as surface boundary conditions,
185 extracted in 1 million year intervals from the plate reconstructions (Bower et al., 2015). The
186 initial condition in models M2-M4 without plumes includes a basal thermochemical layer
187 113 km thick just above the core–mantle boundary (CMB) that consists of material 4.2%
188 denser than ambient mantle, while in model M7 this layer is 10% denser than ambient
189 mantle. This condition effectively suppresses plumes in the model within the time frame

190 covered by our model runs. This setup prevents the formation of upwelling mantle plumes,
191 making it possible to study the interaction of moving continents with subduction-driven
192 mantle downwellings and the associated large-scale mantle return flow in the absence of
193 individual plumes. The initial condition for models with plumes features a basal chemical
194 layer 100 km thick that is 2.5% heavier than the ambient mantle, embedded in a 300 km
195 thick thermal boundary layer - see Hassan et al. (2015) for a more detailed description of the
196 model setup. The thickness of the thermal lithosphere, derived from the age of the oceanic
197 lithosphere and tectono-thermal age of the continental lithosphere, is assimilated into the
198 dynamic model. We use a modified version of the finite element code CitcomS to obtain
199 one-sided subduction, in which the shallow portion of subducting slabs is imposed to a
200 maximum depth of 350 km, below which mantle convection arises dynamically from
201 prescribed time-dependent conditions (Bower et al., 2015). In models M2-M7 air-loaded
202 dynamic topography is calculated from the surface vertical stress resulting from mantle flow
203 in restarts of the main model run in which the surface boundary condition is free-slip and
204 the 350 km uppermost part of the mantle do not contribute to the flow, while lateral
205 viscosity variations are preserved in the whole mantle. In contrast, in model M1 dynamic
206 topography is computed using a no slip surface boundary condition, and only the top 250 km
207 of the mantle are excluded from contributing to surface topography, which partly explains
208 the somewhat greater amplitude of dynamic topography in model M1 as compared to
209 models M2-M7 (see also Flament et al., 2014; Thoraval and Richards, 1997). Finally, the use
210 of models based only on forward calculations versus those through inversion using seismic
211 constraints allows us to evaluate the role of initiation conditions and fits to present-day
212 seismic structure.

213

214 The Rayleigh number that determines the vigour of convection is defined by

215
$$Ra = \frac{\alpha_0 \rho_0 g_0 \Delta T h_M^3}{\kappa_0 \eta_0},$$

216 where α is the coefficient of thermal expansion, ρ the density, g the acceleration of gravity,
217 ΔT the non-adiabatic temperature change across the mantle, h_M the depth of the mantle, κ
218 the thermal diffusivity, and η the viscosity in which the subscript “0” indicates reference
219 values. Key model parameters are listed in Table 1.

220

221 The average model resolution, obtained with $\sim 13 \times 10^6$ nodes and radial mesh refinement, is
222 $\sim 50 \times 50 \times 15$ km at the surface, $\sim 28 \times 28 \times 27$ km at the core–mantle boundary (CMB), and
223 $\sim 40 \times 40 \times 100$ km in the mid-mantle. The modelled dynamic topography through time is
224 computed in the mantle and plate frames of reference for models M1-M7, exploring the
225 parameter space in terms of the alternative plate reconstructions and in assumed mantle
226 viscosities (Table 1). In models M2-M4 temperature and thermal expansion are constant
227 with depth and the dense basal layer has the same thickness as the lower thermal boundary
228 layer. In models M5-M7 thermal expansion decreases with depth, the average mantle
229 temperature increases with depth (adiabat) and the dense basal layer is thinner than the
230 lower thermal boundary layer. In models M2-4, the viscosity varies by 1000 due to
231 temperature-dependence following

232

233
$$\eta(T, r) = A(r) \eta_0 \exp\left(\frac{E_\eta}{T + T_\eta} - \frac{E_\eta}{T_b + T_\eta}\right),$$

234 where r is the radius, where $A(r)$ is the pre-exponential parameter for four layers (Table 2),
235 E_η is the non-dimensional activation energy ($E_{UM} = 9$ in the upper mantle and $E_{LM} = 3$ in the
236 lower mantle, see Table 2 for dimensional values), T is the temperature, $T_\eta = 0.16$ is a

237 temperature offset, and $T_b = 0.5$ is the ambient mantle temperature (see Table 2 for
 238 dimensional values). A similar viscosity parameterization is employed in model M1 (Table 2).
 239 In models M5-7, we use piecewise Arrhenius laws to describe the variation of viscosity with
 240 temperature and depth, which takes the following form:

$$\eta(T, r) = A(r)\eta_0 \exp\left(\frac{E_\eta(r) + (1-r)V_\eta(r)}{T + T_\eta} - \frac{E_\eta(r) + (1-r_{inner})V_\eta(r)}{1 + T_\eta}\right),$$

241 where V_η is the non-dimensional activation volume. For the lower mantle, we use a
 242 dimensional activation energy of 320 KJ mol^{-1} and activation volume of 6.7×10^{-6}
 243 $\text{m}^3 \text{ mol}^{-1}$, corresponding to non-dimensional units of 11 and 26, respectively, which are
 244 comparable to estimates (Karato and Wu, 1993). However, since such viscosity
 245 parameterizations lead to large viscosity variations that cause numerical difficulties, we
 246 adjust the pre-exponential parameter $A(r)$ and the temperature offset T_η (Tackley, 1996) to
 247 limit the viscosity contrast to 3 orders of magnitude. The resulting viscosity profile is similar
 248 to the preferred viscosity profiles of Steinberger and Calderwood (2006). Models M5-M7
 249 also assume an *a priori* mantle adiabat with top and bottom thermal boundary layers
 250 featuring a temperature drop of 1225 K and the initial adiabatic temperature profile has a
 251 potential temperature of 1525 K. Moreover, models M5-7 incorporate viscous dissipation
 252 and adiabatic heating through the extended Boussinesq approximation (c.f. Christensen and
 253 Yuen, 1985; Ita and King, 1994), in addition to internal heating from radiogenic sources. The
 254 Boussinesq and Extended Boussinesq Approximations (BA and EBA, respectively) assume an
 255 incompressible mantle, while the Anelastic Liquid Approximation (ALA), not utilized here, is
 256 used to model mantle compressibility (Ita and King, 1994). However, unlike in the BA,
 257 models based on the EBA allow for the inclusion of a pre-calculated thermodynamic model,

258 featuring depth profiles for adiabatic temperature, thermal expansivity and density – see Ita
259 and King (1994) for more detail. We apply a non-dimensional internal heat generation rate
260 of 100 in these model cases and compute a reference profile for thermal expansion, α , based
261 on analytical parameterizations of Tosi *et al.* (2013), using the *a priori* mantle temperature
262 profile. Additional details on the setup of models M5-M7 can be found in Hassan *et al.*
263 (2015).

264

265 ***Model power spectra***

266 We use spherical harmonic analysis to analyse the spectral characteristics of our global
267 dynamic topography models at present-day for spherical harmonic degrees 1-5, and
268 compare them with two end-member spectra of oceanic residual topography from Hoggard
269 *et al.* (2016, Fig. S2), which capture one of the main sources of uncertainty in computing
270 oceanic residual topography, the flattening of old ocean floor (Fig. 5). The end member
271 spectra are based on Crosby and McKenzie's (2009) plate model (PM) and a thermal
272 boundary layer (TBL) model fitted to ocean floor less than 70 million years in age (Fig. S2 of
273 Hoggard *et al.* (2016). These residual topography models are well suited to evaluate the
274 amplitudes of long-wavelength dynamic topography as we expect long-wavelength dynamic
275 topography predicted from mantle convection models to fall within the residual topography
276 implied by a plate model and a thermal boundary layer model. This is based on the
277 consideration that plate models potentially overestimate the flattening of old ocean floor, as
278 the observed flattening in at least some regions may be due to deep mantle convection-
279 driven dynamic topography, while thermal boundary layer models underestimate it (see
280 Stein and Stein, 2015, for a recent review). Estimates of continental residual topography are

281 much more uncertain (Yang and Gurnis, 2016), reflecting the considerable heterogeneity of
282 continental lithosphere in terms of its temperature, density and thickness, and not used
283 here.

284

285 ***Paleogeography models***

286 In order to evaluate modelled surface dynamic topography against geologically mapped
287 paleogeography we use two alternative sets of published digital paleo-coastlines (Golonka,
288 2007, 2009; Smith et al., 1994) for the period 140-0 Ma, which have been rotated into
289 present-day coordinates (Heine et al., 2015) to allow adaptation to different plate models
290 with ease. We interpolate these paleo-coastlines into discrete paleogeographic maps that
291 delineate the distribution of land and water by first generating points on the sphere that
292 sample equal areas (Gorski et al., 2005) in order to produce binary maps (land or marine) for
293 any given paleo-coastline configuration that is taken to be representative for a certain
294 geological time interval (Heine et al., 2015). We then compute the intersection P of binary
295 maps for a given pair of consecutive time-slices (S_{t-1}, S_t). This intersection is a collection of
296 isolated patches that represent regions being submerged, P_w , or becoming emergent, P_l ,
297 from time S_{t-1} to time S_t that is, $P = P_w \cup P_l$. We then compute a distance map of P to
298 closest land in S_t , based on a spherical distance metric and normalize the distance map such
299 that within each isolated patch distances range between $[0, 1]$ and relate normalized time
300 progressions between S_{t-1} and S_t to the normalized distance map. In other words, as time
301 progresses between times S_{t-1} and S_t , regions featured in the distance map morph from
302 their former state in S_{t-1} to their subsequent state in S_t , based on corresponding values of
303 distances. These transitions proceed such that regions in P_w farthest from land are
304 inundated first, whereas regions in P_l closest to land emerge first. In the absence of detailed

305 continental paleo-elevation constraints, we assume a linear progression of transgressions
306 and regressions, and obtain interpolated paleogeographic maps at one million year intervals.
307 In the following, we refer to these as paleogeographic models, to differentiate them from
308 geodynamic models.

309

310 ***Model topography versus continental paleogeography through time***

311 The observed present-day topography includes dynamic topography, but given the large
312 uncertainties in both estimates of amplitudes of continental residual topography as well as
313 dynamic models (see recent extensive review by Guerri et al., 2016) we depend on using a
314 simple empirical approach to be able to compare our modelled dynamic topography with
315 mapped paleo-coastlines. The dynamic topography from our models cannot be used in its
316 raw form here, reflecting uncertainties in dynamic topography amplitudes as a consequence
317 of our still limited knowledge of mantle physical properties and chemical composition
318 (Guerri et al., 2016).

319 We compare the distribution of land and sea in continental regions predicted by global
320 convection models with those from geologically mapped paleo-coastlines (Heine et al., 2015)
321 (except Antarctica), based on a methodology similar to that outlined in Gurnis (1993). We
322 define a ‘continent function’, $C(t, \theta, \varphi)$, which is a binary map representing the locations of
323 continents through time t and space (θ, φ) . At each time instance t , we apply the ‘continent
324 function’, as a mask to the modelled dynamic topography to obtain continental dynamic
325 topography $h(\theta, \varphi)$. Since the geodynamic models presented in this study exclude a
326 significant portion of the upper mantle – ranging from 250 to 350 km from the surface –
327 while computing dynamic topography, the contribution of shallow mantle processes and the
328 spatial variability of lithospheric thickness toward modulating the computed long-

329 wavelength dynamic topography cannot be modelled. Consequently, the mean computed
330 dynamic topography within individual continents (Fig. 6), at any given time, tends to be
331 discrepant with geological observations. As a simple remedy, we remove the mean dynamic
332 topography within each continent from $h(\theta, \varphi)$, which is a reasonable approach given the
333 current limitations of the models. We then derive binary maps of continental inundation
334 patterns based on the inequality below:

$$h(\theta, \varphi) < S(t) + h_c$$

336
337 where h_c is a constant and $S(t)$ is the independently derived eustatic sea-level at the time
338 (Fig. 1), taken alternatively from Spasojevic and Gurnis (2012), whose sea level curve is
339 based on Müller et al.'s (2008) ocean basin volume estimates, or from sea level curves based
340 on revised age-area distributions in the ocean basins and revised changes in continental area
341 through time for the last 140 million years from Seton et al. (2012) and Müller et al. (2016),
342 including the same approach used in Müller et al. (2008) to account for the effect of oceanic
343 plateaus and deep-sea sediment thickness through time. These estimates are combined
344 with the mean oceanic dynamic topography effect of a given geodynamic model through
345 time (Fig. 7). We do not consider the effect of mean continental dynamic topography on
346 global sea level change (Gurnis, 1993; Conrad and Husson, 2009; Spasojevic and Gurnis,
347 2012), and do not attempt to calculate gravitationally self-consistent sea level changes
348 (Austermann and Mitrovica, 2015; Spasojevic and Gurnis, 2012) back in time, due to our
349 currently lacking ability to construct paleo-digital continental elevation models that are
350 corrected for lithospheric thinning and thickening, erosion and sedimentation back to 140
351 Ma. The choice of ocean volume-derived sea-level curve depends on which plate tectonic

352 model a given geodynamic model is based on. For both the sea level curves based on Seton
353 et al.'s (2012) and Müller et al.'s (2016) oceanic age-area distribution through time, we
354 include the long-term fluctuations in global sea-level from varying continental ice volume
355 since 38 Ma, derived from a zonally averaged energy balance climate model bi-directionally
356 coupled to a one-dimensional ice sheet model (Stap et al., 2016). This approach implies that
357 the continental ice volume between the Early Cretaceous and the Eocene was negligible.
358 There is evidence for the existence of a 5 million year long glacio-eustatic low-stand in the
359 mid-Cretaceous from the flooding record of the Arabian Plate (Maurer et al., 2013) which is
360 not included in our long-term sea level curve. We do not make any ice-volume related
361 adjustment to the sea level curve used by Spasojevic and Gurnis (2012), because eustatic sea
362 level and the regionally fluctuating dynamic topography are self-consistently derived from
363 their geodynamic model and we prefer to retain the relationship between eustasy and
364 dynamic topography as used in their model M5 (here model M1).

365 In the inequality used to derive continental inundation patterns, the constant h_c is
366 empirically constrained for each convection model separately, such that the spatial extent of
367 predicted flooding in the Early Cretaceous (or Late Cretaceous for model M1) is comparable
368 to those inferred in the paleogeographic maps. This simple empirical approach enables us to
369 generate rough estimates of the flooding of continents through time by combining dynamic
370 topography and eustasy. Long wavelength, space-time varying dynamic topography lows are
371 thus adopted as proxies for shallow marine seas that may arise when continents travel
372 above regions where subducted slab material sinks in the mantle, modulated by global sea
373 level fluctuations.

374

375 We compute the fraction of land through time, $\tau(t)$, for each continent, as predicted by the
376 convection models with those from two alternative sets of paleogeography grids (Golonka,
377 2007, 2009; Smith et al., 1994) derived from the digital paleo-coastlines in Heine et al. (2015)
378 (Figs 8-10b, Supp. Figs 6-9b). Land fractions computed for predictions from the convection
379 models are based on the present-day boundaries between continental and oceanic crust
380 (COB) (Müller et al., 2016), while those computed for the paleogeography grids are based on
381 paleo-coastlines. Thus, in order to make meaningful comparisons between these derived
382 land fractions, we normalize each curve by their corresponding maxima, considering that
383 paleo-coastlines record maximum flooding through time. We further compute similarity
384 coefficients through time, $\mu(t)$, which show the degree of spatial overlap of the distribution
385 of land or ocean between model predictions and the two sets of paleogeography grids. The
386 similarity coefficients are also normalized – for similar reasons – by their corresponding
387 maxima, and we refer to this measure in the following as inundation overlap.

388

389 We evaluate the overall quality of our models by computing the difference between
390 similarity coefficients for geodynamically modelled and geologically mapped continental
391 inundation, given in terms of land fraction differences, for individual continents as well as
392 the root mean square difference averaged for all continents. These measures highlight
393 whether the models over- or underestimate inundation for individual continents, and
394 summarise the overall performance of a given model to match geologically mapped
395 inundation through time. These statistics are computed for three time periods, i.e. 0-60 Ma,
396 60-100 Ma (90 Ma for M1) and 100-140 Ma, roughly representing the Cenozoic, Late and
397 Early Cretaceous periods, considering that model matches to continental inundation
398 patterns tend to vary significantly between these periods. In addition we compute the

399 overlaps between our model inundation predictions for individual continents and time
400 periods and geologically modelled inundation patterns. This measure reveals to what extent
401 predicted continental inundation overlaps spatially with geologically mapped inundation,
402 and is a useful measure in addition to land fraction similarities.

403

404 ***Cluster analysis of dynamic topography***

405 We use a *k-means* clustering algorithm (Lloyd, 1982) to obtain objective classifications of
406 geographic regions that share similar uplift and subsidence histories, predicted by the
407 geodynamic models. The *k-means* algorithm partitions data items into k clusters such that
408 the sum of the distances over the data items in each cluster to their cluster centre is
409 minimal, i.e., given a set of observations (x_1, x_2, \dots, x_n) , where each observation can be d
410 dimensional, the algorithm partitions the n observations into $k(\leq n)$ sets S . In our case, each
411 observation x_i is a time-series of dynamic topography estimates at a given location; each set
412 S_j identifies localities on continents that share a common uplift and subsidence history.
413 It should be noted that k , the number of clusters, is chosen *a priori* and we chose a range of
414 values of k to draw out dominant trends within dynamic topography predicted by our
415 mantle convection models. The resulting cluster maps (Figs. 11-12) represent an objective
416 regional summary of dynamic topography trends through time as predicted by a range of
417 geodynamic models based on different plate tectonic models, mantle rheologies and other
418 parameters described in the geodynamic modelling methodology section. We further
419 separately analyse continental passive margins, including all passive margins irrespective of
420 their age, using bands of a fixed width of 200 km. We include the dynamic topography
421 history preceding continental rifting for margins younger than 150 Ma.

422

423 **Results**

424

425 ***Model power spectra***

426 Models M1-4 overestimate oceanic residual oceanic topography at degrees 1-3, while the
427 power of models M5-7 at degrees 2-3 is within the range of end-member values based on a
428 plate versus thermal boundary layer model (Fig. 5). M5-M7 display more power at degree 1
429 than M2-M4, and M1 displays the most power at degrees 1-3 amongst all models used here.
430 In models M5-M7, the Pacific region features a strong “superplume”-like upwelling not
431 found in the Indo-Atlantic (Supp. Fig. 5). This pronounced large-scale Pacific upwelling
432 corresponds low seismic velocities in mantle tomography (French and Romanowicz, 2015).
433 The degree 1 spherical harmonic (Supp. Fig. 5) captures this Pacific upwelling and becomes
434 dominant in the power spectrum. This is the case even in model M7 with an initial 10%
435 denser basal thermochemical layer that suppresses the formation of mantle plumes; but
436 even here, the large-scale Pacific upwelling is more pronounced than in the Indo-Atlantic. As
437 a consequence, the degree 1 spherical harmonic that captures the sub-Pacific upwelling
438 becomes dominant in the power spectrum. It is important to note that this degree 1 peak
439 driven by a Pacific “superplume” does not influence our results for the continents, which do
440 not intersect its periphery, with the exception of the western portion of South America and
441 eastern Australia (Supp. Fig. 5), with the latter having experienced renewed uplift of its
442 eastern highlands in the Late Cenozoic perhaps due to overriding the edge of this upwelling
443 (Müller et al., 2016).

444

445 The r.m.s. amplitude of model M7, as a representative case of models M5-M7, at degree 1 is
446 1050 m, while end member residual topography models (Fig. 5) yield 350m (TBL) and 270m

447 (PM), mainly illustrating that the amplitude of large-scale Pacific mantle upwelling is
448 overestimated. However, at degree 2, the r.m.s. amplitude of model M7 is 740 m, within
449 the range of residual topography end-members at 850 m (TBL) and 530 m (PM). At degree
450 3, the r.m.s. amplitude of M7 is 570 m, also within a plausible residual topography range of
451 770 m (TBL) and 510 m (PM). This analysis indicates that model M7, as well as M5-6, yield
452 dynamic topography amplitudes which are consistent with observations at long
453 wavelengths, with the exception of degree 1, making them promising candidates for
454 understanding continental inundation through time. In models M5-7 viscous heating causes
455 slab interiors to weaken over timescales of a few million years (e.g. Larsen et al. (1995).
456 Consequently, the negative dynamic topography associated with sinking slabs diminishes
457 with the warming of slab material during its descent in the mantle. Although the models
458 here have unrealistically high amplitudes at degrees 1-3, by adjusting the mean of individual
459 models, as outlined above, we ensure that models can be used to investigate the role of
460 dynamic topography on continental inundation patterns through time.

461

462 ***Geodynamically modelled continental inundation versus paleogeography***

463 We evaluate the combination of predicted dynamic surface topography jointly with eustatic
464 sea level curves based on modelled ocean basin volumes from the plate model that was
465 used as surface boundary condition for a given geodynamic model combined with the
466 oceanic dynamic topography predicted by the same geodynamic model (Figs 2-4; supp. Figs
467 1-4) with regard to their match to published paleo-coastline locations in a plate reference
468 frame (Figs 8a-10a; supp. Figs. 6a-9a), first in terms of predicted versus geologically mapped
469 land fractions through time and similarities in the degree of spatial overlap of the

470 distribution of submerged continental regions between model predictions and the two sets
471 of paleogeography grids (Figs 8b-10b; supp. Figs. 6b-9b).

472

473 Geodynamic model M1 (Fig. 2), based on combined backward advection models of the
474 present-day mantle structure and forward subduction modelling (Spasojevic and Gurnis,
475 2012), matches the two paleogeography models for the Cenozoic compared here (Golonka,
476 2009; Smith et al., 1994) fairly well for North and South America, while somewhat less well
477 for the land fractions and inundation overlaps of other continents (Fig. 8b); the overall fit of
478 this model to either paleogeography deteriorates somewhat from the Cenozoic into the Late
479 Cretaceous (Fig. 8b). However, the trend of increased inundation (decreased land fraction)
480 of both Americas, and to some extent Eurasia, from the Early Cenozoic back into the Late
481 Cretaceous is well captured by this model (Fig. 8b). For the Cenozoic, inundation overlaps for
482 model M1 are similar for both paleogeographic reconstructions (Golonka, 2009; Smith et al.,
483 1994), while for the Late Cretaceous period (~100-65 Ma) the two reconstructed
484 paleogeographies diverge more significantly from each other (Fig. 8b). Generally, this results
485 in a better agreement of the geodynamic model with Golonka's (2007) model than with
486 Smith et al.'s (1994) model before 65 Ma, with the exception of North America (Fig. 8b),
487 where this order is reversed.

488

489 Next we evaluate the predicted dynamic topography through time from three forward
490 geodynamic models (M2-4) driven by alternative plate models (Fig. 3, supp. Figs 1 and 2). In
491 terms of modelled land fractions and inundation patterns, models M3 and M4 fit South
492 America relatively well, but only for Golonka's paleogeography (Golonka, 2007, 2009) (Fig. 9,
493 Supp. Figs 6b, 7b). The good fit in this region reflects the modelled inclusion of inferred

494 episodes of flat slab subduction along South America (Flament et al., 2015). Model M2
495 (Supp. Fig. 6b) based on the older (Seton et al., 2012) reconstruction, fits less well, partly
496 reflecting its reduced global sea level amplitude compared with the two other two models
497 (Müller et al., 2008; Müller et al., 2016) as well as a different history of the age of subducting
498 crust through time along the Andes. The evolution and eastward migration of the western
499 interior seaway in North America (Liu et al., 2014), reflecting the effect of the Laramide flat
500 slab (English et al., 2003) on surface dynamic topography, is moderately well captured in
501 these geodynamic models, but again only as compared with Golonka's paleogeography.
502 However, in these geodynamic models North America remains excessively flooded in the
503 early-mid Cenozoic (Fig. 9b), likely reflecting that Laramide slab breakoff and sinking into the
504 lower mantle (potentially resulting in surface rebound) is not appropriately captured in
505 these models.

506

507 Models M2-M4 capture Eurasia's flooding history relatively well over the Late Cenozoic, but
508 tend to underestimate flooding in the Early Cenozoic, likely reflecting the complex tectonic
509 history of Eurasia that results in regional flooding not accounted for in our models. Even
510 though the modelled Cretaceous land fractions match Golonka's paleogeography Golonka
511 (2007) for Eurasia, the inundation overlaps are generally not better than 80% for and
512 between 65 and 75% for Smith et al. (1994) (Fig. 9b). Africa's land fraction is poorly matched
513 for this entire set of models (Fig. 9b, Supp. Figs 6b, 7b). For Australia, models M2-4
514 underestimate Cenozoic flooding of Australia, likely reflecting shortcomings in the modelled
515 subduction history around Australia in all plate models used here. The early Cretaceous
516 flooding followed by Late Cretaceous rebound is generally captured with a 5-15 Myr time lag
517 depending on the plate motion model used (Fig. 9b, Supp. Figs 6b, 7b). Both subduction

518 zone locations east of Australia as well as the absolute plate motion interaction between
519 Australia and sinking slabs in the mantle depend on particular plate models (see Müller et al.
520 (2016) for discussion).

521

522 Models M5-7, being forward subduction models like models M2-4, but with a pressure-
523 dependent viscosity structure that involves a more gradual increase in viscosity between the
524 upper and lower mantle and uses an extended Boussinesq approximation, are characterized
525 by somewhat smaller amplitudes in surface dynamic topography (e.g. see M7, Fig. 4). In
526 addition, models M5 and M6 both include plumes, but their exact location and arrival time
527 at the surface cannot be controlled as the plumes develop dynamically in the lower mantle
528 (Supp. Figs 3, 4). As a consequence, continents that are affected by model plumes through
529 time, particularly Africa and Australia (M5 and M6) show different and sometimes unrealistic
530 inundation patterns (Supp. Figs 8b, 9b) compared with model M7, in which plumes are
531 suppressed (Fig. 10). Hassan et al. (2015) demonstrated that in model M5 (identical to their
532 model M3) that plumes arise in locations near to present day hot spot locations at a
533 statistically significant level. However, the exact arrival time and location of a given plume
534 head at the surface as well as the subsequent lateral plume motion or tilt is dynamic rather
535 than imposed, resulting in a variable match with geological observations. Model M5,
536 characterised by the evolution of an Afar-like plume (Fig. 4), results in the most reasonable
537 Cenozoic flooding history of Africa amongst all the geodynamic models analysed here
538 (Fig. 10b), while the opposite holds for Australia, where the evolution of model mantle
539 upwellings in M5 worsens the fit to paleogeography (Fig. 10b). In contrast, matches
540 between geodynamic model predictions with paleogeography for Cenozoic Eurasia are
541 relatively unaffected by the presence or absence of plumes (Fig. 10b), reflecting that most of

542 continental Eurasia is not affected by mantle plumes since the Cretaceous. However, it is
543 worth noting that all seven geodynamic models perform poorly in terms of Cretaceous
544 patterns of inundation in Eurasia, partly reflecting Eurasia's tectonic complexity (De Grave
545 and Buslov, 2007). Model M7, in which plumes are suppressed, results in the best overall
546 matches to continental flooding amongst models M5-7 (Fig. 10b), with the exception of
547 Africa.

548

549 Reconstructed paleo-coastlines not only reflect eustasy and dynamic topography, but also
550 lithospheric thickening and thinning, factors not considered in the forward models. Because
551 of this, continental inundation models display short- to medium-wavelength mismatches
552 with geologically-reconstructed coastlines, and this is reflected in improved model fits for
553 relatively stable continents (like North America) versus continents that have experienced
554 numerous orogenies and rifts (like Eurasia).

555

556 ***Clusters of dynamic topography evolution***

557 The long-term evolution of continental dynamic topography is primarily driven by their
558 interaction with sinking slabs and large mantle upwellings away from slabs, which represent
559 the large-scale vertical return flow in response to subduction. During the breakup of Pangea
560 and the subsequent dispersal of continents, some continental regions have remained in the
561 vicinity either of a large upwelling (associated with an LLSVP) or of "neutral" mantle away
562 from both large upwellings and subduction zones (Figs 2-4). In these cases continental
563 regions have experienced little change in dynamic topography. In contrast, many other
564 continental regions have moved over subducting slabs after the breakup of Pangea, resulting
565 in these regions being drawn down at least in a particular period during dispersal, affecting

566 both regional as well as global sea level. Some regions are still overlying “slab burial
567 grounds”, while others have moved across subducting slabs and have experienced a change
568 from dynamic subsidence to uplift. There are also regions that have mainly experienced
569 distinct dynamic uplift by either being over a plume at certain time intervals (in those
570 models that include plumes) or by moving towards a large mantle upwelling (sometimes
571 referred to as superplumes) (Figs 2-4). Cluster analysis allows us to segment continental
572 regions into different classes of dynamic topography evolution, but the most appropriate
573 number of classes is not known *a priori*. We explore 3-5 clusters for all continental regions,
574 and then repeat the same analysis for passive margin regions only to assess the differences
575 in which continental interiors and passive margins may have been affected by mantle-driven
576 dynamic topography.

577

578 Different groups of geodynamic models naturally yield different categories of clusters. For
579 instance, models M5-7 differ substantially from all other models, with M5 and M6 moreover
580 including plumes, which are suppressed in models M2, M3 and M7, while model M1 may
581 contain active upwellings assimilated from the S20RTS mantle tomography model (Ritsema
582 et al., 2004). Based on these differences, some models display evolutionary paths of
583 dynamic topography that are not common in others. Considering three clusters
584 demonstrates that this number insufficiently captures the diversity of dynamic topographic
585 evolution of continental regions, with the exception of M1 (Fig. 11), partly reflecting that this
586 model only covers a 90 million year time period. Here three clusters differentiate
587 topographic stability from long-term subsidence in regions overlying slab burial grounds for
588 the entire model period as well as long-term subsidence in more elevated regions initially
589 more distal to subduction zones, but gradually moving over sinking slabs (Fig. 11).

590

591 For all other models, a choice of four clusters (Fig. 12) differentiates topographic stability,
592 long-term pronounced subsidence, initial stability over a dynamic high followed by moderate
593 subsidence and a fourth cluster representing regions proximal to subduction zones, either
594 with initial subsidence followed by uplift (M2-4), or accelerating subsidence through time
595 (M5-7) (Table 3). We find that using five clusters does not add improve the categorisation of
596 continental dynamic topography evolution. In terms of the maximum amplitude of total
597 dynamic topographic change over 140 million years, our favoured model M7 results in up to
598 500 (± 150) m of total dynamic subsidence (TDS) (cluster 1) while the other clusters are
599 limited to total change of the order of 200-300 m (Fig. 12), reflecting that the long-term
600 dynamic topographic change effect in most continental regions is within the range of first-
601 order eustatic sea level fluctuations for this model (Fig. 1). The subsidence clusters in Model
602 M1 result in a maximum of 350 (± 200) m of TDS over 90 million years (Fig. 11), of a similar
603 order of magnitude to model M7, while the subsidence cluster in models M2-4 typically
604 results in TDS of 1000 (± 400) m, which is significantly larger than estimated eustatic sea level
605 fluctuations over this time period (Fig. 1), reflecting that this class of models may
606 overestimate negative dynamic topography.

607

608 The four-cluster categorisation of dynamic topography through time as described for the
609 continents is mirrored by continental margins (Fig. 13), exhibiting similar evolutionary paths.
610 The most commonly observed process is a gradual move of passive margins from dynamic
611 highs towards dynamic lows during Pangea fragmentation, reflecting that many continental
612 passive margins now overlie slabs sinking in the lower mantle. This holds for portions of the
613 eastern margins of North (Spasojevic et al., 2008) and South America (Flament et al., 2014),

614 northern Africa (Barnett - Moore et al., 2017) as well as some segments of Australia's
615 margins, particularly the northeast (DiCaprio et al., 2010), while the margins of eastern
616 Brazil, South Africa as well as southwest Australia are examples where dynamic stability or
617 uplift is predicted in most of our models (Fig. 13). For passive margins, the maximum
618 predicted dynamic topographic change over 140 million years in model M7 is about 350
619 (± 150) m of subsidence, about an order of magnitude smaller than the total tectonic
620 subsidence (Sawyer, 1985) typically caused by rifting, making dynamic signals difficult to
621 detect in tectonic subsidence analyses based on borehole stratigraphy, especially
622 considering typical uncertainties in paleo-water depth (Allen and Allen, 2013). The
623 subsidence clusters in model M1 result in a maximum of 400 (± 200) m of TDS over 90 million
624 years (Fig. 11), while the subsidence cluster in models M2-4 results in TDS of 1000 (± 400) m,
625 similar to estimates for all continents (Fig. 13).

626

627 In a model in which plumes are suppressed, such as M7, passive margins exhibit a more
628 pronounced tendency to be affected by uplift than continental interiors, but the mean
629 amplitude of this effect is of the order of 100 ± 50 m (M7, cluster 4, Fig. 13). In models M2-
630 M4 the magnitude of this effect along passive margins is as large as $\sim 500 \pm 300$ m (model M3),
631 and there are instances, like southeast Australia, where Cenozoic dynamic uplift of about
632 500 m is supported by river profile inversion (Czarnota et al., 2014). Models with plumes
633 result in transient passive margin uplift of about 200 ± 200 m (e.g. M6, cluster 4, Fig. 13), but
634 are mostly characterised by long-term subsidence. A region in which plume-related Late
635 Cenozoic dynamic uplift as modelled in M6 (Fig. 13) is promising in terms of its match to
636 observations is the north eastern Brazilian Borborema Province, where post 50 Ma
637 magmatic plugs have been interpreted as Brazil moving over a hotspot (Mizusaki et al.,

638 2002). Similarly, Late Cenozoic uplift of the south African margin (Roberts and White, 2010)
639 is captured in this model. A detailed comparison of tectonic subsidence derived from
640 exploration wells with our geodynamic models is beyond the scope of this paper. However,
641 we provide interactive geodynamic model access via the GPlates Portal (portal.gplates.org,
642 Müller et al., 2016), where end users can easily extract the predicted dynamic history for any
643 given site, for any of the models presented herein, download the data and evaluate their
644 match with any given tectonic subsidence history.

645

646 **Discussion**

647 Before discussing our results for the dynamic topography for passive margins and their
648 hinterlands through time, we first review model predictions for present-day dynamic
649 topography. Residual oceanic basement depth is perhaps the most useful present-day
650 validation of surface dynamic topography, given the much larger uncertainty of estimating
651 continental residual topography (Colli et al., 2016; Yang and Gurnis, 2016), but it is
652 dependent on a number of assumptions related to the depth-age relationship of “normal”
653 ocean floor (see recent review by Stein and Stein, 2015). Even though oceanic depth-age
654 models are typically constructed by excluding data from hotspot swells (see for instance
655 Crosby and McKenzie, 2009), they are based on the assumption that long-wavelength
656 dynamic topography does not exist or is insignificant, inverting the observations from
657 presumably “normal” ocean floor to derive a global best-fit depth-age relationship.
658 However, just as the amplitude of eustatic sea level fluctuations cannot be gleaned from any
659 single locality (Bond, 1978), it is equally difficult to estimate the anomalous depth of oceanic
660 basement at a given location. What the two problems have in common is the ubiquity of
661 mantle-driven dynamic topography that affects the surface of the Earth at any given site

662 (Figs 2-4). It follows that an inversion for a depth-age curve from sediment-unloaded oceanic
663 basement depths that is expected to solely reflect thermal boundary layer cooling and/or
664 plate-model-related time-dependent small-scale convection beneath oceanic plates will
665 inherit biases from any other process that is ignored. A review of published numerical
666 dynamic topography models (Flament et al., 2013) illustrates that large parts of the ocean
667 basins are particularly affected by positive dynamic topographic anomalies owing to large-
668 scale upwellings – this can also be seen in the models used here (Figs 2-4). While considering
669 that the amplitude of these topographic features is uncertain, this nevertheless suggests
670 that this bias may lead to excess plate model flattening of old ocean floor, affecting the
671 comparison of residual topography and numerically computed dynamic topography, if
672 computed residual topography is based on a reference oceanic-depth age model that has
673 inherited dynamic topography signals. Our spectral analysis of these models demonstrates
674 that if the uncertainty in oceanic depth-age models is considered in this context, then the
675 power spectra of models M5-M7 are too red only at degree 1 (Fig. 5), representing a
676 significant improvement over previous models. At spherical harmonic degrees 2 and 3, the
677 r.m.s. dynamic topography amplitude of model M7 is 740 m, and 570 m, respectively, within
678 the range of residual topography end-members at 850 m (thermal boundary layer (TBL)
679 subsidence only) and 530 m (plate model (PM) subsidence) at degree 2 and 770 m (TBL) and
680 510 m (PM) at degree 3. This suggests that at degree 2, Crosby and McKenzie's (2009) plate
681 subsidence model may be contaminated by about 200 m of deep mantle convection-driven
682 dynamic topography. However, it needs to be kept in mind that our preferred model M7
683 still overestimates the amplitude of dynamic topography at spherical degree 1 as estimated
684 by Hoggard et al. (2016) by at least 700 m, but the magnitude of residual topography at
685 these long wavelengths is still debated (Yang and Gurnis, 2016). Discrepancies between

686 residual topography and dynamically computed topography at long wavelengths possibly
687 reflect the computation of dynamic topography from sources below 350 km depth, well
688 below the depth of continental cratonic keels and below the depth to which slabs are
689 assimilated in the models, as well as the inherent limitations in spherical harmonic
690 expansion of a limited global point data set (Yang and Gurnis, 2016).

691

692 We find that our geodynamic forward model M7 provides the best overall fit to
693 paleogeography-derived continental inundation (Fig. 14). All models generally match
694 geological observations better for the last 60 million years than for earlier times. The mean
695 Cenozoic differences between geodynamically-modelled versus geologically-reconstructed
696 land fractions are within $\pm 5\%$ with the exception of M6, and the land fraction differences
697 based on the two alternative paleogeography reconstructions are overall similar (Fig. 14).
698 This reflects a consensus in the reconstructions of Cenozoic paleogeography and that our
699 models match them well overall. African land fractions are nearly always underestimated
700 with the exception of model hybrid model M1. Forward models that include plumes (M5 and
701 M6) are less useful for modelling continental dynamic topography, as in these models
702 plumes evolve fully dynamically such that neither the time nor the location of their initial
703 arrival at the surface can be well tuned to match the observed occurrence of plume-related
704 uplift. This emphasises the future prospects of models with sequential data assimilation.
705 Model M7, in which plumes are suppressed and which is characterised by relatively
706 moderate dynamic topography amplitudes as compared to models M1-4, overall fits the
707 land fraction for all continents as compared with Golonka's (2007) paleoshorelines, with the
708 exception of the Late Cretaceous flooding of Africa (due to lacking plumes) (Fig. 14e). Using

709 Smith et al.'s (1994) paleoshorelines worsens the fit for Australia and Eurasia, with M7
710 underestimating continental flooding (Figs 14f, j).
711
712 It is important to note that agreement between modelled and inferred land fractions is
713 possible even if the flooded regions only partially coincide spatially. Therefore it is essential
714 to evaluate the model agreement in land fractions jointly with spatial overlaps (Fig. 14).
715 Model M1 performs slightly better than other models for land fraction overlaps, with the
716 best matches in Australia, Eurasia, Africa and North America, while South American
717 paleogeography is less well matched, possibly because Andean flat slabs were not
718 incorporated in Model M1. On the other hand, even though South American flat slabs are
719 included in M2-M7, only M2 and M7 result in reasonable fits to geologically mapped South
720 American inland seas for Golonka's (2009) Cenozoic paleogeography (Fig. 14c), while using
721 Smith et al.'s (1994) paleo-coastlines improves the South American match for most models
722 (Fig. 14h), suggesting that the latter paleogeography might be better constrained for the
723 Cenozoic of South America than the former. By the same token the match in Australia is
724 slightly improved for nearly all models using Smith et al.'s (1994) paleogeography,
725 underlining its greater similarity to the detailed Australian paleogeography by Langford et al.
726 (1995) in the Cenozoic as compared with Golonka's (2009). In contrast, using Golonka's
727 (2009) paleogeography considerably improves the model match for North America as
728 compared with Smith et al.'s (1994), highlighting that currently there is no single preferred
729 global model for paleo-coastlines on all continents.
730
731 Model M7 maximises the combined inundation overlap for Eurasia, Africa, South and North
732 America, but misfits the Cenozoic flooding of Australia (Fig. 14c, d). It has been shown before

733 that the inclusion of a large mantle upwelling straddling East Antarctica is important for
734 modelling Australia's dynamic topography in the Cenozoic, as the continent progressively
735 moved away from this upwelling and towards the dynamic low associated with Southeast
736 Asian subduction zones (DiCaprio et al., 2011). In the models analysed here, the mantle
737 structure associated with this upwelling is only considered in Model M1, and this explains
738 why M1 outperforms all other models in terms of replicating Australia's Cenozoic inundation
739 patterns (Fig. 14c, d). Overall, modelled versus geologically reconstructed land fractions
740 match within 10% for most models, and the spatial overlaps of inundated regions are mostly
741 between 85% and 100% for the Cenozoic, dropping to about 75-100% in the Cretaceous.

742

743 In terms of dynamic topography of passive margins through time, the favoured model M7
744 results in up to 500 (± 150) m of tectonic subsidence of continental interiors while along
745 passive margins the maximum predicted dynamic topographic change over 140 million years
746 is about 350 (± 150) m of subsidence, substantially smaller than the total tectonic subsidence
747 caused by rifting. Typical rates of dynamic topographic change range from ± 10 m/Myr
748 (Model M7, Figs 15, 16). In an extended Boussinesq model in which plumes are suppressed,
749 such as M7, passive margins exhibit a more pronounced tendency to be affected by uplift
750 than continental interiors, but the mean amplitude of this effect is only of the order of 100
751 ± 50 m, because dynamic topography amplitudes in all extended Boussinesq models are
752 smaller overall compared with other models. Other models, such as M3, which shares the
753 same plate model with M7, also perform reasonably well in terms of modelled land fractions
754 and inundation overlaps, but we favour M7 because of its improved match to residual
755 oceanic topography at long wavelengths. Models with plumes can result in more
756 pronounced passive margin uplift of about 200 ± 200 m. This effect is more pronounced along

757 continental margins than interiors because some passive margins have either moved over
758 the periphery of a large mantle upwelling, like eastern Australia (Müller et al., 2016) or have
759 been affected by a mantle plume through time, with the northeast coast of Brazil being a
760 potential example (Mizusaki et al., 2002). Australia and South American are also the two
761 continents exhibiting the largest long-term gradients in mean rates of change in dynamic
762 topography (Fig. 16), with Australian experiencing growing rates of subsidence throughout
763 the Cenozoic, reflecting its northeastward migration towards the Melanesian slab burial
764 ground, while South America experiences a gradual increase in uplift rates over the last 40
765 million years, reflecting an intensification of the large-scale mantle upwelling centered on
766 Africa, straddling the east coast of South America, paired with a rebound of the west coast
767 of South America from being previously drawn down by the sinking Farallon slab (Fig 4).

768

769 **Conclusions**

770 We have carried out a global analysis of mantle convection-driven dynamic surface
771 topography for the last 140 Ma, using seven geodynamic models combined with alternative
772 eustatic sea level curves, and evaluated predicted continental flooding patterns against two
773 alternative sets of geologically-derived paleo-coastlines. After evaluating the power spectra
774 of the dynamic topography models, the match between model predictions and published
775 paleo-coastlines is established based on computing modelled land fractions as well as
776 inundation overlaps through time. We find that forward geodynamic model M7, which is
777 based on an extended Boussinesq approximation and a mantle viscosity profile similar to
778 that of Steinberger and Calderwood (2006), provides the best overall fit to geologically-
779 derived continental inundation. Model M3 also performs fairly well, reflecting that M3 and
780 M7 are based on the same recent plate model. However, for the last 60 million years, model

781 M1 fits best, reflecting a backward-forward modeling approach with an assimilated
782 tomographically-imaged mantle structure into forward models that works well for the recent
783 geological past. For the Cenozoic, model M1 also stands out by performing better than all
784 other models for matching Africa's flooding history, both in terms of land fractions and
785 inundation overlaps. Our model evaluation reveals that our overall best-fit model M7 fits
786 Golonka's (2007, 2009) paleogeography somewhat better than Smith et al.'s (1994) in the
787 Cretaceous, whereas the alternative paleo-coastline reconstructions are roughly equivalent
788 in the Cenozoic, with the exception for Australia and South America, where modelled
789 inundation is better matched by Smith et al.'s (1994) paleogeography.

790

791 We categorise the evolution of modelled dynamic topography in both continental interiors
792 and along passive margins using cluster analysis to investigate how clusters of similar
793 dynamic topography time series are distributed spatially. A subdivision of four clusters is
794 found to best reveal end-members of dynamic topography evolution, differentiating
795 topographic stability, long-term pronounced subsidence and initial stability over a dynamic
796 high followed by moderate subsidence. The fourth cluster represents regions that are always
797 proximal to subduction zones, and exhibits evolutionary paths including initial subsidence
798 followed by uplift, or accelerating subsidence through time. This four-cluster categorisation
799 of continental dynamic topography through time is mirrored by passive margins. The most
800 commonly observed process is a gradual movement of passive margins from dynamic highs
801 towards dynamic lows during the fragmentation of Pangea, reflecting that many continental
802 passive margins now overlie slabs sinking in the lower mantle. This may explain why passive
803 margin highlands are relatively rare.

804

805 The overall match between predicted dynamic topography, modulated by eustasy, to
806 geologically mapped paleo-coastlines through time in terms of land fractions and inundation
807 overlaps suggests that mantle-driven dynamic topography is a critical component of relative
808 sea level change, and indeed the main basis for understanding the patterns of large-scale
809 continental inundation through time. By ground-truthing models using the flooding history
810 of continental interiors, we have established a robust method for evaluating dynamic
811 topographic change along passive continental margins, where dynamic topography signals
812 are more difficult to detect in the geological record.

813

814 Geodynamic forward models that are well calibrated for relatively recent geological periods
815 in terms of their predicted dynamic topography open up the opportunity to model dynamic
816 surface topography in the Early Mesozoic and Paleozoic, as plate models with topologically
817 closing plate boundaries ranging from the Devonian Period to the present (e.g. Matthews et
818 al., (2016) become available. This would lead to an improved understanding of large-scale
819 continental uplift and subsidence and the interplay between shifting coastlines, sediment
820 sources and sinks through time. Coupling this approach with surface process models (e.g.
821 Salles and Hardiman, 2016) would provide a more quantitative understanding of the origin
822 and pathways of sediments that have filled sedimentary basins through time, and would
823 provide genetic insights into the stratigraphy of individual basins and margins.

824

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831 manuscript. An interactive globe on the GPlates Portal (portal.gplates.org) will allow users to
832 explore the detailed predictions of individual models presented here for any given
833 continental and passive margin locations.

834

835 **Figure Captions**

836 Fig. 1. Alternative global sea level curves used in combination with modelled dynamic
837 topography. The red- green- and blue-dotted curves show the sea level curves derived from
838 the oceanic age-area distributions through time from Müller et al. (2016), Seton et al.
839 (2012), and (Spasojevic and Gurnis, 2012), modified from Müller et al. (2008), respectively.
840 The large difference in the amplitude of these curves reflects the uncertainty in
841 reconstructing now subducted ocean floor, particularly in the Pacific and Tethys, including
842 now destroyed Cretaceous Tethyan back-arc basins. The large drop in sea level back in time
843 from 120 to 140 Ma visible in both curves is an artefact reflecting an underestimate in the
844 length of the global mid-ocean ridge system at these times, especially in terms of back-arc
845 basins, as well as our insufficient knowledge of oceanic plateaus before 120 Ma. The blue
846 and black curves show the preferred sea level curves in Spasojevic and Gurnis (2012) and
847 Haq et al. (1987), respectively. The red and green areas show the range of global sea level
848 estimates obtained from combining ocean basin volume-derived sea level with oceanic

849 dynamic topography (DYN) from global convection models with the exception of model M1
850 which is based on a different modelling approach combining the effects of mean oceanic and
851 continental dynamic topography on eustasy. Including dynamic topography effects enhances
852 modelled sea level highs, because the majority of the ocean basins are dominated by large-
853 scale mantle upwellings.

854 Fig. 2. Modelled dynamic topography in model M1 at 10 Myr intervals, with reconstructed
855 continents (Gurnis et al., 2012) overlain.

856 Fig. 3. Modelled dynamic topography in model M3 at 10 Myr intervals, with reconstructed
857 continents (Müller et al., 2016) overlain.

858 Fig. 4. Modelled dynamic topography in model M7 at 10 Myr intervals, with reconstructed
859 continents (Müller et al., 2016) overlain.

860 Fig. 5. Power spectra of present-day dynamic topography for models M1-7. The spectra of
861 residual depth anomalies calculated with the depth-age plate model of Crosby and McKenzie
862 (2009) (light gray) and a half-space cooling model (dark gray) are taken from Hoggard et al.
863 (2016). The scale of the power is given km^2 ; taking the square root of the power at a given
864 degree will provide amplitude, as discussed in the text.

865 Fig. 6. Mean dynamic topography through time for each continental region considered here.

866 Fig. 7. Mean oceanic dynamic topography for models presented in this study through time.
867 The long-term trend in the evolution of oceanic dynamic topography in models in group A
868 (top) shows a sharp contrast with that from models in group B (bottom). This is a
869 consequence of cold subducting slabs playing a more significant role in models M1-M4
870 (group A). In these models slabs do not heat up as they descend into the mantle and thus
871 trigger a larger return flow as compared with models in which descending slabs do heat up.
872 In models M1-M4 the evolution of dynamic topography therefore directly reflects the
873 evolution of the age distribution, and thickness, of subducting oceanic lithosphere through

874 time, which directly controls the buoyancy of slabs. In contrast, in models M5-M7 cold
875 subducting slabs undergo viscous heating as they descend into the mantle, and more
876 importantly, models in this group include plume upwellings (although suppressed in model
877 M7) that could reflect the long-term rise in cumulative large igneous province (LIP) volumes
878 since 150 Ma (Yale and Carpenter, 1998), even though the process of melting and LIP
879 generation is not included in our models. Hence oceanic dynamic topography in these
880 models shows a steady increase towards present-day.

881 Fig. 8. (a (i)) The solid curve (DYN_{M1} -SG12) shows the sea level curve used in (Spasojevic and
882 Gurnis, 2012) combining ocean-basin volume effects with the contribution of oceanic
883 dynamic topography, while the dotted curve (SG12) shows the curve based on the oceanic
884 age-area distribution from Müller et al. (2008), with the shaded region illustrating the
885 contribution of oceanic dynamic topography. We show the effects of sea level curves with
886 and without correction for mean oceanic dynamic topography in all our models, considering
887 that sea level curves without this contribution represent an underestimate, but curves
888 including oceanic dynamic topography may overestimate global sea level amplitudes, given
889 that the geodynamic models overestimate dynamic topography amplitudes at long
890 wavelengths (Fig. 5).

891 (a (ii)) The distribution of oceans and continents predicted in model M1, using the SG12 sea
892 level curve based on (Müller et al., 2008), is shown in the first column at labelled ages (see
893 methods for more details). The second column shows equivalent predictions based on the
894 DYN_{M1} -SG12 sea level curve. The distribution of oceans and continents in the
895 paleogeography models of Smith et al. (1994) and Golonka (Golonka, 2007, 2009) are shown
896 in the third and fourth columns, respectively.

897 (b) Evolution of the fraction of land, $\tau(t)$, over the last 90 Ma as predicted in model M1,
898 based on sea level curves SG12 (blue) and DYN_{M1} -SG12 (black) are shown on the first

899 column for each labelled continent. Evolution of the fraction of land, $\tau(t)$, computed for
900 paleogeography grids from Smith et al. (1994) and Golonka (2007, 2009) are also shown in
901 red and green, respectively, for labelled ages in the first column. The yellow and black curves
902 in the second column show the evolution of spatial overlap, $\mu(t)$, (see methods) between
903 predictions of inundation patterns within each continent in model M1, based on sea level
904 curve SG12, and those from paleogeography grids (Golonka, 2007, 2009; Smith et al., 1994).
905 The cyan and magenta curves show the equivalent result, but for sea level curve DYN_{M1}-
906 SG12.

907 Fig. 9. (a (i)) Same as in Fig. 6a(i) but for model M3 and is based on the sea level curve
908 indicated – see Fig 1. for sea level keys. (a (ii)) Same as in Fig. 6a(ii), b (Golonka, 2007, 2009)
909 but for model M3 and based on sea level curves as indicated. (b) Same as in Fig. 6b, but for
910 model M3 and based on sea level curves as indicated, with the time ranging from 140 Ma to
911 the present.

912 Fig. 10. (a (i)) Same as in Fig. 6a(i) but for model M7 and is based on the sea level curve
913 indicated – see Fig 1. for sea level keys. (a (ii)) Same as in Fig. 6a(ii), b (Golonka, 2007, 2009)
914 but for model M7 and based on sea level curves as indicated. (b) Same as in Fig. 6b, but for
915 model M7.

916 Fig. 11. (a) Cluster analyses of modelled continental dynamic topography in model M1 over
917 the last 90 Ma (see methods and Table 3). (b) Evolution of dynamic topography within each
918 cluster, with $\pm 1\sigma$ envelopes. (c) Same as in (a), but with analyses restricted to passive
919 margin regions only. (d) Same as in (b), but for clusters shown in (c).

920 Fig. 12. (a) Cluster analyses of modelled continental dynamic topography in models M2-M7.
921 (b) Evolution of dynamic topography within each cluster, with $\pm 1\sigma$ envelopes.

922 Fig. 13. (a) Cluster analyses of modelled continental dynamic topography, restricted to
923 passive margin regions only, in models M2-M7. (b) Evolution of dynamic topography within
924 each cluster, with $\pm 1\sigma$ envelopes.

925 Fig. 14. Predictions shown here from models M1-M7 are based on their respective sea level
926 curves, which include contributions from oceanic dynamic topography – see Figs 6-8 and
927 Supp. Figs 8-11 for more details. (a) Deviations of mean fractions of land predicted by
928 models M1-7 from that implied in paleogeography grids in Golonka (2007, 2009) for each
929 continent, between 0 – 60 Ma, are shown in columns 1-5. Blue colours indicate
930 overestimates of continental flooding, whereas red colours indicate excess land areas in our
931 models. Hatched patterns for model M1 indicate absent model outputs in the Early
932 Cretaceous, as this model only spans the time period from 0-90 Ma. The last column shows
933 r.m.s. deviations for all continents over the same period as an indicator of the overall global
934 model-data match for a given time period. (b) Same as in (a), but using paleogeography grids
935 in Smith et al. (2007, 2009). (c) Mean spatial overlaps between predictions of inundation
936 patterns within each continent in models M1-7 and those implied in paleogeography grids in
937 Golonka (1994), between 0 – 60 Ma are shown in columns 1-5. Warm colours indicate larger
938 spatial overlap than cool colours. The last column shows the mean spatial overlap for all
939 continents over the same period. (d) Same as in (c), but using paleogeography grids in Smith
940 et al. (1994). (e-h) Same as in (a-d), but for the time interval between 60 – 100 Ma. (i-l) Same
941 as in (a-d), but for the time interval between 100 – 140 Ma.

942 Fig. 15. Rates of change of dynamic topography from our preferred model M7 in 10 million
943 year intervals from 140 Ma to the present. Blue colours indicate subsidence while red
944 colours indicate uplift.

945 Fig. 16. Mean rates of change of dynamic topography and standard deviations for individual
946 continental regions considered here.

Parameter	Symbol	Value (Model M1)	Value (Models M2-4)	Value (Models M5-7)	Units
Rayleigh Number	Ra	7.5×10^7	7.8×10^7	5×10^8	
Thermal expansion coefficient	α_0	3×10^{-5}	3×10^{-5}	1.42×10^{-5}	K^{-1}
Density	ρ_0	3340	4000	3930	$kg\ m^{-3}$
Gravity acceleration	g_0	9.81	9.81	10	$m\ s^{-2}$
Temperature change	ΔT	2800	2825	3500	K
Mantle thickness	h_M	2867	2867	2867	km
Thermal diffusivity	κ_0	1×10^{-6}	1×10^{-6}	1×10^{-6}	$m^2\ s^{-1}$
Viscosity	η_0	1×10^{21}	1×10^{21}	1×10^{21}	Pa s
Activation energy	E_η	348, upper mantle 348, lower mantle	100, upper mantle 33, lower mantle	233, upper mantle 320, lower mantle	$kJ\ mol^{-1}$
Activation Volume ^a	V_η	N/A	N/A	1.5×10^{-6} , upper mantle 6.7×10^{-6} , lower mantle	$m^3\ mol^{-1}$
Temperature offset	T_η	1400	452	560	K
Dissipation Number ^b	Di	N/A	N/A	0.8	
Background mantle temperature ^c	T_b	1400	1685	Depth-dependent	K
Radius of the Earth	R_0	6371	6371	6371	km

948

949 **Table 1:** Parameters common to all model cases. Subscript "0" denotes reference values.

950 Common parameter values between models in groups A and B are only shown for group A.

951 ^a *Viscosity parameterization in group A models do not require the activation volume*
952 *parameter.*

953 ^b *Only group B models employ the extended Boussinesq approximation and thus require a*
954 *dissipation number.*

955 ^c *Background mantle temperature varies with depth in group B models.*

956

957

Name/Acronym	Plate Model	Viscosity profile	Additional notes
Model 1	Gurnis et al. (2012)	100,0.1,1,60	Combined backward/forward convection model, Spasojevic and Gurnis (2012) Model M2
Model 2	Seton et al. (2012)	1,0.1,1,100	Müller et al. (2016) Model 3
Model 3	Müller et al. (2016)	1,0.1,1,100	Müller et al. (2016) Model 2
Model 4	Van Der Meer et al. (2010)	1,0.1,1,100	(Flament et al., 2017) Model "Case 24"
Model 5	Seton et al. (2012)	0.05, 0.0001, 0.0025, 0.07	Model with plumes, Hassan et al. (2015) Model M3
Model 6	Müller et al. (2016)	0.05, 0.0001, 0.0025, 0.07	With plumes, (Barnett - Moore et al., 2017), Model "Case C1"
Model7	Müller et al. (2016)	0.05, 0.0001, 0.0025, 0.07	As Model 6, but with plumes suppressed

959

960 **Table 2:** Tectonic boundary conditions and depth-dependence of viscosity of the models
961 referred to in this study. For models M2-7, the viscosity profile is given by applying a factor to
962 the reference model viscosity (1×10^{21} Pa s) above 160 km depth (lithosphere), between 160
963 and 310 km depth (asthenosphere), between 310 and 670 km depth (upper mantle) and
964 below 670 km depth (lower mantle) in order to obtain depth-dependent viscosity in addition
965 to temperature-dependent viscosity. For models M2-M7 these values represent the pre-
966 exponential parameter, A , at four depths through the mantle (2007, 2009) for details.

967

Model type	Cluster 1	Cluster 2	Cluster 3	Cluster 4
Model M1 (hybrid backward advection – forward model)	Long-term stability over dynamic high	Slow long-term subsidence over pronounced dynamic low	Slow long-term subsidence over low-amplitude dynamic low	
Geodynamic forward model group A (models M2-4)	Pronounced long-term subsidence	Stability with slow long-term uplift	Initial Early Cretaceous stability followed by gradually accelerating subsidence	Subsidence for most of the Cretaceous (140-80 Ma) followed by uplift
Geodynamic forward model group B (models M5-7)	Pronounced long-term subsidence	Slow long-term subsidence	Long-term stability over dynamic high with variable combinations of initial or late uplift/subsidence	Long-term slow subsidence over dynamic low

968

969 **Table 3.** Cluster categories for alternative groups of geodynamic models using 4 clusters

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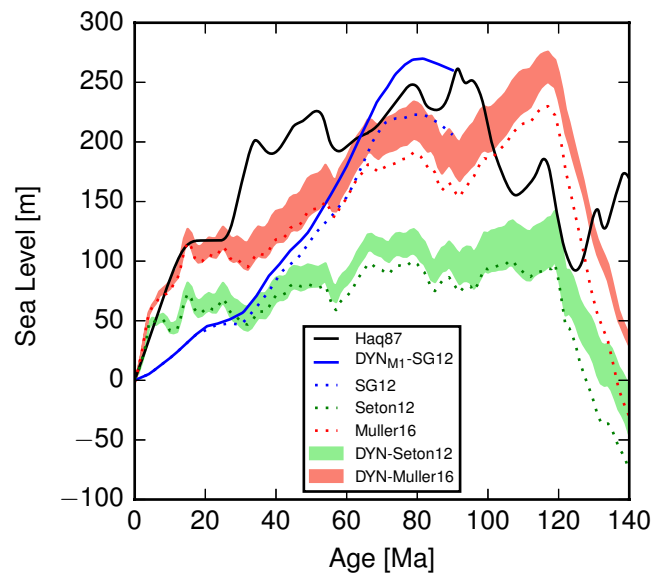
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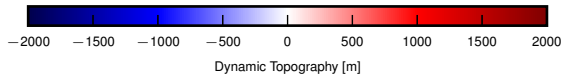
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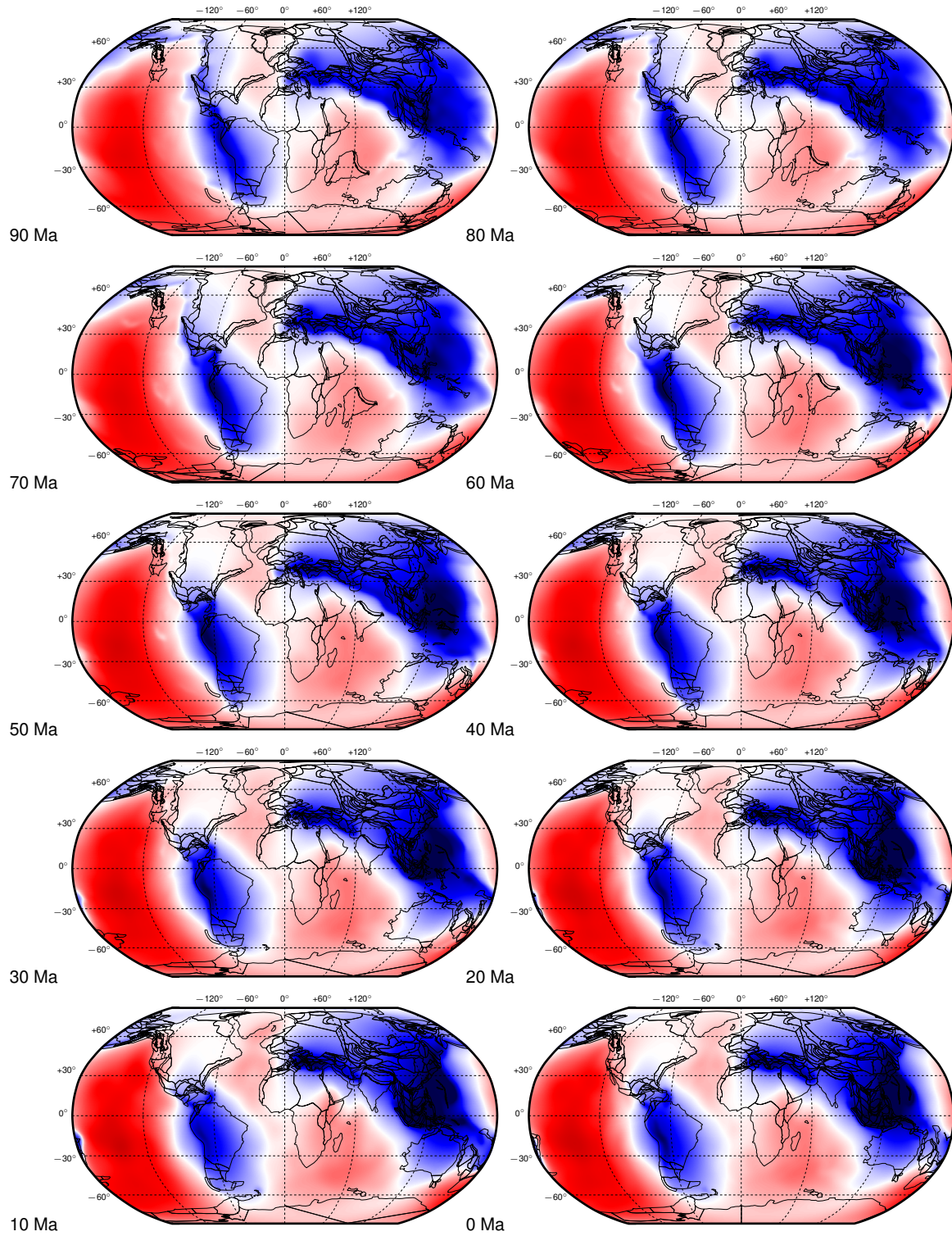
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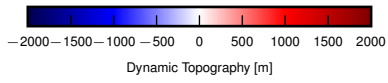
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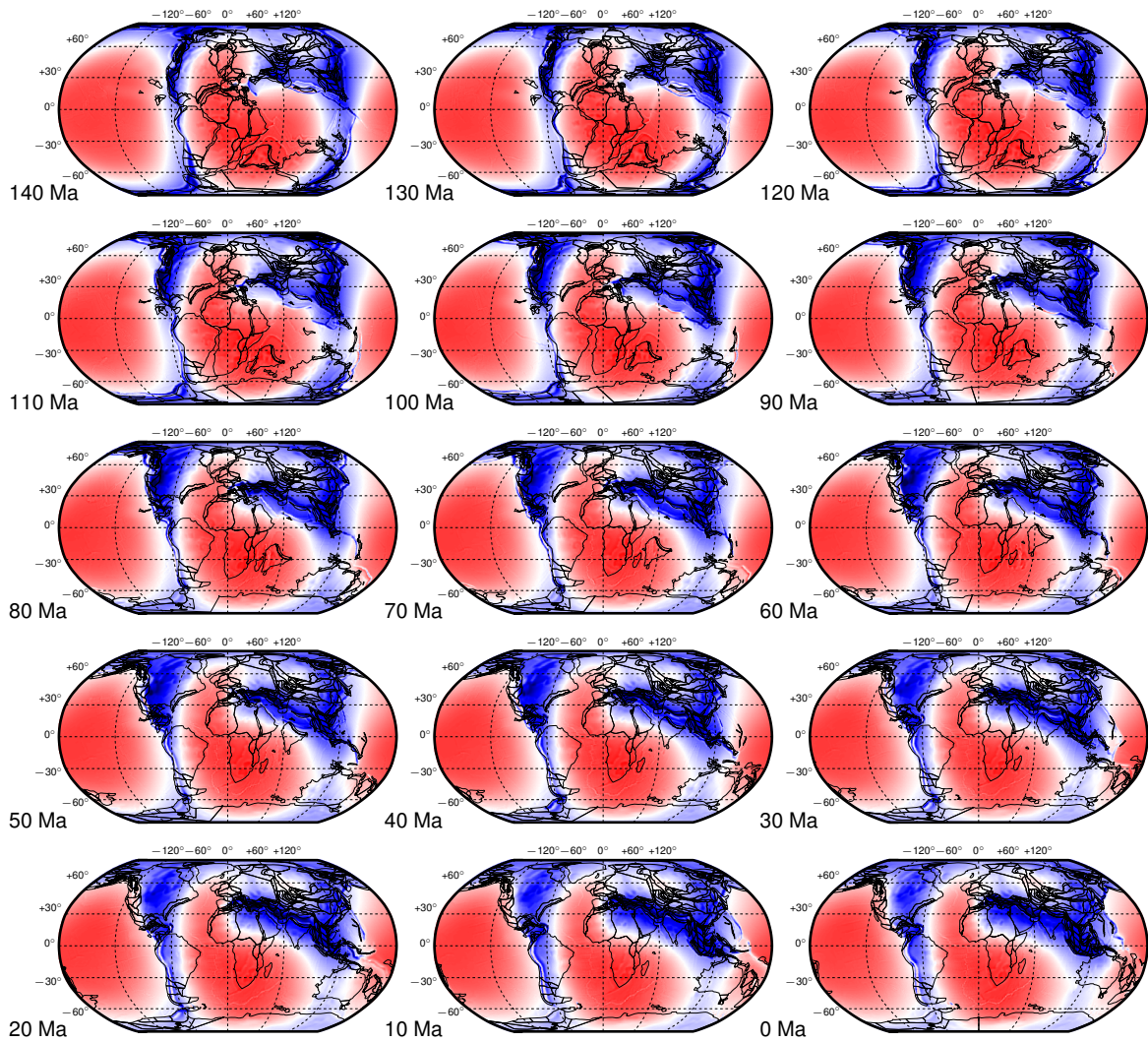


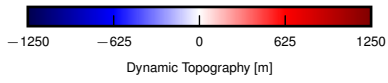
Model: M1



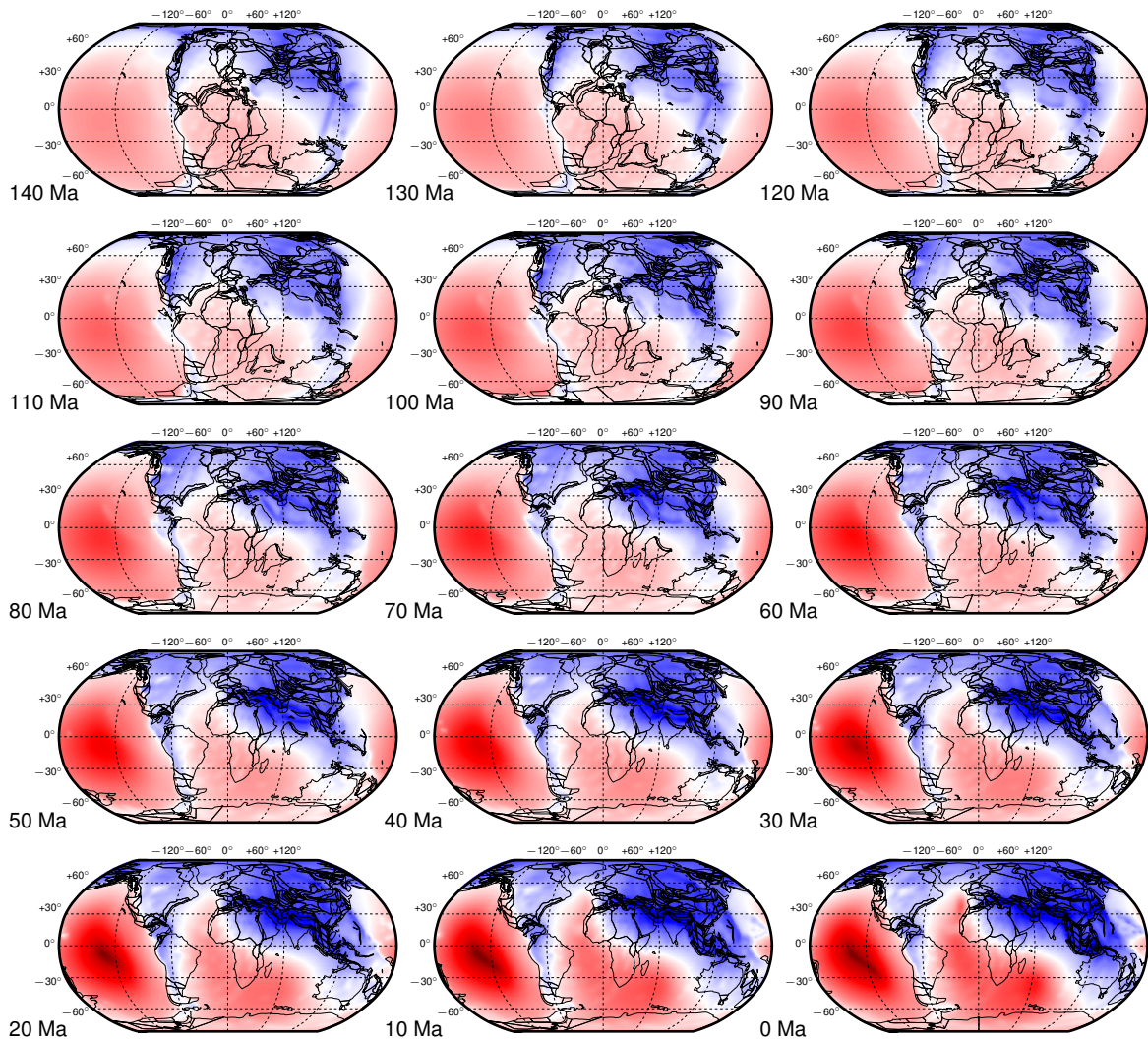


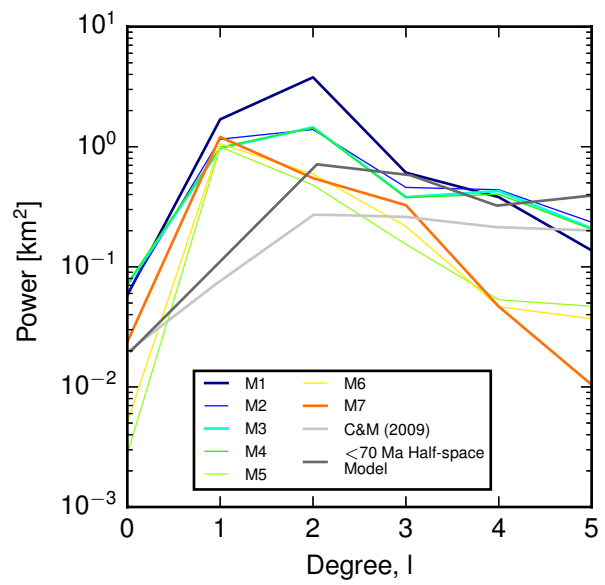
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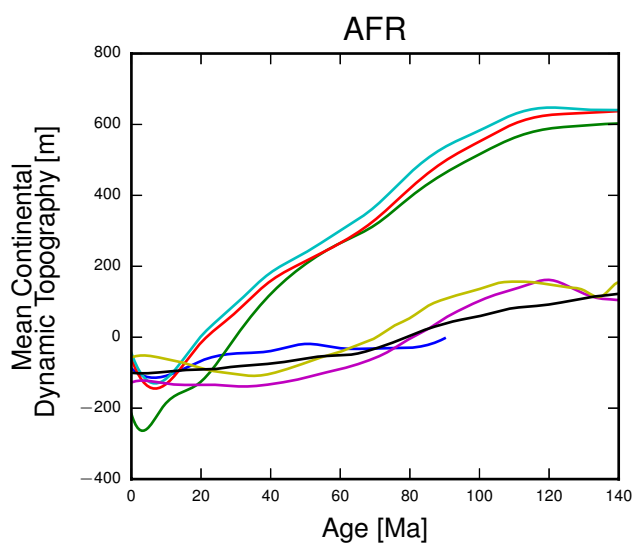
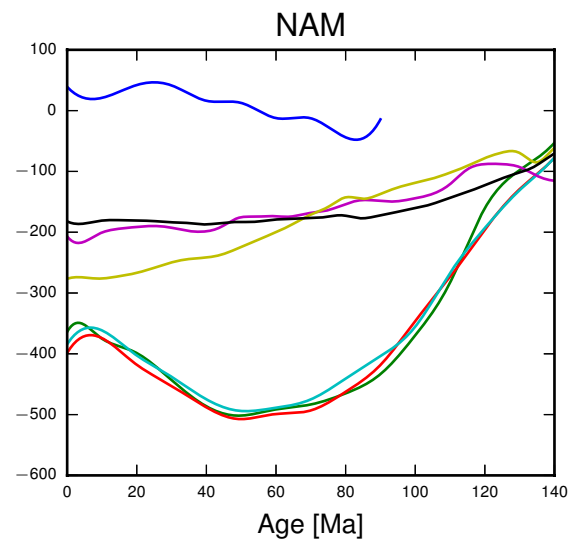
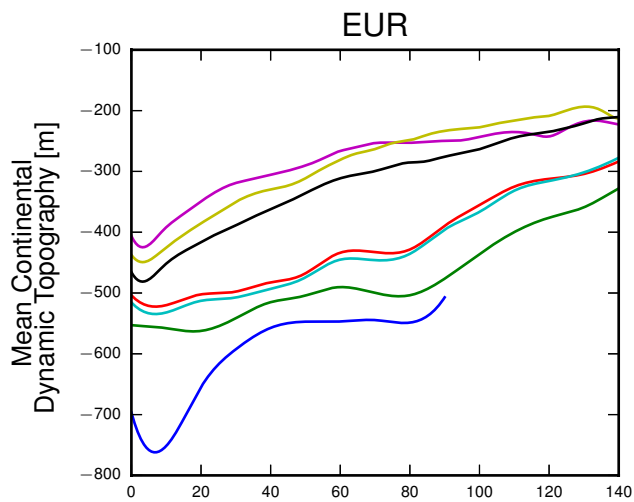
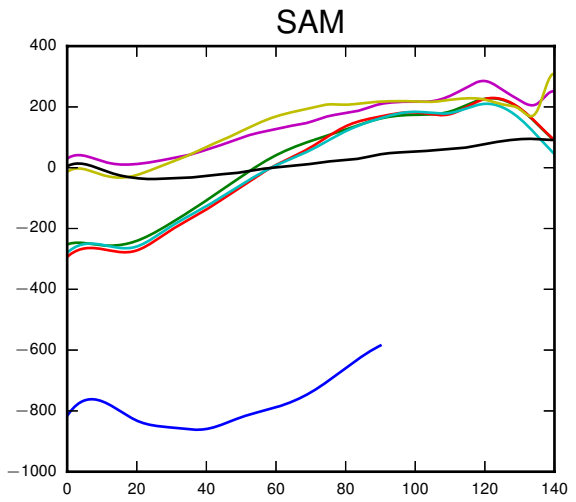
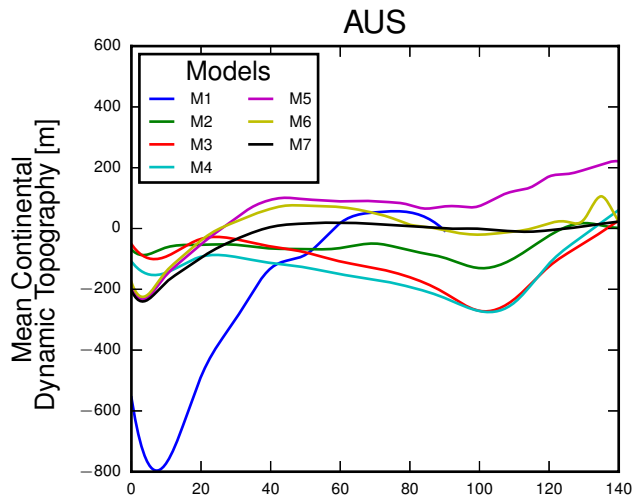


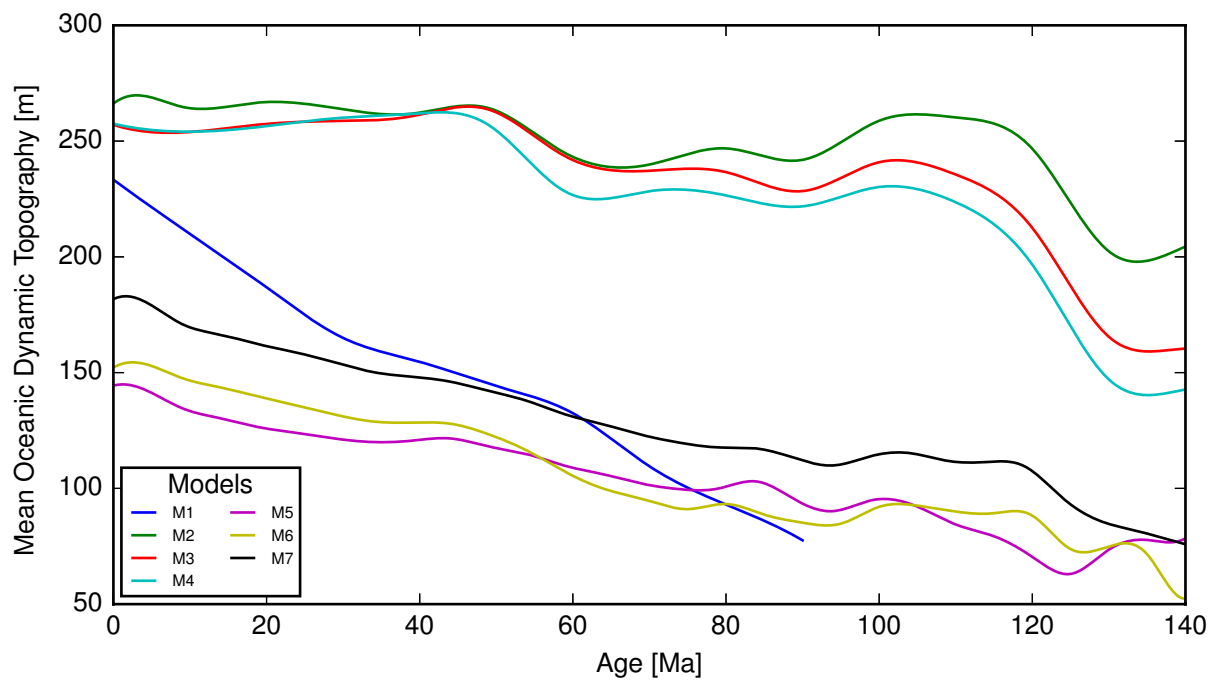


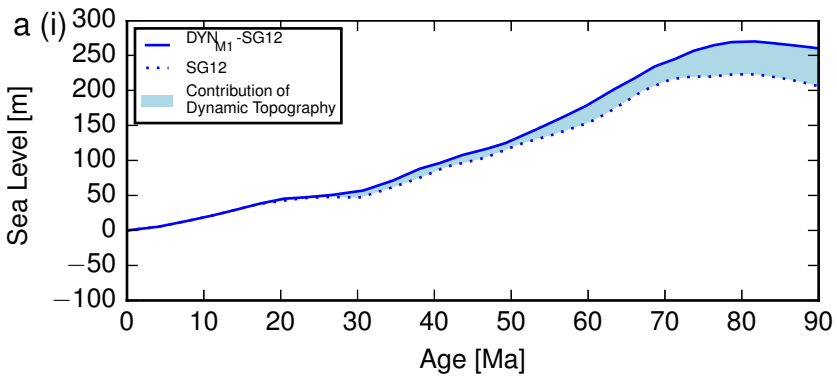
Model: M7





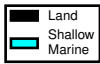






(ii)

$h_c = -325$ m

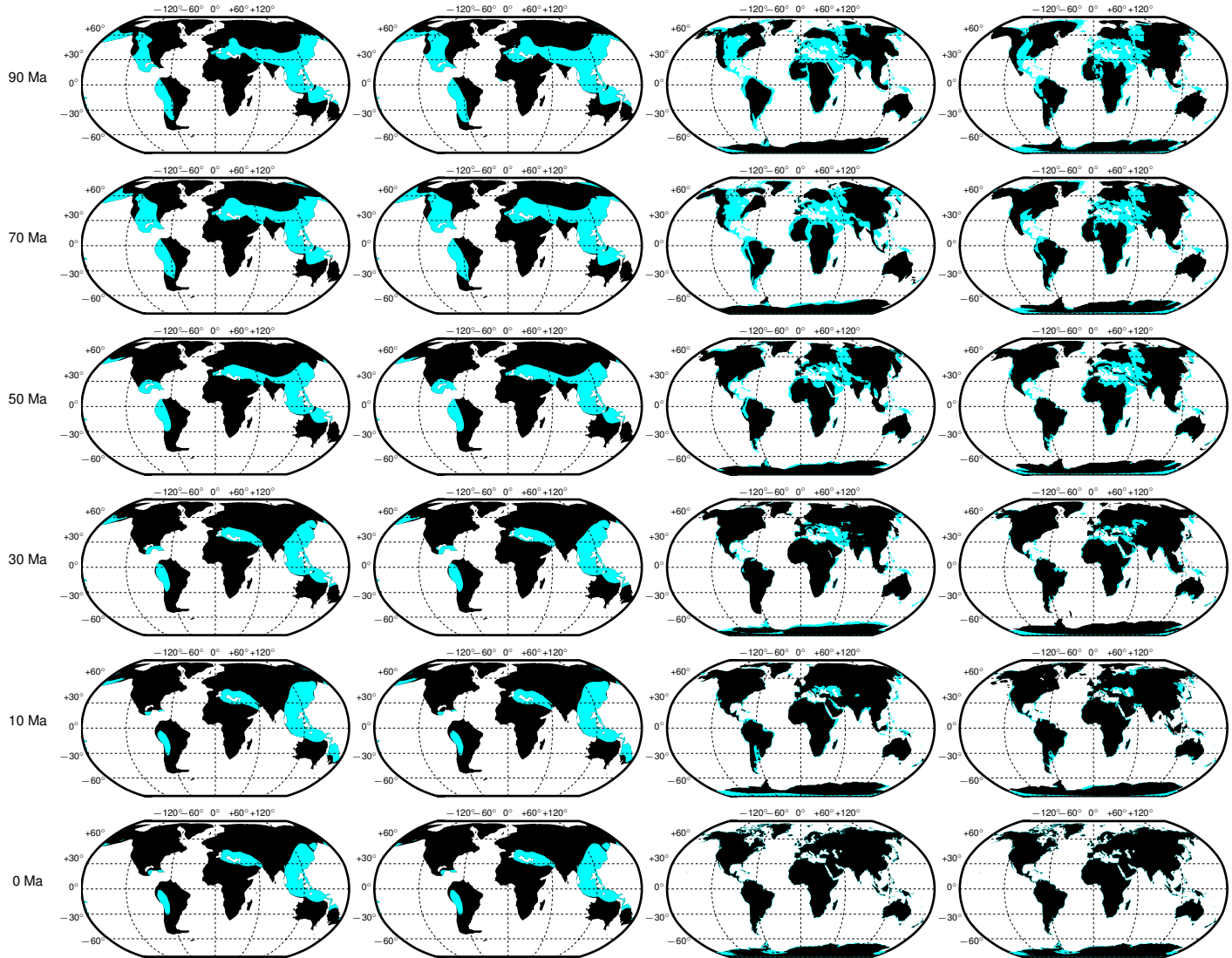


Prediction (SG12)

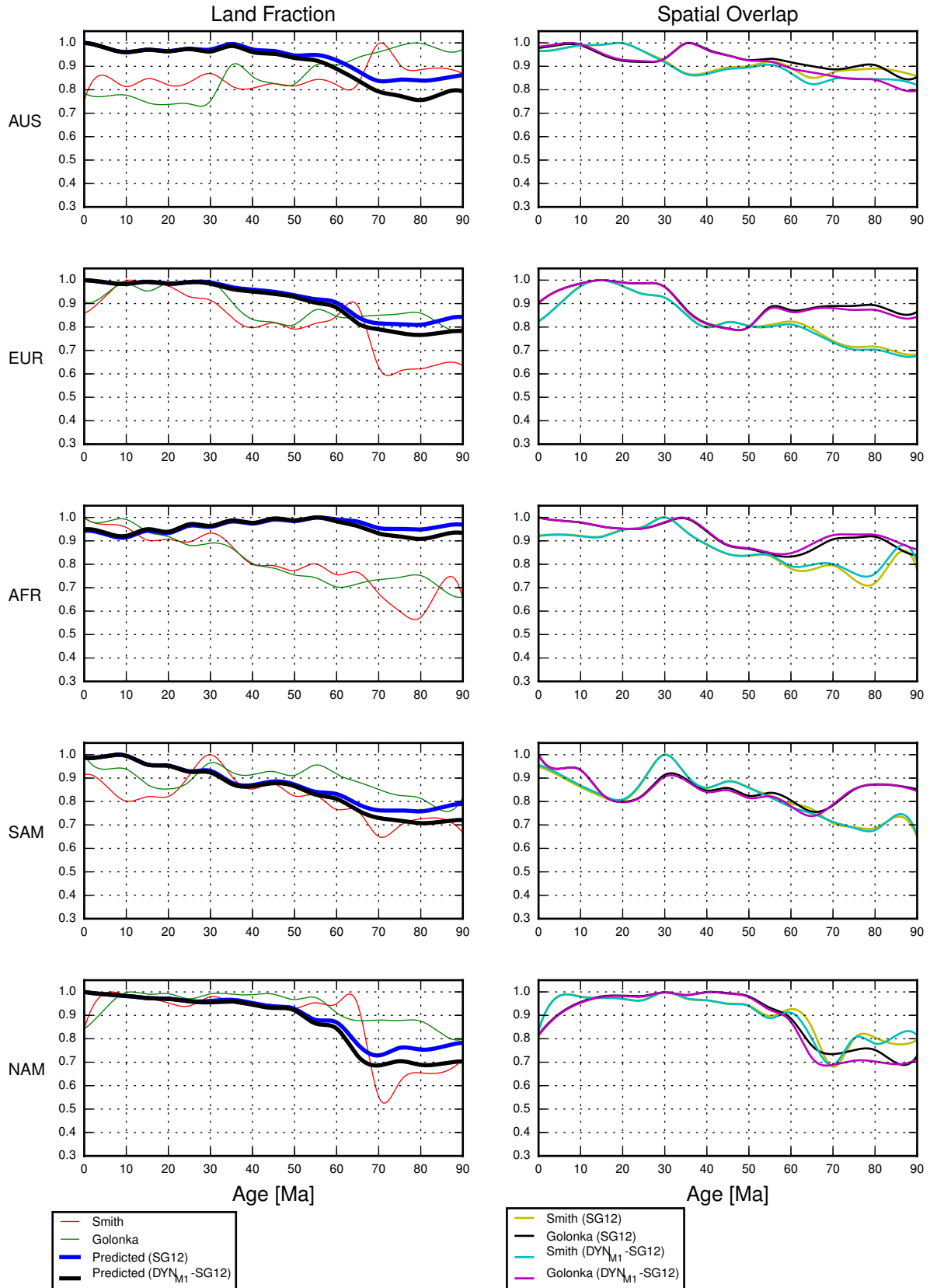
Prediction (DYN_{M1}-SG12)

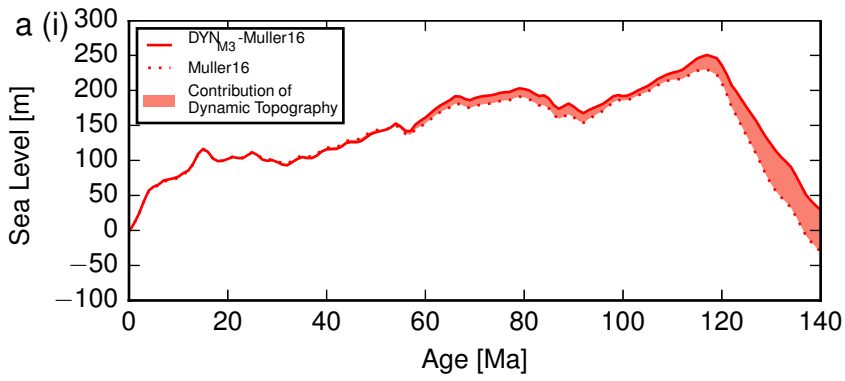
Smith

Golonka

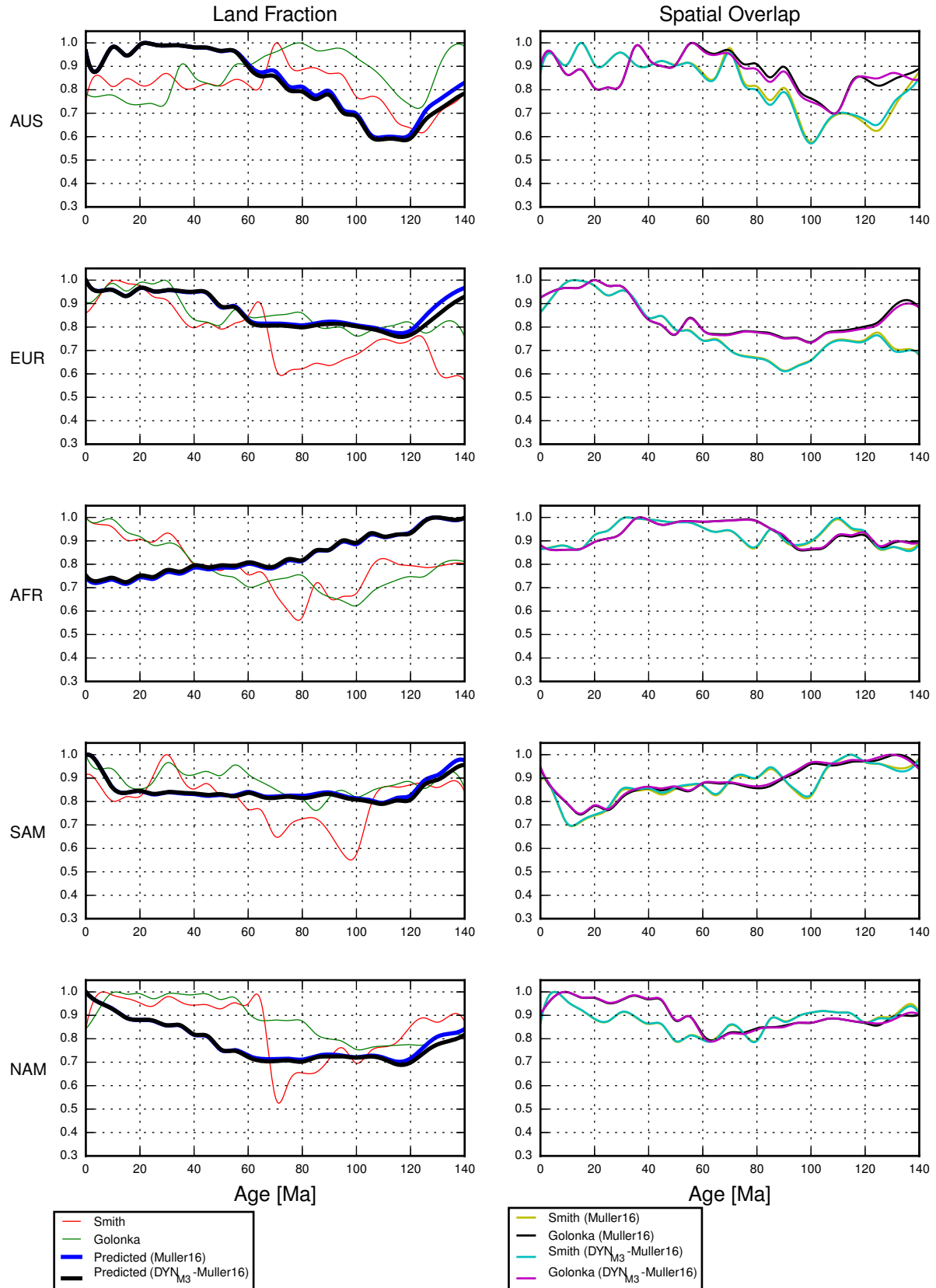


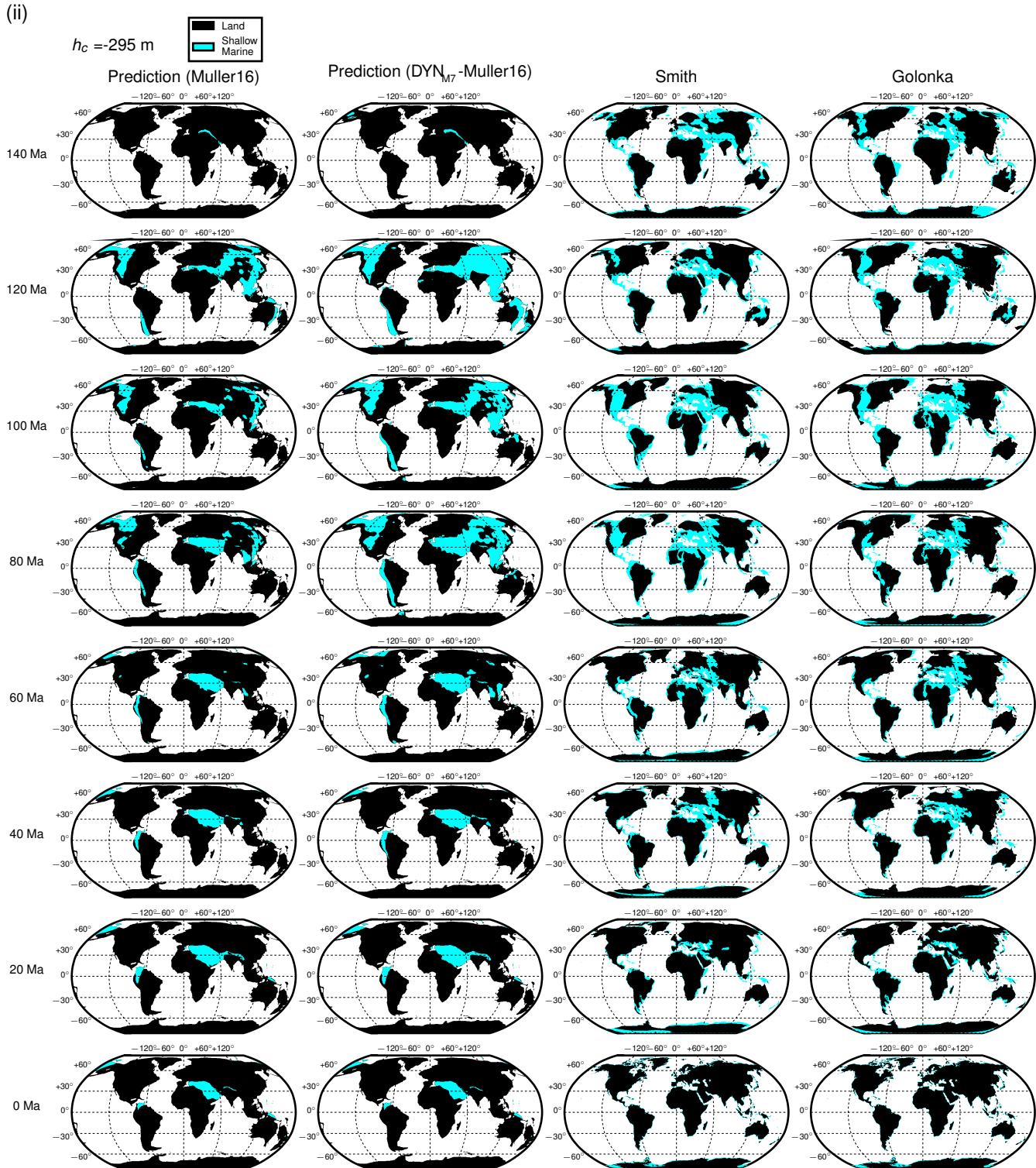
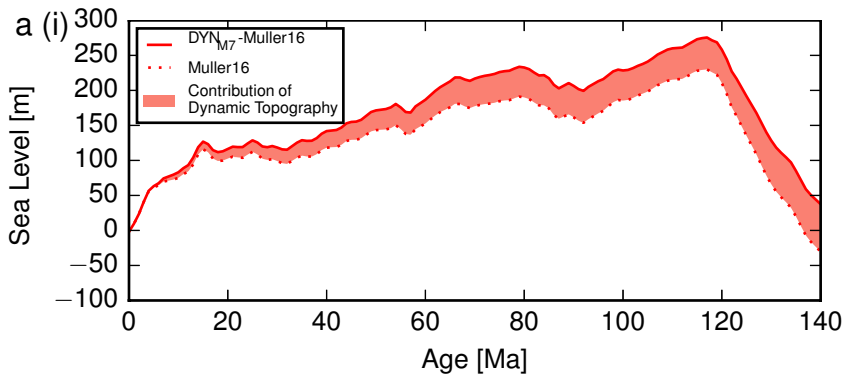
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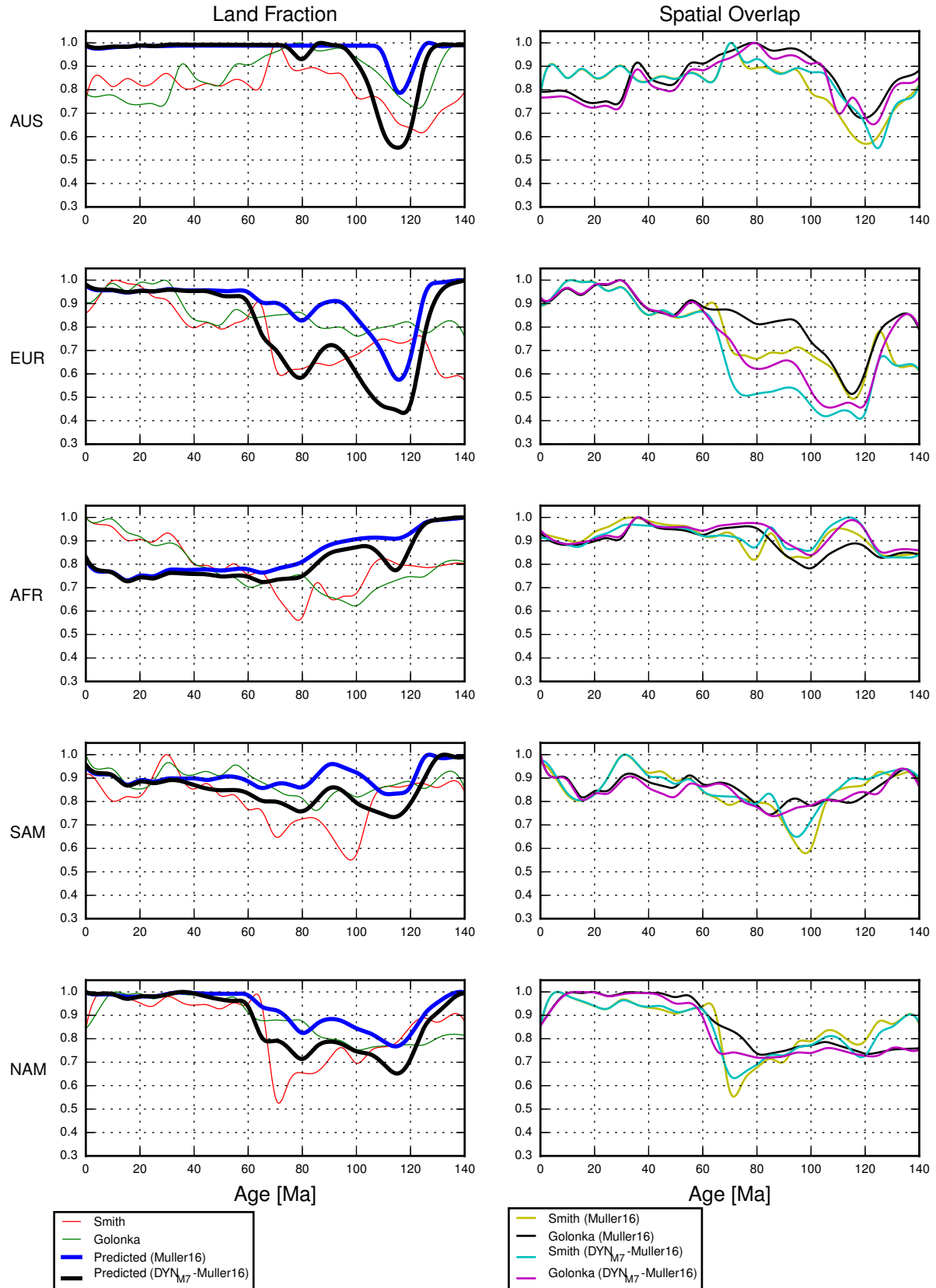


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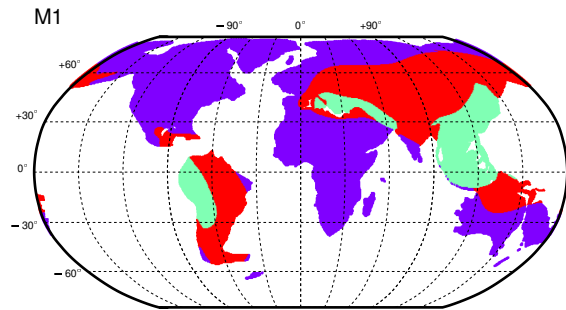




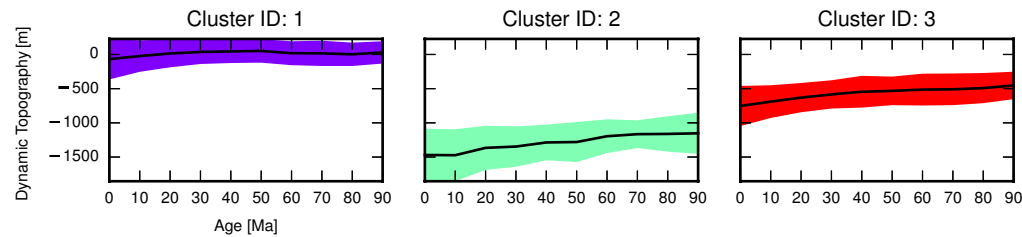
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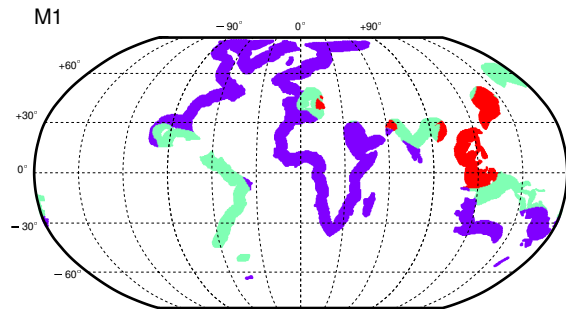
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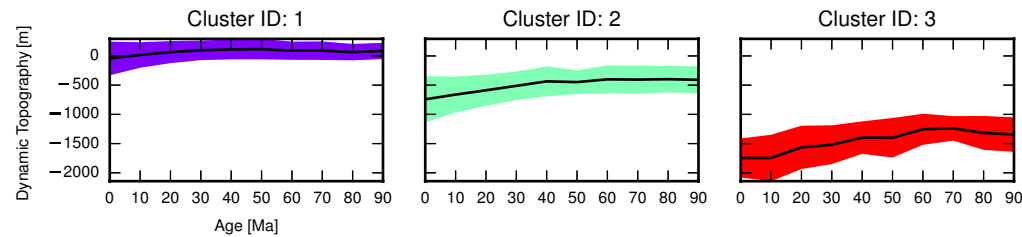
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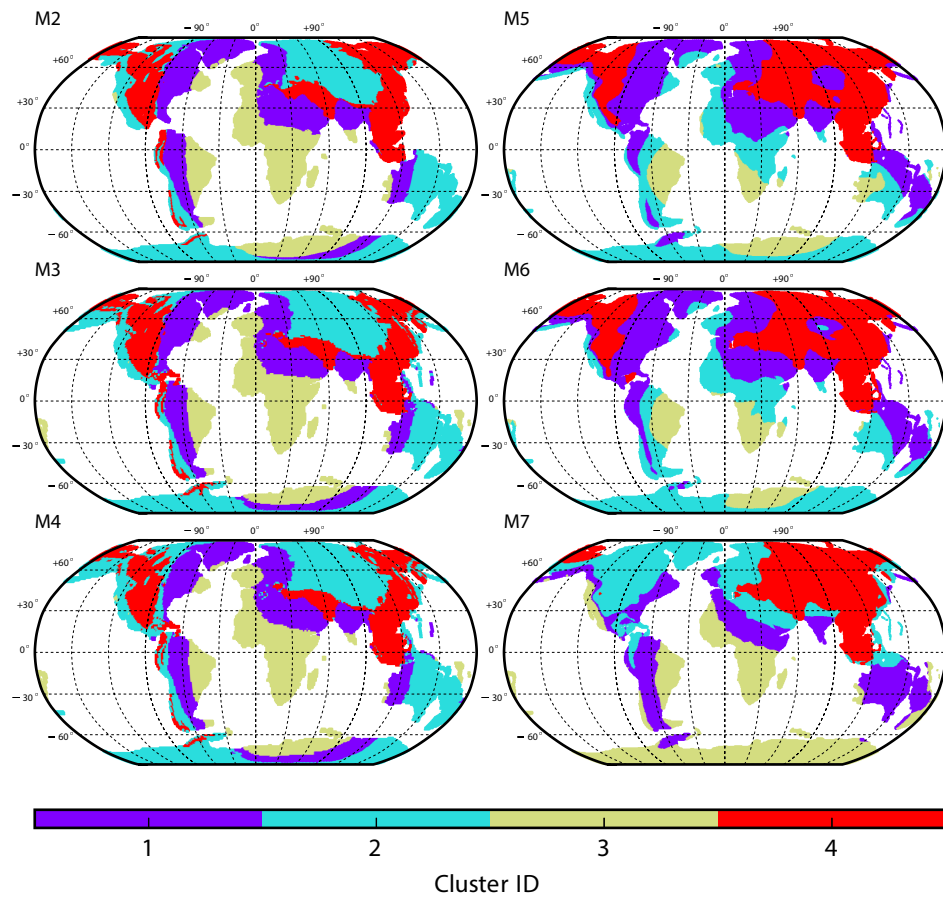
c



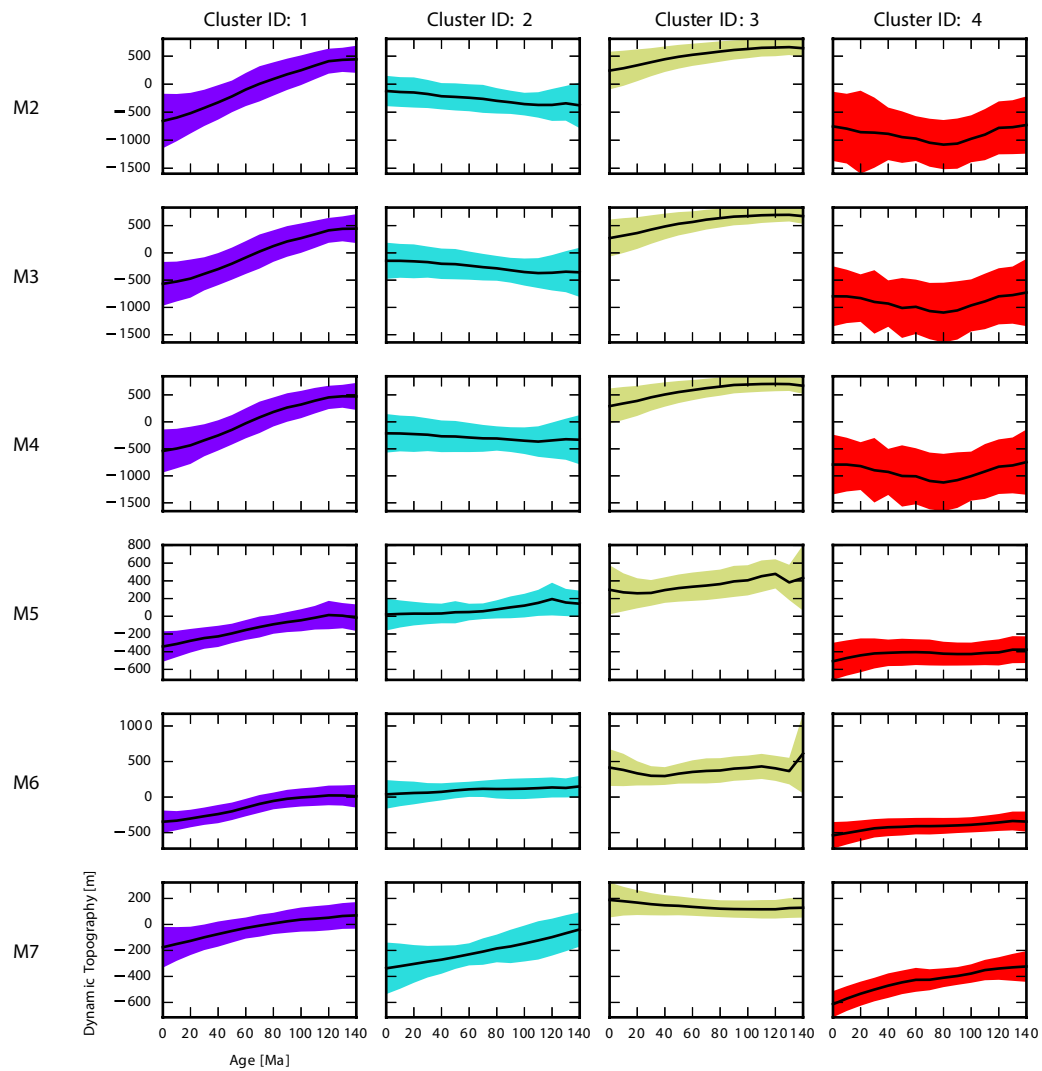
d



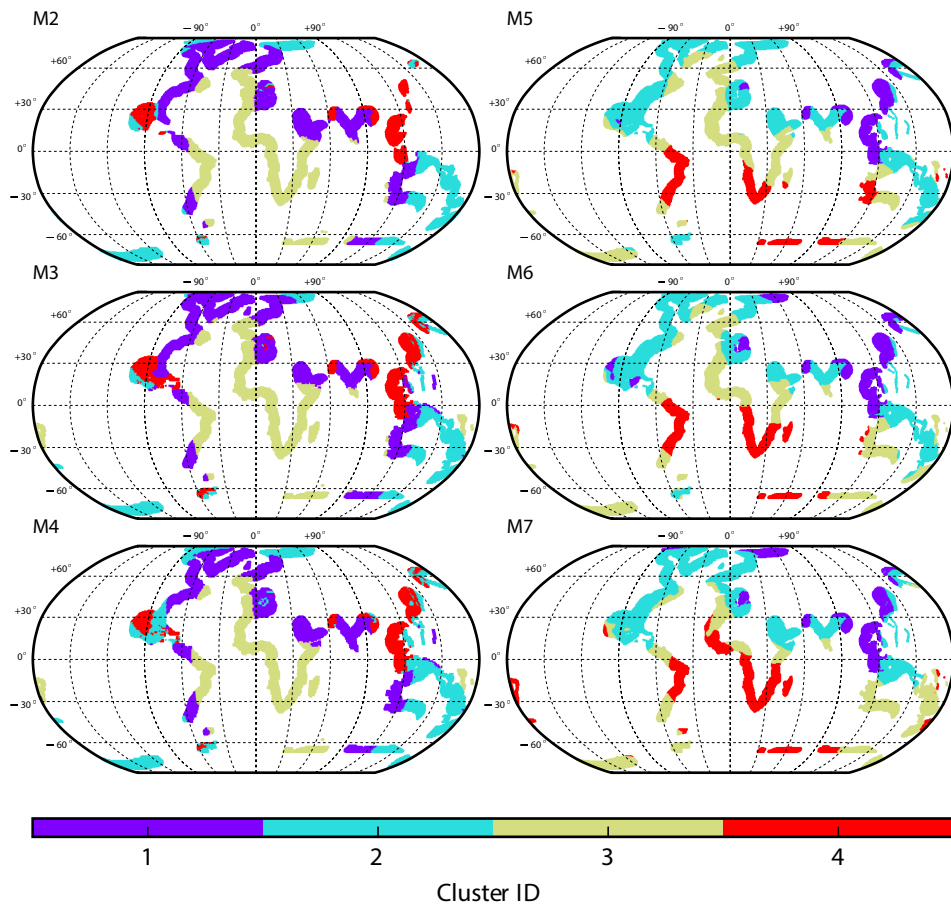
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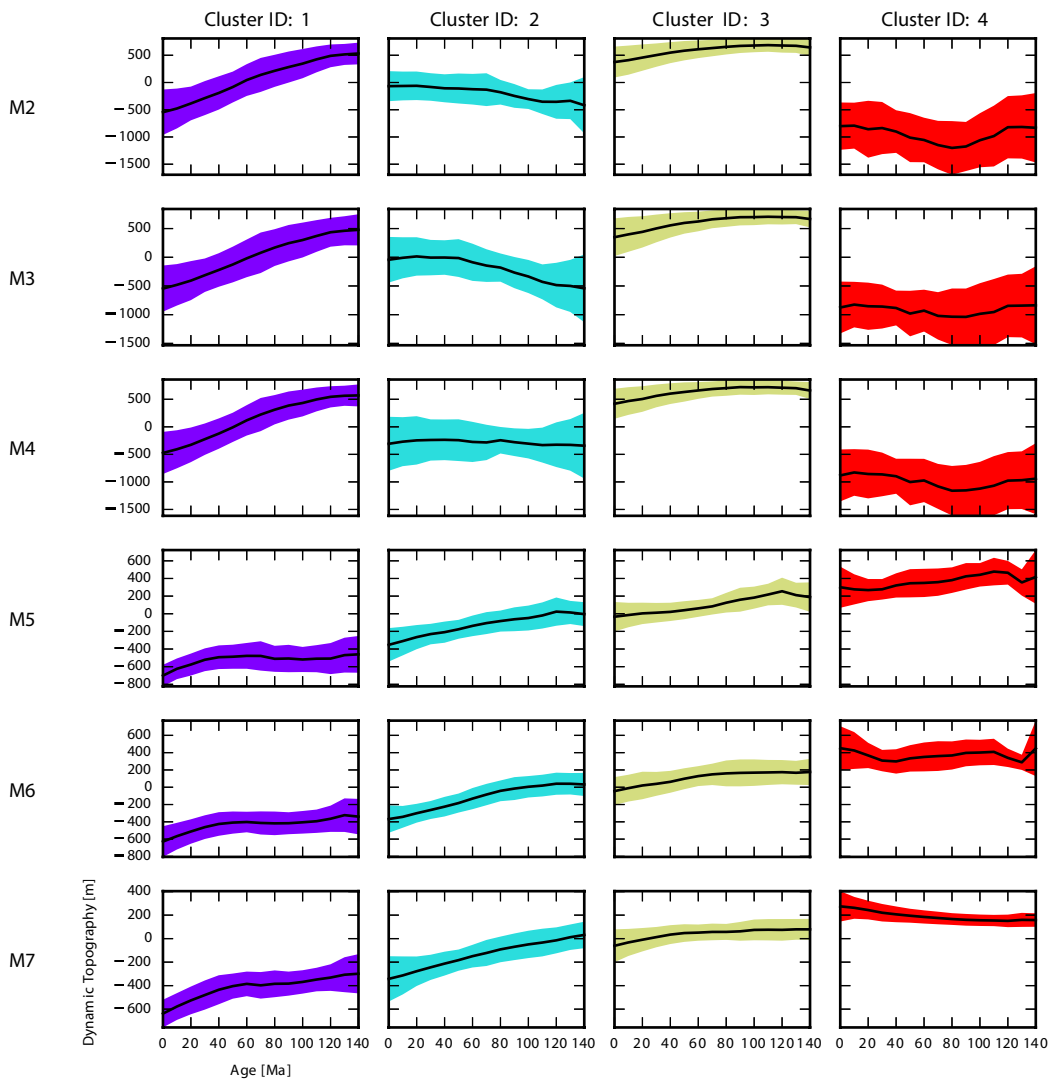
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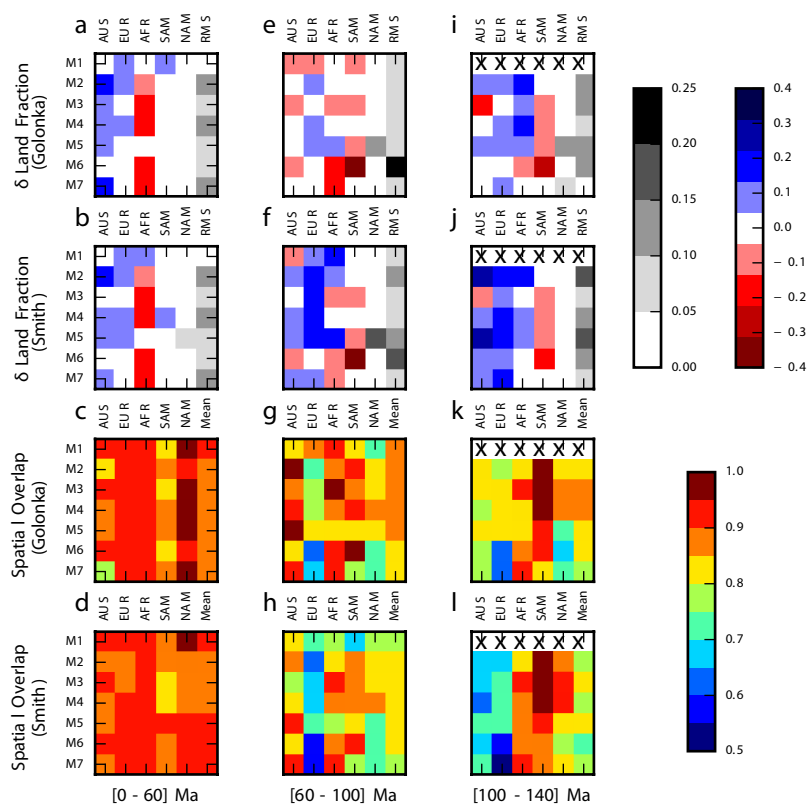


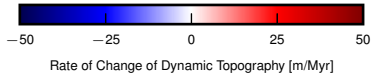
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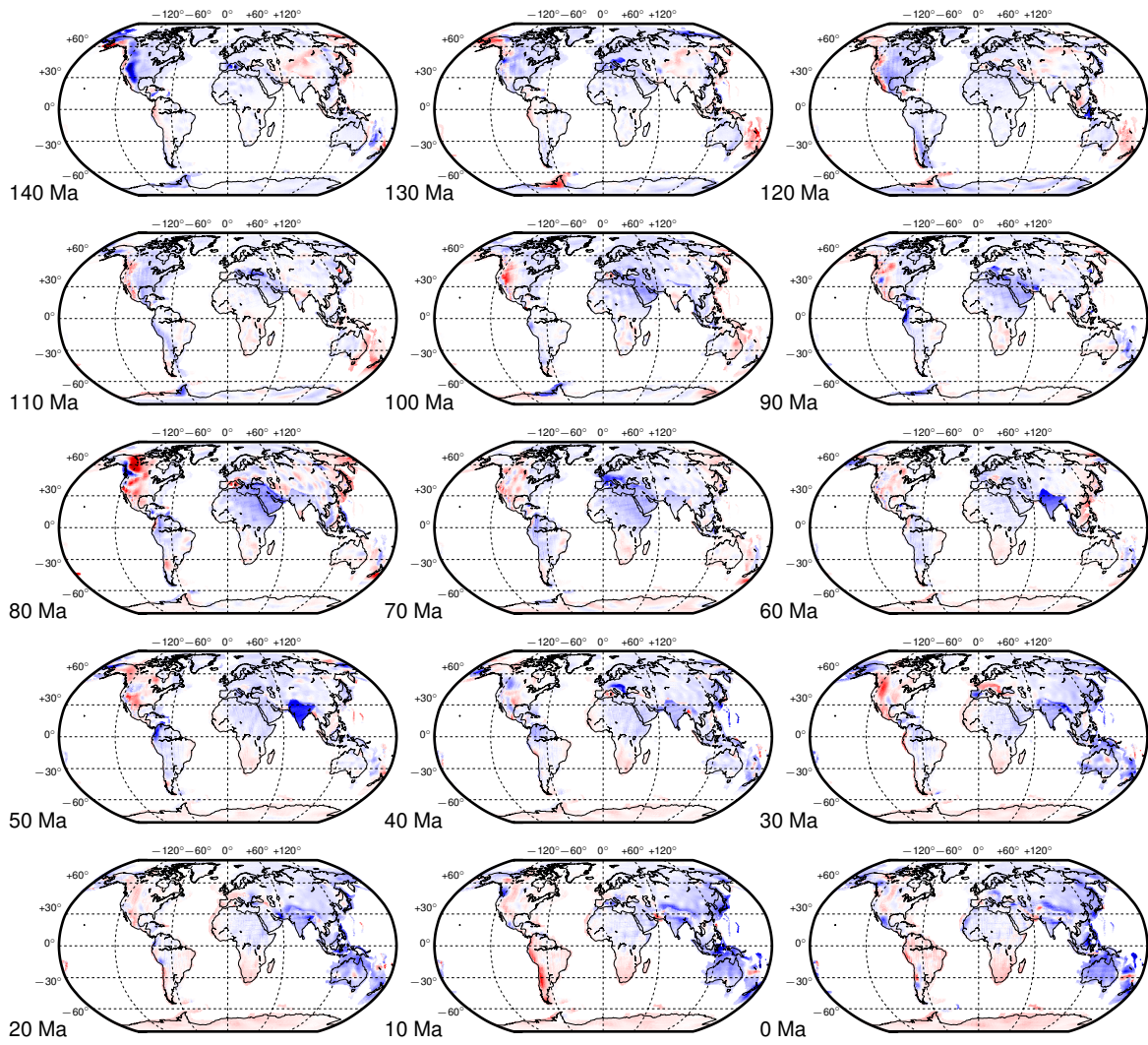
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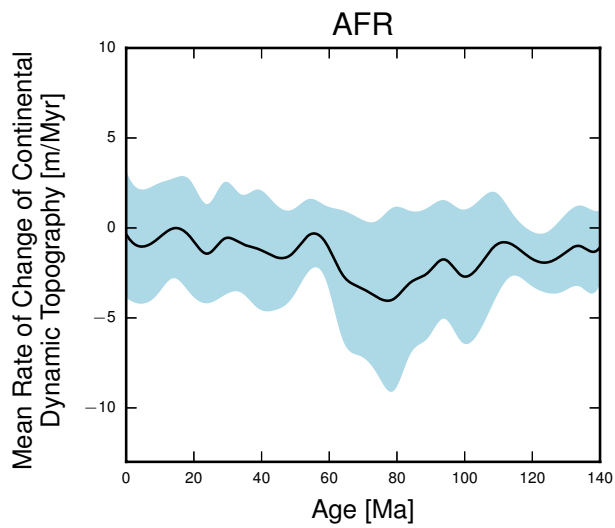
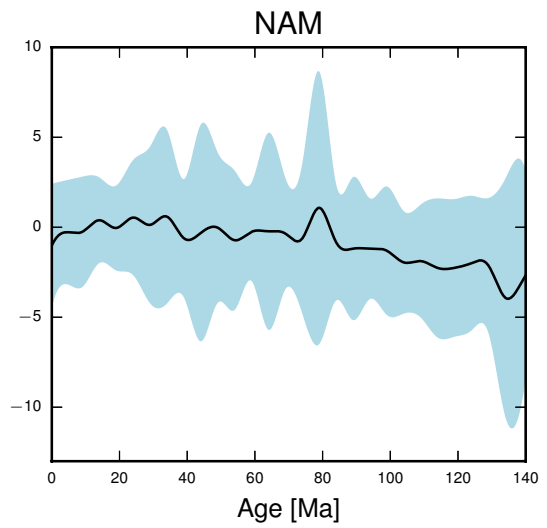
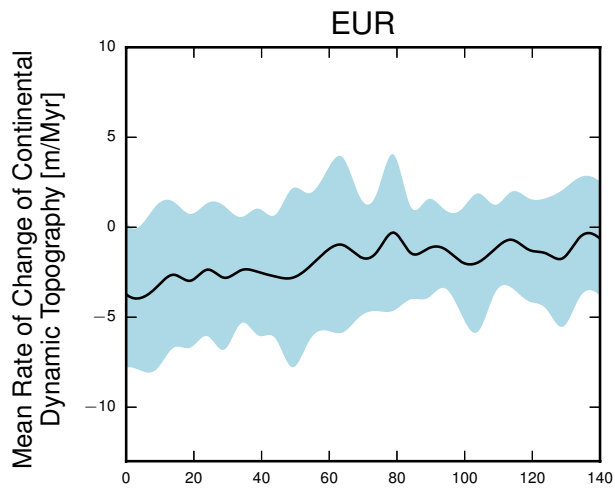
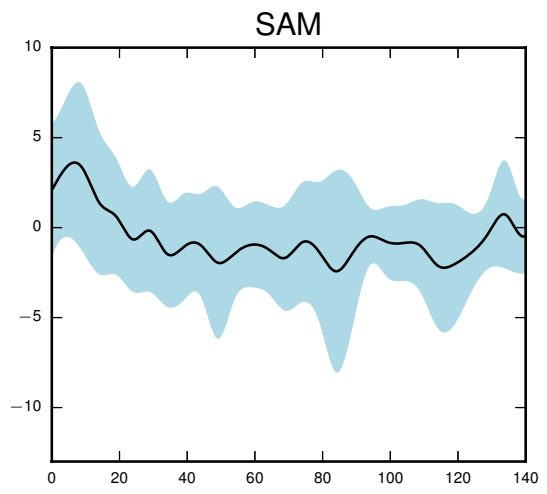
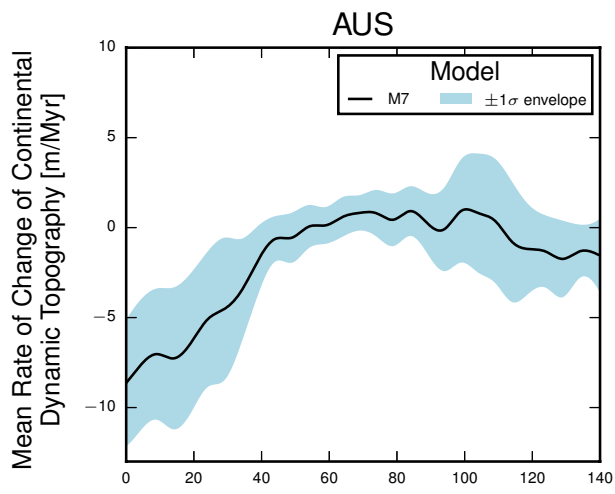




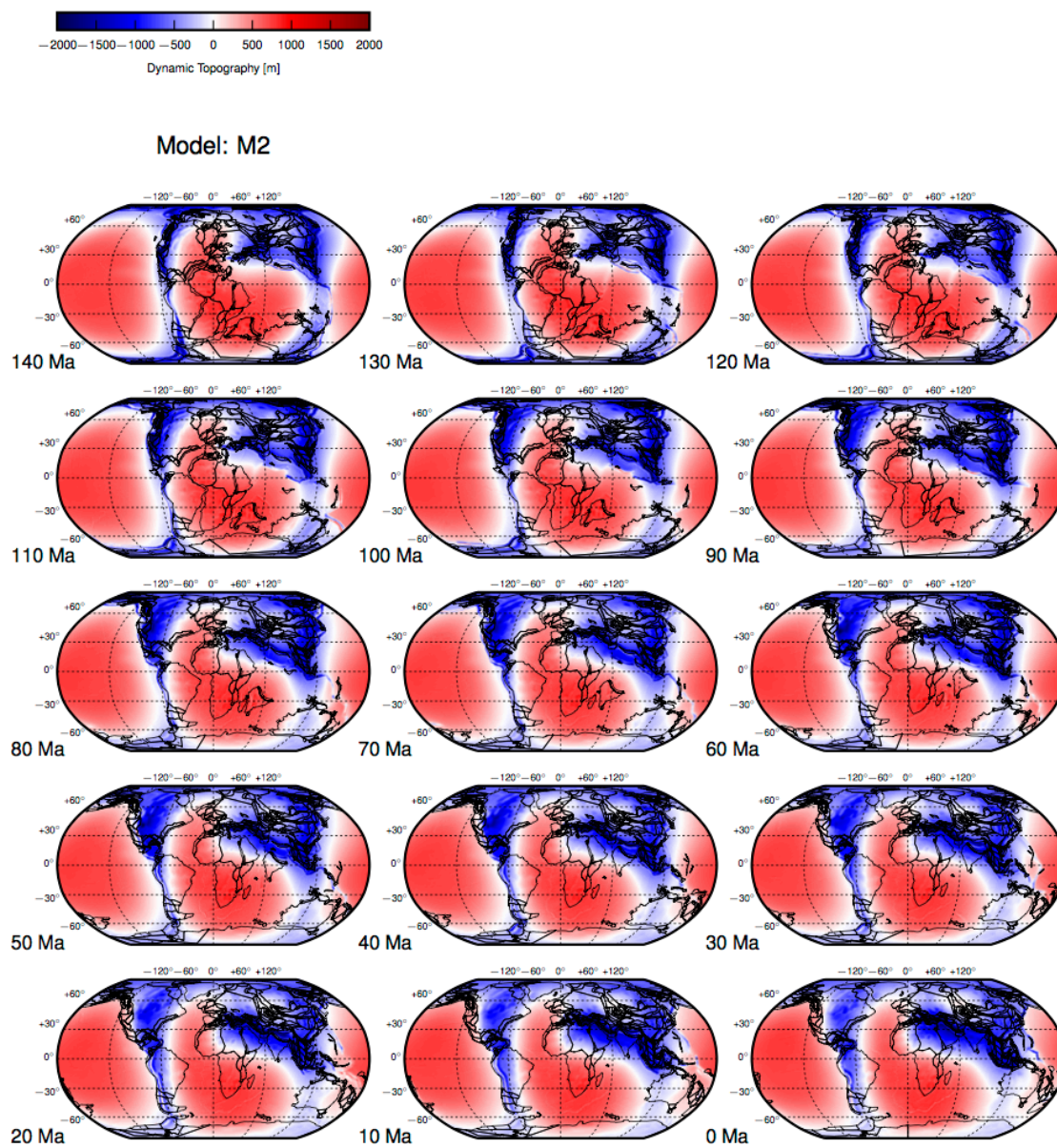


Model: M7

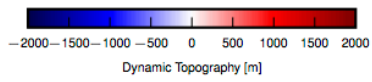




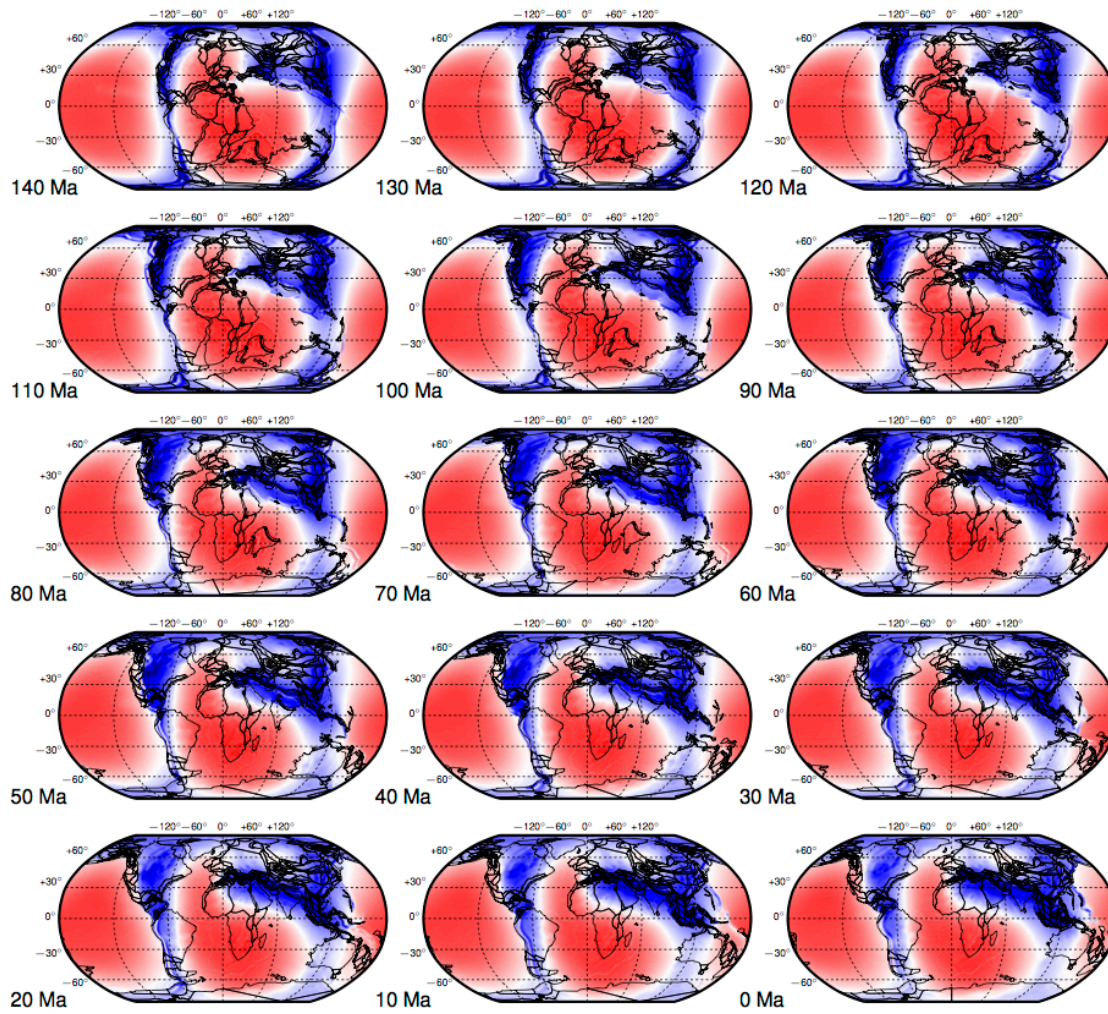
Supplementary figures to “Dynamic topography of passive continental margins and their hinterlands since the Cretaceous” by R.D. Müller, R. Hassan, M. Gurnis, N. Flament, S.E. Williams.



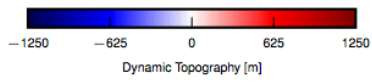
Supp. Fig. 1. Modelled dynamic topography in model M2 at 10 Myr intervals, with reconstructed continents (Seton et al., 2012) overlain.



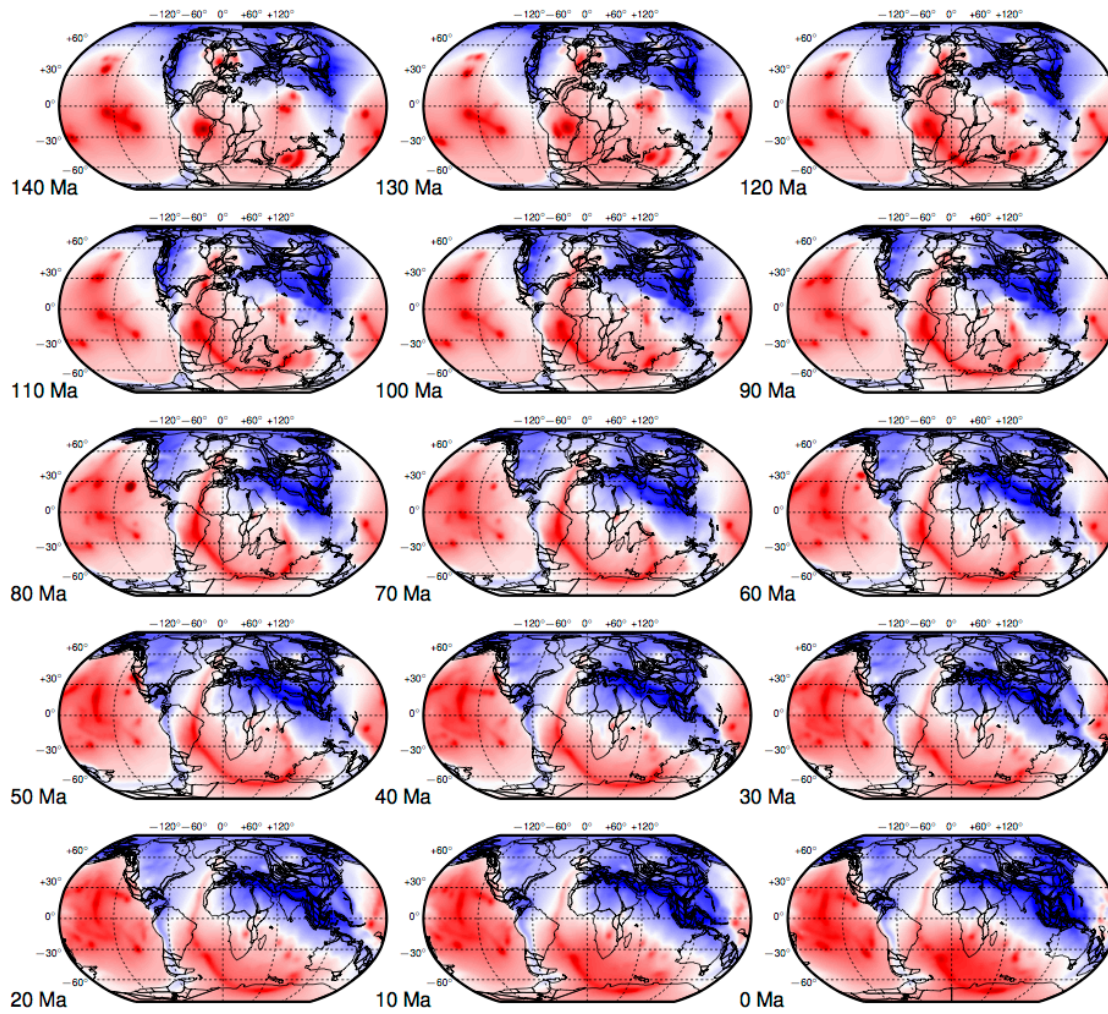
Model: M4



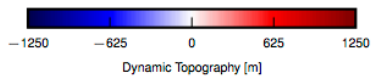
Supp. Fig. 2. Modelled dynamic topography in model M4 at 10 Myr intervals, with reconstructed continents (Van Der Meer et al., 2010) overlain.



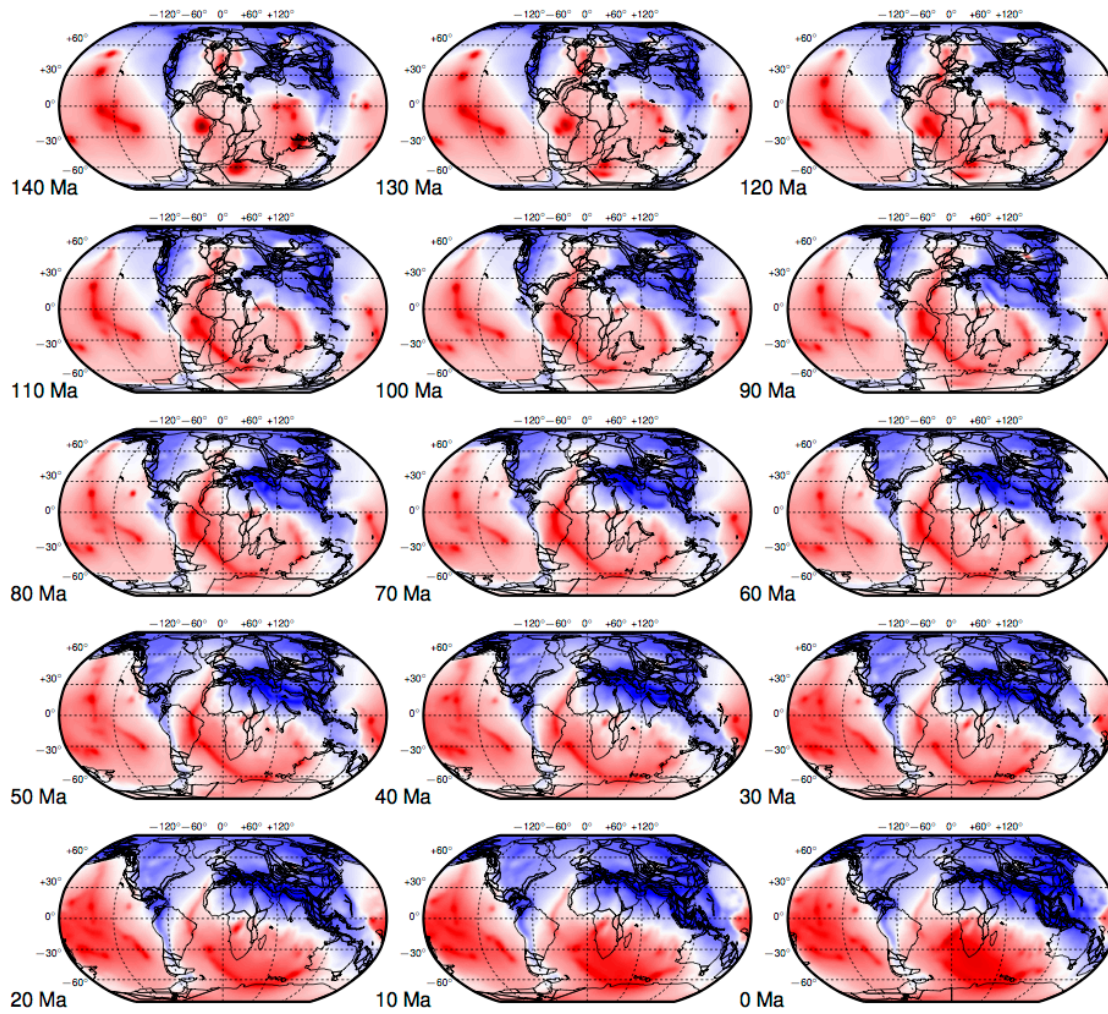
Model: M5



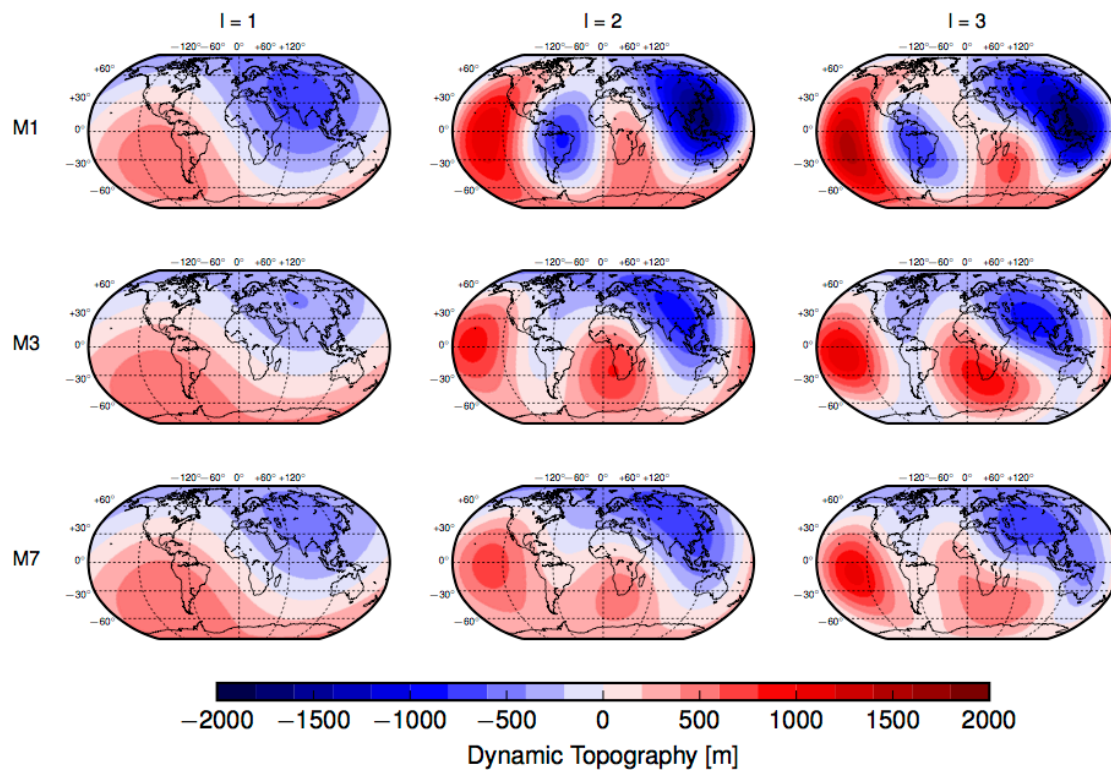
Supp. Fig. 3. Modelled dynamic topography in model M5 at 10 Myr intervals, with reconstructed continents (Seton et al., 2012) overlain.



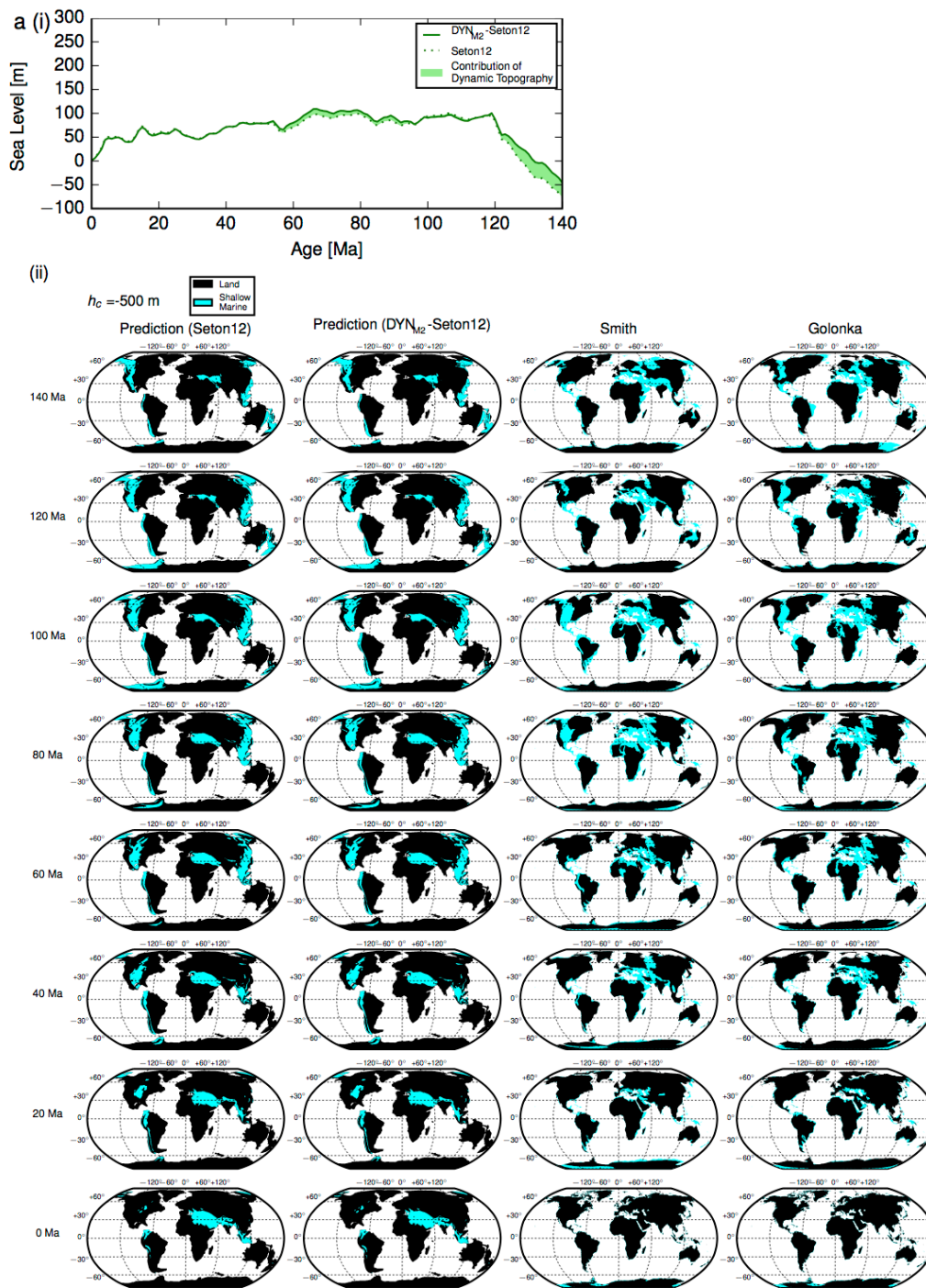
Model: M6



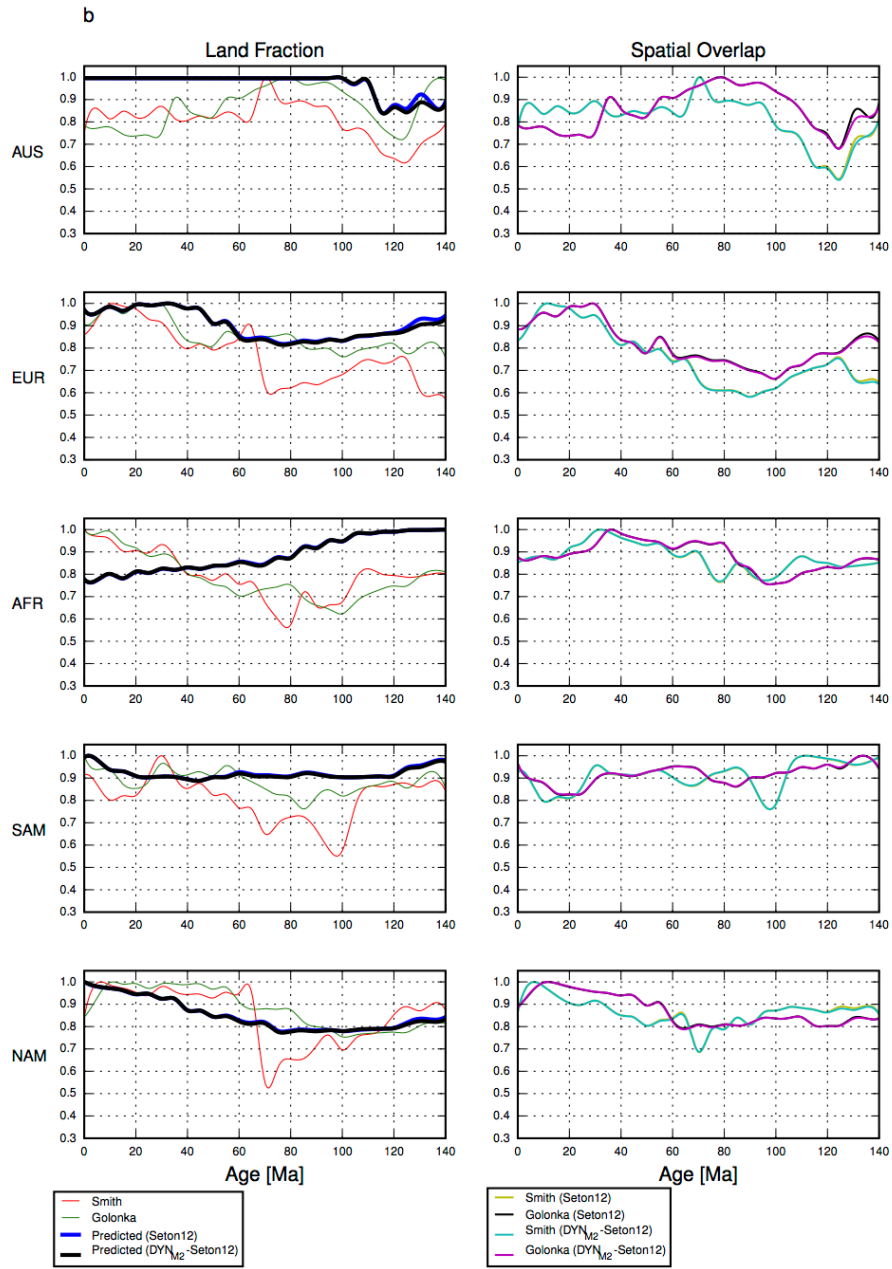
Supp. Fig. 4. Modelled dynamic topography in model M6 at 10 Myr intervals, with reconstructed continents (Müller et al., 2016) overlain.



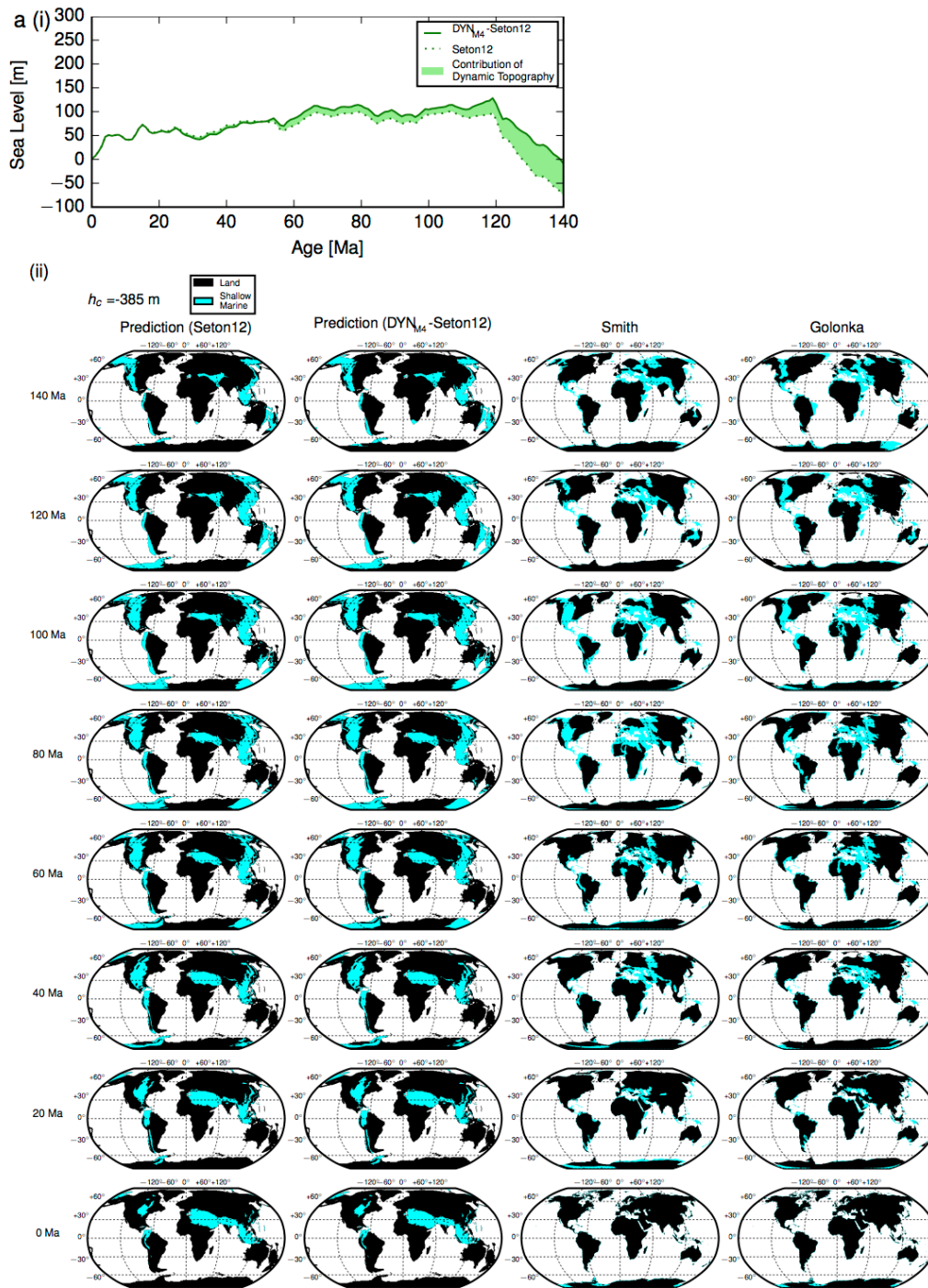
Supp. Fig. 5. Low-degree (1-3) spherical harmonic representations of present-day modelled dynamic topography in models M1, M3 and M7.



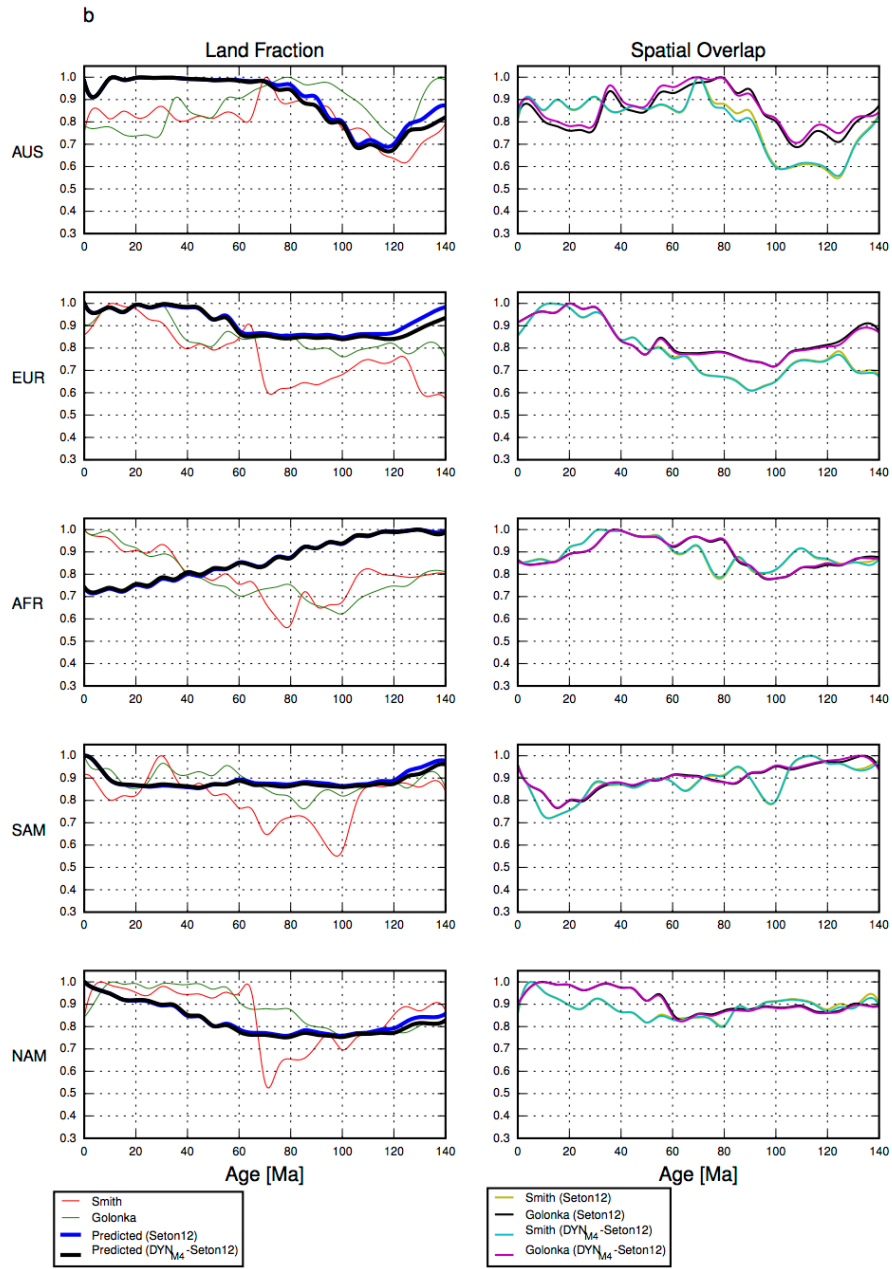
Supp. Fig. 6. (a) The distribution of oceans and continents predicted in model M2, based on the H87 long-term sea level curve, is shown in the first column at labelled ages (see methods for more details). The second column shows equivalent predictions based on the M16 sea level curve. The distribution of oceans and continents in the paleogeography models of Smith et al. (1994) and Golonka (2007, 2009) are shown in the third and fourth columns, respectively.



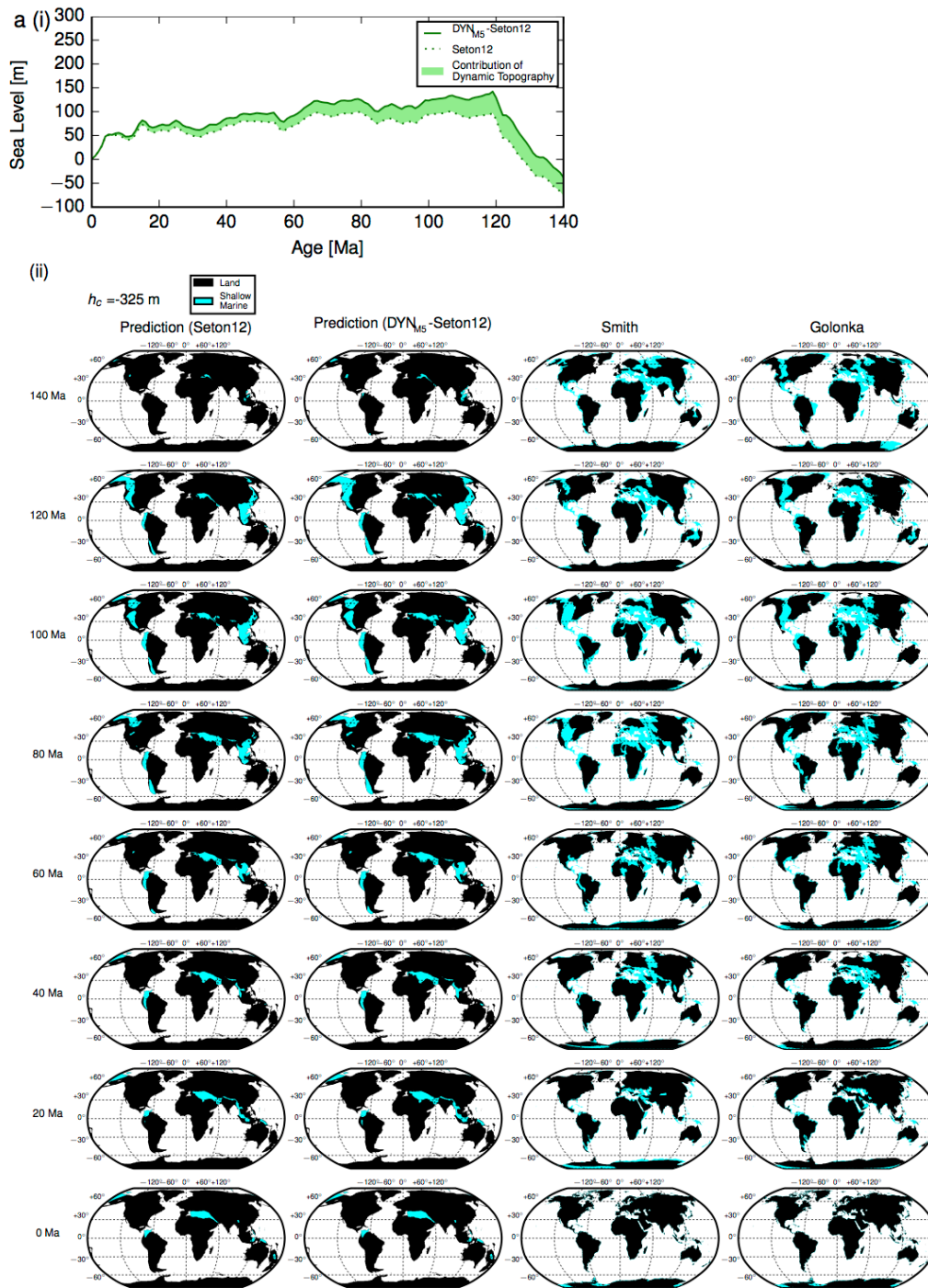
Supp. Fig. 6 (b). Evolution of the fraction of land, $\tau(t)$, over the last 140 Ma as predicted in model M2, based on sea level curves H87 (blue) and M16 (black) are shown on the first column for each labelled continent. Evolution of the fraction of land, $\tau(t)$, computed for paleogeography grids from Smith et al. (1994) and Golonka (2007, 2009) are also shown in red and green, respectively, for labelled ages in the first column. The yellow and black curves in the second column show the evolution of spatial overlap, $\mu(t)$, (see methods) between predictions of inundation patterns within each continent in model M2, based on sea level curve H87, and those from paleogeography grids of Smith et al. (1994) and Golonka (2007, 2009), respectively. The cyan and magenta curves show the equivalent result, but for sea level curve M16.



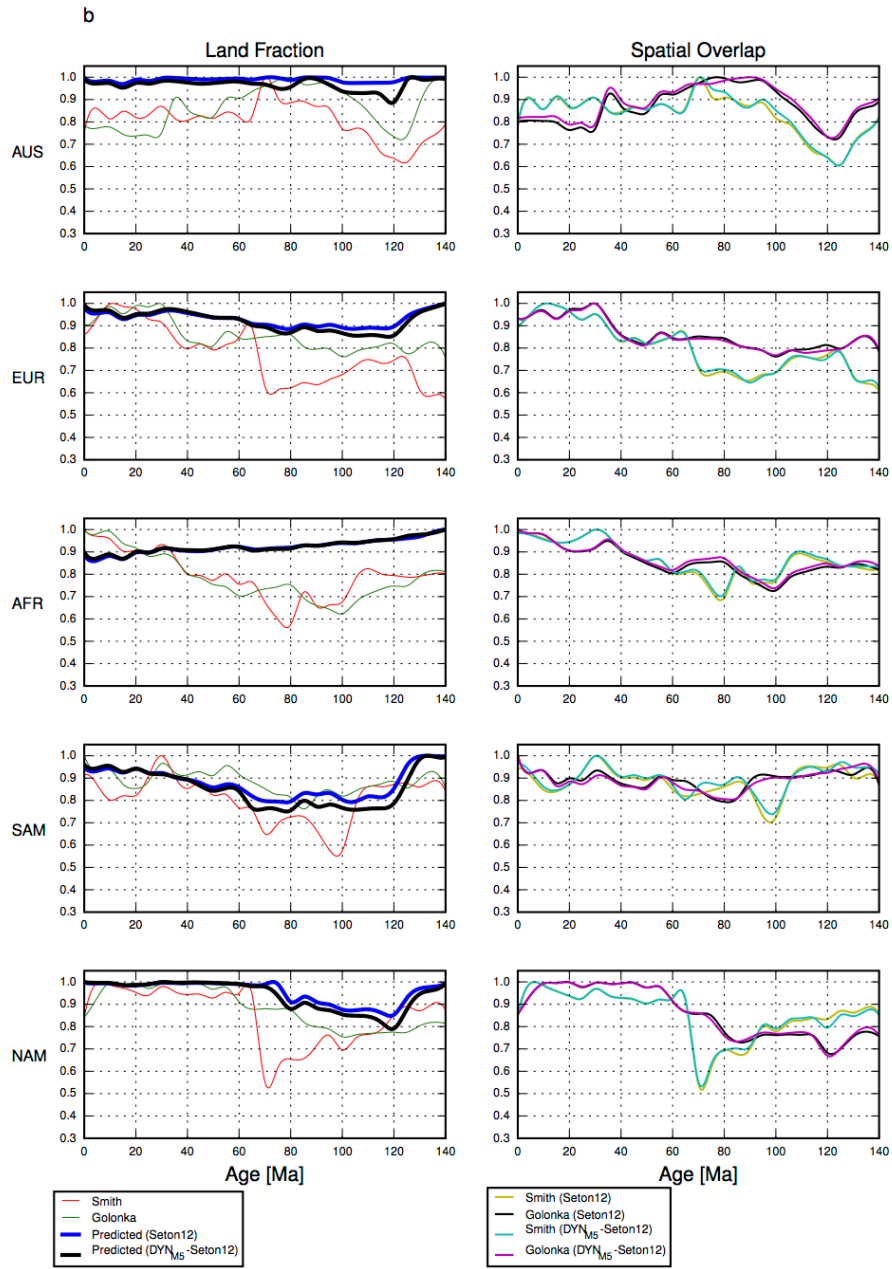
Supp. Fig. 7. (a) The distribution of oceans and continents predicted in model M4, based on the H87 long-term sea level curve, is shown in the first column at labelled ages (see methods for more details). The second column shows equivalent predictions based on the M16 sea level curve. The distribution of oceans and continents in the paleogeography models of Smith et al. (1994) and Golonka (2007, 2009) are shown in the third and fourth columns, respectively.



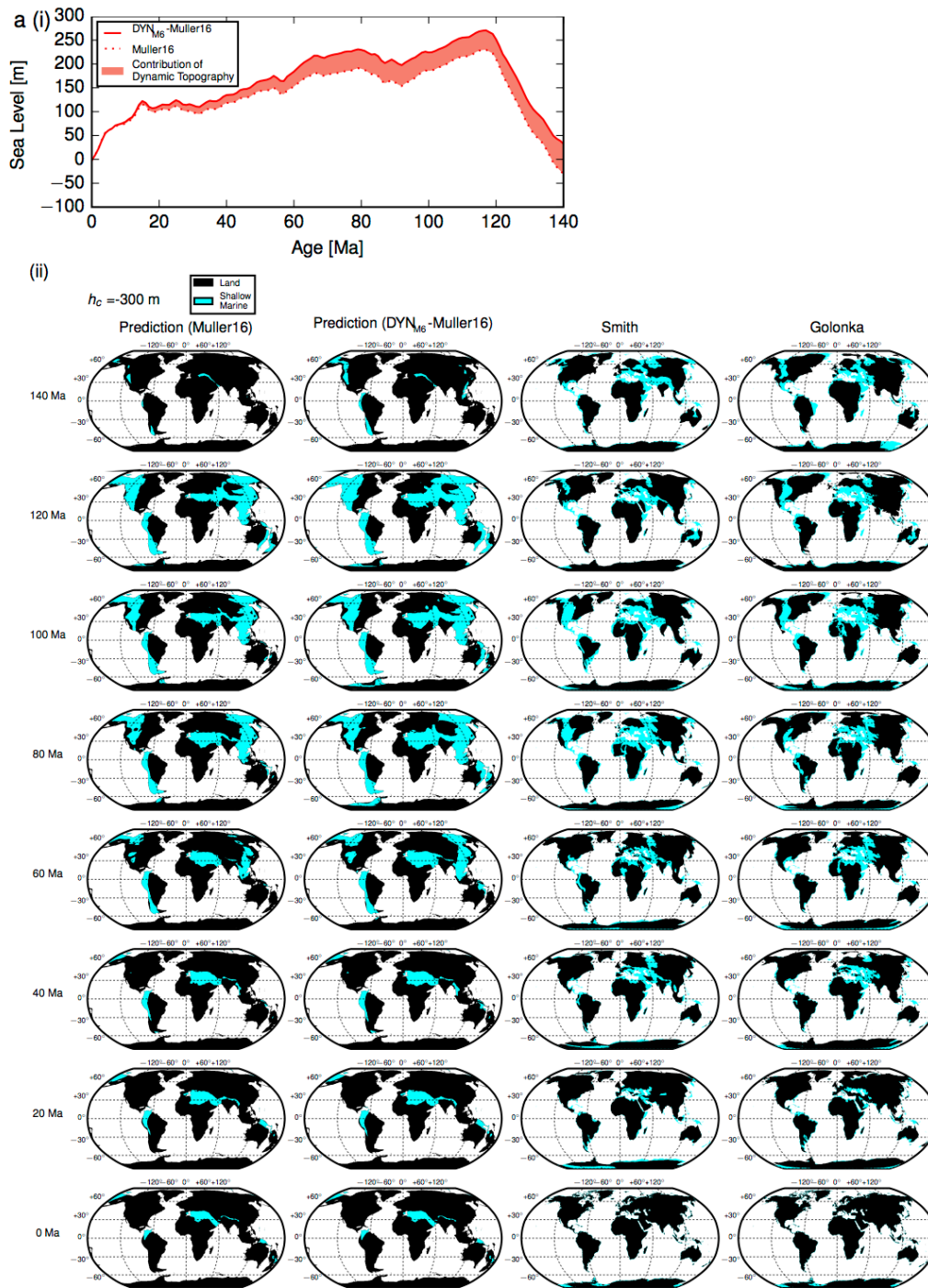
Supp. Fig. 7 (b). Evolution of the fraction of land, $\tau(t)$, over the last 140 Ma as predicted in model M4, based on sea level curves H87 (blue) and M16 (black) are shown on the first column for each labelled continent. Evolution of the fraction of land, $\tau(t)$, computed for paleogeography grids from Smith et al. (1994) and Golonka (2007, 2009) are also shown in red and green, respectively, for labelled ages in the first column. The yellow and black curves in the second column show the evolution of spatial overlap, $\mu(t)$, (see methods) between predictions of inundation patterns within each continent in model M4, based on sea level curve H87, and those from paleogeography grids of Smith et al. (1994) and Golonka (2007, 2009), respectively. The cyan and magenta curves show the equivalent result, but for sea level curve M16.



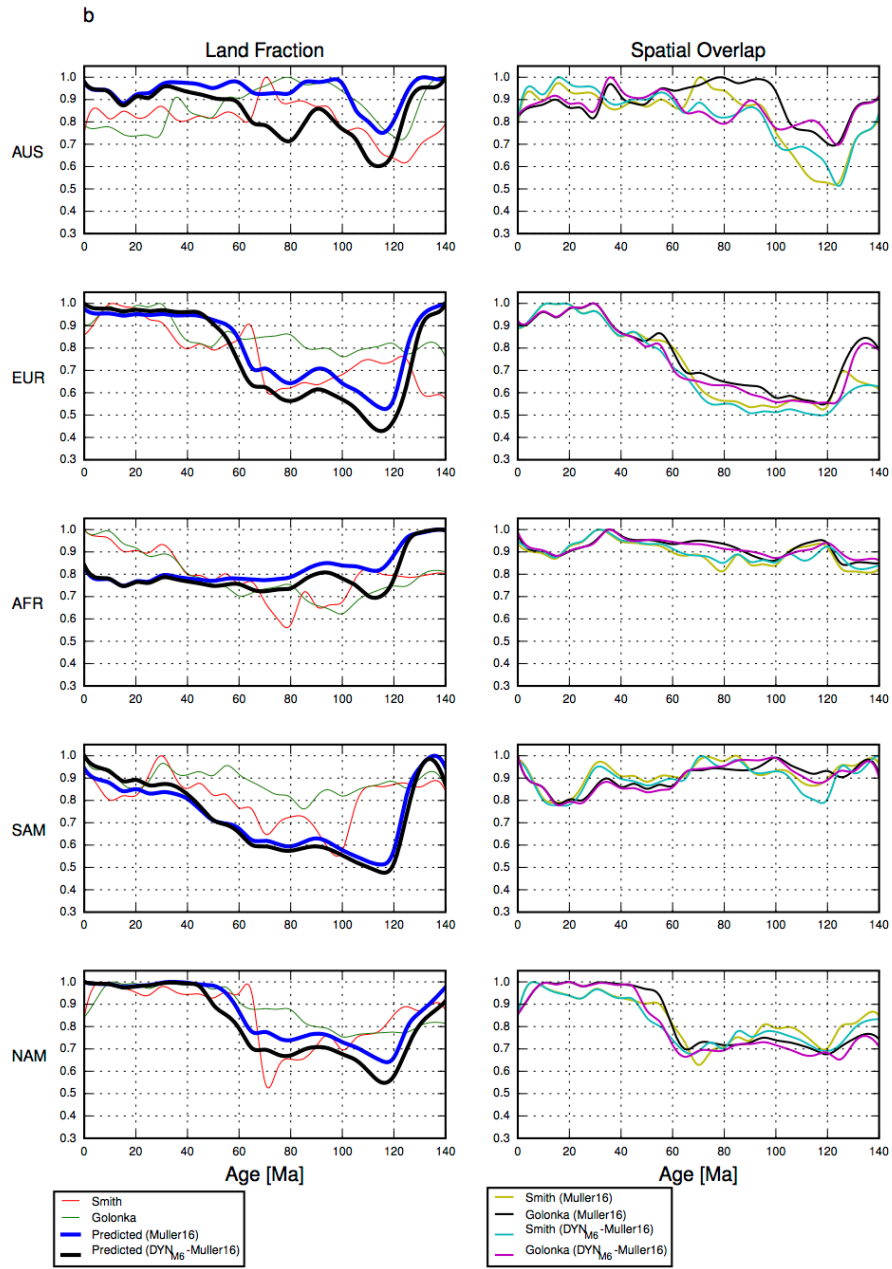
Supp. Fig. 8. (a) The distribution of oceans and continents predicted in model M5, based on the H87 long-term sea level curve, is shown in the first column at labelled ages (see methods for more details). The second column shows equivalent predictions based on the M16 sea level curve. The distribution of oceans and continents in the paleogeography models of Smith et al. (1994) and Golonka (2007, 2009) are shown in the third and fourth columns, respectively.



Supp. Fig. 8 (b). Evolution of the fraction of land, $\tau(t)$, over the last 140 Ma as predicted in model M5, based on sea level curves H87 (blue) and M16 (black) are shown on the first column for each labelled continent. Evolution of the fraction of land, $\tau(t)$, computed for paleogeography grids from Smith et al. (1994) and Golonka (2007, 2009) are also shown in red and green, respectively, for labelled ages in the first column. The yellow and black curves in the second column show the evolution of spatial overlap, $\mu(t)$, (see methods) between predictions of inundation patterns within each continent in model M5, based on sea level curve H87, and those from paleogeography grids of Smith et al. (1994) and Golonka (2007, 2009), respectively. The cyan and magenta curves show the equivalent result, but for sea level curve M16.



Supp. Fig. 9 (a). The distribution of oceans and continents predicted in model M6, based on the H87 long-term sea level curve, is shown in the first column at labelled ages (see methods for more details). The second column shows equivalent predictions based on the M16 sea level curve. The distribution of oceans and continents in the paleogeography models of Smith et al. (1994) and Golonka (2007, 2009) are shown in the third and fourth columns, respectively.



Supp. Fig. 9 (b). Evolution of the fraction of land, $\tau(t)$, over the last 140 Ma as predicted in model M6, based on sea level curves H87 (blue) and M16 (black) are shown on the first column for each labelled continent. Evolution of the fraction of land, $\tau(t)$, computed for paleogeography grids from Smith et al. (1994) and Golonka (2007, 2009) are also shown in red and green, respectively, for labelled ages in the first column. The yellow and black curves in the second column show the evolution of spatial overlap, $\mu(t)$, (see methods) between predictions of inundation patterns within each continent in model M6, based on sea level curve H87, and those from paleogeography grids of Smith et al. (1994) and Golonka (2007, 2009), respectively. The cyan and magenta curves show the equivalent result, but for sea level curve M16.

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