REGULARIZED INNER PRODUCTS AND WEAKLY HOLOMORPHIC HECKE EIGENFORMS

KATHRIN BRINGMANN AND BEN KANE

1. Introduction and statement of results

Hecke operators play a central role in the study of modular forms. The classical theory of newforms and modular *L*-functions relies on the fact that spaces of cusp forms are finite-dimensional and can be diagonalized with respect to the Hecke algebra. Although spaces of weakly holomorphic modular forms are not finite-dimensional, it turns out that Hecke operators still play an important role and there is a natural subspace whose elements can be considered as trivial. Then a natural Hecke structure can be described.

For $\kappa \in \mathbb{Z}$, denote by $M_{2\kappa}^!$ the space of weight 2κ weakly holomorphic modular forms (i.e., those meromorphic modular forms whose only possible pole occurs at $i\infty$) with respect to the full modular group $\mathrm{SL}_2(\mathbb{Z})$. Motivated by applications coming from the theory of p-adic modular forms, Guerzhoy [12] observed that even though the space $M_{2k}^!$ ($k \in \mathbb{N}$) is infinite-dimensional and the Hecke operators T_m (defined in (2.5)) increase the order of the pole, there is a meaningful Hecke theory on this space. Emulating the usual definition of Hecke eigenforms (i.e., for every $m \in \mathbb{N}$ there exists $\lambda_m \in \mathbb{C}$ such that $f|T_m = \lambda_m f$), Guerzhoy defined a weakly holomorphic Hecke eigenform to be any $f \in M_{2k}^! \setminus \mathcal{J}$ (\mathcal{J} is a specific subspace of $M_{2k}^!$ which is stable under the Hecke algebra) which satisfies

$$f|T_m \equiv \lambda_m f \pmod{\mathcal{J}}.$$

Here and throughout, for a vector space \mathcal{S} and a subspace $\mathcal{J} \subset \mathcal{S}$, for $f, g \in \mathcal{S}$ the congruence $f \equiv g \pmod{\mathcal{J}}$ means that $f - g \in \mathcal{J}$. In [12], Guerzhoy chose

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 $\mathcal{J}=D^{2k-1}(M_{2-2k}^!)$, where $D:=\frac{1}{2\pi i}\frac{\partial}{\partial z}$. To understand Guerzhoy's motivation, recall that Bruinier, Ono, and Rhoades [10] showed that D^{2k-1} is an injective map from the space of so-called harmonic Maass forms (non-holomorphic generalizations of modular forms) into the space $M_{2k}^!$ and determined its image. Guerzhoy realized that the image of the distinguished subspace $M_{2-2k}^!$ was a natural object to factor out by. However, it was later determined that $\mathcal{J}=D^{2k-1}(S_{2-2k}^!)$ is a better choice (cf. [7, 14]), where $S_{2\kappa}^!$ is the space of weak cusp forms of weight $2\kappa \in 2\mathbb{Z}$, i.e., those weakly holomorphic modular forms with vanishing constant term. Note that D^{2k-1} and T_m essentially commute (see (2.4)) and $S_{2-2k}^!$ is preserved by T_m , so $D^{2k-1}(S_{2-2k}^!)|T_m \subseteq D^{2k-1}(S_{2-2k}^!)$, and hence $D^{2k-1}(S_{2-2k}^!)$ is stable under the Hecke operators. Thus a weakly holomorphic Hecke eigenform is any $f \in M_{2k}^! \setminus D^{2k-1}(S_{2-2k}^!)$ for which there exist $\lambda_m \in \mathbb{C}$ satisfying (for all $m \in \mathbb{N}$)

$$f|T_m \equiv \lambda_m f \pmod{D^{2k-1}(S_{2-2k}!)}$$
 (1.1)

Functions f satisfying (1.1) are not eigenfunctions of T_m in the usual sense. However a reasonable interpretation is to consider $[f] := f + D^{2k-1}(S_{2-2k}^!)$ as an element of the quotient space $M_{2k}^!/D^{2k-1}(S_{2-2k}^!)$. Functions f satisfying (1.1) may then be viewed as eigenvectors [f] of the Hecke operators in that factor space. As is usual, we exclude $[0] = D^{2k-1}(S_{2-2k}^!)$ in the definition of eigenvector.

Viewed in this light, the definition of weakly holomorphic Hecke eigenforms is perhaps not very enlightening at first glance. If one simply defines "eigenforms" as elements of some quotient space, then one can replace $D^{2k-1}(S_{2-2k}^!)$ with any space \mathcal{J} that is preserved under the action of the Hecke operators. It is hence natural to ask why $D^{2k-1}(S_{2-2k}^!)$ is the correct subspace. There are a few answers to this question. The initial perspective taken in [12] was a p-adic one, because for $f \in D^{2k-1}(S_{2-2k}^!)$ with integral coefficients the $(p^m n)$ -th coefficients become divisible by higher powers of p as m gets larger. Elements of $D^{2k-1}(S_{2-2k}^!)$ have vanishing period polynomials [7] and there is a recent cohomological interpretation as a space of coboundaries announced by Funke. Moreover certain regularized critical L-values are zero for elements in this space (see [6, Theorem 2.5]).

In this paper, we give another reason for the choice $\mathcal{J} = D^{2k-1}(S_{2-2k}^!)$ by viewing these Hecke eigenforms in the framework of a regularized inner product \langle , \rangle , defined in [4] for arbitrary weakly holomorphic modular forms. Note that if $f \in M_{2k}$ is a classical holomorphic Hecke eigenform, then for each $m \in \mathbb{N}$

there exists $\lambda_m \in \mathbb{C}$, for which

$$\langle f|T_m,g\rangle = \lambda_m\langle f,g\rangle \tag{1.2}$$

for all $g \in M_{2k}$. Indeed, since the inner product is non-degenerate on M_{2k} (see [18, Section 5]), (1.2) is equivalent to the usual definition of Hecke eigenforms on M_{2k} .

One may hence treat (1.2) as an alternative definition of Hecke eigenforms. Both from this perspective and from the perspective of investigating the inner product, it is very natural to factor out by the space

$$M_{2k}^{!,\perp} := \left\{ f \in M_{2k}^{!} : \langle f, g \rangle = 0 \text{ for all } g \in M_{2k}^{!} \right\}$$
 (1.3)

of weakly holomorphic forms which are orthogonal to all of $M_{2k}^!$.

The main point of this paper is to give another motivation of the definition in [12] by proving that the space $D^{2k-1}(S_{2-2k}^!)$ is the "degenerate part" of $M_{2k}^!$ in the sense that it is orthogonal to all of $M_{2k}^!$, from which one also conversely concludes that $M_{2k}^{!,\perp}$ is a very large subspace.

Theorem 1.1. A function $f \in M_{2k}^!$ is in $M_{2k}^{!,\perp}$ if and only if $f \in D^{2k-1}(S_{2-2k}^!)$.

As alluded to above, Theorem 1.1 yields another characterization of weakly holomorphic Hecke eigenforms.

Corollary 1.2. A function $f \in M_{2k}^!$ is a weakly holomorphic Hecke eigenform if and only if for every $m \in \mathbb{N}$ there exists λ_m satisfying (1.2) for all $g \in M_{2k}^!$. Remarks.

(1) As shown in [4, Theorem 1.3], the inner product is Hermitian on $M_{2k}^!$. Hence if f is a weakly holomorphic Hecke eigenform, then (1.2) implies that for all $g \in M_{2k}^!$

$$\langle g, f | T_m \rangle = \overline{\lambda_m} \langle g, f \rangle$$
.

Moreover, by Lemma 2.1 below, the Hecke operators are also Hermitian. Thus, taking g = f, we see that $\lambda_m \in \mathbb{R}$ or $\langle f, f \rangle = 0$.

(2) The regularized inner product plays an important role in both mathematics and physics, due to the fact that Borcherds used this regularization to construct theta lifts of weakly holomorphic modular forms on orthogonal groups. For example, such lifts recently occurred in Angelantonj, Florakis, and Pioline's investigation [1] of BPS amplitudes with momentum insertions (see [1, (3.23)]). The same authors evaluate higher-dimensional analogues of these lifts in their continued study of one-loop integrals in [2, Section 2.2].

(3) The regularized inner product that Borcherds used was based on a regularization by physicists Harvey and Moore [13] and the regularization considered in this paper agrees with Harvey and Moore's regularization whenever theirs exists. Moreover, weakly holomorphic modular forms and mock modular forms appear in moonshine (such as the j-function occurring in McKay–Thompson's Monstrous Moonshine) and the theory of quantum black holes (see [11]). Due to these connections, it may be of interest to physicists to understand the structure of the inner product better; specifically, we discuss next whether the space of such functions has some kind of Hilbert space structure. As seen in Theorem 1.1, $M_{2k}^!$ contains a large degenerate subspace $D^{2k-1}(S_{2-2k}^!)$, and thus the space is itself not a Hilbert space. However, combining Theorem 1.1 with [7, Theorem 1.2], there is an isomorphism

$$M_{2k}^! / M_{2k}^{!,\perp} = M_{2k}^! / D^{2k-1} (S_{2-2k}^!) \cong M_{2k} \oplus M_{2k}.$$

In particular, a natural subspace is isomorphic to two copies of S_{2k} . Since S_{2k} is a finite-dimensional Hilbert space under the classical Petersson inner product, there is a naturally occurring Hilbert space of twice the dimension of S_{2k} , and the weakly holomorphic Hecke eigenforms form an orthogonal basis of this space. It may hence be natural to investigate physical properties arising from weakly holomorphic modular forms occurring within this subspace.

The paper is organized as follows. In Section 2, we introduce a family of non-holomorphic modular forms known as harmonic Maass forms which play an important role in the proof of Theorem 1.1 and then we explain how one can define regularized inner products. In Section 3, we prove Theorem 1.1 and Corollary 1.2.

2. Preliminaries

2.1. **Harmonic Maass forms.** We begin by defining harmonic Maass forms, which were first introduced by Bruinier–Funke [9].

Definition. For $\kappa \in \mathbb{Z}$, a harmonic Maass form of weight 2κ is a function $F : \mathbb{H} \to \mathbb{C}$, which is real analytic on \mathbb{H} and satisfies the following conditions:

(1) For every $M = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathrm{SL}_2(\mathbb{Z})$, we have

$$F\left(\frac{az+b}{cz+d}\right) = (cz+d)^{2\kappa}F(z).$$

(2) We have $\Delta_{2\kappa}(F) = 0$, where $\Delta_{2\kappa}$ is the weight 2κ hyperbolic Laplace operator (z = x + iy)

$$\Delta_{2\kappa} := -y^2 \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) + 2\kappa i y \left(\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right).$$

(3) The function F grows at most linear exponentially at $i\infty$.

A weight 2κ harmonic Maass form F has a Fourier expansion of the type $(c_F(n,y) \in \mathbb{C})$

$$F(z) = \sum_{n \in \mathbb{Z}} c_F(n, y) e^{2\pi i n z}.$$
 (2.1)

If F is weakly holomorphic, then $c_F(n,y) = c_F(n) \in \mathbb{C}$ is independent of y. More generally, one may use condition (2) above to determine the dependence on y and conclude that (2.1) naturally decomposes into holomorphic and non-holomorphic parts. Namely, for a harmonic Maass form F of weight 2-2k < 0,

$$F(z) = F^{+}(z) + F^{-}(z)$$
(2.2)

where, for some (unique) $c_F^{\pm}(n) \in \mathbb{C}$, we have

$$F^{+}(z) = \sum_{n \gg -\infty} c_F^{+}(n) e^{2\pi i n z},$$

$$F^{-}(z) = c_F^{-}(0) y^{2k-1} + \sum_{\substack{n \ll \infty \\ n \neq 0}} c_F^{-}(n) \Gamma(2k-1, -4\pi n y) e^{2\pi i n z}.$$

Here the incomplete gamma function is given by $\Gamma(\alpha, w) := \int_w^\infty e^{-t} t^{\alpha-1} dt$ (for $\text{Re}(\alpha) > 0$ and $w \in \mathbb{C}$). We call F^+ its holomorphic part of F and F^- the non-holomorphic part.

We let $H_{2\kappa}^{\text{mg}}$ be the space of weight 2κ harmonic Maass forms. The operators $\xi_{2-2k} := 2iy^{2-2k} \frac{\overline{\partial}}{\overline{\partial z}}$ and D^{2k-1} , defined in the introduction, map H_{2-2k}^{mg} to $M_{2k}^{!}$, and we let H_{2-2k} be the subspace of H_{2-2k}^{mg} consisting of forms which map to cusp forms under ξ_{2-2k} . These differential operators also act naturally on the Fourier expansion (2.2). In particular, as collected in [5, Theorems 5.5 and 5.9],

$$\xi_{2-2k}(F(z)) = (2k-1)\overline{c_F^-(0)} - (4\pi)^{2k-1} \sum_{\substack{n \gg -\infty \\ n \neq 0}} n^{2k-1} \overline{c_F^-(-n)} e^{2\pi i n z},$$

$$D^{2k-1}(F(z)) = -\frac{(2k-1)!}{(4\pi)^{2k-1}}c_F^-(0) + \sum_{n \gg -\infty} n^{2k-1}c_F^+(n)e^{2\pi i n z}.$$

Recall that the operators D^{2k-1} and ξ_{2-2k} are Hecke-equivariant (cf. [5, (7.4) and (7.5)]), i.e., for any harmonic Maass form F

$$\xi_{2-2k}(F|_{2-2k}T_m) = m^{1-2k}\xi_{2-2k}(F)|_{2k}T_m, \tag{2.3}$$

$$D^{2k-1}(F|_{2-2k}T_m) = m^{1-2k}D^{2k-1}(F)|_{2k}T_m. (2.4)$$

Here for $\kappa \in \mathbb{Z}$, the Hecke operators $T_m: H_{2\kappa}^{\text{mg}} \to H_{2\kappa}^{\text{mg}}$ are given on the expansion (2.1) by

$$F(z)|_{2\kappa}T_m := \sum_{n\in\mathbb{Z}} \sum_{d|(m,n)} d^{2\kappa-1}c_F\left(\frac{mn}{d^2}, \frac{d^2y}{m}\right) e^{2\pi i n z}.$$
 (2.5)

2.2. Regularized inner products. For two cusp forms $f, g \in S_{2k}$, Petersson's classical inner product is defined by

$$\langle f, g \rangle := \int_{\mathrm{SL}_2(\mathbb{Z}) \backslash \mathbb{H}} f(z) \overline{g(z)} y^{2k} \frac{dx dy}{y^2}.$$
 (2.6)

This was extended to an inner product over all of $M_{2k}^!$ in a series of steps. The first such attempt to do so appears to be by Petersson himself [17], which was later rediscovered and extended by Harvey–Moore [13] and Borcherds [3] as we describe below. Setting

$$\mathcal{F}_T := \left\{ z \in \mathbb{H} : |z| \ge 1, \ y \le T, \ -\frac{1}{2} \le x \le \frac{1}{2} \right\},$$

the first regularized inner product is defined by

$$\langle f, g \rangle := \lim_{T \to \infty} \int_{\mathcal{F}_T} f(z) \overline{g(z)} y^{2k} \frac{dxdy}{y^2}.$$
 (2.7)

Borcherds extended this regularization by multiplying the integrand by y^{-s} and taking the constant term of the Laurent expansion around s=0 of the analytic continuation in s. The regularization (2.7) coincides with (2.6) whenever (2.6) converges. Unfortunately, the regularization (2.7) and even Borcherds's extension also do not always exist. In particular, one cannot use them to define a meaningful norm for weakly holomorphic forms. This problem was overcome to obtain a regularized inner product for arbitrary $f, g \in M_{2k}^!$ by Diamantis, Ehlen, and the first author [4]. The idea is simple. One multiplies the integrands with a function that forces convergence and then analytically continues.

For this, observe that for $Re(w) \gg 0$, the integral

$$I(f,g;w,s) := \int_{\mathcal{F}} f(z)\overline{g(z)}y^{2k-s}e^{-wy}\frac{dxdy}{y^2}$$

converges and is analytic, while formally plugging in w=0=s yields Petersson's classical inner product (2.6). For every $\varphi\in(\pi/2,3\pi/2)\setminus\{\pi\}$ it has an analytic continuation $I_{\varphi}(f,g;w,s)$ to $U_{\varphi}\times\mathbb{C}$ with $U_{\varphi}\subset\mathbb{C}$ a certain open set containing 0. Then define

$$\langle f, g \rangle_{\varphi} := \operatorname{CT}_{s=0} I_{\varphi}(f, g; 0, s) - i \sum_{n>0} c_f(-n) \overline{c_g(-n)} \operatorname{Im} \left(E_{2-2k, \varphi}(-4\pi n) \right),$$

where $CT_{s=0} F(s)$ denotes the constant term of an analytic function F, c_F are defined as in (2.1), and $E_{r,\varphi}$ is the generalized exponential integral (see [15, (8.19.3)]) defined with branch cut on the ray $\{xe^{i\varphi}: x \in \mathbb{R}^+\}$. In [4] it is shown that for $f, g \in M_{2k}^!$, $\langle f, g \rangle$ exists and is independent of the choice of φ . Furthermore, it equals the regularization (2.7) whenever (2.7) exists.

The following lemma describes that the inner product from [4] commutes with the Hecke algebra.

Lemma 2.1. For $k \in \mathbb{N}$, the Hecke operators are Hermitian on $M_{2k}^!$.

Proof: We split the proof into three cases. First we assume that $f \in S_{2k}^!$ and $g \in M_{2k}^!$, then we use the fact that the inner product is Hermitian (see [4, Theorem 1.3]) for the case that $f \in M_{2k}^!$ and $g \in S_{2k}^!$, and finally we consider $f = g = E_{2k}$; these cover all cases by linearity.

Following [7, (1.15)], one defines, for $f, g \in M_{2k}^!$,

$$\{f,g\}_0 := \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{c_f(-n)c_g(n)}{n^{2k-1}}.$$

The ratio $c_g(n)/n^{2k-1}$ can also be interpreted as the *n*th coefficient of the holomorphic part of any harmonic Maass form G for which $D^{2k-1}(G) = g$. This yields a connection with the *Bruinier-Funke pairing* (cf. [9] for the original definition restricted to certain subspaces)

$$\{f,G\} := \sum_{n \in \mathbb{Z}} c_f(-n)c_G^+(n)$$

between $f \in M_{2k}^!$ and $G \in H_{2-2k}^{mg}$.

We first assume that $f \in \bar{S}_{2k}^{!}$ and $g \in M_{2k}^{!}$. In this case

$$\{f,g\}_0 = \{f,G\}.$$
 (2.8)

This is related to the inner product by [4, Theorem 4.1]; namely,

$$\{f,G\} = \langle f, \xi_{2-2k}(G) \rangle. \tag{2.9}$$

Setting $h := \xi_{2-2k}(G)$, we combine (2.9) with (2.8) and use the fact that the Hecke operators are Hermitian with respect to $\langle \cdot, \cdot \rangle_0$ to obtain

$$\langle f|_{2k}T_m, h\rangle = \{f|_{2k}T_m, G\} = \{f|_{2k}T_m, g\}_0 = \{f, g|_{2k}T_m\}_0.$$

Next, by (2.4), we have

$$D^{2k-1}(G|_{2-2k}T_m) = m^{1-2k}D^{2k-1}(G)|_{2k}T_m = m^{1-2k}g|_{2k}T_m.$$

Plugging this into (2.8) and using (2.9), yields

$$\{f, g|_{2k}T_m\}_0 = m^{2k-1} \{f, G|_{2-2k}T_m\} = m^{2k-1} \langle f, \xi_{2-2k} (G|_{2-2k}T_m) \rangle$$
.

We finally use (2.3) to obtain

$$m^{2k-1} \langle f, \xi_{2-2k} \left(G|_{2-2k} T_m \right) \rangle = \langle f, h|_{2k} T_m \rangle.$$

Altogether, we have therefore shown that

$$\langle f|_{2k}T_m, h\rangle = \langle f, h|_{2k}T_m\rangle. \tag{2.10}$$

Since ξ_{2-2k} is surjective on $M_{2k}^!$, we obtain the claim for arbitrary $f \in S_{2k}^!$ and $h \in M_{2k}^!$.

Next suppose that $f \in M_{2k}^!$ and $g \in S_{2k}^!$. Since the inner product is Hermitian (see [4, Theorem 1.3]), by (2.10) we have

$$\langle f|_{2k}T_m, g\rangle = \overline{\langle g, f|_{2k}T_m\rangle} = \overline{\langle g|_{2k}T_m, f\rangle} = \langle f, g|_{2k}T_m\rangle.$$

Finally, for $f = E_{2k} = g$, we obtain the result directly from the fact that both f and g are Hecke eigenforms with the same real eigenvalues.

3. Degeneracy and Proofs of Theorem 1.1 and Corollary 1.2

In this section, we show that the space $M_{2k}^{!,\perp}$ defined in (1.3) coincides with the space $D^{2k-1}(S_{2-2k}^!)$ which appears in the definition of Hecke eigenforms for weakly holomorphic modular forms. In [4, Corollary 4.5], the space $M_{2k}^{!,\perp}$, defined in (1.3), was explicitly determined. Theorem 1.1 rewrites this in a form which is useful for Hecke eigenforms. In order to prove Theorem 1.1, we also require the following useful lemma about the *flipping operator*

$$\mathfrak{F}_{\kappa}\left(F(z)\right) := -\frac{y^{-\kappa}}{(-\kappa)!} \overline{R_{\kappa}^{-\kappa}\left(F(z)\right)},$$

where

$$R_{\kappa}^{\ell} := R_{\kappa+2\ell-2} \circ \cdots \circ R_{\kappa}$$

is the repeated raising operator with the raising operator defined by

$$R_{\kappa} := 2i \frac{\partial}{\partial z} + \frac{\kappa}{y}.$$

The following lemma follows from the calculation in the proof of [10, Theorem 1.1] and [10, Remark 7]; it may be found in this form in [5, Proposition 5.14].

Lemma 3.1. If $F \in H_{2-2k}^{mg}$, then the operator \mathfrak{F}_{2-2k} satisfies

$$\xi_{2-2k}(\mathfrak{F}_{2-2k}(F)) = \frac{(4\pi)^{2k-1}}{(2k-2)!} D^{2k-1}(F), \tag{3.1}$$

$$D^{2k-1}(\mathfrak{F}_{2-2k}(F)) = \frac{(2k-2)!}{(4\pi)^{2k-1}} \xi_{2-2k}(F). \tag{3.2}$$

Furthermore, $\mathfrak{F}_{2-2k} \circ \mathfrak{F}_{2-2k}(F) = F$.

We now use Lemma 3.1 to prove Theorem 1.1.

Proof of Theorem 1.1: By [4, Corollary 4.5], $f \in M_{2k}^!$ is orthogonal to $M_{2k}^!$ if and only if there exists $F \in H_{2-2k}^{\text{mg}}$ for which $F^+ = 0$ and $\xi_{2-2k}(F) = f$. By (3.2), we have

$$f = \xi_{2-2k}(F) = \frac{(4\pi)^{2k-1}}{(2k-2)!} D^{2k-1}(\mathfrak{F}_{2-2k}(F)). \tag{3.3}$$

Note that $F^+=0$ is equivalent to $D^{2k-1}(F)=0$ and $c_F^+(0)=0$. Since $D^{2k-1}(F)=0$, (3.1) implies that

$$\xi_{2-2k}(\mathfrak{F}_{2-2k}(F)) = \frac{(4\pi)^{2k-1}}{(2k-2)!}D^{2k-1}(F) = 0.$$

We conclude that $\mathfrak{F}_{2-2k}(F) \in M_{2-2k}^!$ because this is the kernel of ξ_{2-2k} . Thus by (3.3), we have $f \in D^{2k-1}(M_{2-2k}^!)$. Moreover, [8, displayed formula before Lemma 3.1] states that (note that in this paper the flipping operator is renormalized)

$$c_{\mathfrak{F}_{2-2k}(F)}^+(0) = -\overline{c_F^+(0)} = 0.$$

Thus $F^+ = 0$ is equivalent to $f \in D^{2k-1}(S^!_{2-2k})$, which is the claim. \square

We conclude the paper by using Theorem 1.1 to show that one may obtain an alternative characterization of weakly holomorphic Hecke eigenforms by requiring (1.2) to hold for all $g \in M_{2k}^!$ instead of (1.1).

Proof of Corollary 1.2: By Theorem 1.1, (1.1) is equivalent to

$$f|T_m \equiv \lambda_m f \pmod{M_{2k}^{!,\perp}}$$
.

Hence f is a weakly holomorphic Hecke eigenform if and only if for every m there exists $\lambda_m \in \mathbb{C}$ such that $h := f|T_m - \lambda_m f \in M^{!,\perp}_{2k}$. Linearity of the regularized inner product implies that for every $g \in M^{!,\perp}_{2k}$,

$$\langle f|T_m,g\rangle = \langle \lambda_m f + h,g\rangle = \lambda_m \langle f,g\rangle + \langle h,g\rangle,$$

which satisfies (1.2) for every $g \in M_{2k}^!$ if and only if $h \in M_{2k}^{!,\perp}$.

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Mathematical Institute, University of Cologne, Weyertal 86-90, 50931 Cologne, Germany

 $E ext{-}mail$ address: kbringma@math.uni-koeln.de

Department of Mathematics, University of Hong Kong, Pokfulam, Hong Kong

E-mail address: bkane@hku.hk